

# **Spectrum of the S-1/2, XXZ Antiferromagnet from Zero to High Longitudinal Magnetic Field**

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## **Spin-1/2 Antiferromagnet**

Heisenberg and Ising limits  
Spinons and bound modes

## **SrCo<sub>2</sub>V<sub>2</sub>O<sub>8</sub> – XXZ Heisenberg-Ising chain**

Free spinons  $T > T_N$   
bound spinons  $T < T_N$   
theoretical models

## **SrCo<sub>2</sub>V<sub>2</sub>O<sub>8</sub> – in Longitudinal Magnetic Field**

B, T phase diagram  
Bethe string states  
Bethe Ansatz calculations  
Neutron observation of Bethe strings  
comparison to theory

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Bhabha Atomic Research Centre, Mumbai 400085, India  
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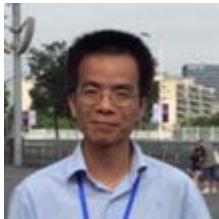
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## Spin-1/2 AFM chains

$$H = J \sum_j S_j^z S_{j+1}^z + \varepsilon (S_j^x S_{j+1}^x + S_j^y S_{j+1}^y)$$

$$H = J' \sum_j (S_j^x S_{j+1}^x + S_j^y S_{j+1}^y) + \Delta S_j^z S_{j+1}^z$$

$$J = J' \Delta \quad \varepsilon = 1/\Delta$$

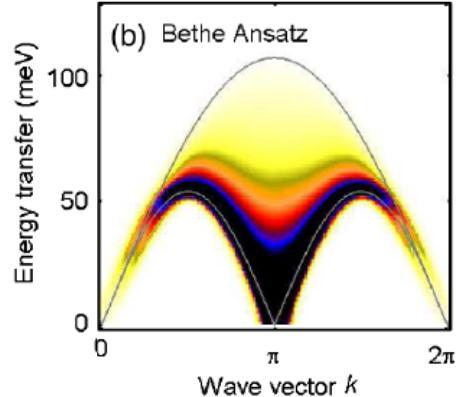
# Spin-1/2, Heisenberg, Antiferromagnetic Chain

$$H_{AFM} = J \sum_j S_j S_{j+1} \quad \begin{matrix} \varepsilon = 0 \\ \Delta \rightarrow \infty \end{matrix}$$

**Hans Bethe, (1931)**  
No long-range order

**Faddeev & Taktajan (1981)**  
Spinon Excitations

**J.-S. Caux, R. Hagemans, J. M. Maillet (2006)** Spectrum

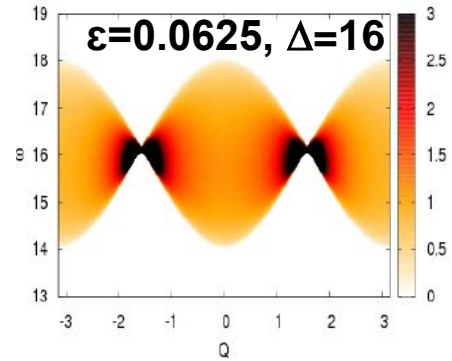
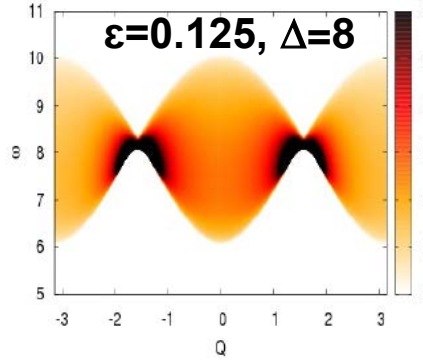
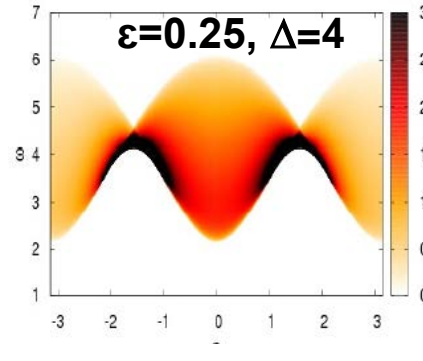
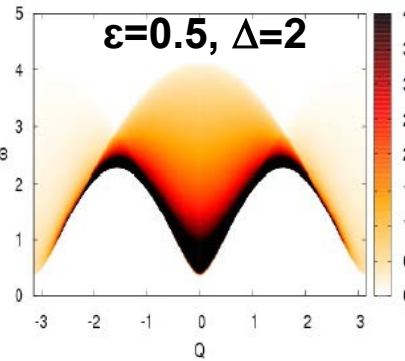


## Spinons

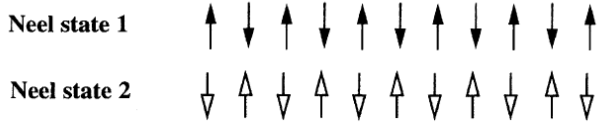
- Fractional S-1/2 particles
- created in multipl pairs
- spinon-pair continuum

$$H = J \sum_j S_j^z S_{j+1}^z + \varepsilon (S_j^x S_{j+1}^x + S_j^y S_{j+1}^y)$$

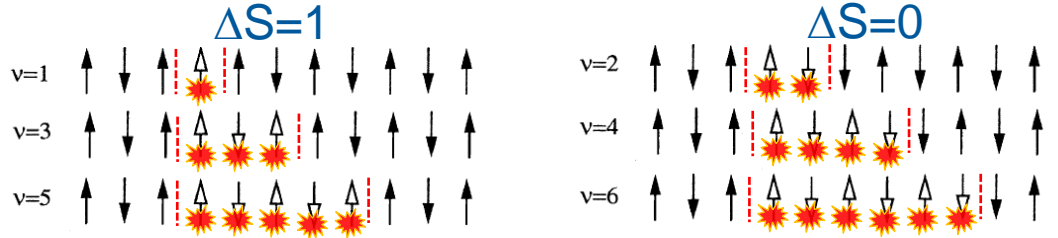
**J..S. Caux J. Mossel, I. P. Castillo, (2008).**  
Spinon Continuum, Ising to Heisenberg  $0 < \varepsilon < 1$



doubly degenerate ground state



highly degenerate first excited states.



# Spin-1/2 Heisenberg Antiferromagnetic Chain

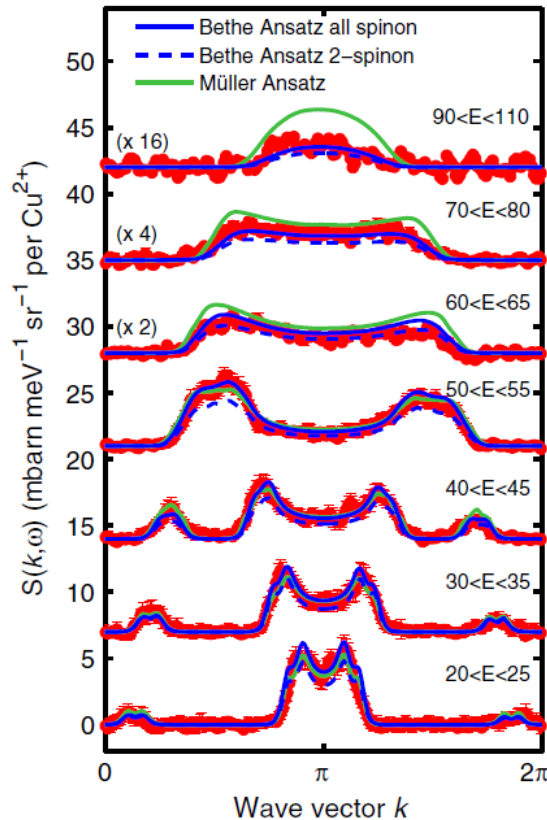
Heisenberg limit ( $\epsilon = 1$ )

$$H = J \sum_j S_j S_{j+1}$$

Gapless Spinon continuum

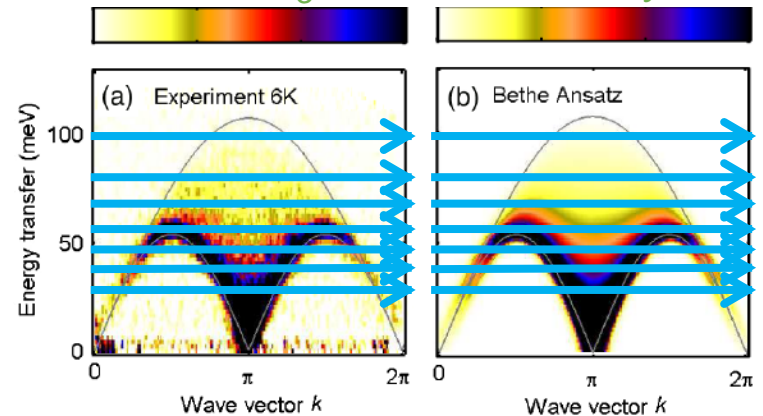
*J.-S. Caux,  
R. Hagemans,  
J. M. Maillet  
(2006)*

Comparison to theories



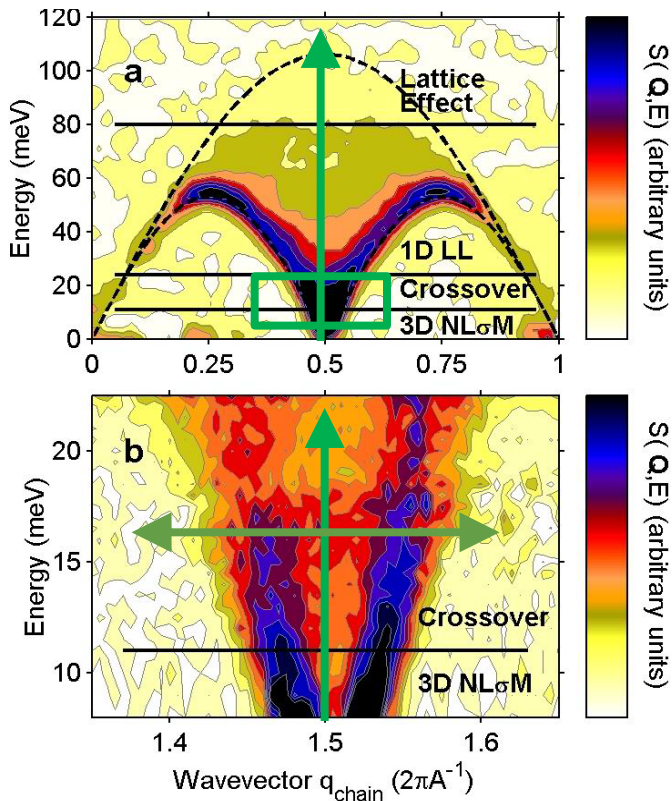
B. Lake; KITP, Oct 2019

MAPS, ISIS  $\text{KCuF}_3$   $T=6\text{K}$  Theory



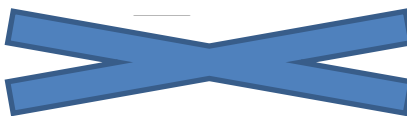
B. Lake et al, Phys. Rev. Lett. (2013)

# Spin-1/2 Heisenberg Antiferromagnetic Chain



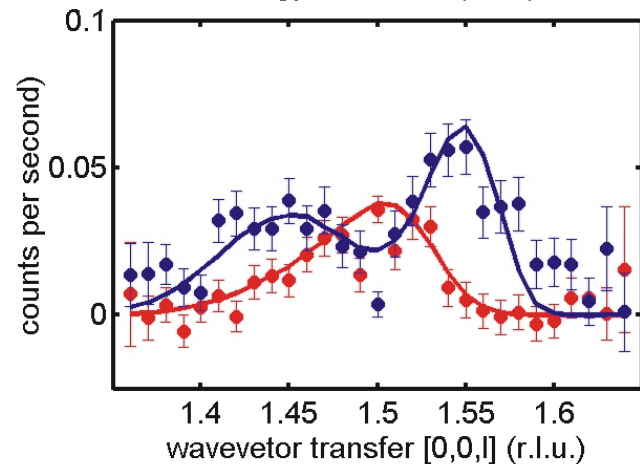
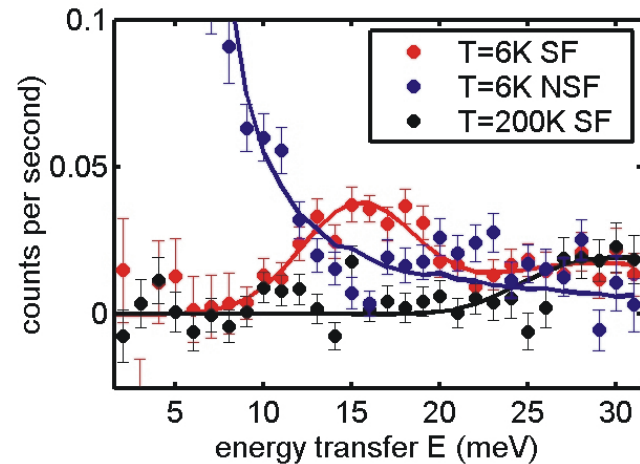
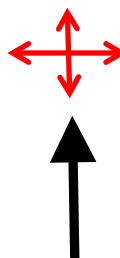
## Free spinon excitations

$$H = J \sum_j S_{j,r} \cdot S_{j+1,r}$$



$$-h^z \sum_j (-1)^j S_j^z$$

$T_N = 39K$   
Polarisation analysis



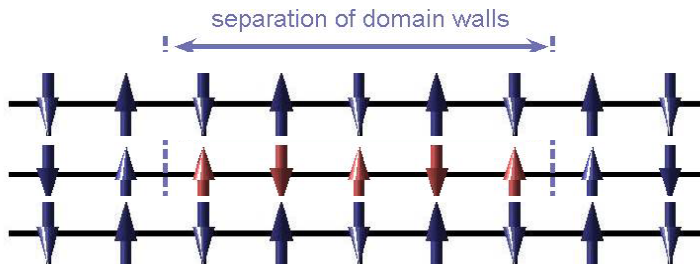
Gapless transverse  $S^{xx}$  spin-waves

Gapped longitudinal  $S^{zz}$  Mode  $\rightarrow$  binding of Spinons

B. Lake, D.A. Tennant, S.E. Nagler et al Phys. Rev. Lett. (2000)

F. H. L. Essler, A.M. Tsvelik, G. Delfino, Phys. Rev. B 56, 11 001 (1997).

Spinons confined by interchain coupling



Longitudinal Mode,

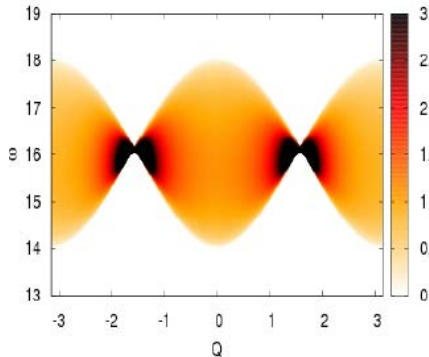
- Observed in  $S^{zz}$
- Broad, finite lifetime

# Spin-1/2 Ising Antiferromagnet Chain

$$H = J \sum_j S_j^z S_{j+1}^z + \varepsilon (S_j^x S_{j+1}^x + S_j^y S_{j+1}^y) - h^z \sum_j (-1)^j S_j^z$$

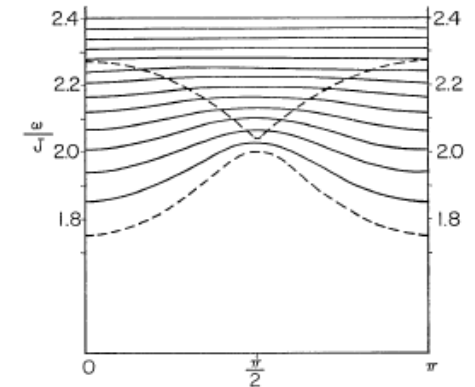
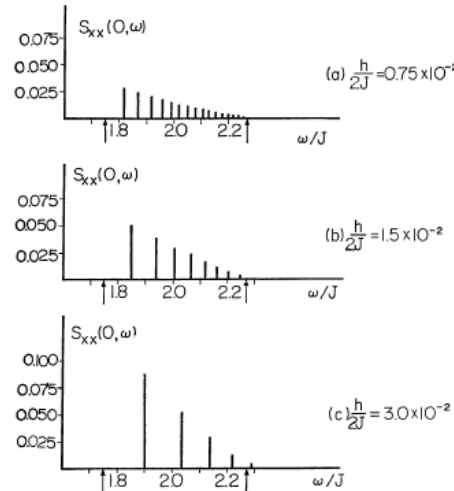
Strong Ising limit  $\varepsilon \ll 1$

Free spinons/kinks



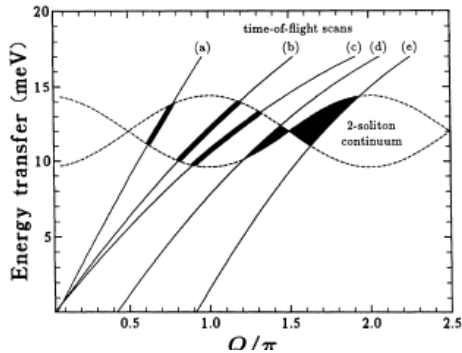
N. Ishimura, H. Shiba JPJS PTP 1980  
J..S. Caux J. Mossel, I. P. Castillo,  
Stat Mech P08006 (2008).

Bound spinons/kinks

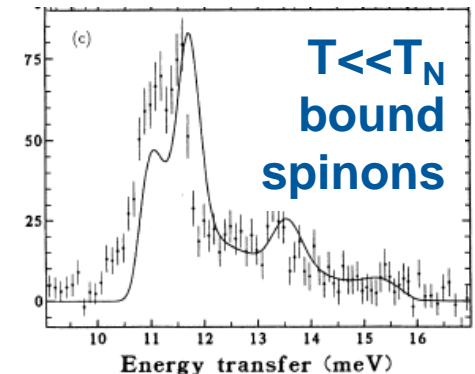
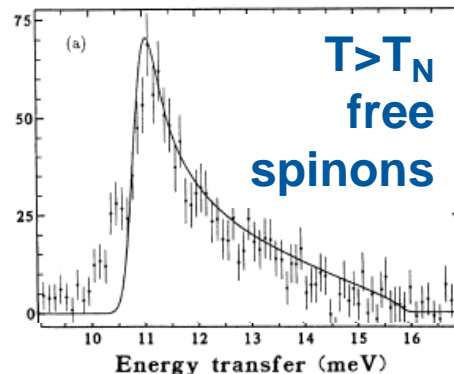


H. Shiba JPJS PTP 1980

**CsCoCl<sub>3</sub>**



J.P. Goff, D.A. Tennant,  
S.E. Nagler PRB 52 15992 (1995)



Transverse bound modes,  $S^{xx}$ , Sharp, long lifetime



# SrCo<sub>2</sub>V<sub>2</sub>O<sub>8</sub> XXZ Spin-1/2 AFM chain

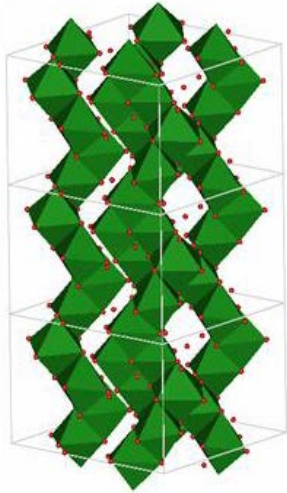
$$H = J \sum_j S_j^z S_{j+1}^z + \varepsilon (S_j^x S_{j+1}^x + S_j^y S_{j+1}^y)$$

$$H = J' \sum_j (S_j^x S_{j+1}^x + S_j^y S_{j+1}^y) + \Delta S_j^z S_{j+1}^z$$

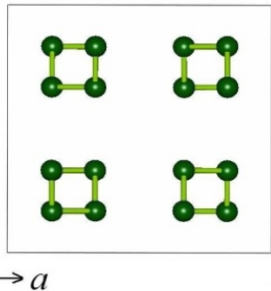
$$\Delta = 2 \quad \varepsilon = 1/2$$

# SrCo<sub>2</sub>V<sub>2</sub>O<sub>8</sub> - Quasi-1D XXZ spin-1/2 Antiferromagnetic

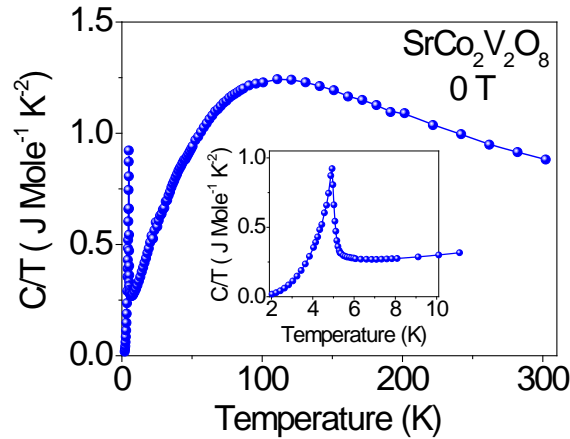
SrCo<sub>2</sub>V<sub>2</sub>O<sub>8</sub>



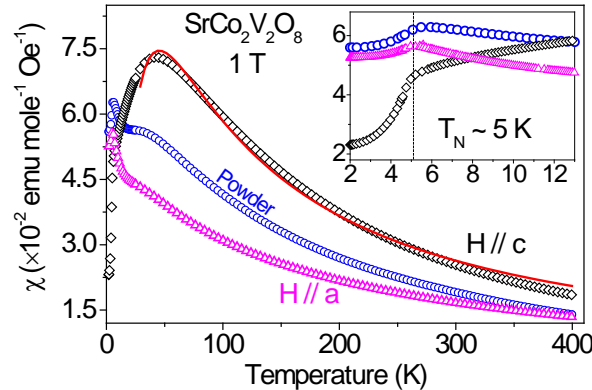
*c*  
*b*  
*a*



Specific heat

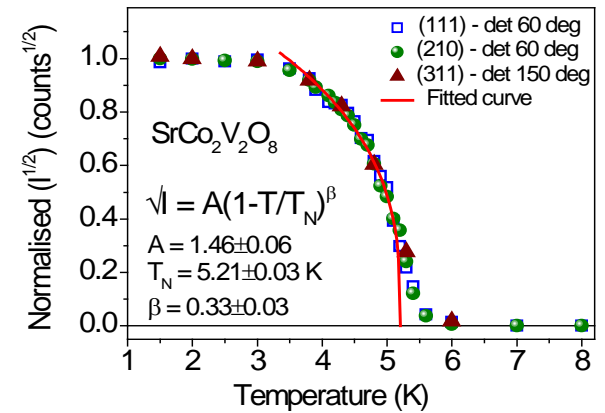


Susceptibility

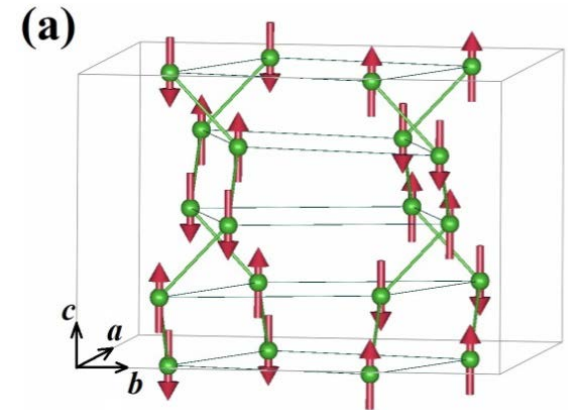


A. K. Bera, B. Lake *et. al.*,  
Phys. Rev. B 89, 094402 (2014)

Neutron diffraction



$M \propto (1-T/T_N)^\beta$   $\beta=0.33$   
Ising Universality Class



Half-integer spin-chain  
Co<sup>2+</sup>; 3d<sup>7</sup>, effective S=1/2  
Kramer doublet - anisotropy

Spin-1/2 screw chains  
along *c* axis.

AFM order below  $T_N = 5.2$  K  
Anisotropic magnetism

Ising universality class  
Spins point along *c* axis

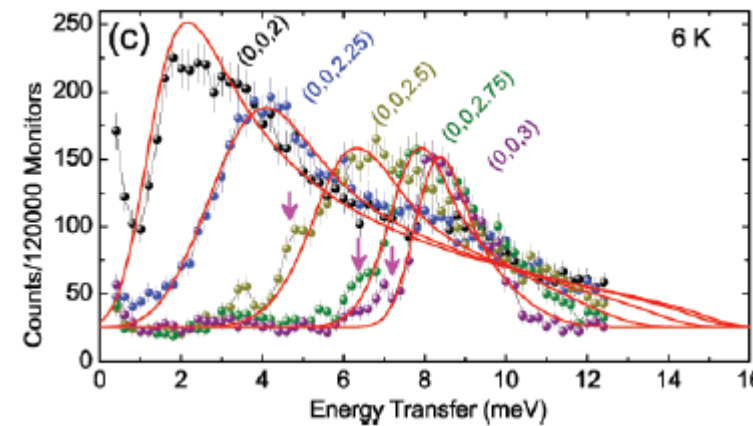
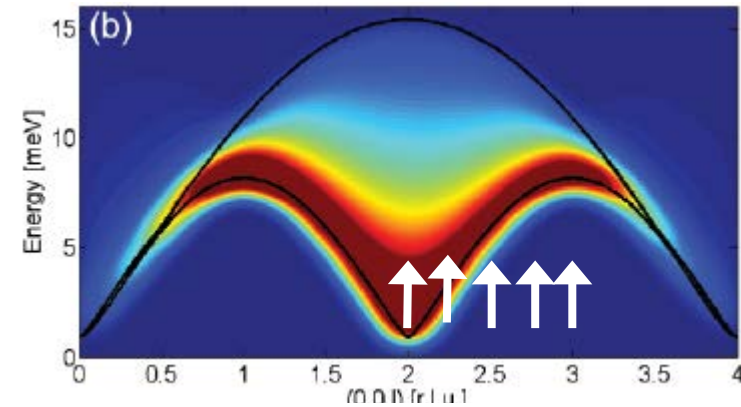
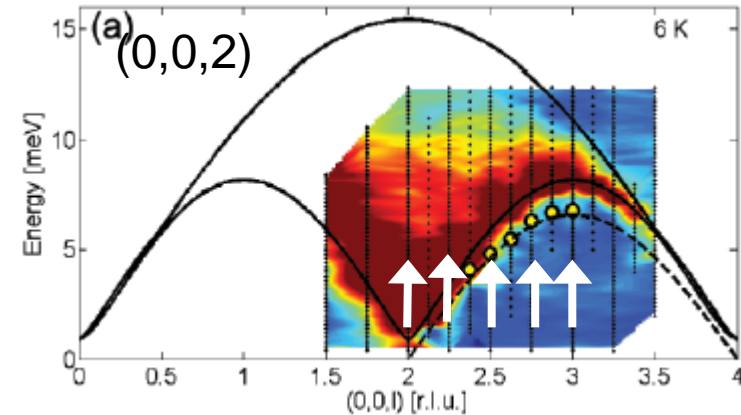
# Magnetic excitations of $\text{SrCo}_2\text{V}_2\text{O}_8$ for $T > T_N$



## Experiment at $T=6\text{K} > T_N=5.2\text{K}$

- Spinon continuum
  - free spinons
- Energy gap,  $\Delta \sim 1.5\text{ meV}$ 
  - XXZ anisotropy
- Dispersions up to  $14.5\text{ meV}$

A. K. Bera, et al. *Phys. Rev. B* 96, 054423 (2017)



## Theory at $T=0\text{K}$

J..S. Caux J. Mossel, I. P. Castillo,  
J. Stat Mech P08006 (2008).

H. Bougourzi, M. Karbach, G. Müller,  
PRB 57, 11429(1998).

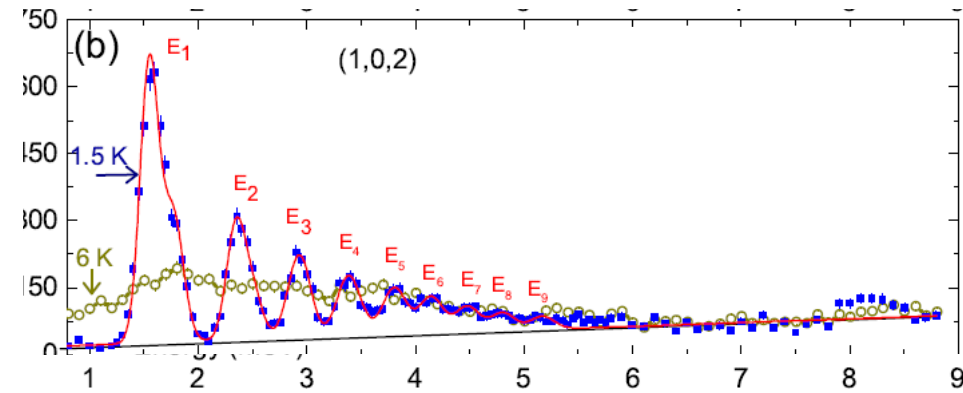
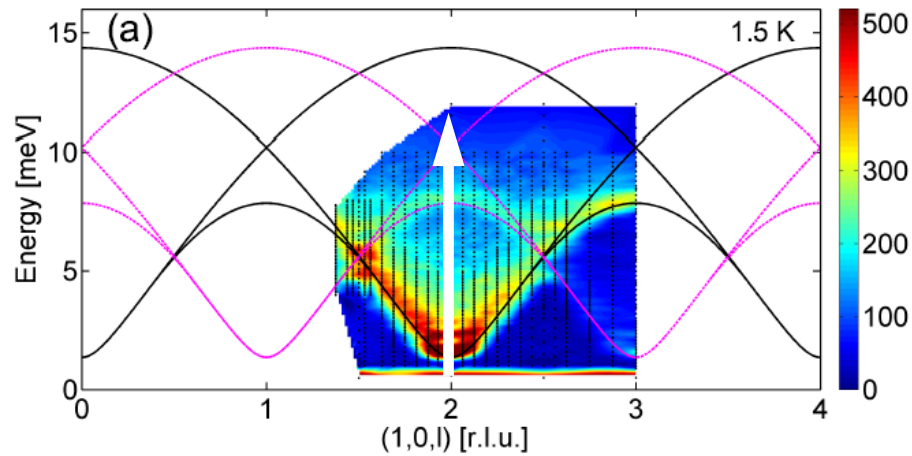
### Fitted values

- $J=7\text{ meV}$
- $\epsilon=0.56$
- Or
- $J=3.9\text{ meV}$
- $\Delta=1.8$

$$H = J \sum_j S_j^z S_{j+1}^z + \epsilon (S_j^x S_{j+1}^x + S_j^y S_{j+1}^y)$$

$$H = J' \sum_j (S_j^x S_{j+1}^x + S_j^y S_{j+1}^y) + \Delta S_j^z S_{j+1}^z$$

# Magnetic excitations of SrCo<sub>2</sub>V<sub>2</sub>O<sub>8</sub> for T < T<sub>N</sub> – Spinon Confinement

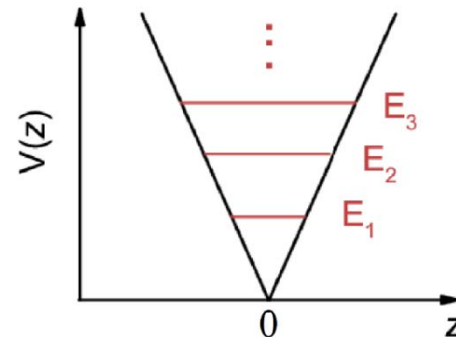
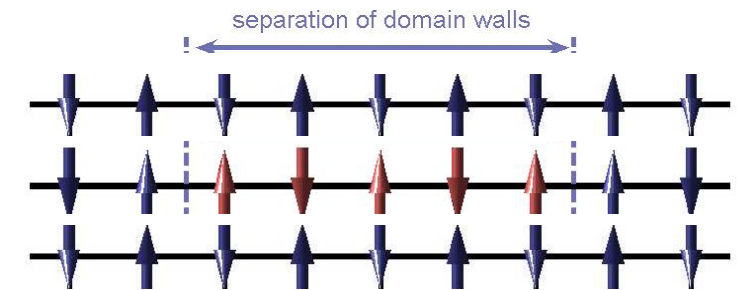


A. K. Bera, B. Lake Phys Rev B 96 054423 (2017)

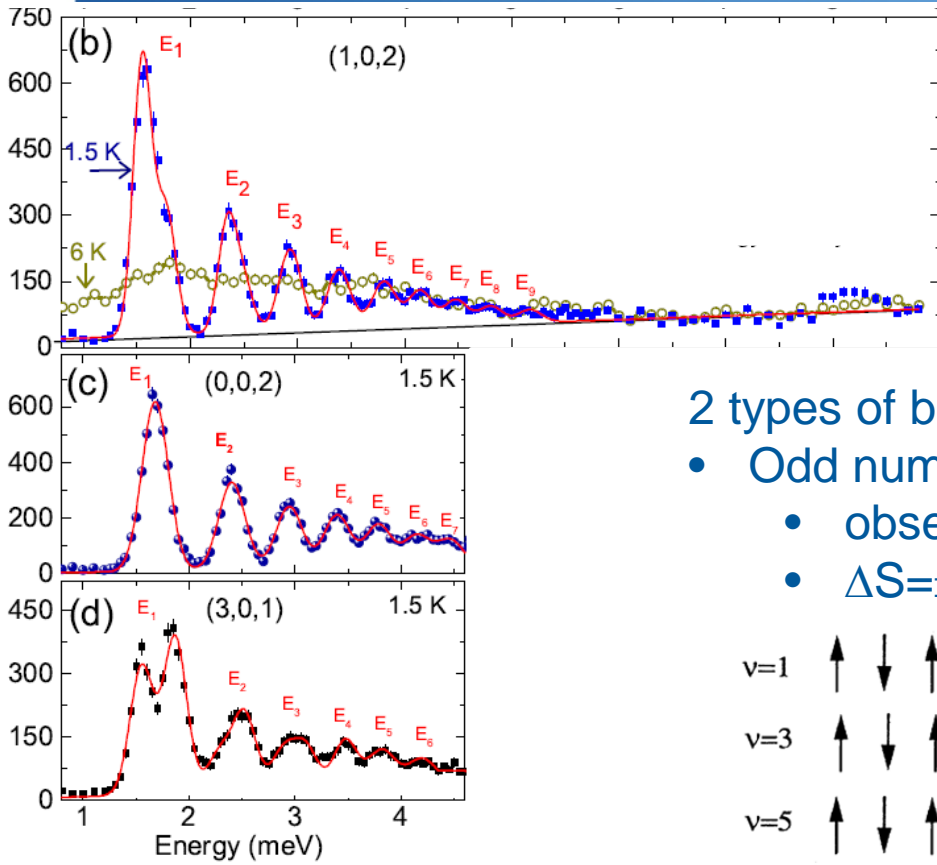
- Hierarchy of discrete excited states

$$H_{AFM} = J \sum_j S_j^z S_{j+1}^z + \varepsilon (S_j^x S_{j+1}^x + S_j^y S_{j+1}^y)$$

- Confinement of pairs of spinons.
- String potential due to interchain coupling.



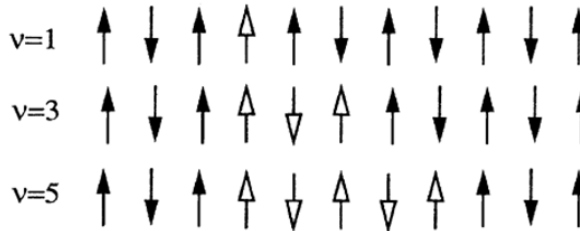
# Magnetic excitations of SrCo<sub>2</sub>V<sub>2</sub>O<sub>8</sub> for T < T<sub>N</sub> – Spinon Confinement



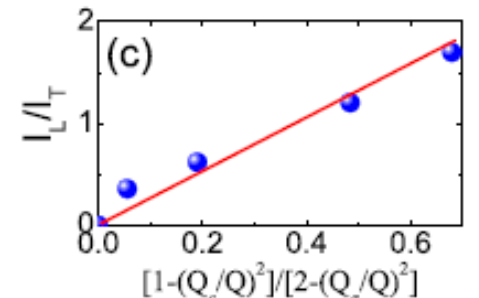
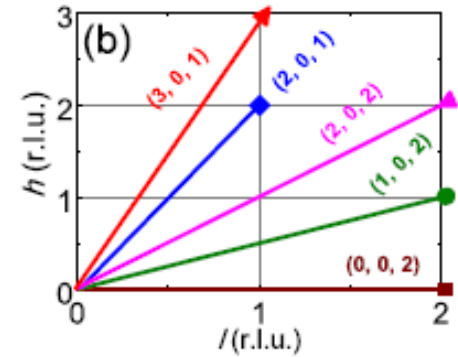
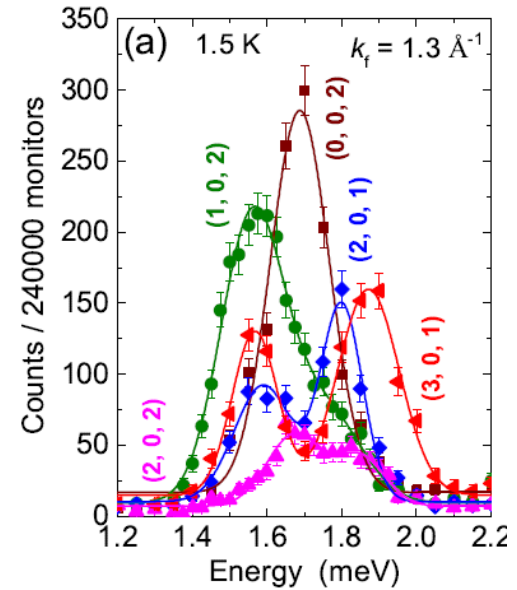
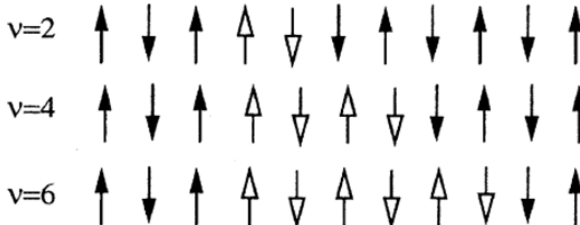
Transverse  
&  
Longitudinal  
excitations

2 types of bound spinons

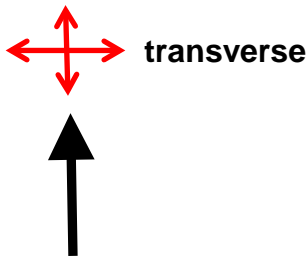
- Odd number of spin flips
  - observed in  $S^{xx}$  and  $S^{yy}$
  - $\Delta S = \pm 1$ , transverse



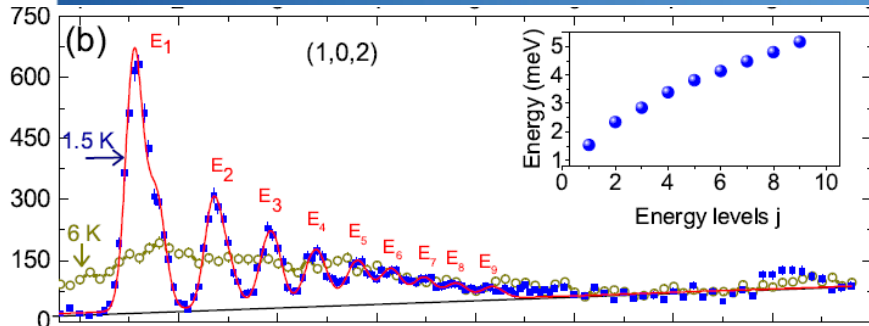
- Even number of spin flips
  - observed in  $S^{zz}$
  - $\Delta S = 0$ , longitudinal



longitudinal



# SrCo<sub>2</sub>V<sub>2</sub>O<sub>8</sub> XXZ, S-1/2 1D AFM – T>T<sub>N</sub> Confined Spinons



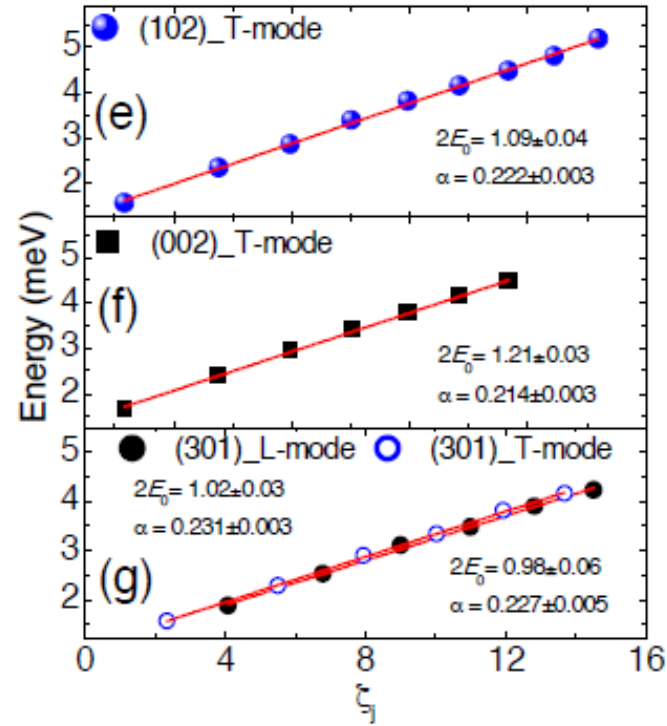
1D Schrödinger equation with linear potential

$$-\frac{\hbar^2}{\mu} \frac{d^2\varphi}{dz^2} + \lambda|z|\varphi = (E - 2E_0)\varphi,$$

Solution: –ve zeros of the Airy function

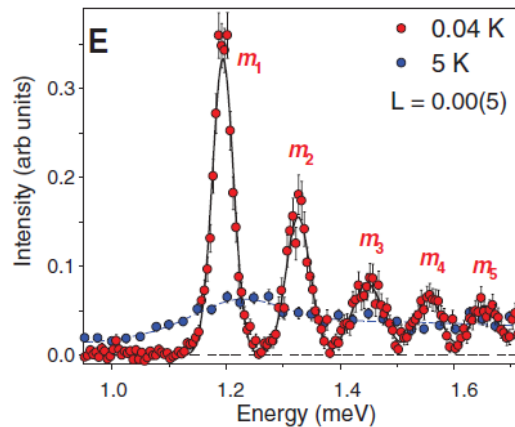
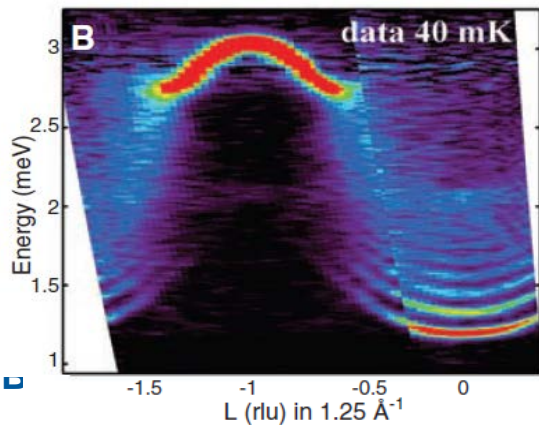
$$E_j = 2E_0 + \alpha\zeta_j \quad j = 1, 2, 3, \dots,$$

$$Ai(-\zeta_j) = 0, \quad \zeta_j = 2.338, 4.088, 5.520,$$



Similar results from BaCo<sub>2</sub>V<sub>2</sub>O<sub>8</sub>

## Spin-1/2 XXZ Ferromagnetic Chain – CoNb<sub>2</sub>O<sub>6</sub>



R. Coldea D.A.Tennant et al science

Bound states only observed in the *transverse* structure factor

S<sup>xx</sup> & S<sup>yy</sup>  
not S<sup>zz</sup>

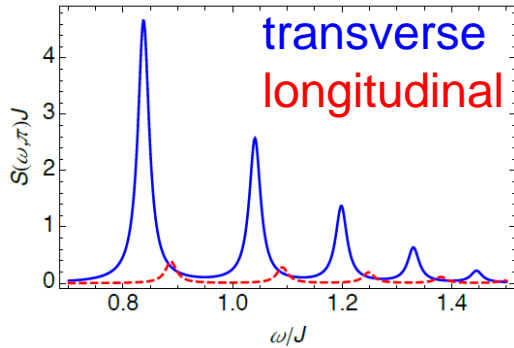
# Theoretical Methods

## Perturbation theory

strong Ising limit  $\varepsilon=0.2$ ,  $h=0.05$

Strong transverse, weak longitudinal

F.H.L. Essler;

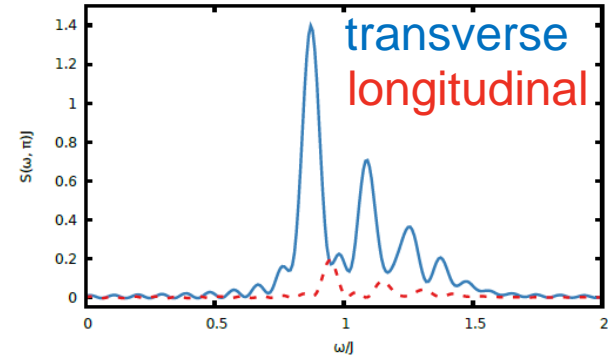


## DMRG

strong Ising limit  $\varepsilon=0.2$ ,  $h=0.05$

Strong transverse, weak longitudinal

C. Hubig, U. Schollwöck

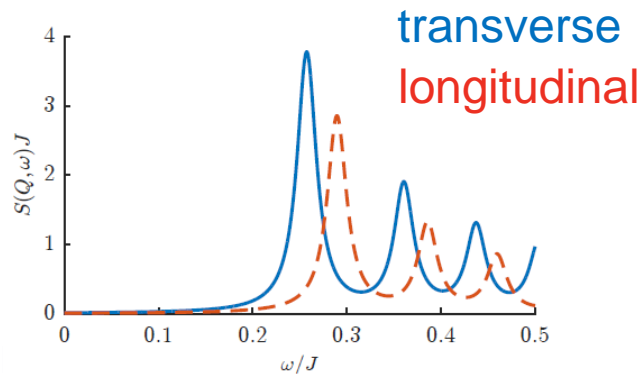


## Tangent Space Matrix Product Space Method

Intermediate XXZ region,  $\varepsilon=0.56$ ,  $h=0.006$

Similar strength Transverse & longitudinal

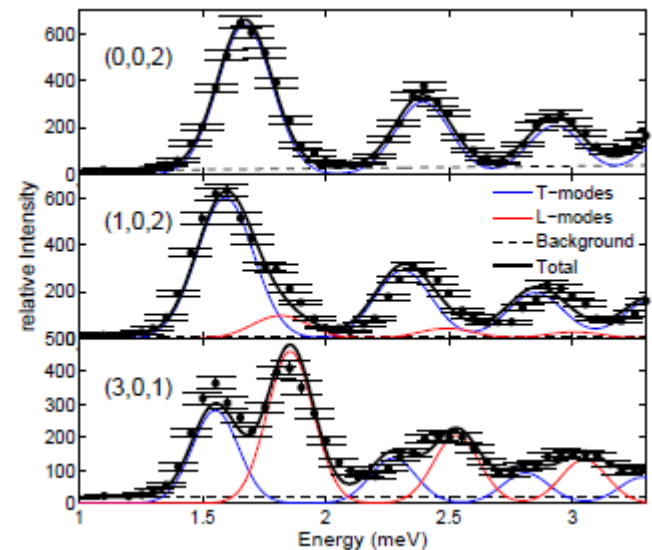
L. Vanderstraeten;



B.

## Comparison to data

$\text{SrCo}_2\text{V}_2\text{O}_8$   $\varepsilon=0.56$ ,  $h=0.006$



# SrCo<sub>2</sub>V<sub>2</sub>O<sub>8</sub> in a Longitudinal Magnetic Field

$$H = J \sum_j S_j^z S_{j+1}^z + \varepsilon (S_j^x S_{j+1}^x + S_j^y S_{j+1}^y) - g \mu_B \sum_i S_i^z B$$

$$H = J' \sum_j (S_j^x S_{j+1}^x + S_j^y S_{j+1}^y) + \Delta S_j^z S_{j+1}^z - g \mu_B \sum_i S_i^z B$$

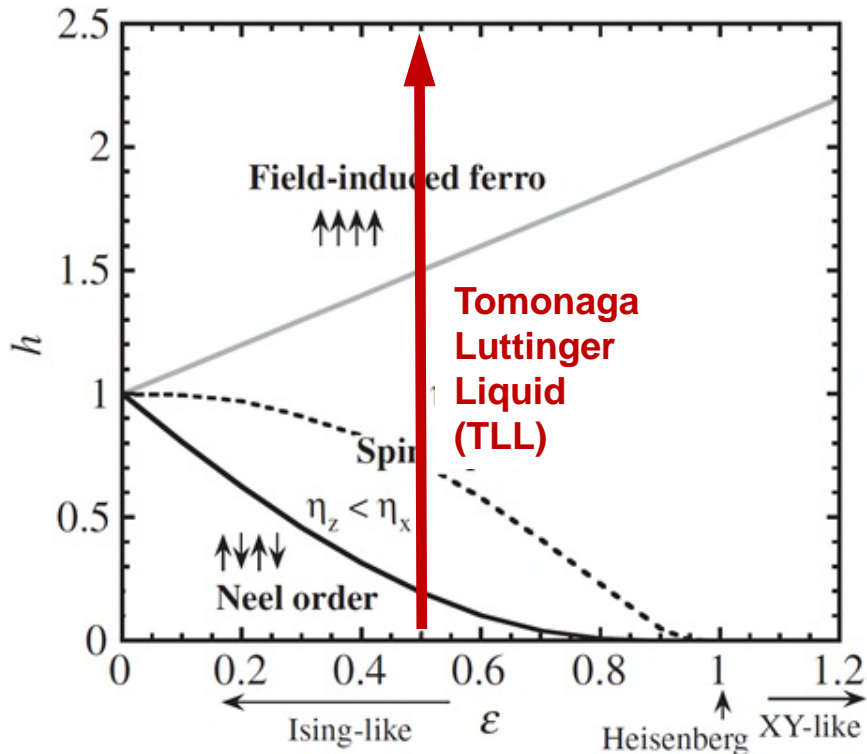
$$\Delta = 2 \quad \varepsilon = 1/2$$



# Spin-1/2 XXZ AFM Chain in a Longitudinal Magnetic Field

$$H = J \sum_j S_j^z S_{j+1}^z + \varepsilon (S_j^x S_{j+1}^x + S_j^y S_{j+1}^y) - g \mu_B \sum_i S_i^z B$$

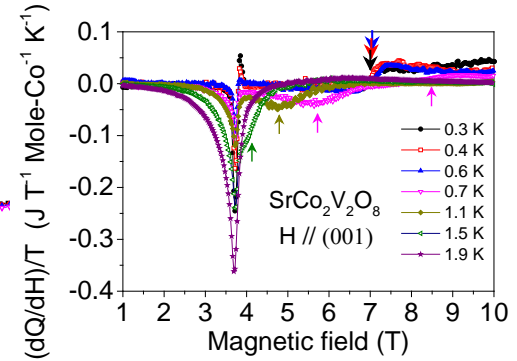
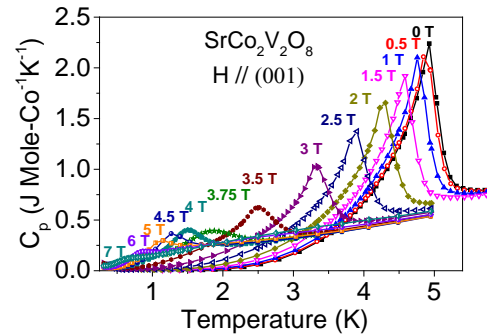
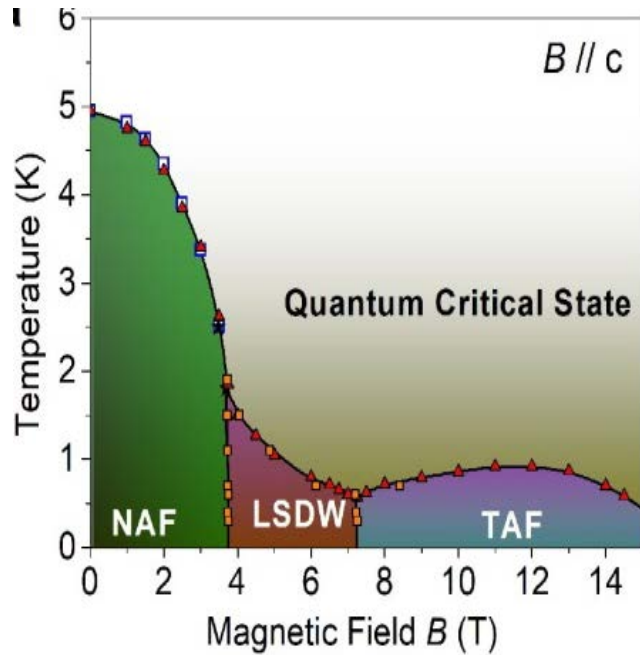
$$H = J' \sum_j (S_j^x S_{j+1}^x + S_j^y S_{j+1}^y) + \Delta S_j^z S_{j+1}^z - g \mu_B \sum_i S_i^z B$$



## Ising-like ( $\varepsilon < 1$ ; $\Delta > 1$ ):

- Field induced quantum critical behavior
- Incommensurate longitudinal spin component
- Staggered transverse component

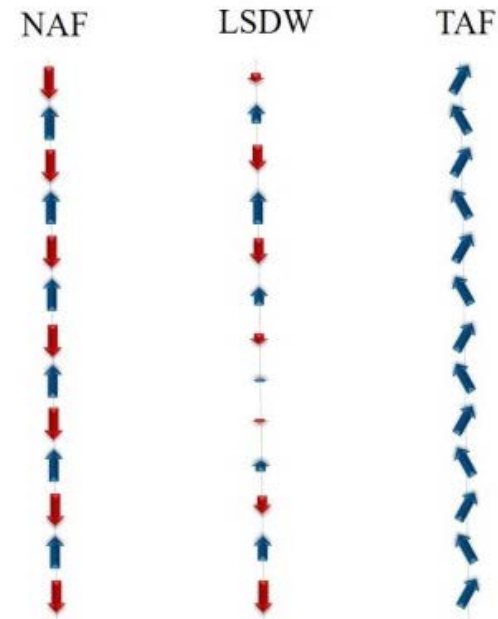
# Phase diagram for $\text{SrCo}_2\text{V}_2\text{O}_8$ in a Longitudinal Field



2 Field-induced phase transitions.  
 $B_{c1} = 3.7$  T:  $B_{c2} = 7.2$  T:

NAF:  $B < B_{c1}$   
 LSDW:  $B_{c1} < B < B_{c2}$   
 TAF:  $B_{c2} < B < B_s$

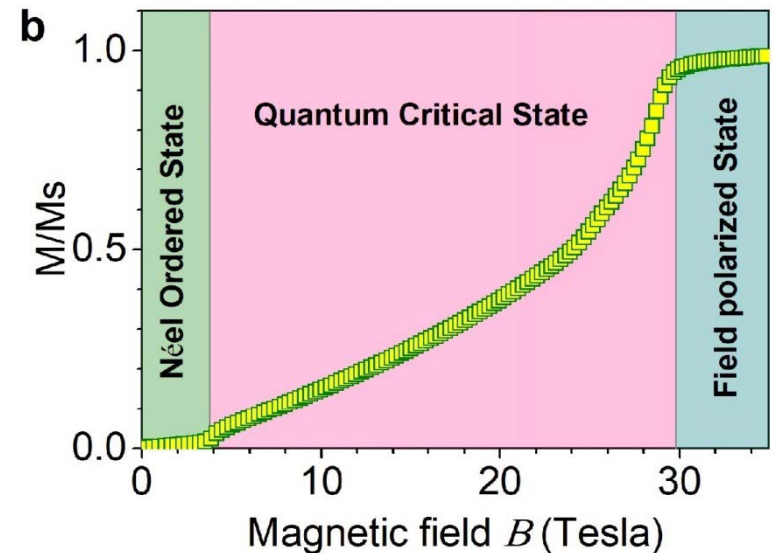
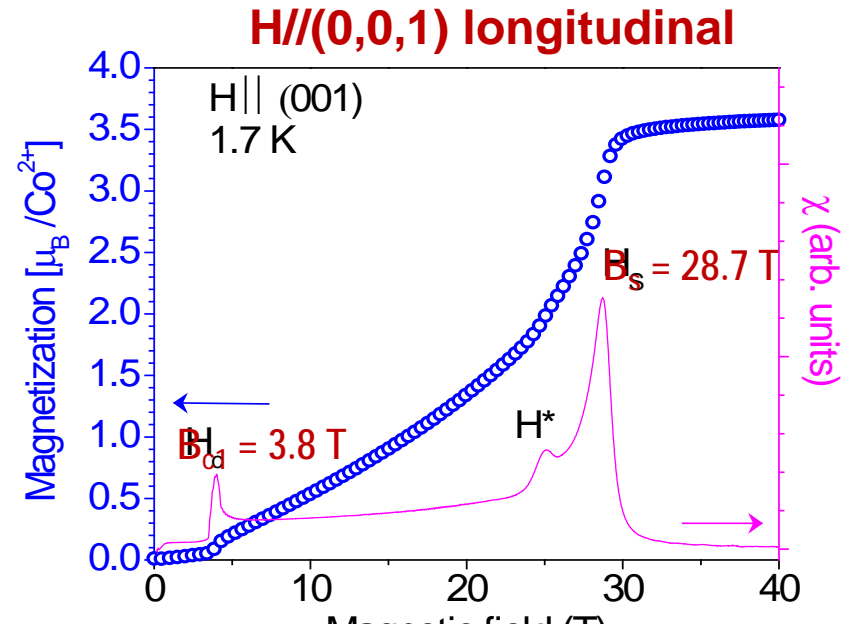
Long-range magnetic order is  
 B. La strongly suppressed above  $B_{c1}$



# Phase diagram for $\text{SrCo}_2\text{V}_2\text{O}_8$ in a Longitudinal Field

## Longitudinal field

- Field induced QPT at  $B=3.8\text{T}$  to quantum critical Tomonaga Luttinger Liquid.
- Non-linear magnetization indicates strong quantum fluctuations
- Saturation at  $B_s=28.7\text{T}$

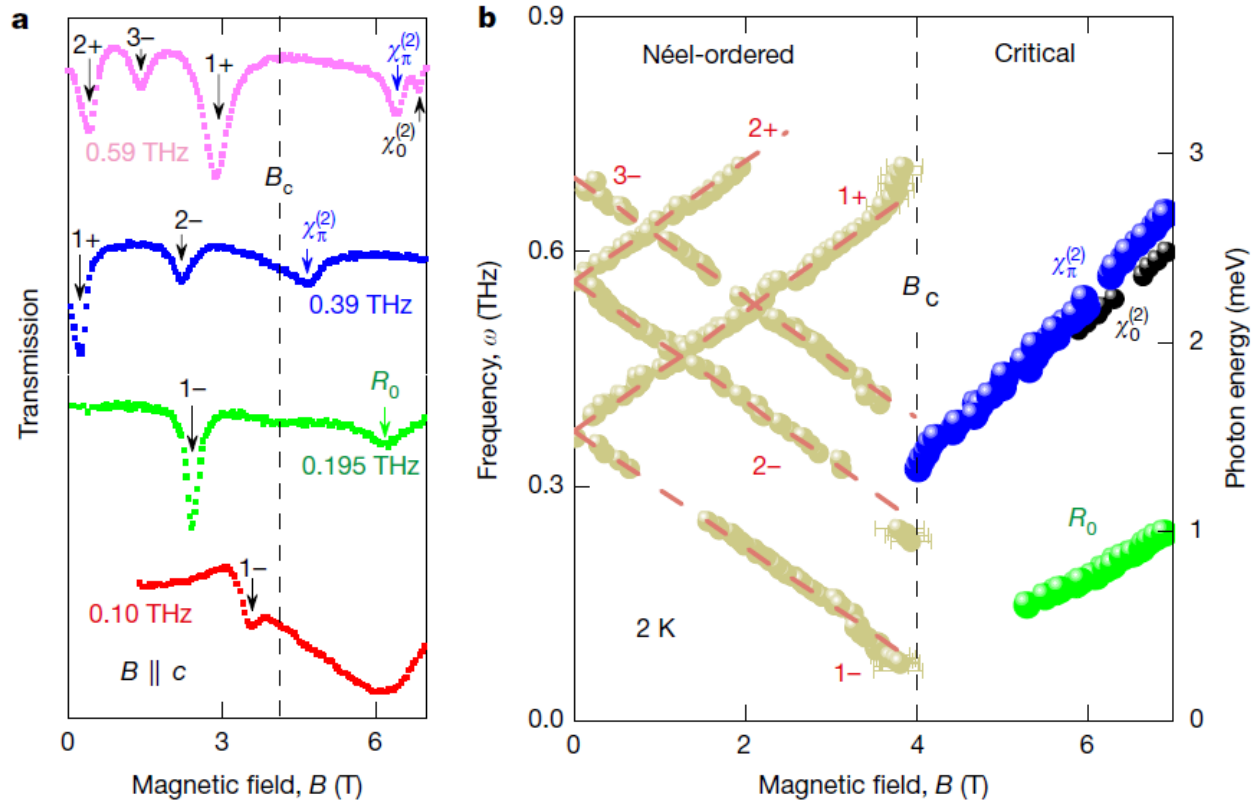


# Excitations of $\text{SrCo}_2\text{V}_2\text{O}_8$ in a small longitudinal field

## Terahertz Optical Spectroscopy

Field-dependent spectrum at  $Q = 0, \pi/2, \pi; 3\pi/2$   
 Transverse excitations only

*Zhe Wang, Jianda Wu,  
 Wang Yang et al  
 Nature 554 221 (2019)*



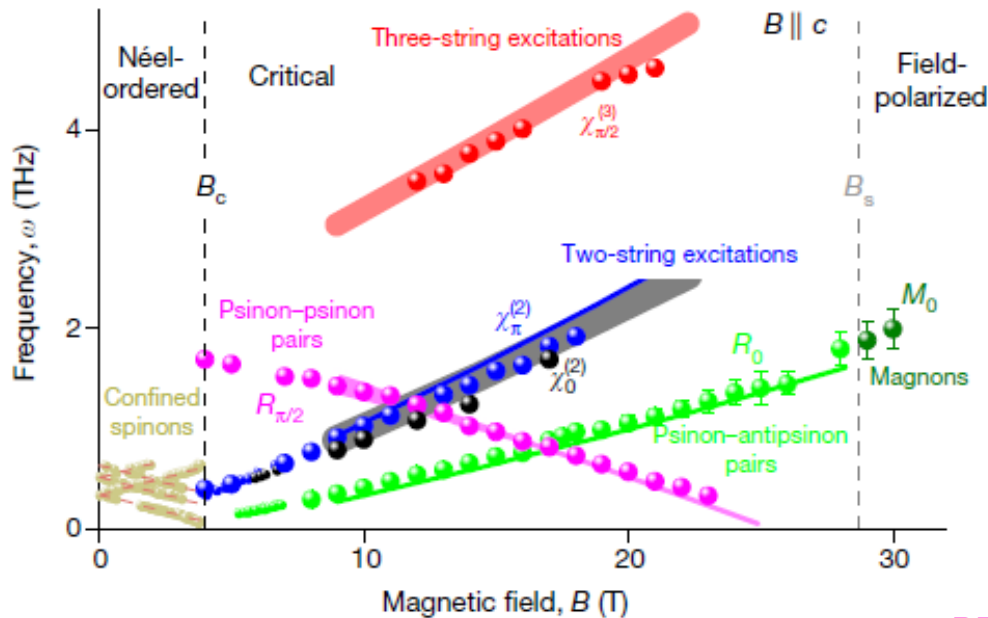
- $B < B_{c1}$  Zeeman splitting of the the bound spinons
- $B = B_{c1}$  Lowest E excitation goes into the ground state
- $B > B_{c1}$  Complete change in the spectrum

# Excitations of $\text{SrCo}_2\text{V}_2\text{O}_8$ in a high longitudinal field

## Terahertz Optical Spectroscopy for $B > B_{c1}$

Field-dependent spectrum at  $Q = 0, \pi/2, \pi$

Transverse excitations only



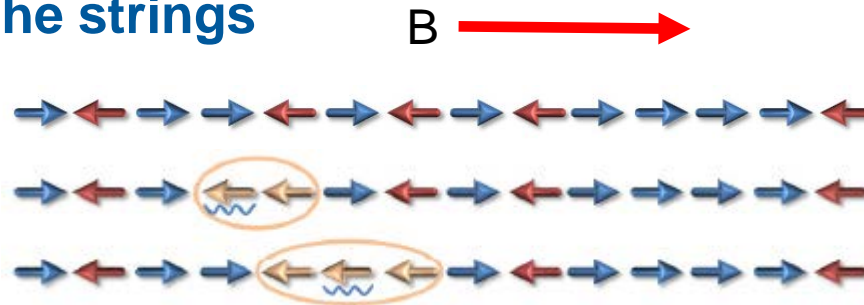
New excitations observed for  $B > B_{c1}$  identified by comparison to Bethe Ansatz

*Zhe Wang, Jianda Wu, Wang Yang et al  
Nature 554 221 (2019)*

$R^{PP}$  psinon-psinon pair  
 $R^{PAP}$  psinon-antipsinon pair  
 $\chi^{(2)}$  2-string Bethe  
 $\chi^{(3)}$  3-string Bethe  
 n-string = addition of n bound magnons  
 (spins pointing against the field)

# Bethe Strings

## History of Bethe strings



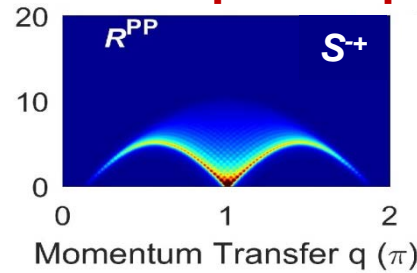
- A sequence of flipped spins or magnons bound together within the strongly correlated background of single spin flips
- 2-string state first predicted for Heisenberg chain in zero field by Hans Bethe in 1931
- Multi-string states later predicted by several theorists stabilized by anisotropy
- Not observed due to their weak signal dominated by the spinon spectrum.
- Observable for XXZ chain in a longitudinal magnetic field which separates them from the other excitations

# Theory for S-1/2 1D XXZ AFM in longitudinal Field

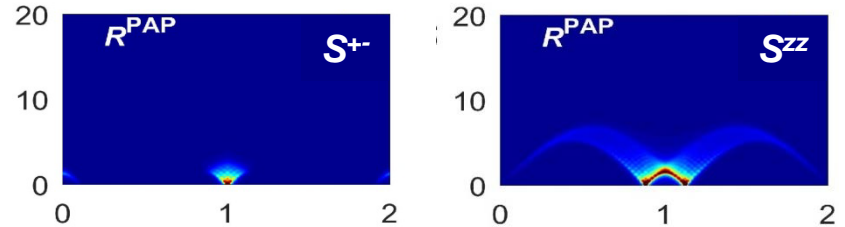
Bethe Ansatz  
 Jianda. Wu,  
 Wang Yang

- $J=3.9$  meV
- $\Delta=1.8$
- $2m=0.12$

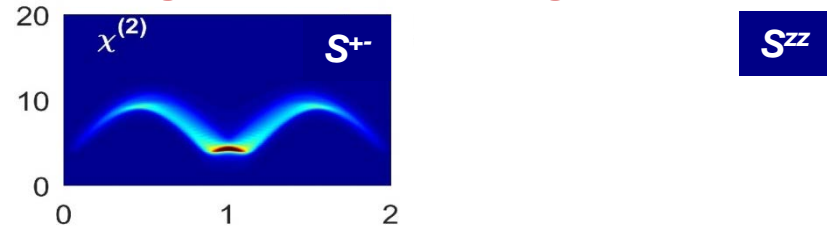
## Psinon-psinon pair



## Psinon-antipsinon pair

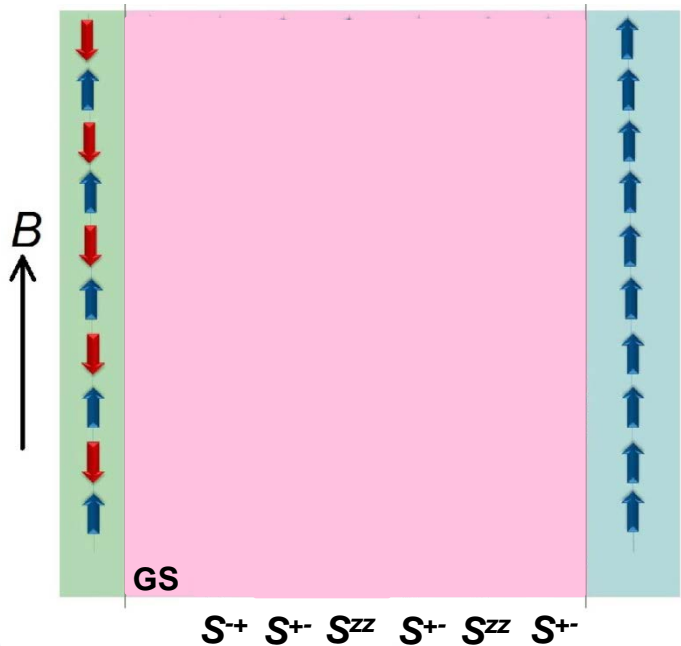
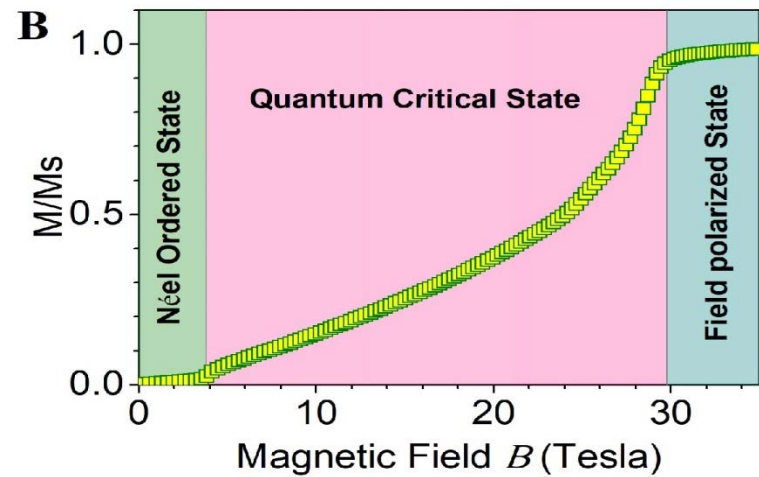


## 2-magnon Bethe String



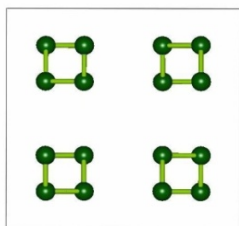
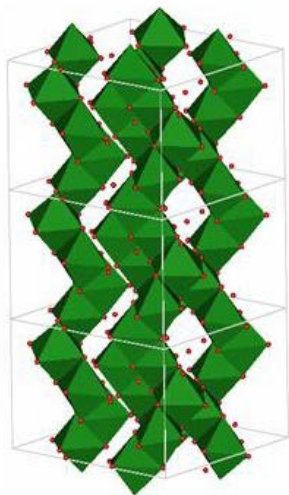
## 3-magnon Bethe String

$S^{+-}$



# Dispersions of Psinon, Anti-Psinon & Bethe Strings

$\text{SrCo}_2\text{V}_2\text{O}_8$

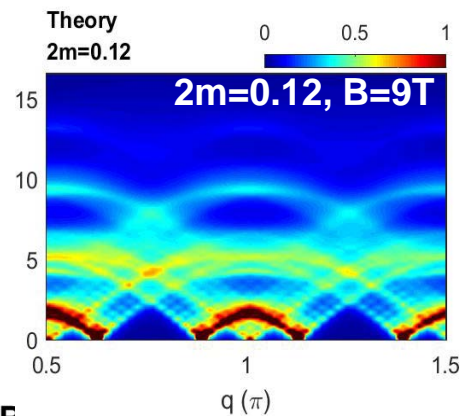


$b$   
 $a$

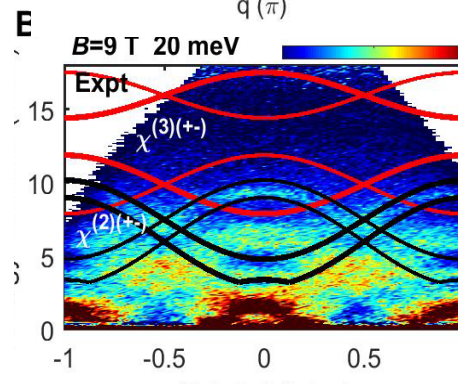
4-fold screw chain



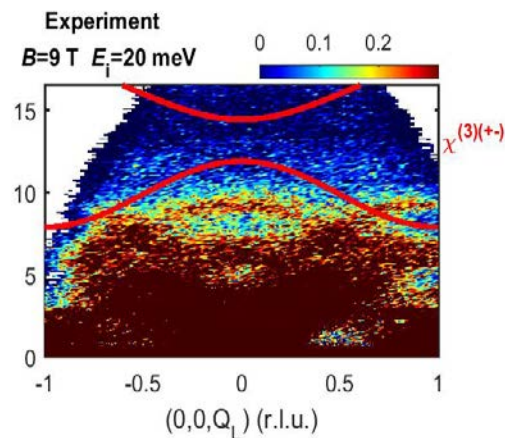
Zone folding



Theory



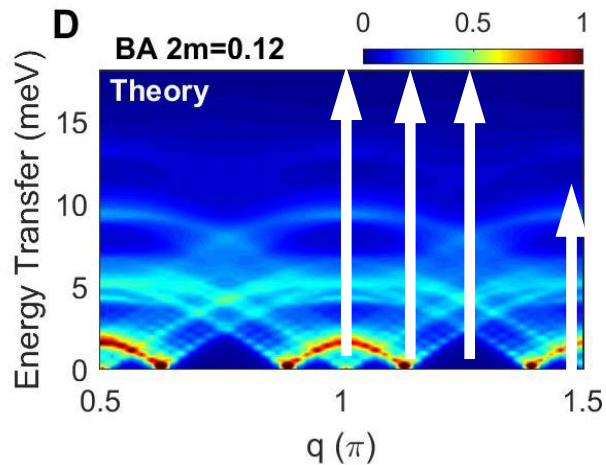
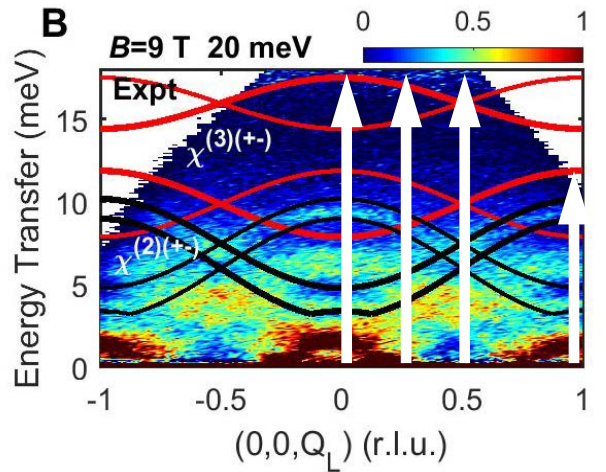
Experiment



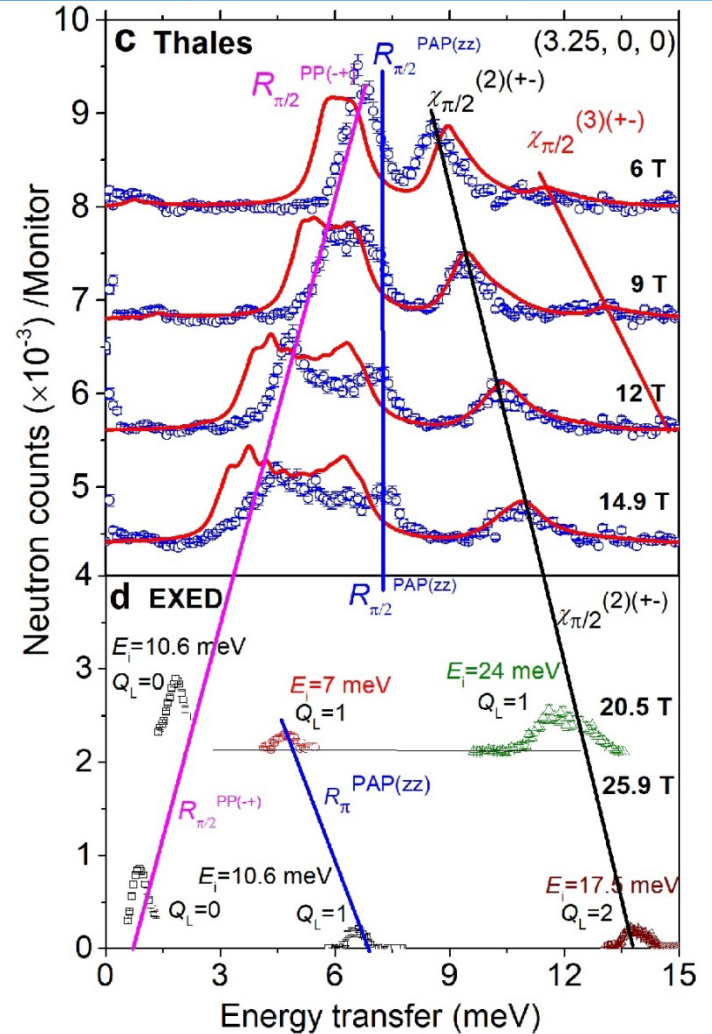
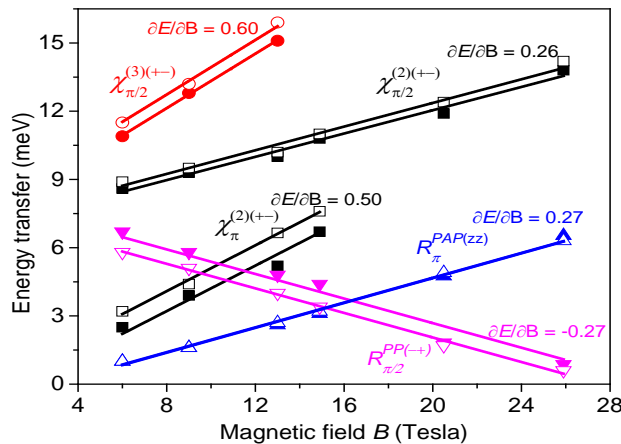


# Dispersions of Psinon, Anti-Psinon & Bethe Strings

Detailed  
comparison to  
Bethe Ansatz



# Field-Dependence of Psinon, Anti-Psinon & Bethe Strings



Good agreement between experiment and theory for the field dependence and intensity

## Spin-1/2 Antiferromagnet

Heisenberg and Ising limits

Spinons and bound modes

## $\text{SrCo}_2\text{V}_2\text{O}_8$ – XXZ Heisenberg-Ising chain

Free spinons  $T > T_N$

bound spinons  $T < T_N$

theoretical comparison

## $\text{SrCo}_2\text{V}_2\text{O}_8$ – in Longitudinal Magnetic Field

B, T phase diagram

Bethe string states -terahertz

Bethe Ansatz calculations

Neutron observation of Bethe strings

Agreement for dispersion, field, intensities,