

Stellar Death and Supernovae Conference

Kavli Institute for Theoretical Physics, Santa Barbara, August 17–21, 2009

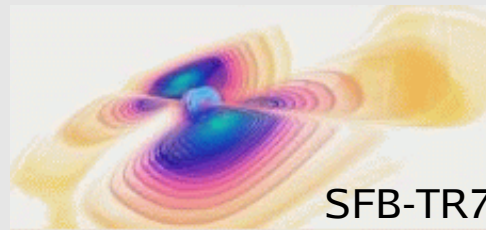
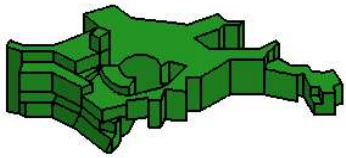
Status of 8-10 Solar Mass Objects and Beyond

Hans-Thomas Janka

(Max Planck Institute for Astrophysics, Garching, Germany)

Student, postdocs, collaborators:

Lorenz Hüpdepohl, Andreas Marek, Bernhard Müller, Nico Hammer, Ewald Müller,
Shinya Wanajo



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Status of Core-Collapse Supernova Modeling in Garching

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(Max Planck Institute for Astrophysics, Garching, Germany)

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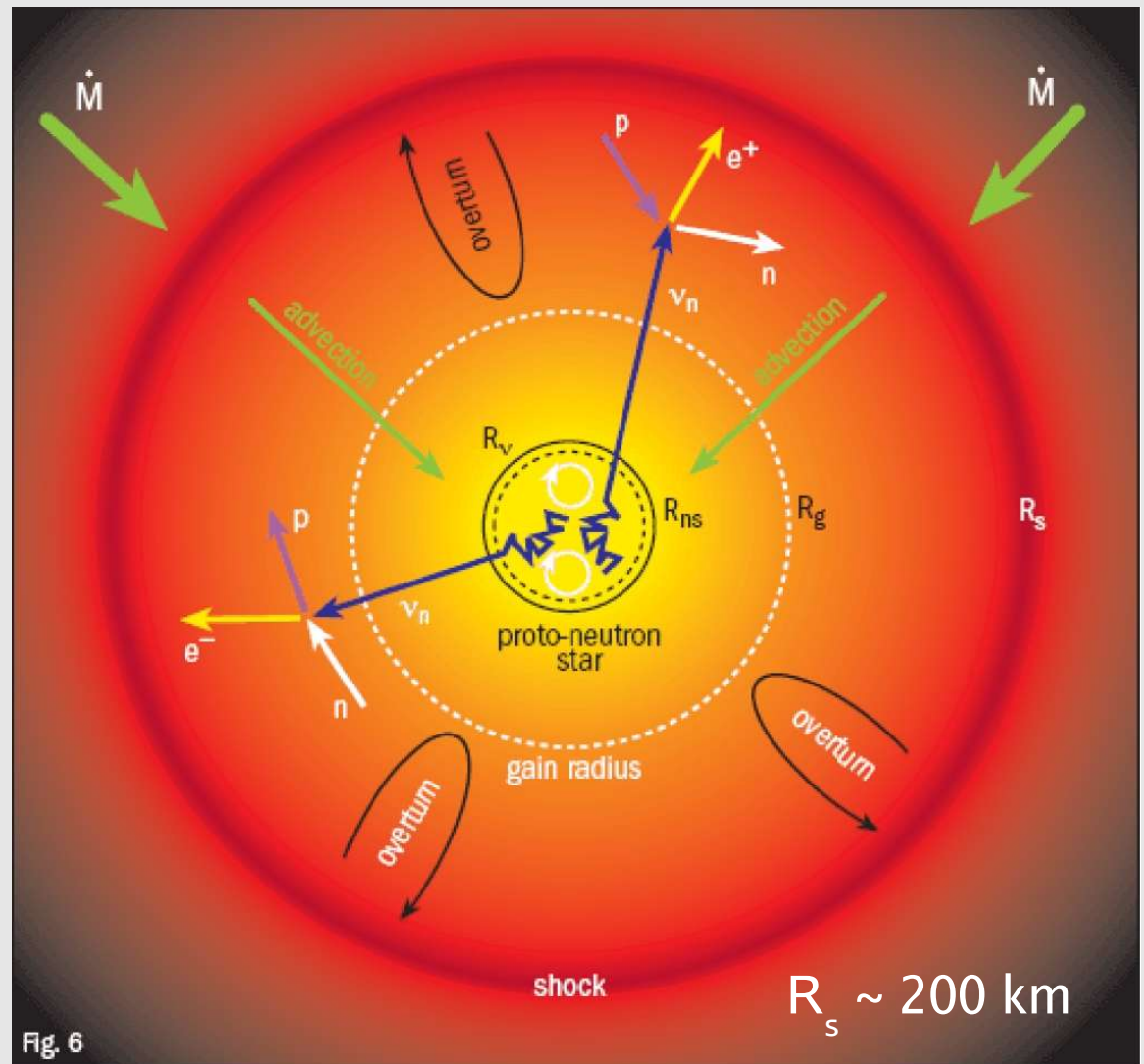
Lorenz Hudepohl, Andreas Marek, Bernhard Müller, Nico Hammer, Ewald Müller,
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Exploration of neutrino-driven explosions

- 3D simulations of **mixing instabilities in supernovae** until shock breakout
- Explosion models for **electron-capture supernovae** of $8-10 M_{\text{sun}}$ stars
- **Neutrino-driven explosion models** of 11 and $15 M_{\text{sun}}$ stars

Neutrinos & Explosion Mechanism

Paradigm: Explosions by the convectively supported neutrino-heating mechanism



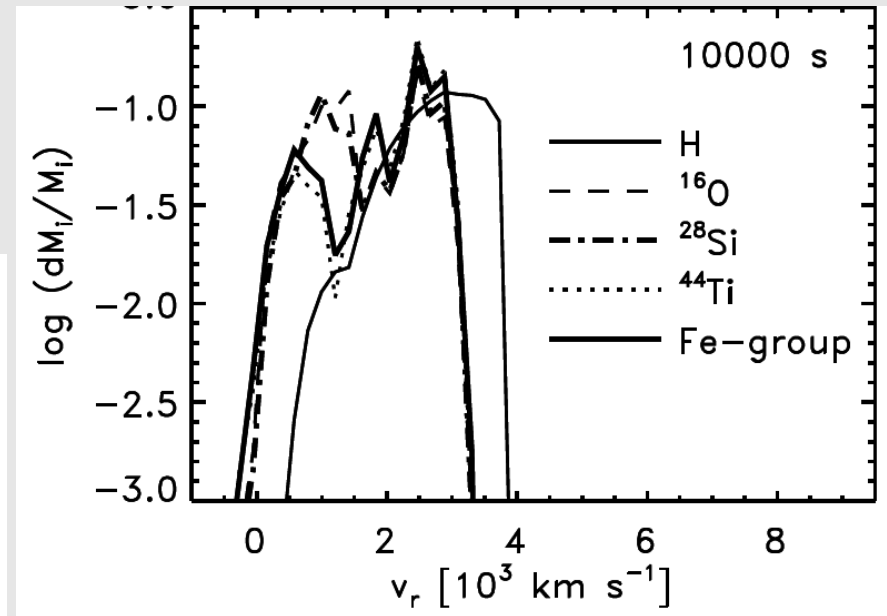
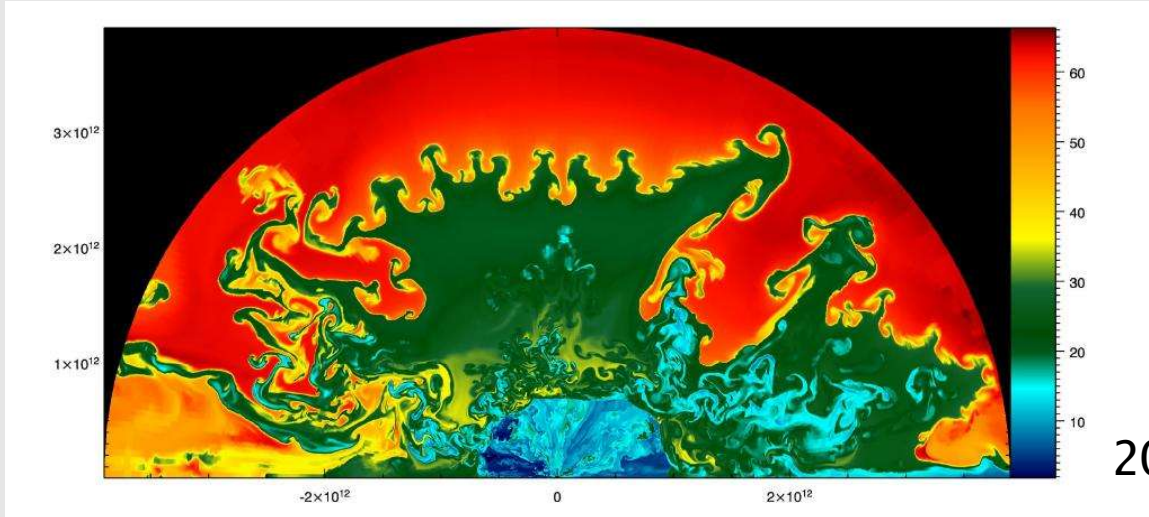
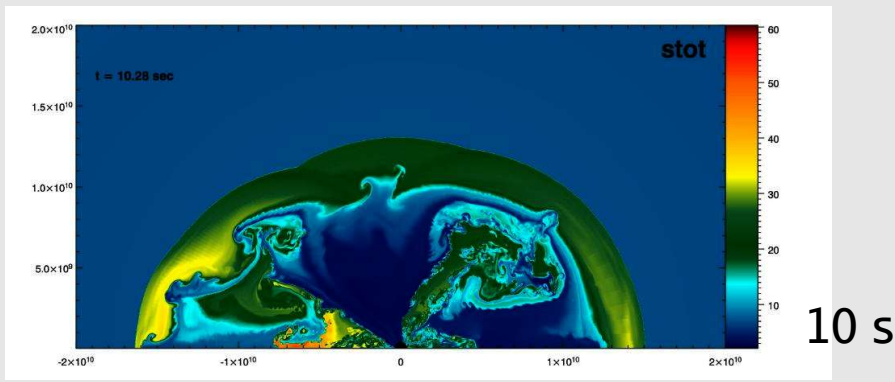
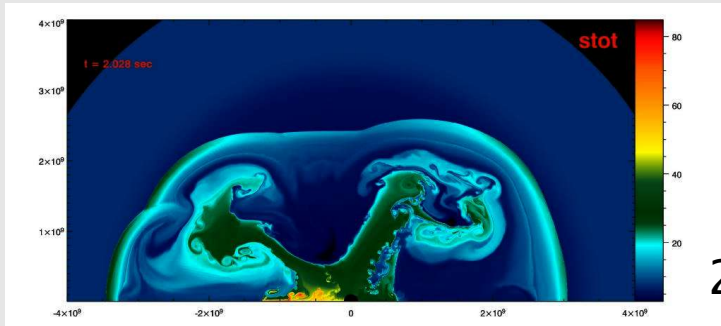
- “Neutrino-heating mechanism”: Neutrinos ‘revive’ stalled shock by energy deposition (Colgate & White 1966, Wilson 1982, Bethe & Wilson 1985);
- Convective processes & hydrodynamic instabilities play an important role (Herant et al. 1992, 1994; Burrows et al. 1995, Janka & Müller 1994, 1996; Fryer & Warren 2002, 2004; Blondin et al. 2003; Scheck et al. 2004,06,08).

3D Models of CCSN Explosions

Mixing Instabilities during SN Explosions

2D models with large initial shock deformation:

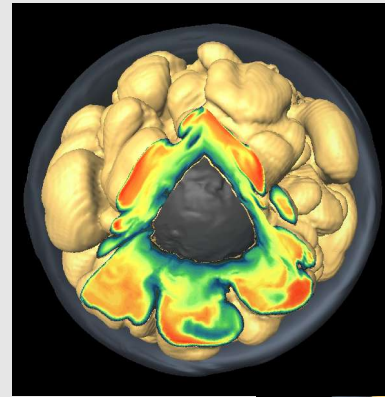
- Strong metal mixing into H envelope
- Strong H mixing deep into He layer
- large asymmetry of metal distribution
- High Ni velocities ($v_{\max} > 3000$ km/s as observed in SN 1987A)



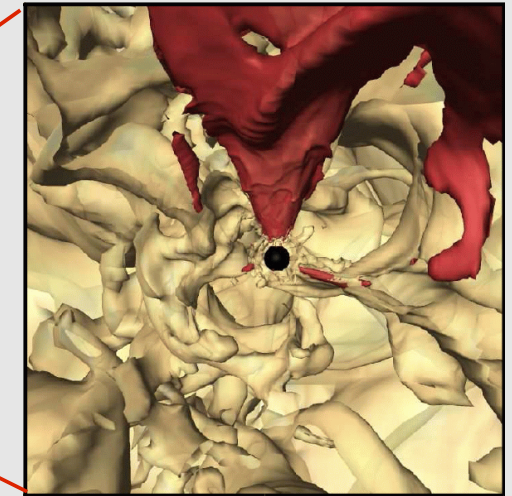
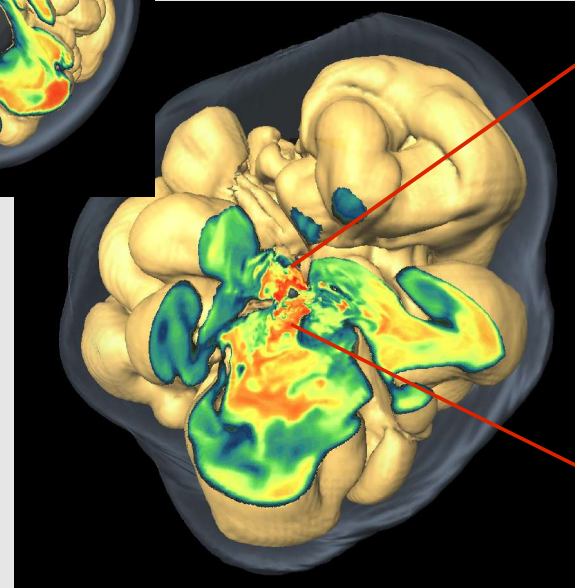
Kifonidis et al., A&A 453 (2005), 661

Parametric Explosion Studies in 3D

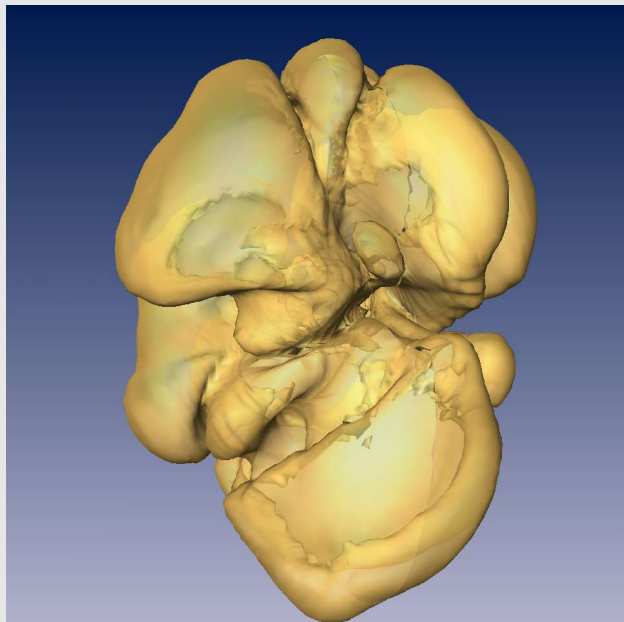
- Explosions in 3D show also very large asymmetries
- Accretion flow to neutron star develops $l = 1$ mode also in 3D
- Should produce neutron star kicks similar to 2D



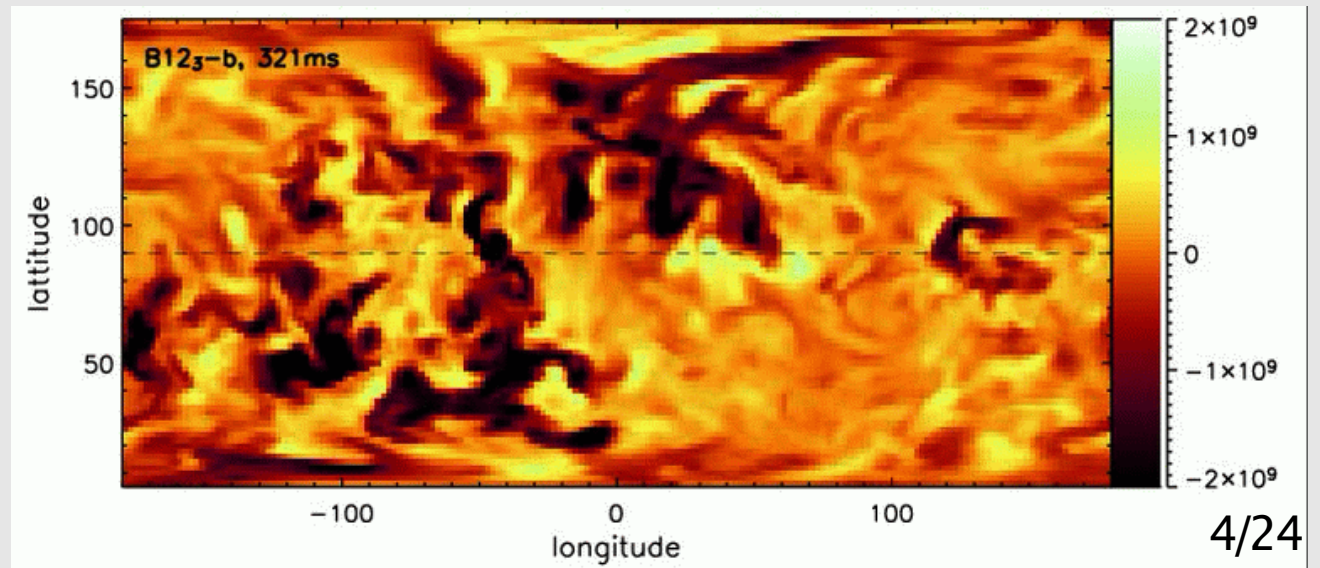
15 M_{sun} star without rotation
(Scheck, PhD Thesis 2006)



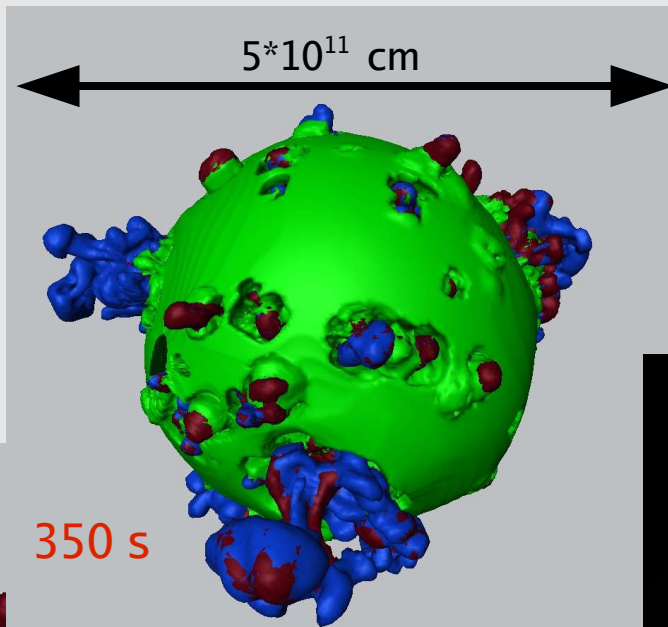
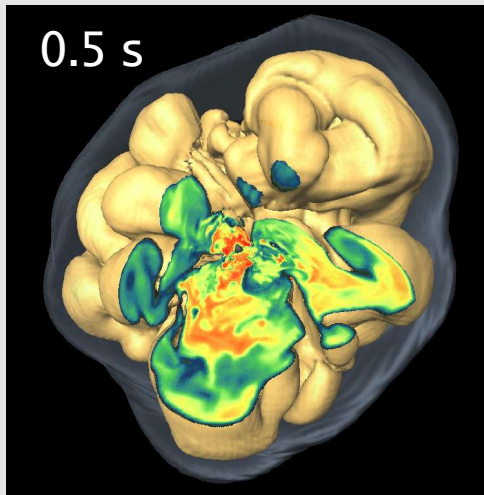
1400 km



3D with rotation (Scheck, PhD Thesis 2006)

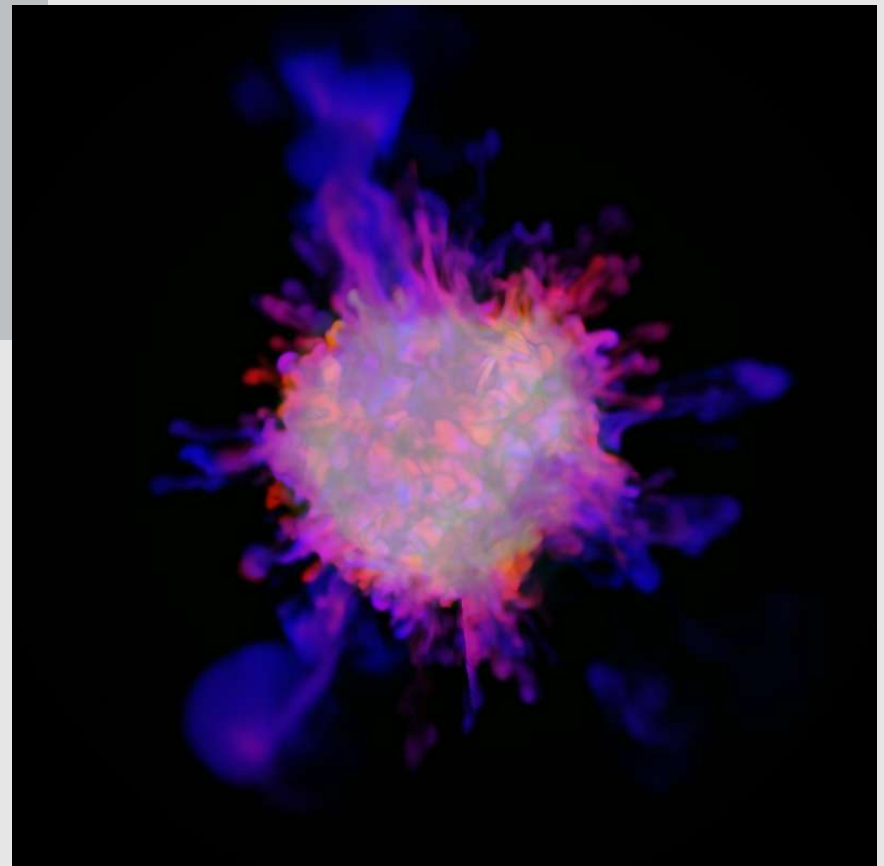
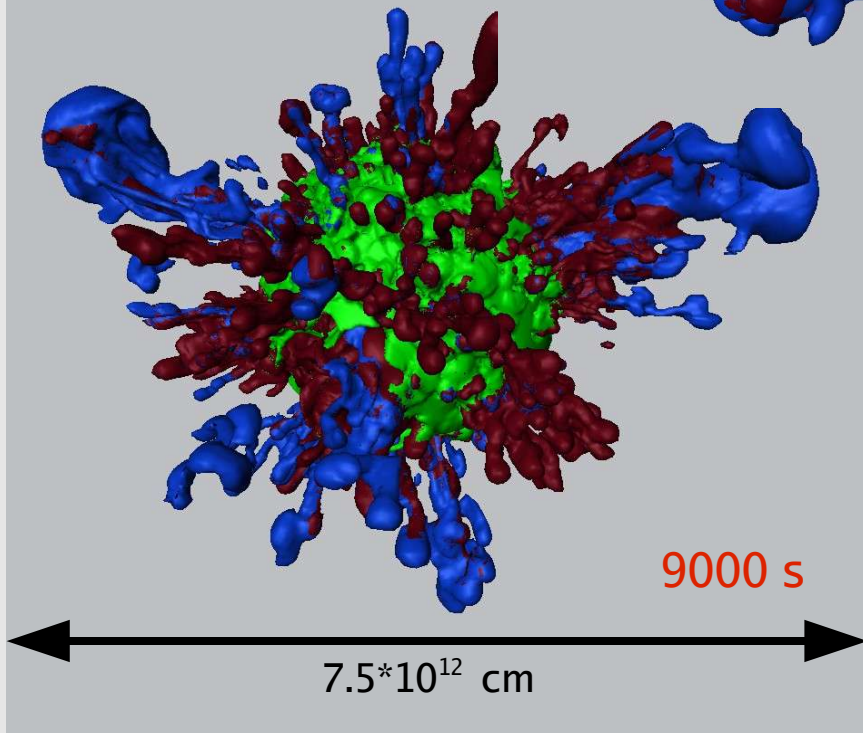


Mixing Instabilities in 3D SN Models



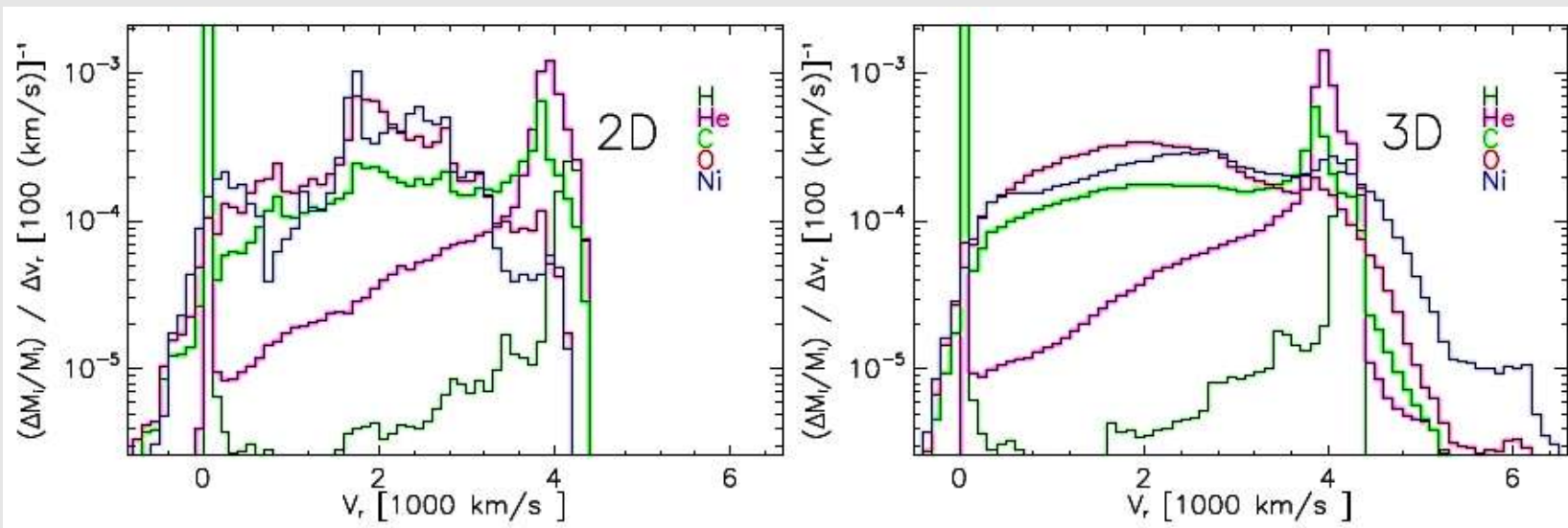
green: carbon
red: oxygen
blue: nickel

350 s

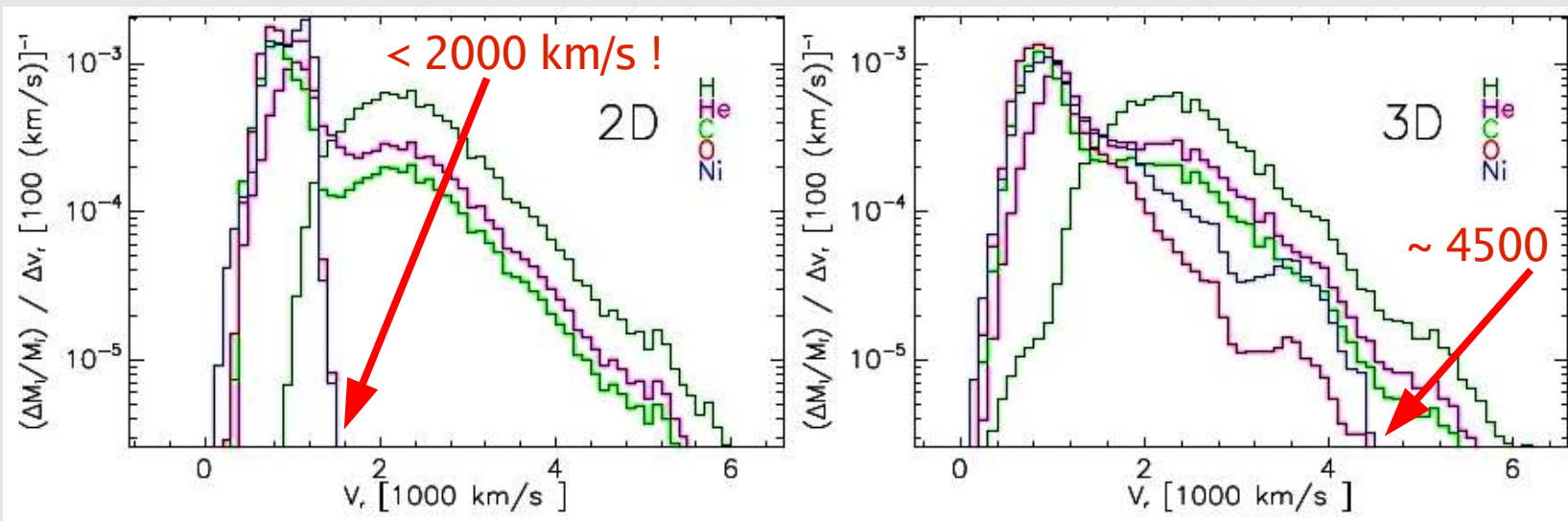


(Hammer, Janka, Müller, submitted)

Mixing Instabilities in 3D SN Models



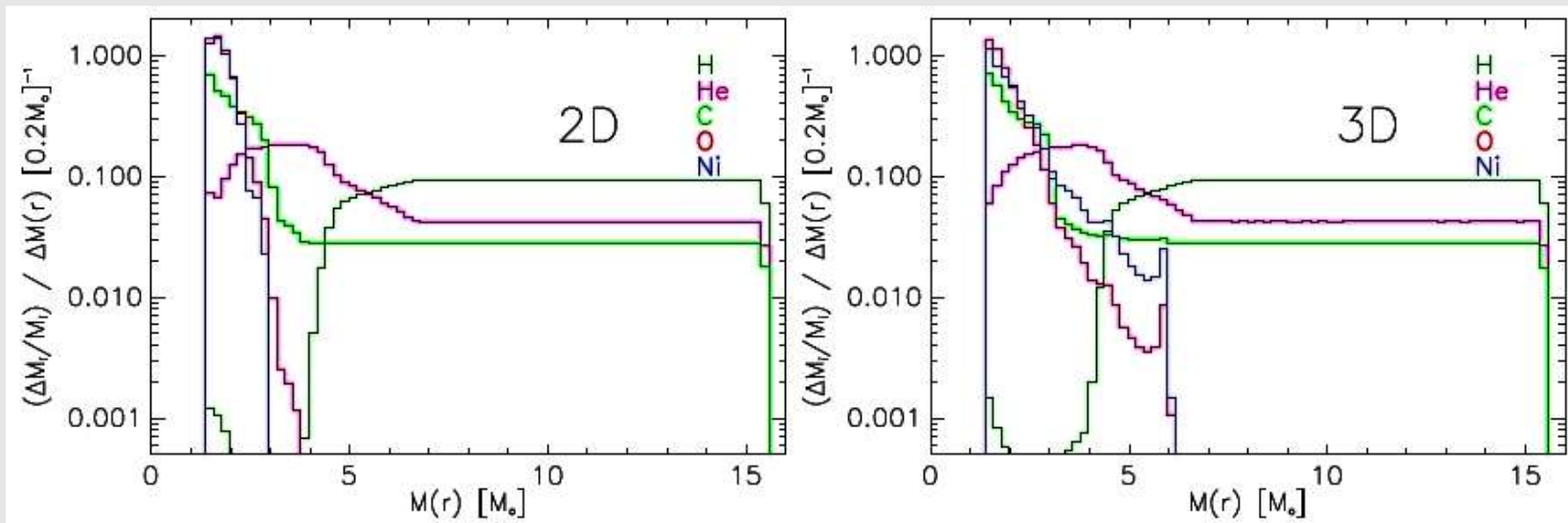
350 s



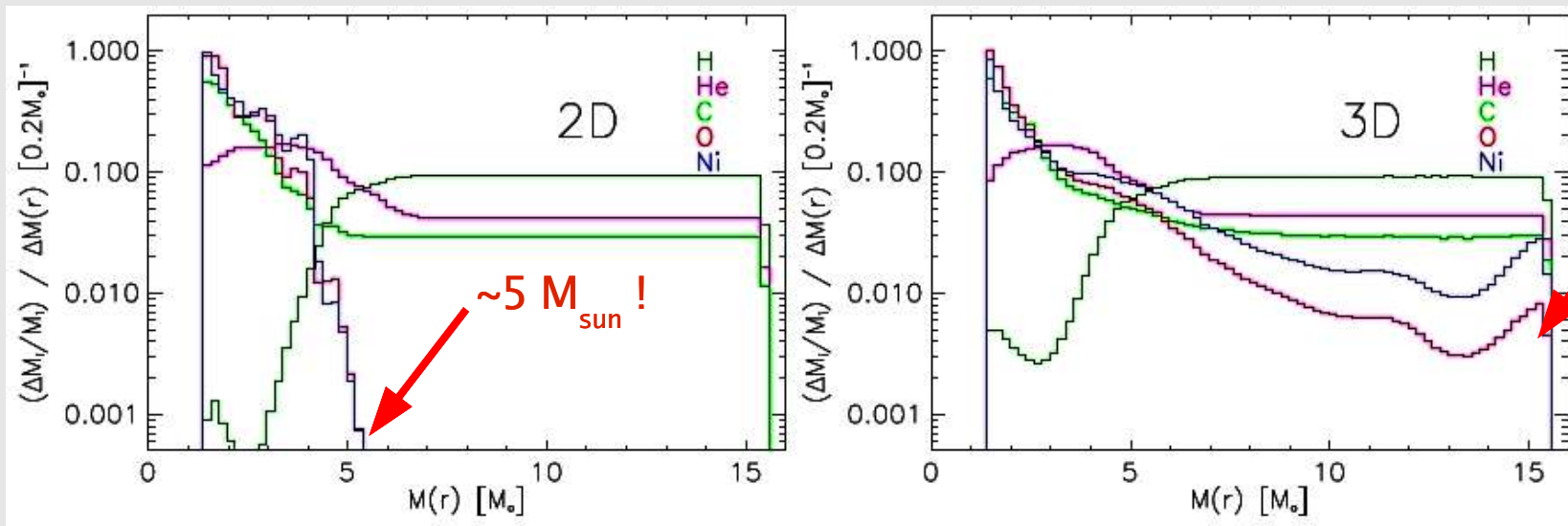
9000 s

- In 3D Rayleigh-Taylor instabilities grow faster \implies clumps with higher initial velocities
- In 3D drag forces are smaller \implies less deceleration of propagating clumps

Mixing Instabilities in 3D SN Models



350 s



~15 M_{sun} !

9000 s

- In 3D mixing much more efficient; very fast metal clumps with up to several $10^{-3} M_{\text{sun}}$

- Onion-shell structure of progenitor is turned over: Fe overtakes O, O faster than C

Explosions of
 $M_{\text{star}} \sim 8-10 M_{\text{sun}}$ Stars

SN Progenitors: Core density profiles

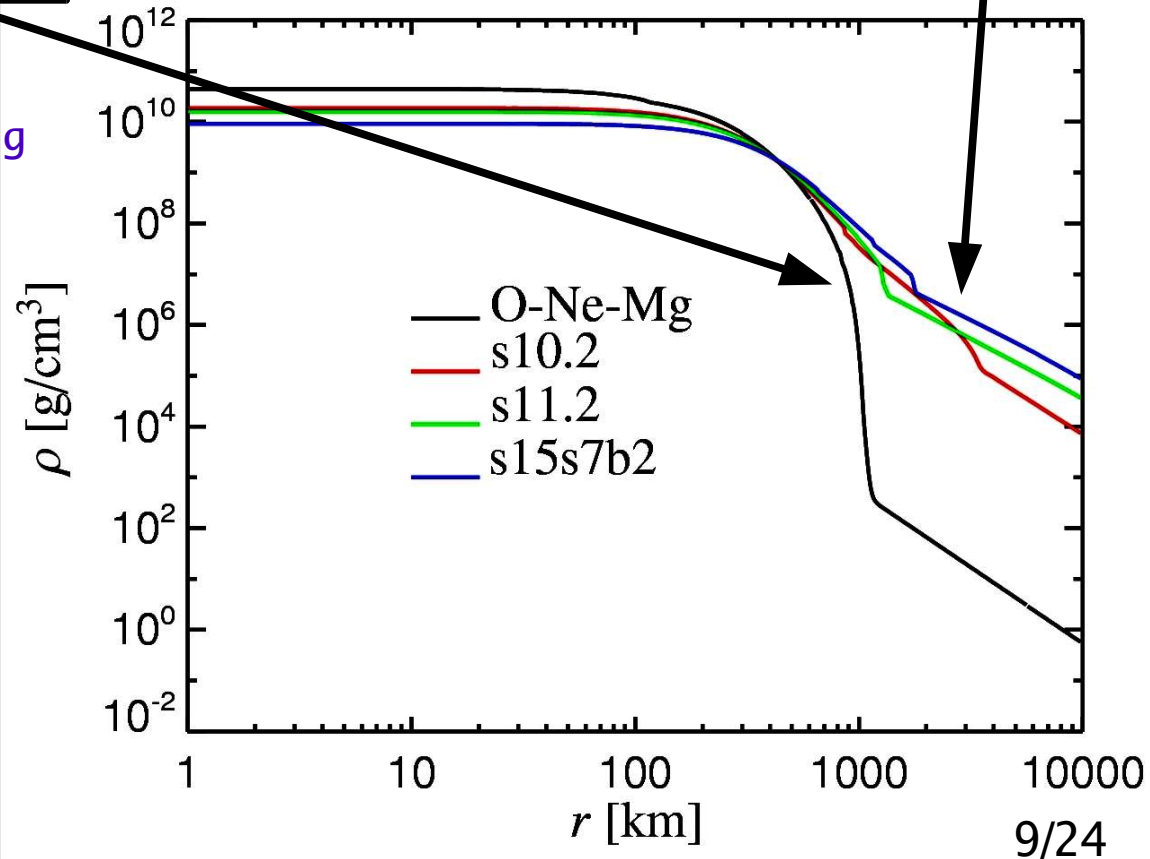
~8–10 M_{sun} (super-AGB) stars have ONeMg cores with a very steep density gradient at the surface
 (====> rapidly decreasing mass accretion rate after core bounce)

>10 M_{sun} stars have much higher densities outside of their Fe cores
 (e.g. Heger et al., Limongi et al., Nomoto et al., Hirschi et al.)
 (====> ram pressure of accreted mass decreases slowly after core bounce)

8.8 M_{sun} progenitor model (Nomoto 1984):
 2.2 M_{sun} H+He, 1.38 M_{sun} C+O, 1.28 M_{sun} ONeMg

~30% of all SNe (Nomoto et al. 1981, 84, 87)

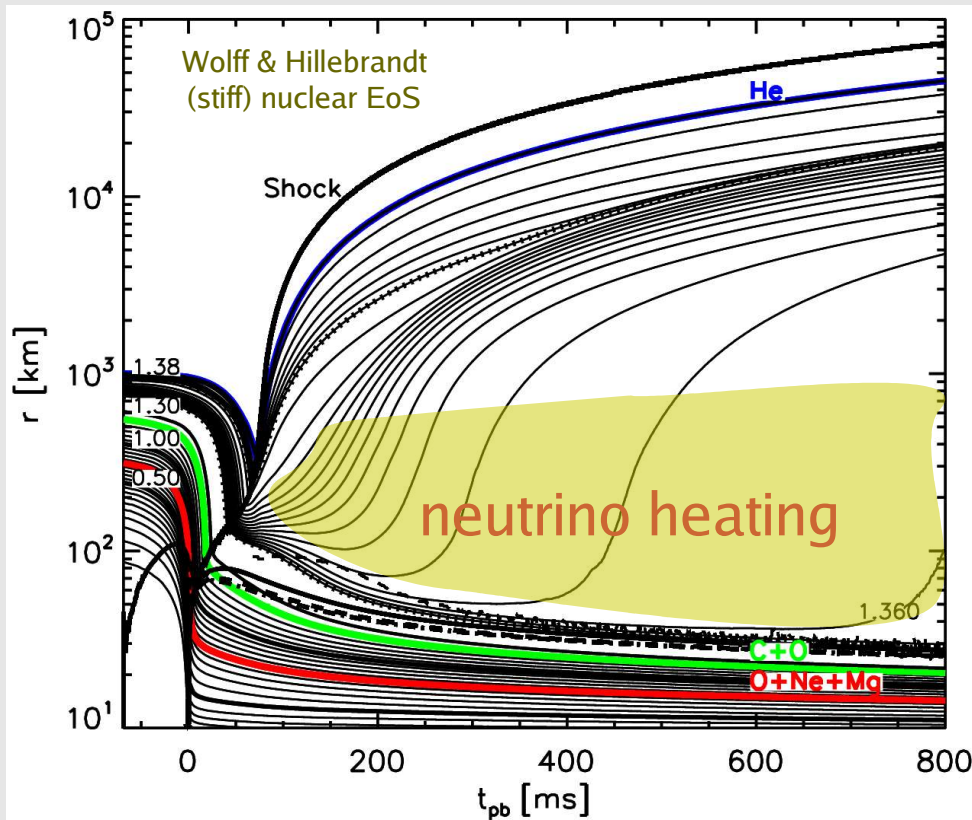
(8.75 $M_{\text{sun}} < M_{\text{ZAMS}} < 9.25 M_{\text{sun}}$: < 20% of all SNe; Poelarends et al., A&A 2006)



SN Simulations:

$M_{\text{star}} \sim 8..10 M_{\text{sun}}$

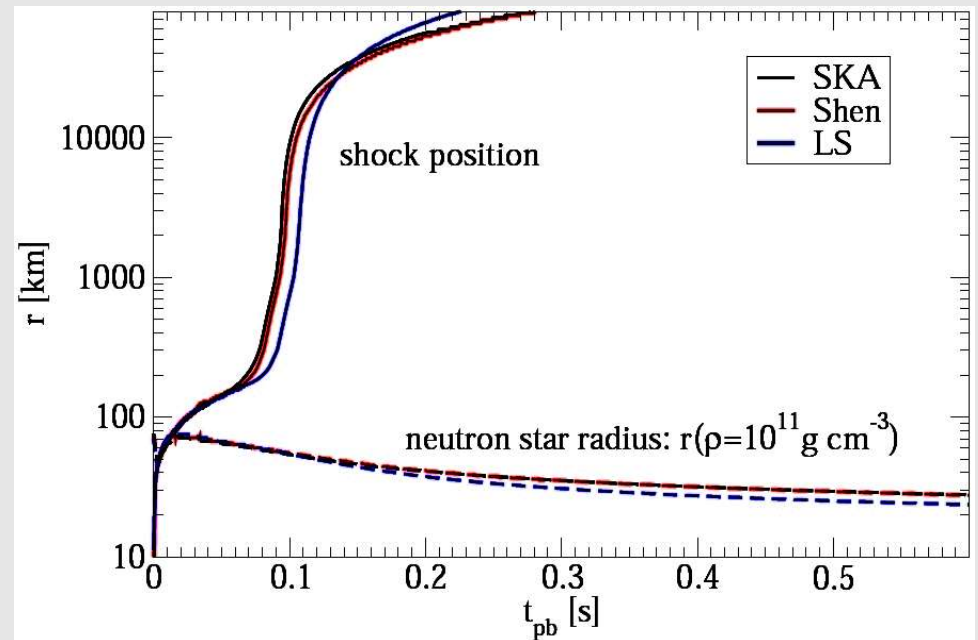
"Electron-capture supernovae"
or "ONeMg core supernovae"



Kitaura et al., A&A 450 (2006) 345;
Janka et al., A&A 485 (2008) 199

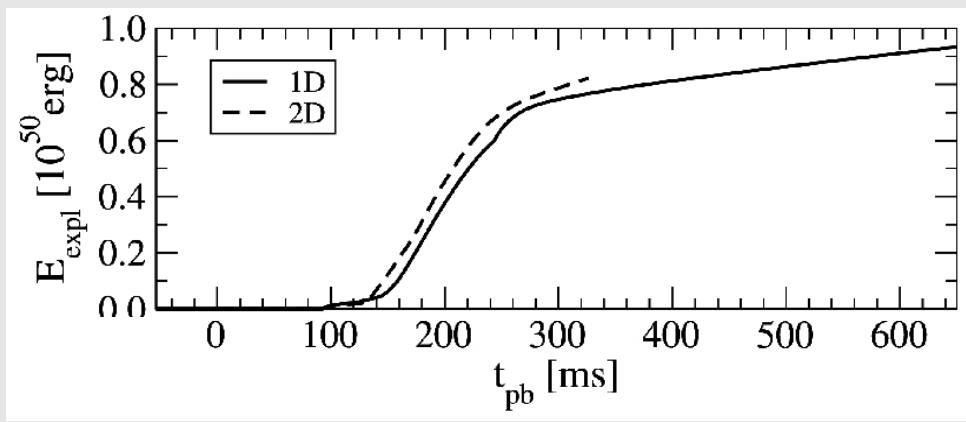
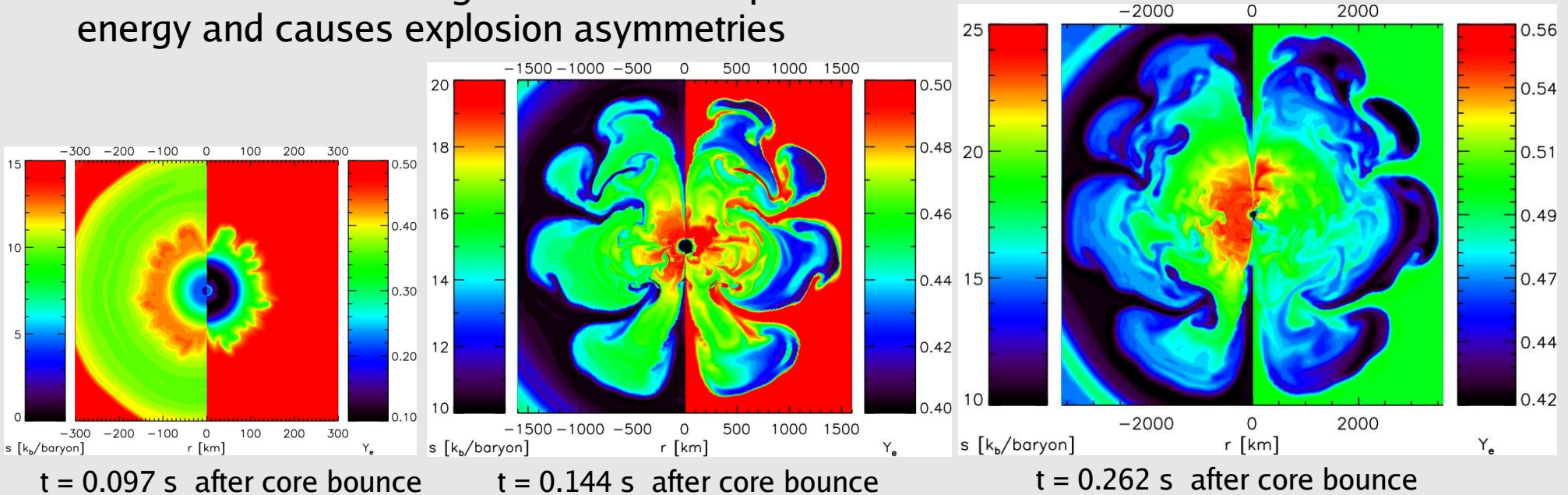
Convection is not necessary for launching explosion
but occurs in NS and in neutrino-heating layer

- **No prompt explosion !**
- Mass ejection by "neutrino-driven wind"
(like Mayle & Wilson 1988
and similar to AIC of WDs;
see Woosley & Baron 1992, Fryer et al. 1999;
Dessart et al. 2006)
- Explosion develops in similar way for
soft nuclear EoS (i.e. compact PNS)
and stiff EoS (less compact PNS)



SN Simulations: $M_{\text{star}} \sim 8..10 M_{\text{sun}}$

Convection leads to slight increase of explosion energy and causes explosion asymmetries

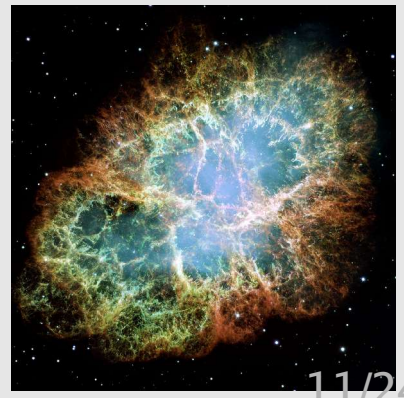


Müller et al. (in preparation)

Low explosion energy and ejecta composition – little Ni, C, O – of CRAB (SN1054) is compatible with ONeMg core explosion

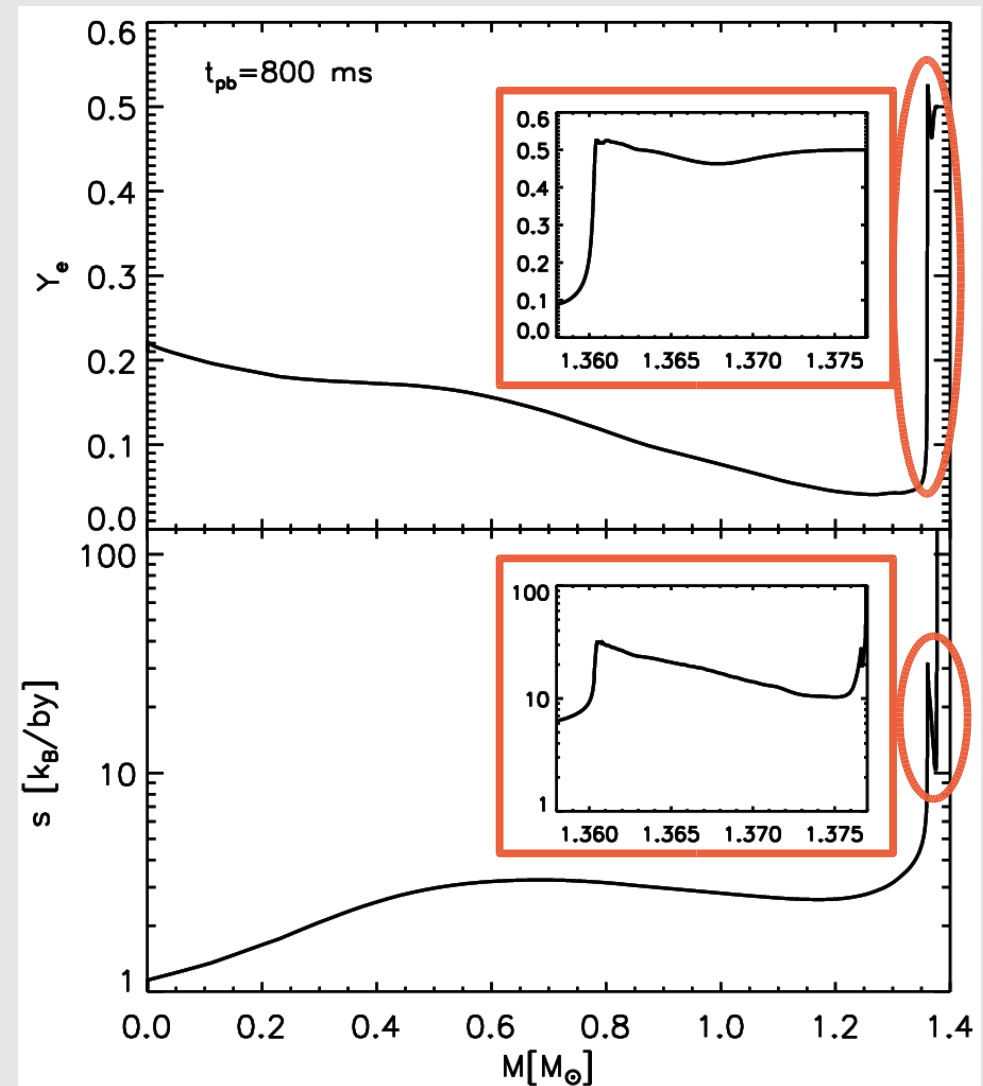
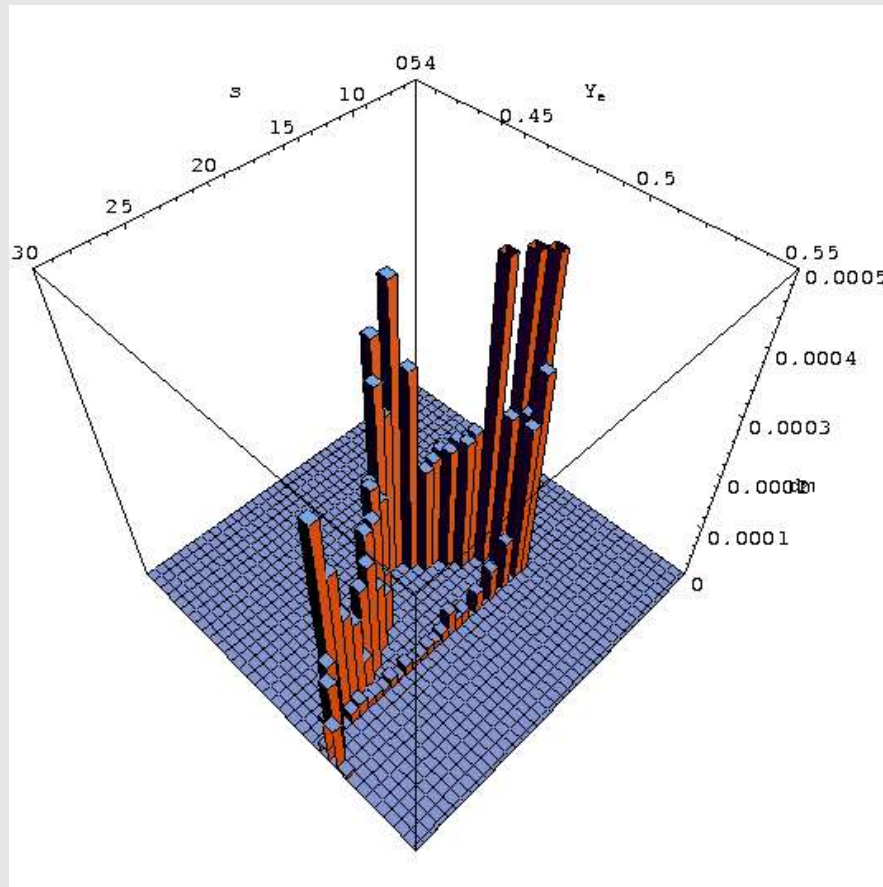
(Nomoto et al., Nature, 1982; Hillebrandt, A&A, 1982)

Might also explain other low-luminosity supernovae (e.g. SN1997D, 2008S, 2008HA)



O-Ne-Mg Core Supernovae: Ejecta

- Early SN ejecta have Y_e around 0.5 and even $Y_e > 0.5$
- Entropies are $\sim 10 k_B$ per nucleon, but **strong increase for ejected matter in steep gradient at core surface**



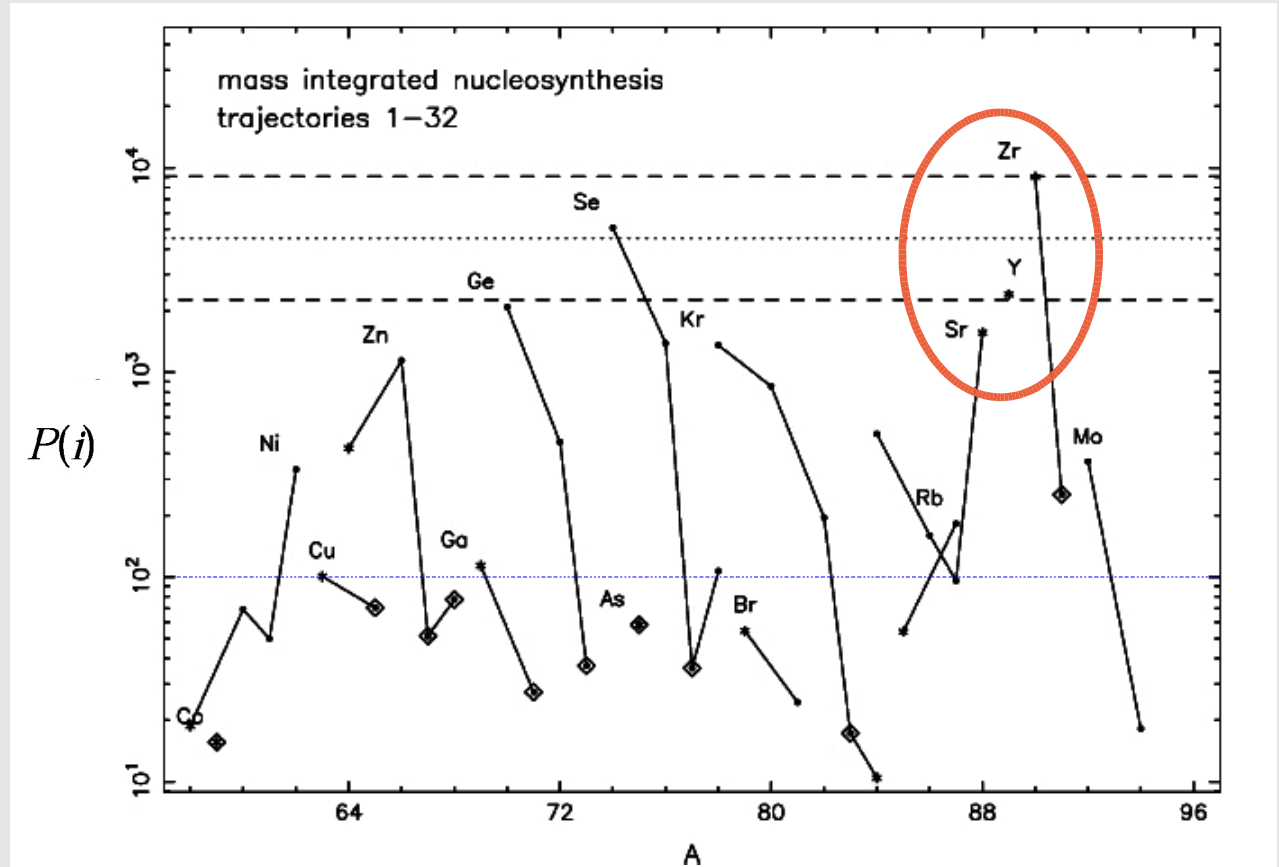
Nucleosynthesis in O-Ne-Mg Core SNe

- **No r-processing in the early ejecta**
(in contradiction to suggestion by Ning et al. 2007)

- Mass cut: $M_{\text{cut}} = 1.363 M_{\text{sun}}$
- Ejecta mass: $M_{\text{ej}} = 1.263 M_{\text{sun}}$
- Nucleosynthesis mass yield:
 $M_{\text{nucsyn}} \sim 1.5 \cdot 10^{-2} M_{\text{sun}}$
- Mass-weighted production factors of nuclides i in all ejecta shells j , normalized to total amount of ejecta:

$$P(i) = \sum_j \frac{M_j}{M^{\text{ej}}} \frac{X_j(i)}{X_{\odot,i}}$$

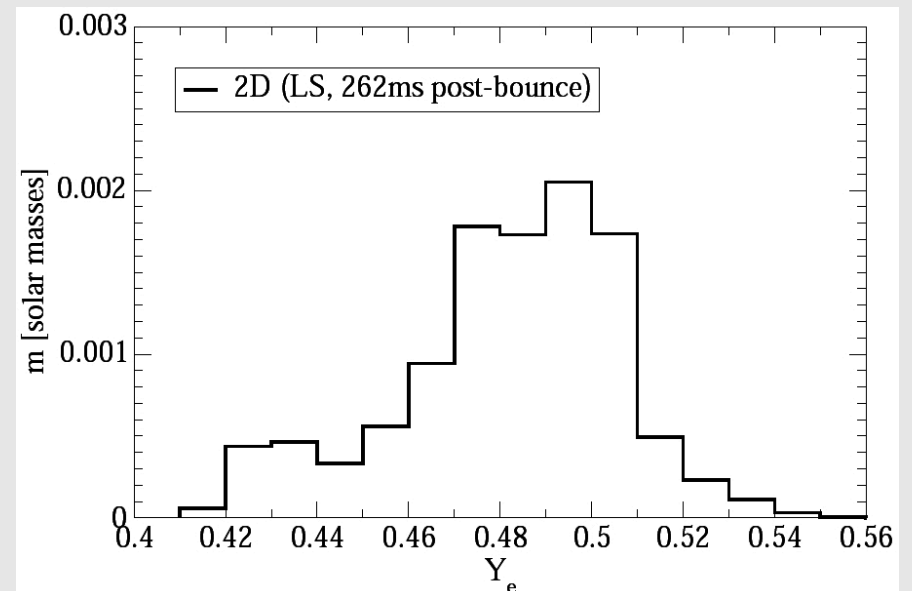
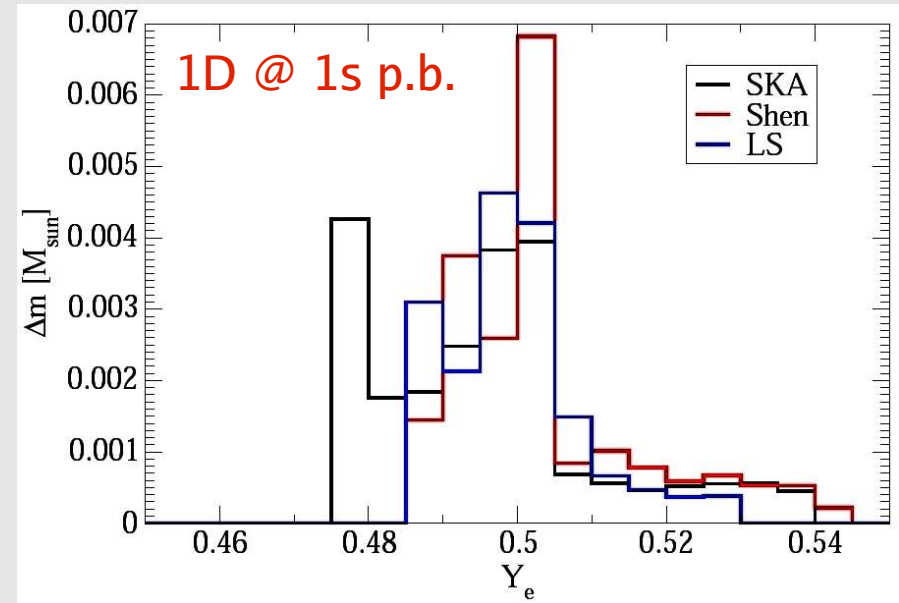
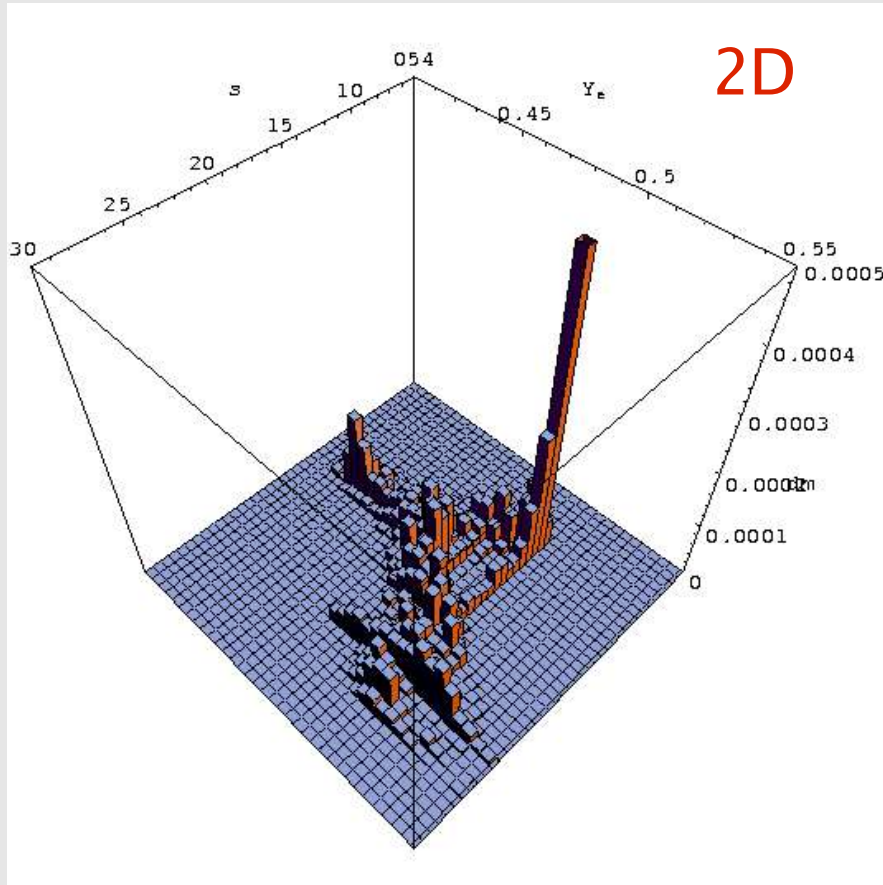
Hoffman, Janka, Müller, ApJL (2008), similar results by Wanajo et al. (2009)



- * To be compatible with Galactic chemical abundances:
 $P(i) < 100$ (if ~10% of all SNe are O-Ne-Mg core collapses).
- * **Massive overproduction** (10–50 times solar value over history of Galaxy) of $N=50$ closed neutron-shell nuclei ^{88}Sr , ^{89}Y , ^{90}Zr , because $\sim 5.5 \cdot 10^{-3} M_{\text{sun}}$ are ejected with $Y_e < 0.47$ and low entropies ($s \sim 20 k_B/\text{nucleon}$).

Nucleosynthesis in O-Ne-Mg Core SNe

- ~2% increase of minimum Y_e removes overproduction problem (Wanajo et al. 2009)
- Slightly improved neutrino treatment indeed causes required Y_e increase!
- **But: 2D effects go in opposite direction!**

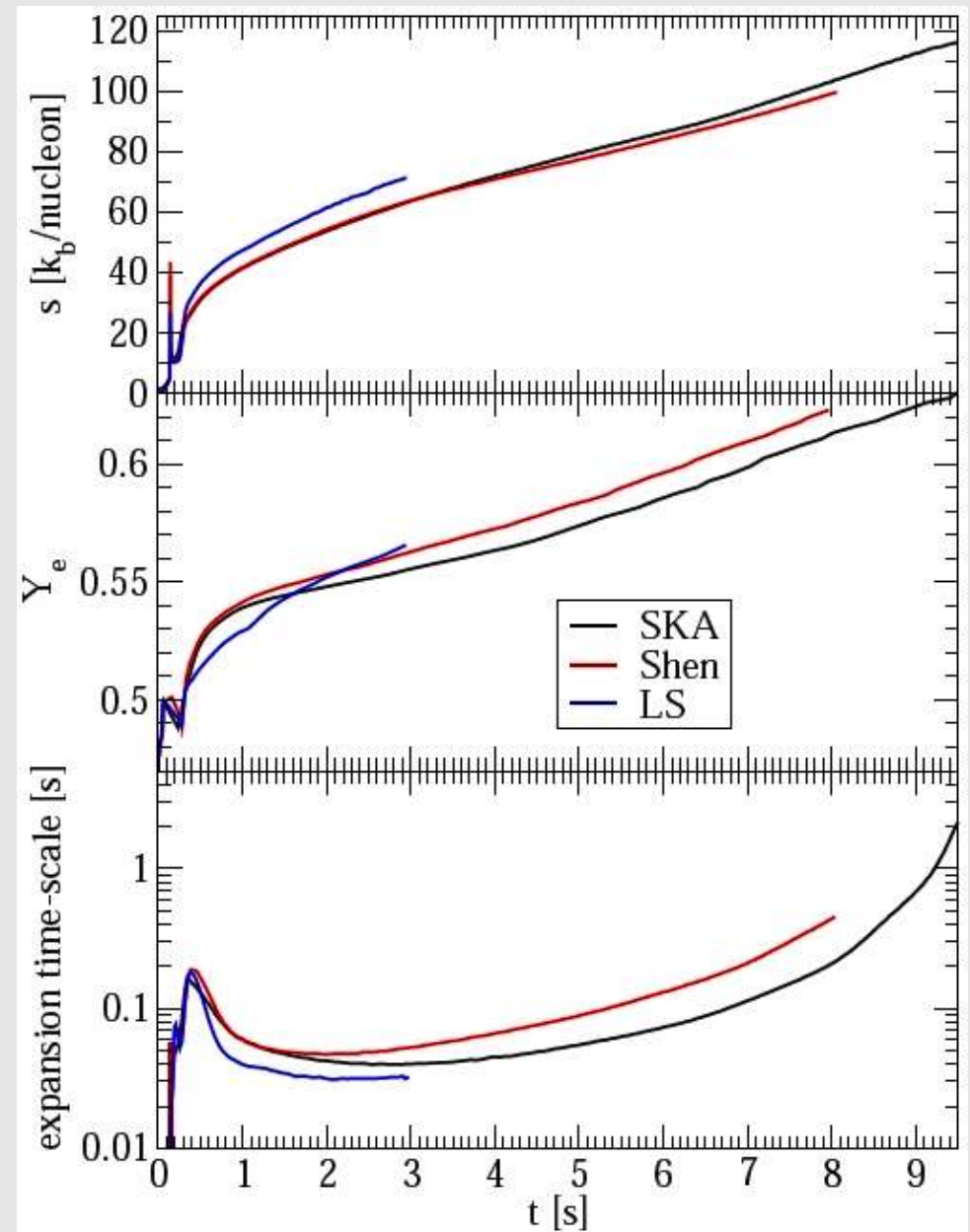
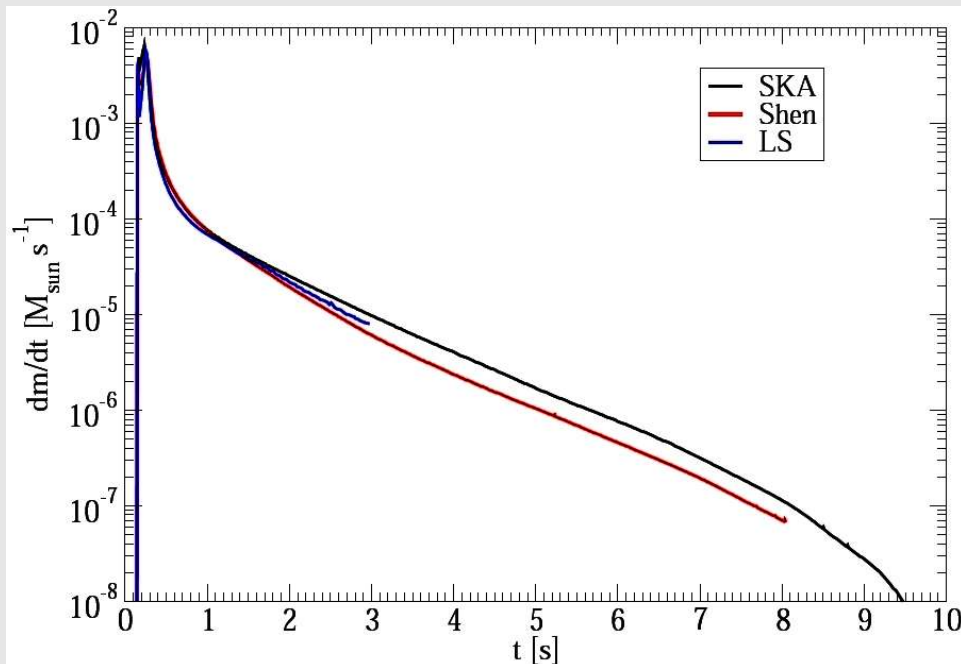


(Müller & Janka, in preparation)

Nucleosynthesis in O-Ne-Mg Core SNe

- Neutrino-driven wind remains p-rich for >10 seconds!
- No r-process in the late neutrino-driven wind!

Hüdepohl (Diploma Thesis 2009)



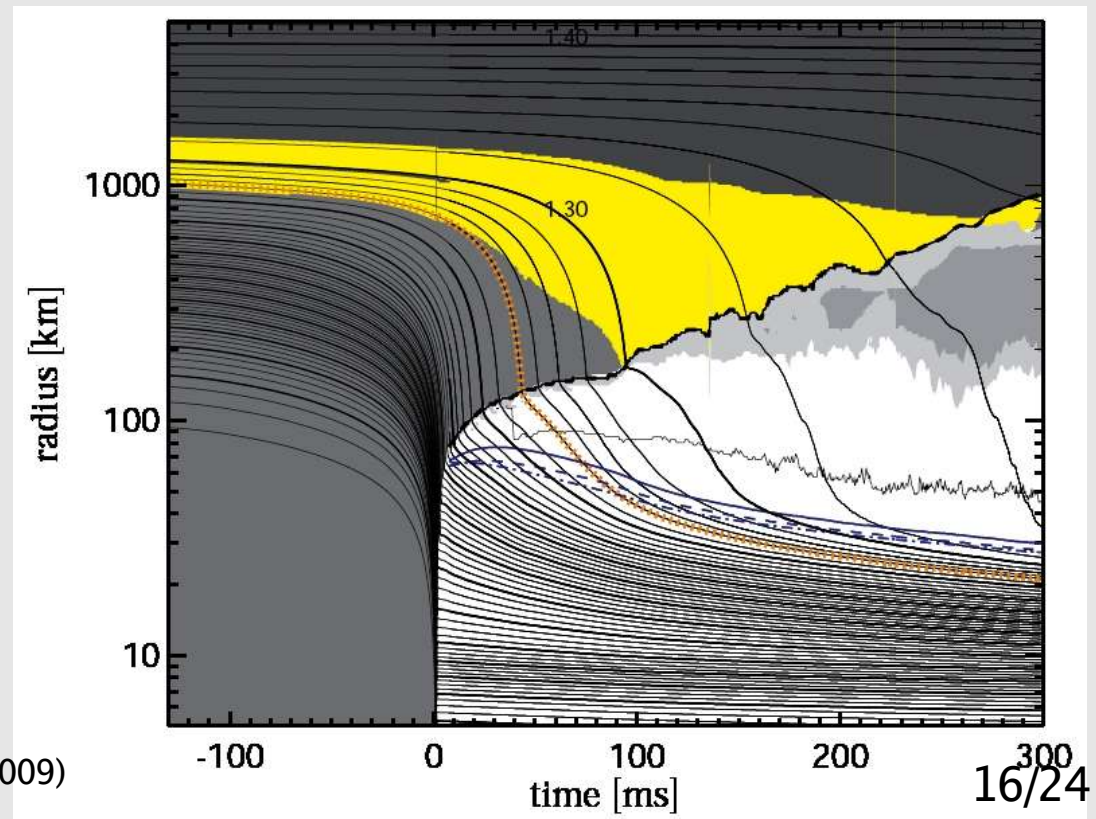
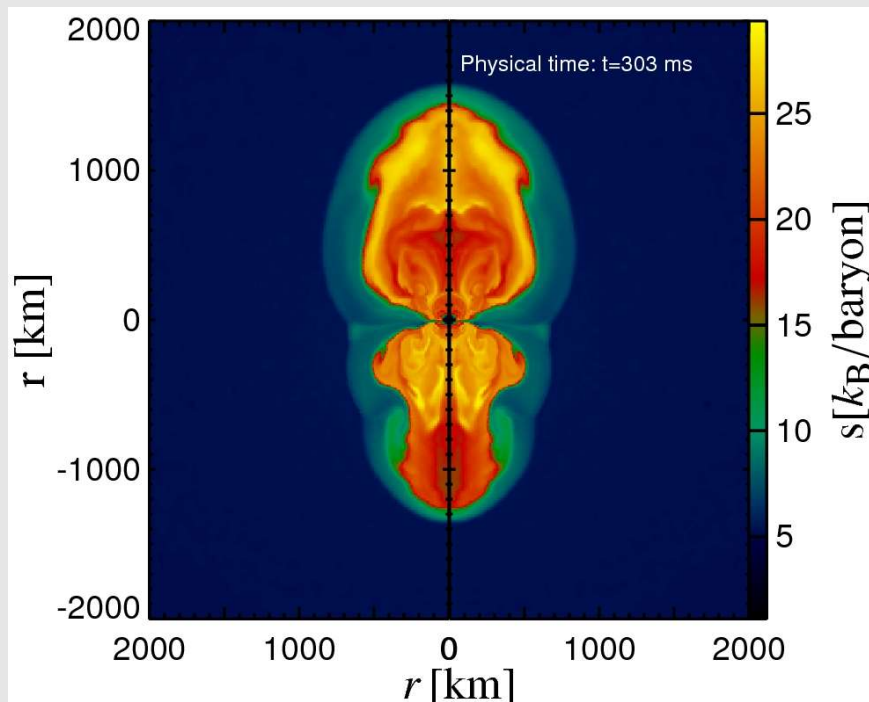
Explosions of
 $M_{\text{star}} > 10 M_{\text{sun}}$ Stars

2D SN Simulations: $M_{\text{star}} \sim 11 M_{\text{sun}}$

For explosions of stars with $M > 10 M_{\text{sun}}$ multi-dimensional effects (nonradial hydrodynamic instabilities) are crucial !

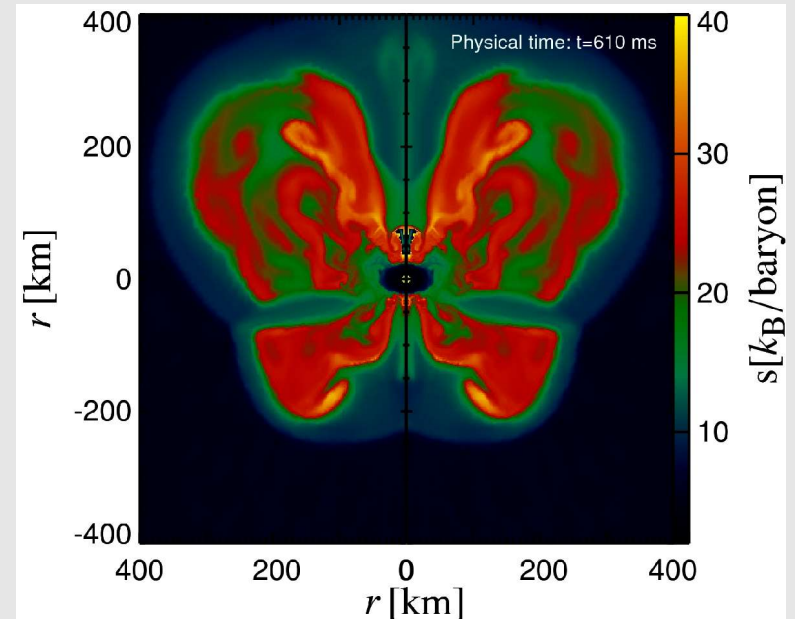
Low-mode nonradial (dipole, $l=1$, and quadrupole, $l=2$) "standing accretion shock instability" ("SASI"; Blondin et al. 2003) develops and pushes shock to larger radii

====> This stretches residency time of matter in neutrino heating layer and thus increases neutrino energy deposition;
Initiation of globally aspherical explosion by neutrino heating even without rotation

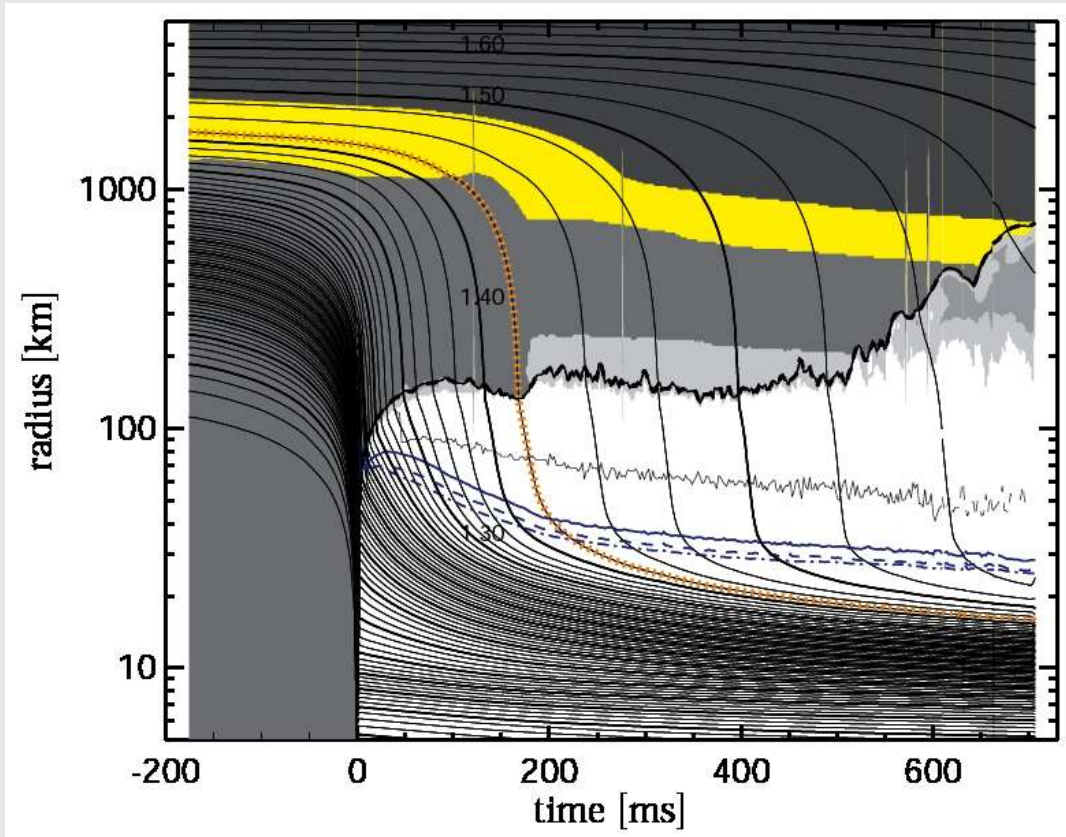


2D SN Simulations: $M_{\text{star}} = 15 M_{\text{sun}}$

Violent SASI oscillations,
 ν -driven explosion sets in
at $t \sim 600$ ms after bounce



(Marek, PhD Thesis 2007;
Marek & THJ, ApJ, 2009)



Explosion Energies and NS masses

$$E_{\text{exp}} \approx \dot{E}_{\nu} \tau_{\text{acc}} + E_{\text{wind}} + E_{\text{burn}} - E_{\text{bind}}$$

$$\dot{E}_{\nu} \sim \zeta \dot{M}_{\text{acc}} \dot{q}_{\nu} \tau_{\text{adv}}$$

$$\sim 2 \times 10^{51} \frac{\text{erg}}{\text{s}} \left(\frac{\zeta}{0.5} \right) \left(\frac{\dot{M}_{\text{acc}}}{0.2 M_{\odot}/\text{s}} \right) \times$$

$$\times \left(\frac{\dot{q}_{\nu} m_{\text{B}}}{300 \text{ MeV}/\text{s}} \right) \left(\frac{\tau_{\text{adv}}}{30 \text{ ms}} \right)$$

(E_{exp} depends on the duration of simultaneous accretion & outflow after onset of explosion: $t_{\text{acc}} \sim 0.5 \text{ sec}$)

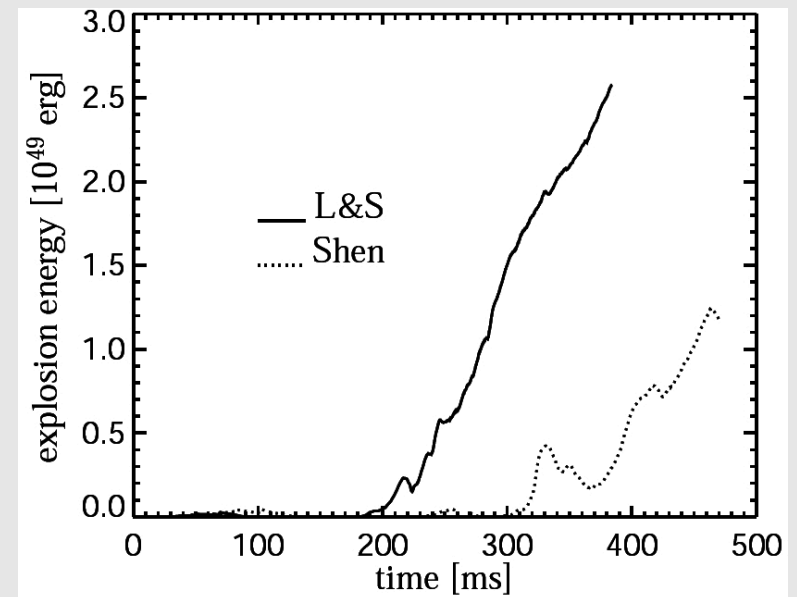
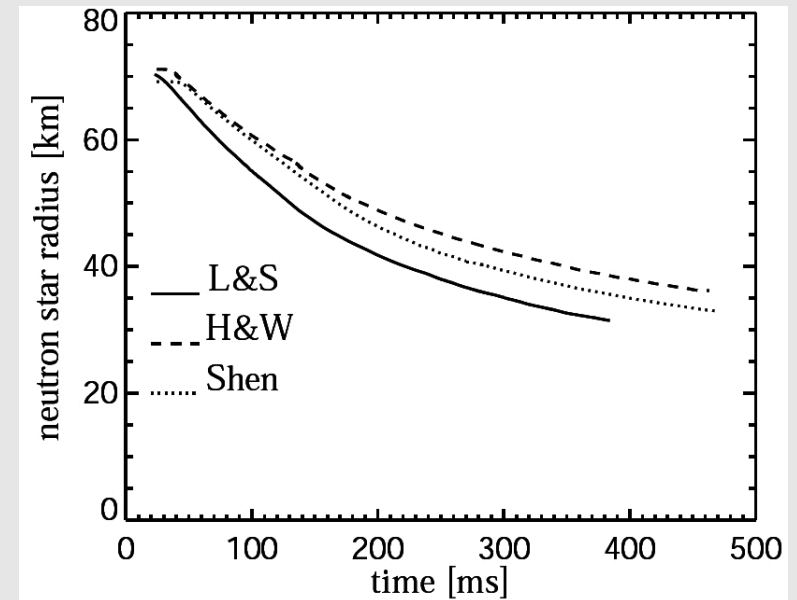
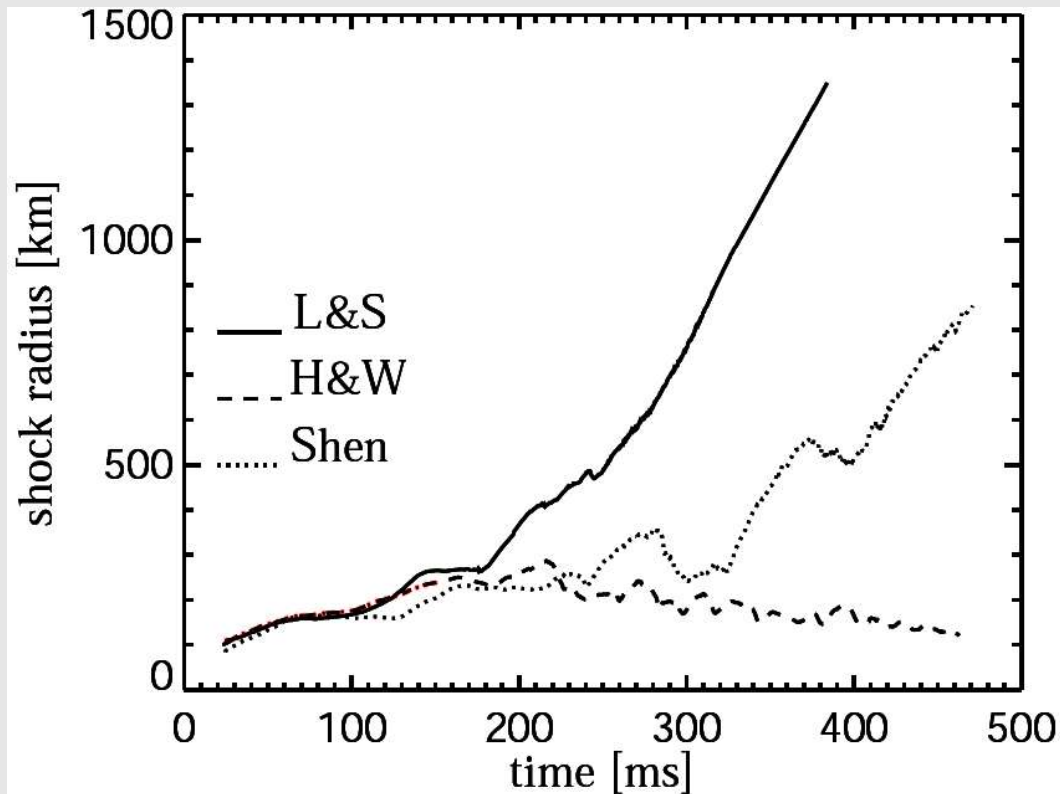
$$\tau_{\text{acc}} \approx \frac{R_{\text{esc}}}{v_{\text{s}}} \sim 0.5 \text{ s } M_{1.5} v_{\text{s},9}^{-3}$$

Stellar mass [M_{sun}]	t_{exp} [ms]	ΔM_{gain} [M_{sun}]	E_{exp} [B]	M_{ns} (baryonic) [M_{sun}]
8 – 10	150	< 0.01	0.1 – 0.2	1.35
~11	250	0.01	0.2 – 0.4	1.30
15	620	0.08	~ 1.0	1.55

NOTE: The stellar properties do not vary monotonically with the progenitor mass (cf. Woosley, Heger, & Weaver 2005)

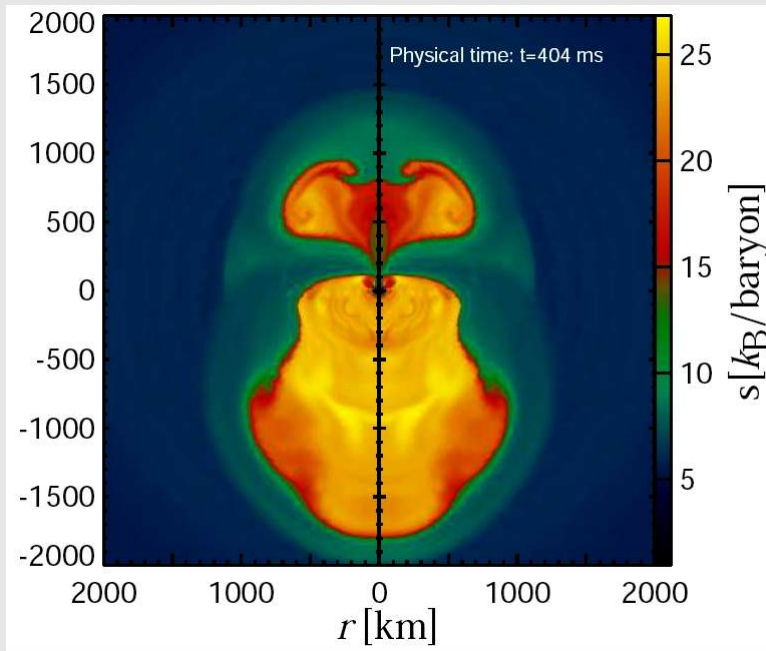
2D Explosions of $11.2 M_{\text{sun}}$ star : Test of EoS Influence

- **Simulations for 3 different nuclear EoSs:** Lattimer & Swesty (L&S), Hillebrandt & Wolff (H&W), Shen et al.
- “Softer” (L&S) EoS and thus more compact PNS leads to earlier explosion

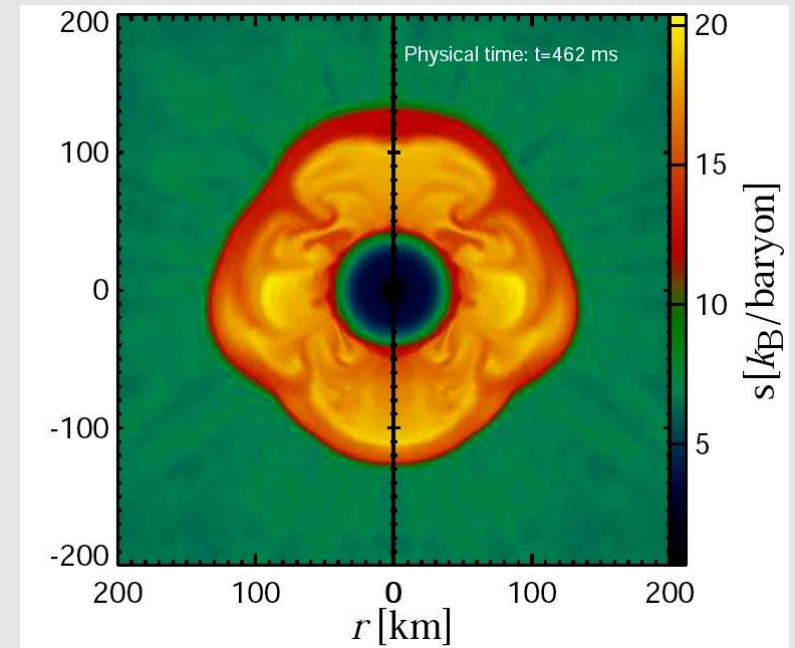


(Marek & THJ, 2009, in preparation)

2D Explosions of $11.2 M_{\text{sun}}$ star : Test of EoS Influence

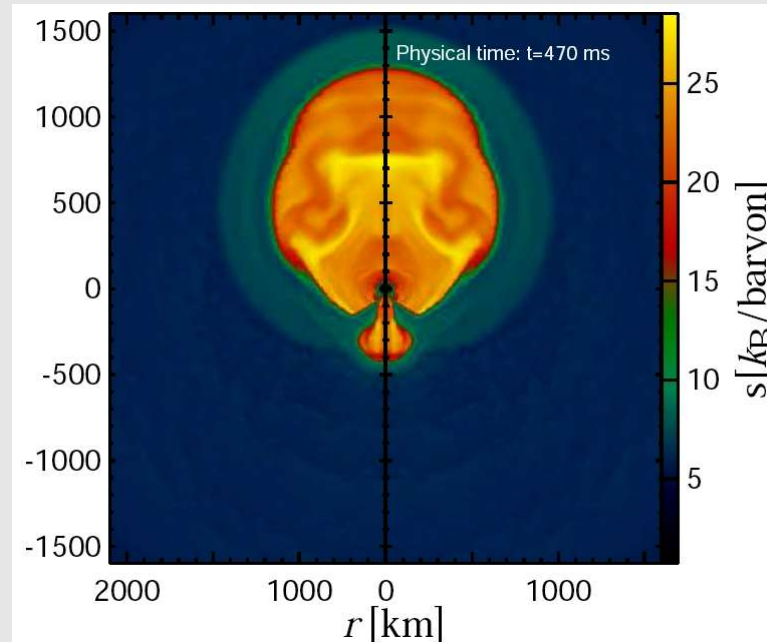


L&S EoS, $t \sim 400$ ms p.b.



H&W EoS, $t \sim 460$ ms p.b.

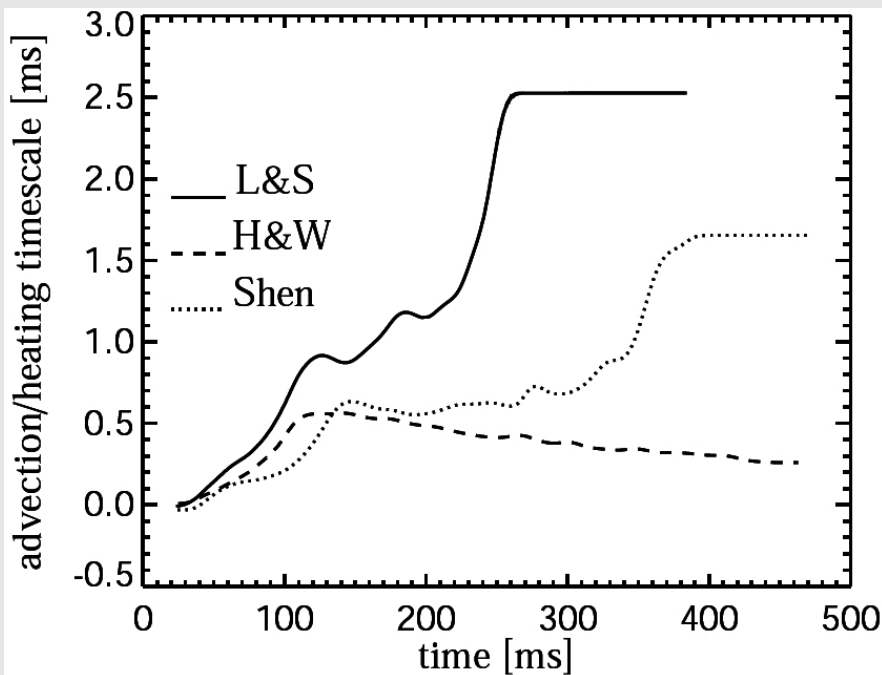
Shen EoS,
 $t \sim 470$ ms p.b.



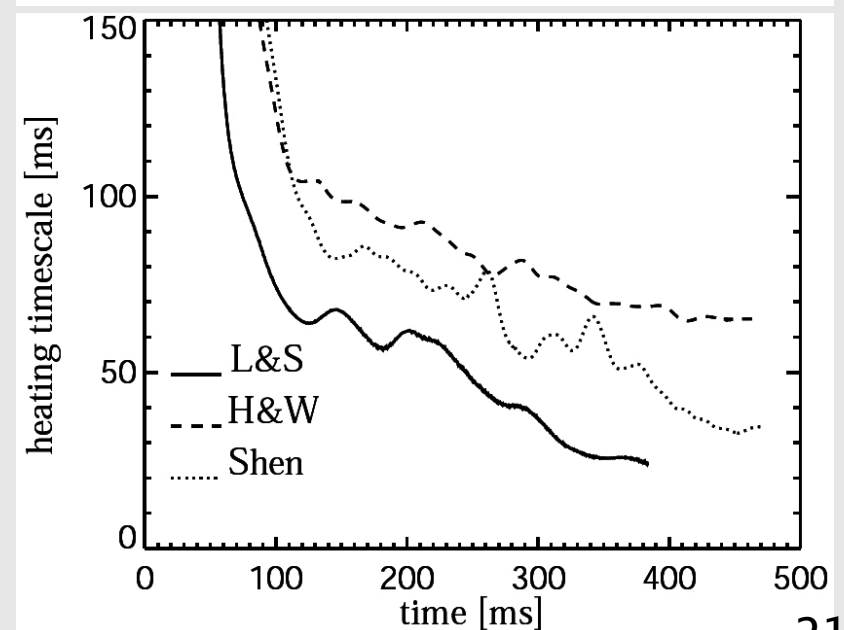
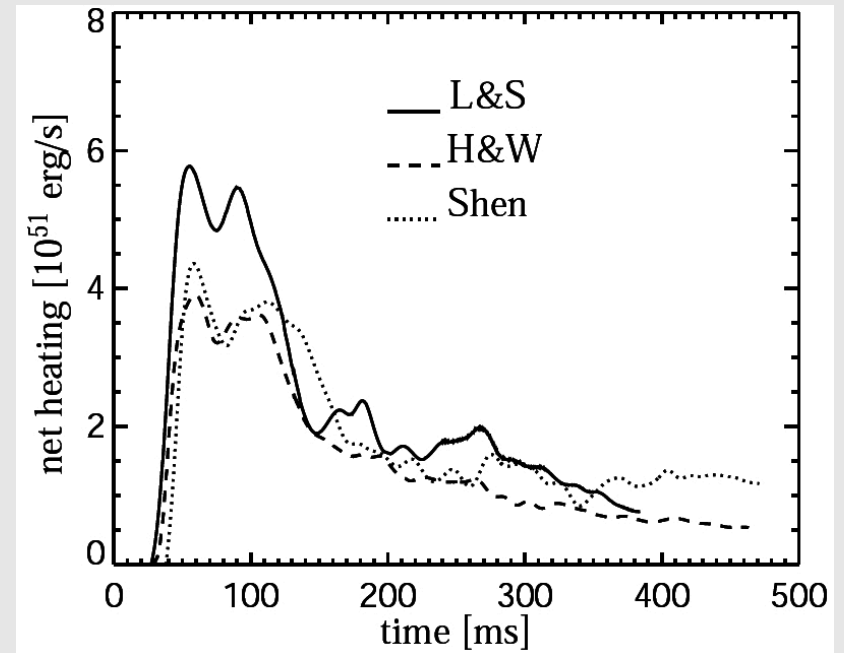
(Marek & THJ, 2009,
in preparation)

2D Explosions of $11.2 M_{\text{sun}}$ star : Test of EoS Influence

- Neutrino-heating timescale decreases with time.
- Advection timescale of matter from shock to neutron star surface increases with time.
- Both evolve favorably for explosion at later times after core bounce.



(Marek & THJ, 2009, in preparation)



2D GR Models with Neutrino Transport

- **New 2D general relativistic supernova code** with ray-by-ray 2D neutrino transport was developed by **Bernhard Müller**
(combining CoCoNuT CFC hydrodynamics code of Dimmelmeier with VERTEX neutrino transport program of Rampp & Janka, Buras)
- Excellent agreement of results in 1D with **AGILE-BOLTZTRAN** (Liebendörfer & Mezzacappa) and **VERTEX**

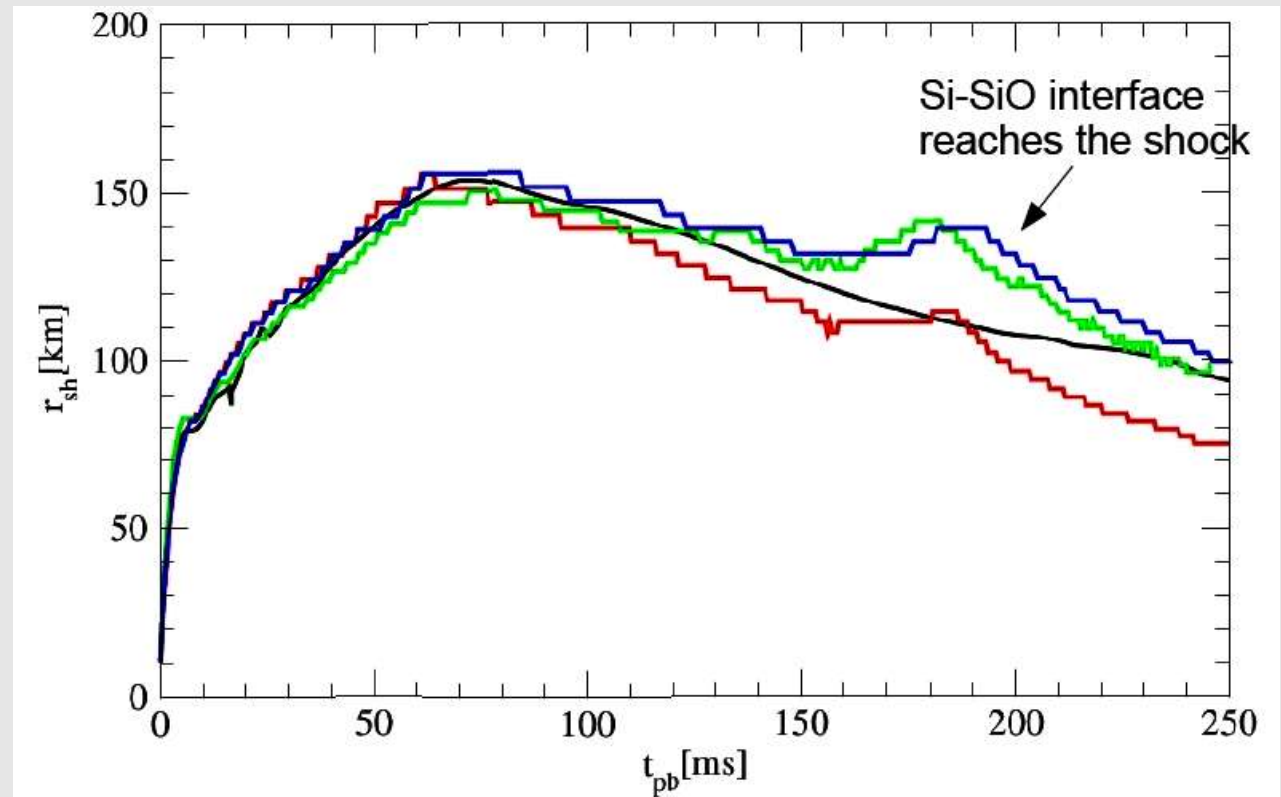
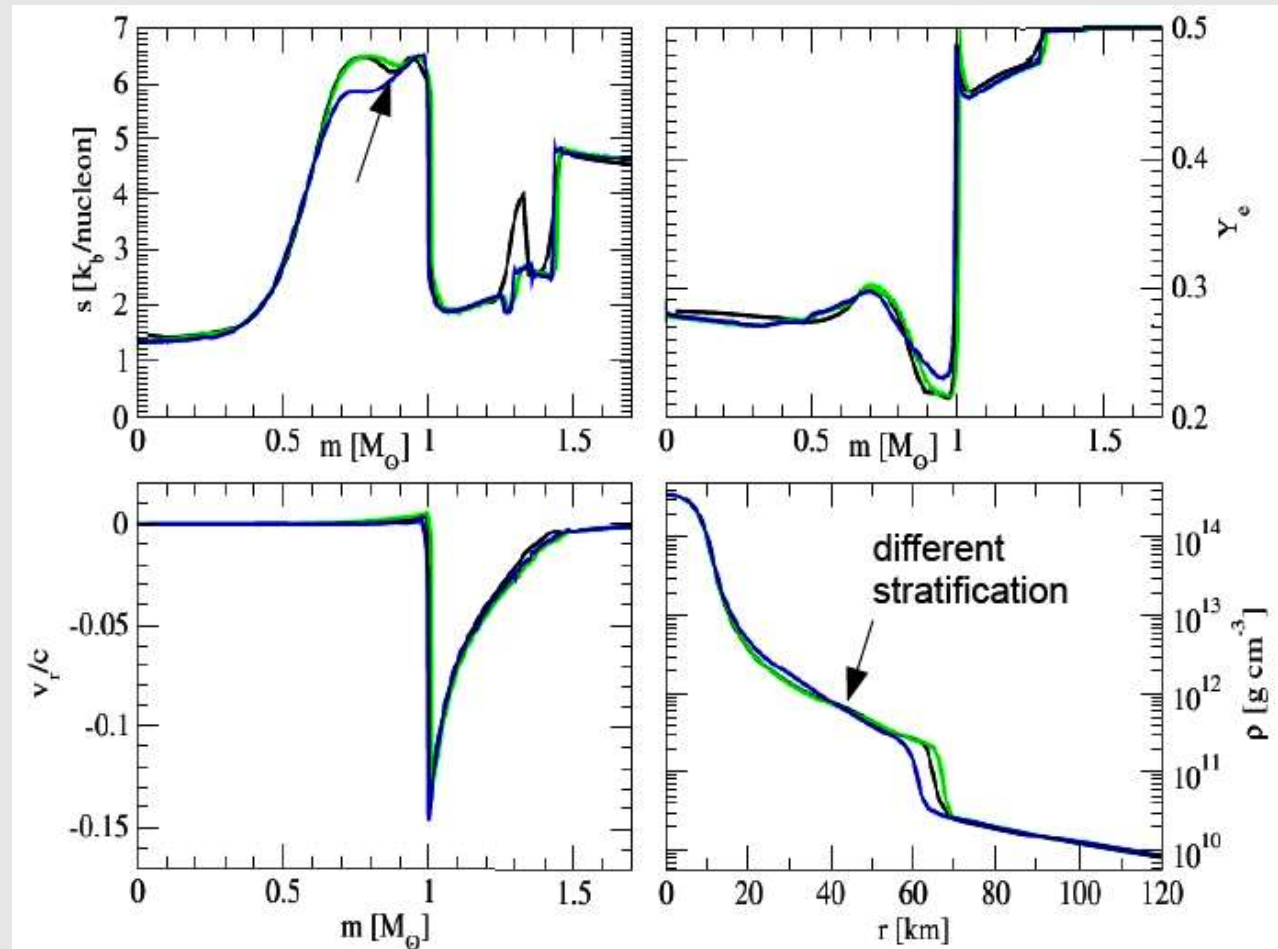


Figure 4.9.: Time evolution of the shock position for model G15 in AGILE-BOLTZTRAN (black), VERTEX-CoCoNuT (green), and VERTEX-PROMETHEUS with the effective potentials A and R (blue, red).

(B. Müller, PhD Thesis, 2009)

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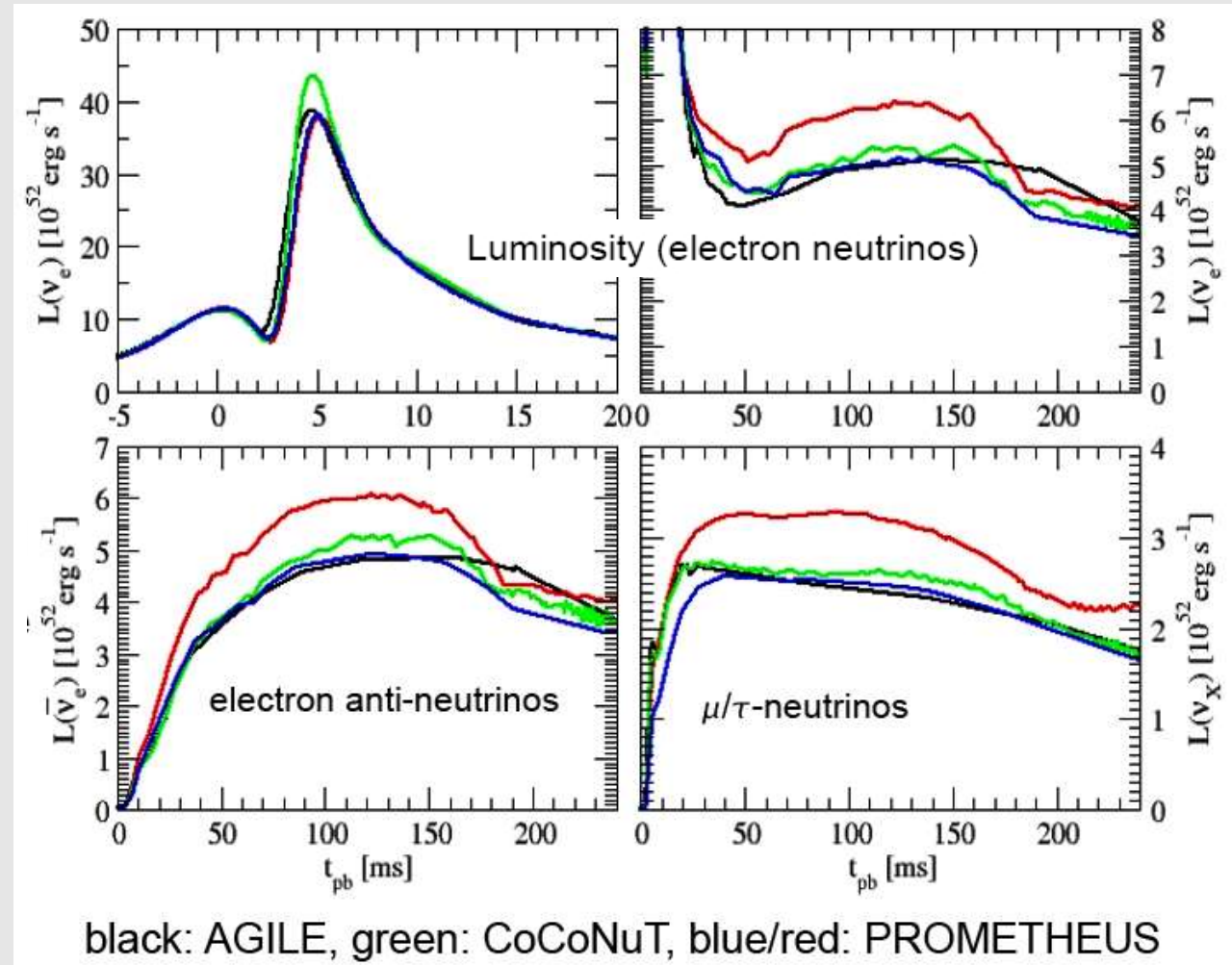


Profiles entropy, electron fraction, velocity and density 3ms after bounce obtained with AGILE (black), CoCoNuT (green), and PROMETHEUS (blue)

(B. Müller, PhD Thesis, 2009)

2D GR Models with Neutrino Transport

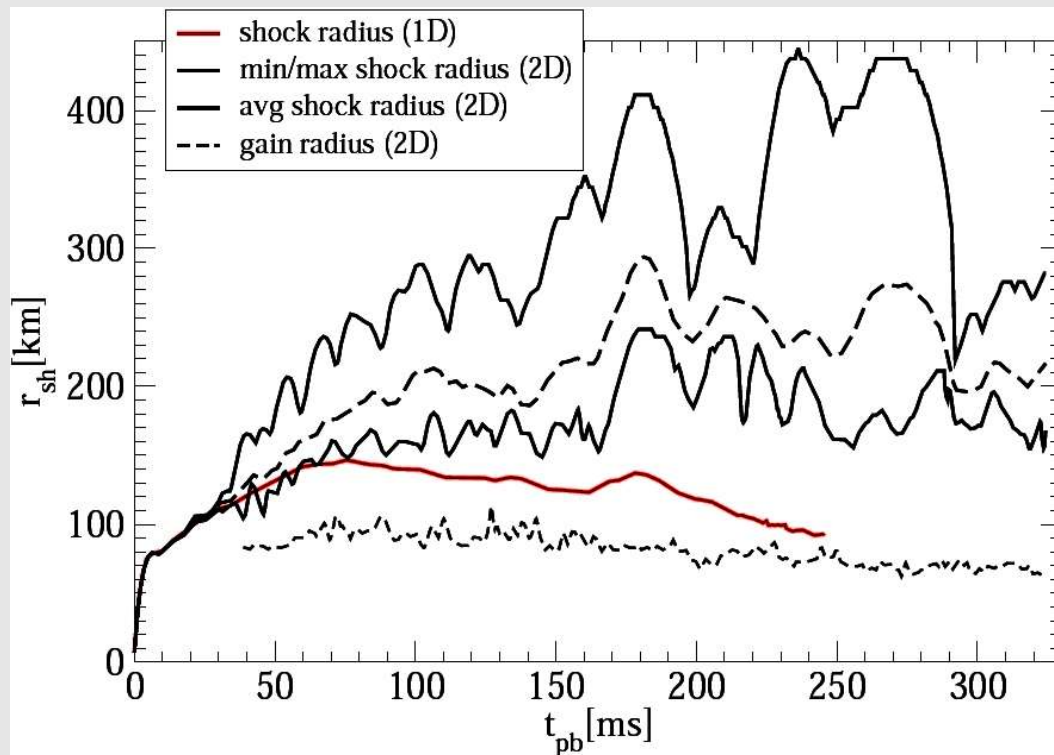
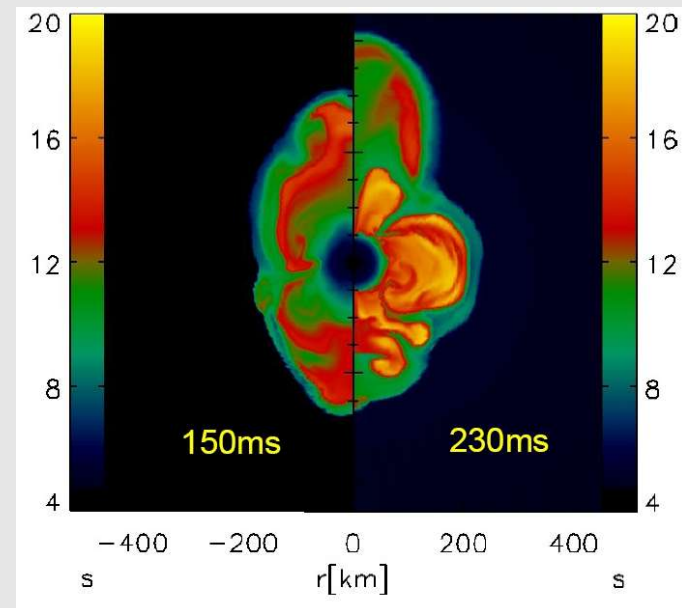
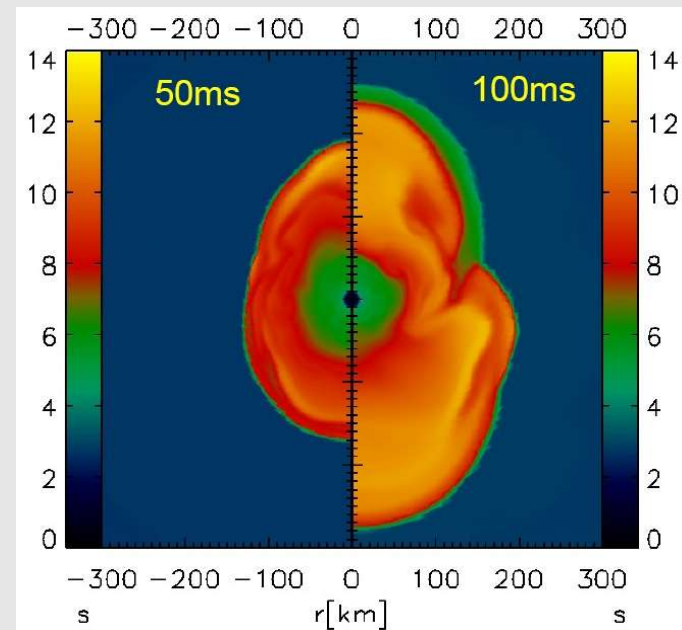
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(B. Müller, PhD Thesis, 2009)

2D GR Models with Neutrino Transport

- $15 M_{\text{sun}}$ progenitor (s15s7b2 of Weaver & Woosley) does not show any explosion within 350 ms of post-bounce evolution in 2D GR simulation
- GR simulations with new hydrodynamics code and improved neutrino scheme qualitatively confirm results by Marek & Janka (2009)!



Conclusions

- Neutrino heating can power explosions of 8–10 M_{sun} stars with ONeMg cores
-----> Crab-like supernovae (agreement of results from different groups).
- Electron-capture supernovae don't provide conditions for r-processing.
- Our most sophisticated present 2D models show that SASI & convectively supported neutrino-driven mechanism may work at least for 11–15 M_{sun} stars.
Explosions occur fairly late after core bounce
Need to verify robustness and need independent confirmation!
- 3D simulations are needed! Hydrodynamic instabilities in 3D supernova explosion models lead to higher velocities of metal clumps and stronger mixing.