

Staggered current Mott insulator of kinetically frustrated bosons

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Discussions

Bhanu Das, Matthew Fisher, Joseph Thywissen, David Huse

arXiv:1109.5703 + stuff in preparation

Support:



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Strongly interacting bosons

Bose-Hubbard Hamiltonian: Simplest model of strong correlation effects

$$H = - \sum_{i,j} t_{i,j} b_i^\dagger b_j + \frac{U}{2} \sum_i n_i(n_i - 1)$$

M.P.A. Fisher et al, PRB 1989

For large fillings, this is closely related to a Josephson junction array (JJA) model
 $U \sim$ charging energy, $t \sim$ Josephson coupling

S. Chakravarty, et al, PRL 1986

Mooij group, PRL 1989

Clarke group, PRL 1997

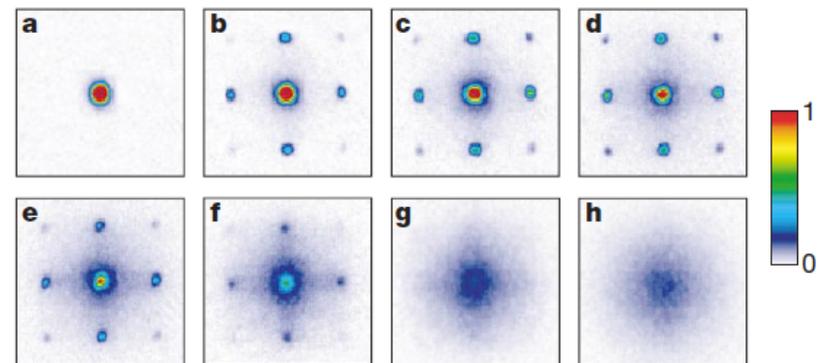
Integer filling

Superfluid \bullet Mott insulator \rightarrow U/t

$(d+1)$ XY

He-4, JJAs, Cold Atoms

Mott insulators of Rb-87 atoms



Greiner et al, Nature 2002

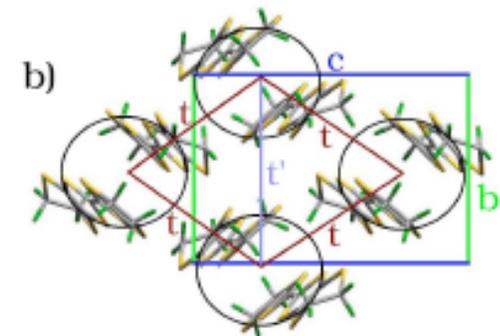
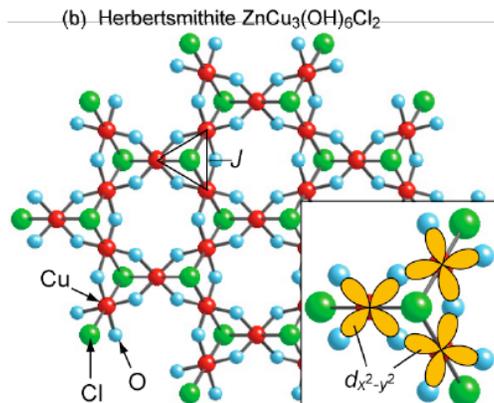
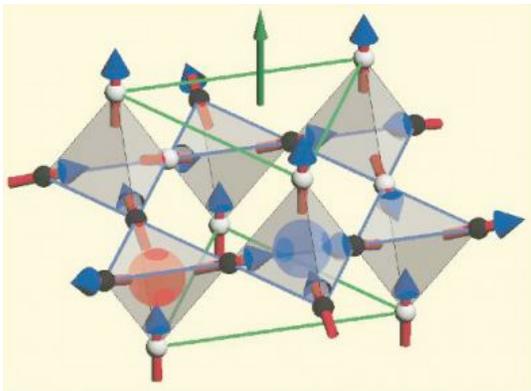
Frustration

Why is frustration interesting?

- . Many many ground states at classical level (Entropy grows with size)
- . Quantum effects have to select unique ground state (“zero” entropy)
- . Source of interesting new states: FQHE, quantum spin liquids

Examples of frustrated spin liquid-like materials?

- . Pyrochlore [spin ice] materials
- . Kagome magnets [Herbertsmithite]
- . Organic antiferromagnets [kappa-BEDT, Pd(MIT)2] - *near Mott transition*



Bosons: Frustration + Strong Correlations

Fully frustrated Bose-Hubbard Hamiltonian

$$H = - \sum_{i,j} t_{i,j} b_i^\dagger b_j + \frac{U}{2} \sum_i n_i (n_i - 1)$$

Pick the kinetic term to include fluxes

Φ	Φ
Φ	Φ

For $\Phi = p/q$ *D. R. Hofstadter (PRB 1976)*

Get q bands with q -degenerate minima in the dispersion

Bosons are thus **kinetically frustrated** and have various ways in which to condense

This kinetic frustration must be resolved by interaction effects

What is the phase diagram as we vary U/t ?

Why is this interesting or relevant?

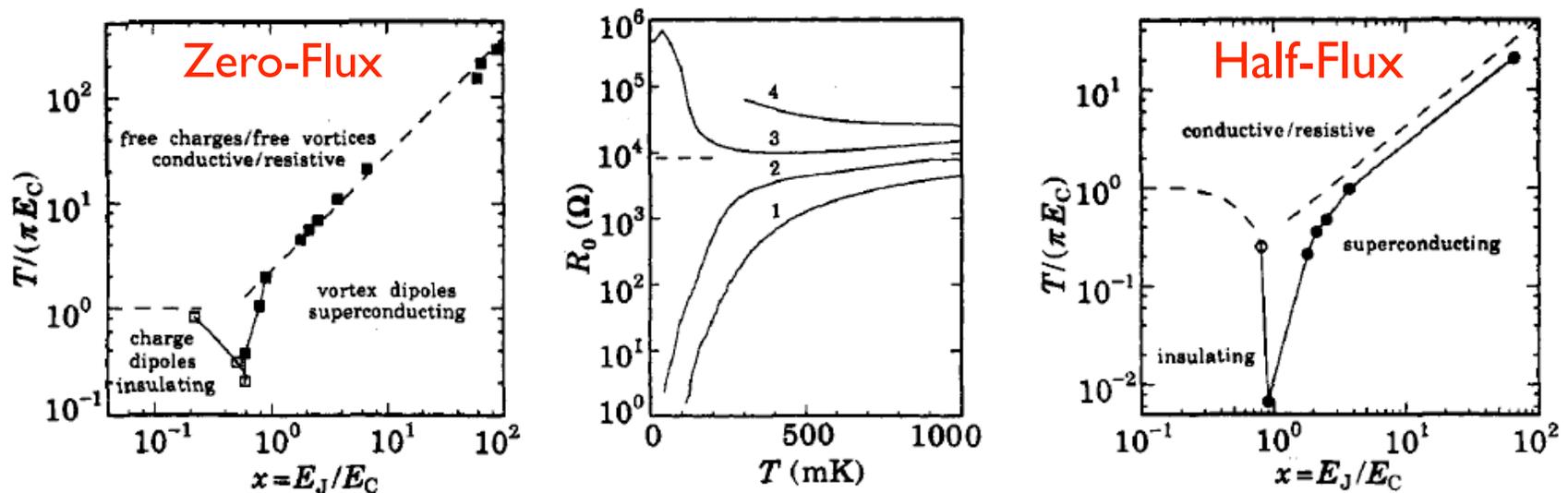
I. Classical variant ~ Fully frustrated Josephson junction array

Teitel and Jayaprakash (PRB 1983); M.Y. Choi and S. Doniach (PRB 1985)

Olsson (PRL 1995) - indicates **two** thermal transitions in 2d

What happens with increasing charging energy?

Fully frustrated Josephson junction array - experiment



Mooij group (EPL 1992, PRB 1996)

Why is this interesting or relevant?

I. Fully frustrated Josephson junction array - some theory

For p/q flux, we have q boson “flavors”

Generalize to an n -component field

- $n = \infty$ decoupling: **O(n) transition** *H. Kawamura - 1988, 1989*
- Assuming valid for $n=2$, expect $(d+1)$ XY criticality *Granato, Kosterlitz - 1990*

- Small system $L=20$ simulations indicate critical exponents \sim XY
- Universal conductivity at the transition for p/q seems to be roughly

$$\sigma_q^* = q \times \sigma_1^*$$

Granato, Kosterlitz - 1990; M.-C. Cha, Girvin - 1994

Epsilon expansion of $O(n)$ model hints at a **chiral fixed point** for

$$n > 21.8 - 23.4\epsilon \quad (H. Kawamura - 1988, 1989)$$

Large system [$L > 150$] numerical studies on stacked triangular lattice indicates a very **weak first order** transition (*H. Diep, PRE 2008, 2010*)

Why is this interesting or relevant?

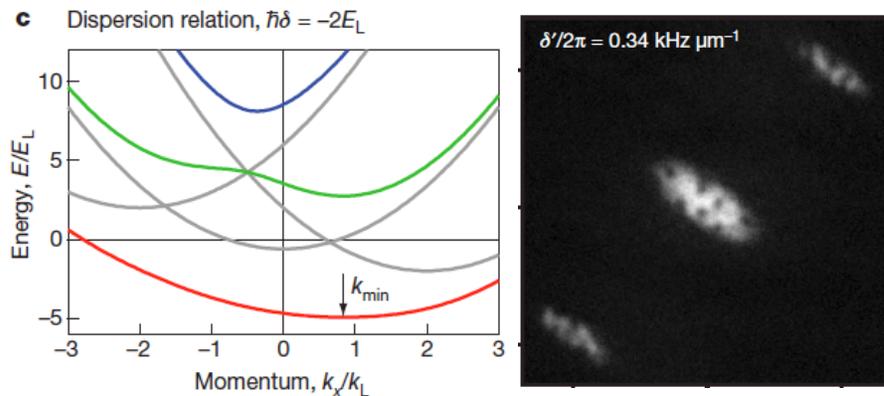
2. Cold atoms with frustration are beginning to be explored in the weakly correlated regime

Synthetic magnetic fields

Spielman, et al (Nature 2009)
Observation of few vortices

Higher band bosons

Hemmerich group (Nat. Phys. 2010)
Observation of “T-breaking” superfluidity



Frustrated hoppings

Sengstock group (Science 2011)
Observation of “120-magnetic order”
in triangular array of tubes

Natural extension is to explore Mott (correlated) regime

Why is this interesting or relevant?

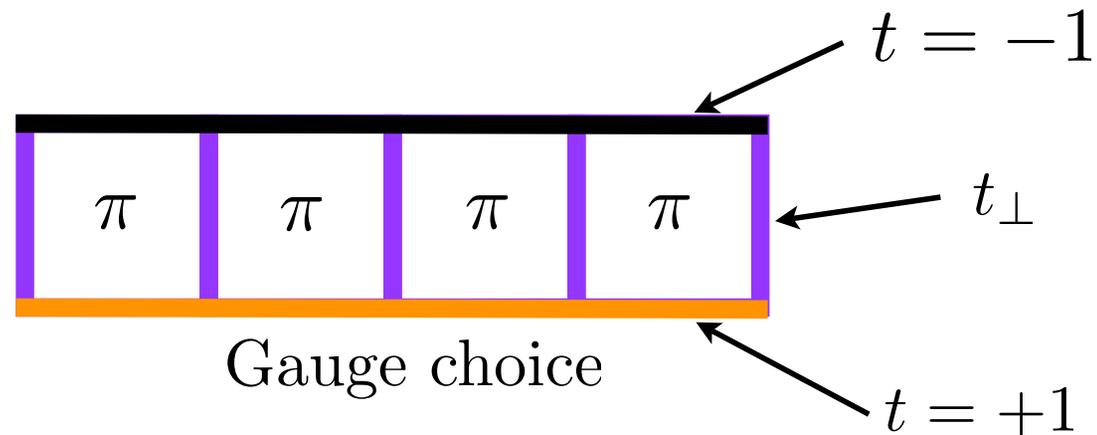
3. Field theories of multiple bosons arise in various different contexts

Eg: Deconfined critical point theories (XY variant) have multiple types of vortices (bosons) interacting with $U(1)$ gauge fields
Senthil et al (Science 2004)

What happens in the absence of gauge fields?

4. Unusual connection to frustrated polyelectrolyte hexagonal bundles with JJA phase \sim displacement of counterions along rodlike DNA (rod direction \sim imaginary time)
Grason and Bruinsma (PRL 2006)

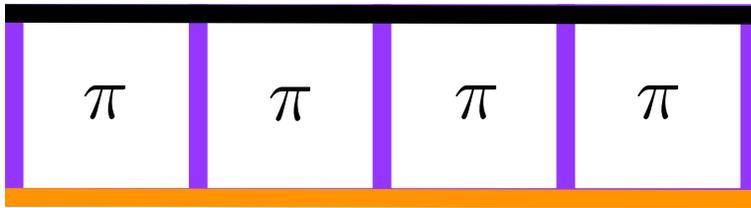
Fully Frustrated ladder



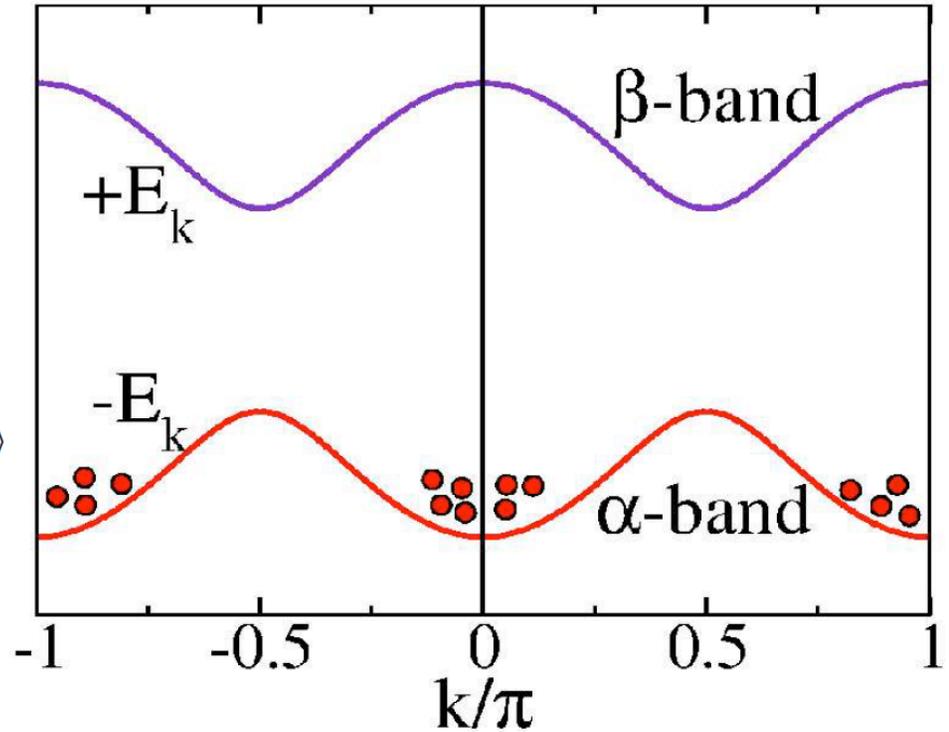
Study the phase diagram as a function of interactions and perpendicular hopping

- . Bilayer classical model (JJA or quantum FFXY)
- . DMRG study of bosons
- . Physical picture

Fully Frustrated ladder: Boson dispersion



$$|\psi\rangle = \frac{1}{\sqrt{N!}} \left[e^{i\varphi} (\alpha_0^\dagger + e^{i\theta} \alpha_\pi^\dagger) \right]^N |0\rangle$$



“Landau theory”

$$E_{\text{low}}^{\text{mft}} = (-E_0 - \mu) \sum_{i=0,\pi} |\varphi_i|^2 + U(u_0^4 + v_0^4) \sum_{i=0,\pi} |\varphi_i|^4 + 8Uu_0^2v_0^2|\varphi_0|^2|\varphi_\pi|^2$$

$$+ 2Uu_0^2v_0^2(\varphi_0^{*2}\varphi_\pi^2 + \varphi_\pi^{*2}\varphi_0^2)$$

Favors relative phase = $\pi/2$

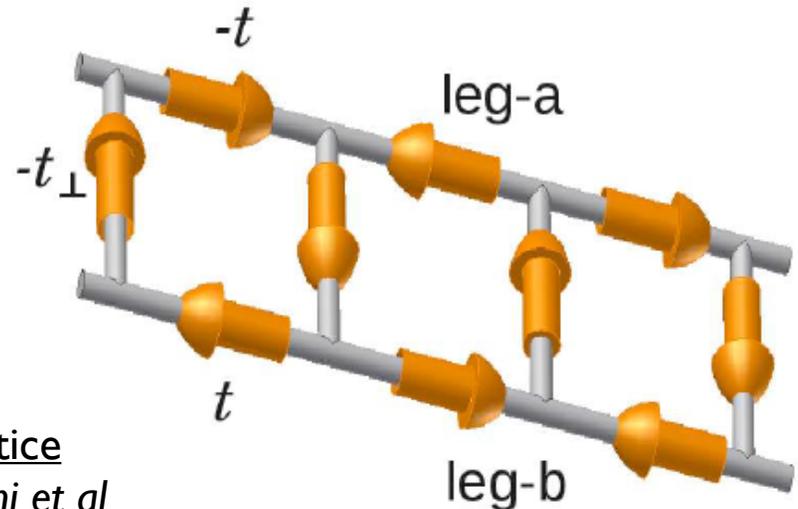
Favors equal amplitude

Fully Frustrated ladder

Mean field state for weak interactions

Resulting mean field superfluid has uniform density but staggered current pattern:

“Chiral” Superfluid (CSF)



Known from various earlier works on square lattice

Eg: L. K. Lim, C. M. Smith, A. Hemmerich (PRL 2008); M. Polini et al (PRL 2005); G. Moller and N. R. Cooper (PRA 2010); S. Sinha and K. Sengupta (EPL 2011); S. Powell, et al (PRL 2010)

Reminiscent of electronic d-density wave metal
proposed for cuprate pseudogap state

S. Chakravarty, R. Laughlin, D. Morr, C. Nayak (PRB 2001)
[Affleck, Marston staggered flux state (1988); Nori, Abrahams, Zimanyi (1990)]
Numerical studies of extended t_j models with many extra interactions
J. B. Marston, et al (PRL 2002); U. Schollwock, et al (PRL 2003); C. Weber et al (PRL 2009)

Fully Frustrated ladder

What happens for very strong interactions?

At integer filling and large U , the bosons must localize into a conventional Mott insulator

Mean field theory (2D: L. K. Lim et al, PRL 2008) suggests a direct continuous CSF-MI transition

MFT on ladder

$$\frac{1}{\sqrt{4t^2 + t_{\perp}^2}} = \frac{n}{\mu - U(n-1)} + \frac{n+1}{Un - \mu}$$

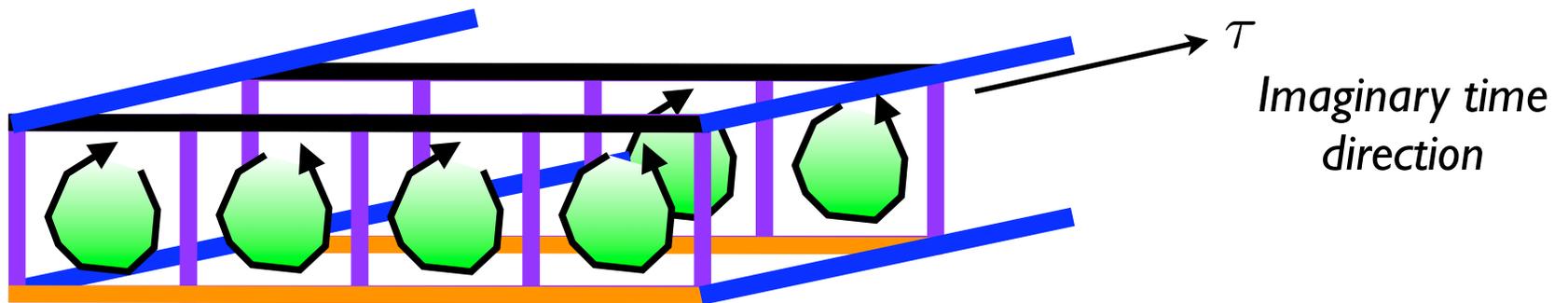
$$\frac{U^{\pi\text{-flux}}_{c,ladder}}{U^{0\text{-flux}}_{c,ladder}} = \frac{\sqrt{4t^2 + t_{\perp}^2}}{2t + t_{\perp}} < 1$$

Frustration makes it easier to undergo the Mott transition

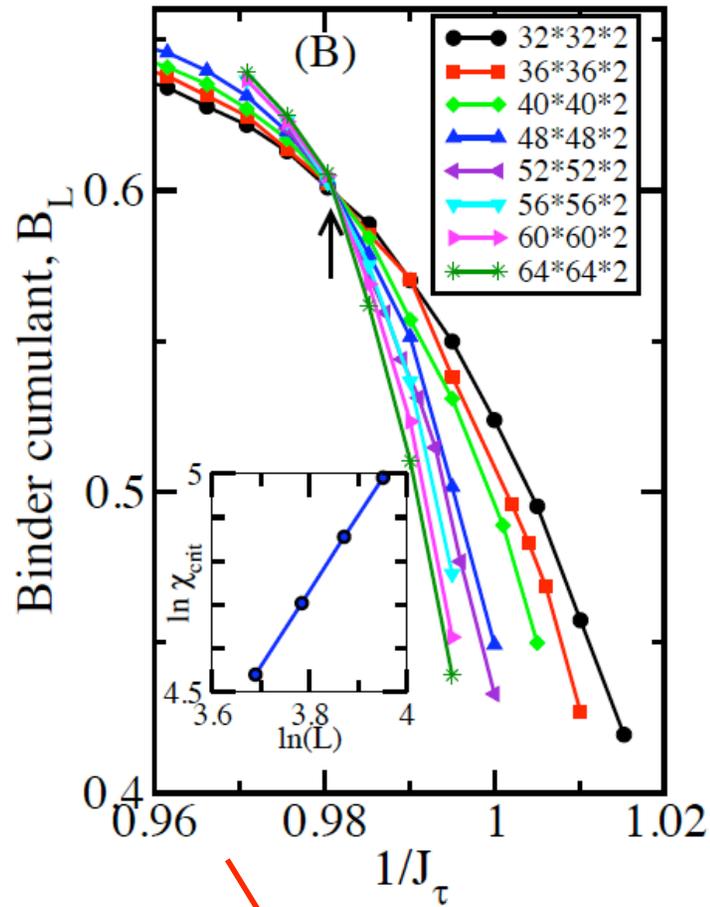
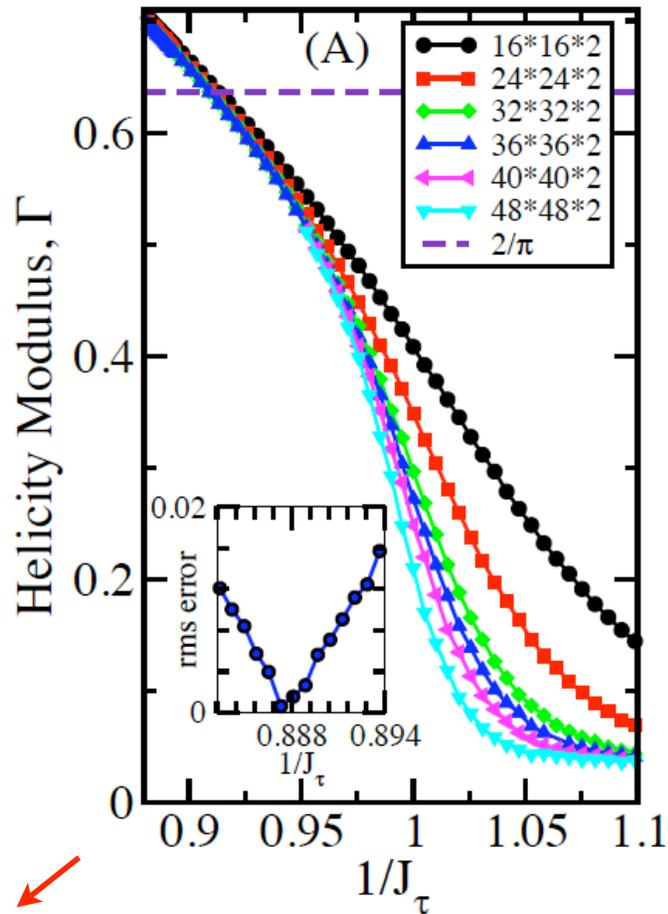
But the CSF breaks $U(1)$ [phase] and $Z(2)$ [Time reversal]
Do these get “unified”?

Fully Frustrated ladder Classical Bilayer Model

$$H_{XY} = - \sum_{i, \delta} J_{\delta} \cos(\varphi_i - \varphi_{i+\delta})$$



Fully Frustrated ladder Classical Bilayer Model

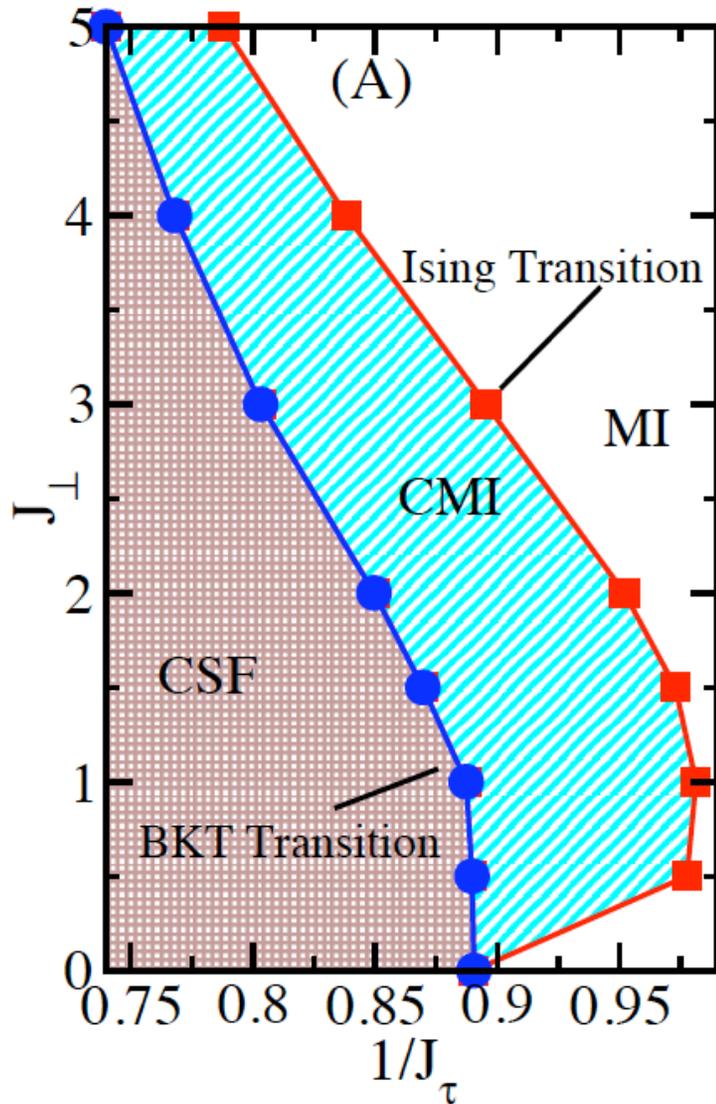


Helicity modulus: BKT transition
 Finite size scaling form: $\Gamma(L) = A \left(1 + \frac{1}{2 \log L + C} \right)$

Weber and Minnhagen (PRB 1988)

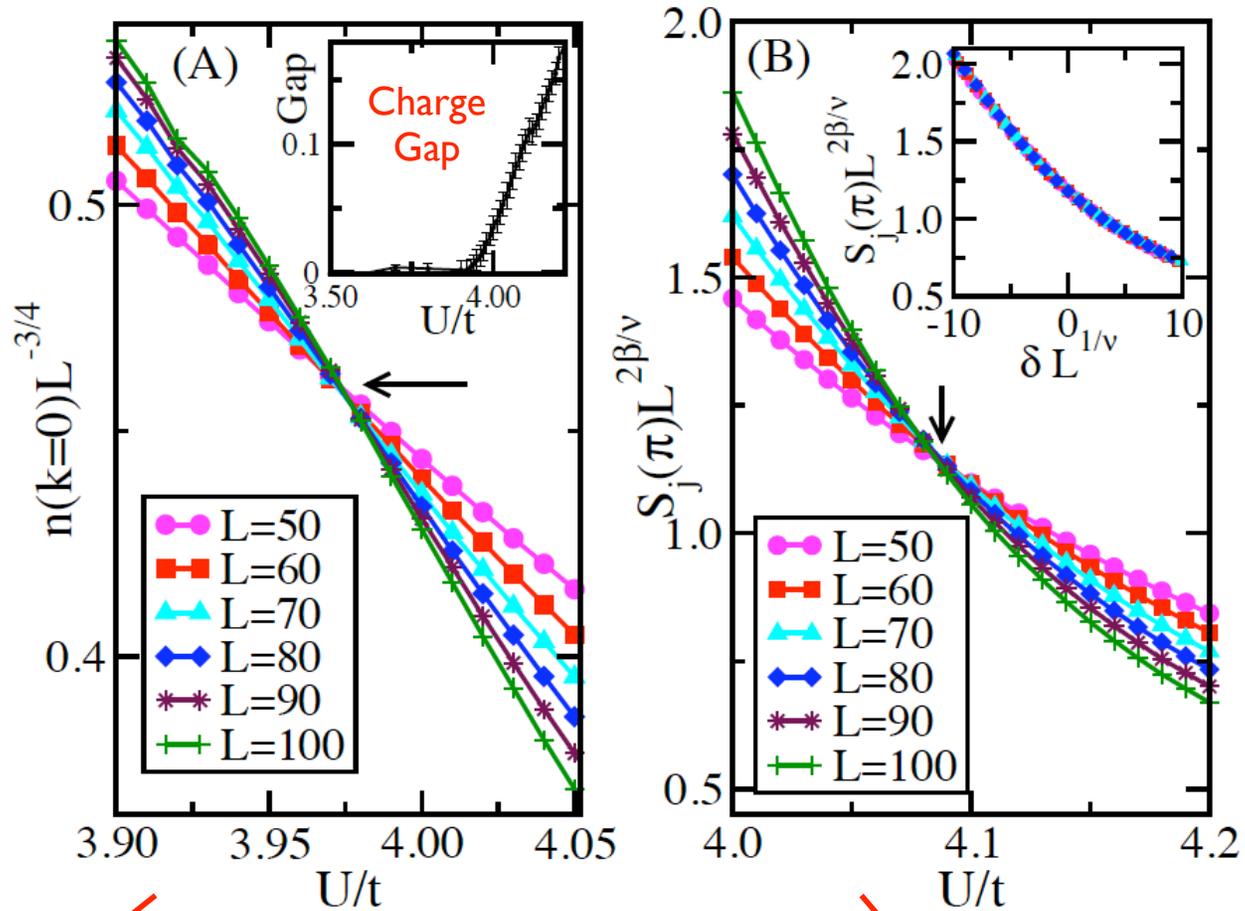
Binder cumulant for the staggered current order

Fully Frustrated ladder Classical Bilayer Model



- . CSF-MI transition splits into two transitions
- . Intermediate CMI phase
- . CMI=paramagnet with staggered currents on the ladder (interlayer)

Fully Frustrated ladder DMRG study of the Hubbard model

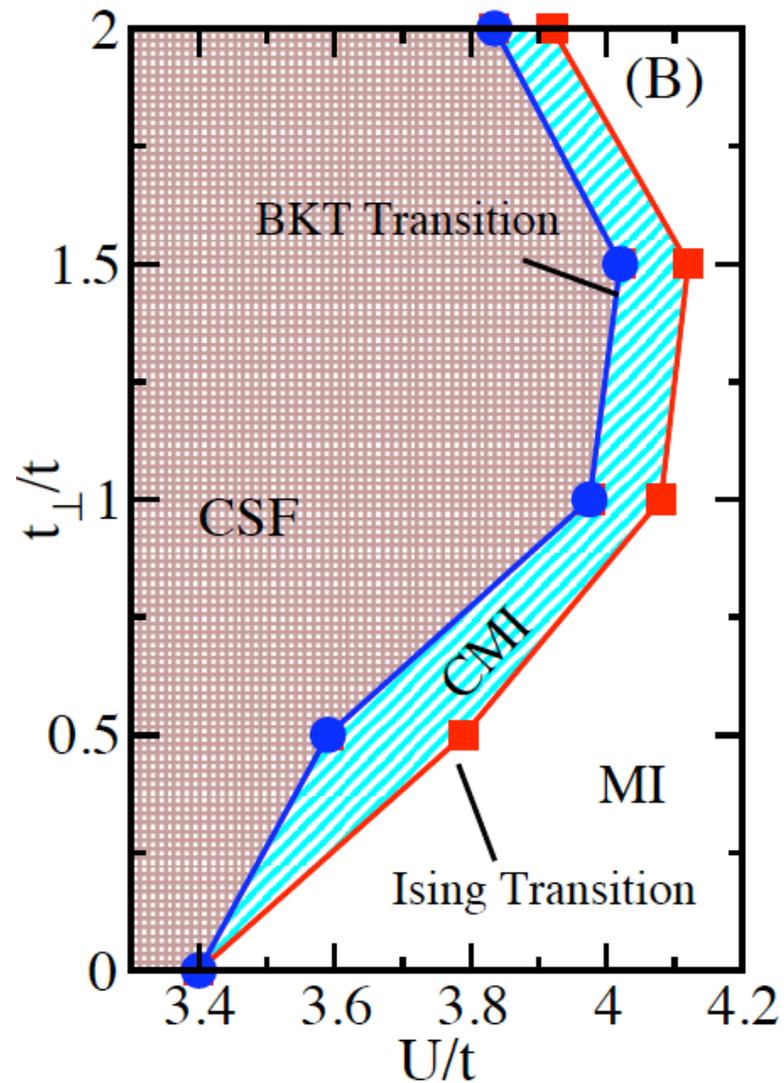


BKT transition implies universal scaling

$$n(k \rightarrow 0) \sim |k|^{-3/4}$$

Staggered Current
Order parameter - Ising scaling

Fully Frustrated ladder DMRG study of the Hubbard model

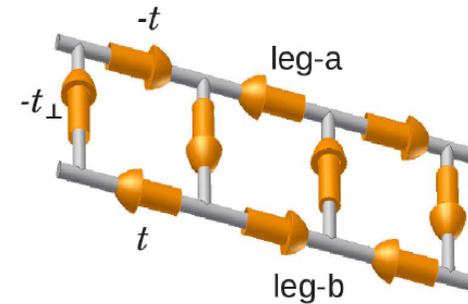


- . CSF-MI transition splits into two transitions
- . Intermediate CMI phase
- . CMI= insulator with staggered currents

Simple pictures for the “chiral” Mott insulator

Chiral superfluid = Vortex Crystal

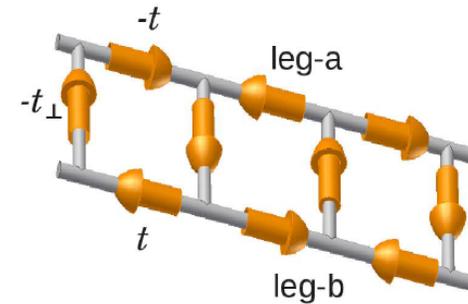
- . Flux nucleates vortex or antivortex
- . Vortex-vortex interaction is repulsive
- . Equal number of V/AV
- . “Antiferromagnetic” crystal



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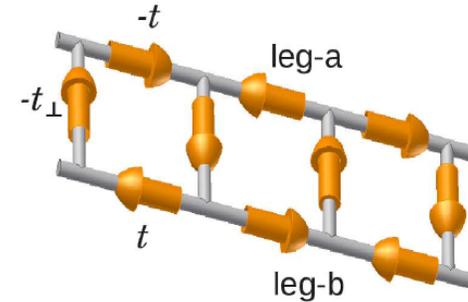
Regular Mott insulator = Vortex Superfluid (D. Haldane; Halperin/Dasgupta; Fisher/Lee)

- . Dual - proliferated quantum phase slips

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Chiral Mott insulator = Vortex Supersolid

- . Defect in crystal: Extra vacancy/interstitial vortex/antivortex
- . Proliferating and condensing dilute defects: Vortex superfluid
- . Background current pattern preserved: Vortex crystal

Simple pictures for the “chiral” Mott insulator

Excitations of a Conventional Mott insulator

- . Gapped Particles: “double occupancy”
- . Gapped Holes: “vacancy”
- . Dispersing particles/holes: Like a “semiconductor”

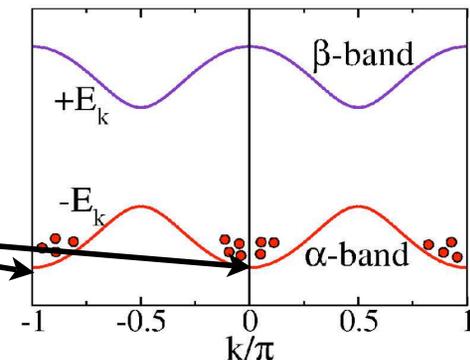
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Excitations of a Conventional Mott insulator **with flux**

- . Gapped Particles: “double occupancy”
- . Gapped Holes: “vacancy”
- . Dispersing particles/holes **with multiple minima**
Like a “semiconductor” with multiple valleys



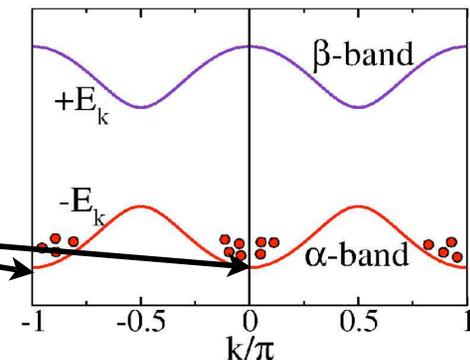
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Chiral Mott insulator: Indirect Exciton condensate!

Wavefunction for the “chiral” Mott insulator

$$\Psi(r_1, r_2, \dots, r_N) = e^{-\sum_{i,j} \tilde{v}(r_i - r_j)} \Psi_{MF}(r_1, r_2, \dots, r_N)$$

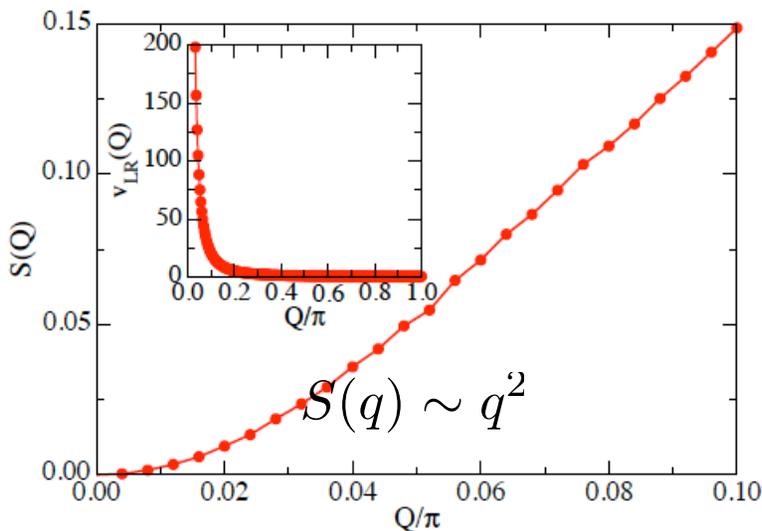
Pick mean field state to be the chiral Bose condensed superfluid

Choose a singular long range Jastrow: $v(q \rightarrow 0) \sim 1/q^2$

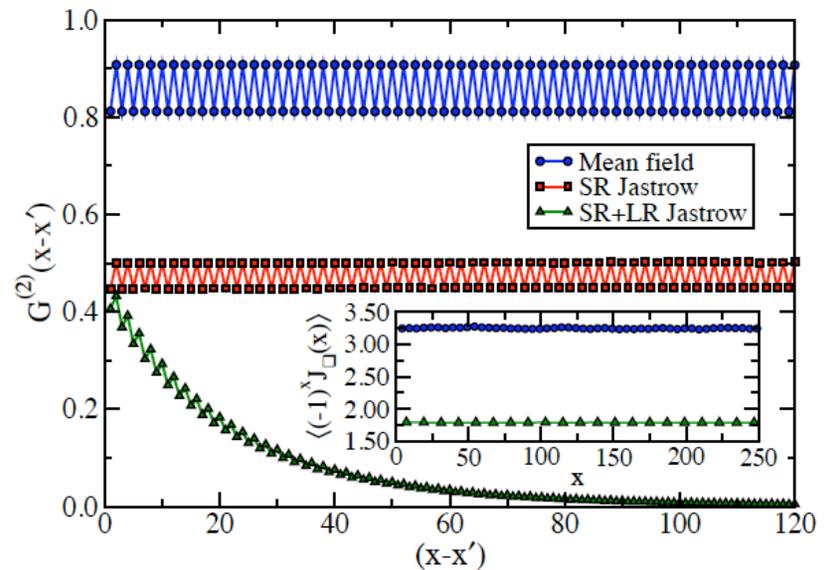
$$|\psi\rangle = \frac{1}{\sqrt{N!}} \left[e^{i\varphi} (\alpha_0^\dagger + e^{i\theta} \alpha_\pi^\dagger) \right]^N |0\rangle$$

Disorder Preserve

Study using variational MC



Density structure factor
suggestive of gap (Feynman single mode)



Boson density matrix
becomes short-ranged

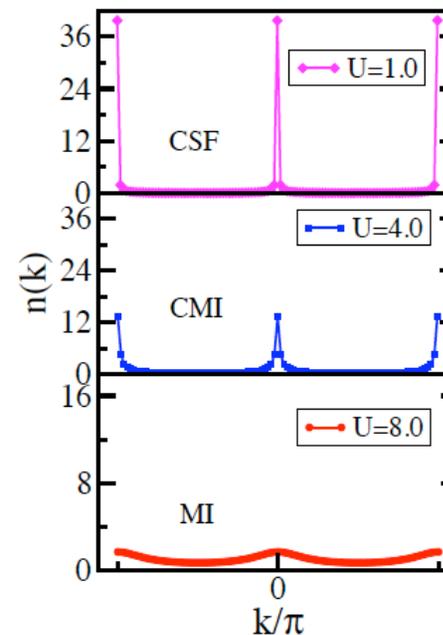
Experimental signatures

1. Josephson junction array realization:

- . Insulator in transport
- . Spontaneous staggered fields (SQUID microscopy)
 $J \sim 1K$ coupling; $1\mu m \times 1\mu m$ cell; $1nT$ fields

2. Cold atom realization:

- . No sharp peaks in $n(k)$
- . Look for residual interference between the two peaks



Summary

1. Fully frustrated Bose Hubbard model supports a **staggered current Mott insulator** on a ladder
2. Simple pictures for the “chiral” Mott insulator
3. Field theory understanding:
Coupled sine-Gordon model - In progress
4. Higher dimensional generalizations: In progress
5. Other models of frustration and their Mott limits?