

# Limiting lifetimes for localized spins in GaAs

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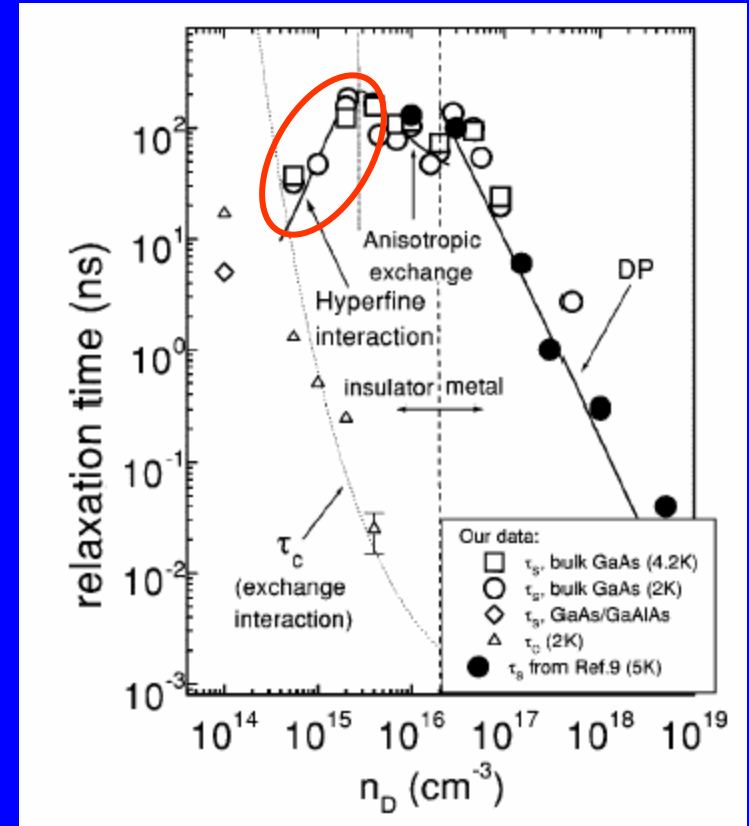
Acknowledgements to Wei Hao (SEAP 2005)  
& Josh Caldwell

# Outline

1. Background
  1. Long lifetimes for spins in non-magnetic semiconductors
  2. Work with donor ensembles
2. Magnetic resonance detected through Faraday Rotation
  1. Method
  2. Nuclear effects
  3. Probe light effects
  4. Higher concentration sample
  5. g-factor and Lifetime
3. Summary and Future Work

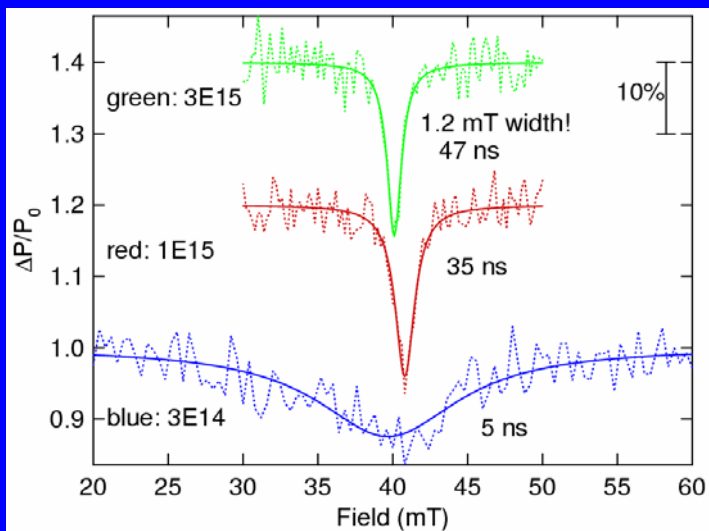
## Background

- Long lifetimes in non-magnetic semiconductors
  - Kikkawa and Awschalom, PRL and Science
- Move to localized spins—qubits and quantum information
  - In III-V's—GaAs donors provide a good ensemble
  - Metal-insulator transition is at  $2 \times 10^{16} \text{ cm}^{-3}$
  - Isolated donors for  $N_D < 4 \times 10^{14} \text{ cm}^{-3}$  for  $B = 0 \text{ T}$

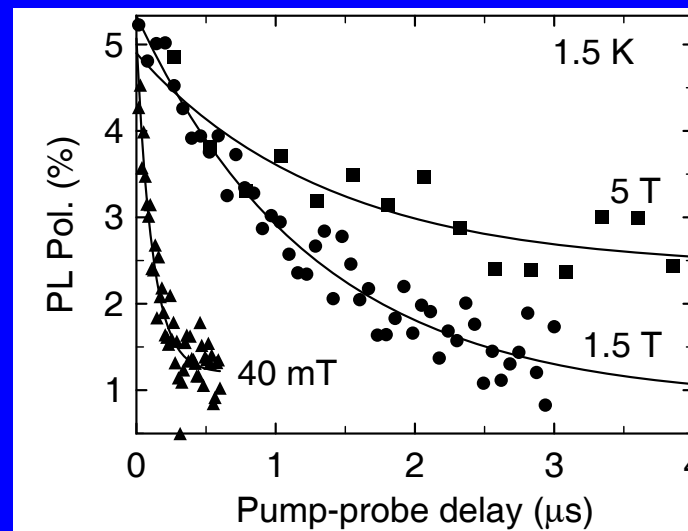


Dzhioev *et al.*, Phys. Rev. B 66, 245204 (2002)

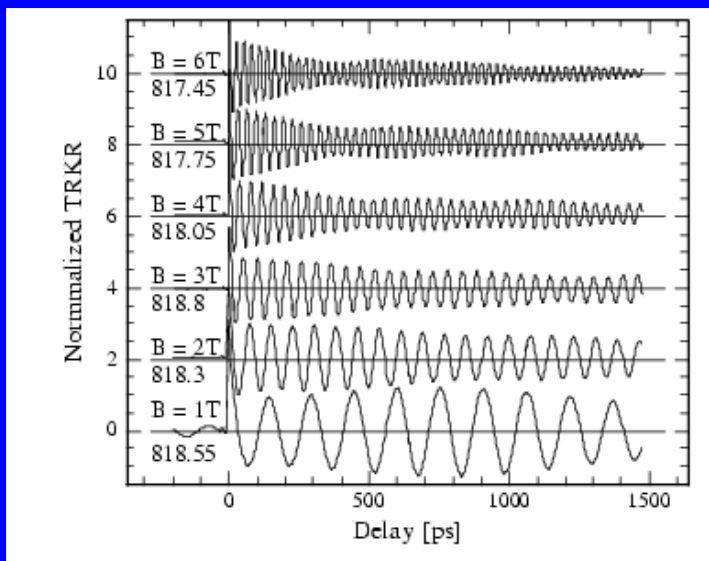
# Spin lifetimes for donors in GaAs



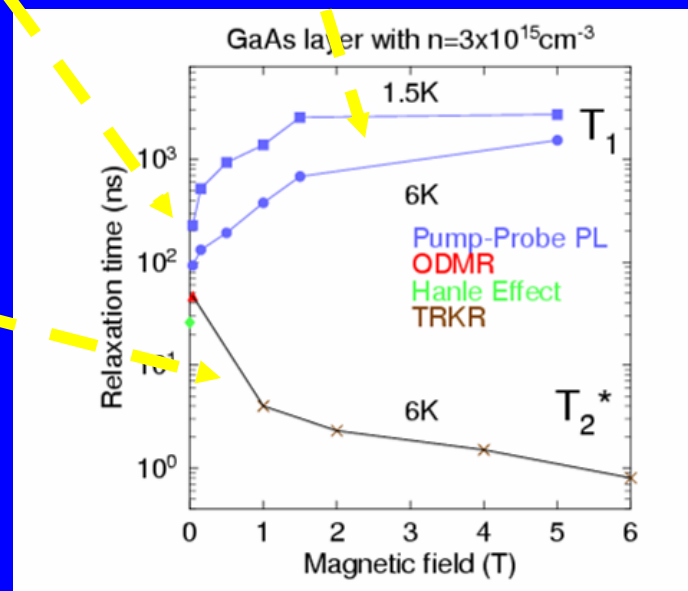
J.S. Colton *et al.*, SSC 132, 613 (2004)



J.S. Colton *et al.*, PRB 69, 121307 (2004)



Michael Scheibner *et al.*, to be pub.



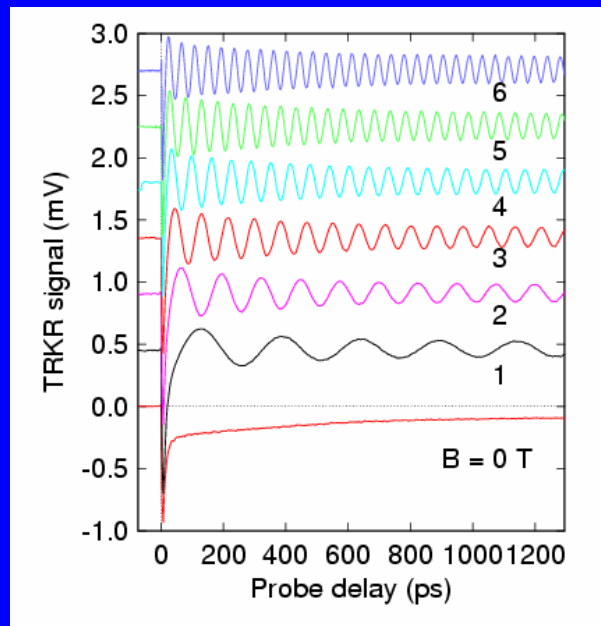
JSC:  $T_1 = 24 \mu\text{s}$  for 1E15 at 7T

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## Previous work leads to

Wide Quantum Well—  
Simplified energy levels  
through HH/LH splitting

Goal of Measuring  $T_2$ —Use  
pulsed magnetic resonance



### Two challenges

- Low spin concentration
- Well defined nuclear spin state--unpolarized

T.A. Kennedy *et al.*, PRB 73,  
045307 (2006)

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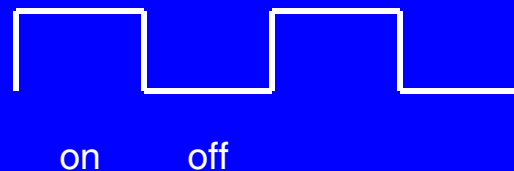
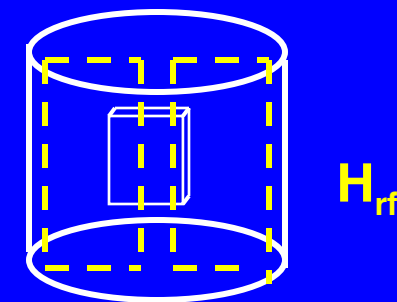
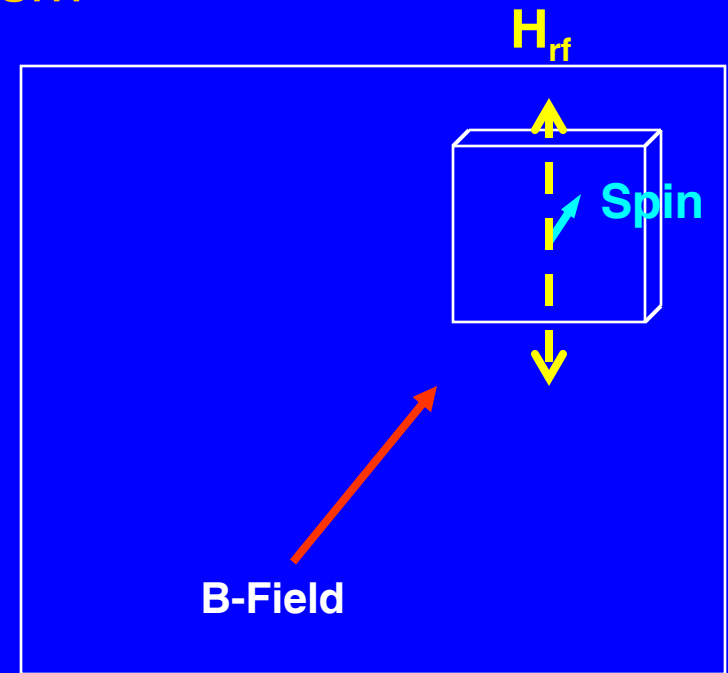
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## Background to This Approach

- Faraday Rotation and EPR
  - R. Romestain (1980): donors in CdS
- GaAs donors
  - Seck, Potemski and Wyder (1997): EPR—DNP enhanced
  - Karasyuk et al. (1994): Donor-bound excitons
  - Kai-Mei Fu (Stanford), Spin-flip Raman Scattering
- Nuclear Effects: ESR pinning in 2DEG's
  - Olshanetsky et al., Physica B **373**, 182 (2006)
  - Hillman and Jiang, PRB **64**, 201308 (2001)
  - Dobers et al., PRL **61**, 1650 (1988)
- ESE in GaAs
  - Petta et al. Science **309**, 2180 (2005)— $T_2$  of 1 to 10  $\mu\text{s}$
  - Loss (Basel), Sham (UCSD), Whaley (UCB), Das Sarma (Maryland), ETH (Zurich), SUNY (Buffalo) and others

## ODMR mechanism

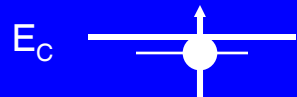
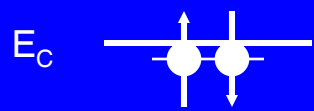
- Polarization
  - Thermal:  $B=6\text{ T}$  and  $T=1.5\text{ K}$ —  
 $g\beta H \sim kT$
  - Off-cycle of microwaves long with respect to  $T_1$
  - Thermalized electron spins:  
 $\langle S_T \rangle$
- Resonance with microwaves reduces the polarization
  - $h\nu = g\beta B$
  - Frequency  $\nu = 35\text{ GHz}$
  - Decreases the  $\langle S \rangle$  from  $\langle S_T \rangle$
- Detection...



35 GHz  
cavity



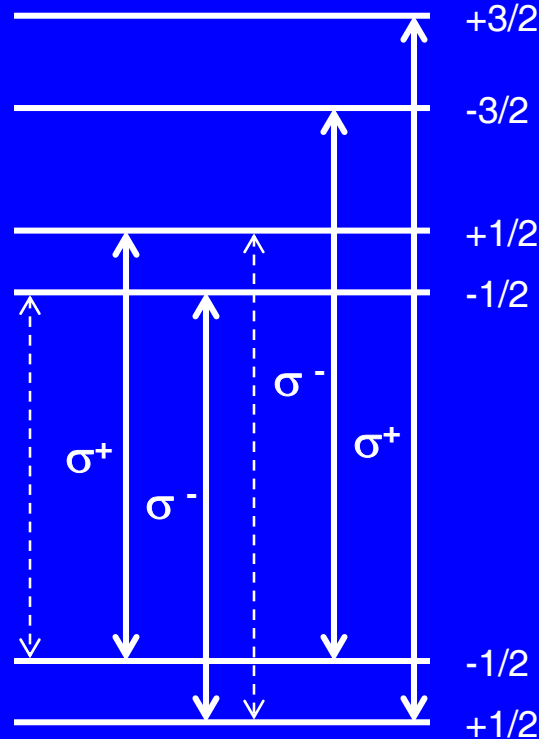
# Detection by Kerr Rotation



Single-Particle States

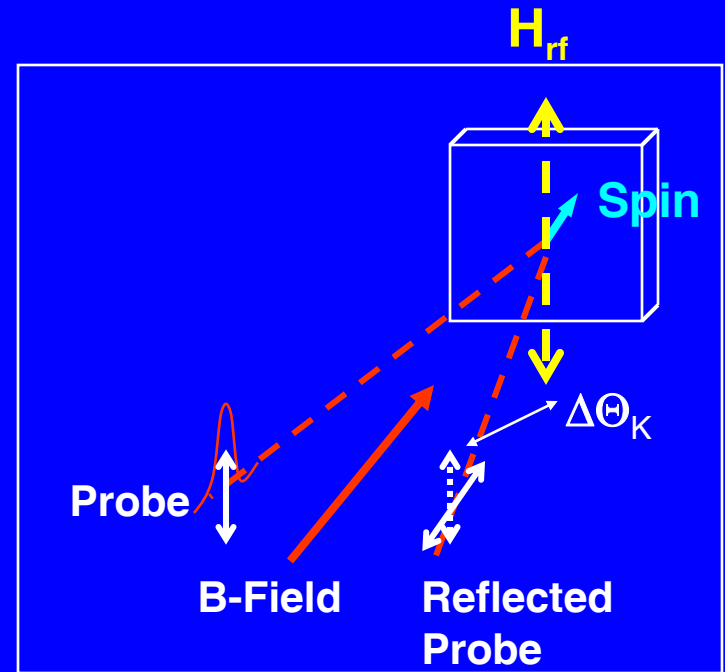
$D^{0X}$

$D^0$



Simplified Optical Transitions

cf. Karasyuk et al., PRB 49,16381 (1994)

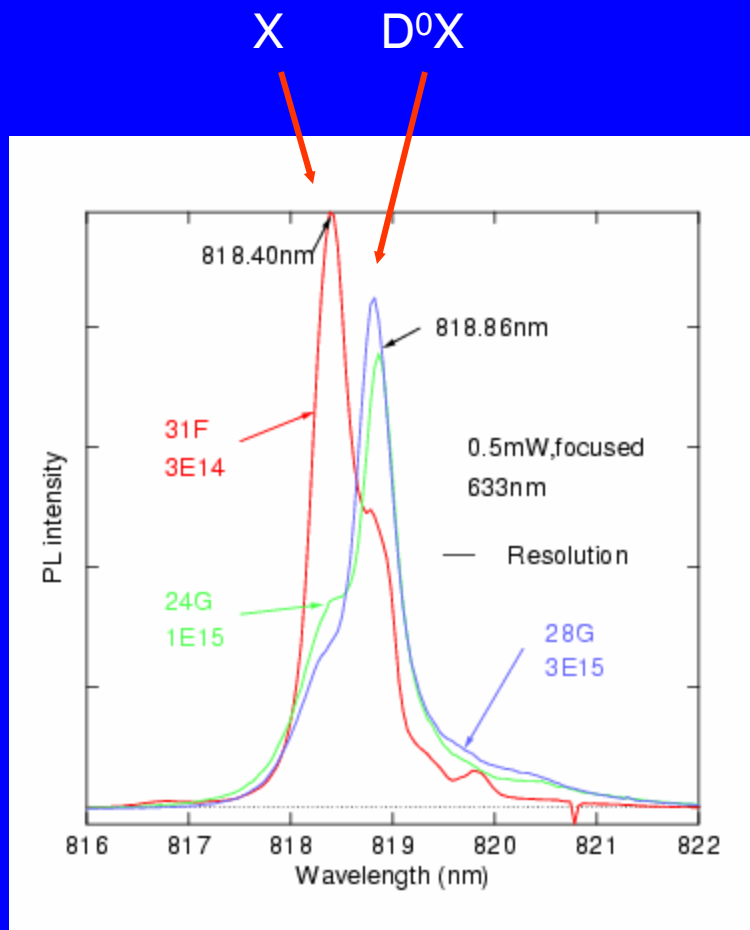


$$\Delta\Theta = -\pi(d/\lambda)(n_+ - n_-)$$

(in transmission)

Kerr rotation is sensitive to  $D^0$  spin through SO-coupling of hole in  $D^{0X}$

## Experimental Details



- Samples
  - 1  $\mu\text{m}$  GaAs surrounded by doped and undoped AlGaAs for flatbanding
  - Concentrations of  $3E14$ ,  $1E15$  and  $3E15$
  - (Metal-insulator transition at  $2 \times 10^{16} \text{ cm}^{-3}$ )
- Equipment
  - Oxford 7 T magneto-cryostat
  - Spectra Physics ps Ti:Sa laser ( $\Delta\nu \sim 1 \text{ meV}$  (0.5nm))
  - Agilent 250kHz to 40GHz signal generator

# Outline

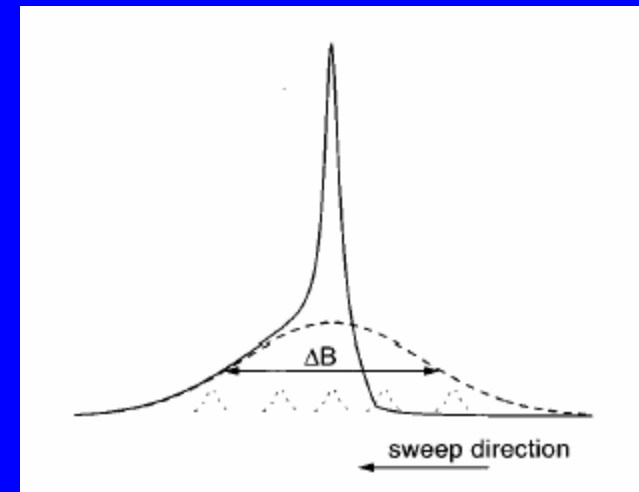
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## Nuclear Spin Effects

- Without nuclei, the Bloch equations describe the saturation of the ESR amplitude ( $A$ ) with microwave power ( $P$ ):

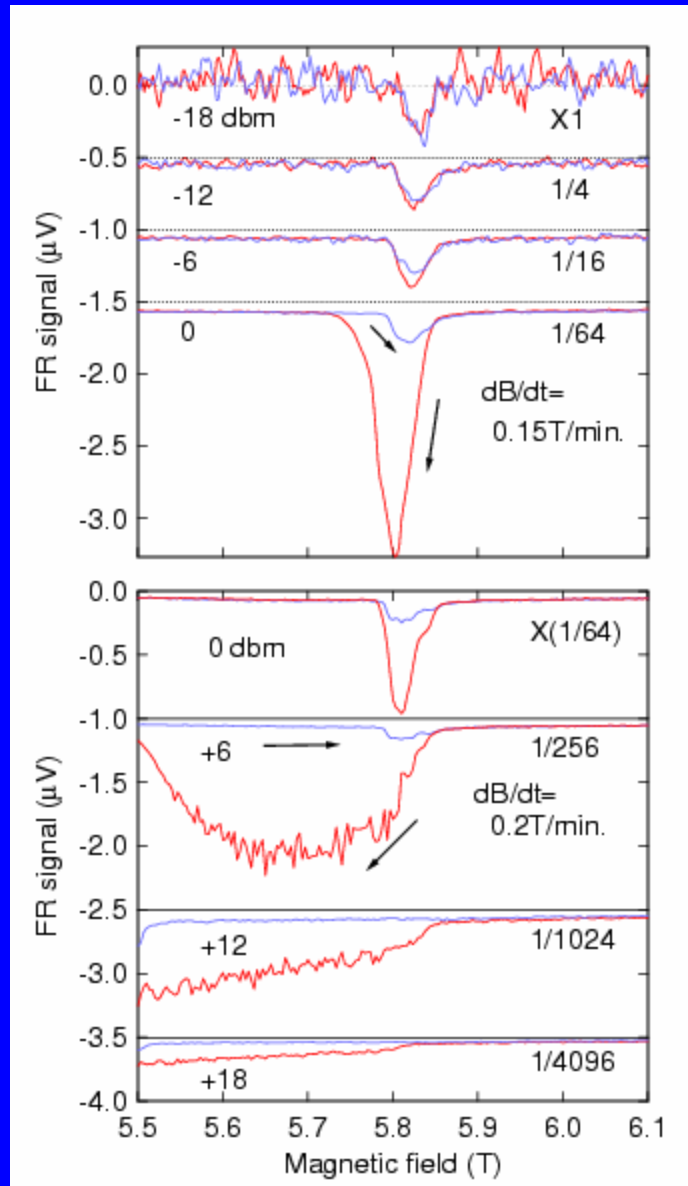
$$A = (\alpha P) / (1 + \beta P)$$

- With nuclei, Dynamic Nuclear Polarization arises from the Overhauser effect: Saturated e-spins try to relax through the nuclear spins using the  $(I^+ S^- + I^- S^+)$  part of the hyperfine interaction
- Sign of the nuclear field:  $B_N$  adds to the external field ( $B_{\text{ext}}$ )
  - $\langle I \rangle = [ I (I + 1) / S (S + 1) ] * [ \langle S \rangle - \langle S_T \rangle ] < 0$
  - $B_N = A \langle I \rangle / g_e \beta ; g_e < 0; B_N > 0$
  - $h\nu = g\beta (B_{\text{ext}} + B_N)$
- Enhancement and pinning of the ESR can occur:
  - Compare  $dB_N/dt$  with  $dB_{\text{ext}}/dt$
  - $\text{DNP} \propto P$
  - Nuclear relaxation time  $\sim 1$  minute
  - Strong effects occur for downward field sweeps



Seck et al., 1997

# Dependence on Microwave Power



$$dB_N/dt \propto \text{Saturation} \propto P$$

- Unsaturated limit

- Down & up are the same
- Amplitude  $\propto$  power

- DNP enhanced:  $dB_N/dt \approx | -dB_{ext}/dt |$

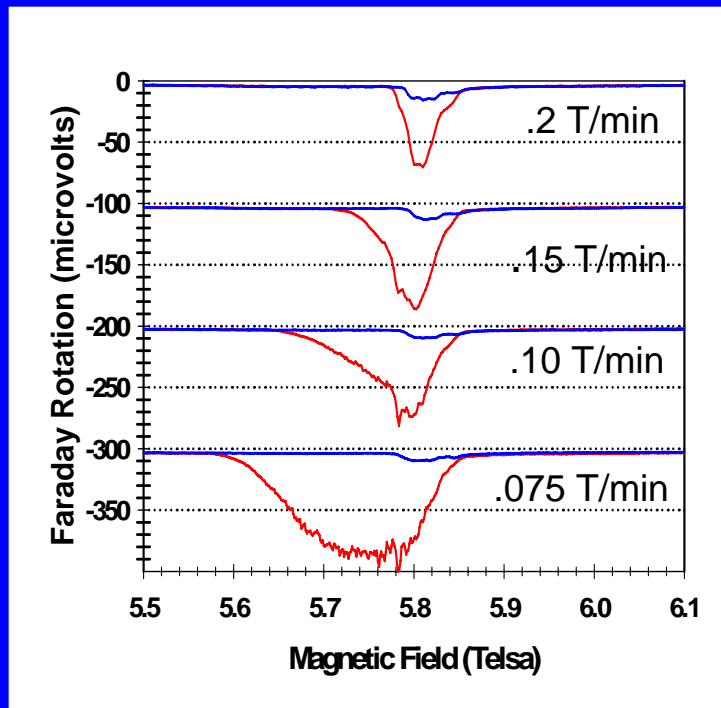
- Strong for down sweeps
- Broadened & shifted
- Up sweeps unaffected

- ESR pinning:  $dB_N/dt > | -dB_{ext}/dt |$

- Resonance not achieved
- $B_N$  becomes  $> 0.3$  T

817.8 nm (near res.), 9/20 & 9/19 data, 31F, T = 1.5 K, 35 GHz

## Dependence on Rate of Change of External Magnetic Field



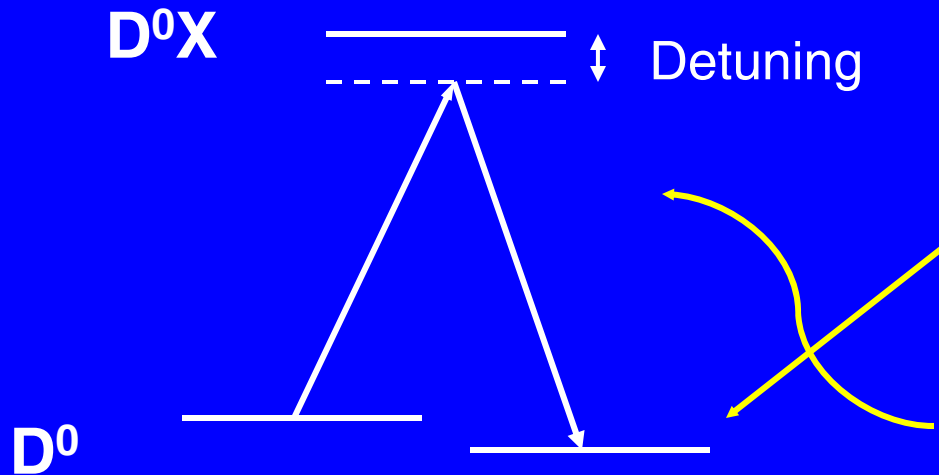
- Dynamics is again controlled by the relative size of  $dB_N/dt$  and  $-dB_{ext}/dt$
- Rate of DNP induced  $B_N$  is constant since microwave power (  $P$  ) is constant.
- Decreasing the sweep rate of magnetic field ( $-dB_{ext}/dt$ ) changes the response from DNP enhanced to ESR pinning.

- 0 dbm, 9/17/05, 817.8 nm

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# Effects of the Probe Light

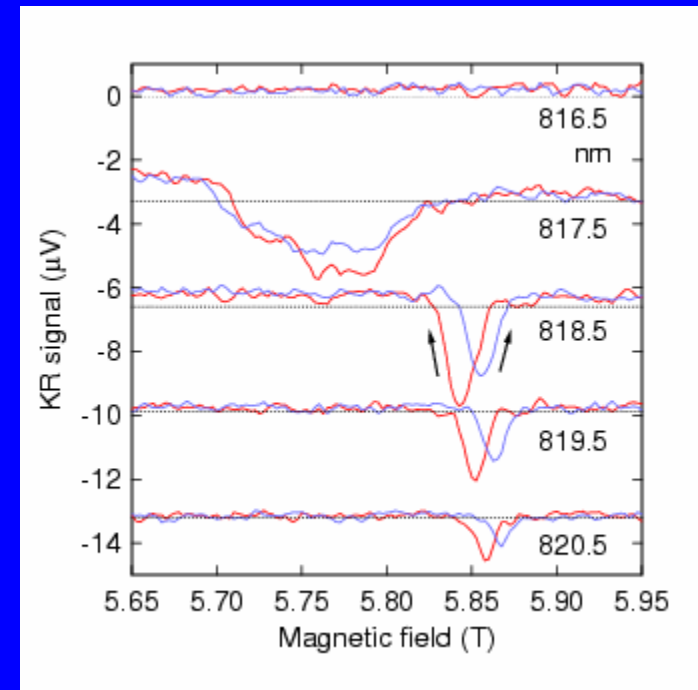
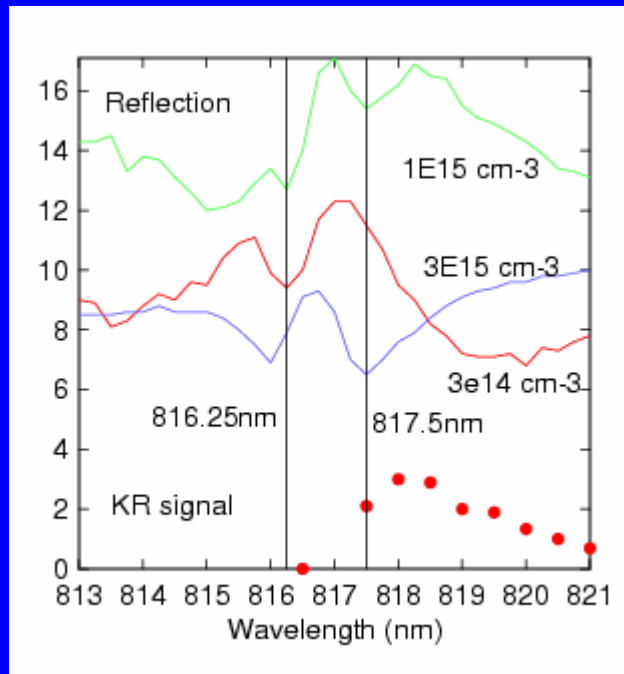


## The experiment is a double resonance

- Microwave resonance in the ground state—tuned by changing the magnetic field
- Optical resonance to the excited state—tuned by changing the laser wavelength

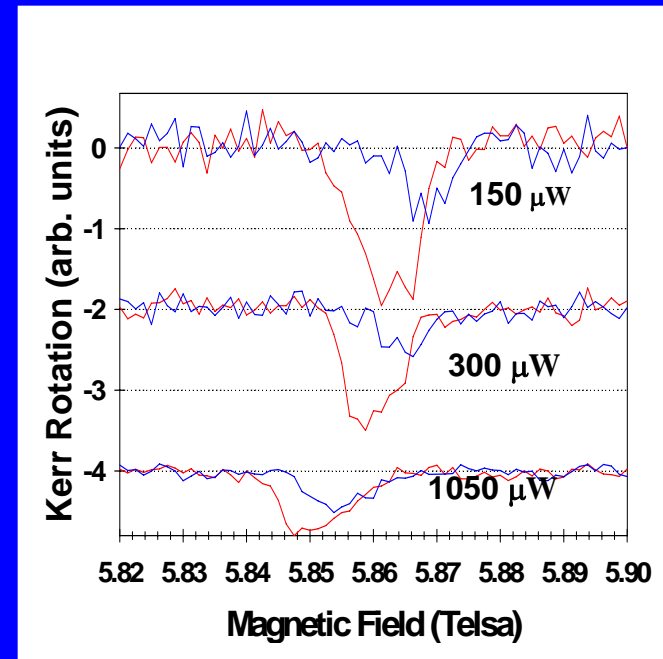
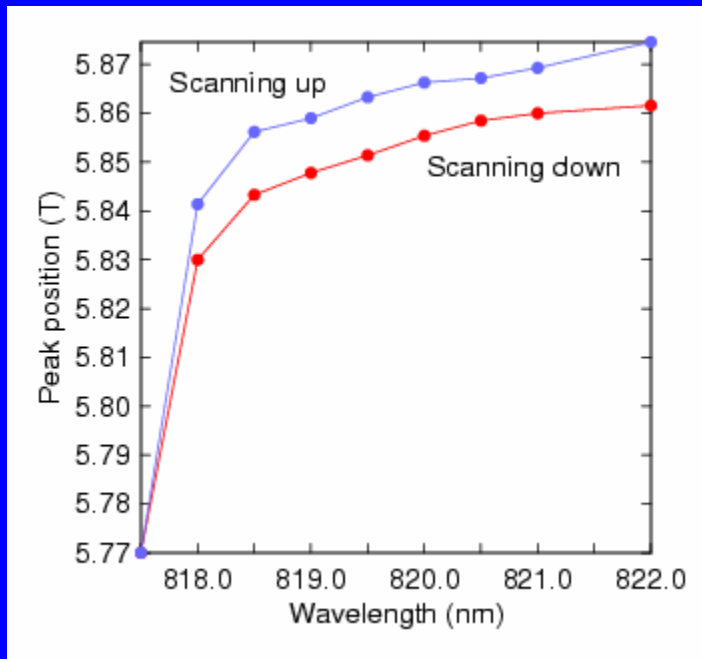


## Dependence on Wavelength of the Probe Light



- Optical Resonances for  $B = 6T$ 
  - $D^0X$  at 817.5 nm
  - $X$  at 816.25 nm
- 6dbm, 34.694, 300 uW probe, 9/20/05 data
- On resonance
  - Large shift
  - Extra line
- Detuned to lower  $E$  (longer  $\lambda$ )
  - Sharp single line
  - Decreasing amplitude
  - Slow shift

# Peak Position versus probe wavelength and power



## Position versus wavelength

- Approaches a limit off-resonance

9/20/05 data—31F, light unfocussed

## Position versus power

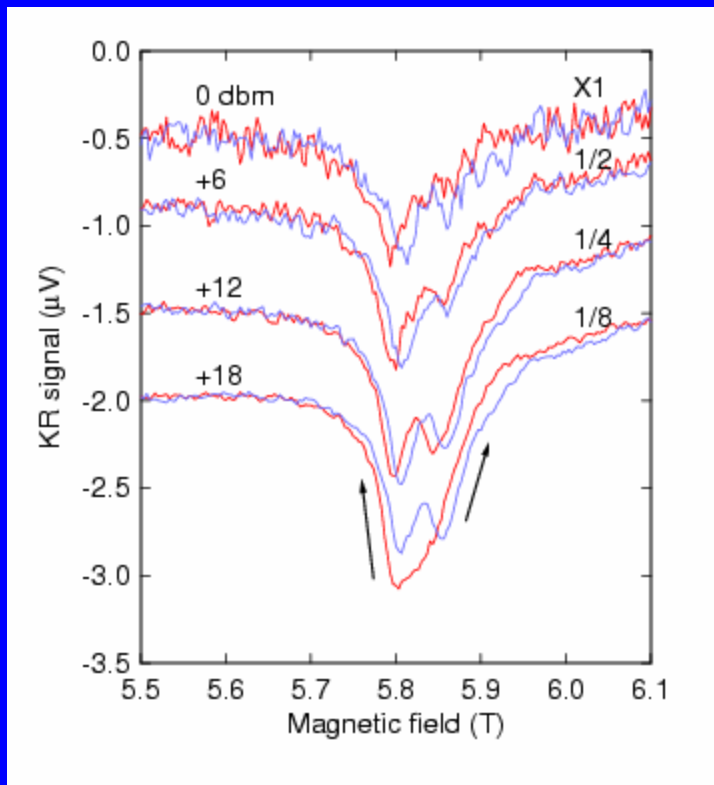
- Shifts to lower fields with increasing power
- DNP by slight depolarization of  $\langle S \rangle$  by the linearly polarized light

9/21/05 data, 820.5nm, -6dbm

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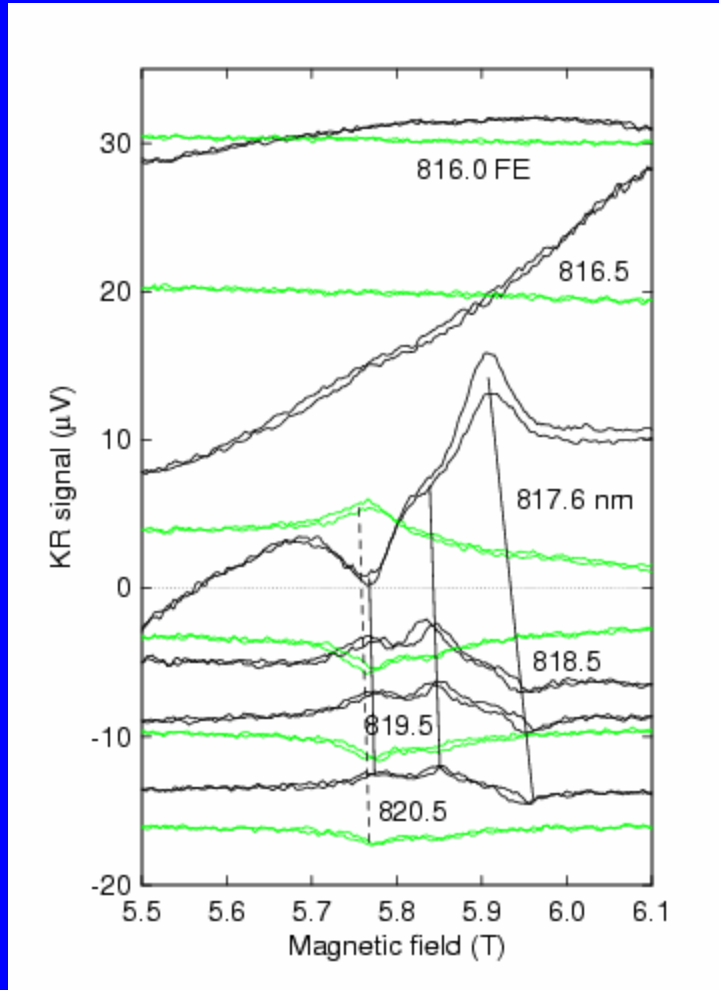
## Dependence on Microwave Power for 3E15



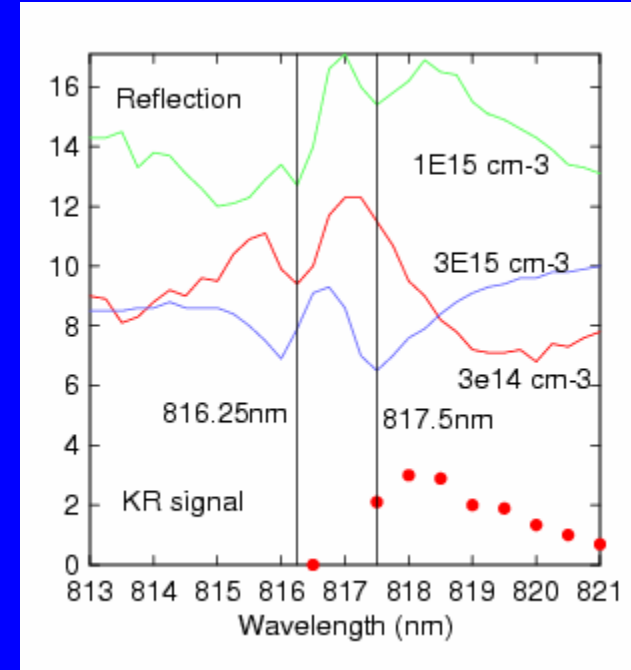
- Different scaling
- Two lines
- Strong signals in quadrature
- DNP enhanced for high-field line at highest power
- Localized and de-localized electrons

28G, 818.6nm, 8/9/05, no offset added

## Dependence on wavelength of probe



24 G,  $3\text{E}15$ , 8/11/05,  
+6dbm, 0.15 T/min

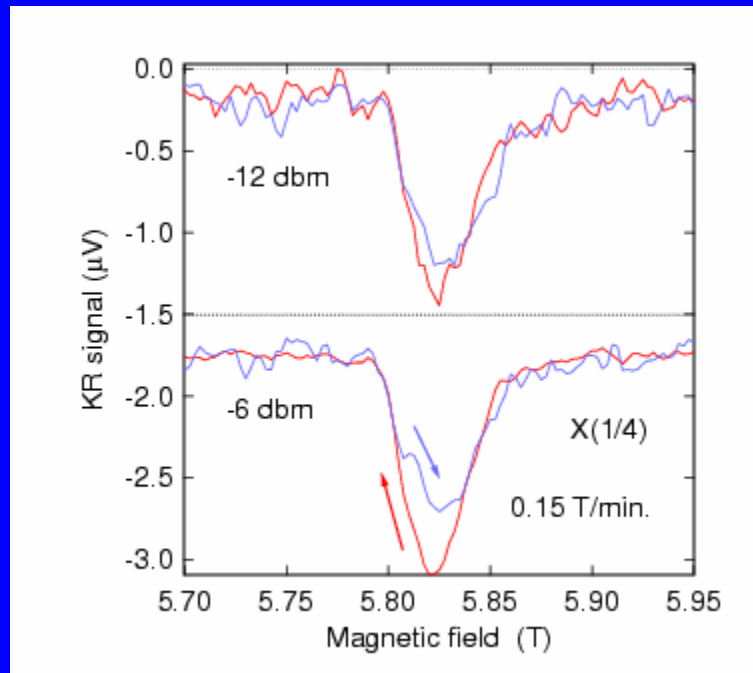


- Black: in phase; Green: in quadrature
- In Phase
  - 2 or 3 resonances
  - phase changes
- In Quadrature
  - one resonance
  - reveals dynamics  $\sim 3 \text{ kHz}$

# Outline

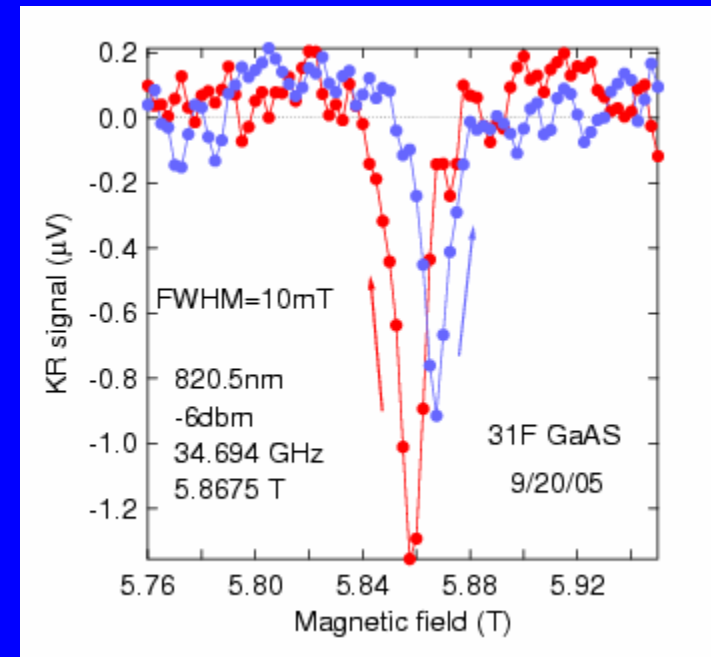
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## g-factor and linewidth



### Near resonance—817.8 nm

- $g = 0.431$ ; FWHM = 49 mT
- Real transitions
  - $D^0X$ : holes
  - X: electrons and holes

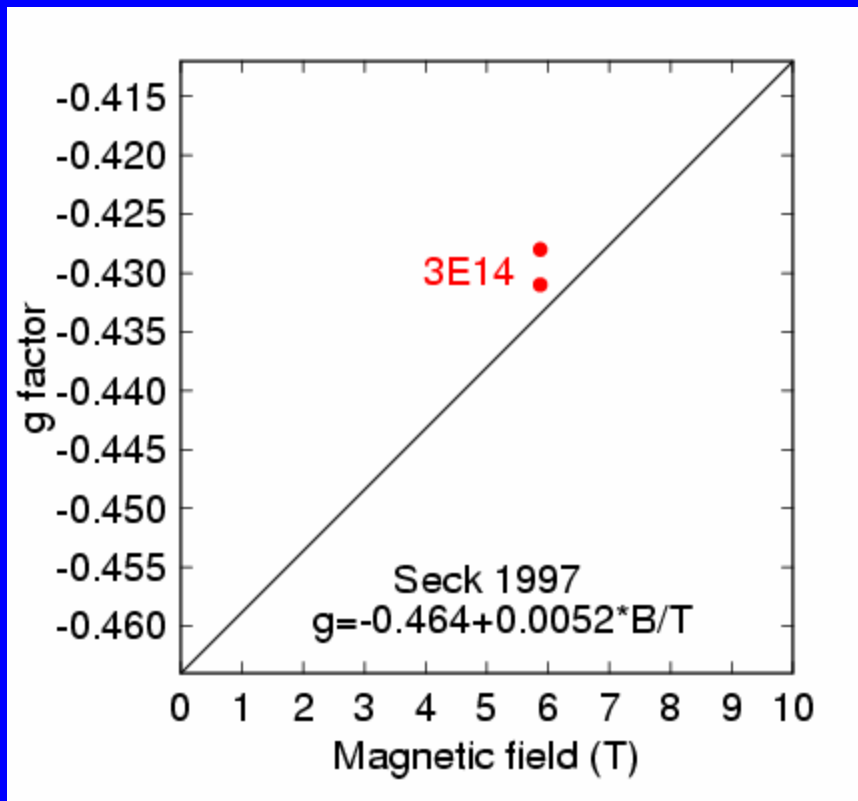


### Off resonance—820.5 nm

- $g = 0.428$ ; FWHM = 10 mT
- Dispersive part of the index

Sweep rate 0.15 T/min

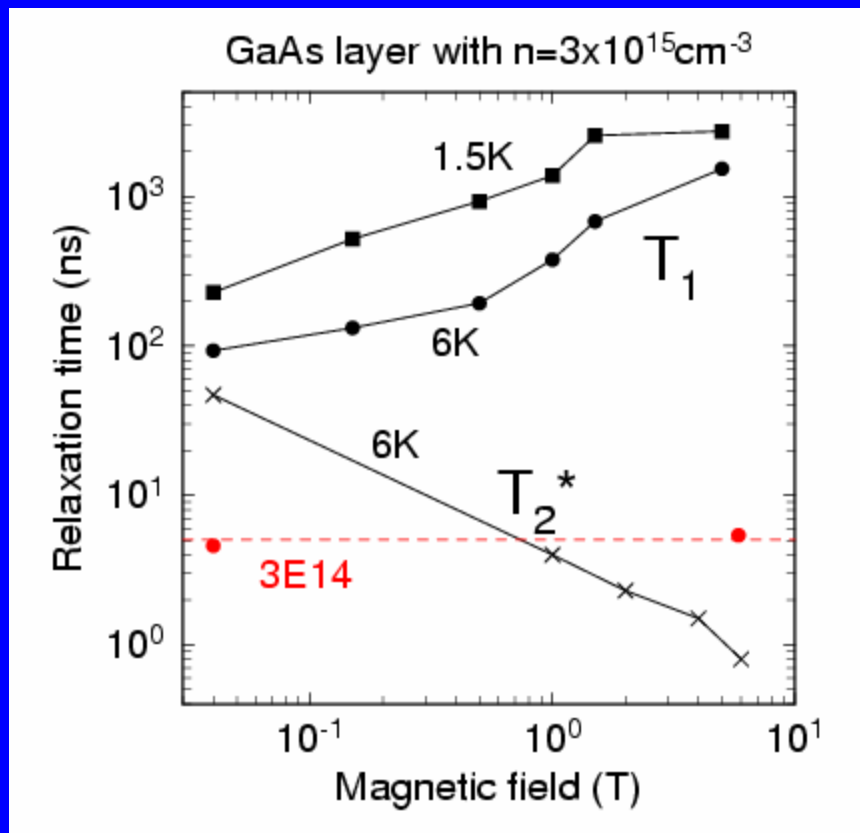
## Comparison of g-factors with other work



- Negative sign added
- $^{31}\text{F}$  shown for near resonance & off-resonance



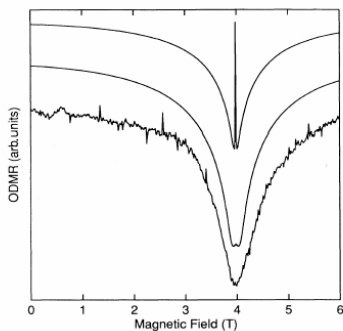
## $T_2^*$ for donors in GaAs



- ~5 ns for low and high magnetic fields
- Limited by fluctuation in nuclear spin
- Good starting point for measuring  $T_2$

## Lifetimes for donors at high magnetic fields

	Sample	$\Delta B$ FWHM (mT)	$T_2^*$
This work	$N_D - N_A = 3 \times 10^{14} \text{ cm}^{-3}$	10	5.4 ns
Kai-Mei Fu et al. (06)	$N_D - N_A = 5 \times 10^{13} \text{ cm}^{-3}$	--	2.3 ns
Seck et al. (1997)	$N_D - N_A = 6.6 \times 10^{14} \text{ cm}^{-3}$	50	1.1 ns
Kikkawa & Awschalom (1998)	Semi-insulating	--	100 ps
Trombetta & Kennedy (1993)	Semi-insulating	900	60 ps



g-broadening is small, possibly negligible

March 2006

## Summary and Future Work

- **Kerr rotation provides adequate sensitivity to work with isolated donors**
- **Nuclear effects**
  - Resonance without DNP enhancement—unpolarized nuclei
  - DNP enhancement—good for finding ESR
  - ESR-pinning—possibly good for polarizing nuclei
- **Linewidth near or at the hyperfine limit—little or no g-broadening**
  
- **Measuring  $T_2$  with electron-spin-echo in ensembles**