

Topological antiferromagnetic spintronics and the crystal Hall effect

Jairo Sinova

Johannes Gutenberg Universität Mainz



23rd of October 2019
KITP, Santa Barbara



European
Research
Council



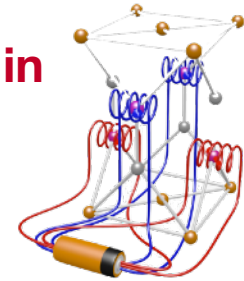
Topological antiferromagnetic spintronics and the crystal Hall effect

I. From Spin-Orbit Torque in FM to Néel Spin-Orbit Torques in Antiferromagnets:

Zelezny, Gao, JS, Jungwirth PRL (2014)

Zelazny, Gao, et al. PRB (2016)

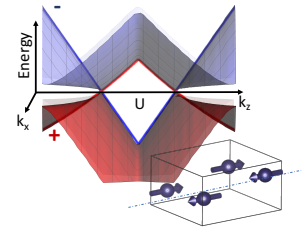
Bodnar, Smejkal, Jourdan, Kläui, Jungwirth, JS, et al, Nat. Com. (2018)



II. Topological Dirac Fermion + Antiferromagnets + Neel SOTs

Smejkal, Zelezny, Sinova, Jungwirth PRL (2017)

Smejkal, Jungwirth, Sinova PSS (2017)

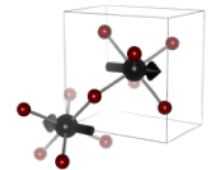


III. AHE in collinear AFMs: Crystal Hall Effect

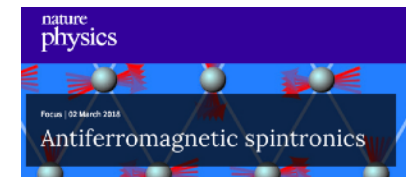
Jungwirth, JS, et al, Nature Physics (AFM spintronics Reviews) (2018)

Surprises of the Spin Hall Effect, Physics Today 70, 7, 38 (2017)

Smejkal, Gonzales, Jungwirth, JS, arXiv 1901.00445 (2019)



Jungwirth, JS, et al, Nature Physics (AFM spintronics Reviews) (2018)



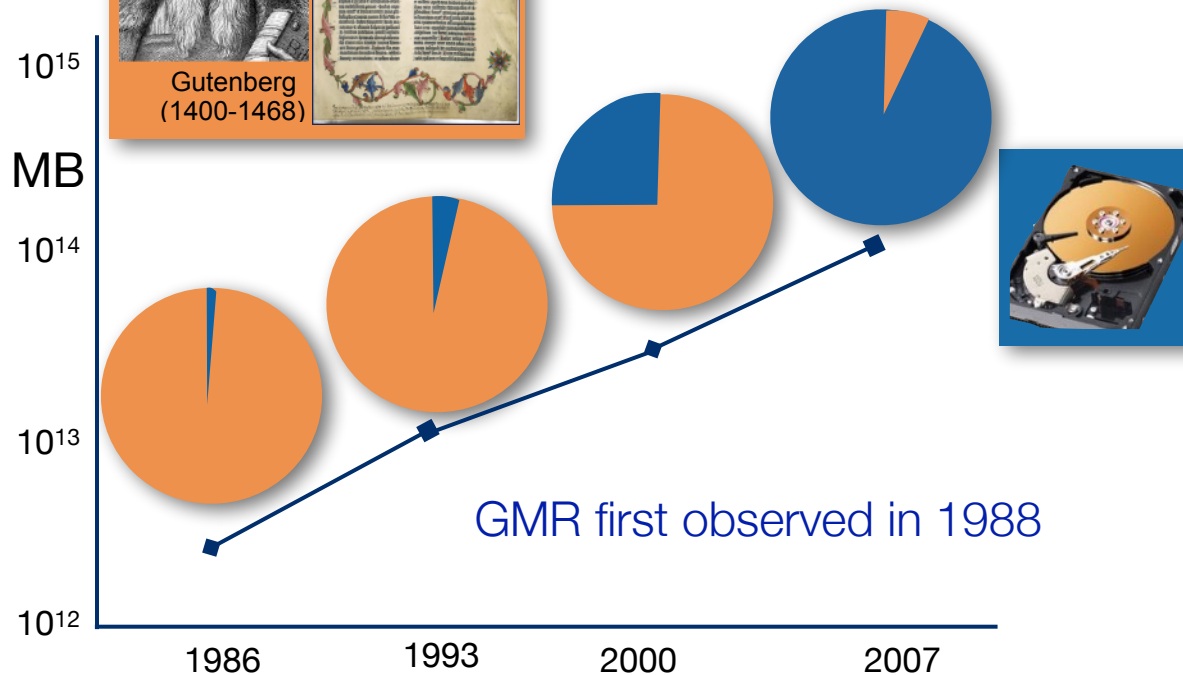
Analog to digital = Ink to Spin

Analog: books, video/film, ...

Digital: Hard-disks, DVDs, ...



Analog to Digital



Hilbert et al. Science (2011)

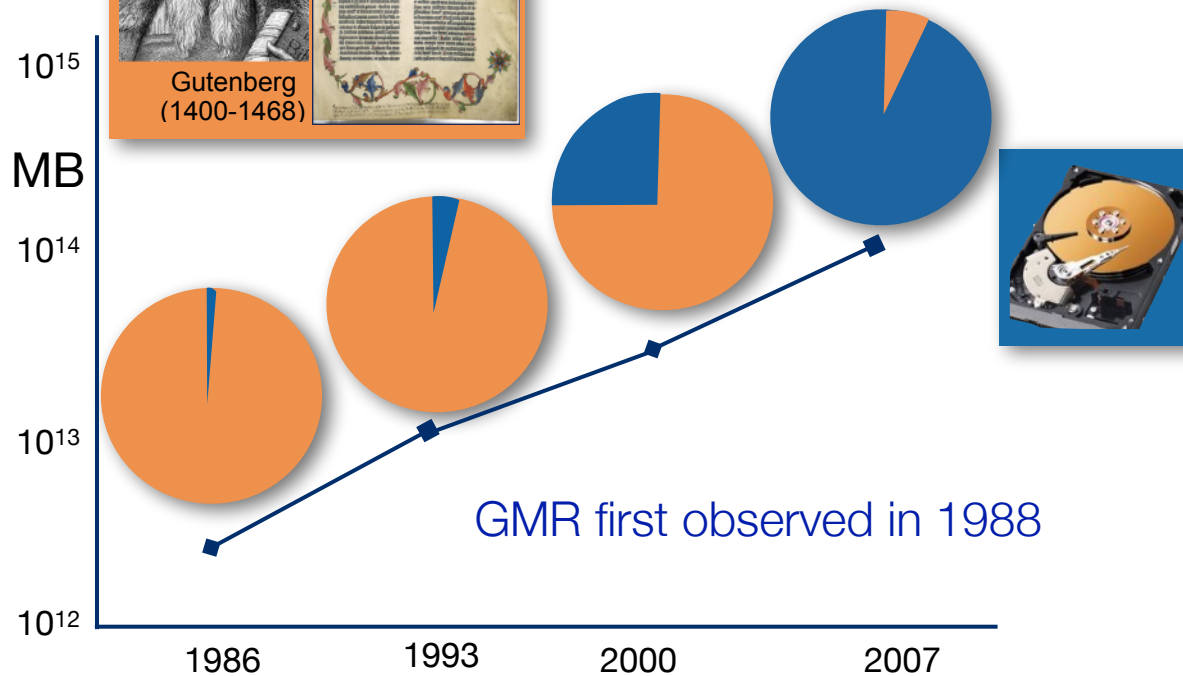
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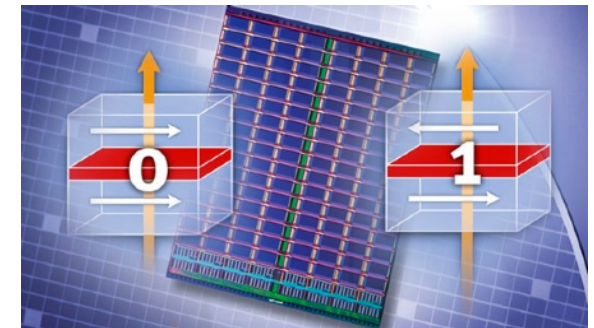


Analog to Digital



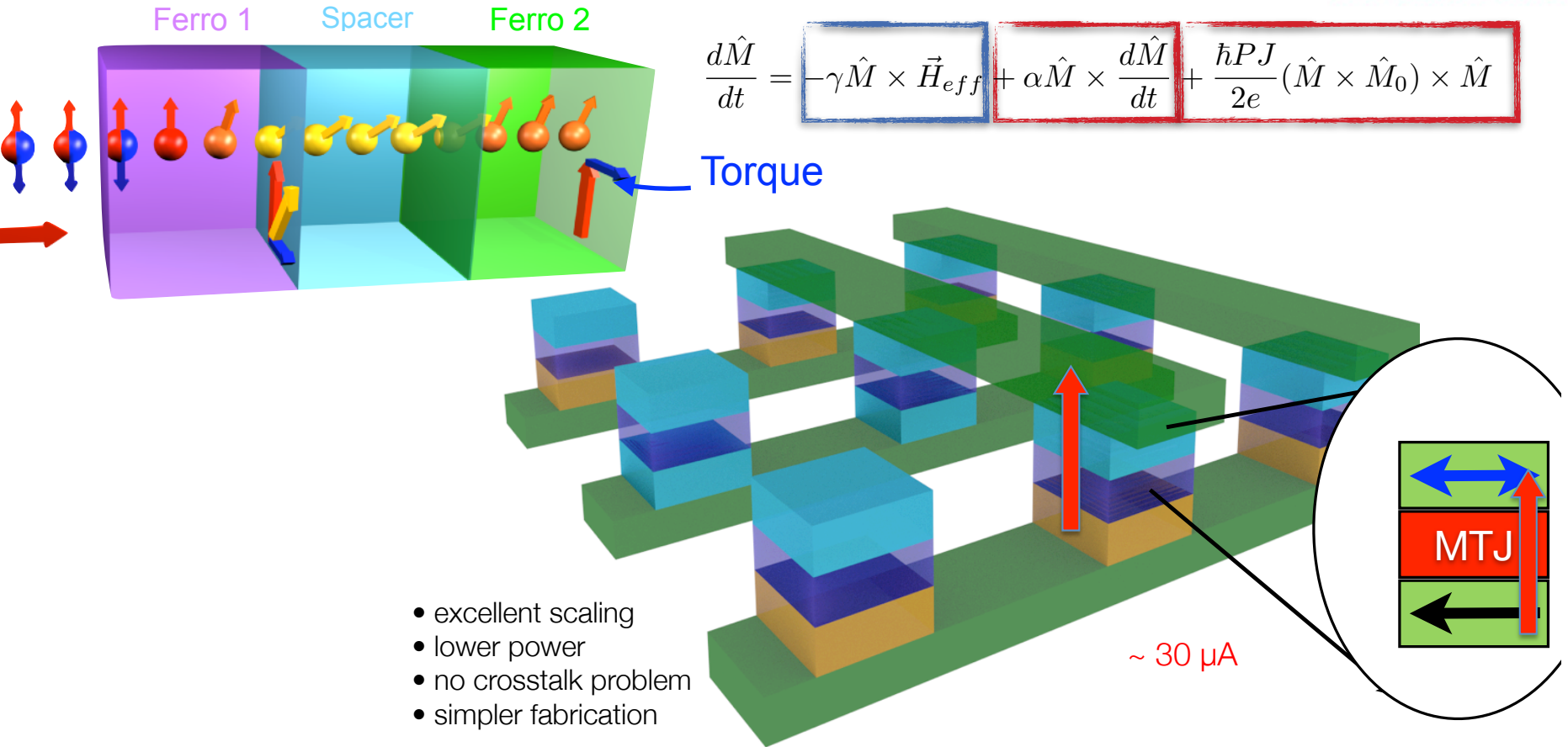
What is next in memory storage?

Spin-Transfer MRAM (2015)



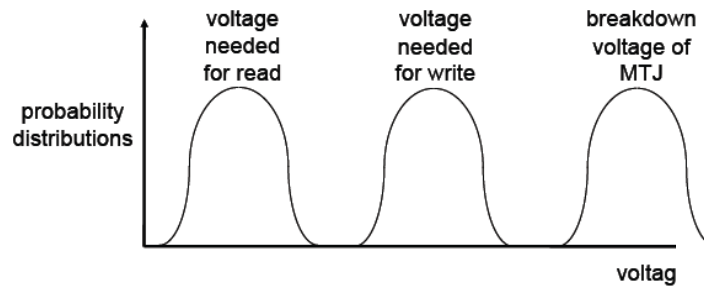
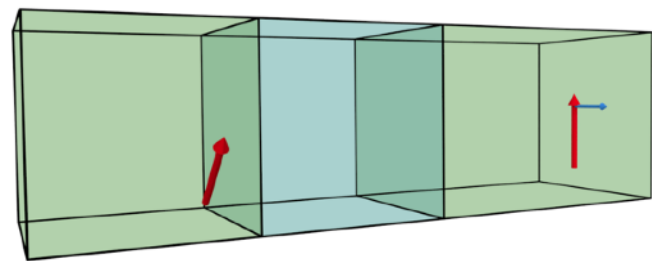
Hilbert et al. Science (2011)

Spin-Transfer-Torque MRAM

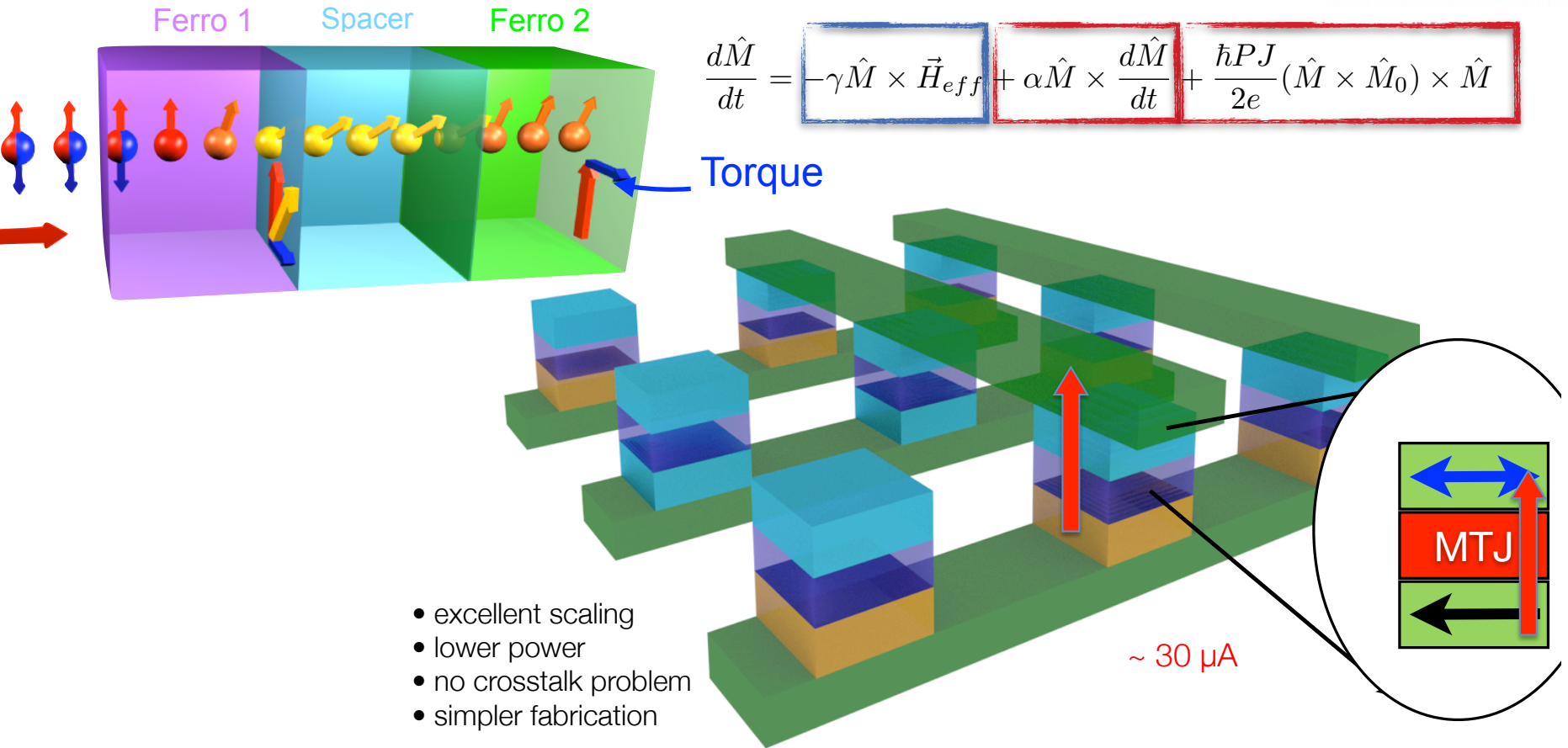


$$\frac{d\hat{M}}{dt} = -\gamma\hat{M} \times \vec{H}_{eff} + \alpha\hat{M} \times \frac{d\hat{M}}{dt} + \frac{\hbar PJ}{2e} (\hat{M} \times \hat{M}_0) \times \hat{M}$$

- excellent scaling
- lower power
- no crosstalk problem
- simpler fabrication

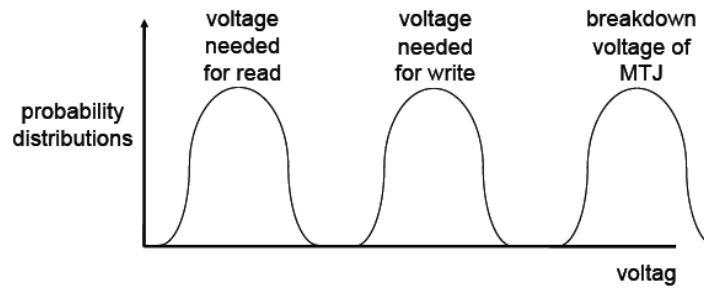
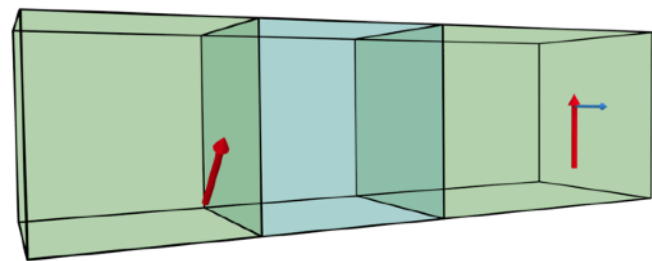


Spin-Transfer-Torque MRAM



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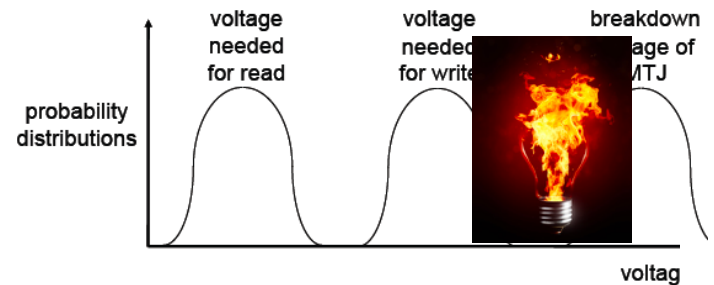
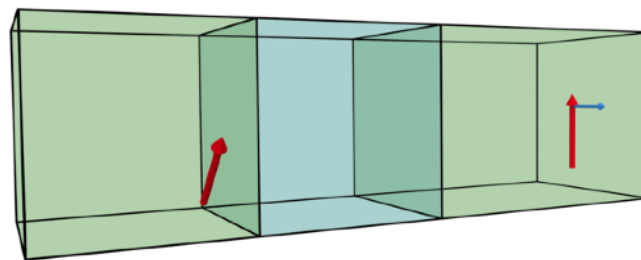
Spin-Transfer-Torque MRAM

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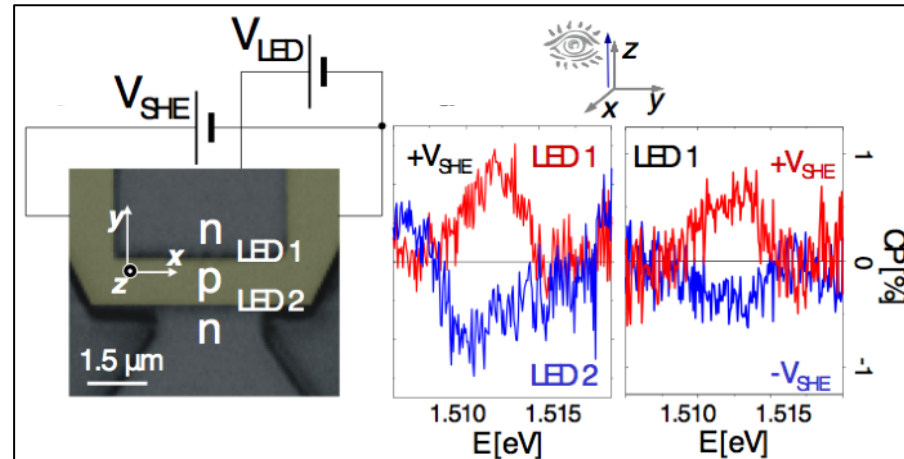
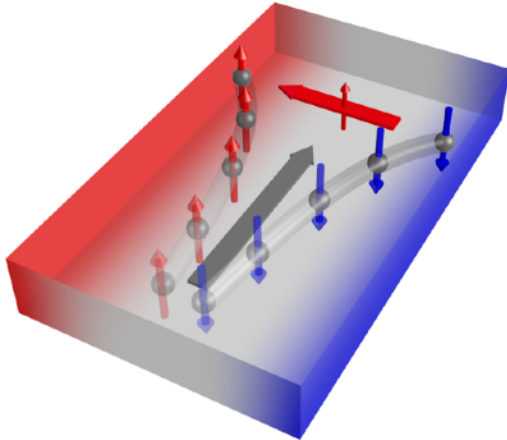
Torque

- excellent scaling
- lower power
- no crosstalk problem
- simpler fabrication

~ 30 μ A



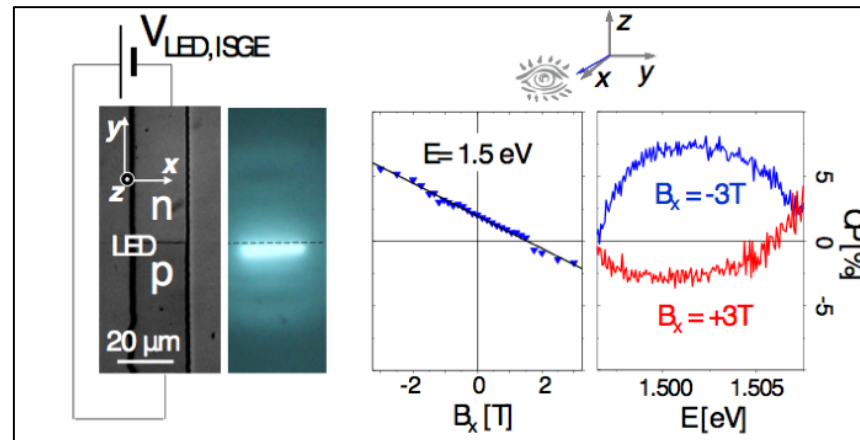
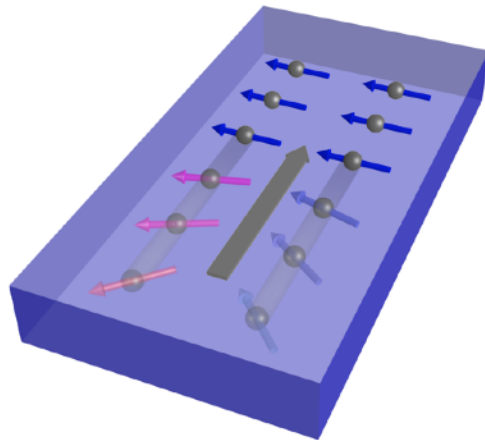
Spin Hall Effect in p-GaAs



Wunderlich et al. arXiv '04, PRL '05

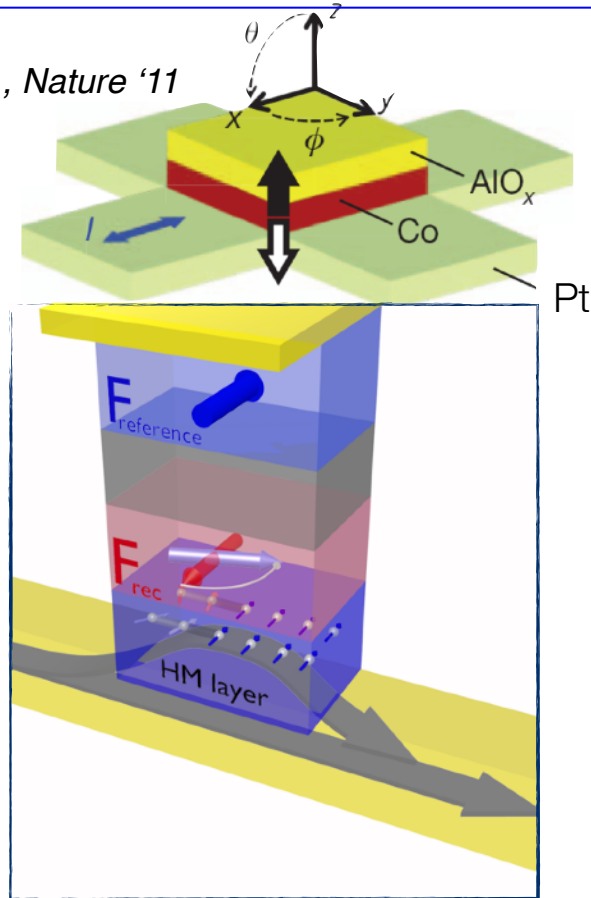
Inverse Spin Galvanic Effect or Edelstein Effect

(Reverse process of circular photo-galvanic effect, Ganichev et al., 2001)



spin-orbit torque at PM/FM interface

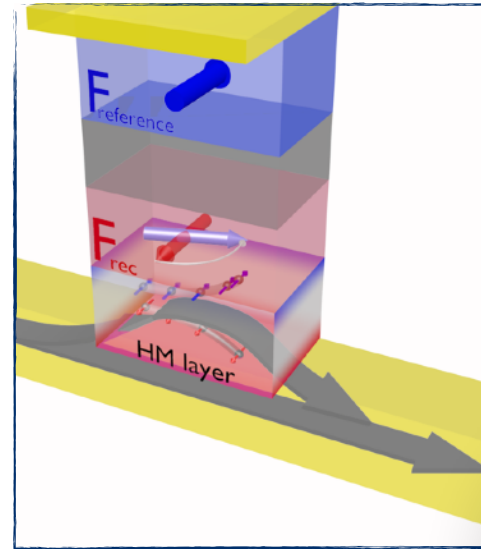
Miron et al., Nature '11



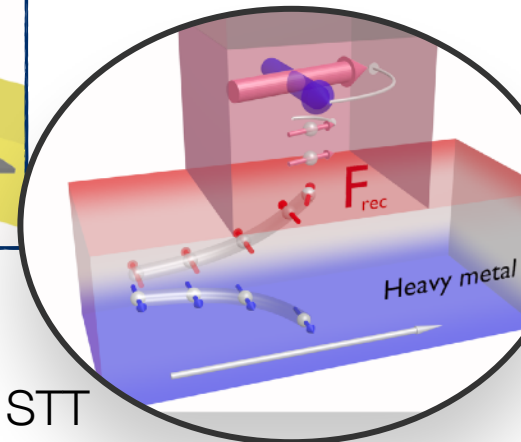
$$H_{ex} = J_{ex} \vec{M} \cdot \delta \vec{s} \quad \left(\frac{d\vec{M}}{dt} \right)_{SOT} = \frac{J_{ex}}{\hbar} \vec{M} \times \delta \vec{s}$$

$$h_{SOT} \parallel z \times J$$

SHE as spin-current generator + STT



Buhrman, et al., Science '12

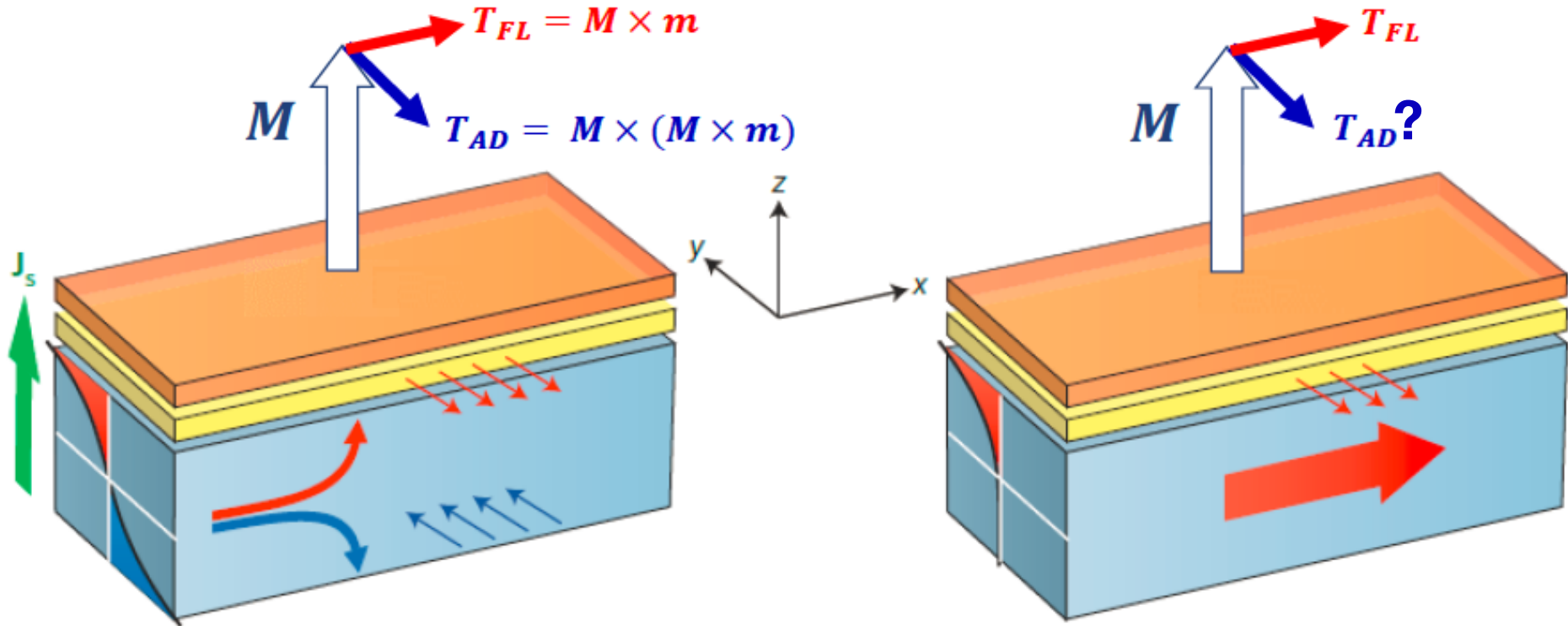


intrinsic SHE + STT

$$\left(\frac{d\vec{M}}{dt} \right)_{SHE-STT} = P \hat{M} \times (\hat{n} \times \hat{M})$$

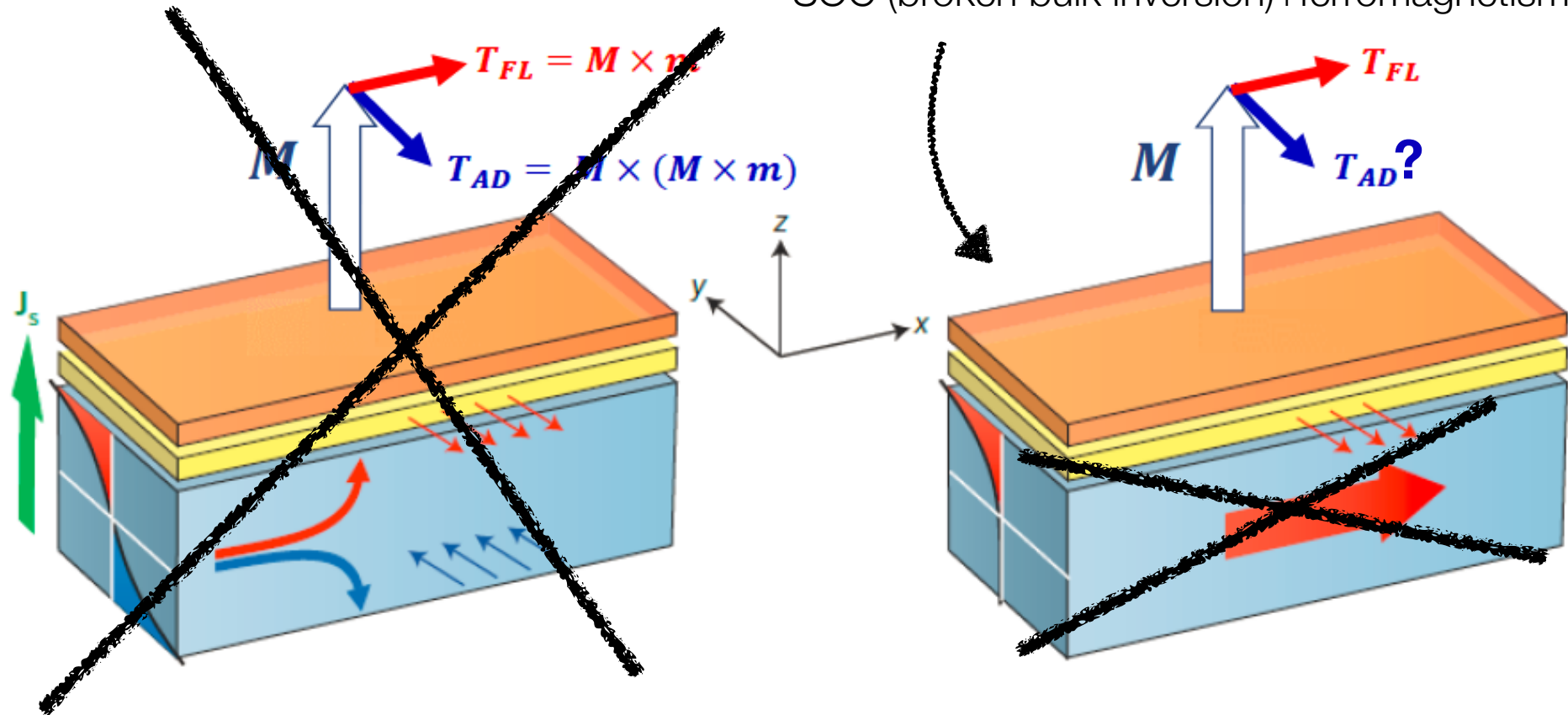
Jairo Sinova and Tomas Jungwirth, Physics Today (2017)

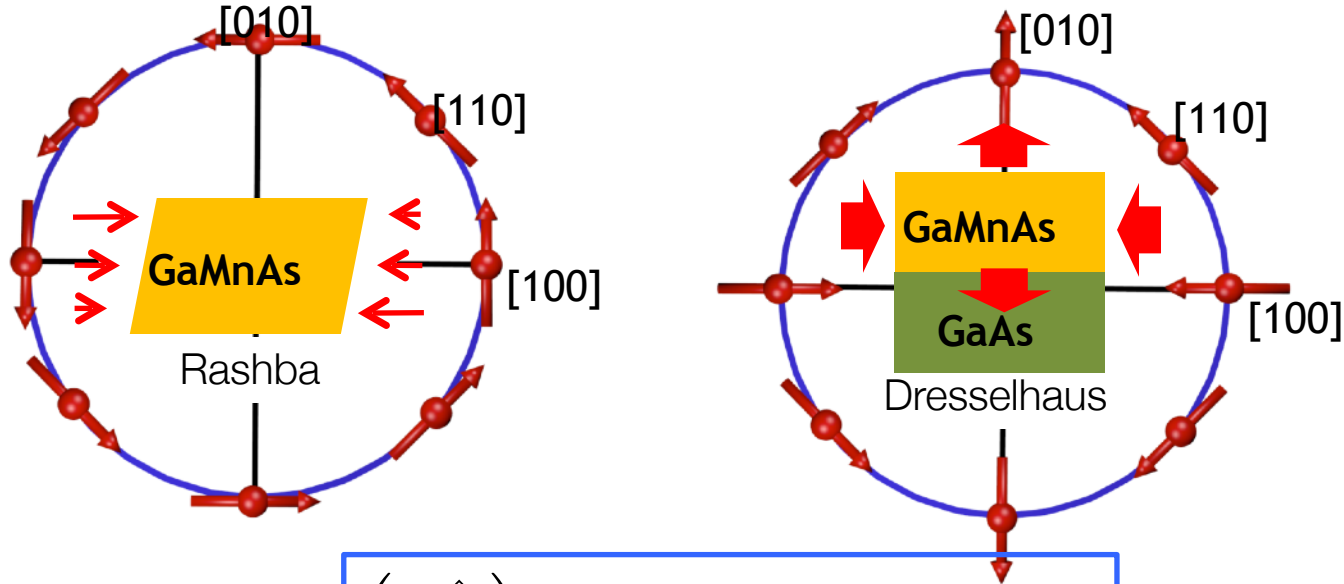
Spin-orbit Torques in Bilayer Systems



Spin-orbit Torques in Bilayer Systems

Make a ferromagnet behave like a cat:
SOC (broken bulk inversion)+ferromagnetism





$$\left(\frac{d\hat{M}}{dt} \right)_{SOT} = \hat{M} \times \delta s_z(\theta_{\mathbf{M}-\mathbf{E}}) \hat{z}$$

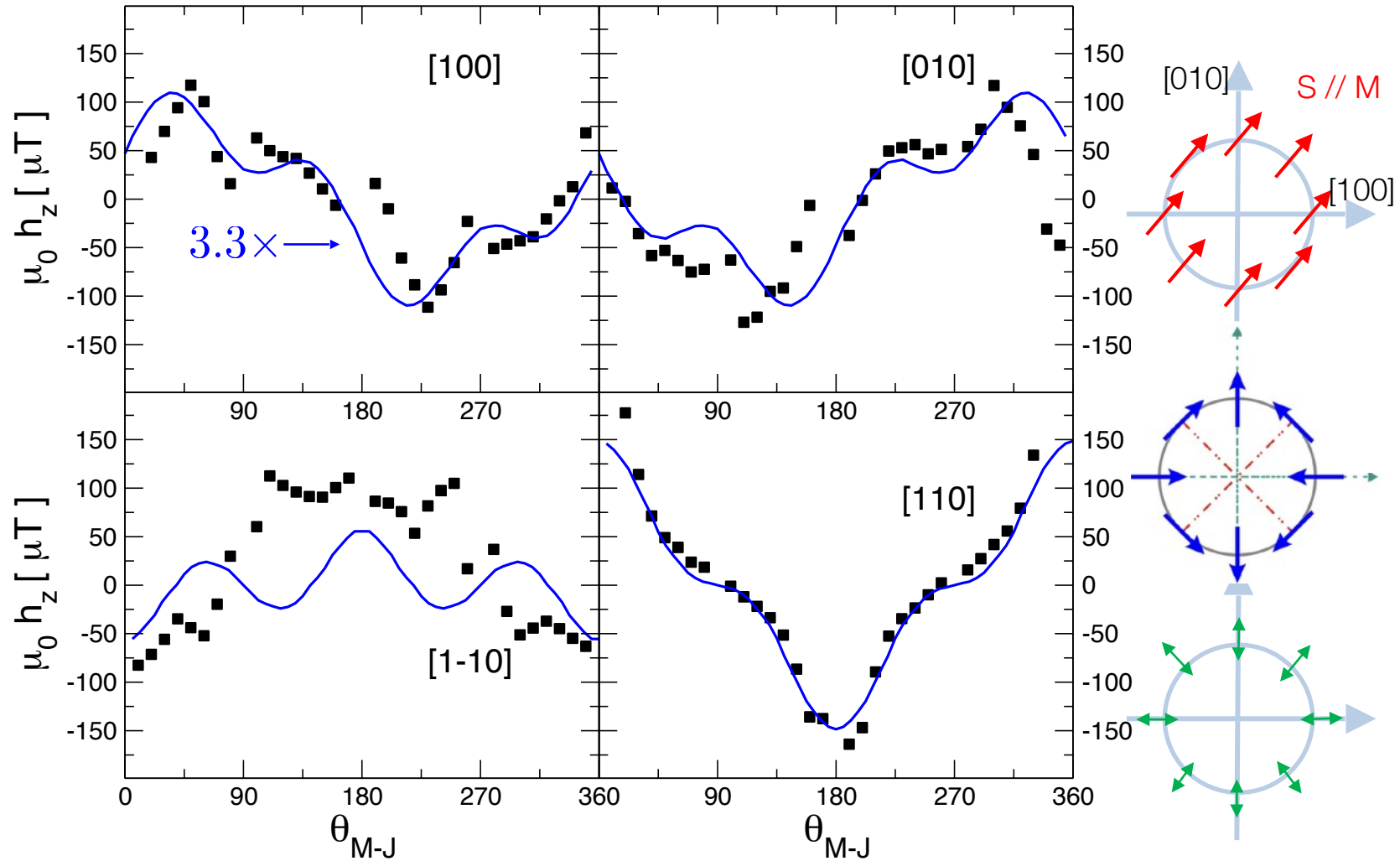
angle between \mathbf{M} and current direction

Kurebayashi, JS, et al.,
Nat. Nanot. (2014)

current direction	Rashba: $\delta s_{z,M} \sim$	Dresselhaus: $\delta s_{z,M} \sim$
$\mathbf{E} \parallel [100]$	$\cos \theta_{\mathbf{M}-\mathbf{E}}$	$\sin \theta_{\mathbf{M}-\mathbf{E}}$
$\mathbf{E} \parallel [010]$	$\cos \theta_{\mathbf{M}-\mathbf{E}}$	$-\sin \theta_{\mathbf{M}-\mathbf{E}}$
$\mathbf{E} \parallel [110]$	$\cos \theta_{\mathbf{M}-\mathbf{E}}$	$\cos \theta_{\mathbf{M}-\mathbf{E}}$
$\mathbf{E} \parallel [1 - 10]$	$\cos \theta_{\mathbf{M}-\mathbf{E}}$	$-\cos \theta_{\mathbf{M}-\mathbf{E}}$



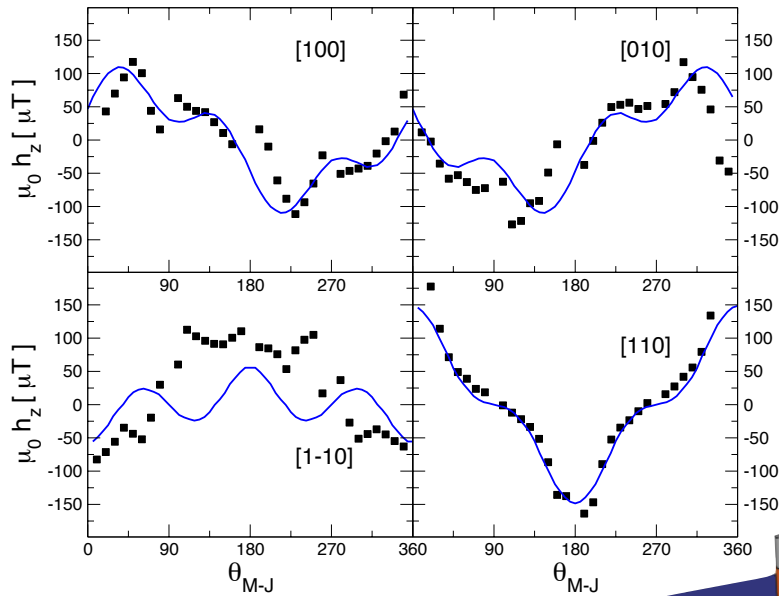
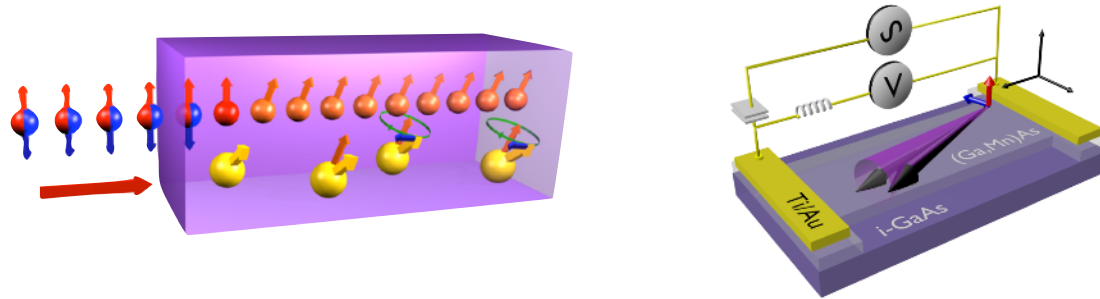
Solid line: Calculations with H_{KL} (captures higher harmonics)



Kurebayashi, JS, et al., Nat. Nanot. (2014)

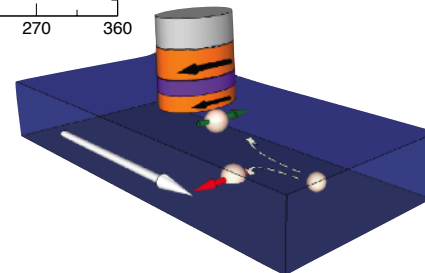


We have achieved a “cat” magnet



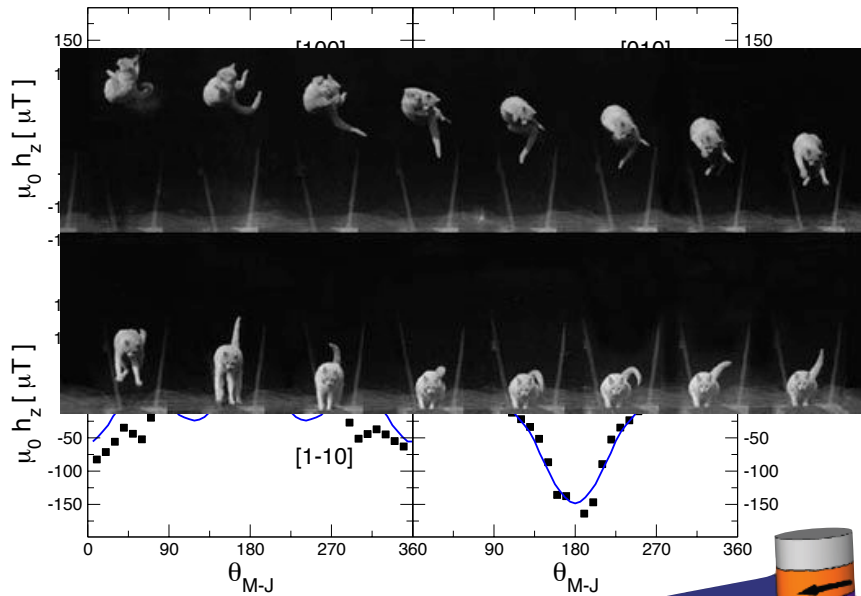
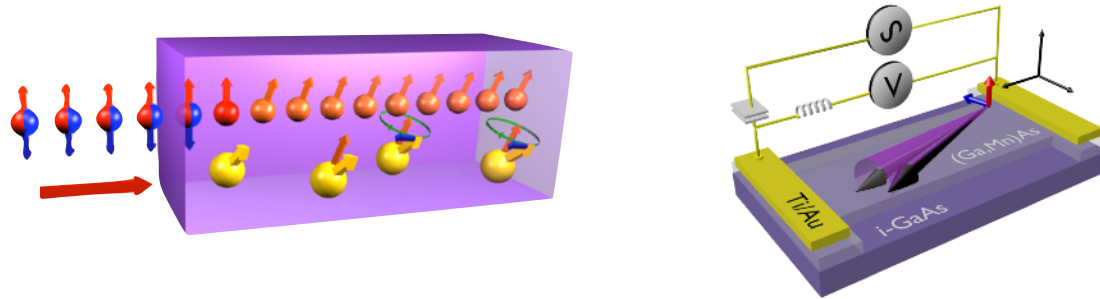
Kurebayashi, JS, et al.,
Nat. Nanot. (2014)

MRAMs without melting



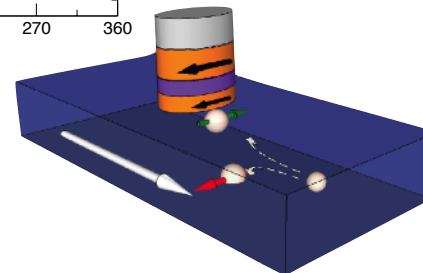
2011-2014

We have achieved a “cat” magnet



Kurebayashi, JS, et al.,
Nat. Nanot. (2014)

MRAMs without melting



2011-2014



Ordered spins
Non-volatile

Spin not charge based
Radiation-hard

No net moment
**Insensitive to magnetic fields,
no fringing stray fields**

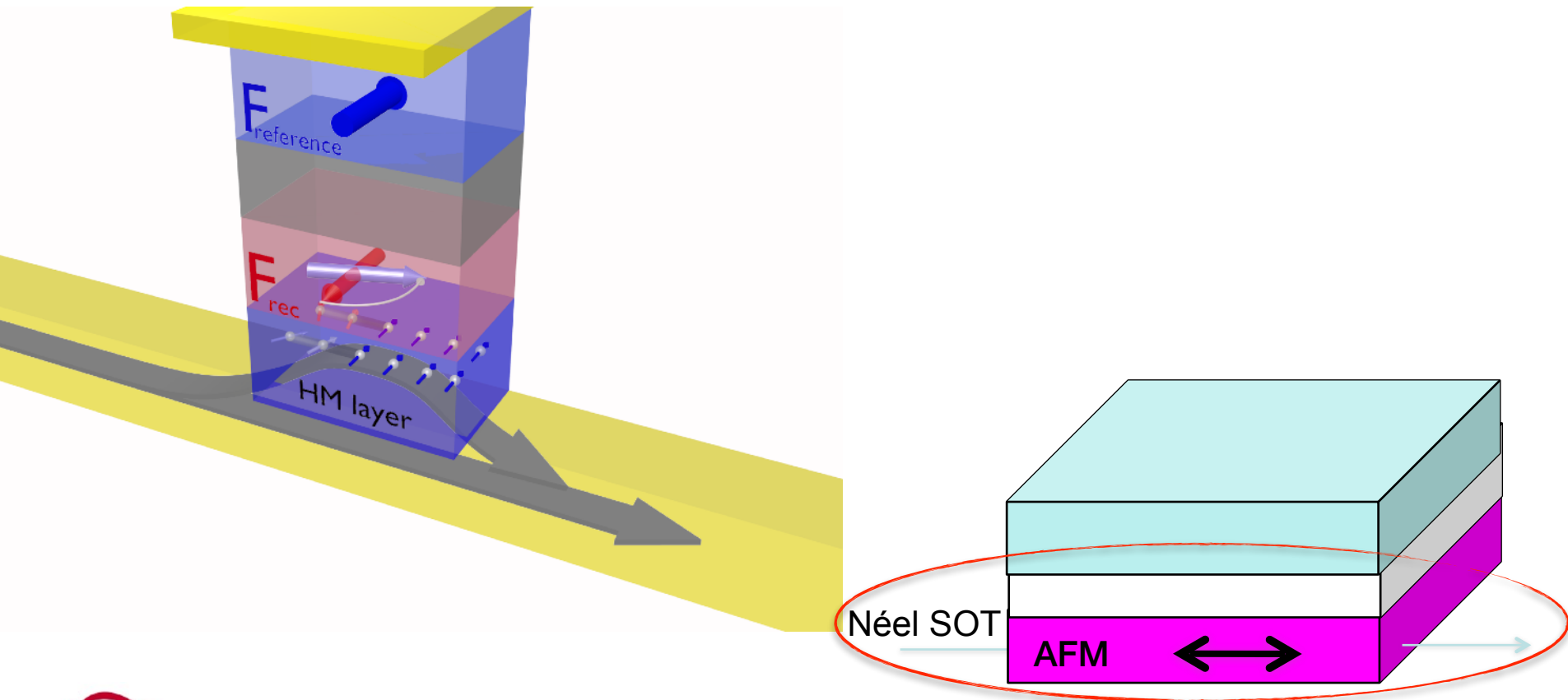
THz dynamics
Ultra-fast switching

Multiple-stable domain configurations
Memory-logic bit cells

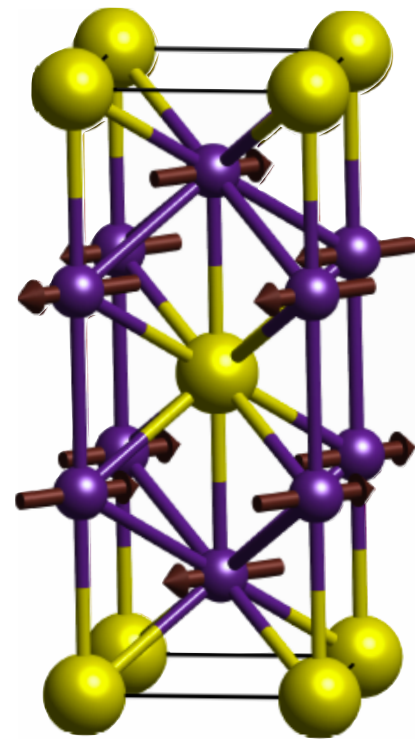
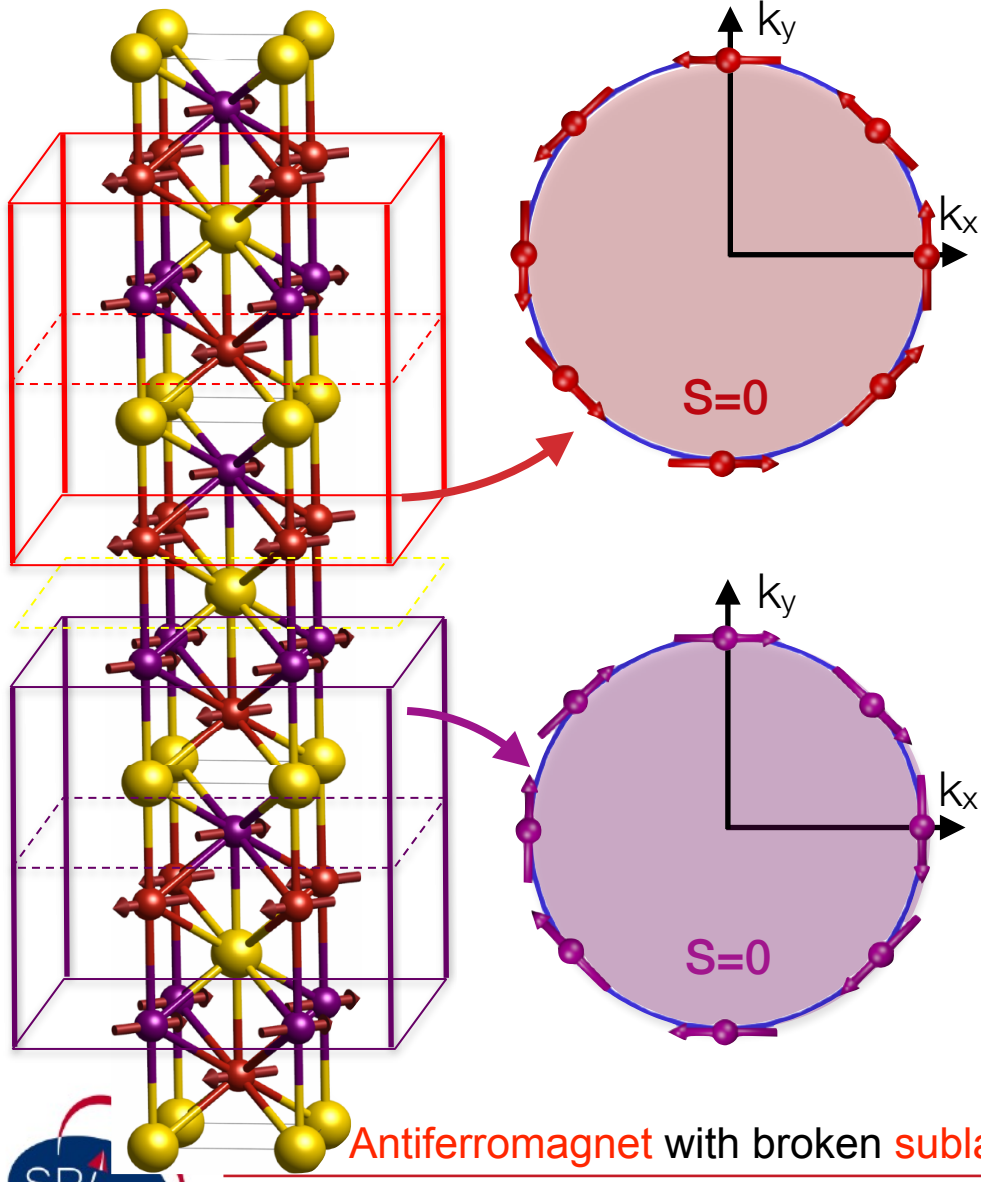
Materials range
**Insulators, semiconductors, semimetals,
metals, superconductors**

Writing by spin-orbit torque in a single-layer ferromagnet

Magnet reversing itself : SOT



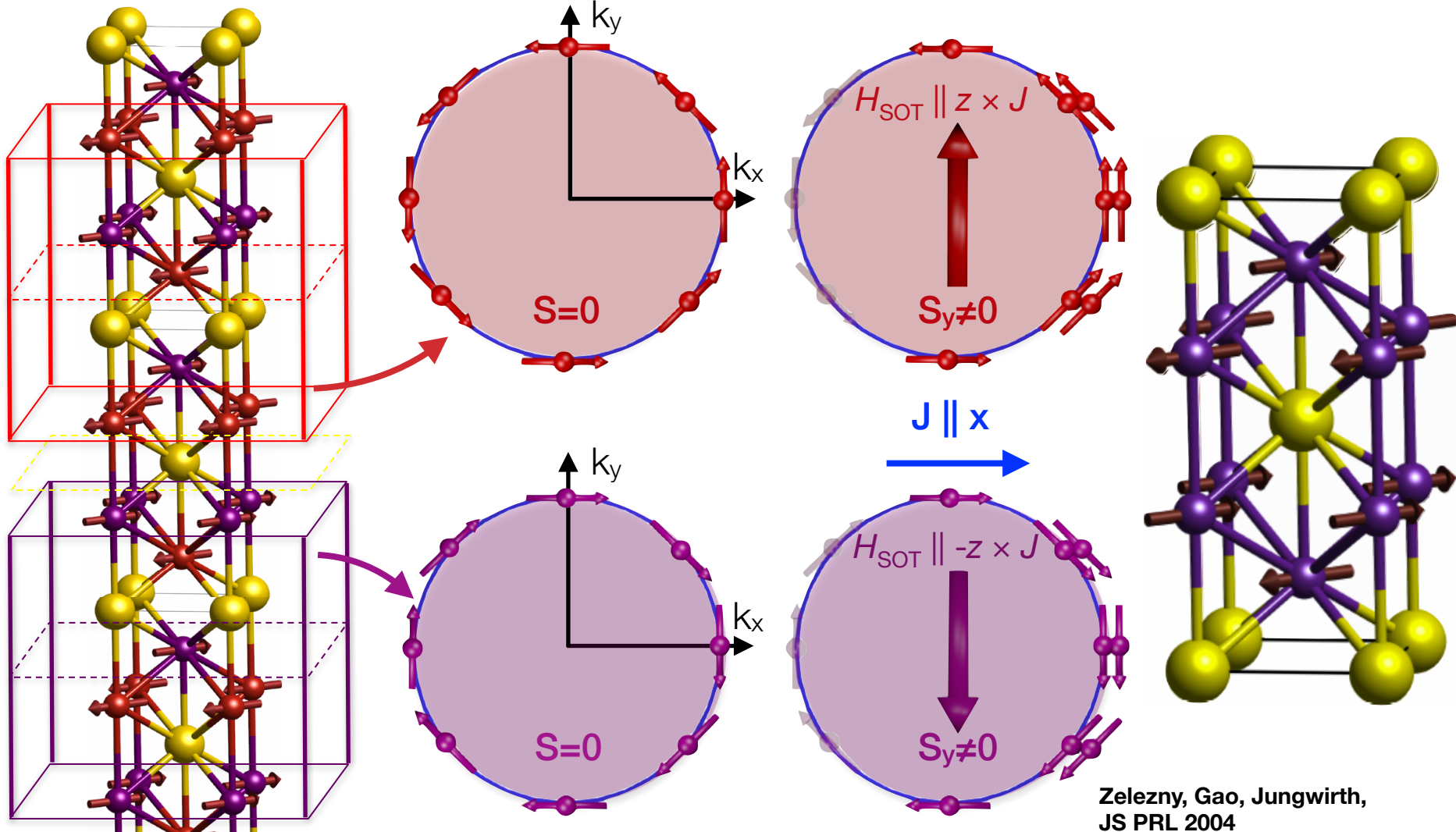
J. Zelezny, H. Gao, K. Vyborny, J. Masek, J. Zemen, A. Manchon, J. Sinova, and T. Jungwirth, *PRL* (2014)



Zelezny, Gao, Jungwirth,
JS PRL 2004

Antiferromagnet with broken sublattice space-inversion symmetry: (Mn₂Au)



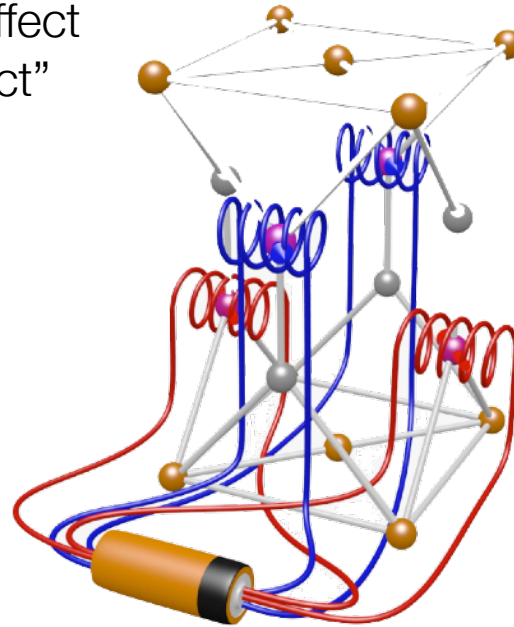
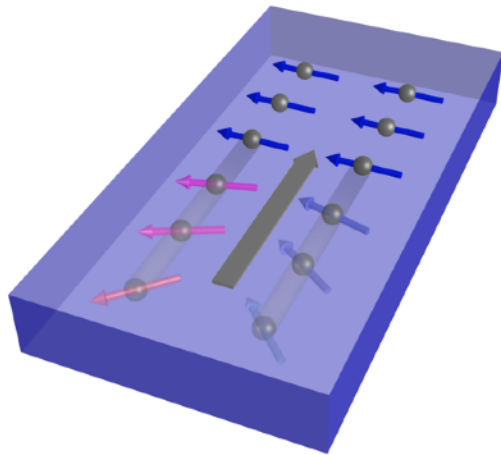


Zelezny, Gao, Jungwirth,
JS PRL 2004

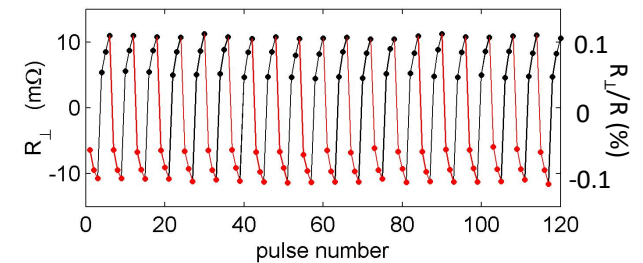
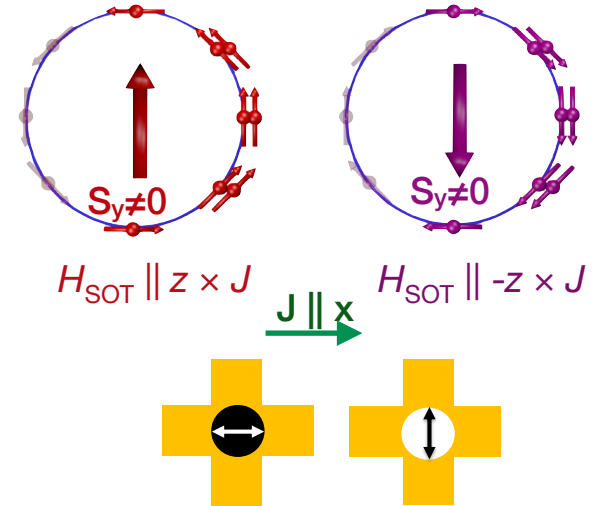
Antiferromagnet with broken sublattice space-inversion symmetry: (Mn₂Au)



Inverse spin galvanic effect (ISGE) “Edelstein effect”



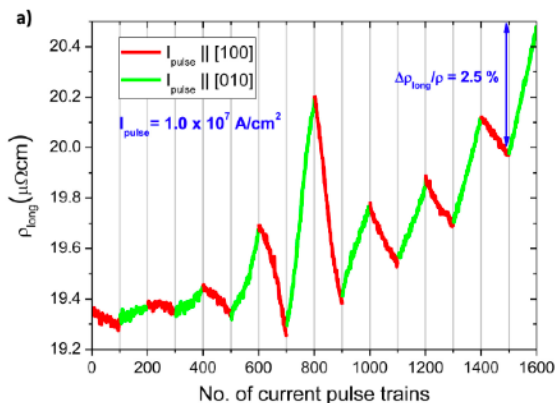
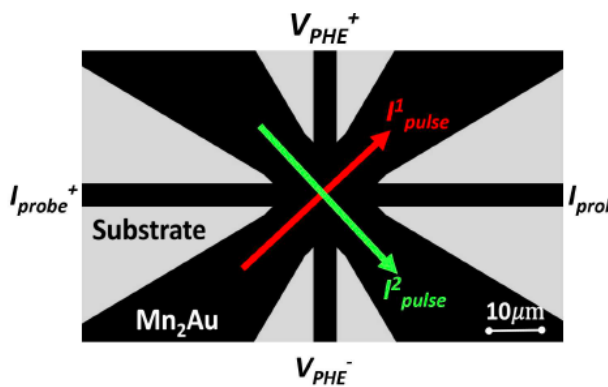
J. Zelezny, J. Sinova, T. Jungwirth et al., Phys. Rev. Lett. (2014)



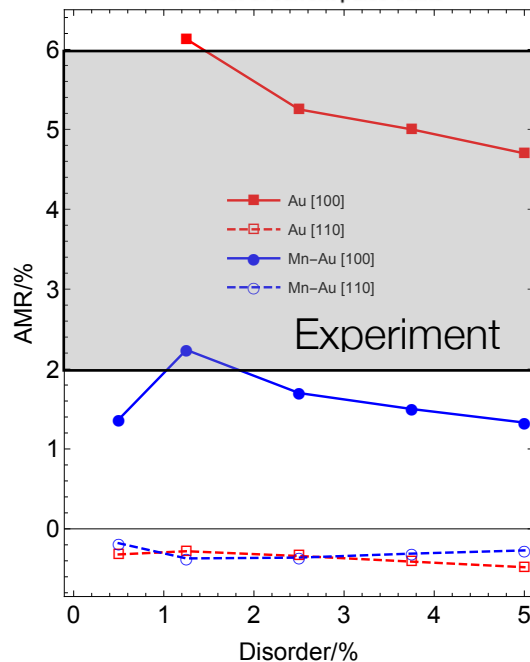
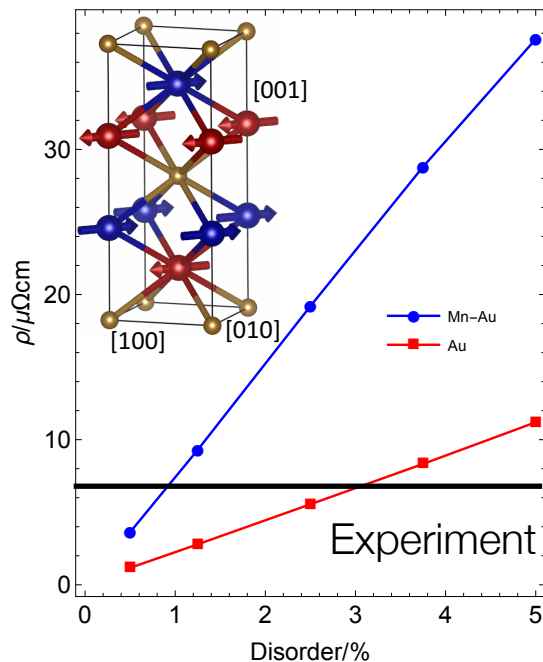
P. Wadley, T.J., et al., Science (2016)
Bodnar, Smajkal, T.J., J.S., et al. Nat. Comm. (2018),

Functionalities based on electrical manipulation of antiferromagnetic order





Relativistic DFT:



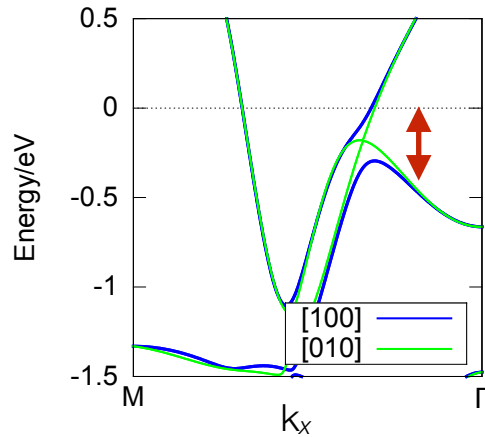
- Experiment:
- *S. Bodnar et al.**, (2.5-6 %),
- *M. Meinert et al.*, *arXiv. 1706.06983* (5 %D: *App. Phys*)
- *M. Jourdan et al.*, *JoP.* (2015): **8 μΩcm**
- *H. Wu et al.*, *AFM* (2016) (1-2.5 %)

0.6±0.3% Au rich,
Mn-Au <5%

Bodnar, Smejkal, Jourdan, Kläui, Jungwirth, JS, et al, *Nature Communications* (2018)

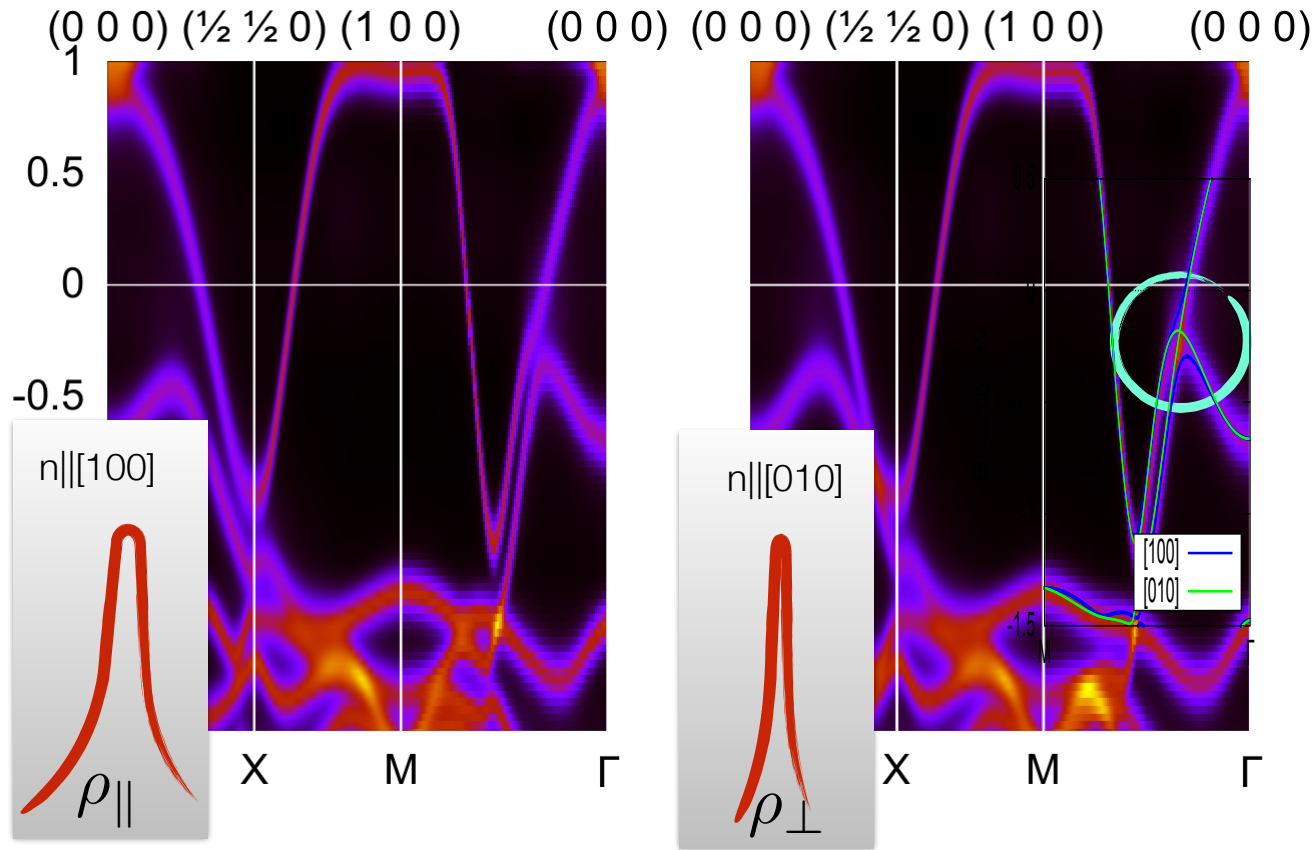


$$E(\mathbf{k}) \rightarrow A(E, \mathbf{k}) = -\frac{1}{\pi} \text{Im} \bar{G}$$



$$\rho_{\parallel} > \rho_{\perp}$$

$$\text{AMR} > 0$$



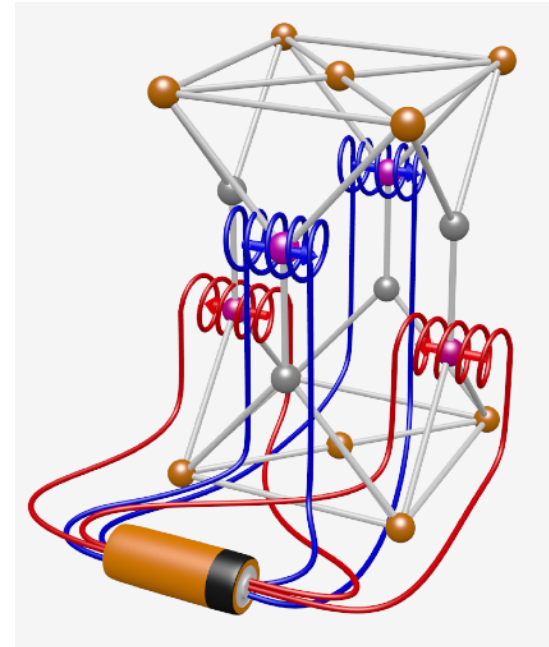
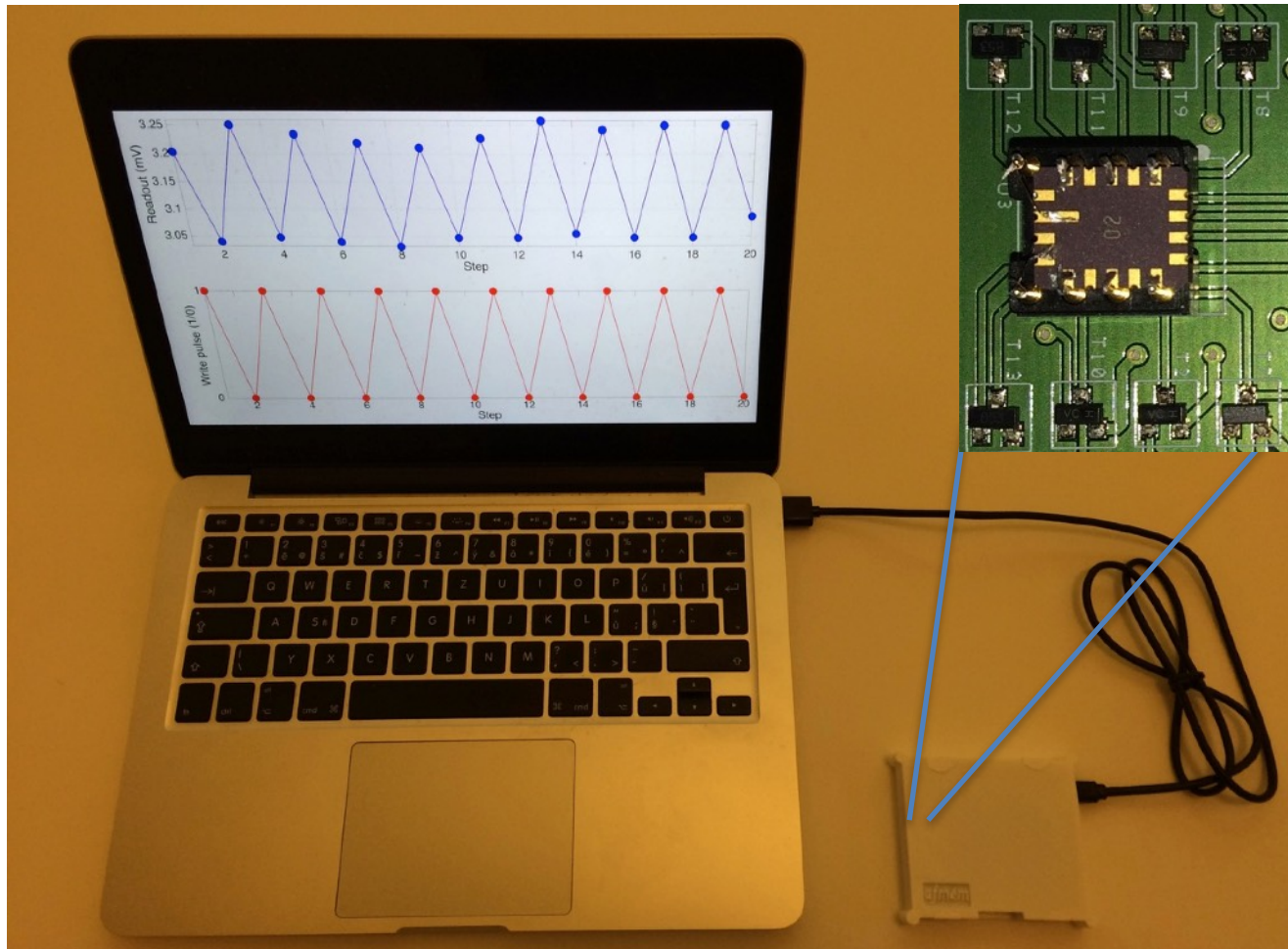
Angular dependence of Bloch spectral function

- CPA preserves the symmetry prevented hybridisations
- Disorder propagates Dirac point to Fermi level and shifts the Fermi level

5% Au-rich **AMR 6%**

S. Y. Bodnar, L. Smajkal, T. Jungwirth, J. Sinova, M. Kläui, M. Jourdan, et al, Nat. Comm. (2018)

From prediction, to observation, to device in 1 one year!!



**Works like this but
not done like this**

Electrical read/write antiferromagnetic memory

Wadley, Jungwirth et al. *Science* '16, Jungwirth, Marti, Wadley, Wunderlich, *Nature Nanotech.* '16



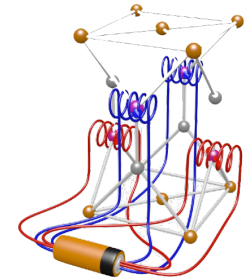
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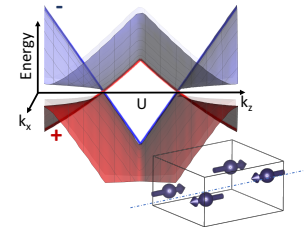
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Smejkal, Zelezny, Sinova, Jungwirth PRL (2017)

Smejkal, Jungwirth, Sinova PSS (2017)

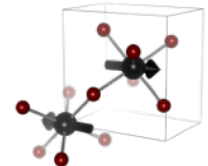


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Jungwirth, JS, et al, Nature Physics (AFM spintronics Reviews) (2018)

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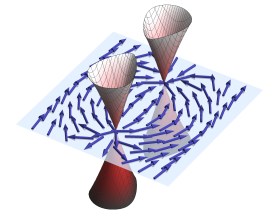
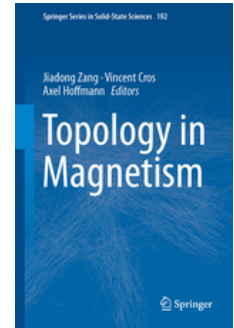


Can we control the relativistic fermions electrically?

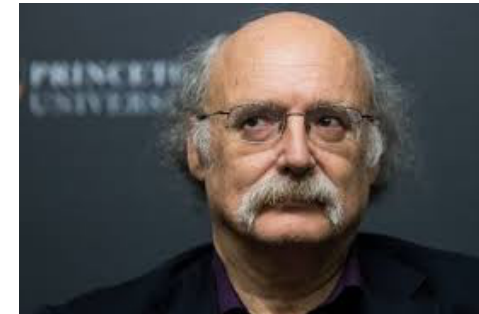


?

+



*“Just knowing the correct laws of quantum mechanics does not mean that we understand all the strange phenomena that it allows”
Haldane, Nobel Lecture 2016*



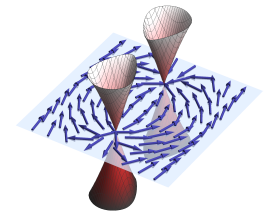
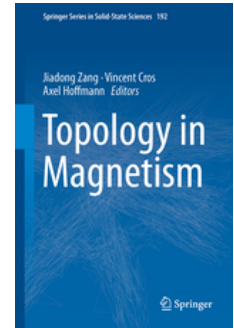
Libor Smejkal, Zelezny, Sinova, Jungwirth PRL (2017)

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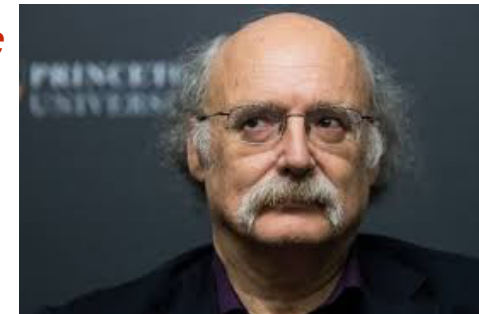


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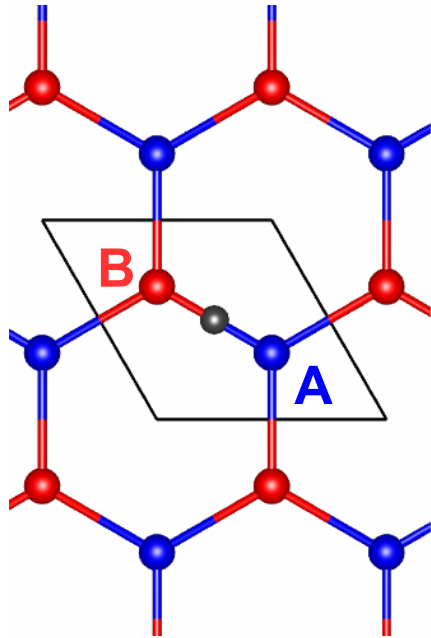
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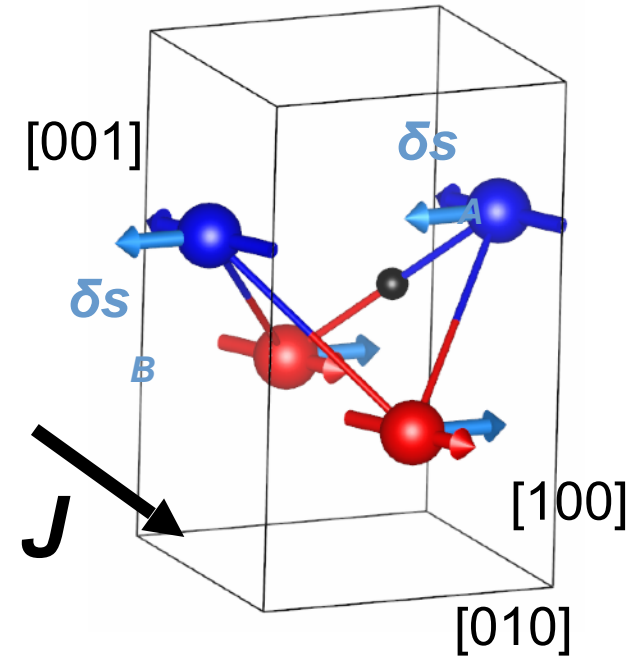
*“Just knowing the correct laws of” **relativistic**
“quantum mechanics does not mean that we
understand all the strange” **and useful**
“phenomena that it allows”*

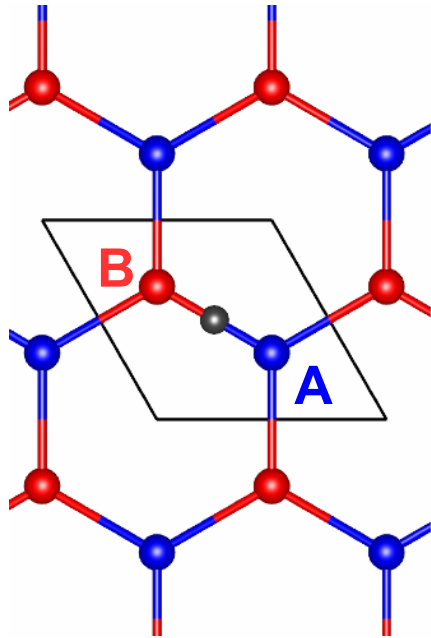


Libor Smejkal, Zelezny, Sinova, Jungwirth PRL (2017)



Dirac fermions + AF spintronics



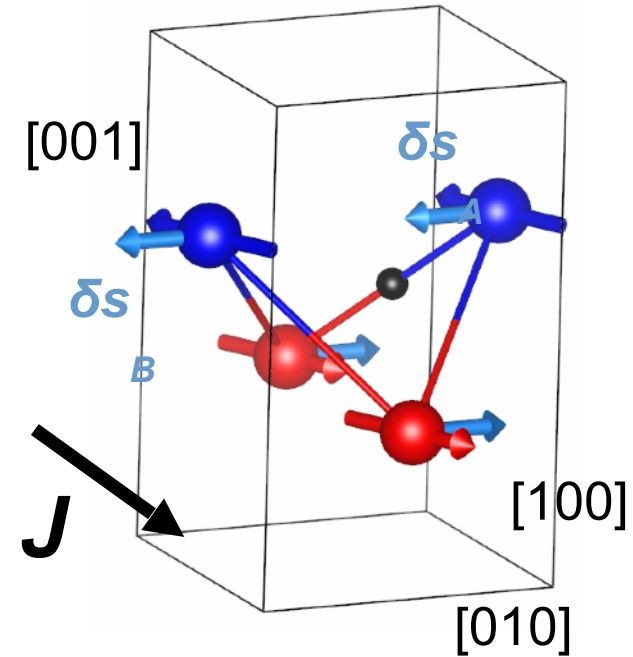


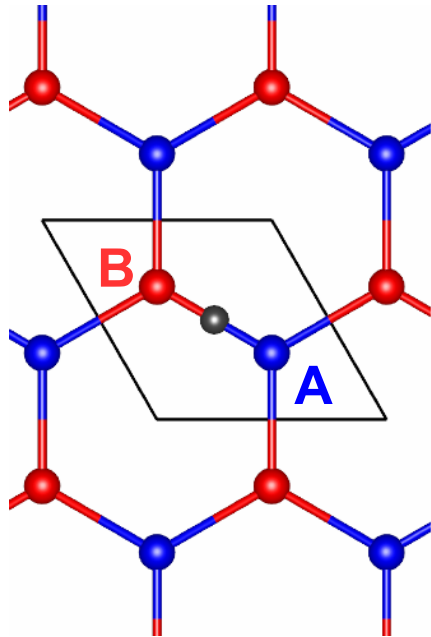
?

Dirac fermions + AF spintronics

YES!

Overlap of symmetry conditions





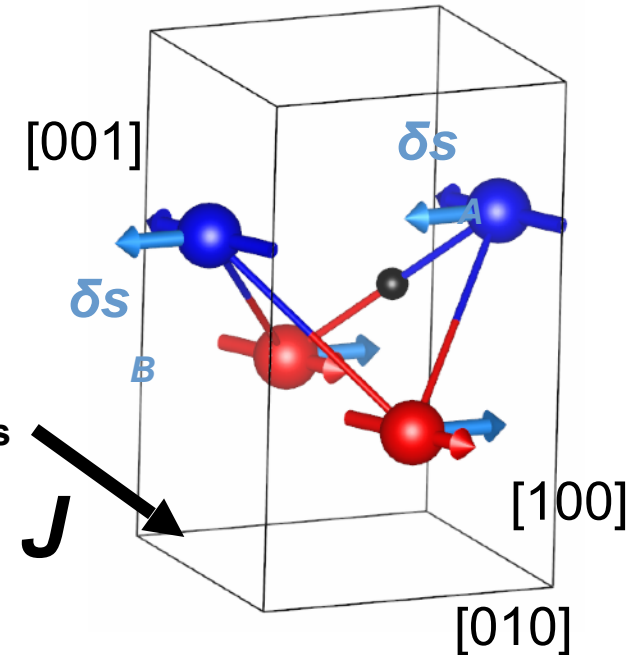
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Dirac fermions + AF spintronics
YES!

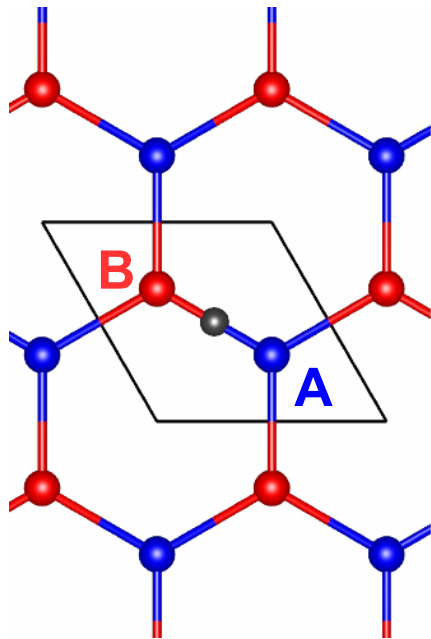
Overlap of symmetry conditions

1. Two sites in unit cell

band crossing

inversion-partner sites
→ staggered field





?
Dirac fermions + AF spintronics
YES!

Overlap of symmetry conditions

1. Two sites in unit cell

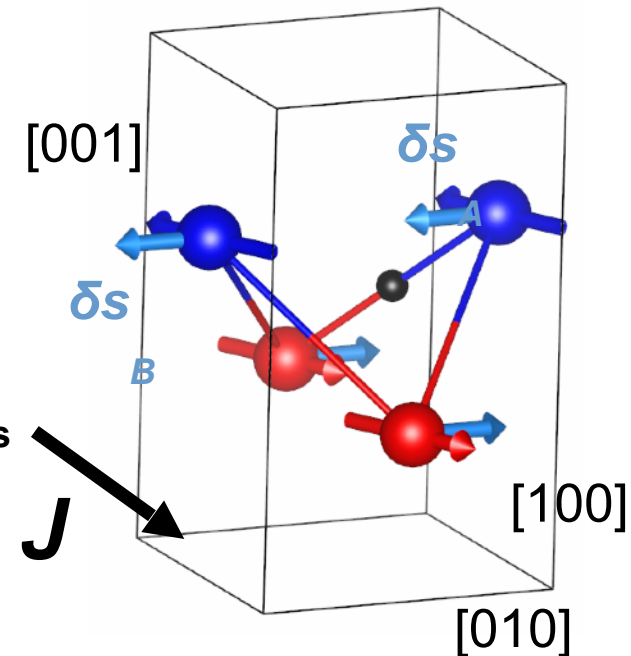
band crossing

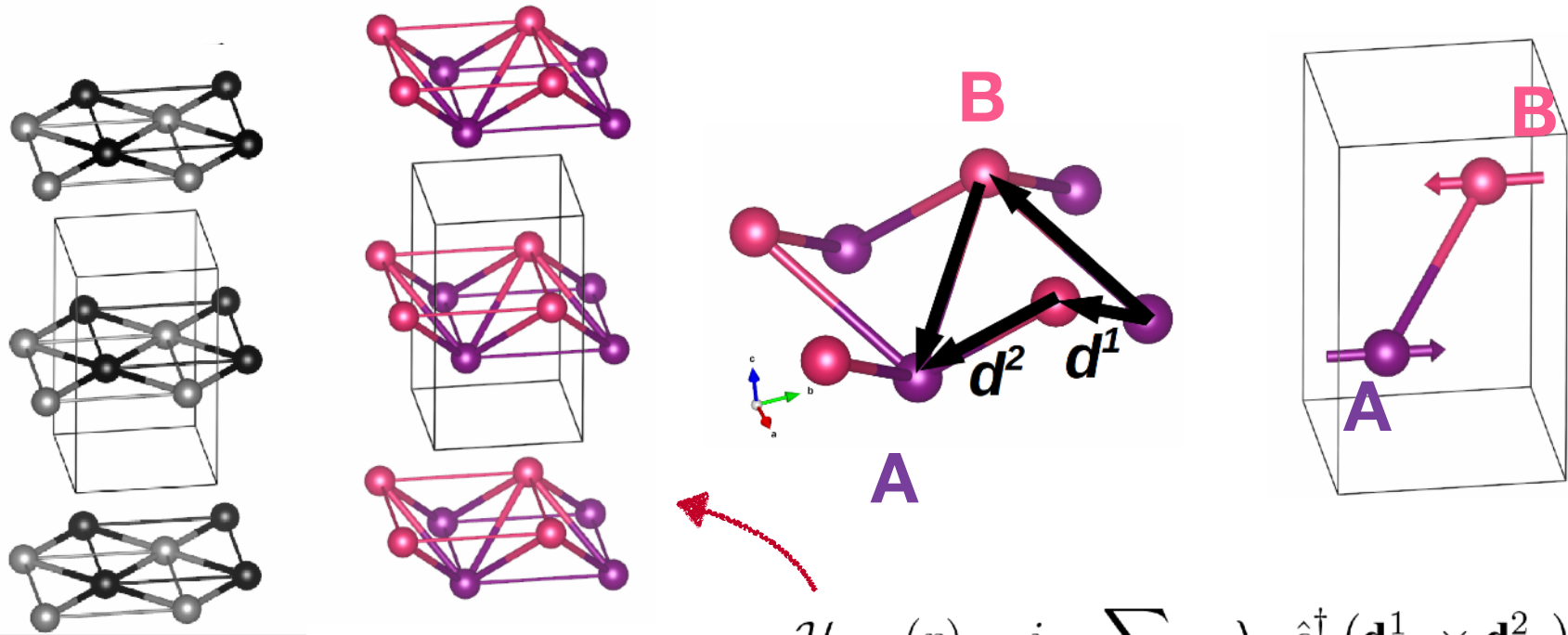
inversion-partner sites
→ staggered field

2. \mathcal{PT} symmetry

Double band degeneracy
→ Dirac point

AF spin-sublattices
at inversion partner sites





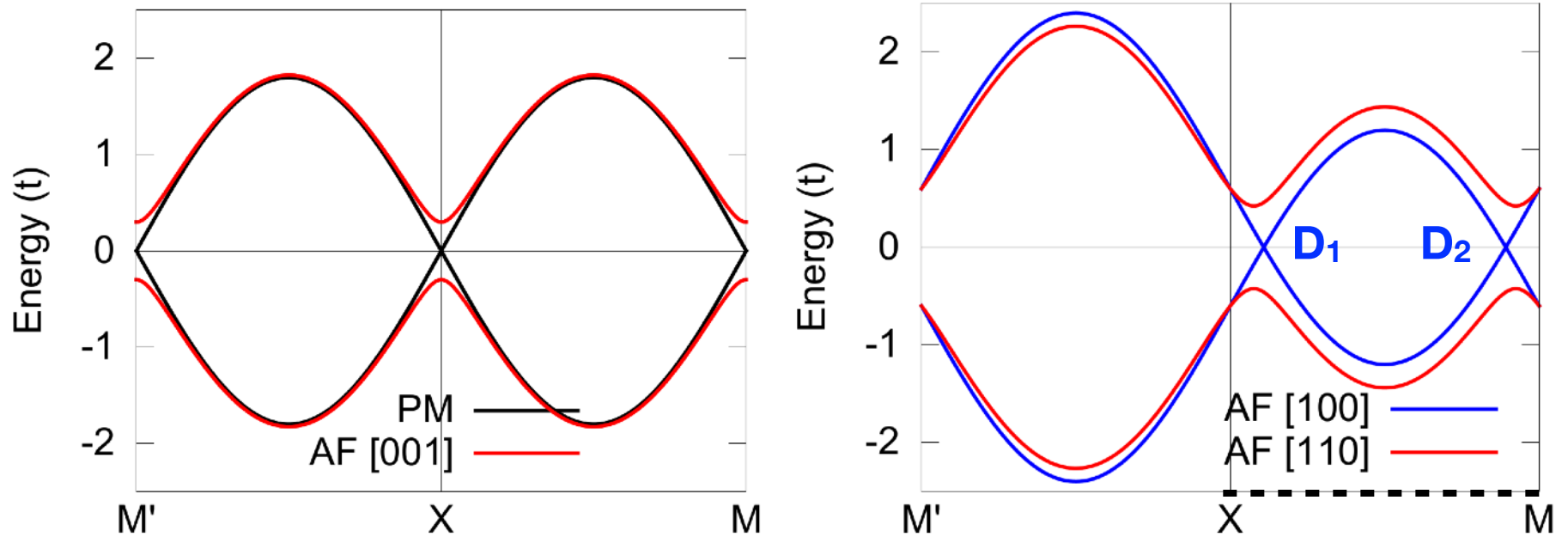
$$\mathcal{H}_{\text{SOC}}(r) = i \sum_{\langle\langle i,j \rangle\rangle, \langle k \rangle} \lambda_{ij} \hat{c}_i^\dagger (\mathbf{d}_{ik}^1 \times \mathbf{d}_{kj}^2) \cdot \boldsymbol{\sigma} \hat{c}_j$$

$$\mathcal{H} = \sum_{\langle i,j \rangle, \langle\langle i,j \rangle\rangle} t_{ij} \hat{c}_i^\dagger \hat{c}_j + \sum_i J_i \hat{c}_i^\dagger \mathbf{n} \cdot \boldsymbol{\sigma} \hat{c}_i$$

Kane-Mele spin-orbit coupling
 Young, Kane, *Phys.Rev.Lett.*(2015)

*L.Smejkal, J.Zelezny, J.Sinova, and
 T.Jungwirth, Phys.Rev.Lett. (2017)*

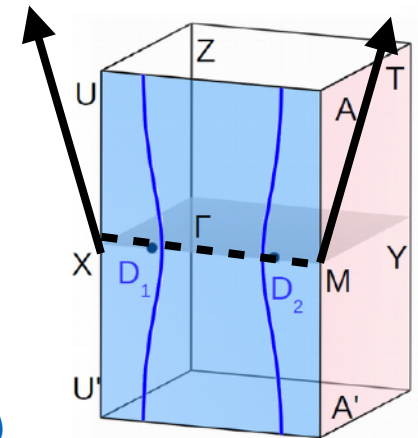
Antiferromagnetic Dirac semimetal



$$H_{\mathbf{k}} = -2t\tau_x \cos \frac{k_x}{2} \cos \frac{k_y}{2} - t' (\cos k_x + \cos k_y) + \lambda\tau_z (\sigma_y \sin k_x - \sigma_x \sin k_y) + \tau_z J_n \boldsymbol{\sigma} \cdot \mathbf{n}$$

Renormalization: 2D Dirac points \rightarrow 3D nodal lines

L.Smejkal, J.Zelezny, JS, and T.Jungwirt, Phys.Rev.Lett. (2017)

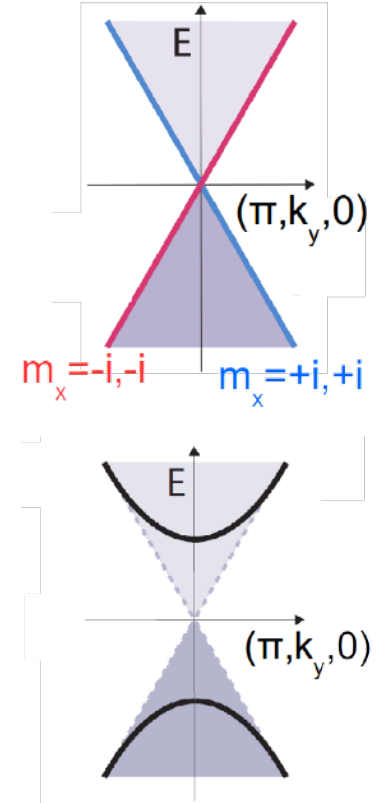
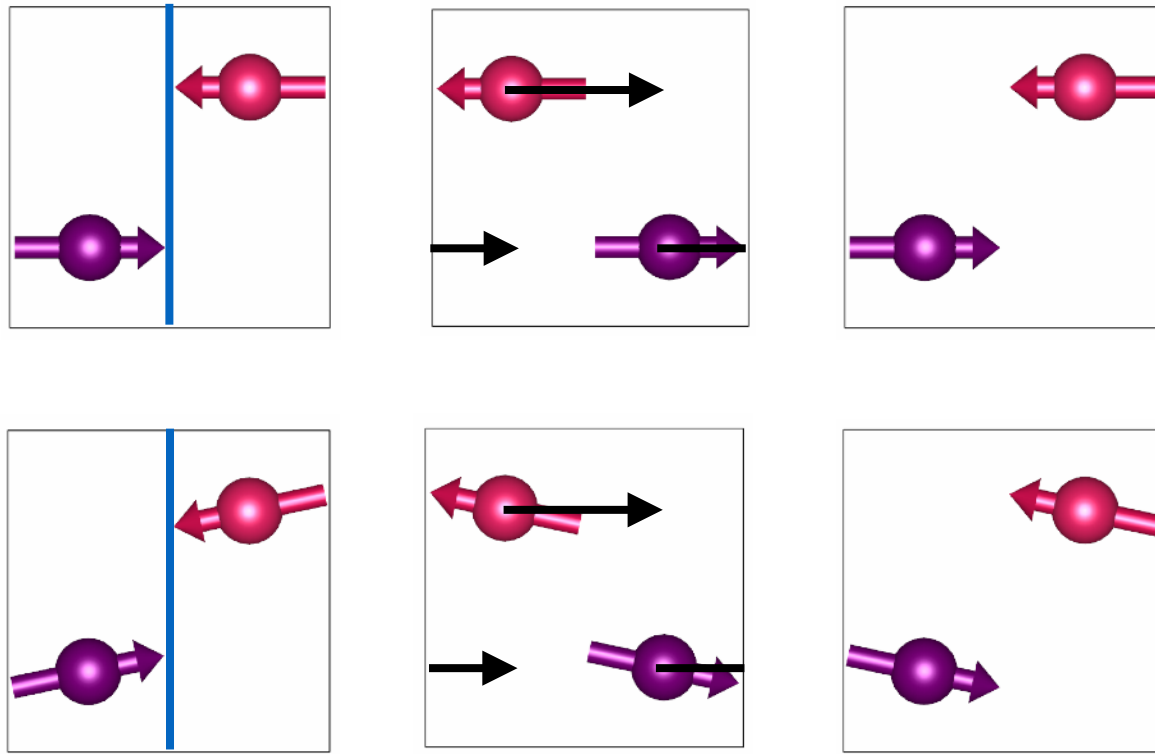


Symmetry protection of Dirac points

Mirror

Translate

at \mathcal{G}_x invariant
subspace



$m_x = -i, -i$ $m_x = +i, +i$

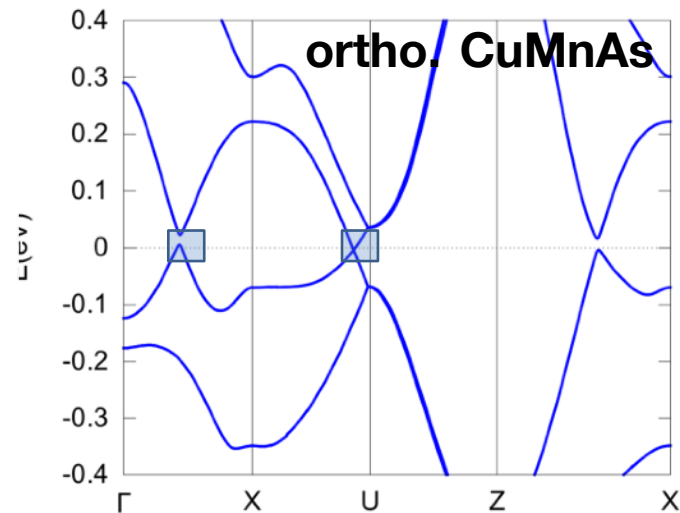
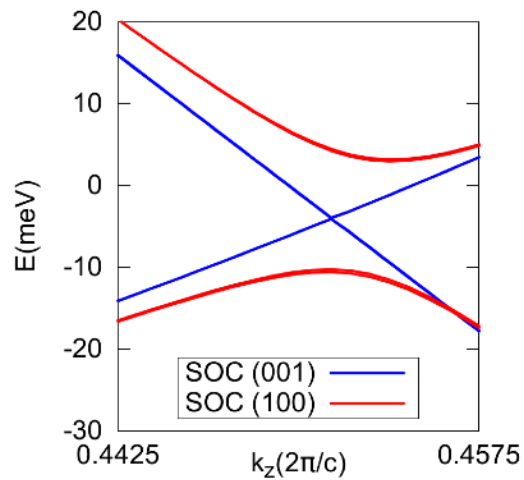
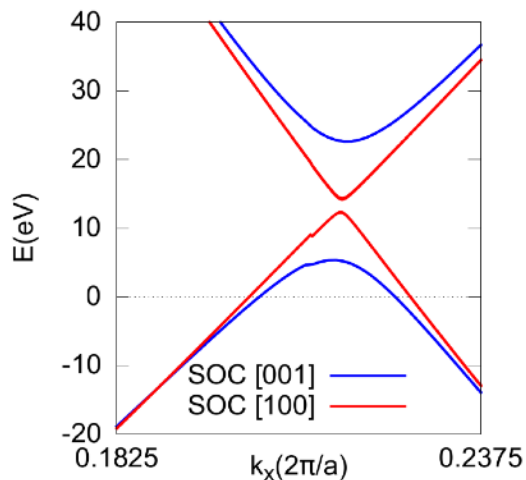
[010]

[100]

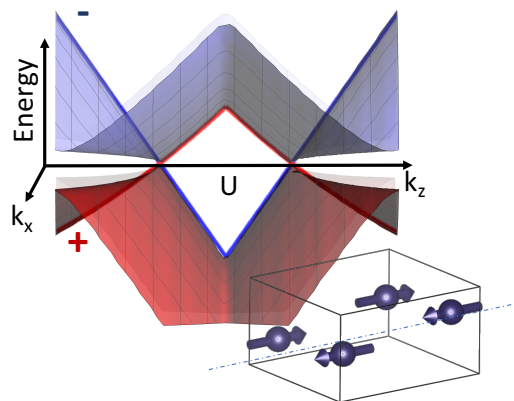
nonsymmorphic symmetry = point group + nontrivial translation

glide mirror plane $\mathcal{G}_x = \{M_x / (1/2, 0, 0)\}$ $M_x = i\sigma_x \tau_z$

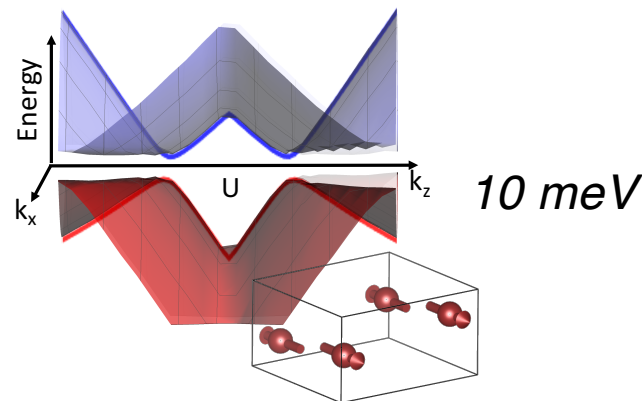




Nonsymmorphic symmetry:
Screw axis+Glide plane

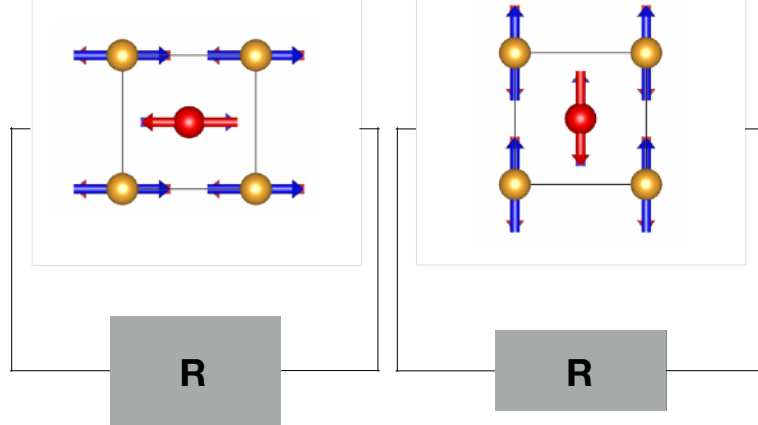


$$[001] Pn'm'a' \left\{ C_{2z} \left| \left(\frac{1}{2}, 0, \frac{1}{2} \right) \right. \right\}$$

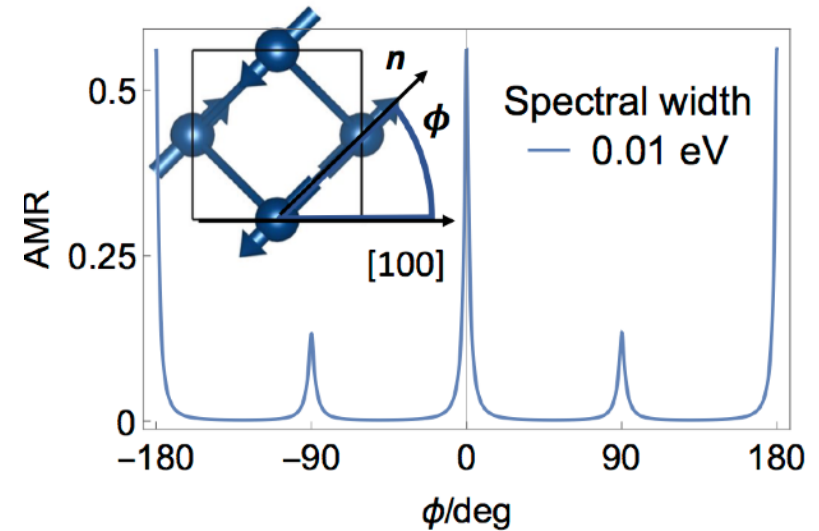


$$[100] Pnm'a \left\{ C_{2z} \left| \left(\frac{1}{2}, 0, \frac{1}{2} \right) \right. \right\}$$

Anisotropic magnetoresistance is even in magnetisation and scattering related



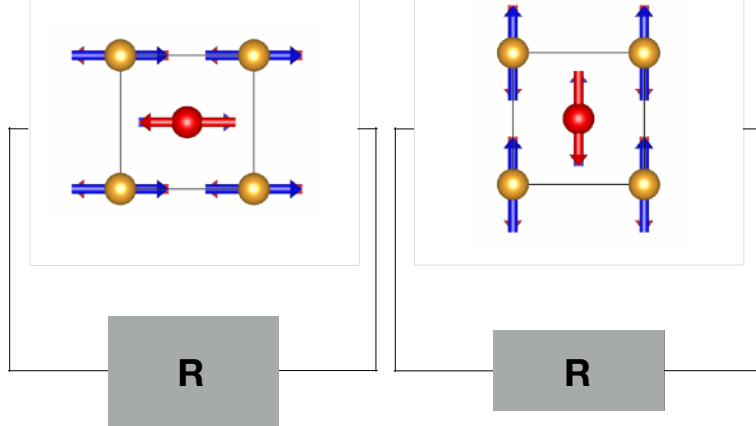
Marti, Jungwirth et al. Nat.Mat.(2014)



$$\Lambda_{100} \equiv \frac{\rho_{L,100}(\phi=0) - \rho_{L,100}(\phi=90^\circ)}{\bar{\rho}}$$

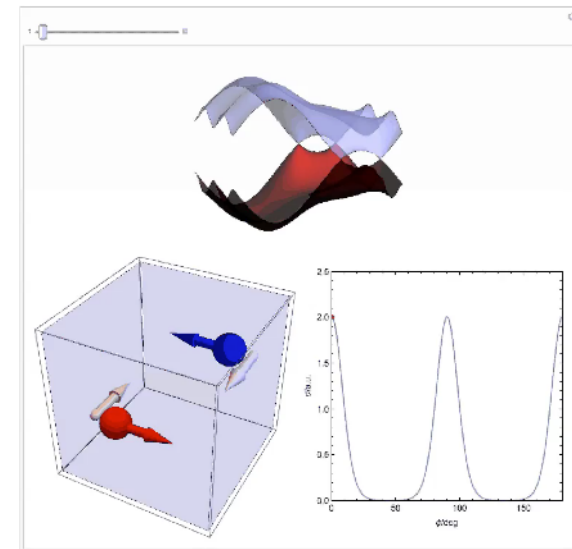
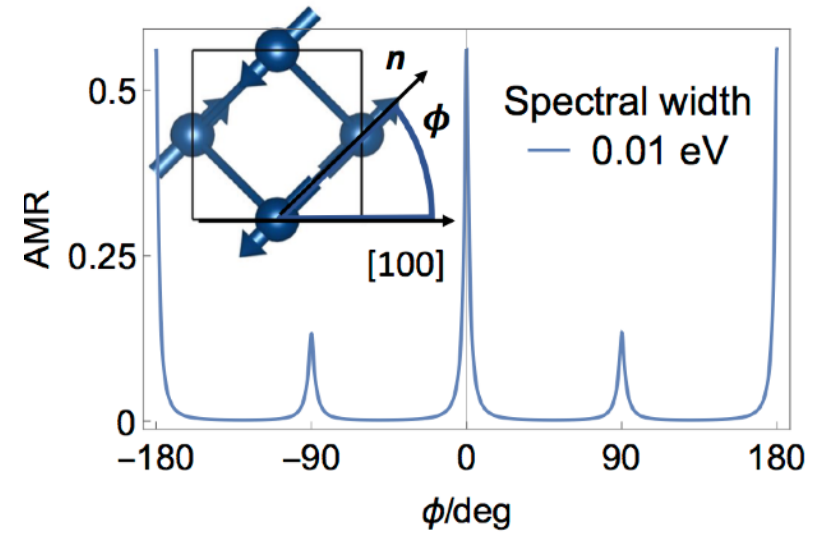


Anisotropic magnetoresistance is even in magnetisation and scattering related



Marti, Jungwirth et al. Nat.Mat.(2014)

$$\Lambda_{100} \equiv \frac{\rho_{L,100}(\phi=0) - \rho_{L,100}(\phi=90^\circ)}{\bar{\rho}}$$



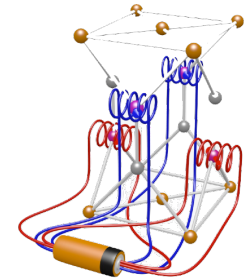
Topological antiferromagnetic spintronics and the crystal Hall effect

I. Néel Spin-Orbit Torques in Antiferromagnets:

Zelezny, Gao, JS, Jungwirth PRL (2014)

Zelazny, Gao, et al. PRB (2016)

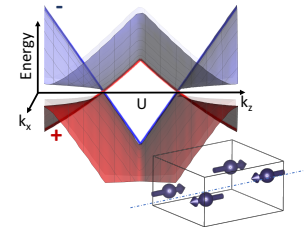
Bodnar, Smejkal, Jourdan, Kläui, Jungwirth, JS, et al, Nat. Com. (2018)



II. Topological Dirac Fermion + Antiferromagnets + Neel SOTs

Smejkal, Zelezny, Sinova, Jungwirth PRL (2017)

Smejkal, Jungwirth, Sinova PSS (2017)

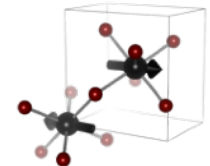


III. AHE in collinear AFMs: Crystal Hall Effect

Jungwirth, JS, et al, Nature Physics (AFM spintronics Reviews) (2018)

Surprises of the Spin Hall Effect, Physics Today 70, 7, 38 (2017)

Smejkal, Gonzales, Jungwirth, JS, arXiv 1901.00445 (2019)





Physica
Volume 21, Issues 6-10, 1955, Pages 877-887



The spontaneous hall effect in ferromagnetics

I

J. Smit

Show more

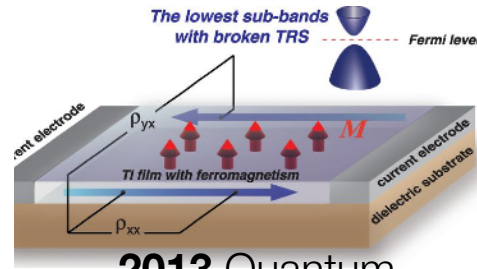
[https://doi.org/10.1016/S0031-8914\(55\)92596-9](https://doi.org/10.1016/S0031-8914(55)92596-9)

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Synopsis

Apart from the normal Hall voltage a magnetized ferromagnetic material usually shows a relatively large extra voltage in the same direction, which can be found by linear extrapolation to $B=0$. It is shown that this spontaneous Hall effect cannot exist in a perfectly periodic lattice. Measurements at different temperatures suggest that the effect is closely related with the electrical-resistivity ρ of the material.

Existing theories on the origin of the effect are shown to be invalid, and it is shown that the explanation has to be based on the anisotropic scattering, caused by spinorbit interaction, of the conducting electrons against the imperfections of the lattice.



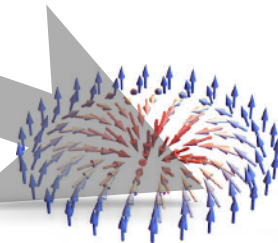
2013 Quantum anomalous Hall in topological insulators

1988 Haldane: quantum Hall effect **without Landau levels**

2002 Berry curvature - **topological properties of wave functions**

1980 Quantum Hall effect

1996 "Topological" Hall effect **without spin-orbit coupling**



Spontaneous Hall effect without complicated antiferromagnetic order?

2015 Anomalous Hall effect in non collinear antiferromagnets **without magnetisation**

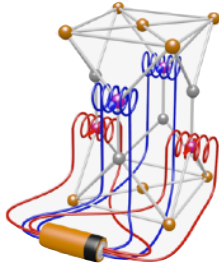
1955+ Role of spin-orbit coupled impurities

1953/58+ Kohn-Luttinger

1879 Hall effect from field

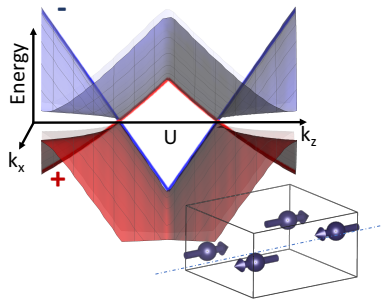
1881 Hall effect from magnetisation





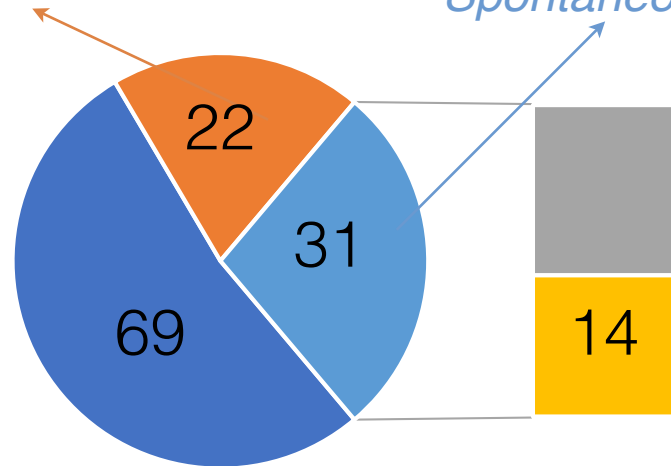
magnetic symmetries: 122/1651

Dirac antiferromagnets



Šmejkal, Zelezny, JS, TJ,
Phys.Rev.Lett. (2017)

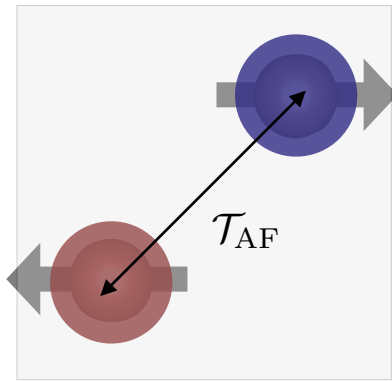
Spontaneous Hall antiferromagnets



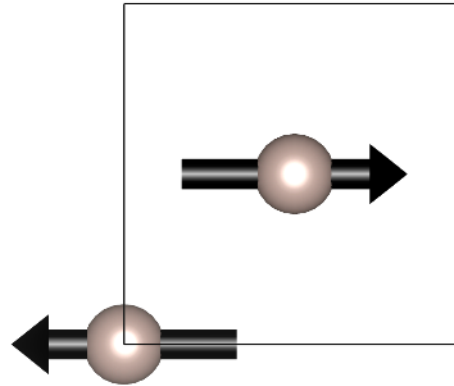
**collinear
antiferromagnets**

spin degenerate bands

(IrMn, CuMnAs, Mn₂Au, ...)

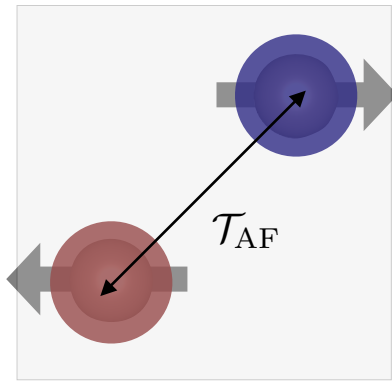


*LŠ, J.Železny, JS, TJ,
Phys.Rev.Lett. (2017)*

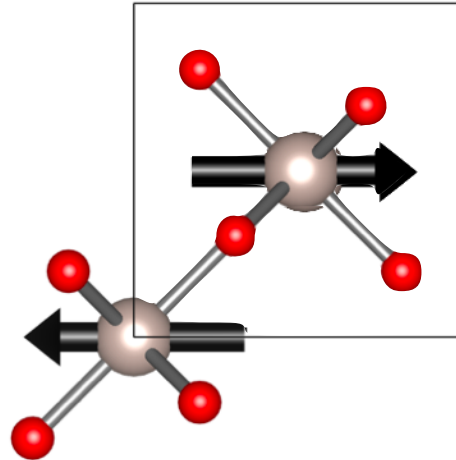


spin degenerate bands

(IrMn, CuMnAs, Mn₂Au, ...)

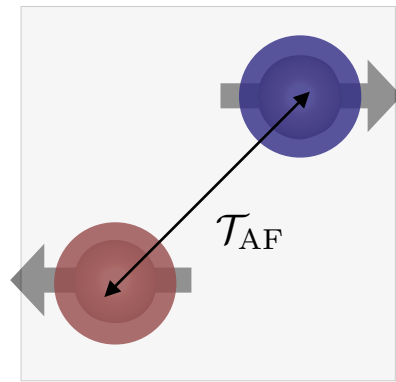


*LŠ, J.Železny, JS, TJ,
Phys.Rev.Lett. (2017)*

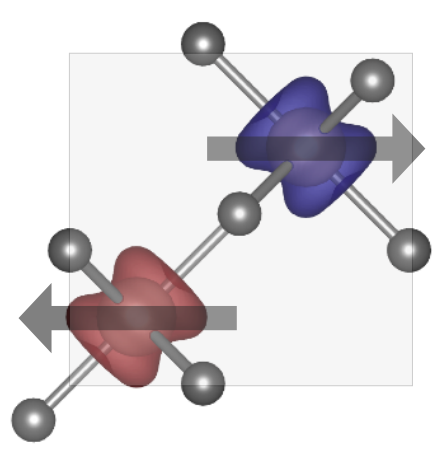


spin degenerate bands

(IrMn, CuMnAs, Mn₂Au, ...)

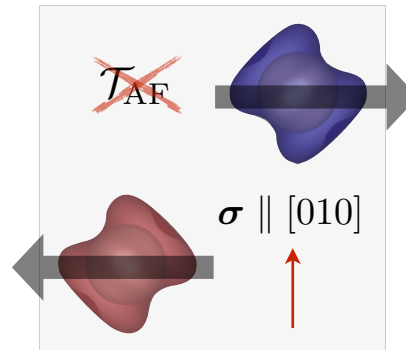
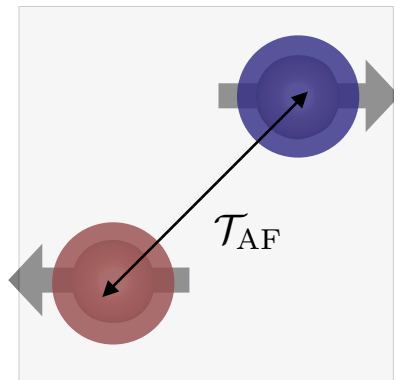


*LŠ, J.Železny, JS, TJ,
Phys.Rev.Lett. (2017)*



spin degenerate bands

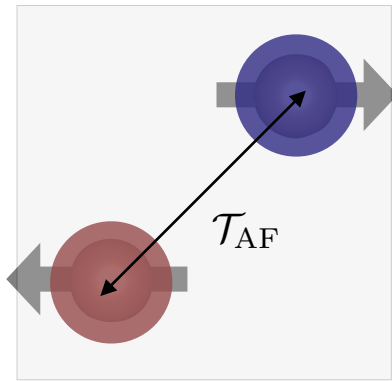
(IrMn, CuMnAs, Mn₂Au, ...)



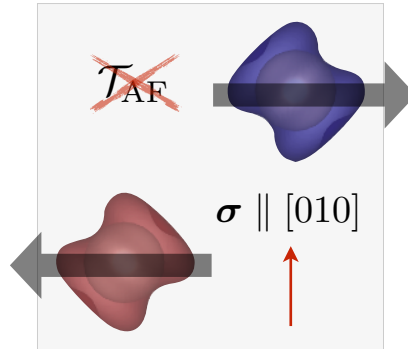
*LŠ, J.Železny, JS, TJ,
Phys.Rev.Lett. (2017)*

Magnetisation density spontaneously breaks symmetries and allows Hall vector! (RuO₂, ...)

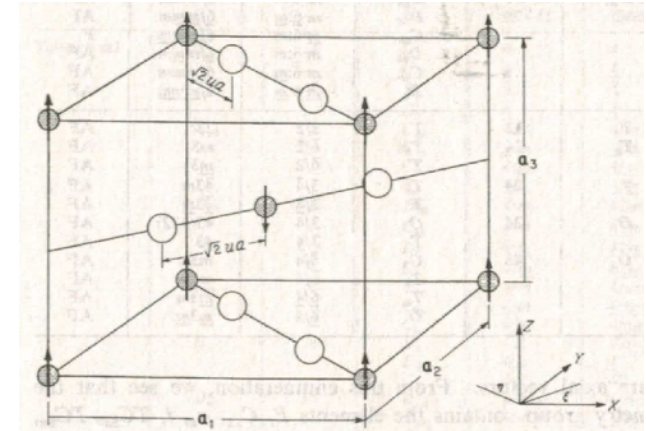
spin degenerate bands
(IrMn, CuMnAs, Mn₂Au, ...)



LŠ, J.Železny, JS, TJ,
Phys.Rev.Lett. (2017)



LŠ, R.H.Gonzales, JS, TJ,
arXiv 1901.00445 (2019)



**Notorious
antiferromagnets**

**Analysing magnetisation projection
vectors (black arrows) is incomplete!**

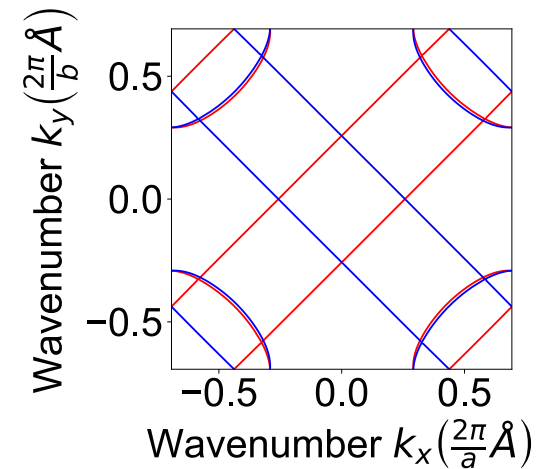
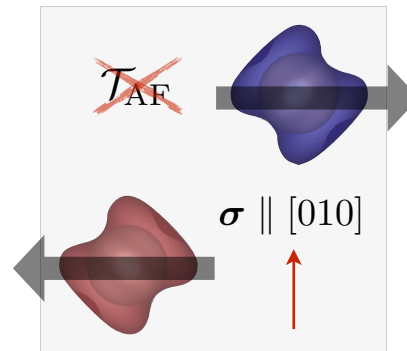
Landau, Lifshitz, Electrodynamics of Continuous Media (1972)

*Bradley & Cracknell, The Mathematical
Theory of Symmetry in Solids* (1972)

*Tinkham, Group theory and quantum
mechanics* (1964)

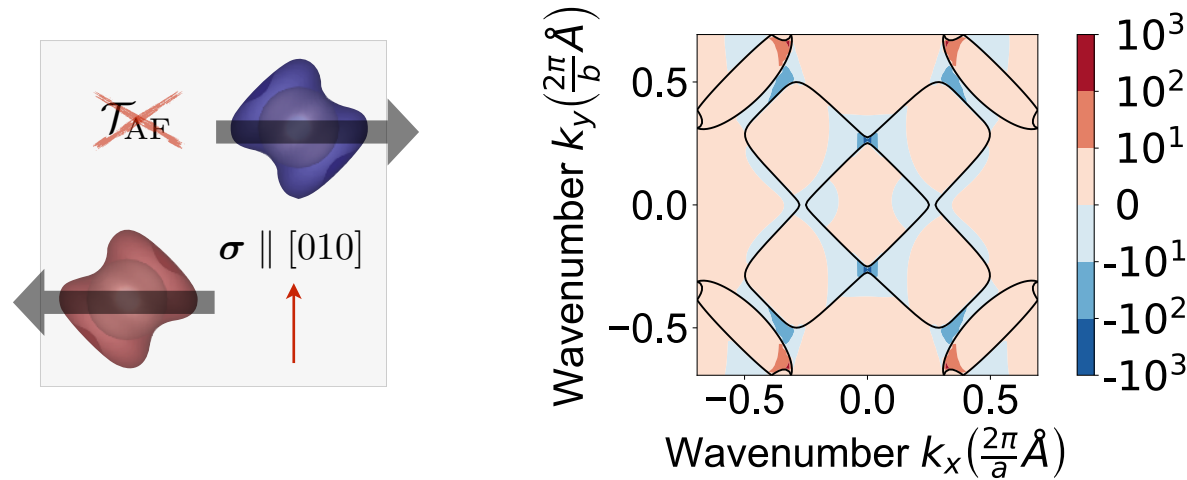


Spontaneous symmetry breaking is strong already without SOC!



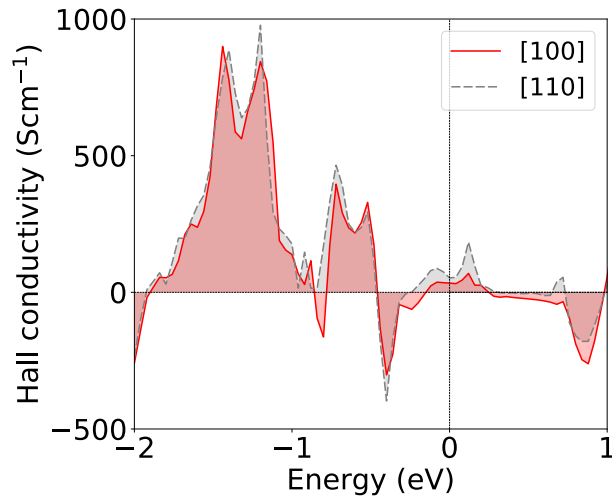
LŠ, R.H.Gonzales, JS, TJ, arXiv 1901.00445 (2019)

Spontaneous symmetry breaking is strong already without SOC!

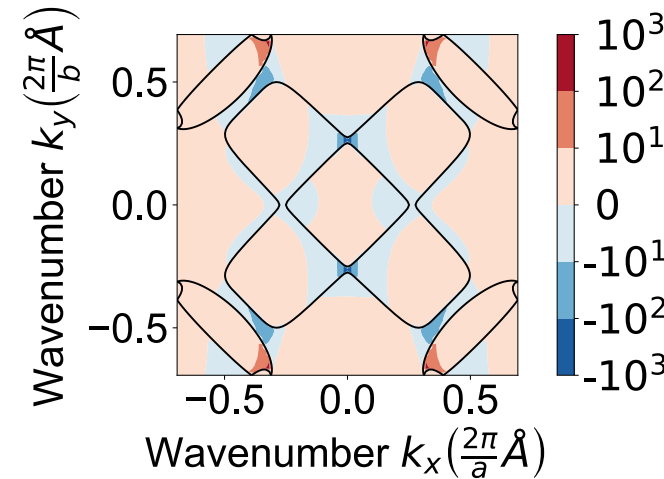
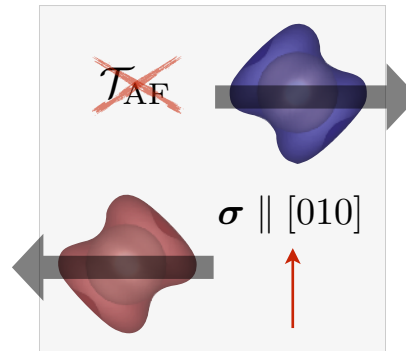


LŠ, R.H.Gonzales, JS, TJ, arXiv 1901.00445 (2019)

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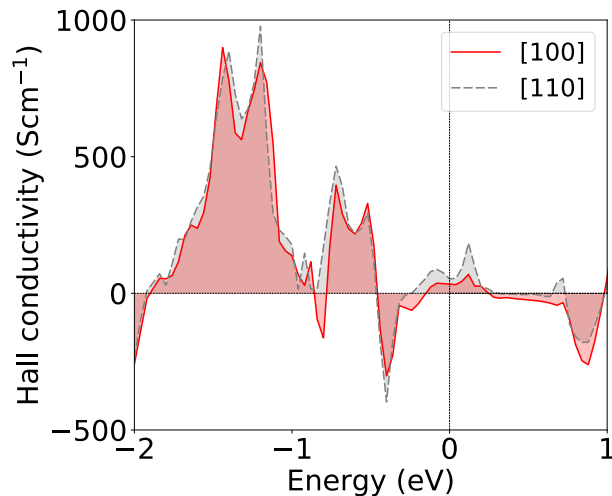


$$\sigma_{xy} = \frac{e^2}{h} \int_{\text{BZ}} \frac{d\mathbf{k}}{(2\pi)^3} \sum_n f(\mathbf{k}) b_z^n(\mathbf{k})$$

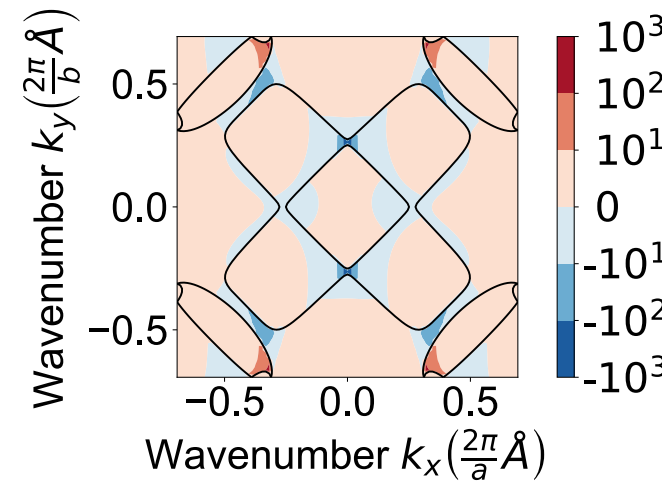
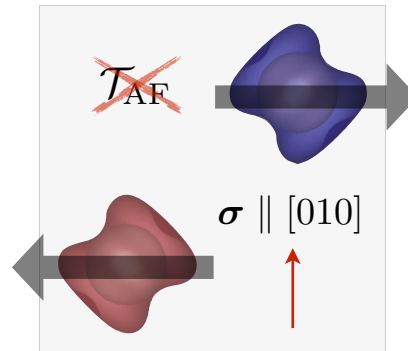


LŠ, R.H.Gonzales, JS, TJ, arXiv 1901.00445 (2019)

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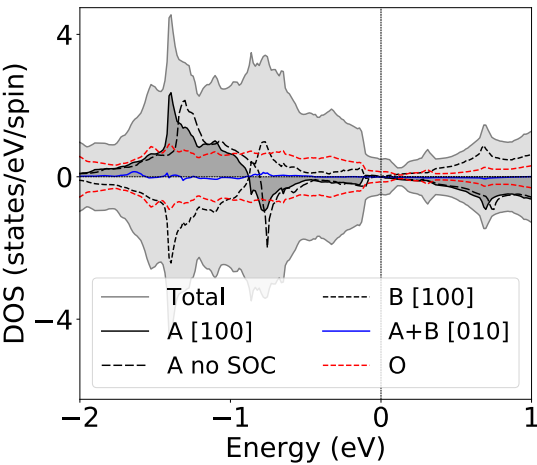


$$\sigma_{xy} = \frac{e^2}{h} \int_{\text{BZ}} \frac{dk}{(2\pi)^3} \sum_n f(\mathbf{k}) b_z^n(\mathbf{k})$$



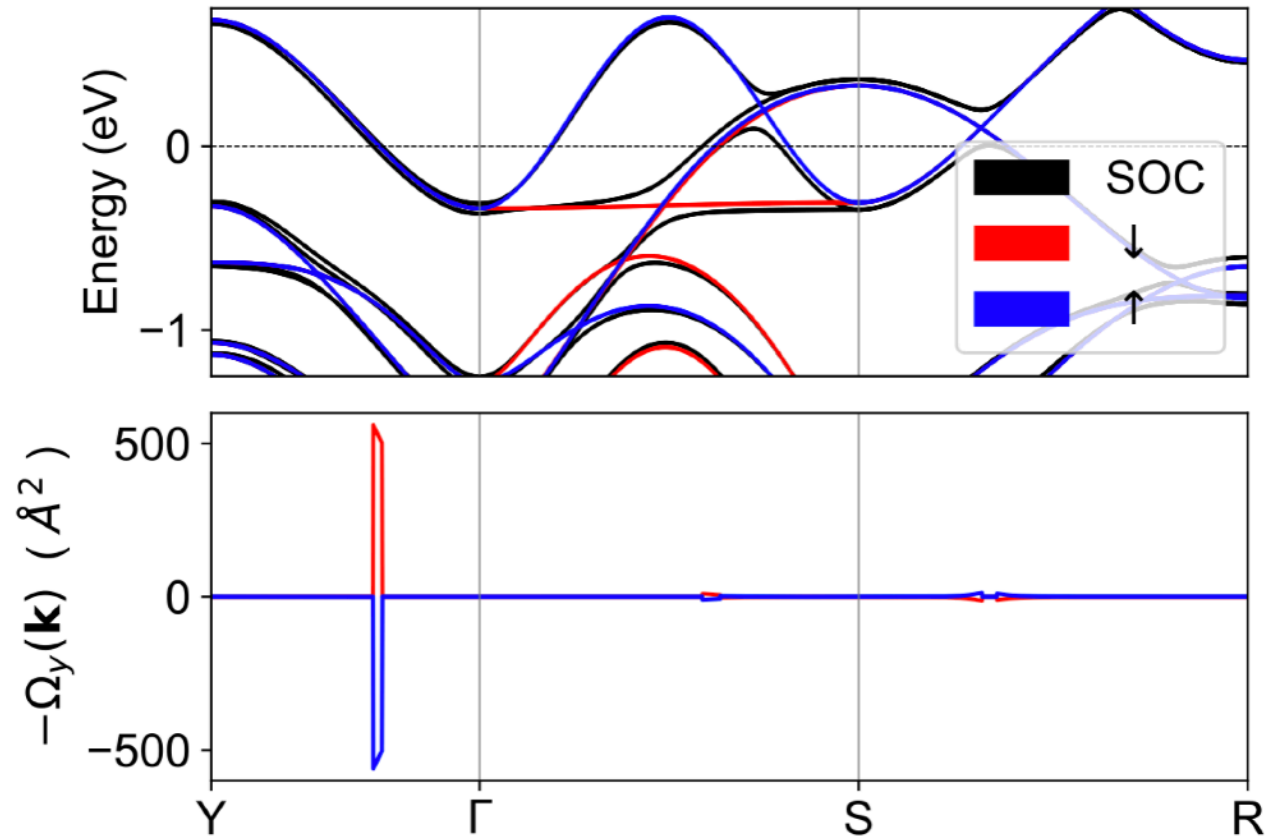
LŠ, R.H.Gonzales, JS, TJ, arXiv 1901.00445 (2019)

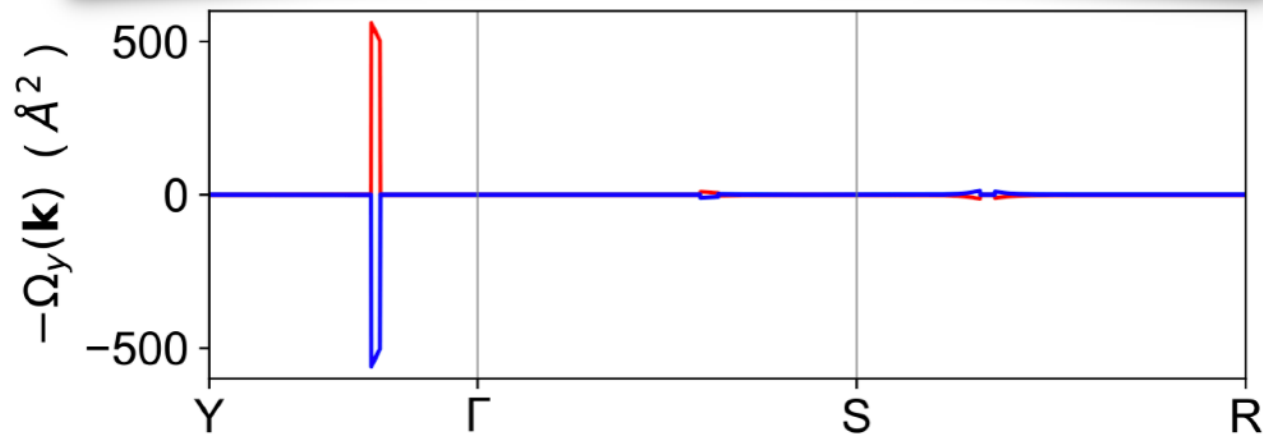
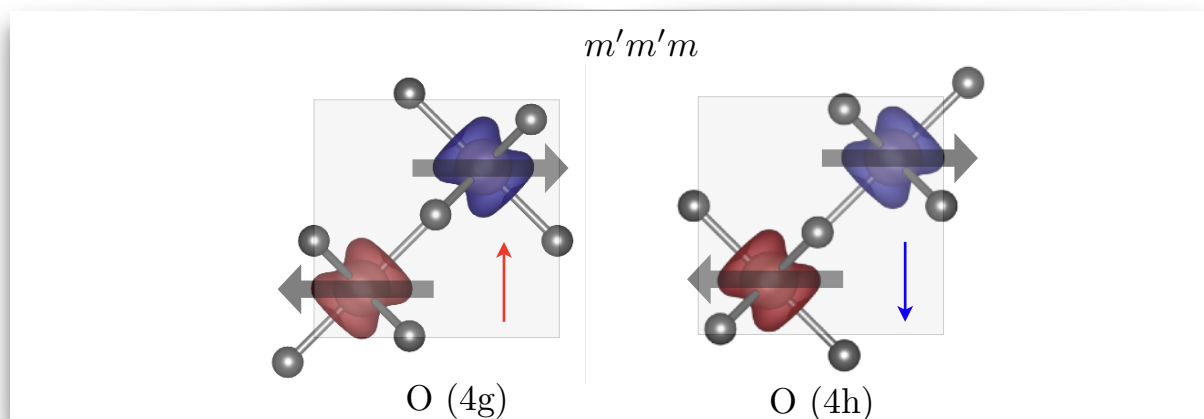
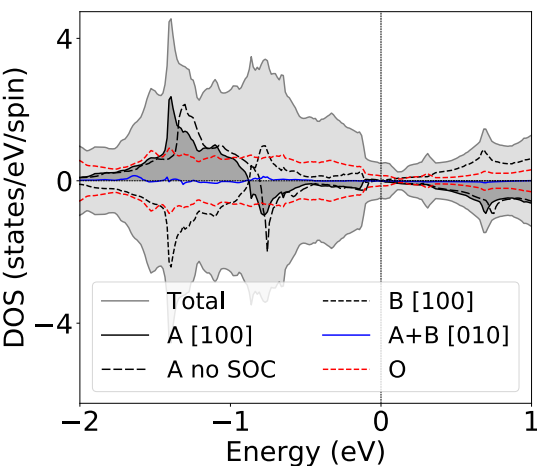
Magnitude can be as large as record values in ferromagnets



Correlations stabilise antiferromagnetism

Berlijn et al., Phys. Rev. Lett. (2017)

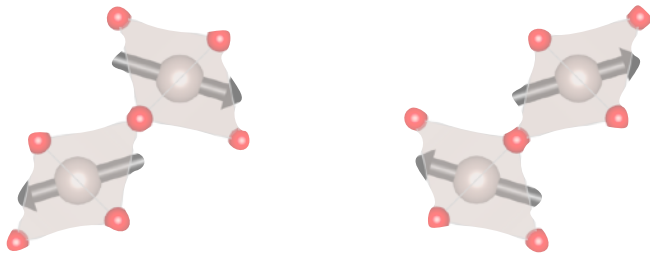




Correlations stabilise antiferromagnetism

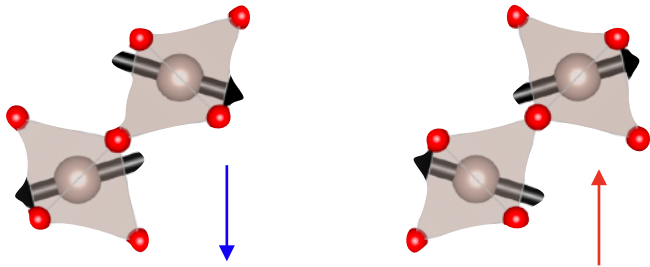
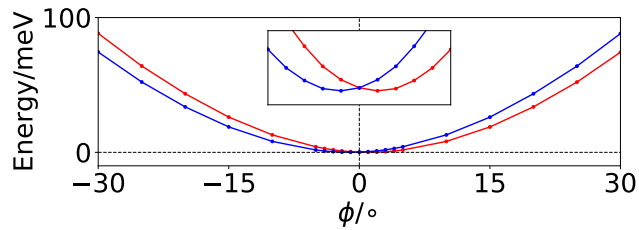
Berlijn et al., *Phys. Rev. Lett.* (2017)

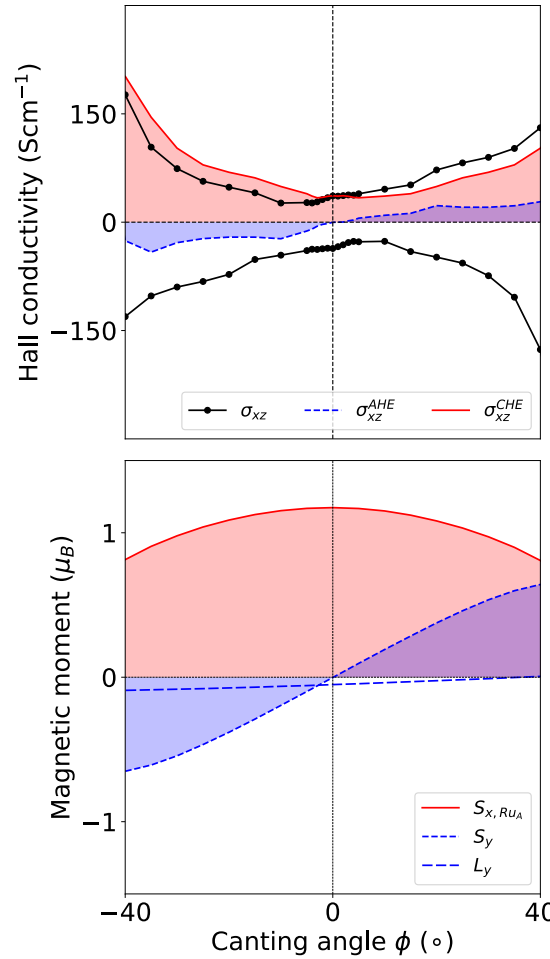
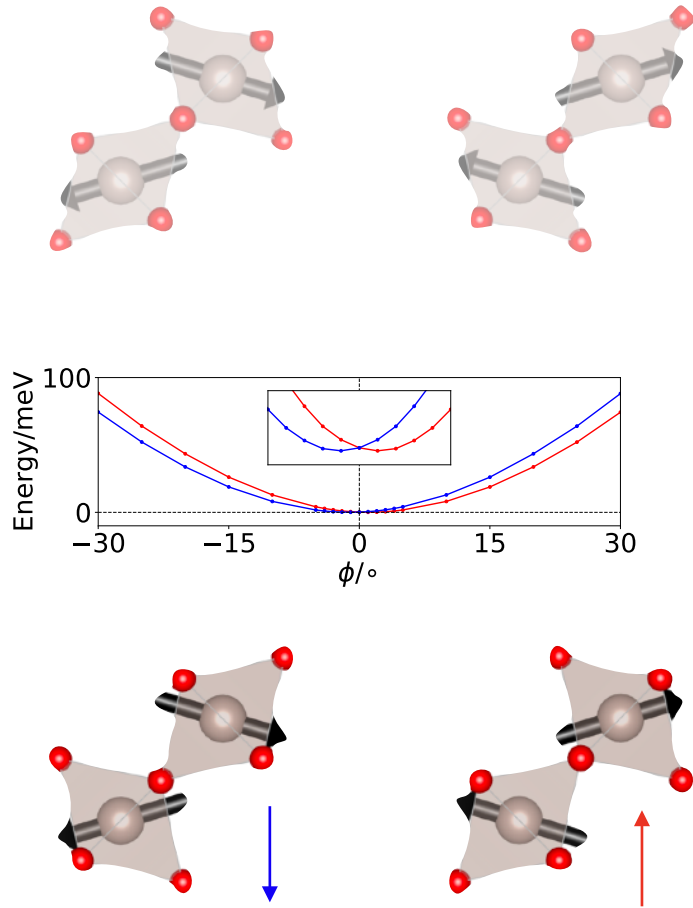




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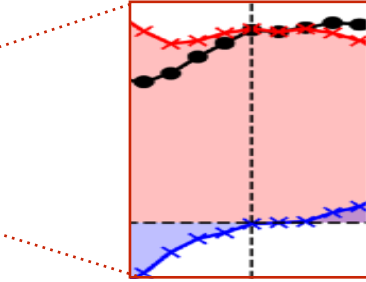
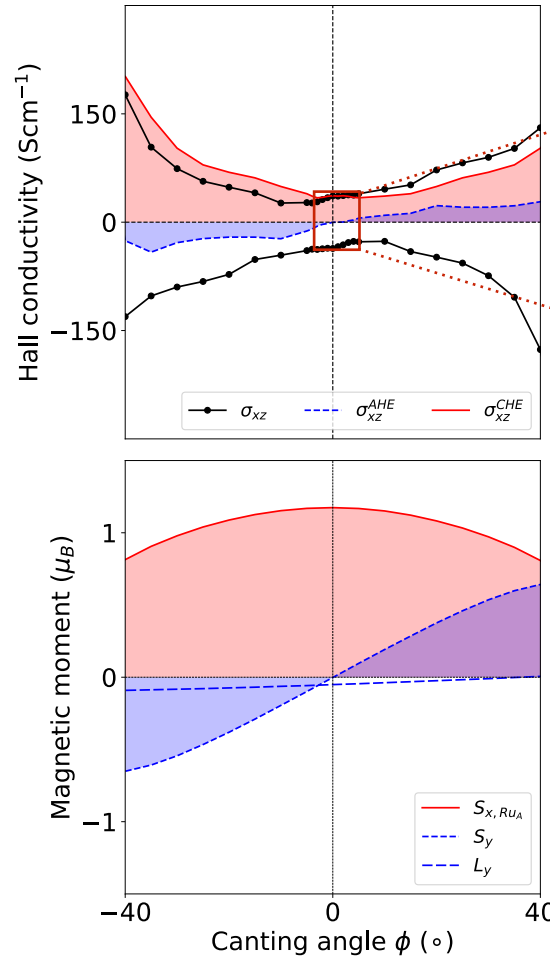
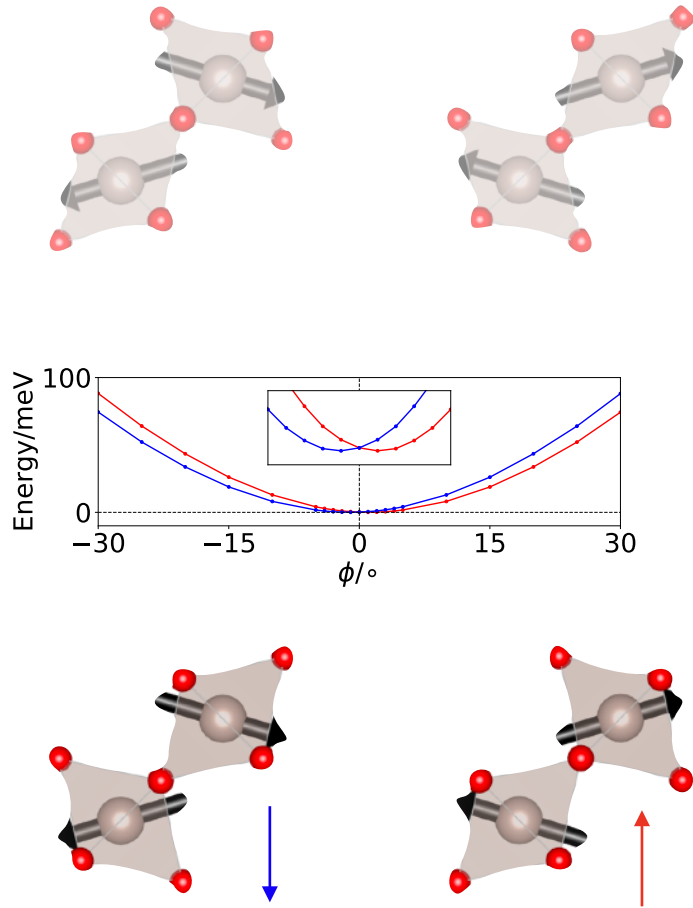




$$\sigma_{xz}^{\text{CH}} = \frac{\sigma_{xz}(\mathbf{n}, \mathbf{m}) + \sigma_{xz}(\mathbf{n}, -\mathbf{m})}{2}$$

$$\sigma_{xz}^{\text{AH}} = \frac{\sigma_{xz}(\mathbf{n}, \mathbf{m}) - \sigma_{xz}(\mathbf{n}, -\mathbf{m})}{2}$$

$$\sigma_{xz}(\mathbf{n}, \mathbf{m}) = -\sigma_{xz}(-\mathbf{n}, -\mathbf{m})$$



$$\sigma_{xz}^{\text{CH}} = \frac{\sigma_{xz}(\mathbf{n}, \mathbf{m}) + \sigma_{xz}(\mathbf{n}, -\mathbf{m})}{2}$$

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$$\sigma_{xz}(\mathbf{n}, \mathbf{m}) = -\sigma_{xz}(-\mathbf{n}, -\mathbf{m})$$

Anisotropic Crystal Hall effect

n	MSG	m (μ_B)	σ_H (Scm^{-1})
[001]	$P4_2'/mnm'$	0	0
[100]	$Pnn'm'$	$\parallel [010]$ 0.05	36
[110]	$Cmm'm'$	$\parallel [\bar{1}10]$ 0.0075	58



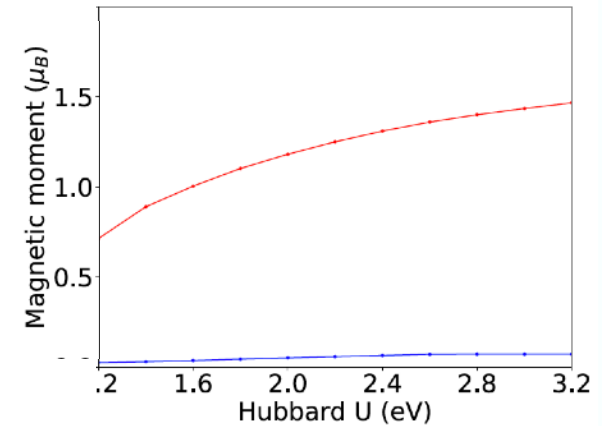
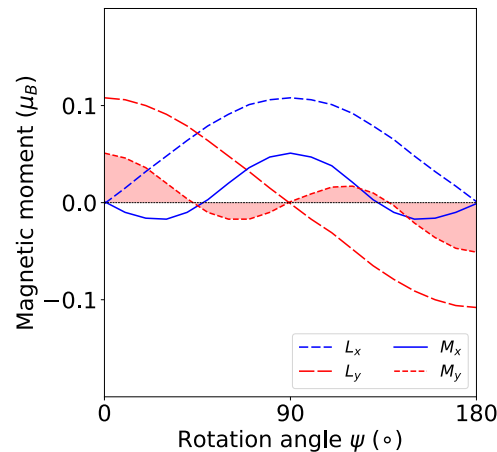
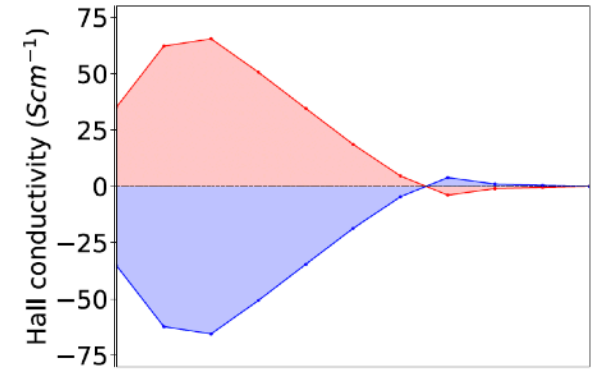
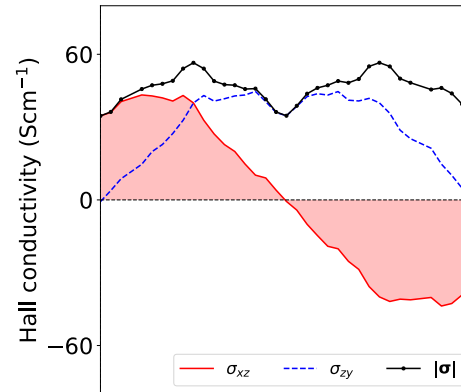
Anisotropic Crystal Hall effect

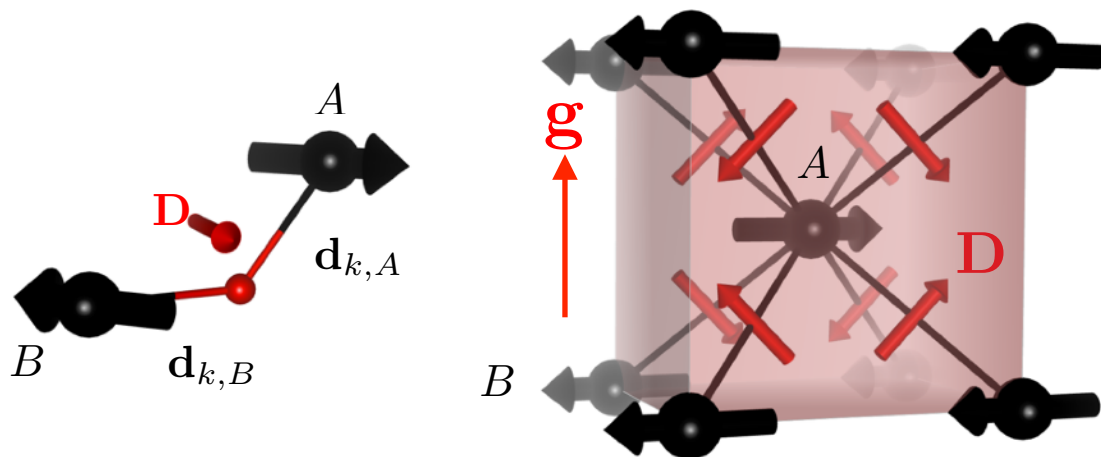
n	MSG	m (μ_B)	σ_H (Scm $^{-1}$)
[001]	$P4'_2/mnm'$	0	0
[100]	$Pnn'm'$	$\parallel [010]$ 0.05	36
[110]	$Cmm'm'$	$\parallel [\bar{1}10]$ 0.0075	58



\mathbf{n}	MSG	\mathbf{m} (μ_B)	σ_{II} (Scm^{-1})
[001]	$P4'_2/mnm'$	0	0
[100]	$Pnn'm'$	$\parallel [010]$ 0.05	36
[110]	$Cmm'm'$	$\parallel [\bar{1}10]$ 0.0075	58

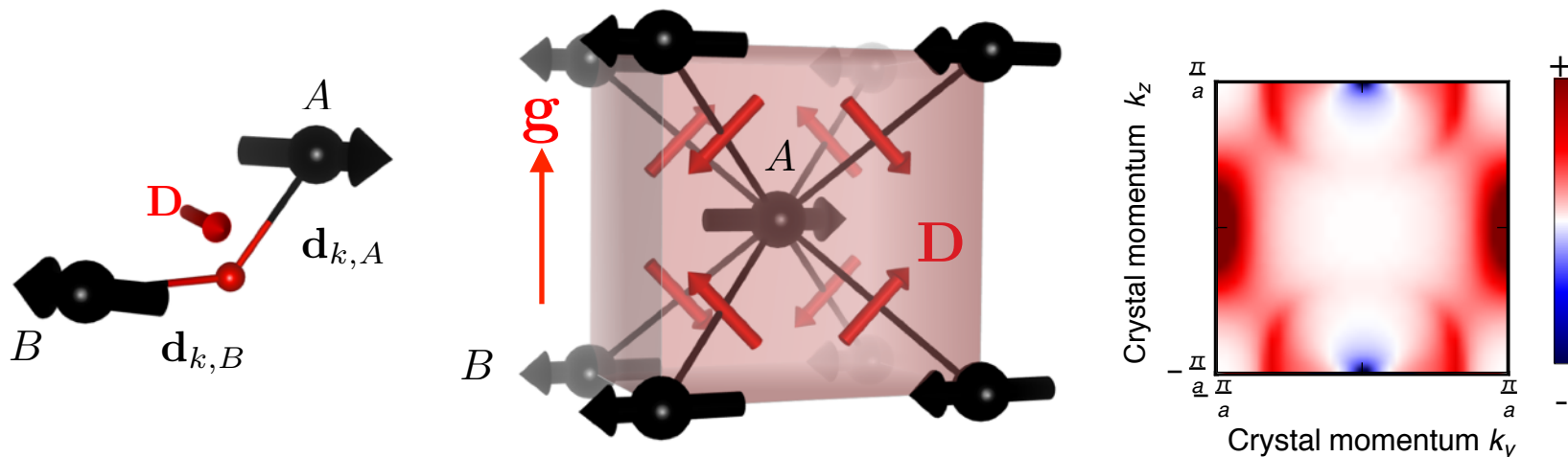
**Hall vector is not
proportional to
Sublattice magnetisation
Net magnetisation**





$$H = t \sum_{ij} c_i^\dagger c_j + J_n \sum_i \mathbf{u}_i \cdot \mathbf{s} c_i^\dagger c_i + \lambda \sum_{ij,k} \mathbf{D}_{ij}^k \cdot \mathbf{s} c_i^\dagger c_j$$

asymmetric spin-orbit coupling

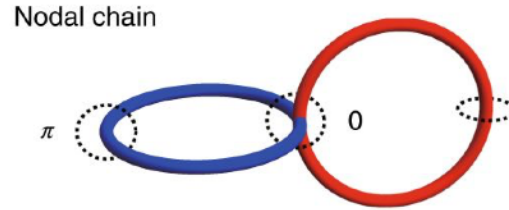


$$H = t \sum_{ij} c_i^\dagger c_j + J_n \sum_i \mathbf{u}_i \cdot \mathbf{s} c_i^\dagger c_i + \lambda \sum_{ij,k} \mathbf{D}_{ij}^k \cdot \mathbf{s} c_i^\dagger c_j$$

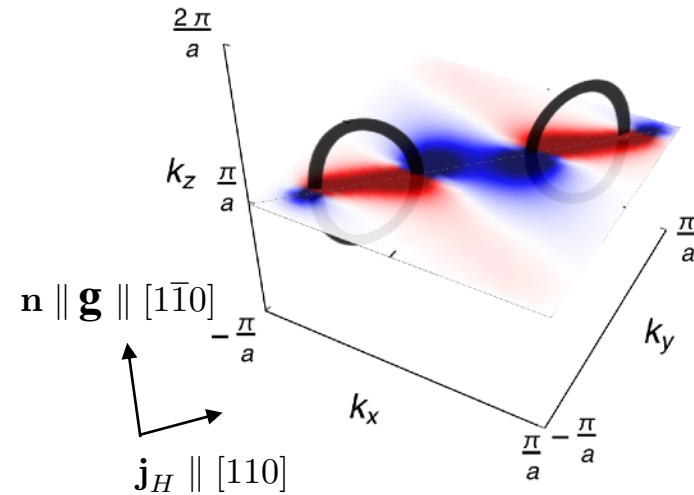
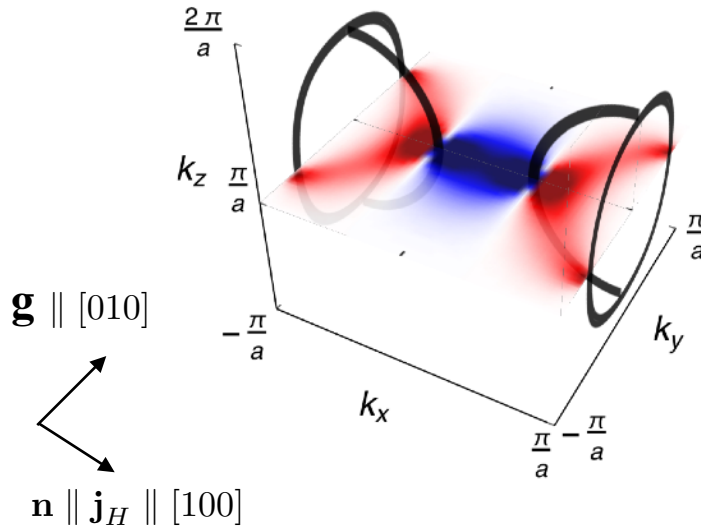
$$H_{\mathbf{k}} = -4t\tau_x \cos \frac{k_x}{2} \cos \frac{k_y}{2} \cos \frac{k_z}{2} + \tau_z J_n \boldsymbol{\sigma} \cdot \mathbf{n} + 4i\lambda \sin \frac{k_z}{2} \left[\sigma_{xy}^{(-)} \sin \frac{k_x + k_y}{2} + \sigma_{xy}^{(+)} \sin \frac{k_x - k_y}{2} \right]$$

asymmetric spin-orbit coupling

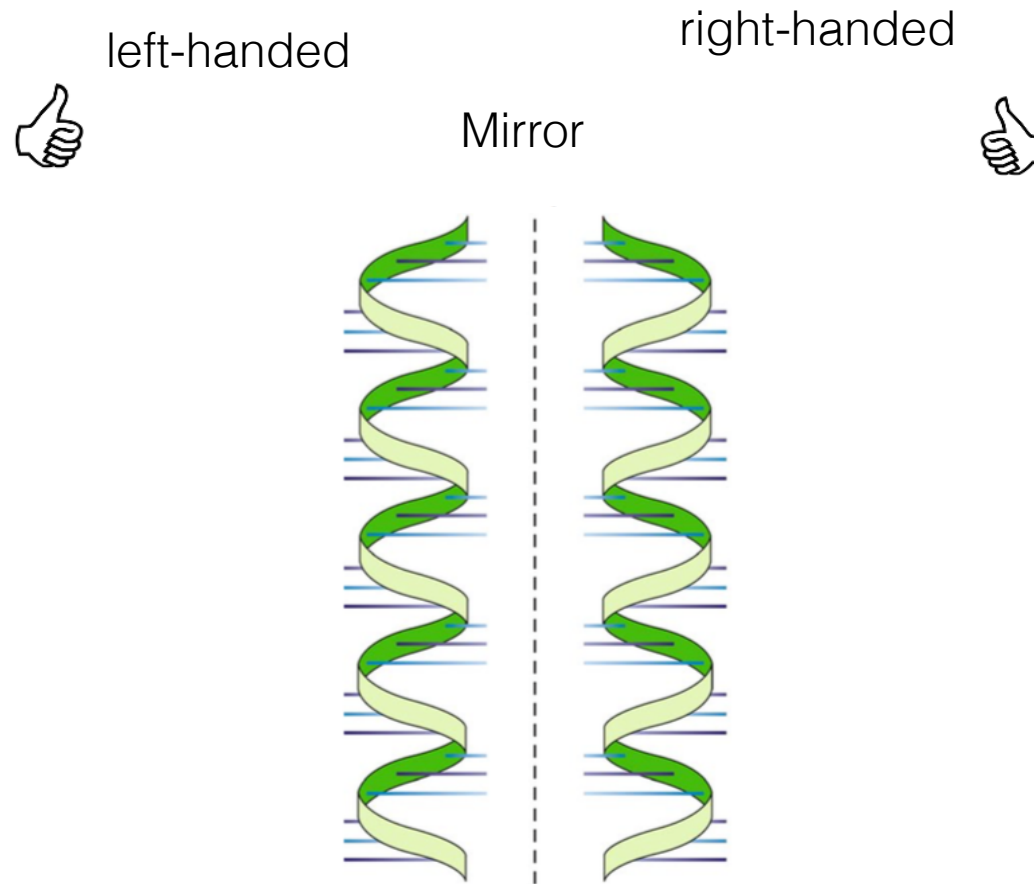
Bzdusek et al., Nature (2016)
Yan et al., Nat. Phys. (2018)

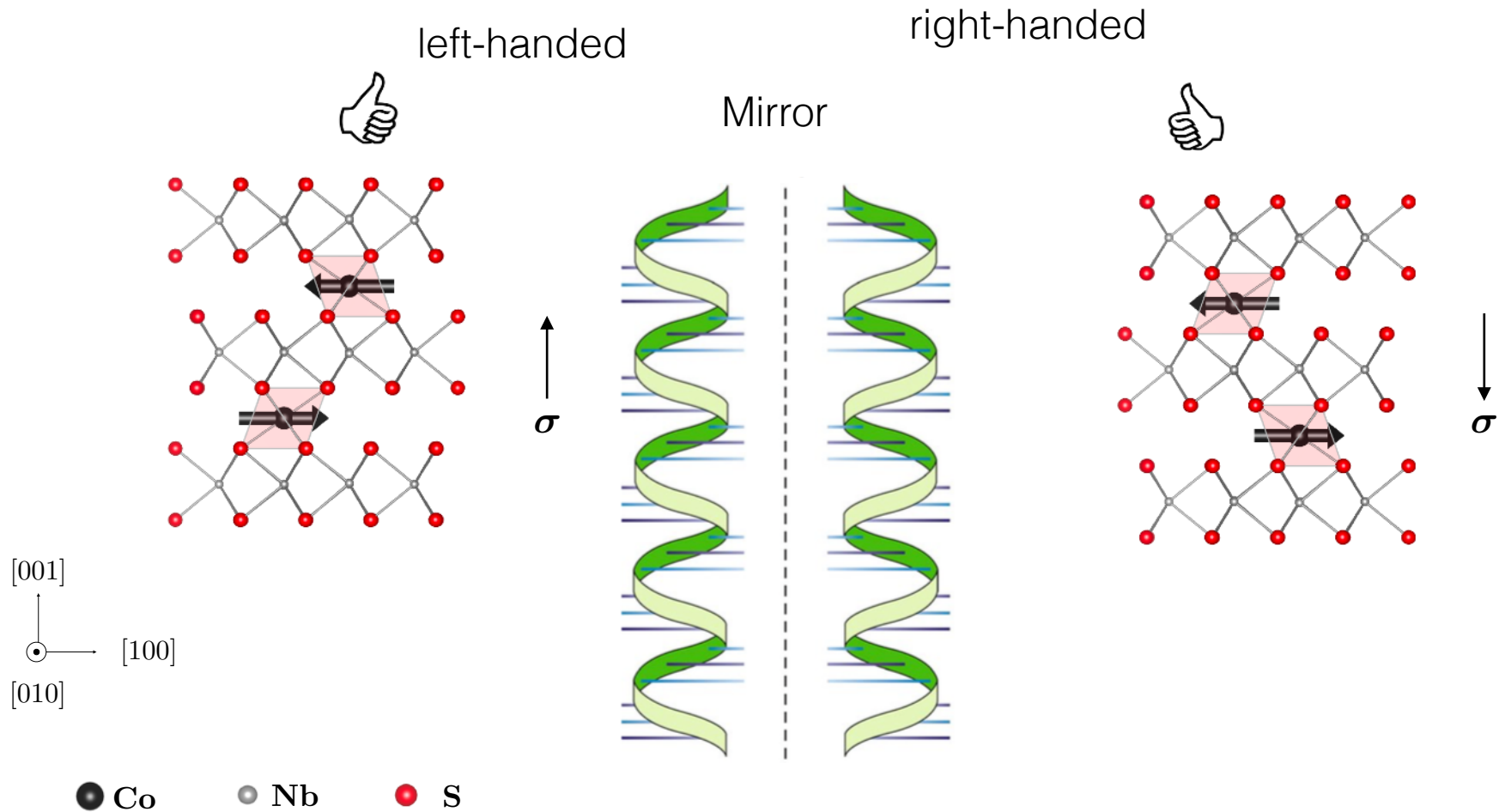


multiple mirror symmetries



If Néel vector points in high symmetry lines, nodal chains are allowed by symmetry, giving rise to high Berry's phase contribution to the crystal Hall effect





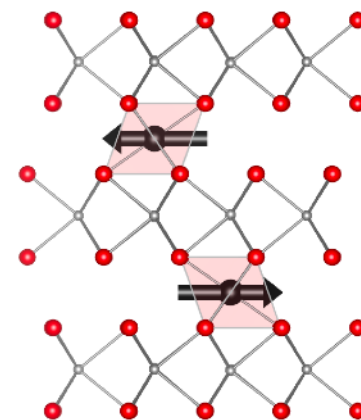
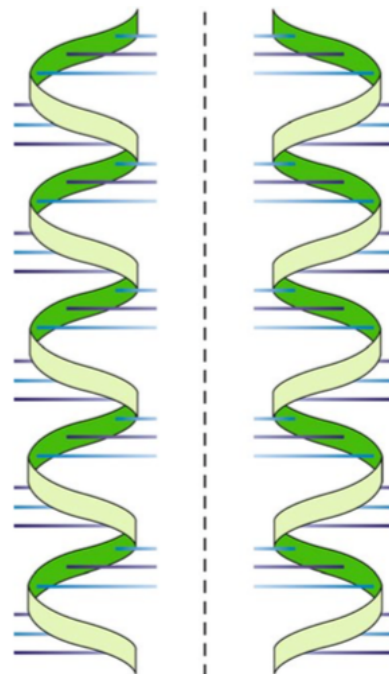
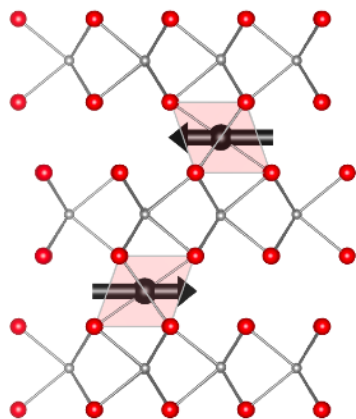
Global chirality in CoNb_3S_6 antiferromagnet

left-handed

right-handed

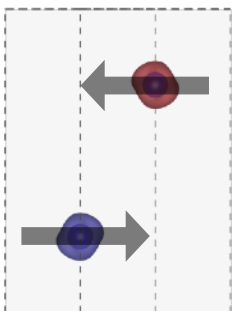


Mirror



σ

σ



[001]
[100]
[010]

● Co ● Nb ● S



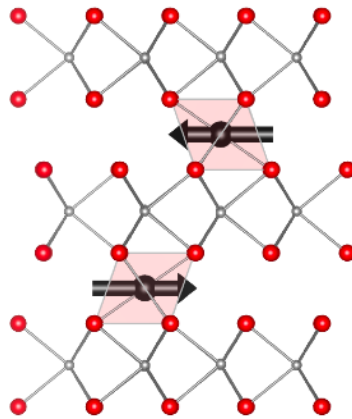
Global chirality in CoNb_3S_6 antiferromagnet

left-handed

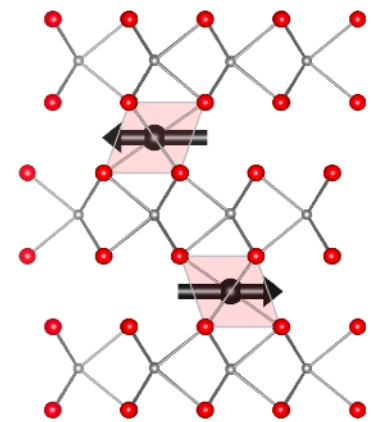
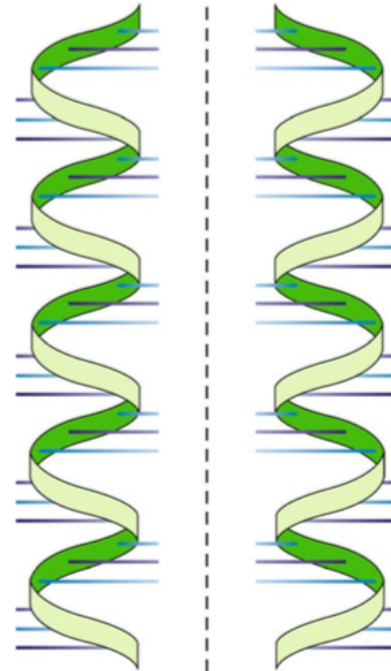
right-handed



Mirror



σ



σ

Experiment: $\sim 27 \text{ S/cm}$
Ghimire et al. Nat. Comm. (2018)
Our theory: $\sim 80 \text{ S/cm}$

[001]
[100]
[010]

● Co ● Nb ● S

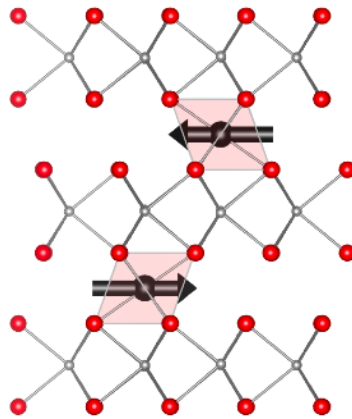


left-handed

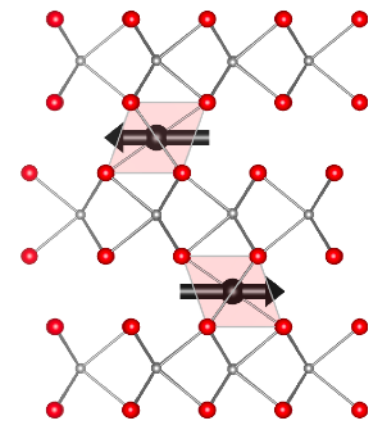
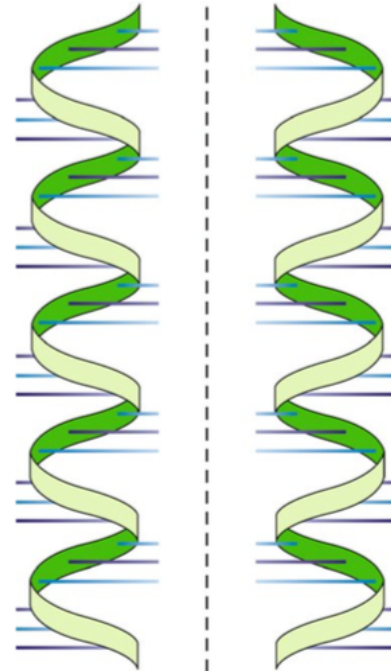
right-handed



Mirror



σ



σ

Experiment: ~ 27 S/cm
Ghimire et al. Nat. Comm. (2018)

Our theory: ~ 80 S/cm

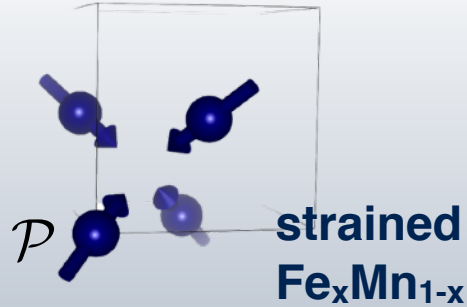
[001]
[100]
[010]

● Co ● Nb ● S

Global crystal chirality controls Hall sign!



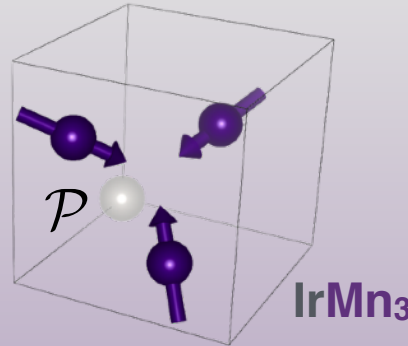
“topological” Hall effect



Onode et al. PRL (2001)
J.P. Hanke et al. Sci.Rep. (2017)

noncoplanar+spin-chirality

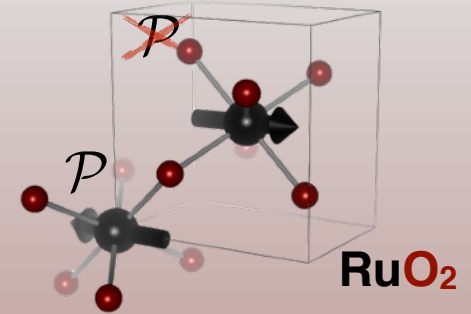
anomalous Hall effect



Hua Chen et al. PRL(2014)

noncollinear+symmetric SOC

crystal Hall effect



Smejkal, Hernandez-Gonzales, TJ, JS, arXiv 1901.00445 (2019)

collinear+asymmetric SOC

spin-orbit coupling complexity

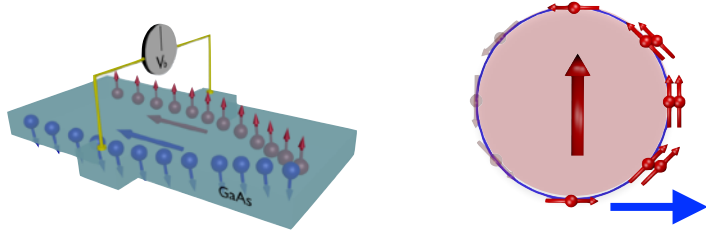
magnetic order complexity

Key signature of CHE:

The sign will change if the broken inversion symmetry (arising from the non-magnetic SOC atoms) is reversed - while retaining the same magnetic order

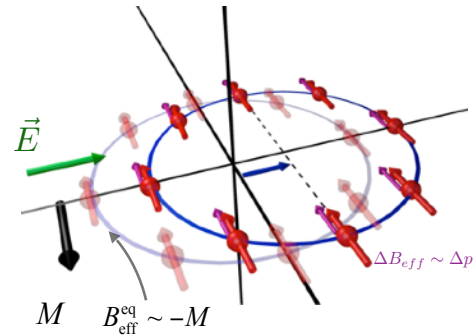
NCAF candidates quite rare but CHE in 1/10 of magnetic database materials!

SHE and ISGE



JS,Valenzuela, Wunderlich, Back,
Jungwirth RMP (2015)

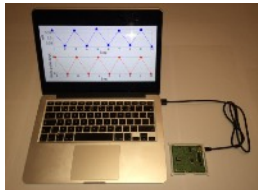
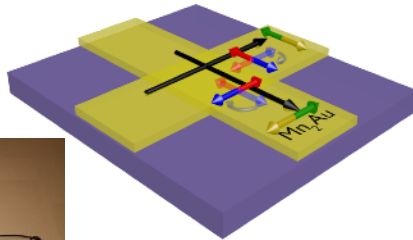
SOT in a single-layer ferromagnet



Kurebayashi, et al.,
Nature Nanotech (2014)

Kurebayashi, et al.,
Nature Physics (2016)

Néel SOT in a single-layer antiferromagnet



J. Zelezny, et al, PRL (2014)
J. Zelezny, et al, PRB (2016)
O. Gomonay, et al, PRL (2016)
Yu, et al. arXiv: 1706.02482 (2017)
Wadley, et al Science (2016)

Jungwirth, JS, et al, Nature Physics (AFM spintronics Reviews) (2018)

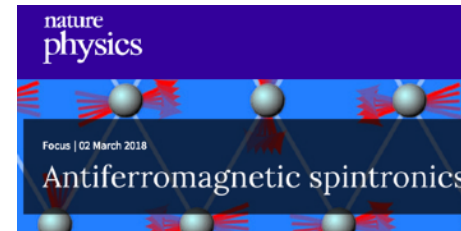
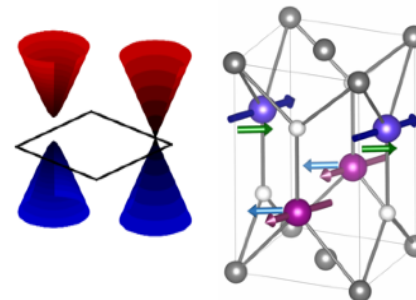
Topological Antiferromagnetic Spin-orbitronics

Topological Dirac

Semi Metal+ AFM (i)

Neel SOT physics (ii)

AHE in RuO2 (col-AFM)(iii)



Libor Smejkal, et al PRL (2017) and PSS (2017)



Kläui, Gomonay



Jungwirth



Zelezny



Institute of Physics Prague



Univ. of Nottingham



Univ. of Cambridge





Kläui, Gomonay



Jungwirth



Zelezny



Institute of Physics Prague



Univ. of Nottingham



Univ. of Cambridge

