

**Quantum Group Symmetry Breaking  
and  
Bose Condensation in Nonabelian Hall States**

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Joint work with Sander Bais

- Bais, Schroers, JKS. PRL 89:18601, 2002
- Bais, Schroers, JKS. JHEP 05:068, 2003
- Bais, Mathy, cond-mat/0602101, cond-mat/0602109, cond-mat/0602115

## Topological Symmetry Breaking and Bose Condensation

- Can describe topological order by extended “symmetry” concepts: TQFTs, Tensor Categories, Hopf Algebras, **Quantum Groups**

Particle types	↔	Irreducible representations
Fusion	↔	Tensor Product
Braiding	↔	R-matrix
Twist	↔	Ribbon Element

- IDEA:  
Relate topological phases by “**Symmetry Breaking**”
- Mechanism? **Bose Condensation!**  
Break the Quantum Group to the “Stabiliser” of the condensate’s order parameter

## Example: $Z_N$ Gauge Theory/ Toric Code

Particles are:

- Electric charges, labelled by representations of the gauge group.

Under gauge transformations:

$$(\alpha \in Z_N) : |q\rangle \mapsto e^{\frac{2\pi i q \alpha}{N}} |q\rangle$$

- Magnetic fluxes, labelled by monodromies (Wilson loop)  $0, 1, \dots, N-1$   
Can think of these as carriers of representations of a **dual group** (also  $Z_N$ )
- Dyons, with flux and charge (transform under  $Z_N \times Z_N$ ).

Topological interactions:

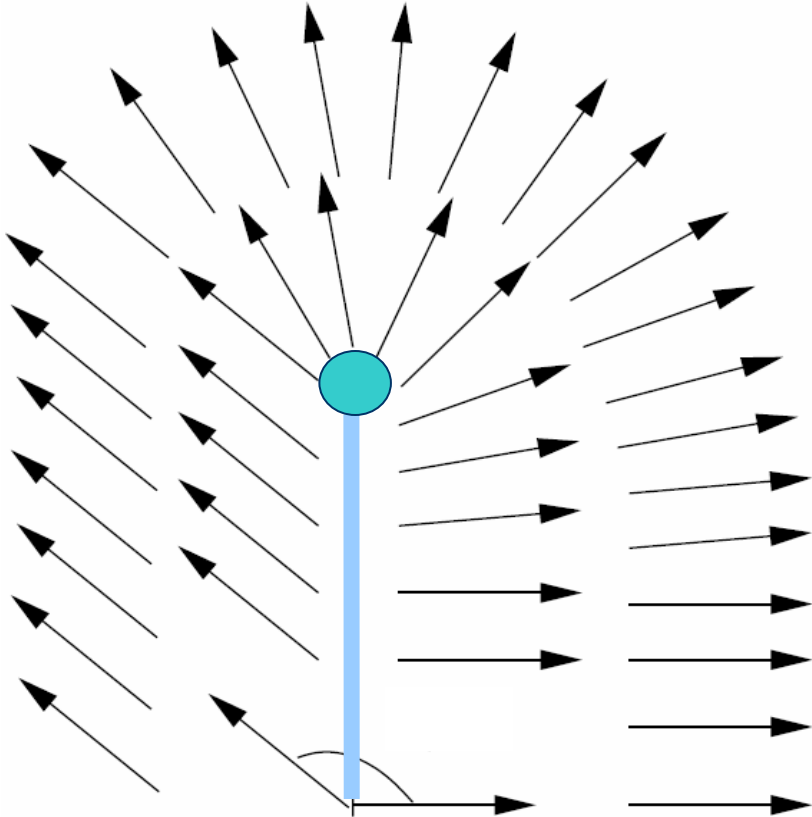
- Fusion (described by tensor product of  $Z_N \times Z_N$  reps)
- Aharonov-Bohm effect, phase factors are

$$\exp\left(\frac{2\pi i (q_1 p_2 + p_1 q_2)}{N}\right)$$

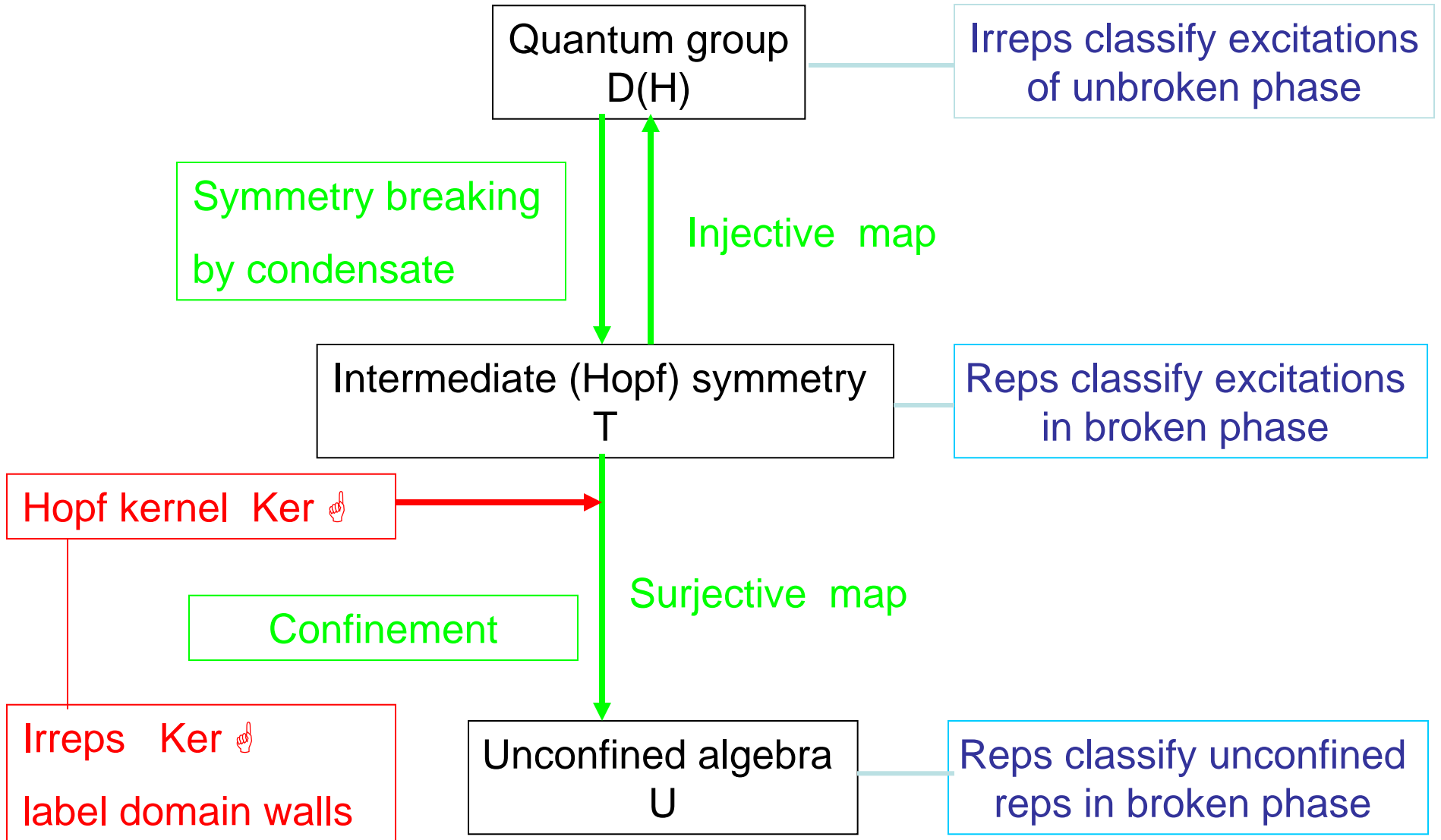
Condensation described by breaking of the full symmetry group (incl dual).

On top of that, have **confinement, from AB-interactions**

**Confined defect**

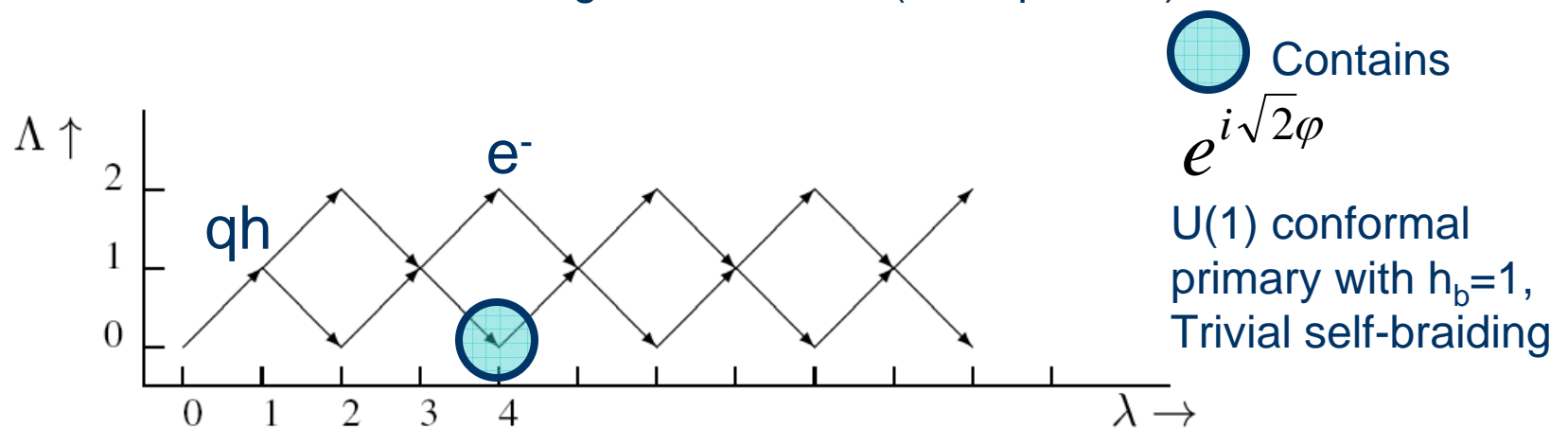


Symmetry breaking scheme

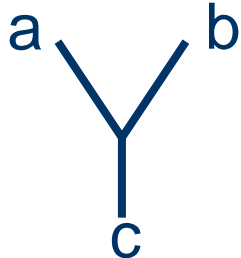




## Bosons in the Read-Rezayi States


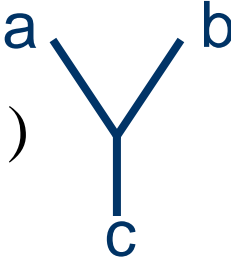
- What is a boson? A particle with
  - **trivial twist factor**/ integer conformal weight
  - **trivial self braiding** in at least on fusion channel, i.e. at least one of the fusion products also has trivial twist/integer weight
- Have a boson in the Pfaffian state (below) and lots of bosons in the higher RR-states (k=4 upwards)



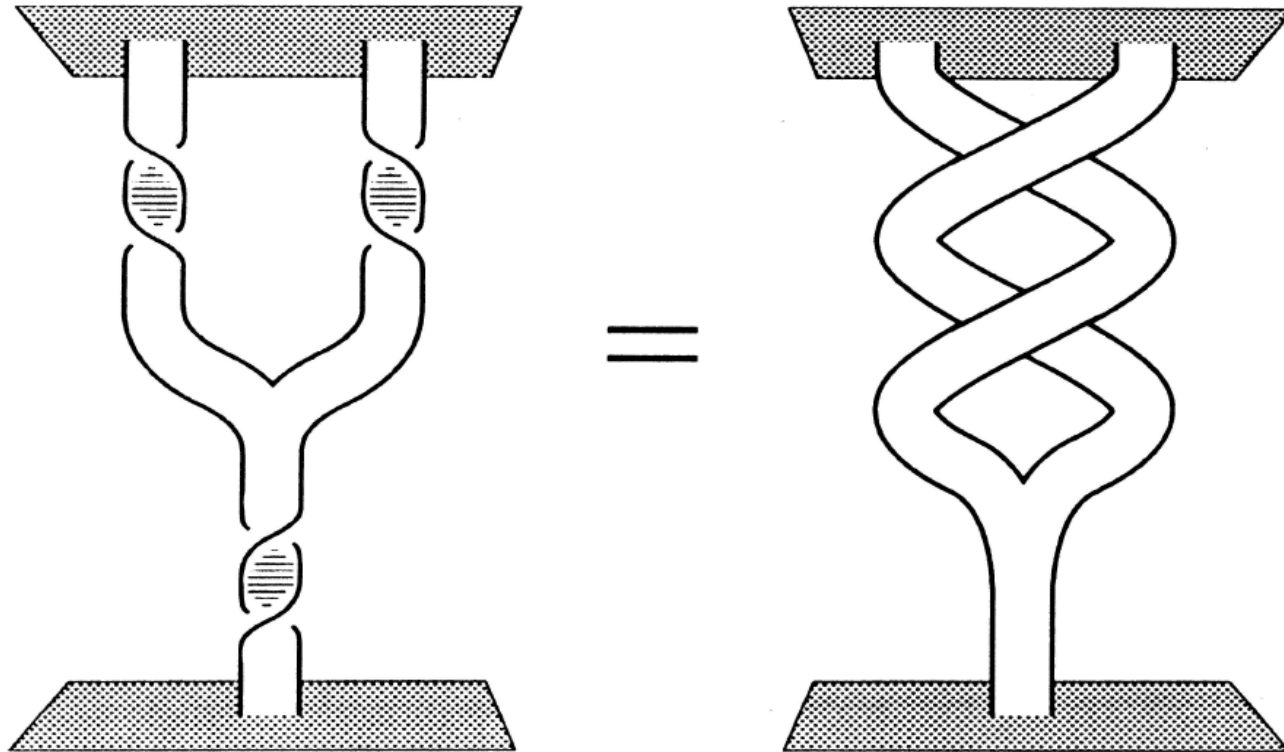
## TQFT reminder

Fusion   $a \times b = \sum_c N_c^{ab} c$

Twist   $= e^{2\pi i h_a}$  

Monodromy   $= e^{2\pi i (h_c - h_a - h_b)}$  

## Braiding and Twisting

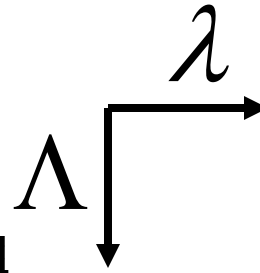




Another RR-boson

k=2

0	X	$\frac{3}{4}$	X	0	X	$\frac{3}{4}$	X	0	
X	$\frac{1}{8}$	X	$\frac{5}{8}$	X	$\frac{5}{8}$	X	$\frac{1}{8}$	X	
$\frac{1}{2}$	X	$\frac{1}{4}$	X	$\frac{1}{2}$	X	$\frac{1}{4}$	X	$\frac{1}{2}$	



$$h_{\lambda}^{\Lambda} = \frac{\Lambda(\Lambda + 2) - \lambda^2}{4(k + 2)} + \text{integer}$$

k=3

0	X	$\frac{4}{5}$	X	$\frac{1}{5}$	X	$\frac{1}{5}$	X	$\frac{4}{5}$	X	0
X	$\frac{1}{10}$	X	$\frac{7}{10}$	X	$\frac{9}{10}$	X	$\frac{7}{10}$	X	$\frac{1}{10}$	X
$\frac{2}{5}$	X	$\frac{1}{5}$	X	$\frac{3}{5}$	X	$\frac{3}{5}$	X	$\frac{1}{5}$	X	$\frac{2}{5}$
X	$\frac{7}{10}$	X	$\frac{3}{10}$	X	$\frac{1}{2}$	X	$\frac{3}{10}$	X	$\frac{7}{10}$	X

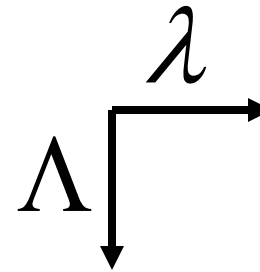
k=4

0	X	$\frac{5}{6}$	X	$\frac{1}{3}$	X	$\frac{1}{2}$	X	$\frac{1}{3}$	X	$\frac{5}{6}$	X	0
X	$\frac{1}{12}$	X	$\frac{3}{4}$	X	$\frac{1}{12}$	X	$\frac{1}{12}$	X	$\frac{3}{4}$	X	$\frac{1}{12}$	X
$\frac{1}{3}$	X	$\frac{1}{6}$	X	$\frac{2}{3}$	X	$\frac{5}{6}$	X	$\frac{2}{3}$	X	$\frac{1}{6}$	X	$\frac{1}{3}$
X	$\frac{7}{12}$	X	$\frac{1}{4}$	X	$\frac{7}{12}$	X	$\frac{7}{12}$	X	$\frac{1}{4}$	X	$\frac{7}{12}$	X
0	X	$\frac{5}{6}$	X	$\frac{1}{3}$	X	$\frac{1}{2}$	X	$\frac{1}{3}$	X	$\frac{5}{6}$	X	0

Even more RR-bosons

k=5

0	X	$\frac{6}{7}$	X	$\frac{3}{7}$	X	$\frac{5}{7}$	X	
X	$\frac{1}{14}$	X	$\frac{11}{14}$	X	$\frac{3}{14}$	X	$\frac{5}{14}$	
$\frac{2}{7}$	X	$\frac{1}{7}$	X	$\frac{5}{7}$	X	0	X	
X	$\frac{1}{2}$	X	$\frac{3}{14}$	X	$\frac{9}{14}$	X	$\frac{11}{14}$	
$\frac{6}{7}$	X	$\frac{5}{7}$	X	$\frac{2}{7}$	X	$\frac{4}{7}$	X	
X	$\frac{3}{14}$	X	$\frac{13}{14}$	X	$\frac{5}{14}$	X	$\frac{1}{2}$	



k=6

0	X	$\frac{7}{8}$	X	$\frac{1}{2}$	X	$\frac{7}{8}$	X	0	
X	$\frac{1}{16}$	X	$\frac{13}{16}$	X	$\frac{5}{16}$	X	$\frac{9}{16}$	X	
$\frac{1}{4}$	X	$\frac{1}{8}$	X	$\frac{3}{4}$	X	$\frac{1}{8}$	X	$\frac{1}{4}$	
X	$\frac{7}{16}$	X	$\frac{3}{16}$	X	$\frac{11}{16}$	X	$\frac{15}{16}$	X	
$\frac{3}{4}$	X	$\frac{5}{8}$	X	$\frac{1}{4}$	X	$\frac{5}{8}$	X	$\frac{3}{4}$	
X	$\frac{1}{16}$	X	$\frac{13}{16}$	X	$\frac{5}{16}$	X	$\frac{9}{16}$	X	
$\frac{1}{2}$	X	$\frac{3}{8}$	X	0	X	$\frac{3}{8}$	X	$\frac{1}{2}$	

$$h_{\lambda}^{\Lambda} = \frac{\Lambda(\Lambda + 2) - \lambda^2}{4(k + 2)} + \text{integer}$$

$\lambda$

k=7

0	X	$\frac{8}{9}$	X	$\frac{5}{9}$	X	0	X	$\frac{2}{9}$	X	
X	$\frac{1}{18}$	X	$\frac{5}{6}$	X	$\frac{7}{18}$	X	$\frac{13}{18}$	X	$\frac{5}{6}$	
$\frac{2}{9}$	X	$\frac{1}{9}$	X	$\frac{7}{9}$	X	$\frac{2}{9}$	X	$\frac{4}{9}$	X	
X	$\frac{7}{18}$	X	$\frac{1}{6}$	X	$\frac{13}{18}$	X	$\frac{1}{18}$	X	$\frac{1}{6}$	
$\frac{2}{3}$	X	$\frac{5}{9}$	X	$\frac{2}{9}$	X	$\frac{2}{3}$	X	$\frac{8}{9}$	X	
X	$\frac{17}{18}$	X	$\frac{13}{18}$	X	$\frac{5}{18}$	X	$\frac{11}{18}$	X	$\frac{13}{18}$	
$\frac{1}{3}$	X	$\frac{2}{9}$	X	$\frac{8}{9}$	X	$\frac{1}{3}$	X	$\frac{5}{9}$	X	
X	$\frac{13}{18}$	X	$\frac{1}{2}$	X	$\frac{1}{18}$	X	$\frac{7}{18}$	X	$\frac{1}{2}$	

k=8

0	X	$\frac{9}{10}$	X	$\frac{3}{5}$	X	$\frac{1}{10}$	X	$\frac{2}{5}$	X	$\frac{1}{2}$	
X	$\frac{1}{20}$	X	$\frac{17}{20}$	X	$\frac{9}{20}$	X	$\frac{17}{20}$	X	$\frac{1}{20}$	X	
$\frac{1}{5}$	X	$\frac{1}{10}$	X	$\frac{4}{5}$	X	$\frac{3}{10}$	X	$\frac{3}{5}$	X	$\frac{7}{10}$	
X	$\frac{7}{20}$	X	$\frac{3}{20}$	X	$\frac{3}{4}$	X	$\frac{3}{20}$	X	$\frac{7}{20}$	X	
$\frac{3}{5}$	X	$\frac{1}{2}$	X	$\frac{1}{5}$	X	$\frac{7}{10}$	X	0	X	$\frac{1}{10}$	
X	$\frac{17}{20}$	X	$\frac{13}{20}$	X	$\frac{1}{4}$	X	$\frac{13}{20}$	X	$\frac{17}{20}$	X	
$\frac{1}{5}$	X	$\frac{1}{10}$	X	$\frac{4}{5}$	X	$\frac{3}{10}$	X	$\frac{3}{5}$	X	$\frac{7}{10}$	
X	$\frac{11}{20}$	X	$\frac{7}{20}$	X	$\frac{19}{20}$	X	$\frac{7}{20}$	X	$\frac{11}{20}$	X	
0	X	$\frac{9}{10}$	X	$\frac{3}{5}$	X	$\frac{1}{10}$	X	$\frac{2}{5}$	X	$\frac{1}{2}$	

## Problem with the old scheme

Problem is: need to select a state inside a representation as order parameter.  
Problem because:

- Quantum Group Symmetry is often “hidden”:  
internal labels of particles are not (fully) physical  
→ what does the “order parameter” mean?
- Particles’ “internal spaces” may have non-integer quantum dimensions...
- Tensor product in many theories (incl. the ones relevant to the Hall effect) is truncated. Product states as expected for Bose condensates usually do not survive the truncation

## The New Symmetry Breaking: Forget the Algebra!

No algebra, No states, Just Labels and Branching:

$$a \rightarrow \sum_i n_{a,i} a_i$$

Three requirements:

1. The new labels themselves form a fusion algebra (need associativity, vacuum and charge conjugation)
2. Branching and fusion are compatible,

$$a \otimes b \rightarrow \left( \sum_i n_{a,i} a_i \right) \otimes \left( \sum_i n_{b,i} b_i \right)$$

3. Not more branching/identification than necessary for 1. and 2. (want the full “stabiliser”)

# Breaking $U_q(\mathfrak{su}(4))$

## $SU(2)_4$

0	$d_0 = 1$	$h_0 = 0$
1	$d_1 = \sqrt{3}$	$h_1 = 1/8$
2	$d_2 = 2$	$h_2 = 1/3$
3	$d_3 = \sqrt{3}$	$h_3 = 5/8$
4	$d_4 = 1$	$h_4 = 1$

$$1 \times 1 = 0 + 2$$

$$1 \times 2 = 1 + 3 \quad 2 \times 2 = 0 + 2 + 4$$

$$1 \times 3 = 2 + 4 \quad 2 \times 3 = 1 + 3 \quad 3 \times 3 = 0$$

$$1 \times 4 = 3 \quad 2 \times 4 = 2 \quad 3 \times 4 = 1 \quad 4 \times 4 = 0$$

# Condensate, splitting and identification

Assume a bosonic condensate forms in the 4 rep of  $SU(2)_4$ :

$$2 \times 2 = 0 + 2 + 4 = 0 + 2 + 0$$

$$\Rightarrow 2 := 2_1 + 2_2 \quad \text{possible because } d_2=2$$

$$2_1 \times 2_1 + 2_1 \times 2_2 + 2_2 \times 2_1 + 2_2 \times 2_2 = 0 + 2_1 + 2_2 + 0$$

$$\Rightarrow 2_1 \times 2_2 = 0$$

$$\text{if } 2_1 \times 2_1 = 2_1$$

$$\text{then } 2_2 \times (2_1 \times 2_1) = 2_2 \times 2_1 = 0$$

$$(2_2 \times 2_1) \times 2_1 = 0 \times 2_1 = 2_1$$

$$\Rightarrow 2_1 \times 2_1 = 2_2 \quad \text{and} \quad 2_2 \times 2_2 = 2_1$$

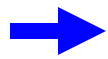
$$1 \times 1 = 0 + 2_1 + 2_2$$

$$1 \times 3 = 0 + 2_1 + 2_2$$

$$\Rightarrow 1 \Leftrightarrow 3$$

0

1

  $2 := 2_1 + 2_2$

$3 \Leftrightarrow 1$

$4 \Leftrightarrow 0$

$$1 \times 1 = 0 + 1$$

$$1 \times 2_1 = 1 \quad 2_1 \times 2_1 = 2_2$$

$$1 \times 2_2 = 1 \quad 2_1 \times 2_2 = 1 \quad 2_2 \times 2_2 = 2_2$$

## Confinement and Braiding

To see which of the particles in the broken theory are confined, look at braiding with the condensed particle.

**How?** For particle  $a_i$ , look in all channels of the old theory that cover  $a_i \times 1 = a_i$

Now notice: Fields that cover 1 have trivial twist factor (condensate is bosonic). Hence braiding with the vacuum is trivial and  $a_i$  is not confined precisely when all the fields that branch to  $a_i$  have equal twist factors (or conformal dimensions that differ by integers).

The non-confined particles all have well defined monodromies with each other, given by their twist factors (which are unambiguously defined from the branching).

# Confinement

From fusion rules (of the original algebra)  
with the condensate 4 and conformal weights one finds that  
the 1 and 3 are confined.

The unconfined algebra becomes  $SU(3)_1$ :

$$\begin{array}{l} 2_1 \times 2_1 = 2_2 \\ 2_1 \times 2_2 = 0 \\ 2_2 \times 2_2 = 2_1 \end{array} \longleftrightarrow \begin{array}{l} 3 \times 3 = \bar{3} \\ 3 \times \bar{3} = 1 \\ \bar{3} \times \bar{3} = 3 \end{array}$$



# Relation to Conformal Embedding

Central charges satisfy  $c(G) = c(H) \implies c(G/H) = 0$

Coset algebra is trivial.

$\implies$  Finite branching of inf. Dim. KM representations

Example:  $SU(2)_4 \implies SU(3)_1$  ( $c=2$ )

$SU(3)_1$  Irreps:

1	$d_1 = 1$	$h_1 = 0$
3	$d_3 = 1$	$h_3 = 1/3$
$\bar{3}$	$d_{\bar{3}} = 1$	$h_{\bar{3}} = 1/3$

$3 \times 3 = \bar{3}$
$3 \times \bar{3} = 1$
$\bar{3} \times \bar{3} = 3$



branching

$1 \rightarrow 0 + 4$
$3 \rightarrow 2$
$\bar{3} \rightarrow 2$

## Some results for the Read-Rezayi states

- $k=2$ :  $\text{Ising} \times U(1)_8 \rightarrow \text{SO}(3)_2 \times U(1)_2$  or  $(\text{Ising}/Z_2) \times U(1)_2$   
Interpretation not obvious (superconductor?!).  
Possibly connected to Fradkin-Nayak-Schoutens '98?
- $k=4$ :  $\text{Pf}_4 \times U(1)_{24} \rightarrow (U(1)_6 \times U(1)_6)/Z_2$   
Expect a Hall state at filling  $2/3$  (neutral condensate)  
Get Abelian topological order with right behavior for the quasiholes.  
More precise connection difficult (naïve ground state gives  $\nu=1/3$  Laughlin)
- $k=6$ : Much like  $k=2$ , but stays nonabelian (contains  $\text{SO}(3)_6$ )
- $k=7$ : Stays nonabelian (contains “twisted” version of  $\text{Pf}_7$ )
- $k=8$ : Stays nonabelian and contains Fibonacci  $\times$  Fibonacci (chiral version)

# Conformal embedding of KM algebra's

The case  $SO(5)_1$

Irreps:	1	$d_1 = 1$	$h_1 = 0$	$4 \times 4 = 1 + 5$
	4	$d_4 = \sqrt{2}$	$h_4 = 5/16$	$5 \times 5 = 1$
	5	$d_5 = 1$	$h_5 = 1/2$	$4 \times 5 = 4$

Conformal embedding:  $SU(2)_{10} \rightarrow SO(5)_1$  ( $c = 5/2$ )

Branching rules:  $0 \Rightarrow 0+6$   
 $4 \Rightarrow 3+7$   
 $5 \Rightarrow 4+10$

# The condensate

Condensation of the 6 demands splitting  $6 := 6_1 + 6_2$  (with  $\dim 6_1 = 1$  and  $6_1 \times 6_1 = 0$ )  
 Because of condensate we get further splitting and identifications:

$$\begin{aligned}
 d_0 &= d_{10} = 1 \\
 d_1 &= d_9 = \sqrt{2 + \sqrt{3}} \\
 d_2 &= d_8 = 1 + \sqrt{3} \\
 d_3 &= d_7 = \sqrt{2} + \sqrt{2 + \sqrt{3}} \\
 d_4 &= d_6 = 2 + \sqrt{3} \\
 d_5 &= 2\sqrt{2 + \sqrt{3}}
 \end{aligned}$$

$$\begin{aligned}
 0 \\
 1 \\
 2 \\
 3 &:= 3_1 + 3_2 \\
 4 &:= 4_1 + 4_2 \Leftrightarrow 4_1 + 2 \\
 5 &:= 5_1 + 5_2 \Leftrightarrow 1 + 3_1 \\
 6 &:= 6_1 + 6_2 \Leftrightarrow 0 + 2 \\
 7 &:= 7_1 + 7_2 \Leftrightarrow 1 + 4_2 \\
 8 &\Leftrightarrow 2 \\
 9 &\Leftrightarrow 3_1 \\
 10 &\Leftrightarrow 4_1
 \end{aligned}$$

$$\begin{aligned}
 d_0 &= 1 \\
 d_1 &= \sqrt{2 + \sqrt{3}} \\
 d_2 &= d_8 = 1 + \sqrt{3} \\
 d_{3_1} &= \sqrt{2 + \sqrt{3}} \\
 d_{3_2} &= \sqrt{2} \\
 d_{4_1} &= 1
 \end{aligned}$$

# Intermediate model

Fusion rules:

$$1 \times 1 = 0 + 2$$

$$1 \times 2 = 1 + 3_1 + 3_2 \quad 2 \times 2 = 0 + 2 + 2 + 4$$

$$1 \times 3_1 = 2 + 4_1 \quad 2 \times 3_1 = 1 + 3_1 + 1 + 3_2 \quad 3_1 \times 3_1 = 0 + 2$$

$$1 \times 3_2 = 2 \quad 2 \times 3_2 = 1 + 3_1 \quad 3_1 \times 3_2 = 2 \quad 3_2 \times 3_2 = 0 + 4_1$$

$$1 \times 4_1 = 3_1 \quad 2 \times 4_1 = 2 \quad 3_1 \times 4_1 = 1 \quad 3_2 \times 4_1 = 3_2 \quad 4_1 \times 4_1 = 0$$

# Braiding and confinement

$SO(5)_1$   $SU(2)_{10}$

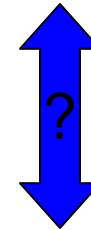
$$\begin{aligned} h_1 &= 0 \\ h_4 &= 5/16 \\ h_5 &= 1/2 \end{aligned}$$

$$\begin{aligned} h_0 &= 0 \\ h_1 &= \frac{1}{16} \\ h_2 &= \frac{1}{6} \\ h_3 &= \frac{5}{16} \\ h_4 &= \frac{1}{2} \\ h_5 &= \frac{35}{48} \\ h_6 &= 1 \\ h_7 &= \frac{21}{16} \\ h_8 &= \frac{5}{3} \\ h_9 &= \frac{33}{16} \\ h_{10} &= \frac{5}{2} \end{aligned}$$

$$\begin{aligned} 0 \\ 1 \\ 2 \\ 3 &:= 3_1 + 3_2 \\ 4 &:= 4_1 + 4_2 \Leftrightarrow 4_1 + 2 \\ 5 &:= 5_1 + 5_2 \Leftrightarrow 1 + 3_1 \\ 6 &:= 6_1 + 6_2 \Leftrightarrow 0 + 2 \\ 7 &:= 7_1 + 7_2 \Leftrightarrow 1 + 3_2 \\ 8 &\Leftrightarrow 2 \\ 9 &\Leftrightarrow 3_1 \\ 10 &\Leftrightarrow 4_1 \end{aligned}$$

$$\begin{aligned} 0 \\ 3_2 \\ 4_1 \end{aligned}$$

$$\begin{aligned} 3_2 \times 3_2 &= 0 + 4_1 \\ 3_2 \times 4_1 &= 3_2 \\ 4_1 \times 4_1 &= 0 \end{aligned}$$



$$\begin{aligned} 4 \times 4 &= 1 + 5 \\ 4 \times 5 &= 4 \\ 5 \times 5 &= 1 \end{aligned}$$

(Algebra of Ising or  $SU(2)_2$   
have different conformal weights!)

## Summary and Outlook

### Results

- Extended Topological symmetry breaking to TQFTs with non-integer quantum dimensions
- Found connection to conformal embeddings
- Had a first go at application to nonabelian FQH states

### Questions/Future Work

- Found Fusion and twist factors.  
How to determine the rest of the TQFT (half-braidings, F-symbols...) ?  
Note: often fixed by consistency (always?)
- Work suggests conformal embeddings of coset chiral algebras.  
Interesting CFT problem...
- Further Physical applications....