Dynamics of dense soft colloids: shape and osmotic effects

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Outline:

Ingredients for (invariant) constitutive equations (link to microstructure) Softness, shape & entroping mixing: routes to tailor flow Molecular design of soft composites with tailored flow properties

Compare stars and microgels (with hard spheres)

Grafted cylinders vs. grafted spheres in dense state Packing and order control dynamics & flow

Soft spheres + big linear polymers From colloidal to polymeric response Polymer-mediated arrest of particles

Discussion with Wilson Poon (March 7)



<u>What is the role of softness?</u> ** Prabhu Nott (80 µm PDMS) DST, reduced effect

Guy et al. PRL 2015 Lin et al. PRL 2015 Krishnamurthy et al. JOR 2005

Main experimental test





Bonn et al. Rev. Mod. Phys. 2017 Christopoulou et al. Phil. Trans. R. Soc. A 2009

The "standard" picture: YSF

Memory not necessary Viscoplasticity



hysteresis (thixotropy), yield stress, wall slip, shear banding

Balmforth et al., Annu. Rev. Fluid Mech. 2014 Oppong, de Bruyn, Rheol. Acta 2011

Coussot et al., PRL 2002

Colloidal glass/paste vs entangled polymer



<u>Yielding in polymers?</u> Memory Viscoelasticity



S. Q. Wang, Macromolecules 2013 F. Snijkers, 2011

Topological constraints in small (nm scale) systems: caging



Edwards, deGennes, Doi, 1970s



MCT, SGR Götze, Cates, Fuchs, Fielding, Sollich



Frenkel, 1941; Pusey, 1991; Cates, 2003

Softness: from polymers to colloids



Softness: state transitions



Challenge: determine volume fraction

$$\phi_{eff}=c/c_{h}^{*}=nV_{0}$$

Poon et al. SM 2012 Conley et al. Sci. Adv. 2017 Bouhid de Auiar et al. Sci. Rep. 2017 van der Scheer et al. ACS Nano 2017

<u>Paddy Royall (March 9)</u> Gardner transition (thermal glass to Gardner phase to jammed glass)

Charbonneau et al., Annu. Rev. Condens. Matter Phys. 2017

Dense Soft Colloids: shape adjustment and interpenetration (with Michel Cloitre, ESPCI)

Microgels and similar



Microgels Emulsions Vesicles

Stars and similar











Cloitre & Bonnecaze, Adv. Polym. Sci. 2010 Seth et al., Nat. Mat. 2011 Helgeson et al., Nat. Mater. 2012 ; JOR 2014 Peng et al., Nat. Mat. 2015 Mohanty et al., PRX 2015; Sci. Rep. 2017 Conley et al., Sci. Adv. 2017 Semenov et al., Langmuir 1999 DV & G. Fytas, Adv. Polym. Sci. 2010 Zhang et al., PRL 2014 Senff, Richtering, Langmuir 1999

Compare stars and microgels

Star polymers in squalene

Number and size of arms

- 800 arms
- Each arm 5kg/mol ($M_w/M_n < 1.1$)
- Carbosilane dendrimer core
- Polybutadiene arms
- Radius 25-30nm



Microgels in water

Number and distance between crosslinks

- Spherical cross-linked polymer network with 128 units/crosslink
- 35% methacrylic acid
- 64% ethylacrylate
- 1% difunctional crosslinker
- No dangling chains
- Radius 305nm (swollen)

Roovers et al., *Macromolecules* 1993 Gauthier and Munam, *Macromolecules* 2010

Pellet and Cloitre, Soft Matter 2016

Partial overview of entropic mixtures



Softness: from repulsive glass to gel





Truzzolillo et al. Macromolecules 2011



Polymer mediated melting of soft glass



Stiakakis et al. EPL 2005 ; JPCM 2011

Aging



microgels, hard spheres

some stars

Aging







<u>Temperature as control parameter</u>







scale with R_h

 $\varphi_{eff}=c/c*_{h}$

<u>fragile vs strong analogy</u>: Mattsson et al. Nature 2009 van der Scheer et al. ACS Nano 2017



Cloitre et al. PRL 2003 Seth et al. Nat. Mat. 2011



Rogers et al. PRL 2008, JOR 2010

Also, Sid Yip: simulations on metallic glasses



Soft glasses with detactable alpha relaxation

Erwin et al. JOR 2010



 $\mathbf{\Pi} \approx kT\boldsymbol{\xi}^{-3}$

Erwin et al. SM 2010



Divoux et al. SM 2011 Ovarlez, Manneville

carbopol slips

PBD stars in squalene do not slip @ stainless fixtures



Similar scaled flow curves: microgels, micelles, stars, but not hard spheres?



Hard spheres band – when they do not slip ; repulsive, attractive (Poon, Cates, Petekidis, Dhont) Stars may band -- role of transient forces ? (Briels, Dhont, JCP 2010; PRF 2017 ; Rogers et al. PRL 2008) Repulsive microgels do not band, attractive do (Cloitre)

Yield stress and gradient banding scenario



Stars Weak dependence Strong shear thinning



Repulsive microgels do not band Attractive mirogels band (Cloitre) Link engagements of arms to attractions

Banding: Olmsted, Fielding, Dhont

Linear viscoelastic spectra

Microgels (c*=1% wt)

Stars (c*=16.5% wt)



Moduli compare to HS glasses (which exhibit stronger φ–dependence) Higher for stars can reach melt!

Plateau modulus



	A [kPa]	C* [%wt]	C _m [%wt]
Stars	24.7	15.6	9.2
Microgels	0.8	1	1.5

Plateau modulus



(Eq. HS volume fraction=1 corresponds to the close packing)

<u>Yielding: start-up in simple shear</u> <u>Influence of waiting time (at 0.2s⁻¹)</u>

Microgels 5% wt (5c*)

Stars 75% wt (4.5c*)



Helgeson et al. JOR 2007 Pham et al. EPL 2006, JOR 2008 Koumakis and Peterkidis, SM 2011

Stars exhibit two peaks: gradual yielding

Peak strain in entangled polymers



Costanzo et al., Macromolecules 2016

Yielding: start-up in simple shear Influence of shear rate for stars c=75% wt



Yield stress peaks reflect dual nature of particles (colloidal, polymeric)? Koumakis et al. SM 2011 Helgeson et al. SM 2014 Kim et al. JOR 2014

Yield strain for microgels and stars

Microgels





Dhont et al. PRF 2017

Star and microgel flow curves: yield stress



Velocity profiles in microgel suspensions: wall slip



To avoid need surface treatment: Seth et al. SM 2012 ; Cloitre and Bonnecaze, Rheol. Acta 2017

Flow instabilities at higher concentrations: microgels





Conclusions I: soft colloids are tunable

Soft (stars, microgels) vs Hard (HS) colloids in glassy state:

Jamming ambiguous in soft ; elastic interactions important <u>Yielding</u> is gradual, depends on internal microstrure <u>Shear banding</u>: present in HS and stars, not in microgels <u>Wall slip</u>: mainly microgels, HS (can control) <u>Flow curves</u>: universal scaling for soft (HB exponent of ½) More than one time scale (η/G) (in stars) Osmotic interactions control flow of mixtures: particle compression and depletion

Expect weaker shear thickening Explore analogies to molecular glasses

Soft particles at contact: engagements, facets, bonding Hard particles at contact: friction Hydrodynamics

Shape matters

<u>isotropic</u>

<u>crystal</u>



concentration

Alder, Pusey, Lekkerkerker, Frenkel, Dhont, Philipse, Fraden, Dogic

Soft colloidal systems of varying shape



Cellular actin filaments

Particle shape





Leibler, Fredrickson, Matsen

Polymer Grafted Nanoparticles with controlled microstructure



poly(3-(Triethoxysilyl)propyl methacrylate)

<u>Shell: PS</u>



PGNOs	core			M,	Mn of PS hairs		graft density	
	shape	size ^a (nm)	wt% core	(kg/mol) (chains/r		chains/nm ²	2)	
lam-PS111k	lamella	16.0	21.1		111		0.118	
cyl-PS111k	cylinder	27.0	21.1		111		0.101	
sph-PS111k	sphere	30.0	21.1		111		0.075	

Same Mw per arm

Similar grafting density

TEM and Light Scattering characterization



RH=49 nm (Light Scattering) PDI=1.04 (TEM) Mw=3x10⁷ g/mol f=212 C*(RH)=0.10 g/mL Lн=630 nm (Light Scattering) Lp=775 nm (Light Scattering) Aspect ratio (L/D)= 8 PDIlength=1.3 (TEM) Mw=6x10⁹ g/mol f=3.6E4 C*(Lн)=0.026 g/mL

Polymer-grafted nanoparticles in diethyl phthalate (good solvent) at 13%wt

$$\varphi_{eff} = \frac{c}{c^*} = 1.3$$

$$\phi_{eff}=\frac{c}{c^*}=5.0$$

SAXS & SANS at the same mass fraction (13wt%)



short rods, AR~6



long rods, AR~100



Cubatic phase - glass

Nematic glass – orientational texture

Dhont, Kang PRL 2013 ; SM 2013

LVE spectra at the same mass fraction (13wt%)

T=25°C 10^{4} DFS + Creep $\phi_{eff} = \frac{c}{c^*} = 1.3$ 10³ $\varphi_{\rm eff} = \frac{c}{c^*} = 5$ **G**"~ω G',G" [Pa] 10^{2} **10**¹ 10⁰ **Spheres** $G' \sim \omega^2$ Cylinders **10**⁻¹ $10^{-6} 10^{-5} 10^{-4} 10^{-3} 10^{-2} 10^{-1} 10^{0} 10^{1} 10^{2}$ ω [rad/s]

Soft glasses with detactable alpha relaxation

8mm Cone +Peltier on Ares Solution (LVE-NLVE) and Anton Paar MCR 702 (Creep) Rejuvenation (DSST) +aging (DtST 30 min) before any measurements

Shape-independent high-frequency dynamics



Trappe, Weitz, PRL 2000

Dynamics of spheres and lamellas: caging effect



Cylinders



thickness: 104 nm





Dynamics of cylinders: fluctuations and tube renewal

13wt%



Cylinder fluctuations inside the tube

$$\tau_{\beta}^{LVE} = \frac{1}{\omega(G_{min})} = 1s$$

$$\mathbf{D}_{\perp} = \frac{(R_{core+shell})^2}{\tau_{\beta}^{LVE}} = 4 * 10^{-15} \frac{m^2}{s}$$

Tube renewal

$$\tau_{\alpha}^{LVE} = \frac{1}{\omega(G_c)} = 67s$$

$$\mathbf{D}_{||} = \frac{(L_H)^2}{\tau_{\alpha}^{\text{LVE}}} = 6 * 10^{-15} \frac{m^2}{s}$$

Perpendicular and parallel diffusion coefficients are very similar at this concentration: weak/absent shell interpenetration

Edwards, Evans, J. Chem. Soc., Far. Trans. 1982 Doi, Edwards, The Theory of Polymer Dynamics, 1986 Teraoka, Hayakawa, JCP 1988

Zhao, Wang, Polymer 2013

Dynamics of cylinders: fluctuations and tube renewal

20wt%



Cylinder fluctuations inside the tube

$$\tau_{\beta}^{LVE} = \frac{1}{\omega(G_{min})} = 12s$$
$$\mathbf{D}_{\perp} = \frac{(R_{core+shell})^2}{\tau_{\beta}^{LVE}} = 2 * 10^{-16} \frac{m^2}{s}$$

Tube renewal

$$\tau_{\alpha}^{LVE} = \frac{1}{\omega(G_c)} = 7692s$$

$$\mathbf{D}_{||} = \frac{(L_H)^2}{\tau_{\alpha}^{\text{LVE}}} = 6 * 10^{-17} \ \frac{m^2}{s}$$

The parallel diffusion coefficient is now slower: shell interpenetration

Edwards, Evans, J. Chem. Soc., Far. Trans. 1982

Doi, Edwards, The Theory of Polymer Dynamics, 1986

Teraoka, Hayakawa, JCP 1988

Zhao, Wang, Polymer 2013

Cylinders: LVE from 10 to 13 to 20%wt.





Cylinders: LVE from 10 to 13 to 20%wt.







Some consequences of softness: viscosity & self-diffusion



scale with R_h

 $\phi_{eff}=c/c_{h}^{*}$

COCIS 2014 ; Adv. Polym. Sci. 2010 ; Phys. B 2001

Nonlinear rheology: start-up shear







Comp

۱D

Ext

10°

Koumakis et al. JoR 2017



Shape shifts pear to higher Pe



Yield strain: Compare cylinders at different concentrations



Conclusions II

Shape effects:

Grafted cylinders: tube-renewal and re-organization, 2 times Yielding reflects different packing efficiency (small aspect ratio) Large yield strain

Measure N1 and N2 Elongational flows