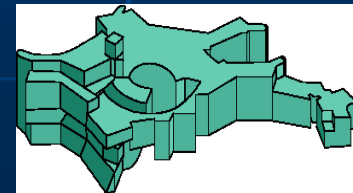


Modeling Turbulent (Thermonuclear) Combustion

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Garching



KITP, UCSB,
February 22, 2006





MPI für Astrophysik
Simulation: W. Hillebrandt, F. Röpke
Visualisierung: R. Bruckschen



MAX-PLANCK-GESELLSCHAFT

Time(sec): 0.00 Size(km): 2029.9

The history of SuCCESs

(*Supernova Combustion Code for Explosion Simulations*)

- *Jens Niemeyer (1994 -)*
- *Martin Reinecke (1996 - 2002)*
- *Wolfram Schmidt (2001 -)*
- *Fritz Röpke (2001 -)*
- *Michael Fink (2006 -)*

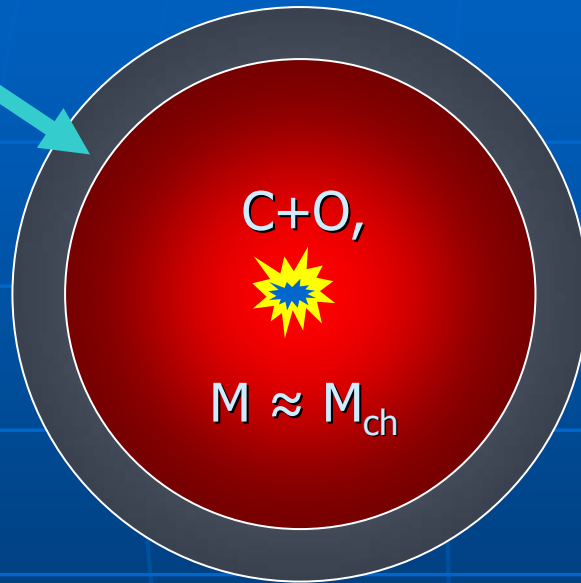
The “standard model”



- White dwarf in a binary system
- Growing to the Chandrasekhar mass by mass transfer

How does the model work?

He (+H)
from binary
companion



Density $\sim 10^9 - 10^{10}$ g/cm

Temperature: a few 10^9 K

Radii: a few 1000 km

Explosion energy:

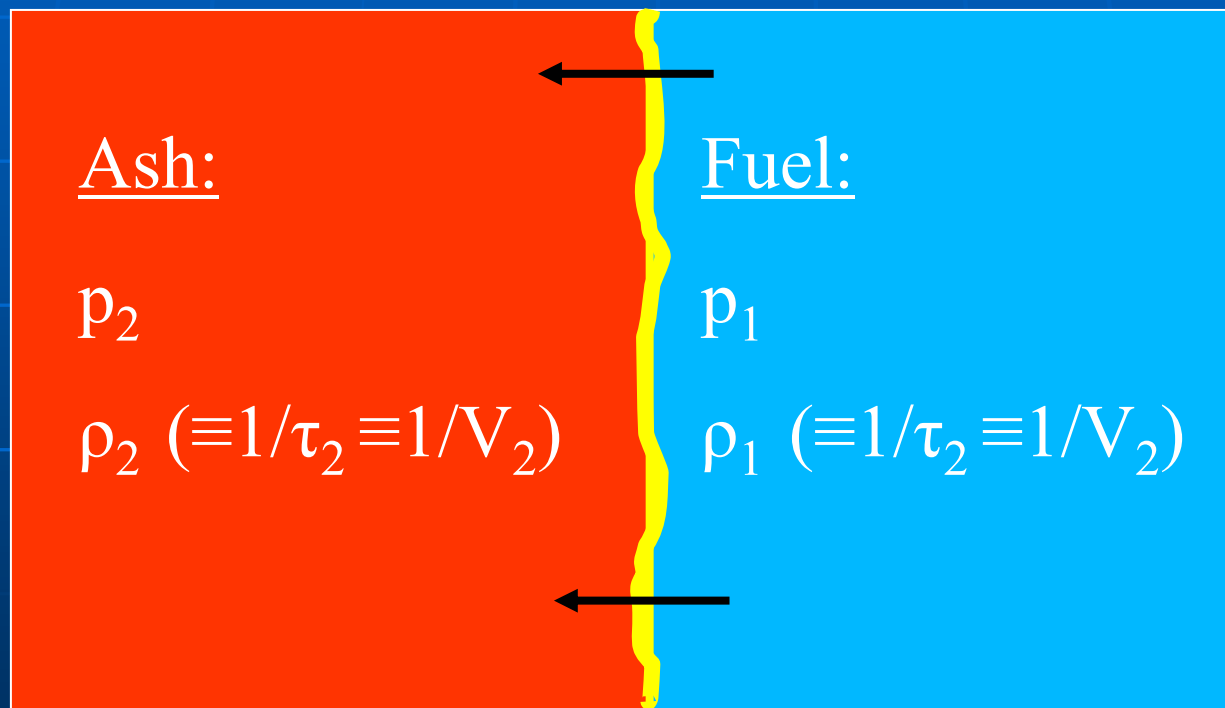
*Fusion C+C, C+O,
O+O \rightarrow "Fe"*

Laminar burning
velocity:

$U_L \sim 100$ km/s $\ll U_S$

Too little is burned!

Some fundamentals of combustion theory



The “Hugoniot-function” for the burned gas,

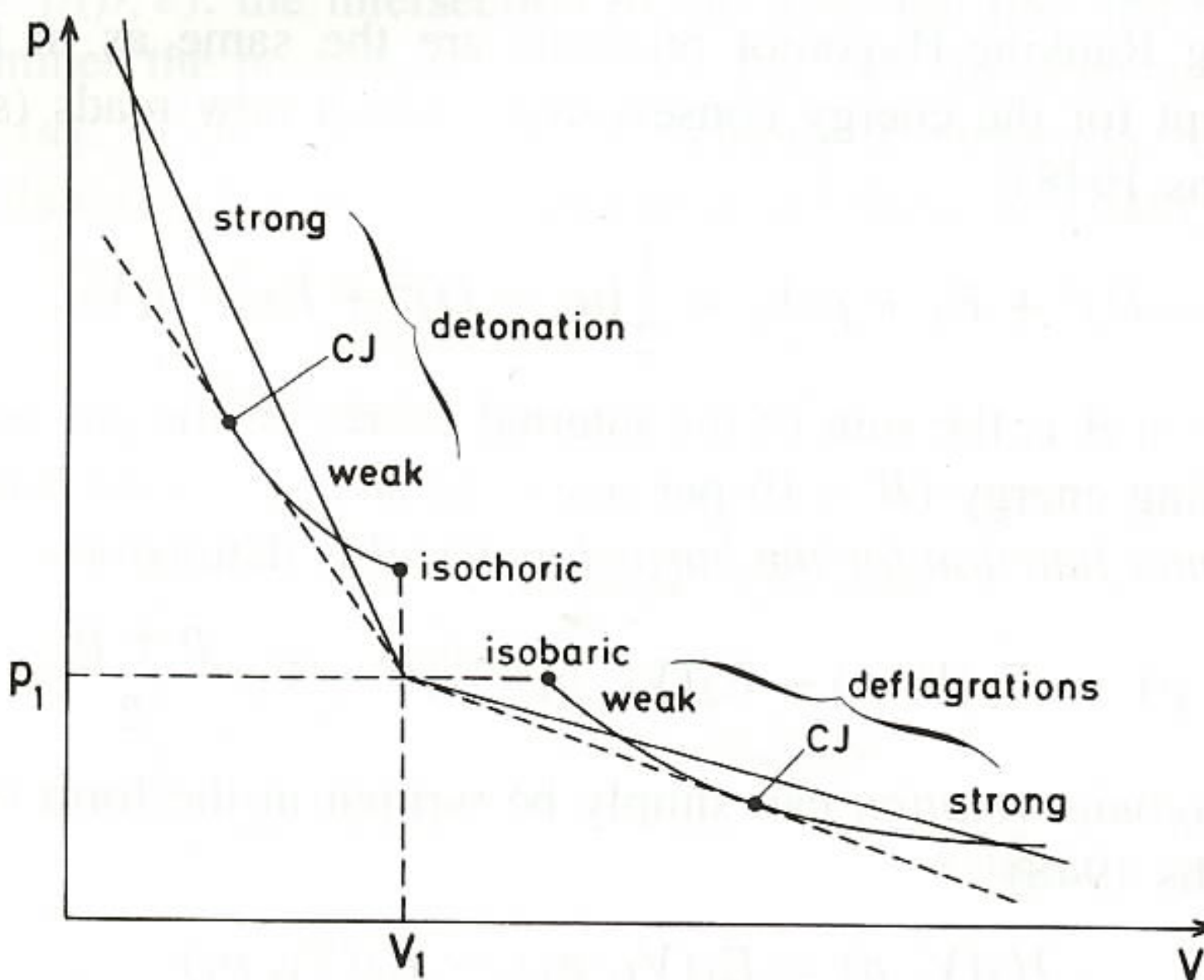
$$H_2(\tau, p) \equiv E_2(\tau, p) - E_2(\tau_1, p_1) + (\tau - \tau_1)(p + p_1)/2$$

and the “Rayleigh-condition”

$$v_B^2 = - (p_2 - p_1)/(\tau_2 - \tau_1); \quad p_2 - p_1 < \tau_2 - \tau_1$$

($\tau = 1/\rho$, “1” = unburned state, “2” = burned state)

*(“Jump conditions” from conservation laws;
analogous to shock waves)*



Observed in “real” combustion experiments:

Only weak deflagrations and Chapman-Jouguet detonations!

What is the mode of nuclear burning in SNe Ia?

➤ “Detonation”:

(Super-) Sonic front;

heating to ignition by a shock wave.

➤ “Deflagration”:

Subsonic front;

heating to ignition by heat diffusion.

Strong Si-lines at maximum light:

Pure detonations are excluded (Arnett, 1969)!

(But possibly at lower densities: DDT ???)

The physics of turbulent combustion

- Everydays experience:
Turbulence increases the burning velocity.
- In a star:
Reynoldsnumber $\sim 10^{14}$!
- In the limit of strong turbulence: $U_B \sim V_T$!
- Physics of thermonuclear burning is very similar to premixed chemical flames.



A couple of definitions:

Kolmogorov (length) scale

$$\eta := (v^3/\varepsilon)^{1/4}$$

(Turbulent) Reynolds number

$$Re := v'/s_L \cdot l/l_F$$

(Turbulent) Damköhler number

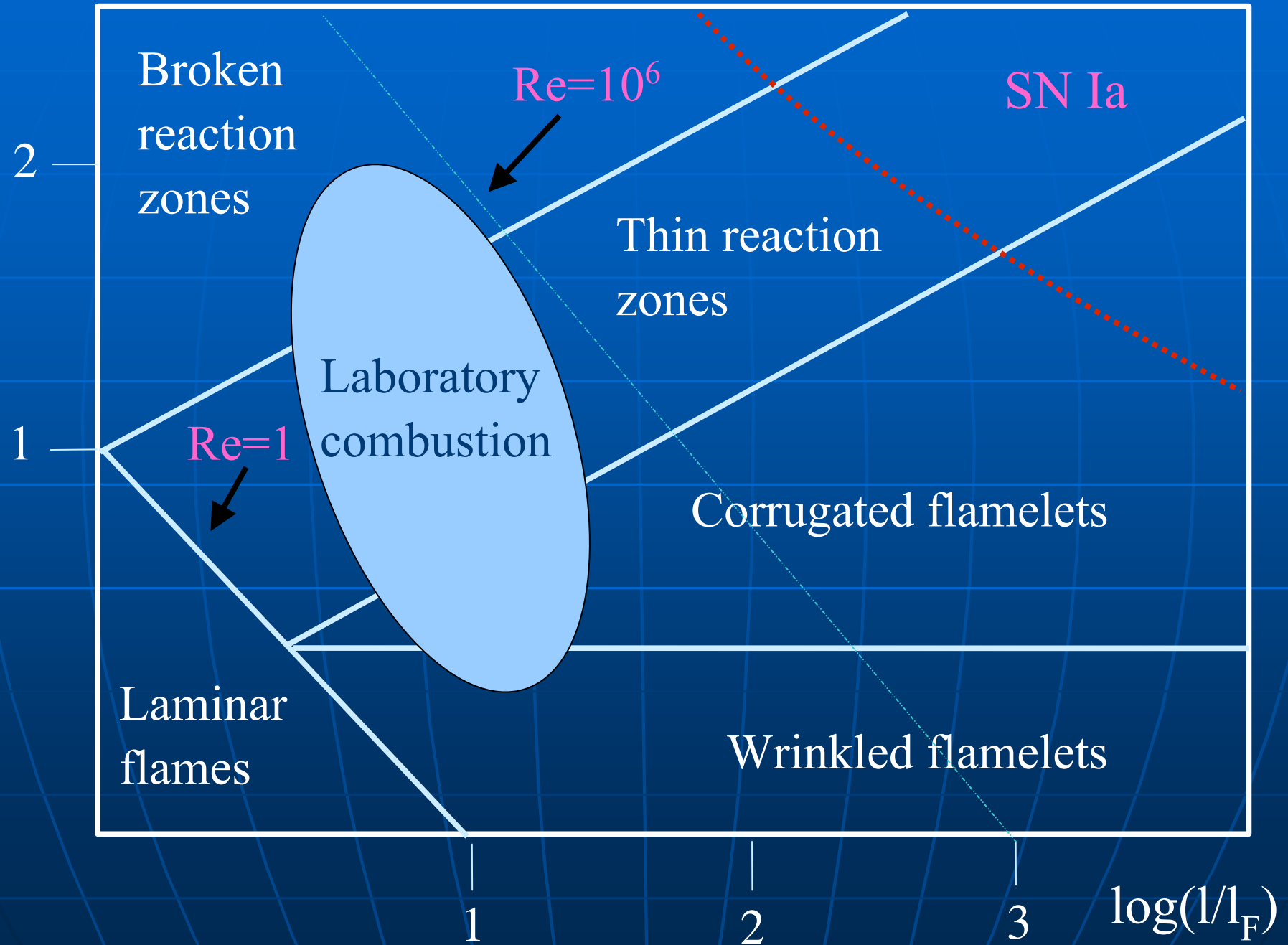
$$Da := s_L/v' \cdot l/l_F$$

(Turbulent) Karlovitz number

$$Ka := l_F^2/\eta^2$$

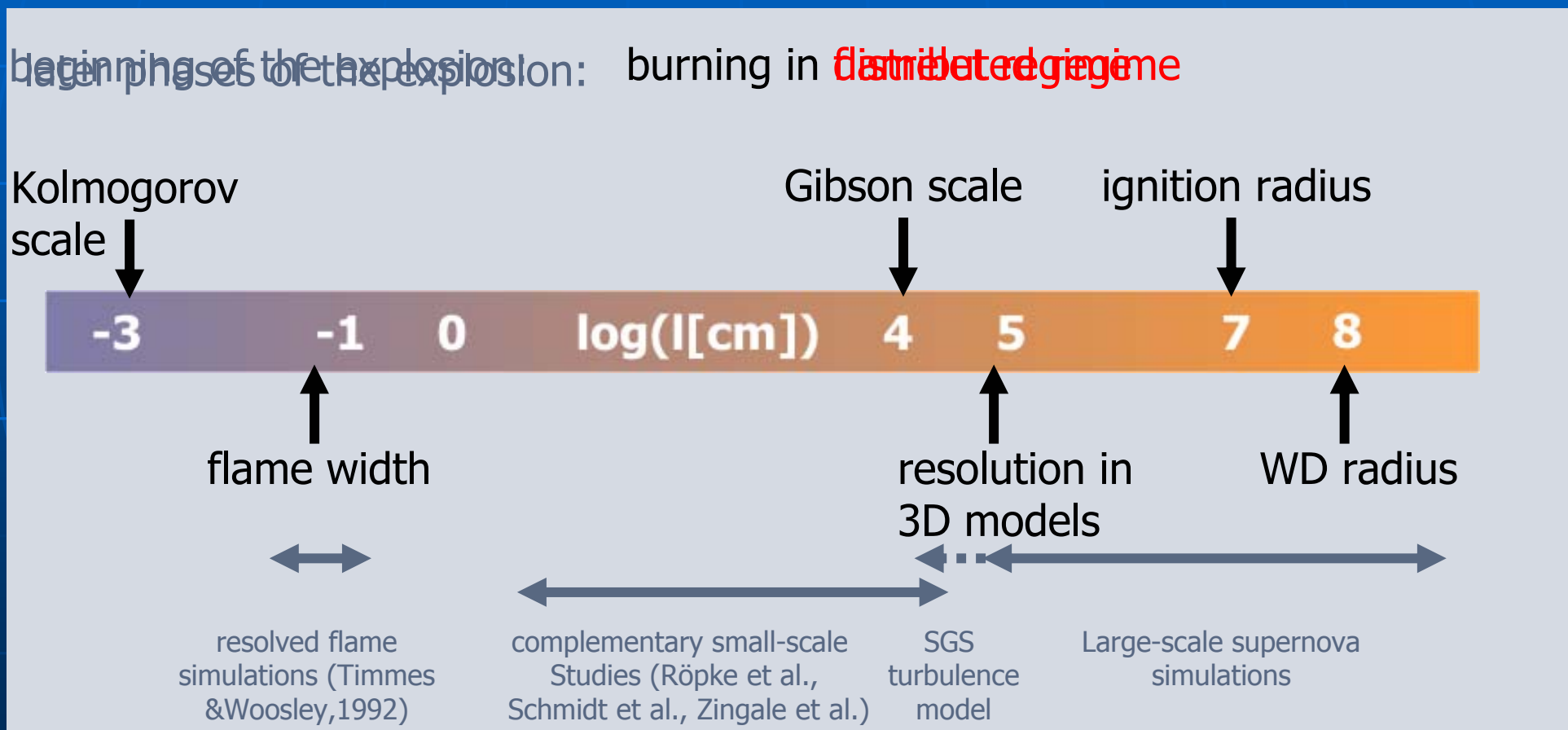
$$\Rightarrow Re = Da^2 \cdot Ka^2$$

$\log(v'/s_L)$



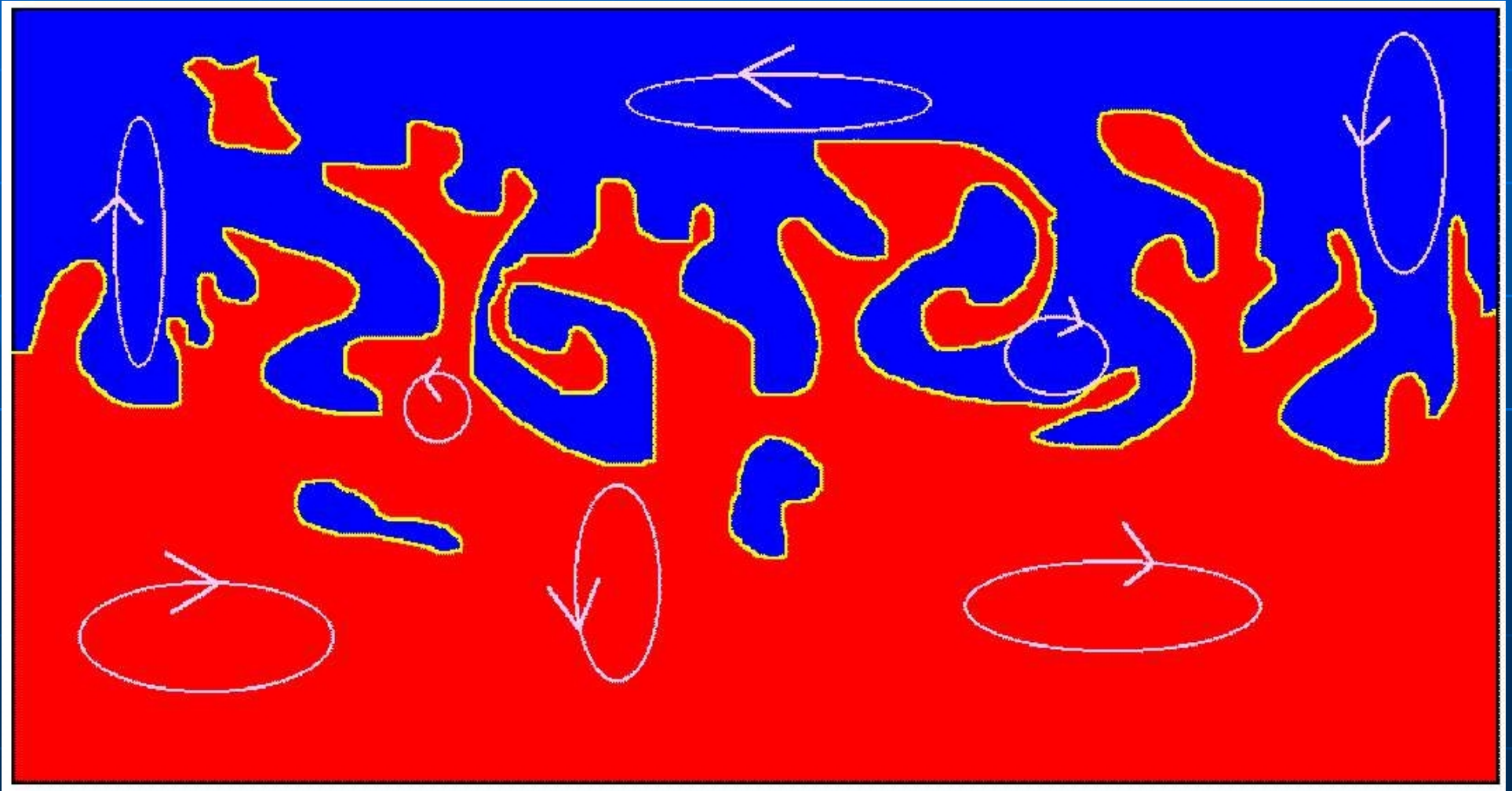
Simulating the relevant scales

- Gibson scale $s_L = v'$: below turbulence does not affect flame propagation



Burning regimes of pre-mixed flames

1. Cellular burning, wrinkled flamelets



Burning regimes of pre-mixed flames

1. Cellular burning, wrinkled flamelets

$$u_{\text{cell}} = s_L [1 + \varepsilon(\mu)] ; \mu = \rho_b / \rho_u ,$$

$$\varepsilon(\mu) \approx 0.41 (1 - \mu)^2$$

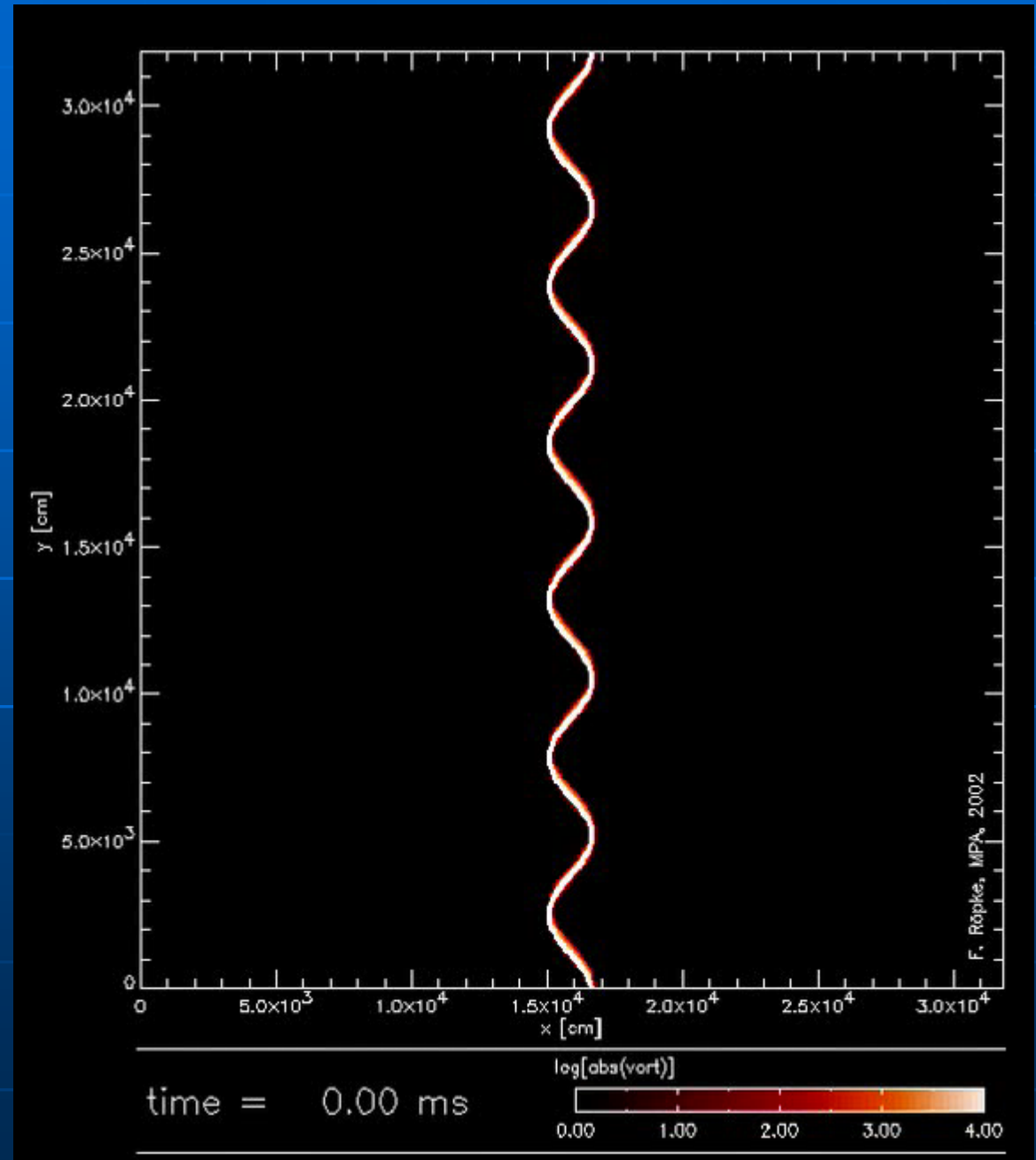
Or: “Fractal model”

$$u_{\text{cell}}(1) = s_L (1/l_{\text{crit}})^{D-1}$$

The Landau-Darrieus instability and its interaction with turbulence:

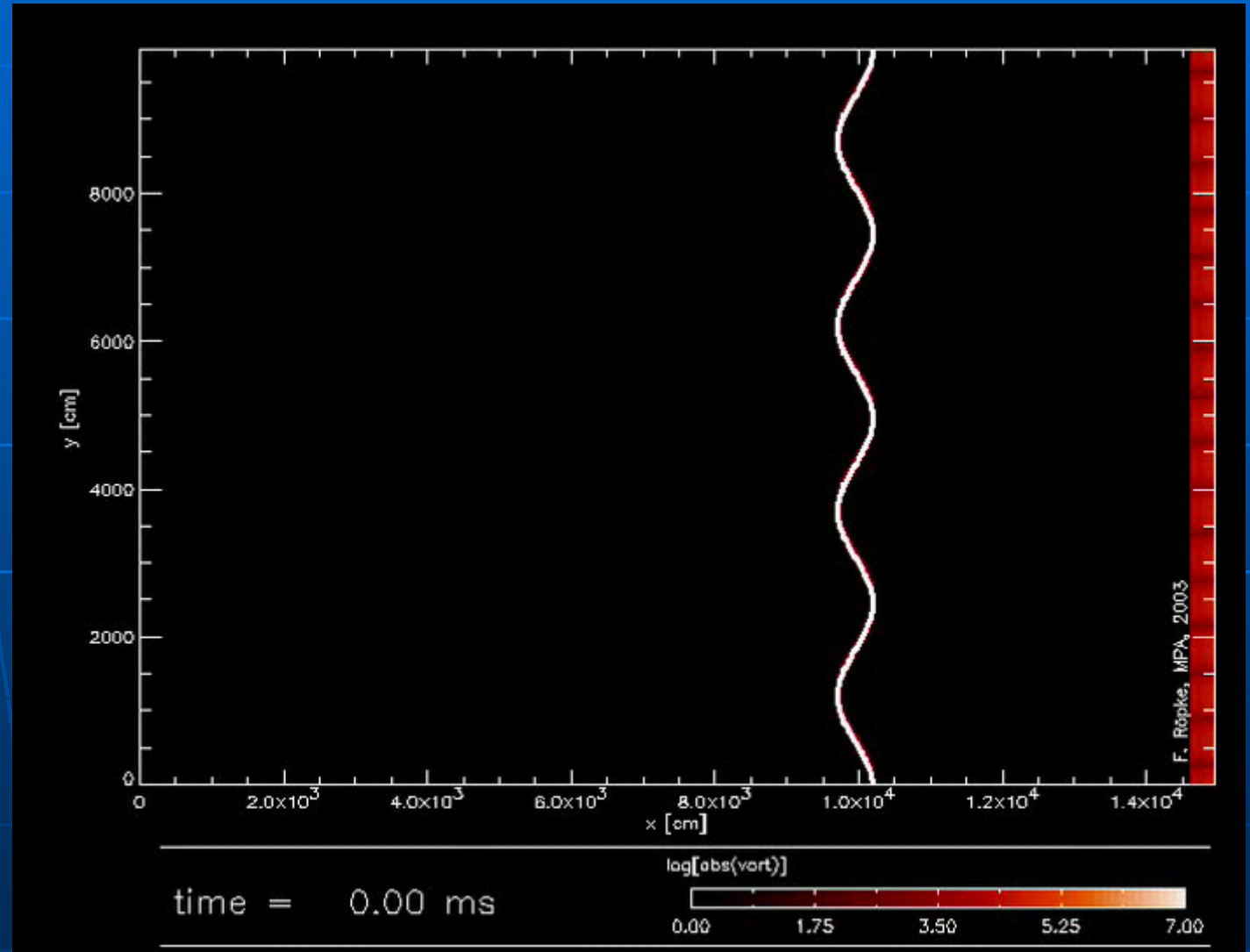
Quiescent fuel

(Röpke et al., 2003a)



The Landau-Darrieus instability and its interaction with turbulence:

*Strong vortical
flow*



(Röpke et al.,
2003b)

Burning regimes of pre-mixed flames

2. The corrugated flamelet regime

Transition at the “Gibson scale”:

$$v(l_{\text{Gibs}}) = u_{\text{cell}}(l_{\text{Gibs}})$$

In the limit of strong turbulence:

$$s_{\text{turb}}(l) \approx v'(l), \quad l > l_{\text{Gibs}} \quad (\text{independent of } s_L!!!)$$

$$d_{\text{turb}} \approx l \quad (\text{“turbulent flame brush”})$$

Fully developed turbulence?

3-D “direct”
numerical simulations
of flames moving in
white dwarf matter:

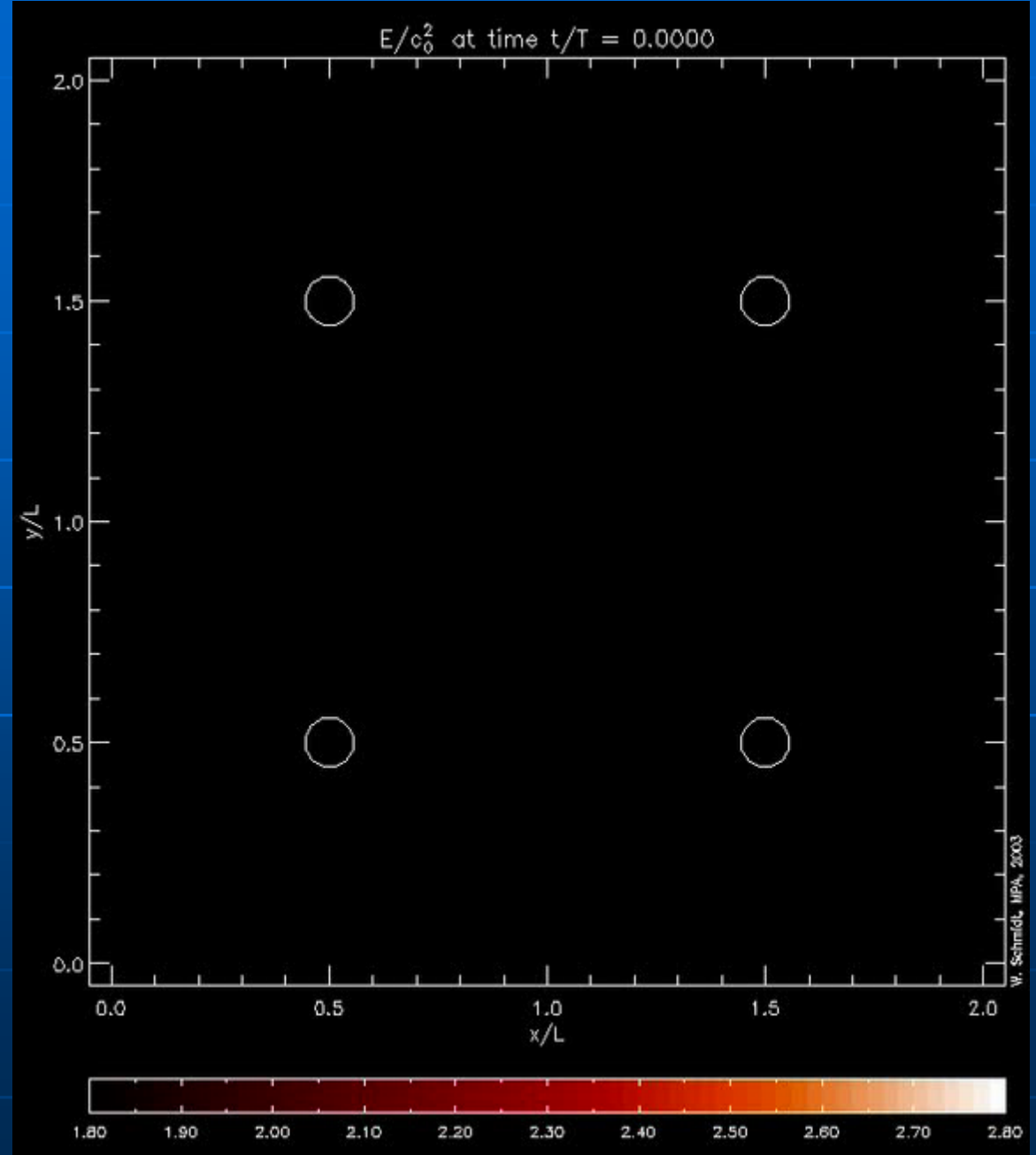
Energy

$$\rho = 2.9 \cdot 10^9 \text{ gcm}^{-3}$$

$$V/s_{lam} = 4$$

$$V/c_0 = 0.043$$

(Schmidt et al., 2004)



Fully developed turbulence?

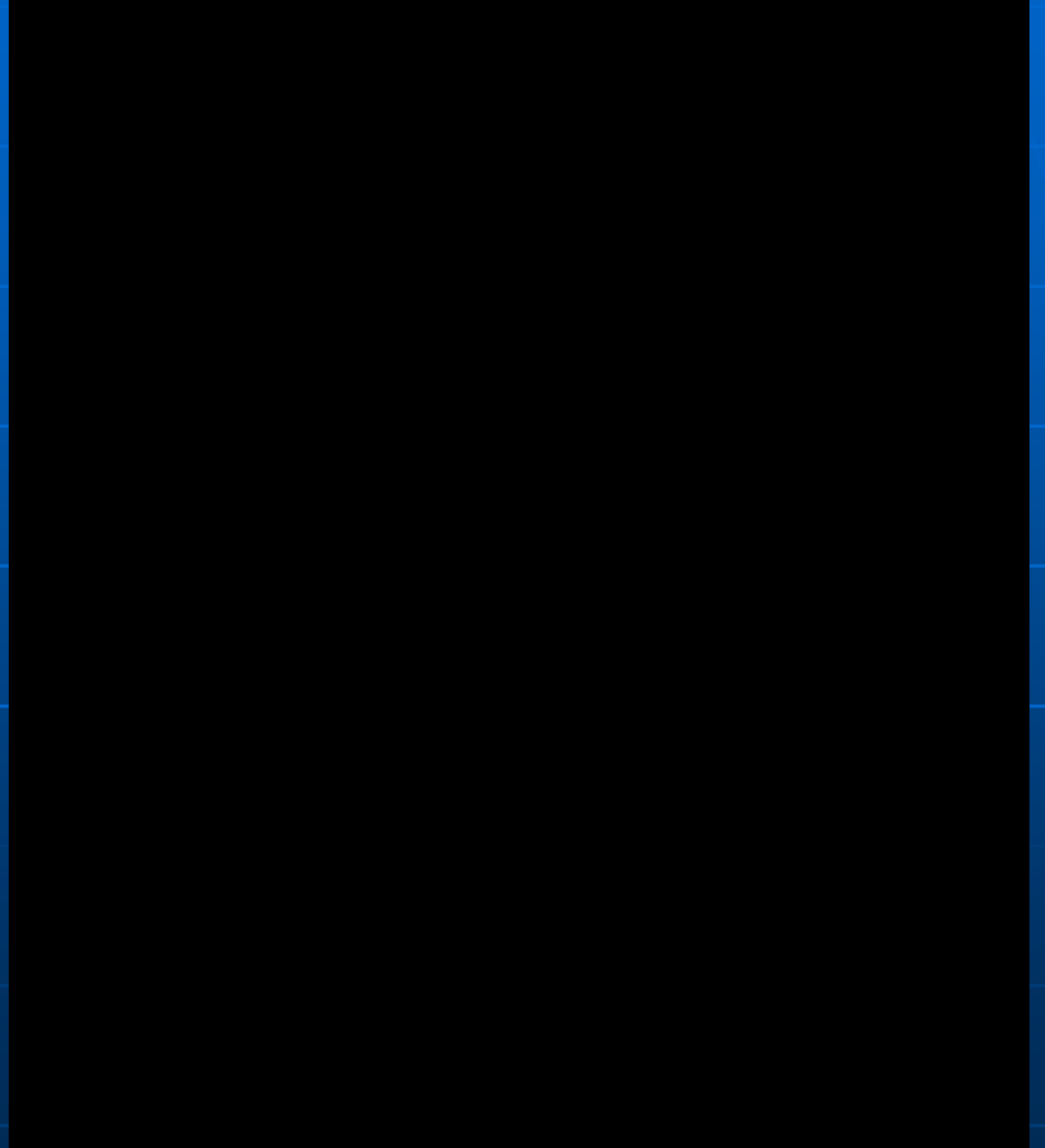
3-D “direct”
numerical simulations
of flames moving in
white dwarf matter:
Vorticity

$$\rho = 2.9 \cdot 10^9 \text{ gcm}^{-3}$$

$$V/s_{\text{lam}} = 4$$

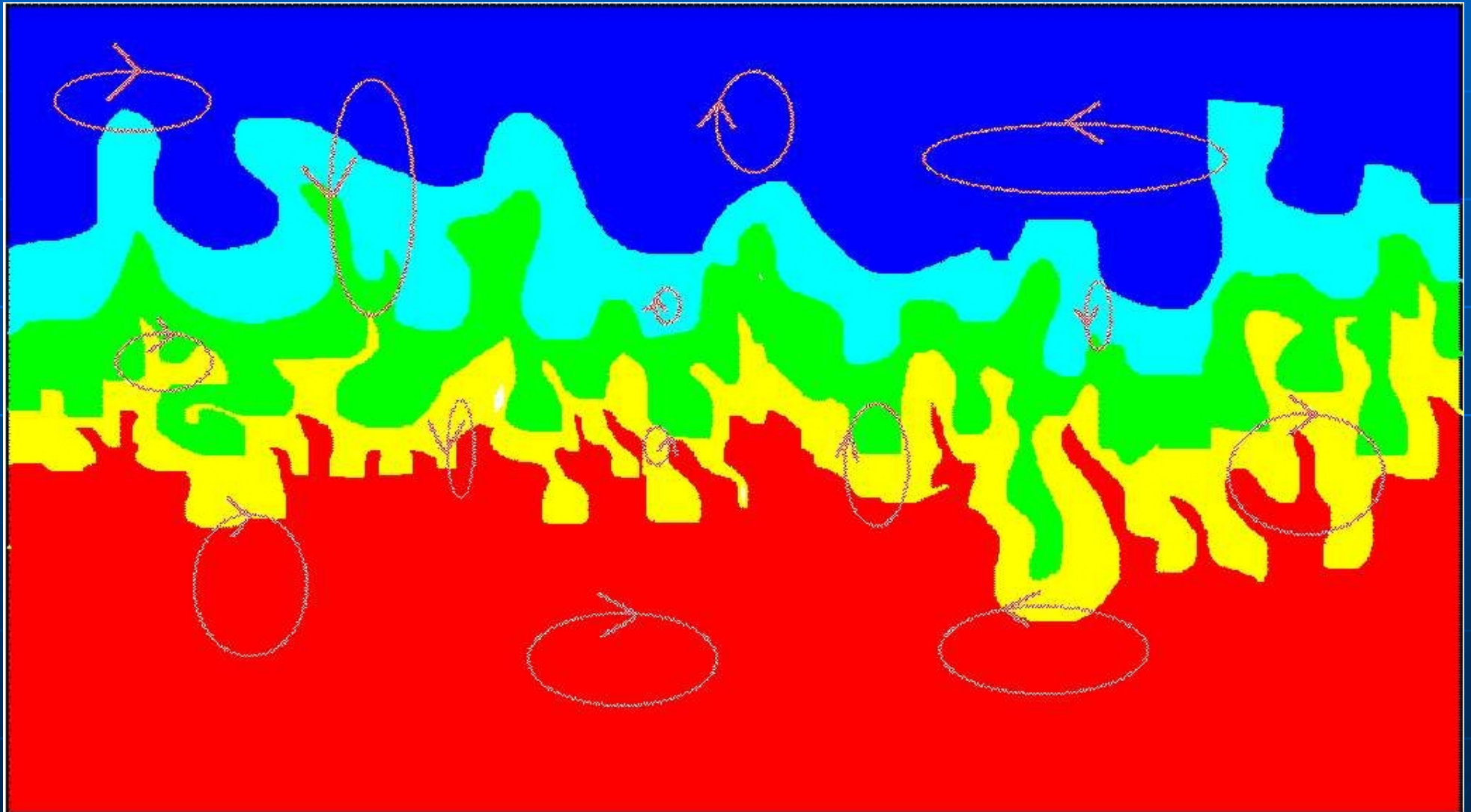
$$V/c_0 = 0.043$$

(Schmidt et al., 2004)



Burning regimes of pre-mixed flames

3. The distributed-burning



Burning regimes of pre-mixed flames

3. The distributed-burning

Turbulent eddies interact with the flame:

$$l_F \geq l_{Gibbs}$$

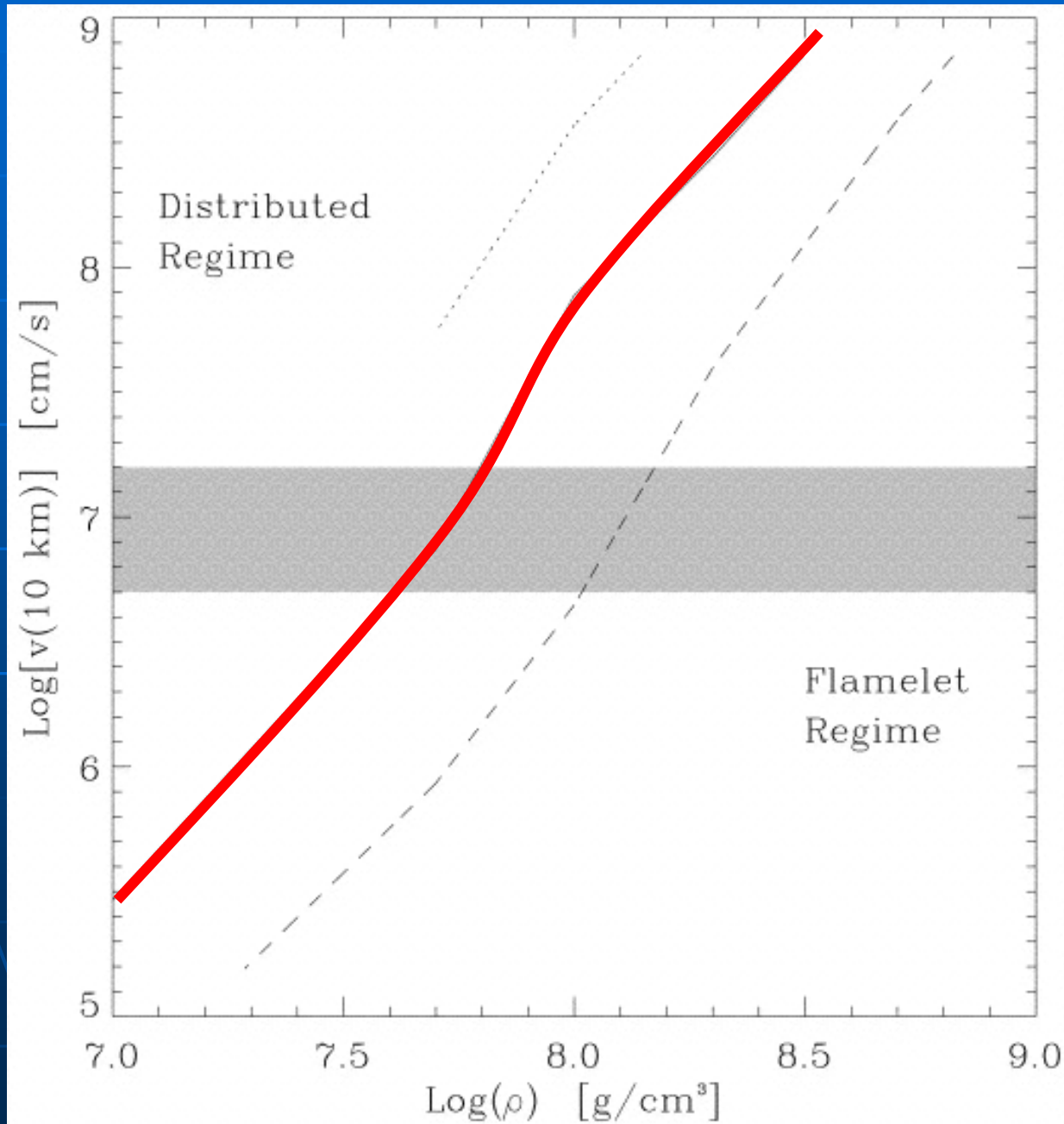
Rough estimate (“Damköhler scaling”):

$$s_{\text{turb}}/s_L \approx \text{const} (D_t/D)^{1/2} \text{ (dependent on } s_L \text{ !!!)}$$

$$\text{const} = O(1)$$

Transition to detonation possible???

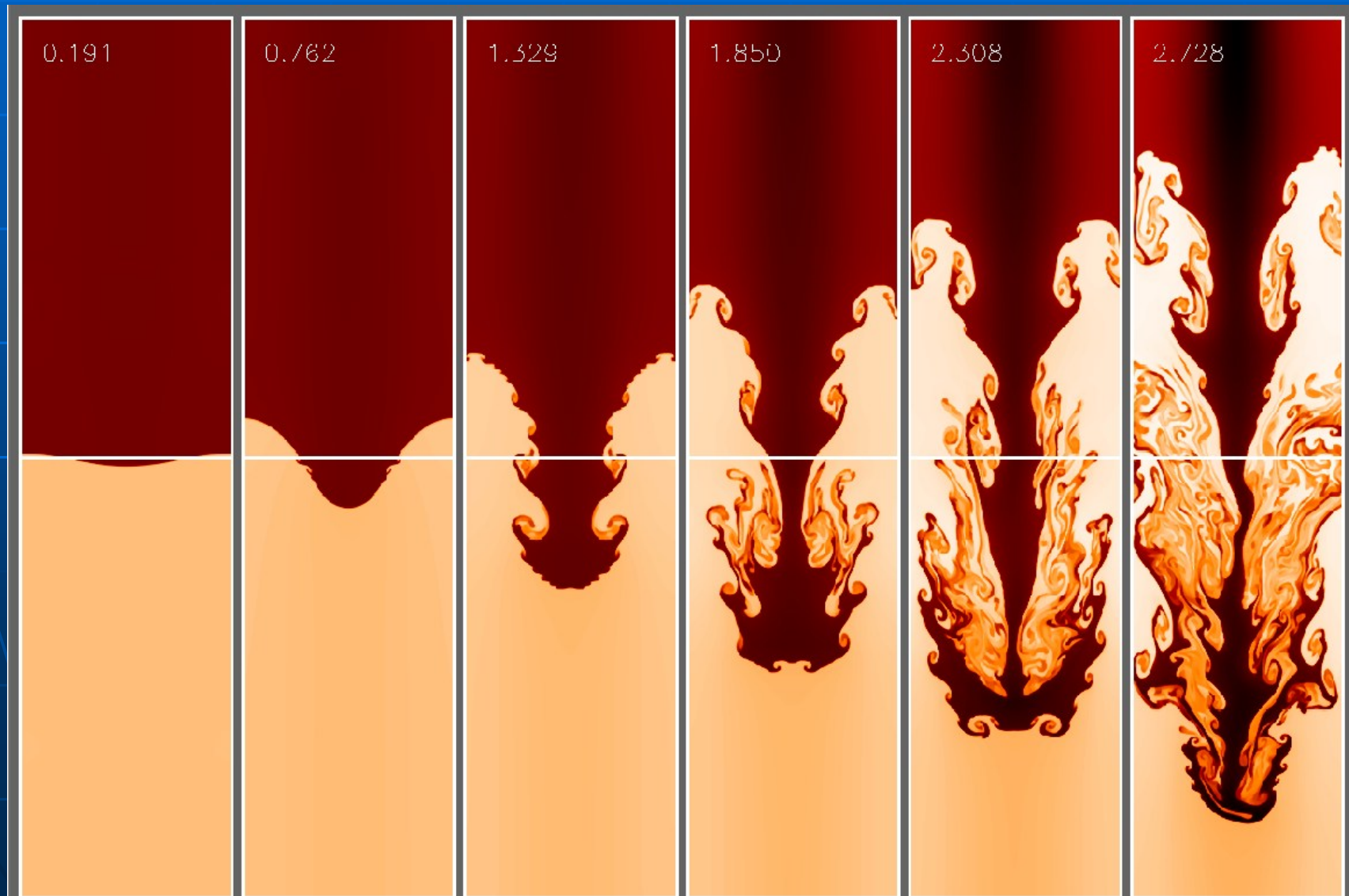
Application to type Ia supernova



Niemeyer &
Woosley (1997)

Burning regimes of pre-mixed flames

4. The Rayleigh-Taylor regime



Burning regimes of pre-mixed flames

4. The Rayleigh-Taylor regime

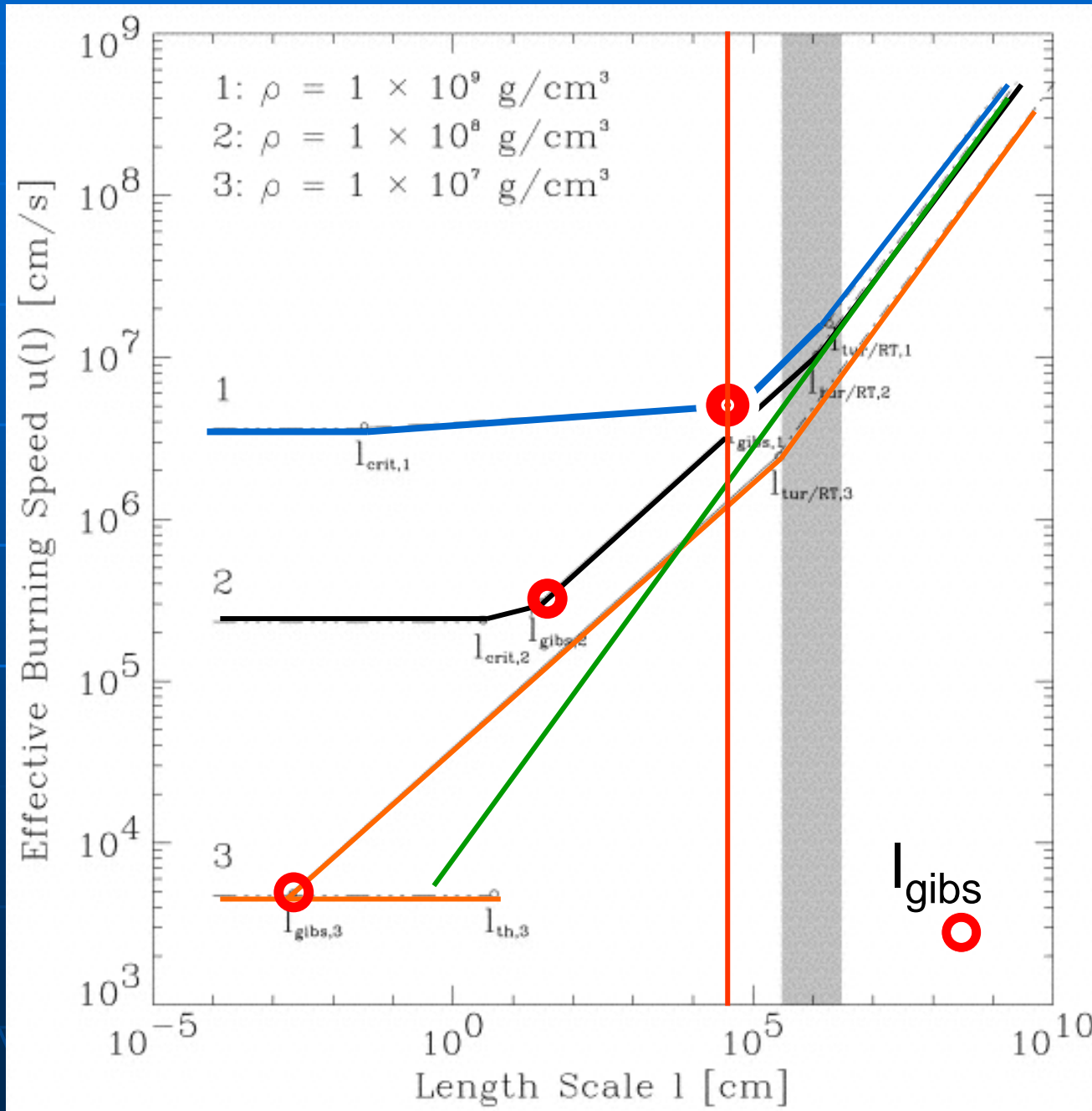
$$v_{RT} = B \sqrt{g_{eff} l} ; B \approx 0.5 ; g_{eff} = At \cdot g$$

Sharp-Wheeler model:

$$r_{sw} \approx 0.05 g_{eff} t^2 ; v_{sw} \approx 0.1 g_{eff} t ;$$

$$l_{tur/RT} \approx 10^6 \text{ cm}$$

Effective burning velocities in SN Ia



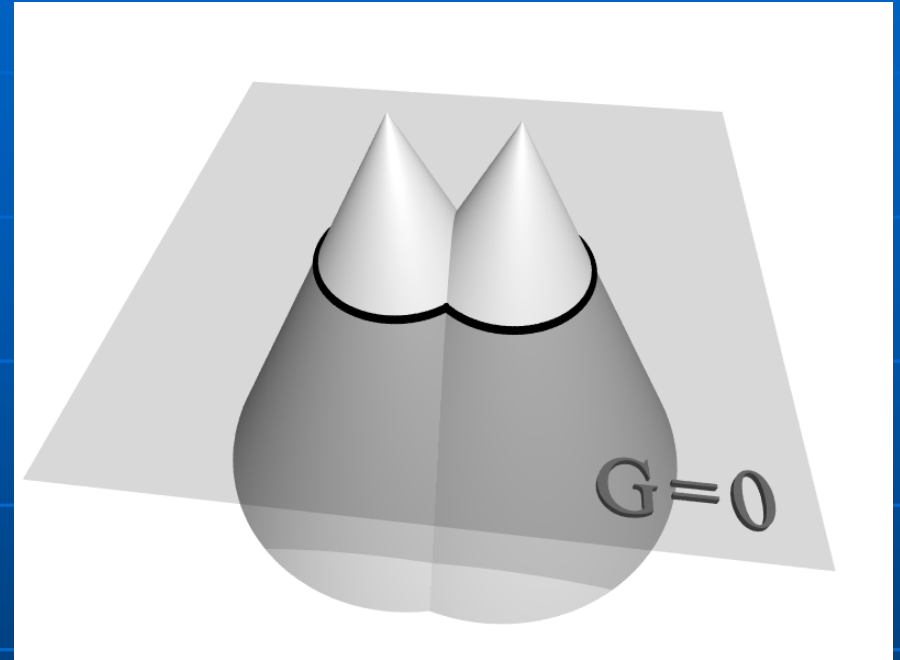
Niemeyer
& Woosley
(1997)

How to model thermonuclear flames?

- The "flames" cannot be resolved numerically.
- The amplitudes of turbulent velocity fluctuations in the length scale of the flame are determined on the integral scale.



"LES" + "Level Set"

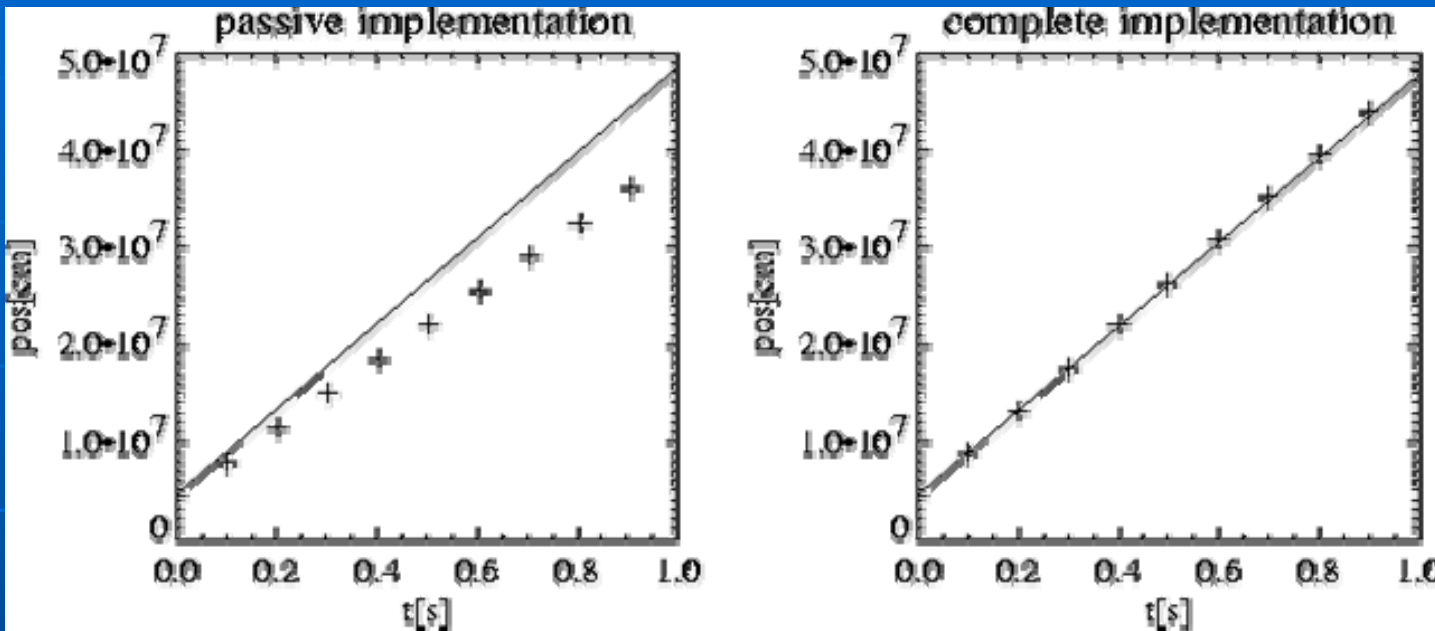


$$\partial G / \partial t = -\mathcal{D}_f \nabla G$$

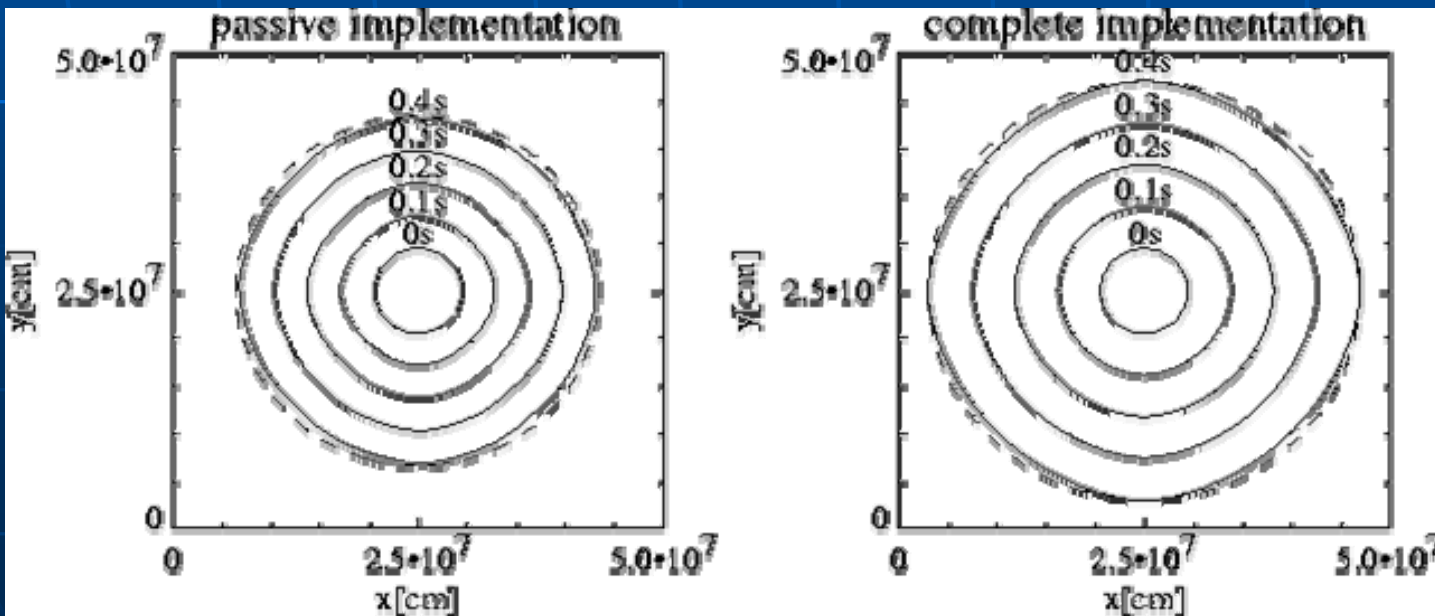
$$\mathcal{D}_f = \mathbf{v}_u + s_{\text{tur}} \mathbf{n}; \quad |\nabla G| = 1$$

Some test of the code

Planar flame

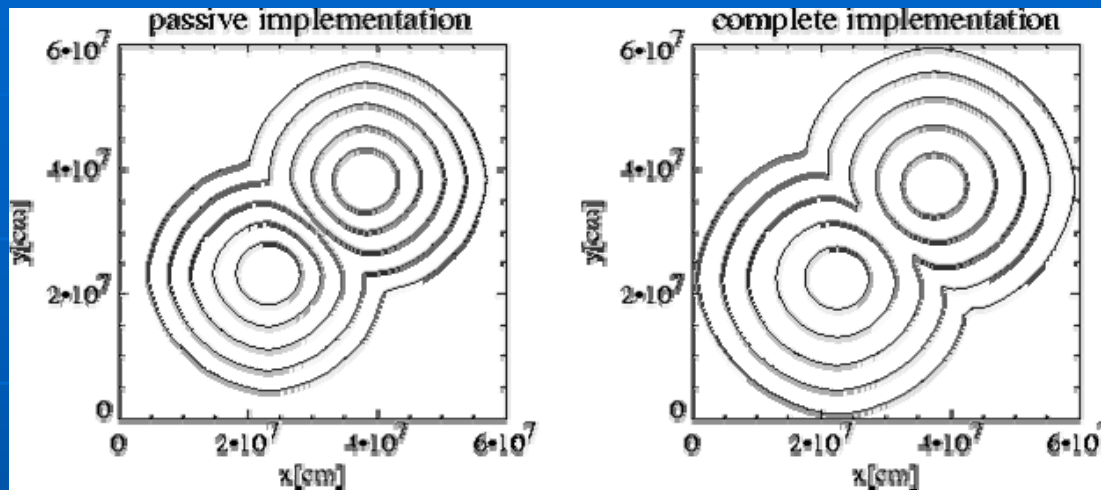


Circular flame

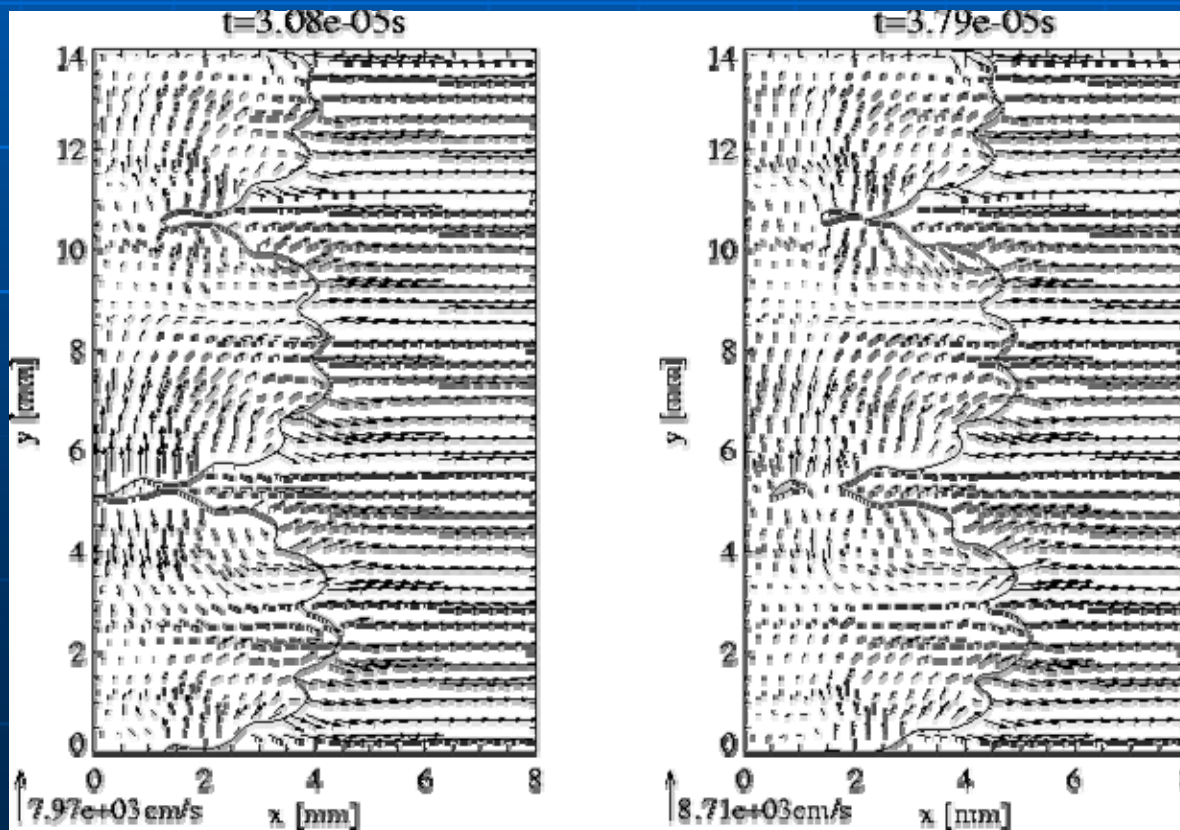


Reinecke et al.
(1999)

Some test of the code (ctn.)



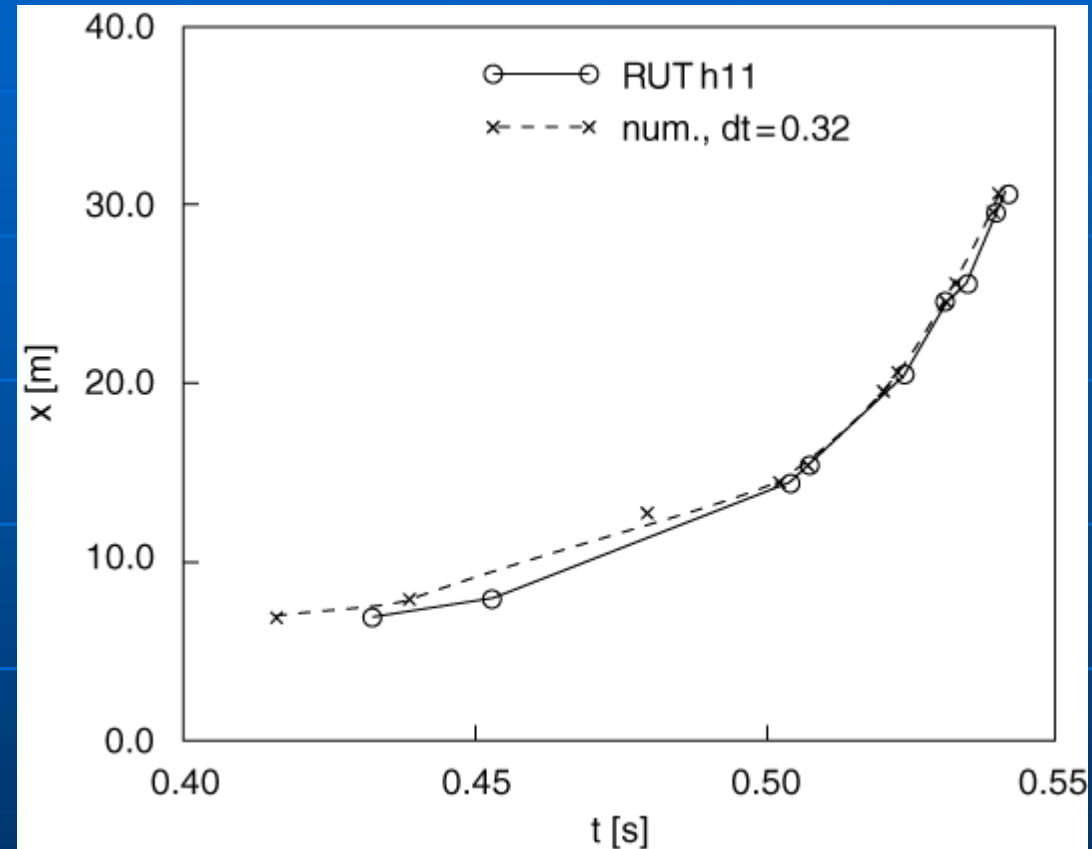
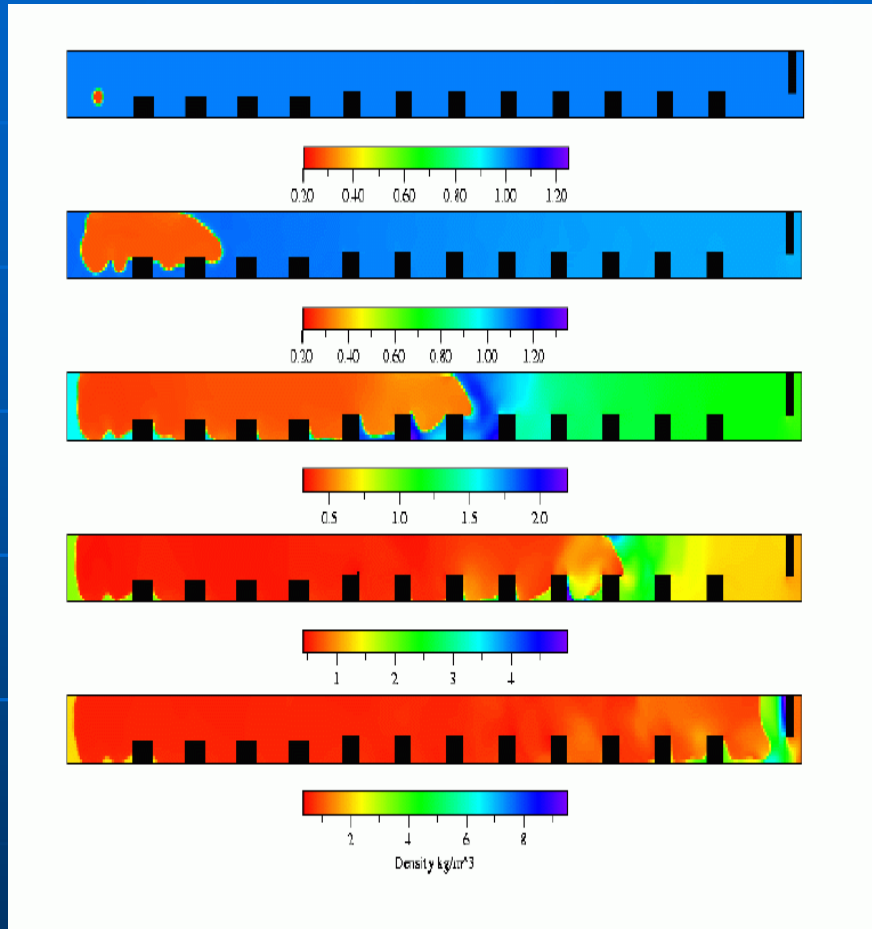
Merging circular flames



Hydrogen-in-air flames

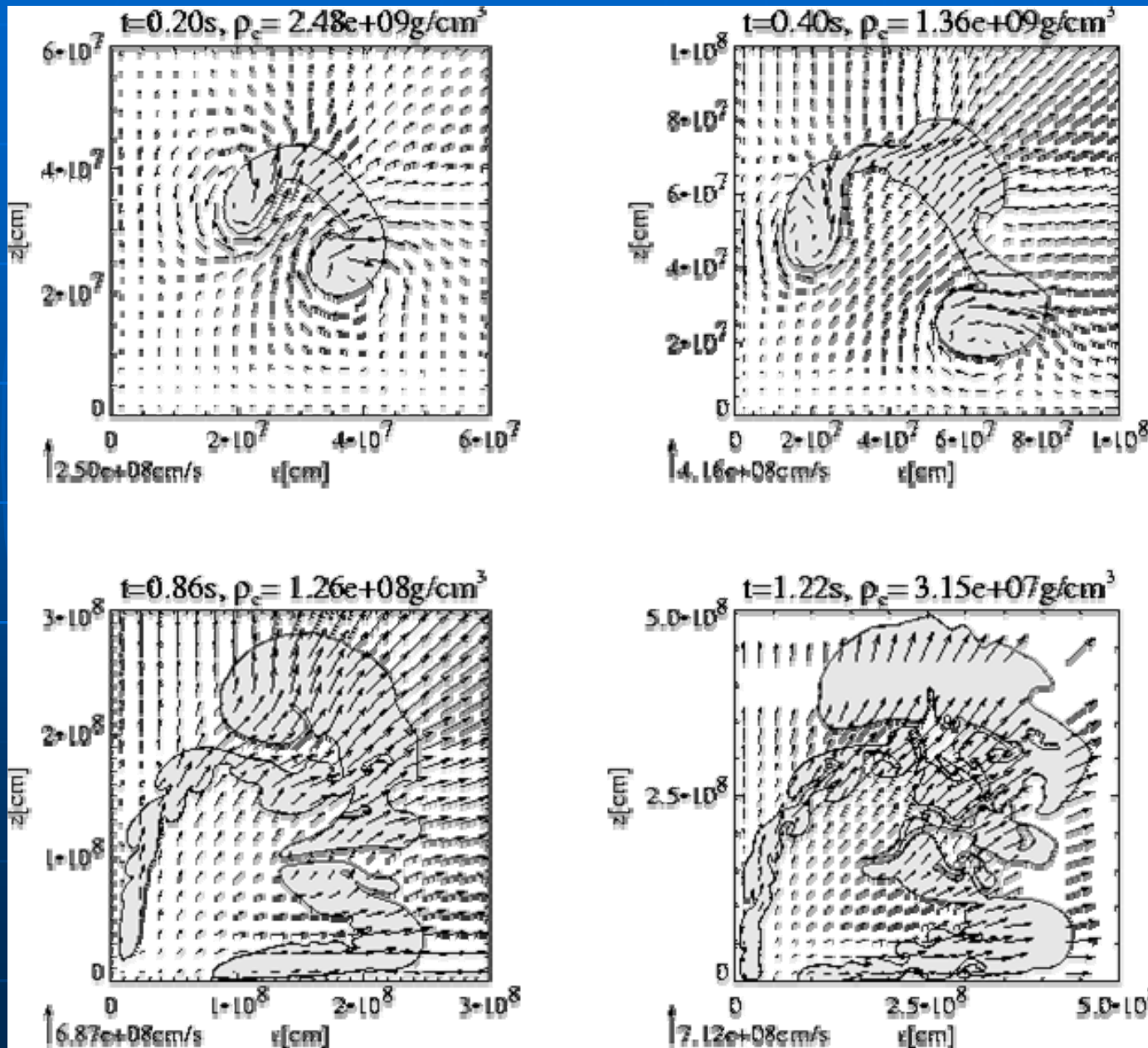
Reinecke et al. (1999)

Application to laboratory flames (hydrogen in air)



The method can reproduce terrestrial experiments well!
(Smiljanowski et al. 1997)

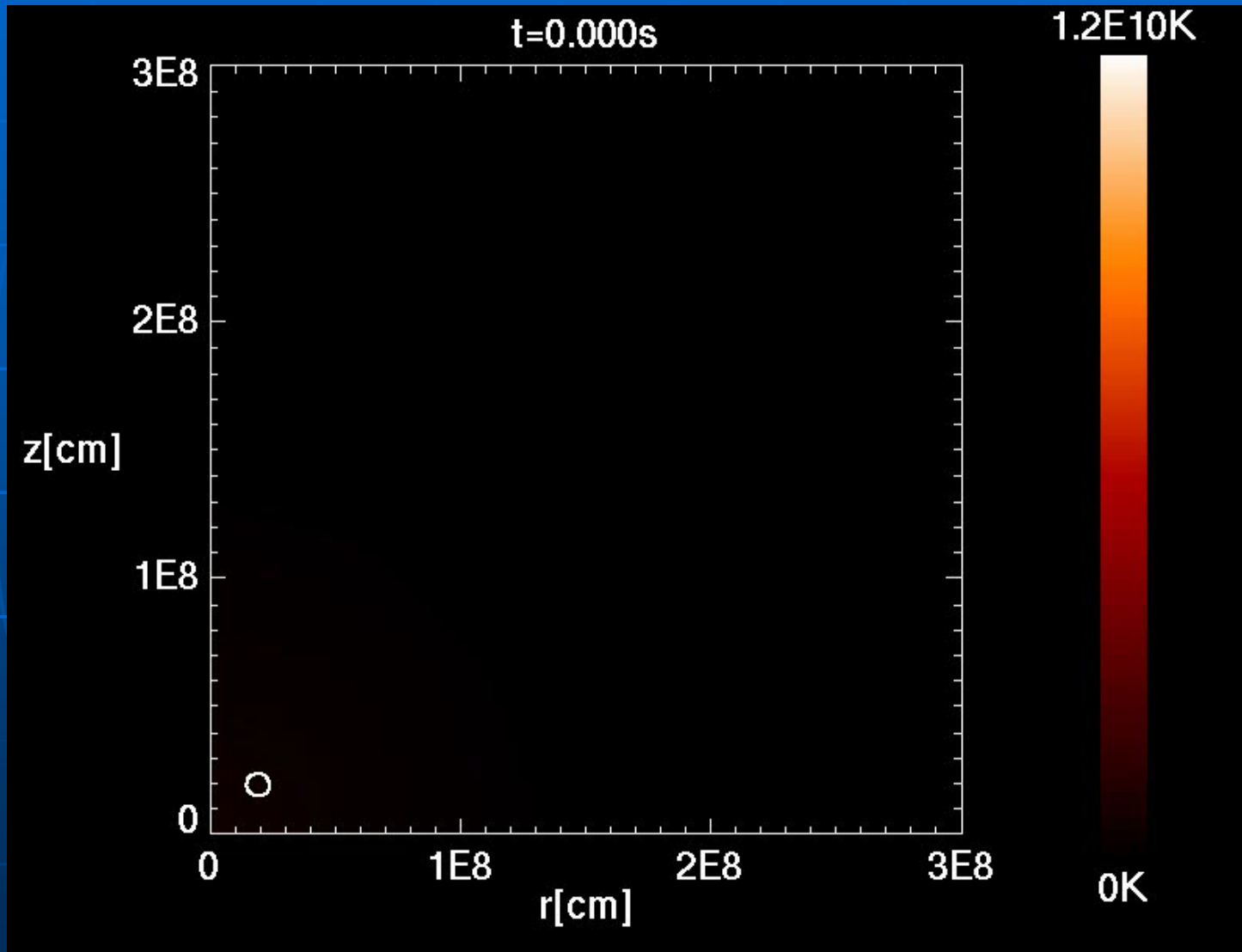
Application to the SN Ia problem



One rising
blob (in 2D)

Reinecke et al.
(1997)

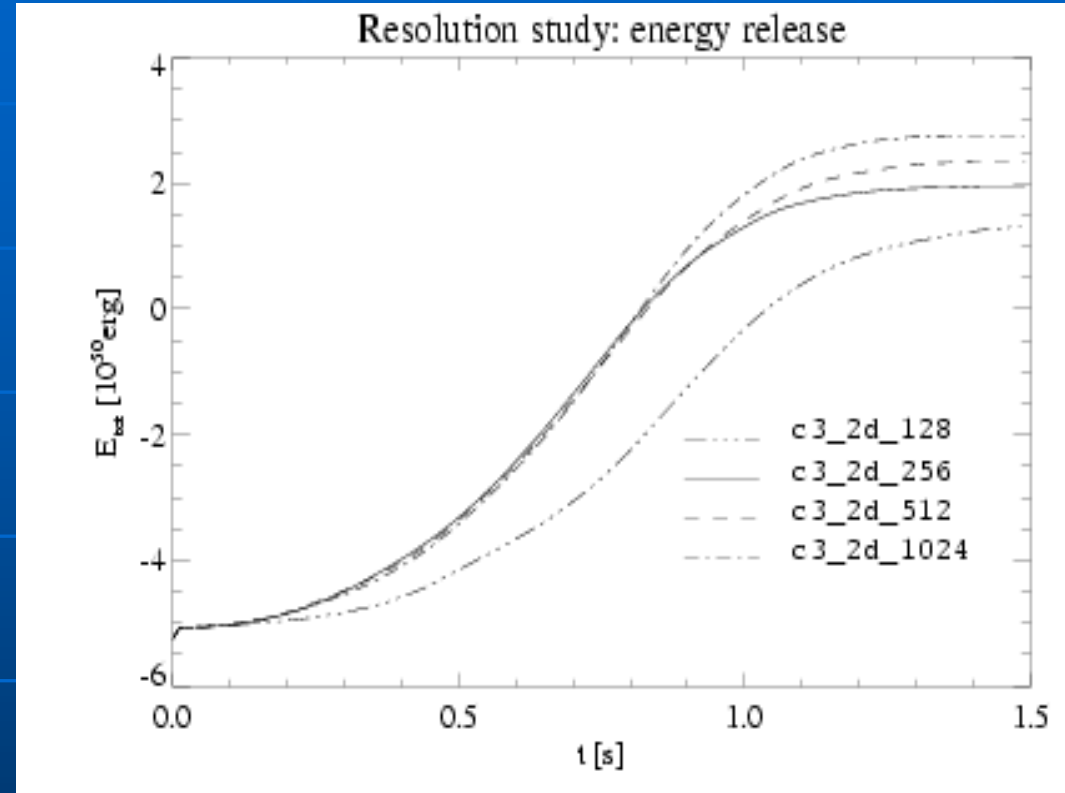
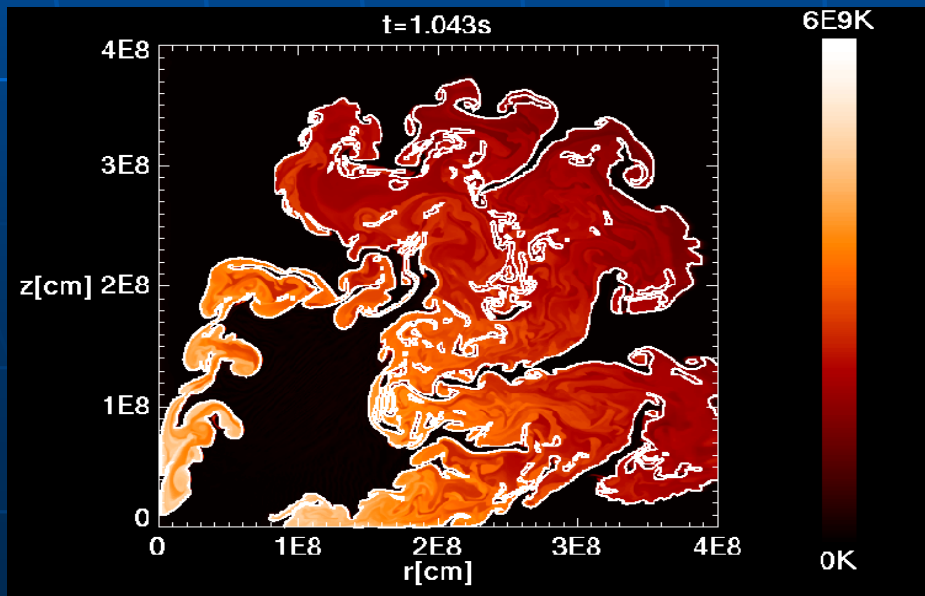
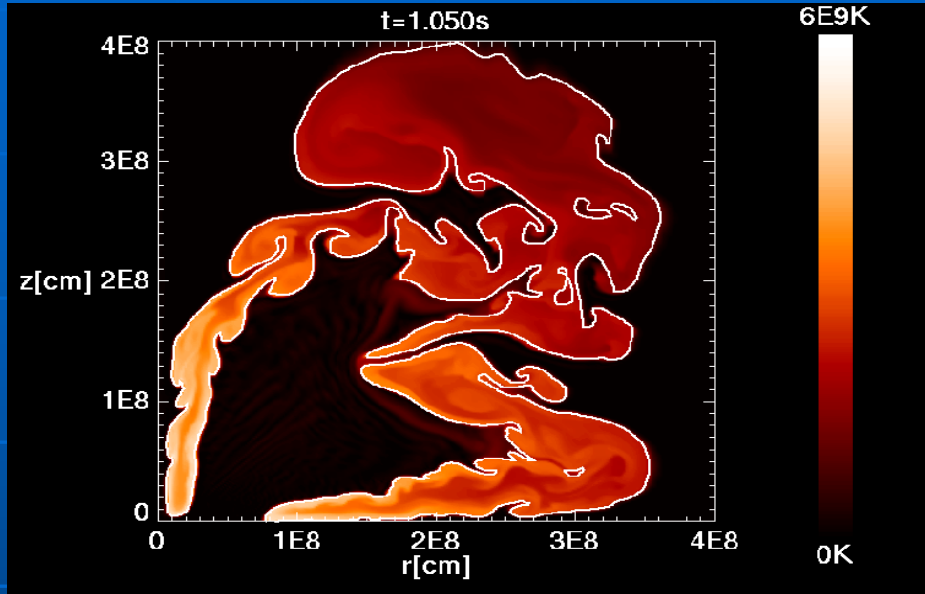
Application to the SN Ia problem



One rising
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Reinecke et al.
(1997)

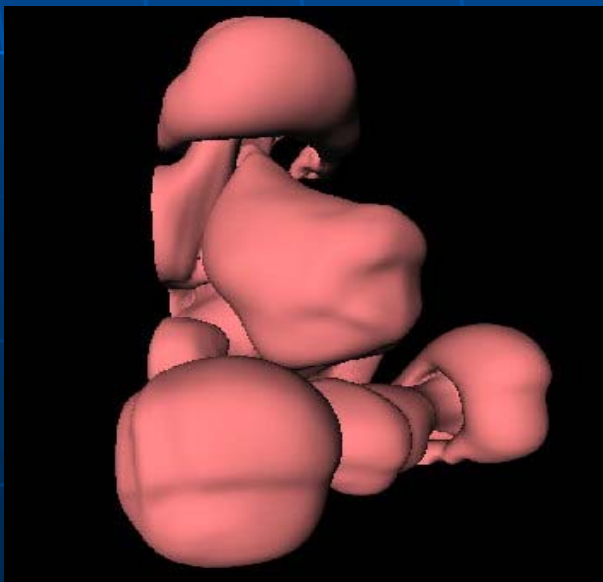
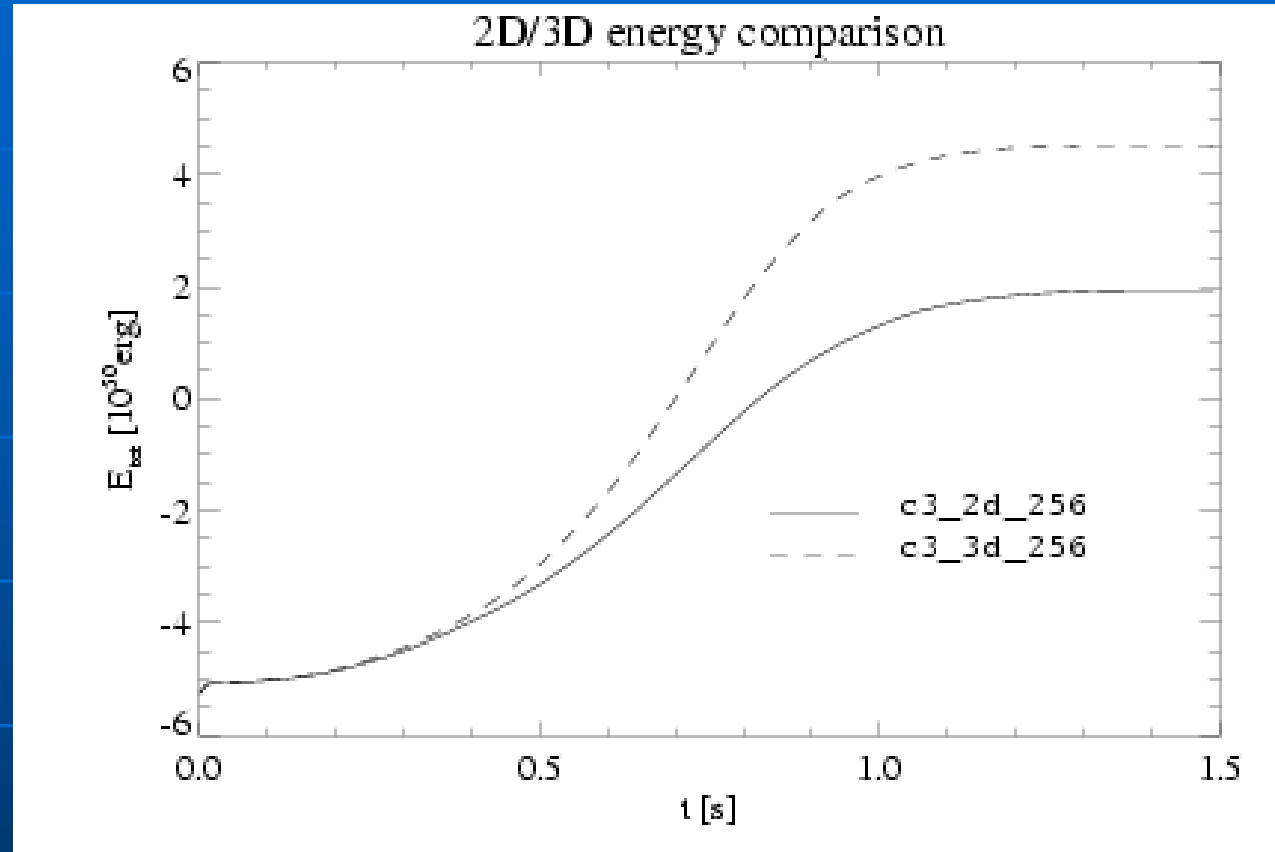
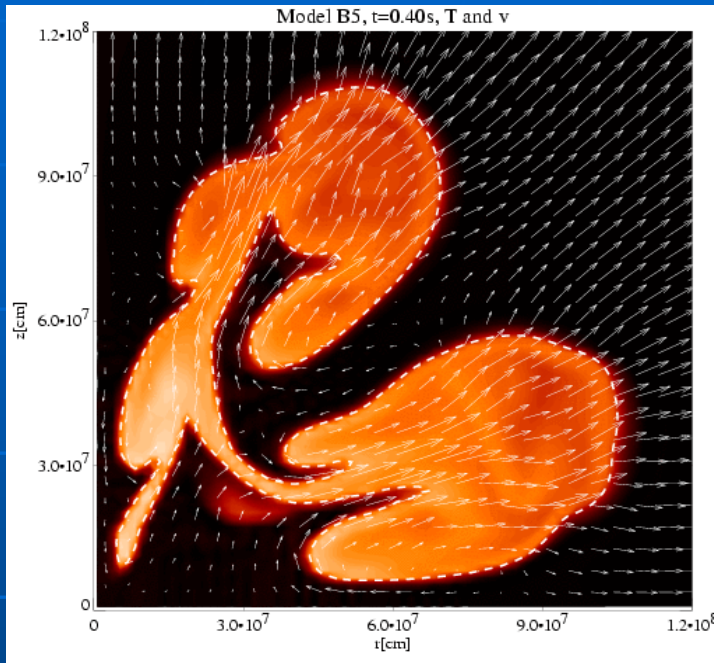
Convergence tests in 2D



Global results are independent of the numerical resolution!

Reinecke et al. (1999, 2002)

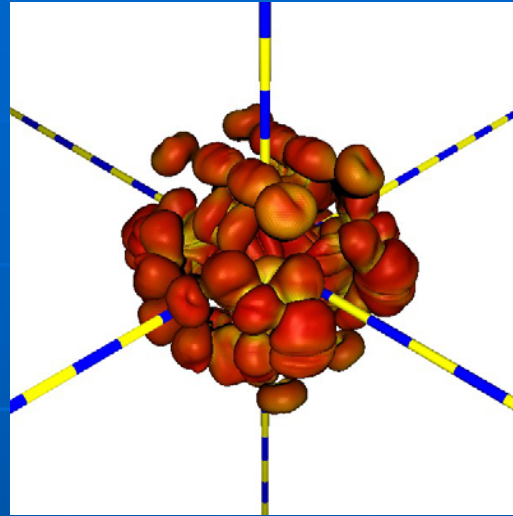
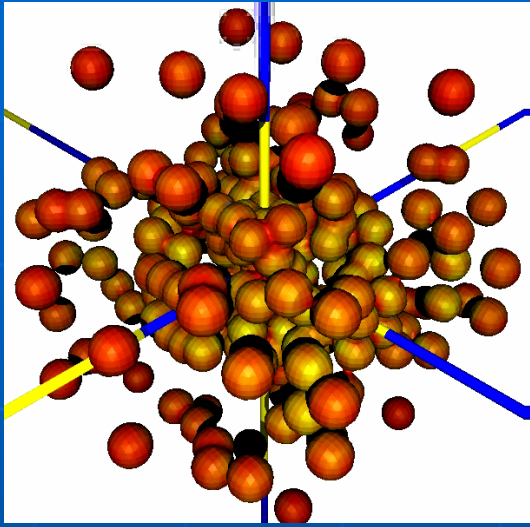
2D \Rightarrow 3D



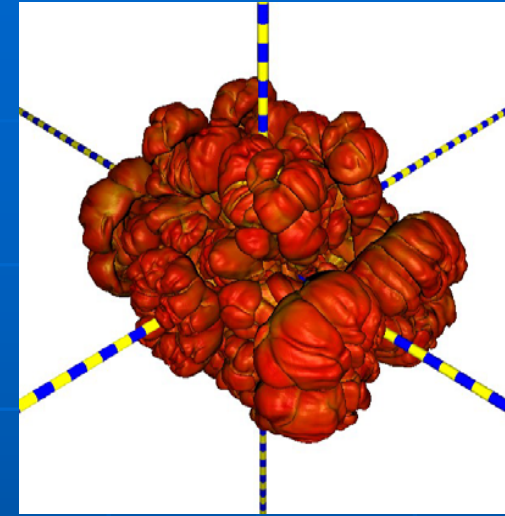
Because of larger surface area:
More energy is produced!

Reinecke et al. (2001)
(See also Gamezo et al., 2003)

3D models: The best we could do (until recently):

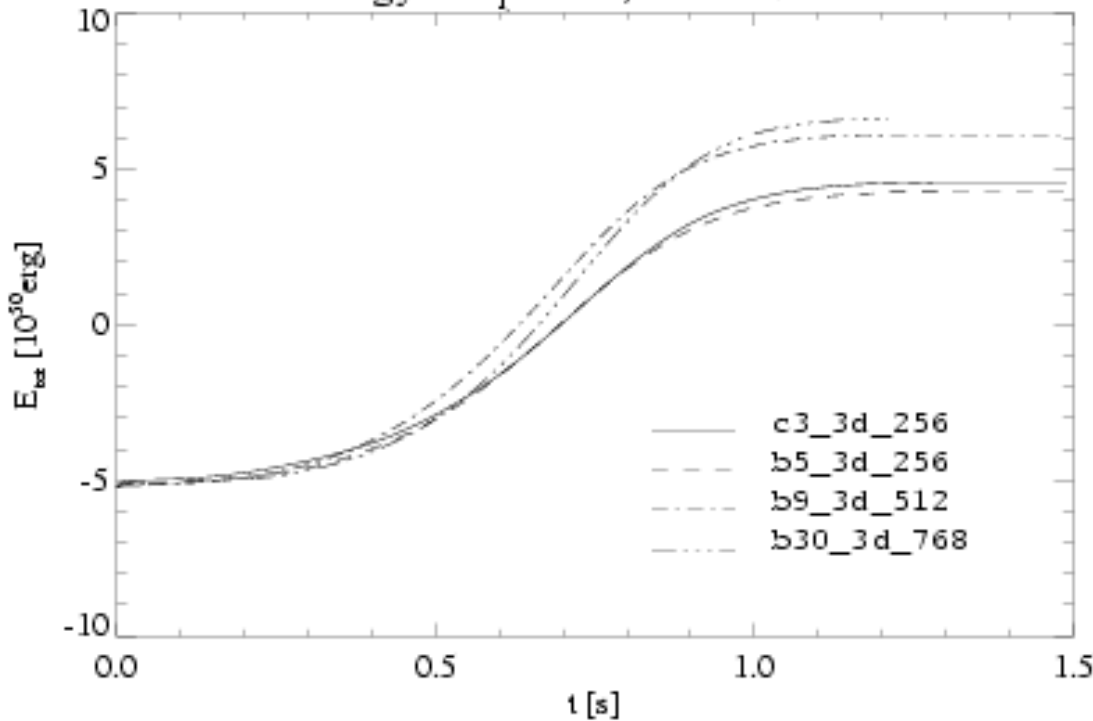


0.25s



0.6s

Energy comparison, 3D simulations



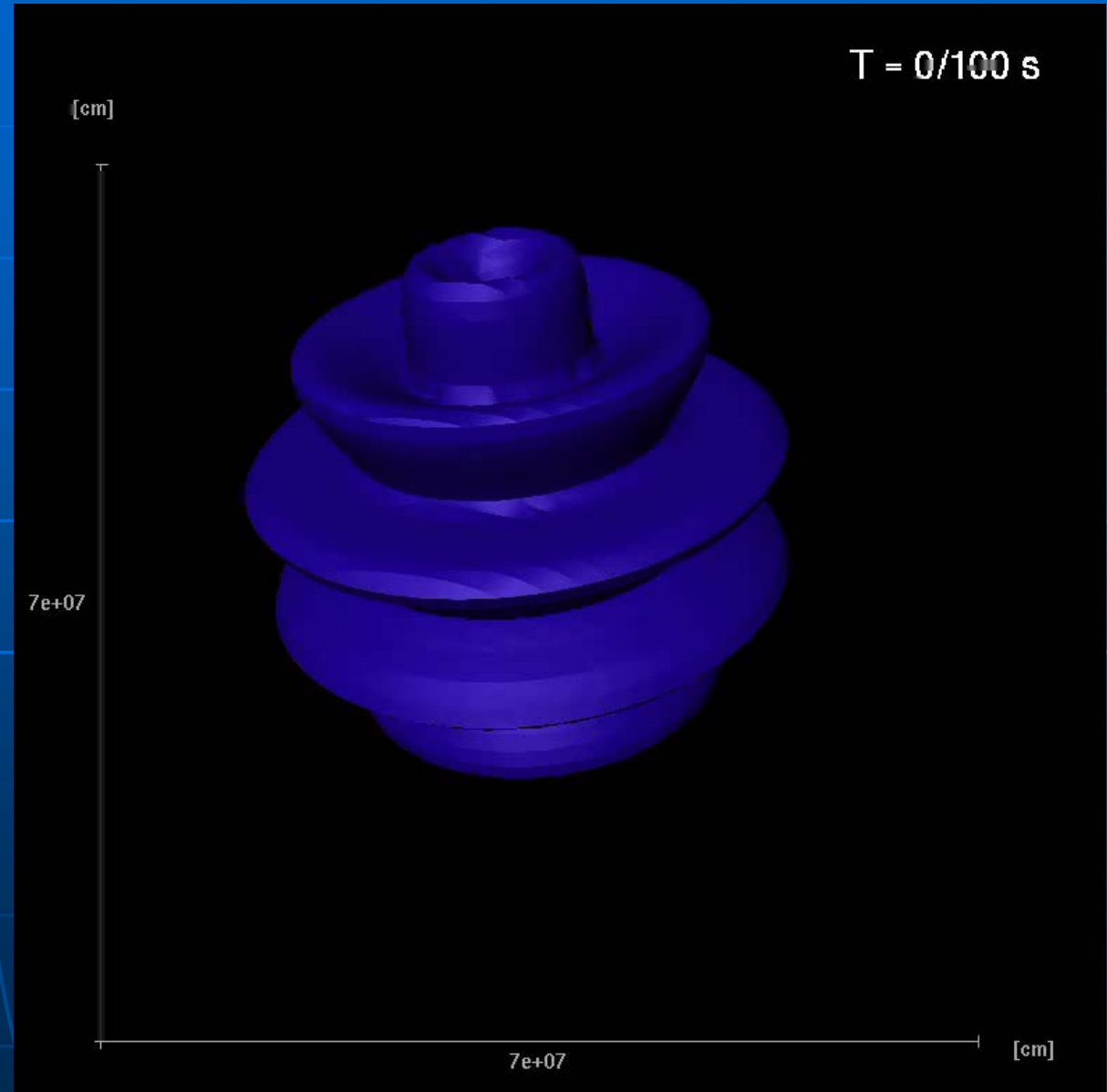
Mod b30_3d

(Reinecke et al., 2003)

Recent modifications of the code:

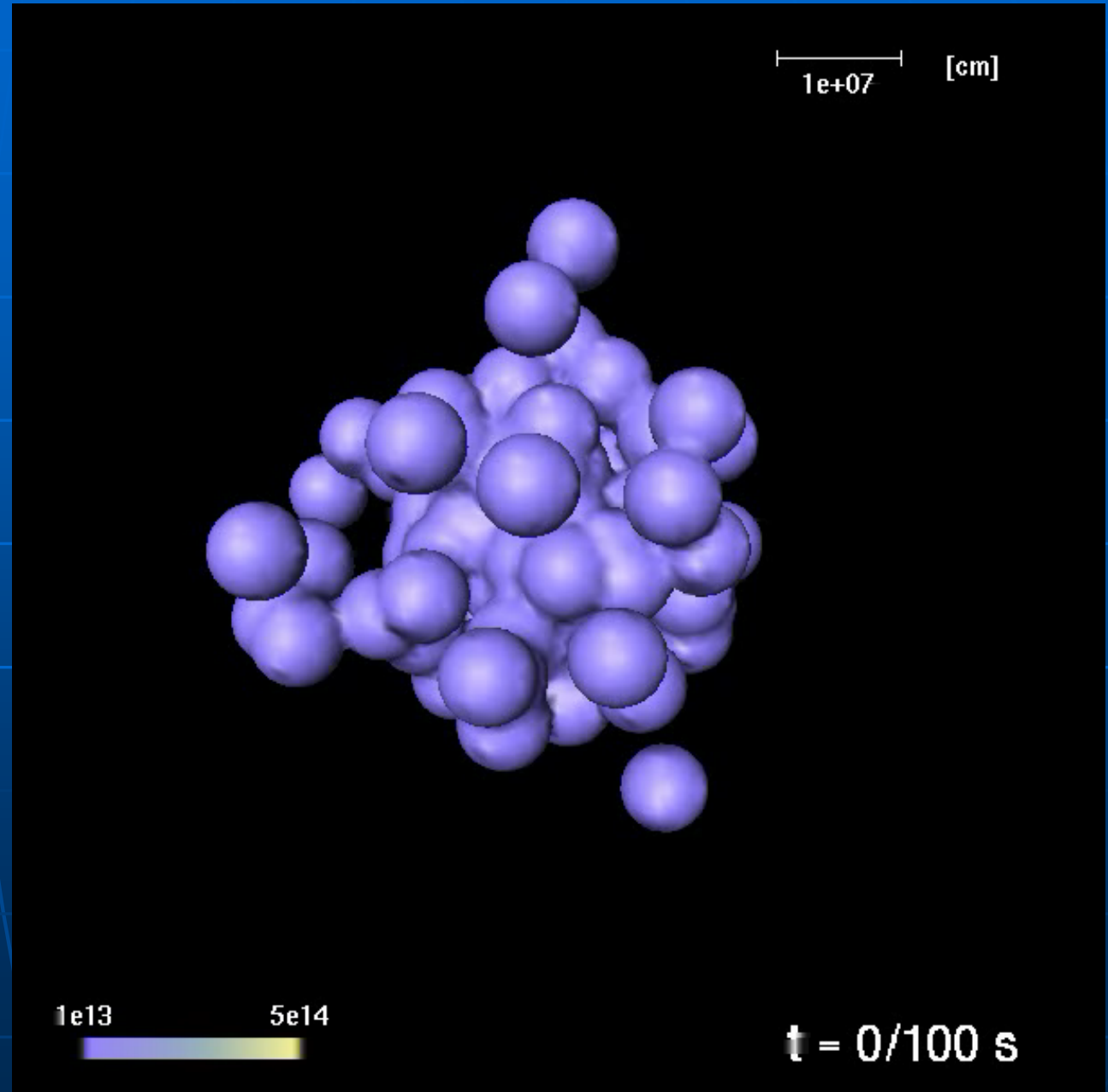
1. Moving grid

Röpke (2004)



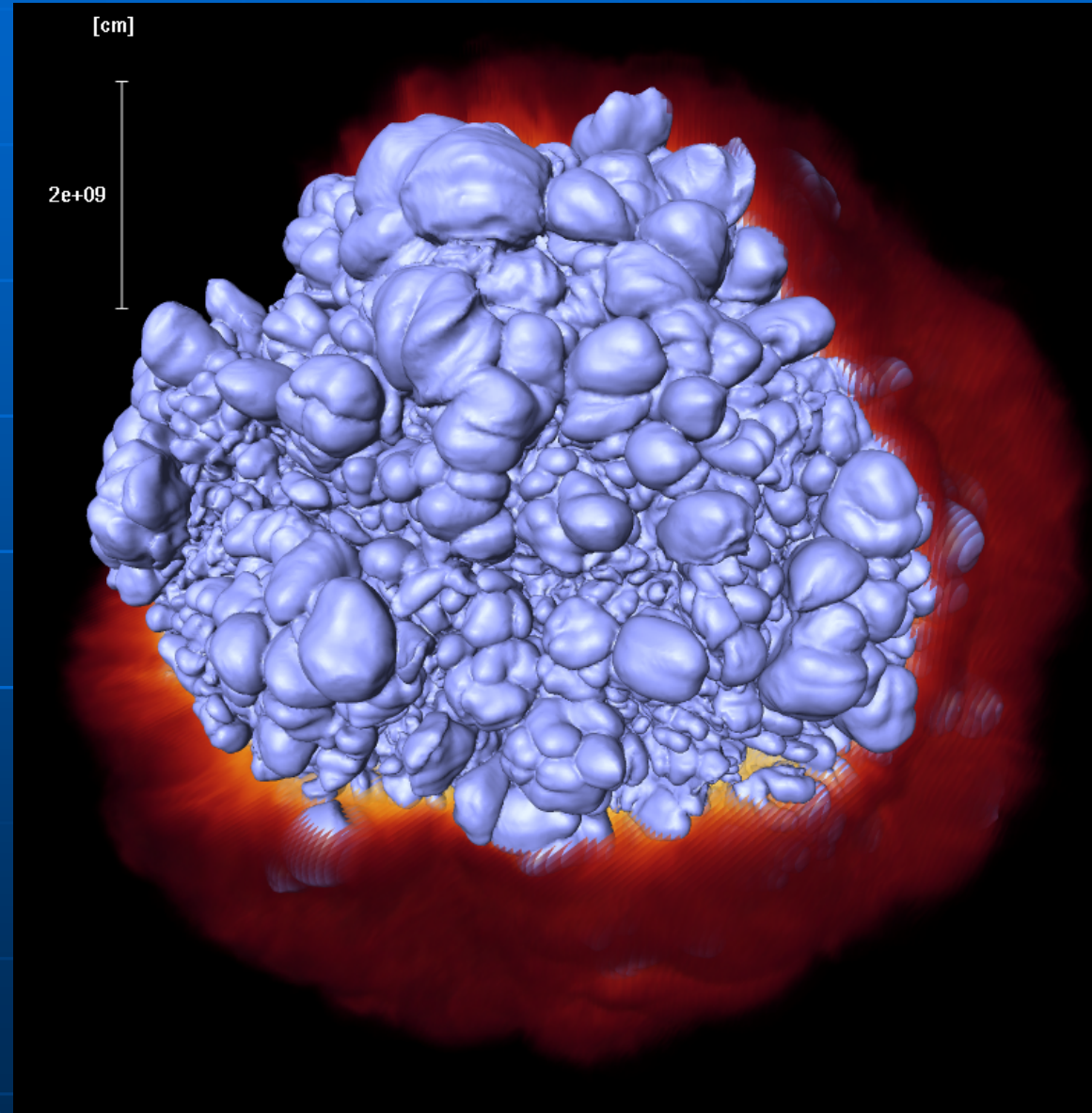
2. Full star (“ 4π ”)

Röpke & Hillebrandt
(2004)

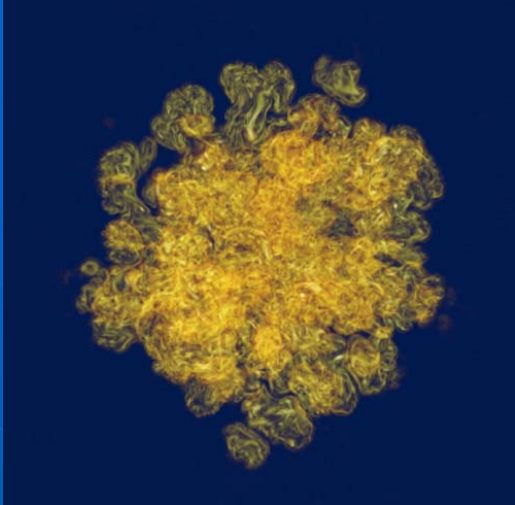


A high-resolution model (“the SNOB run”)

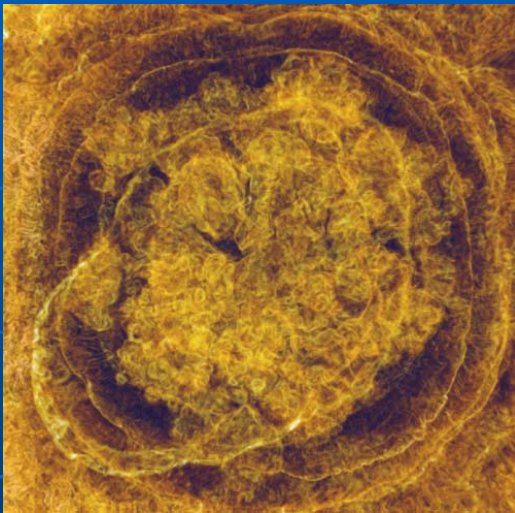
- “ 4π ”
- 1024^3 grid
- initial resolution near the center $\approx 800\text{m}$
- moving grid
- Local & dynamical sgs-model
- ~ 1000 h on 512 processors, IBM/Power4, at RZG



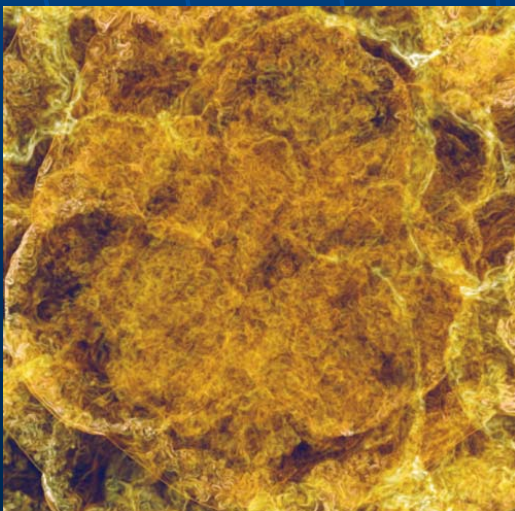
Turbulence?



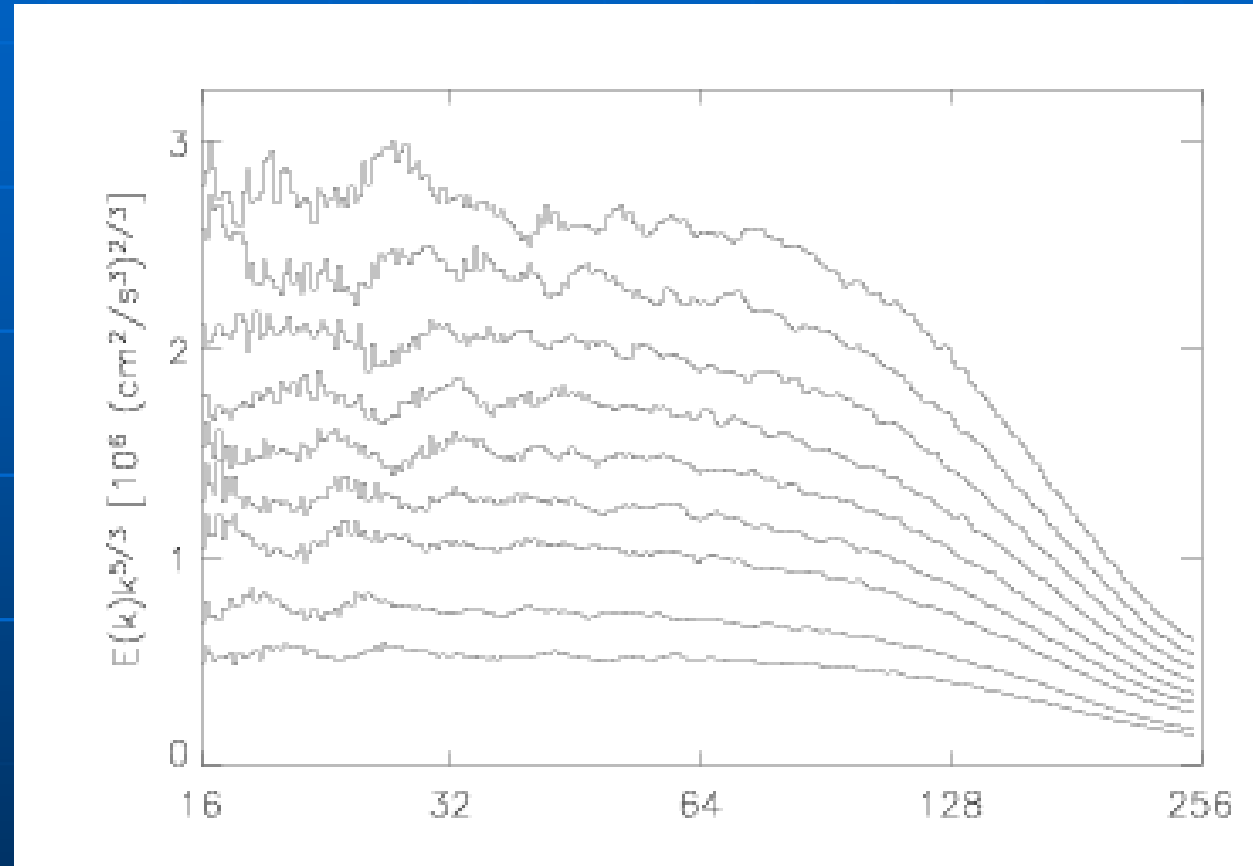
0.25s



0.50s



0.75s



Schmidt et al. (in preparation)

Some (preliminary) results:

- $E_{\text{kin}} = 8.1 \cdot 10^{50}$ erg
- Iron-group nuclei: $0.61 M_{\text{sun}}$ ($\sim 0.41 M_{\text{sun}} {}^{56}\text{Ni}$)
- Intermediate-mass nuclei: $0.43 M_{\text{sun}}$ (from hydro)
- Unburnt C+O: $0.37 M_{\text{sun}}$ (from hydro)
(less than $0.08 M_{\text{sun}}$ at $v < 8000 \text{ km/s}$)
- $V_{\text{max}} \approx 17,000 \text{ km/s}$

Good agreement with observations!

Questions and challenges (to theory)

➤ Ignition conditions:

How do WDs reach M_{Ch} ? Center/off-center ignition?
One/multiple “points”?

➤ Combustion modeling:

Interaction of nuclear flames with turbulence;
“distributed burning”; “active turbulent combustion” ?
Deflagration/detonation transition: Does it happen? Is it
“needed”?

➤ New generation of “full-star” models:

Light curves? Spectra?

➤ Other progenitors:

Mergers? Sub-Chandrasekhars?