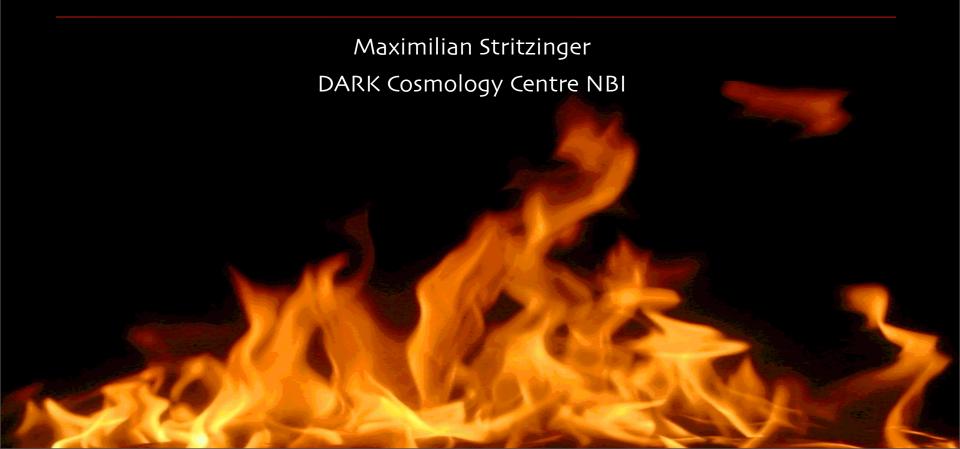
Global Parameters of Type Ia Supernovae



Global parameters

What do we know from observations about

- energy source, 56Ni mass?
- total mass?
- explosion energy ?
- density ?
- element distribution ?

Stritzinger & Leibundgut, 2005, A&A, 431, 423
Stritzinger, Leibundgut, Walch & Contardo, 2006, A&A, 450, 241
Blinnikov, et al., 2006, A&A, 453, 229
Stritzinger, Mazzali, Sollerman & Benetti, 2006, A&A, 460, 793



Usual procedure

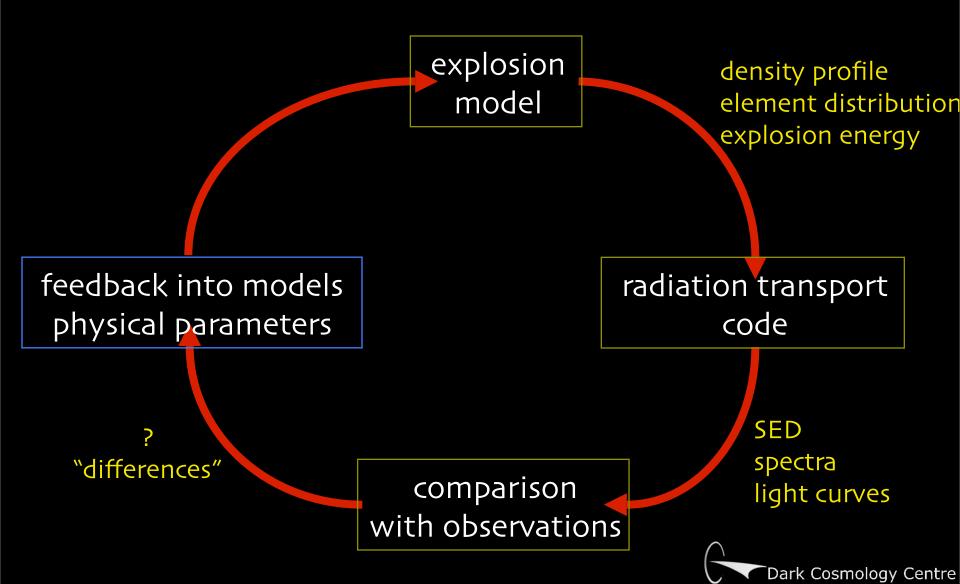
Take explosion model

density and element distribution and
explosion energy
and calculate the emerging spectrum and SED

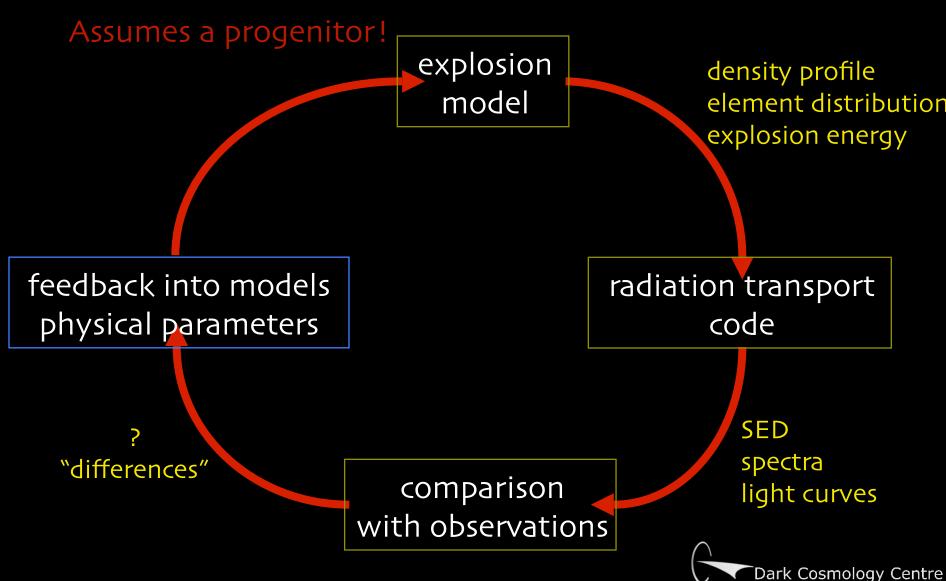
always assume that you know the progenitor usually Chandrasekhar mass white dwarf



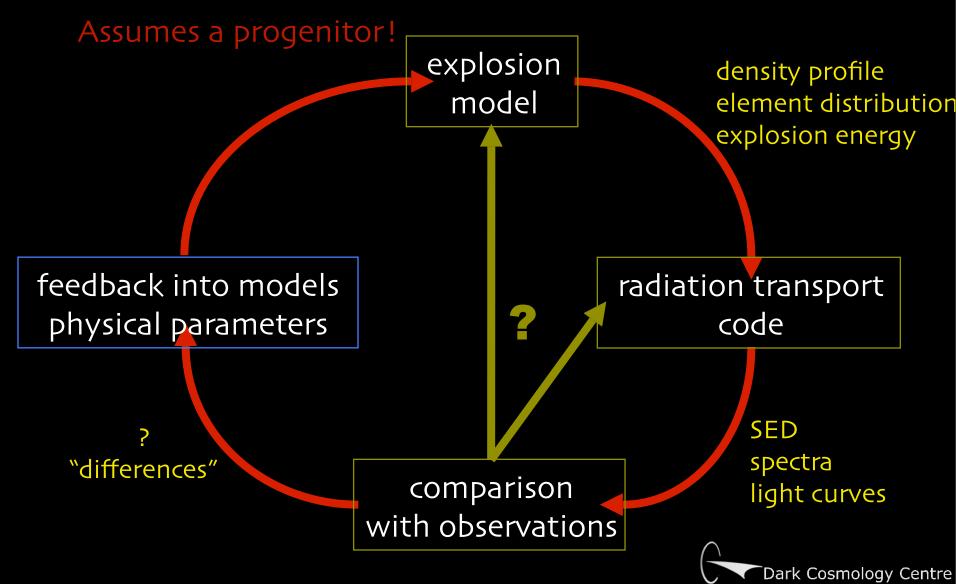
Understanding SNe la



Understanding SNe la



Understanding SNe la



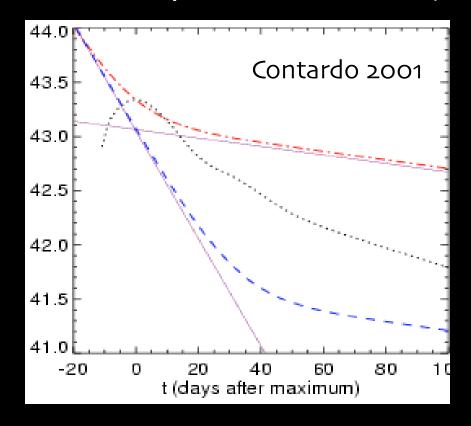
Can we do better?

Determine 56Ni from the peak luminosity

Arnett's law

Requires a bolometric luminosity

- multi filter observations
- distance





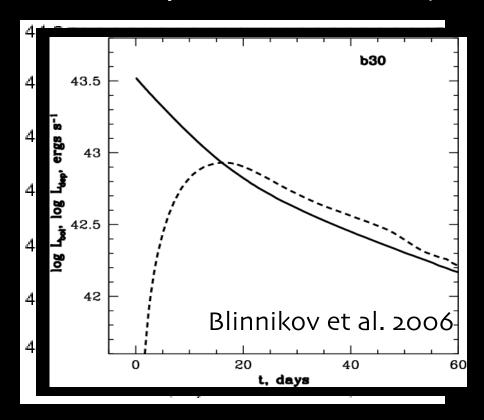
Can we do better?

Determine 56Ni from the peak luminosity

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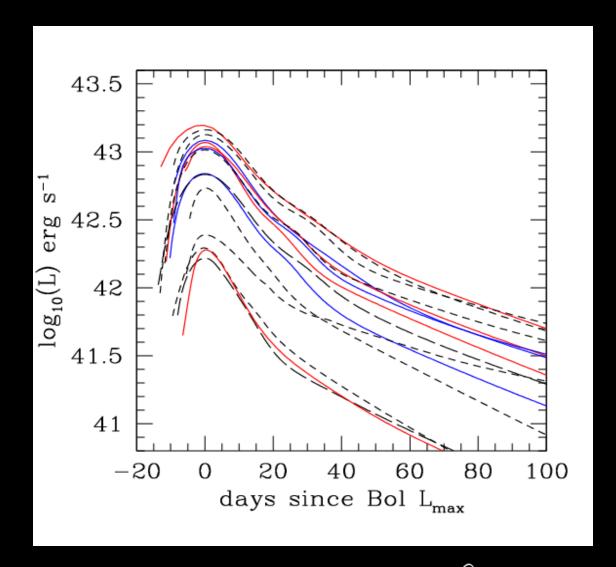
Requires a bolometric luminosity

- multi filter observations
- distance

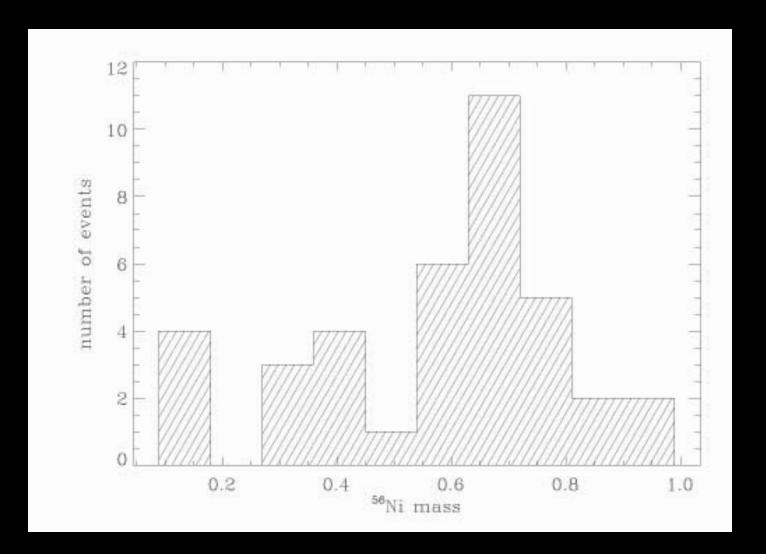




Bolometric light curves

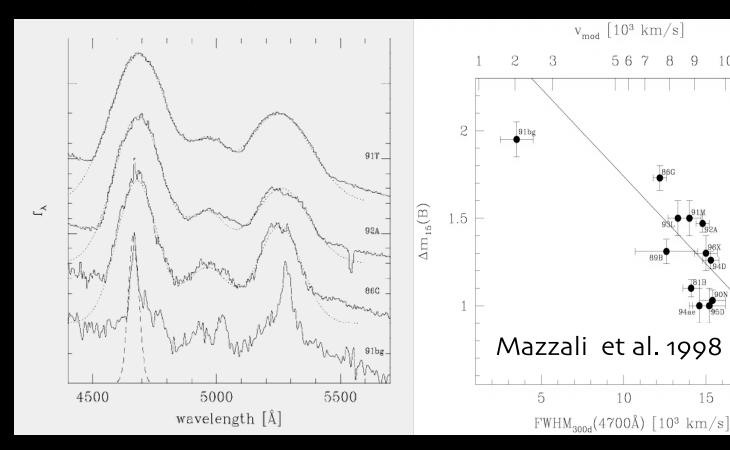


56Ni masses from light curves



Check with a different method

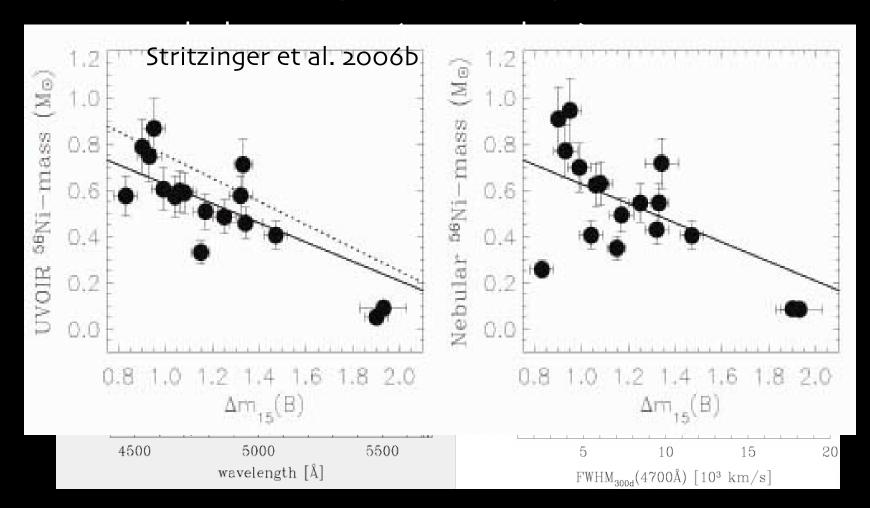
Ni masses from the emission line in nebular spectra (t~300 days)





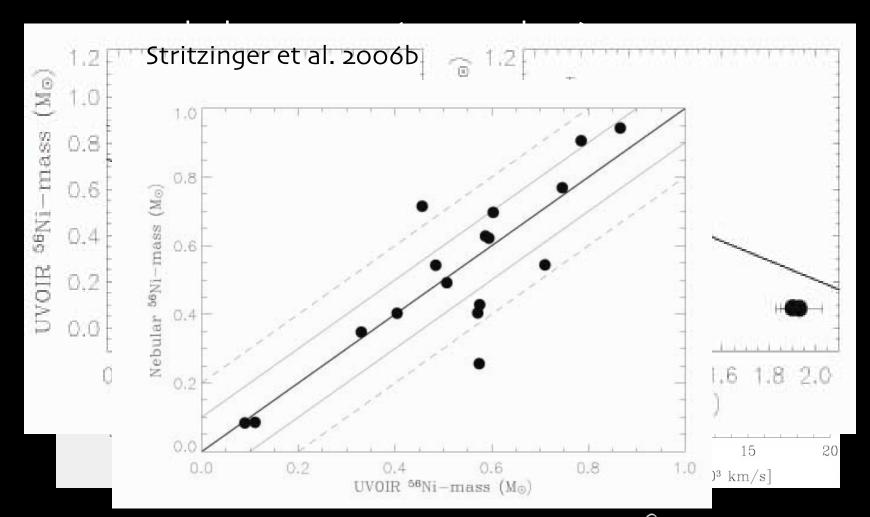
Check with a different method

Ni masses from the emission line in



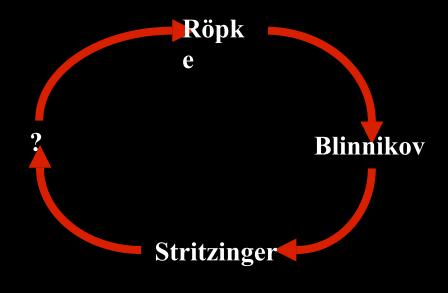
Check with a different method

Ni masses from the emission line in

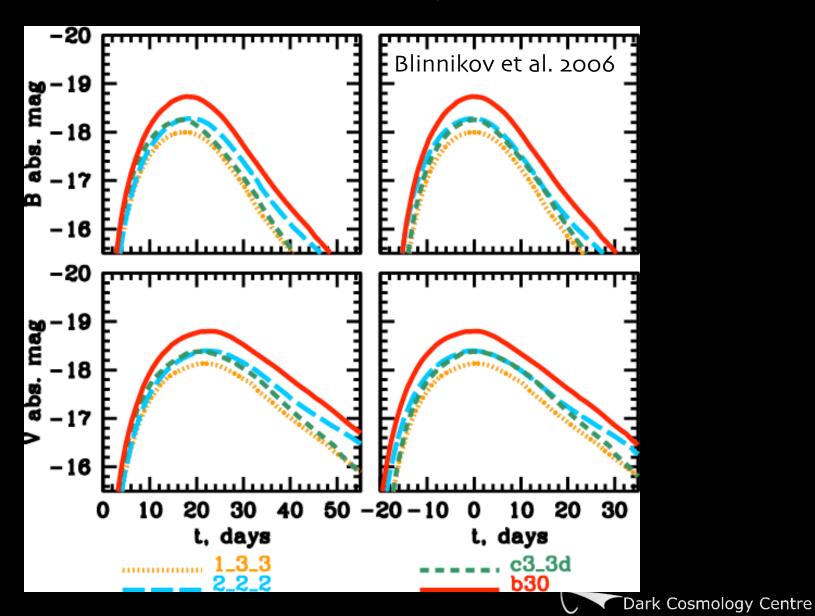


Check with models

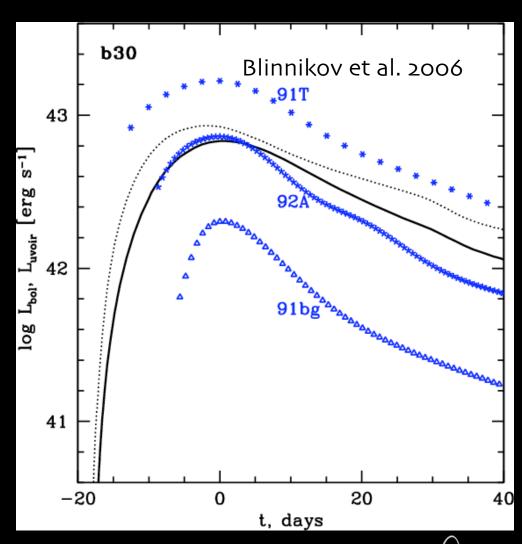
Calculate the emission from explosion models (Röpke et al. 2004-2006) with a radiation transport code (STELLA – Blinnikov et al. 1998) and then derive the parameters from these "observations."



Check with models



Comparison with real data

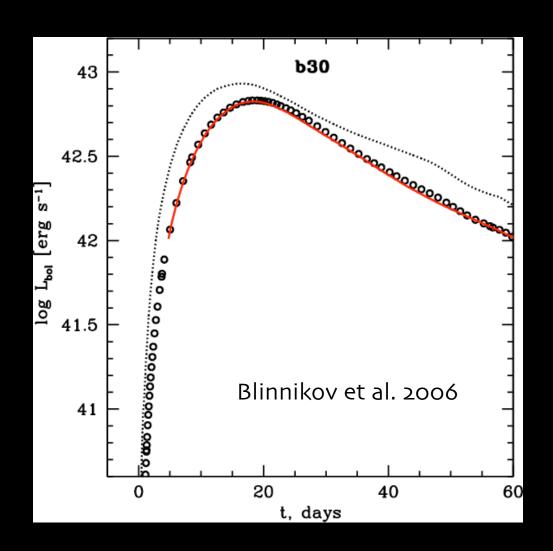


Check ...

UVOIR light curve reproduced very well – validation of procedure in SLo5

True bolometric light curve offset by about 15%

10% correction used is SLo5 not quite right ...





Lower limits on H_o from models of SNe Ia



Determining H_o from models

Hubble's law

$$D = \frac{v}{H_0} = \frac{cz}{H_0}$$

Luminosity distance

$$D_L = \sqrt{\frac{L}{4\pi F}}$$

Ni-Co decay

$$E_{Ni} = \frac{\lambda_{Ni}\lambda_{Co}}{\lambda_{Ni} - \lambda_{Co}} \left\{ \left[Q_{Ni} \left(\frac{\lambda_{Ni}}{\lambda_{Co}} - 1 \frac{1}{J} - Q_{Co} \right) e^{-\lambda_{Ni}t} + Q_{Co} e^{-\lambda_{Co}t} \right\} N_{Ni,0} \right\}$$

$$H_0 = \frac{cz}{D} = cz\sqrt{\frac{4\pi F}{L}} = cz\sqrt{\frac{4\pi F}{\alpha E_{Ni}}} = cz\sqrt{\frac{4\pi F}{\alpha \epsilon(t) M_{Ni}}}$$

 α : conversion of nickel energy into radiation (L= α E_{Ni})

 ϵ (t): energy deposited in the supernova ejecta



$$H_0 = \frac{cz}{D} = cz\sqrt{\frac{4\pi F}{L}} = cz\sqrt{\frac{4\pi F}{\alpha E_{Ni}}} = cz\sqrt{\frac{4\pi F}{\alpha \epsilon (t) M_{Ni}}}$$
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Arnett's rule

 α : conversion of nickel energy into radiation (L= α E_{Ni})

 ϵ (t): energy deposited in the supernova ejecta



$$H_0 = \frac{cz}{D} = cz\sqrt{\frac{4\pi F}{L}} = cz\sqrt{\frac{4\pi F}{\alpha E_{Ni}}} = cz\sqrt{\frac{4\pi F}{\alpha \epsilon(t) M_{Ni}}}$$

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$$H_0 = \frac{cz}{D} = cz\sqrt{\frac{4\pi F}{L}} = cz\sqrt{\frac{4\pi F}{\alpha E_{Ni}}} = cz\sqrt{\frac{4\pi F}{\alpha \epsilon (t) M_{Ni}}}$$
Ni-Co decay and rise time

 α : conversion of nickel energy into radiation (L= α E_{Ni})

 ϵ (t): energy deposited in the supernova ejecta



$$H_0 = \frac{cz}{D} = cz\sqrt{\frac{4\pi F}{L}} = cz\sqrt{\frac{4\pi F}{\alpha E_{Ni}}} = cz\sqrt{\frac{4\pi F}{\alpha \epsilon(t) M_{Ni}}}$$

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 ϵ (t): energy deposited in the supernova ejecta



$$H_0 = \frac{cz}{D} = cz\sqrt{\frac{4\pi F}{L}} = cz\sqrt{\frac{4\pi F}{\alpha E_{Ni}}} = cz\sqrt{\frac{4\pi F}{\alpha \epsilon(t) M_{Ni}}}$$

 α : conversion of nickel energy into radiation (L= α E_{Ni})

 ϵ (t): energy deposited in the supernova ejecta

Need bolometric flux at maximum F and the redshift z as observables



Assumptions

Rise time (15-25 days)

→ about 10% uncertainty

Arnett's rule

```
energy input at maximum equals radiated energy (i.e. \alpha \approx 1, \epsilon(t_{max}) \approx 1)
```

Nickel mass from models

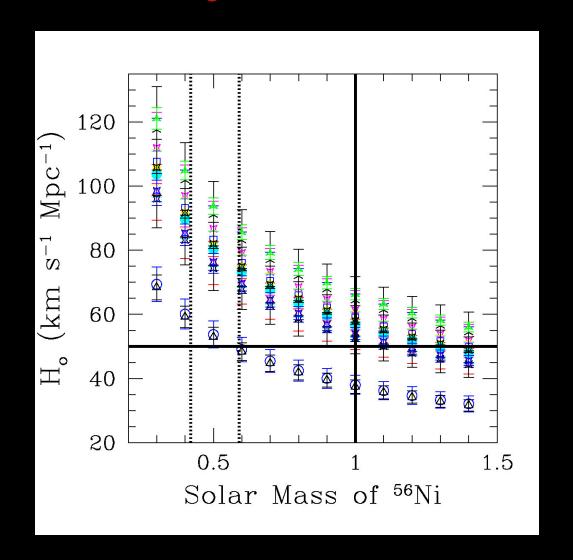
→ uniquely defines the bolometric peak luminosity



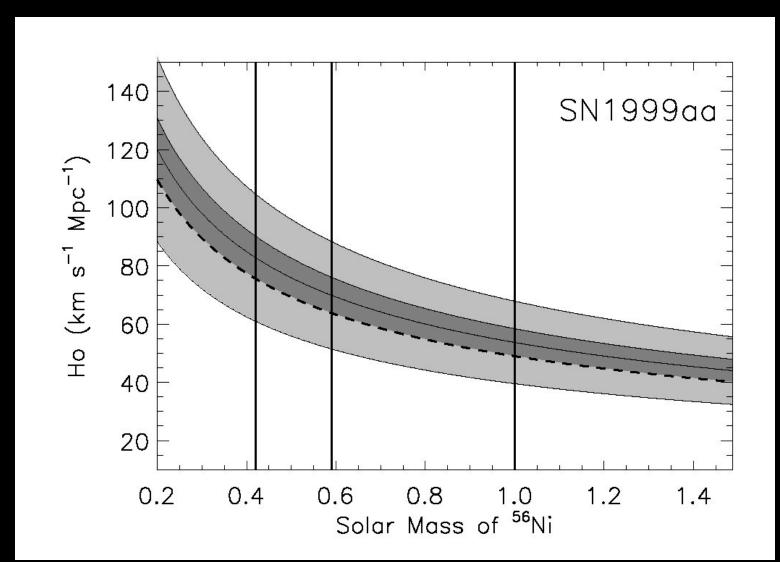
Ho and the Ni mass

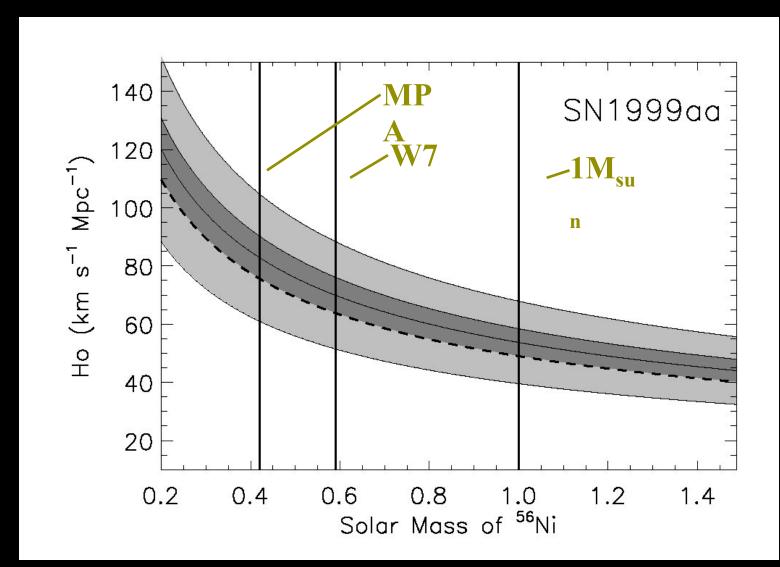
Problem:

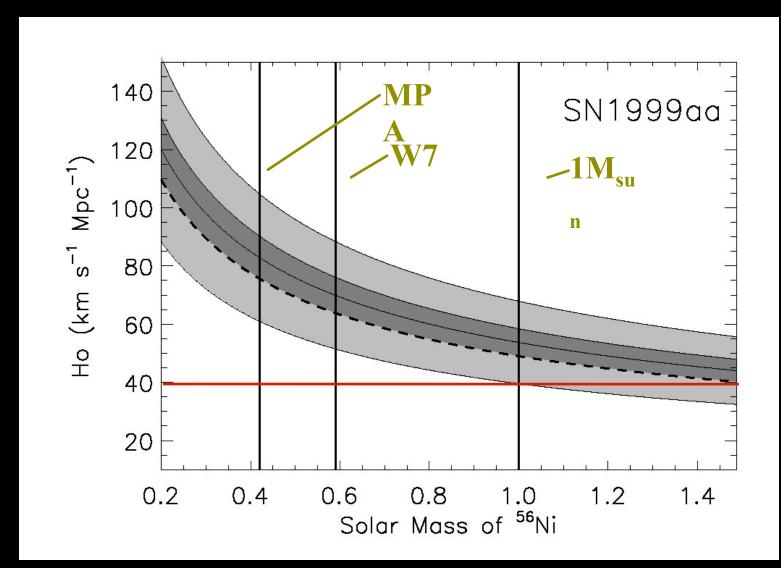
Since SNe Ia have individual Ni masses it is not clear which one to apply!

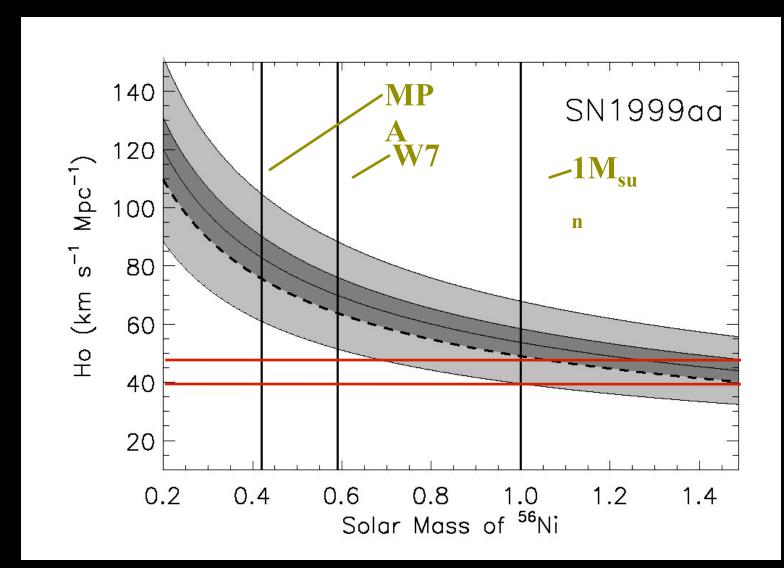


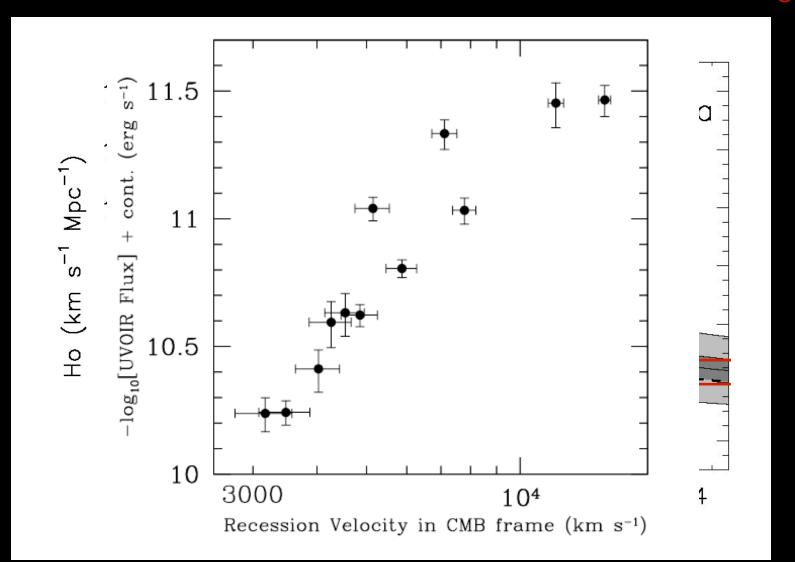












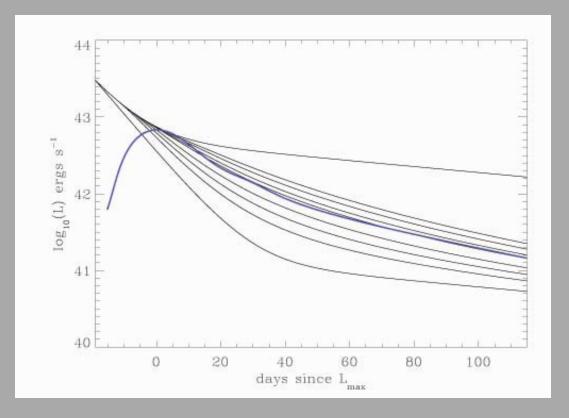
Ejecta masses from light curves

γ-ray escape depends on the total mass

of the ejecta

$$M_{ej} = \frac{8\pi}{\kappa q} t_0^2 v^2 \propto \frac{v^2}{\kappa q}$$

- v: expansion velocity
- κ: γ-ray opacity
- q: distribution of nickel





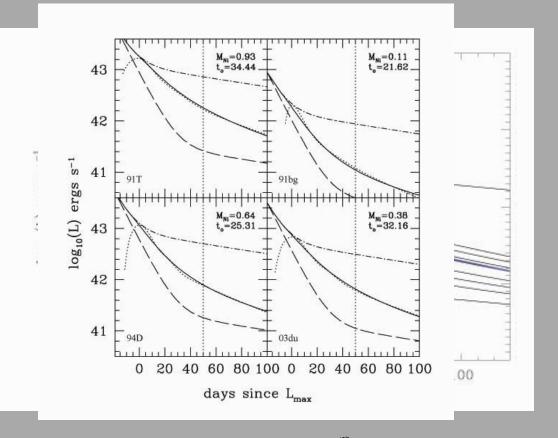
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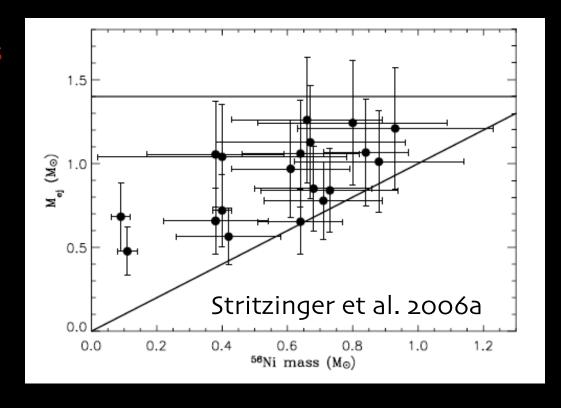




Ejecta masses

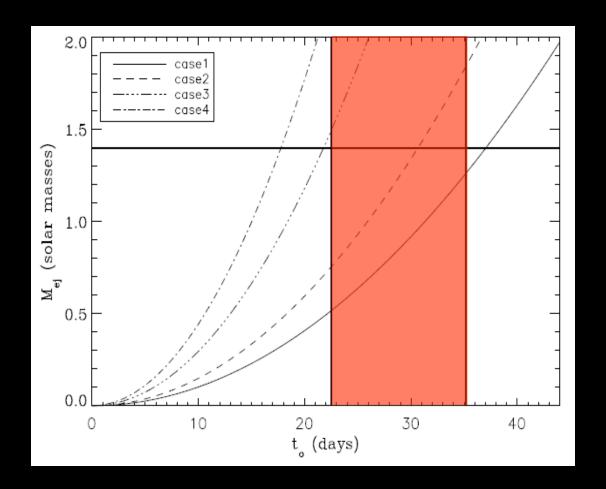
Large range in nickel and ejecta masses

- no ejecta mass at 1.4M_{sun}
- factor of 2 in ejecta masses
- some rather small differences between nickel and ejecta mass



Dependence on explosion parameters

```
case 1: (fiducial)
   v=3000 km/s),
   \kappa=0.025 cm<sup>2</sup>/g and q=1/3
case 2:
   v=3625 km/s
case 3:
   v=3625 km/s,
   \kappa=0.0084 cm/g and
   q=1/2
case 4:
   v=3625 km/s,
   \kappa=0.0084 cm/g and
   q = 1/3
```





Summary

Arnett's rule is astonishingly good

- → determine the nickel mass in the explosion from the peak luminosity
- → large variations (up to a factor of 10)
- \rightarrow lower limit on H_o based on models: 50

Ejecta masses appear to scatter considerably

- → implications for progenitors?
- → model too simple ?
- → how can we improve?



Summary (cont.)

SNe Ia may be more varied than we like... Differences in the explosions?

- density
- ignition
- asymmetries
- progenitor mass
- metallicity

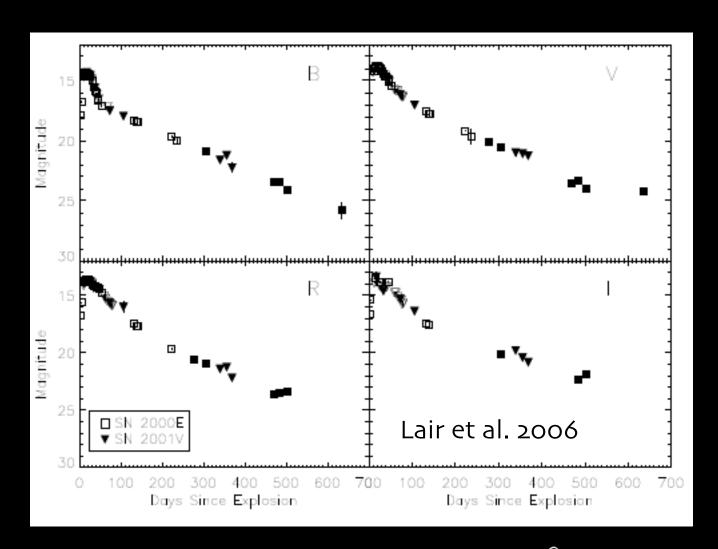


Very Late-time Emission

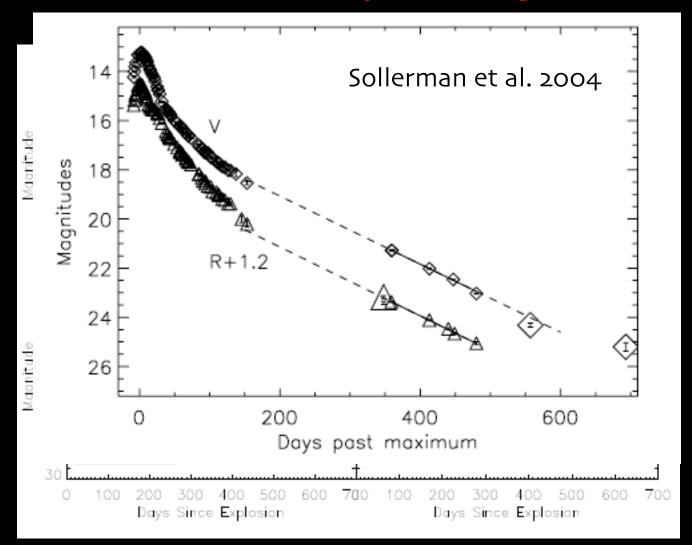
- Nucleosynthesis of radioactive and stable elements and their distribution
- Unique opportunity to probe the magnetic field structure/ configuration
- A weak/radially combed magnetic field
 no deposition of positrons
- A strong/tangled magnetic field
 - --> depostion of positrons
- ♦ Constrain positron's contribution to the Galactic 511 KeV line
- ♦ Study physics of freeze-out and the IR-catastrophe



Late-time optical light curves

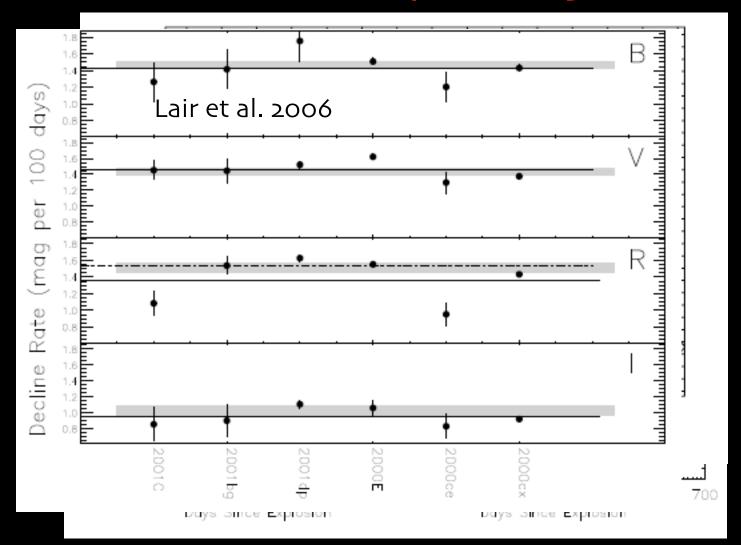


Late-time optical light curves

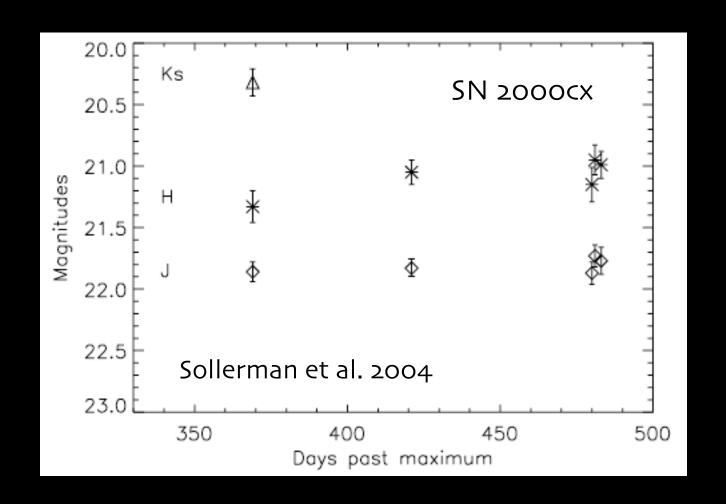




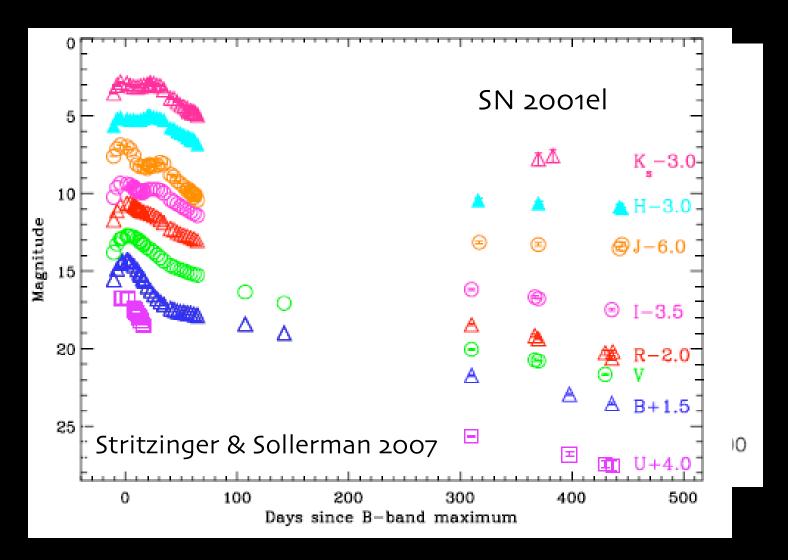
Late-time optical light curves



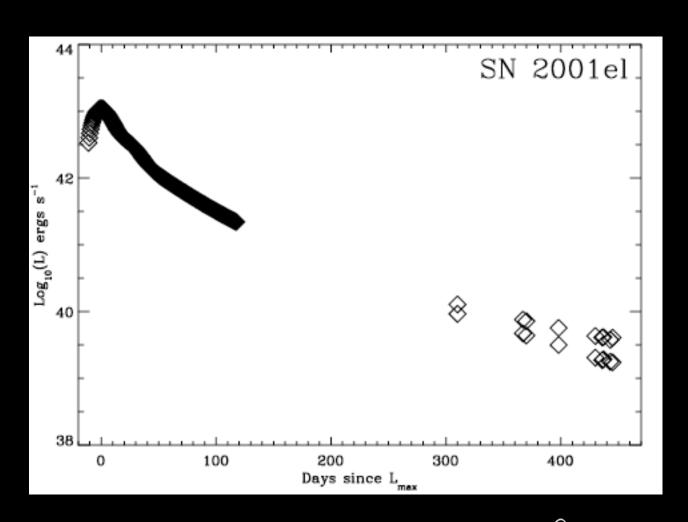
Late-time emission in IR



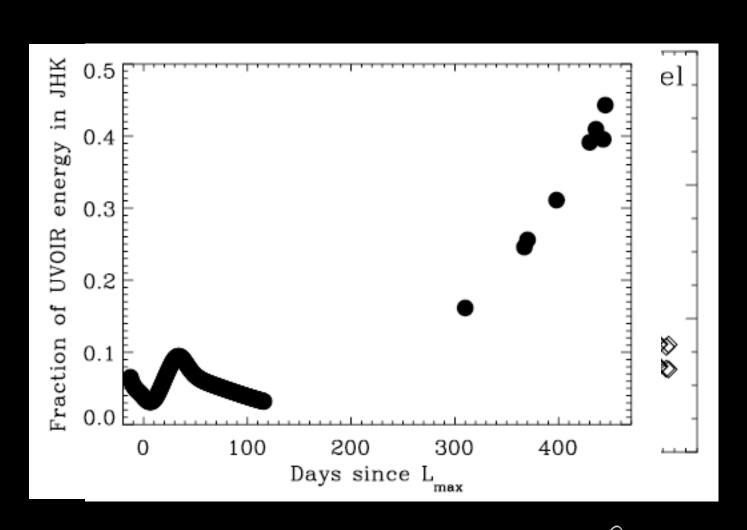
Late-time emission in IR



Late-time energetics



Late-time energetics



Summary

- increase importance of the IR during late-times is a generic feature of normal la's
- pervious conclusions of positron escape based on optical photometry alone are incorrect