

QSiM19

SCHWARZSCHILD BLACK HOLES AS
MACROSCOPIC QUANTUM SYSTEMS

PAOLA VERRUCCHI

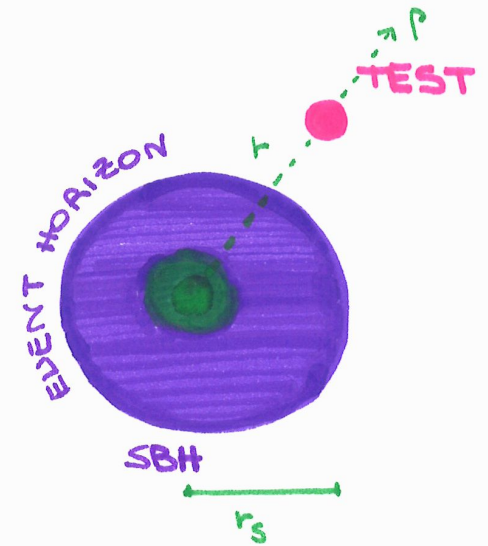
WITH

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INSTITUTE FOR NUCLEAR PHYSICS INFN
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THE SCHWARZSCHILD BLACK HOLE

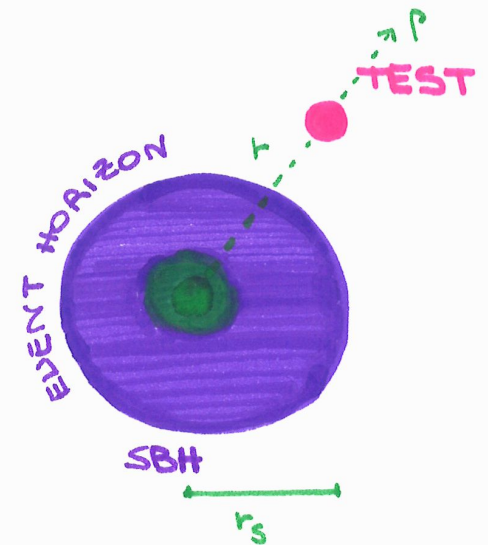
$$h_{\text{SBH}}^{\text{TEST}}(p, r) = \frac{p^2}{2} + \frac{1}{2} \left(\frac{L^2}{r^2} + \mu^2 \right) \left(1 - \frac{r_s}{r} \right)$$



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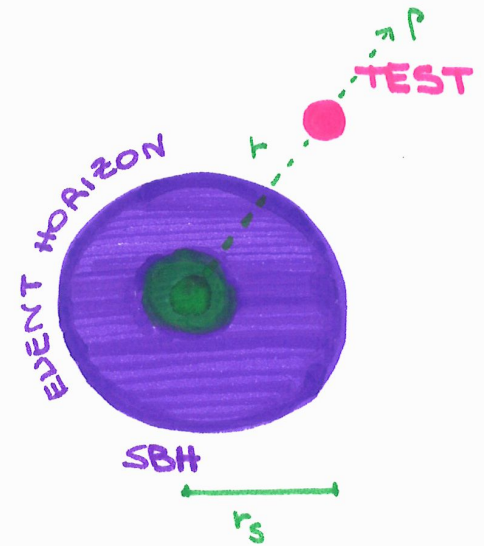
r_s Schwarzschild radius
 L modulus of the conserved angular momentum
 p, r radial momentum and position $\{p, r\}_{\text{PB}} = 1$
 $\mu^2 = 1, 0$ mass²



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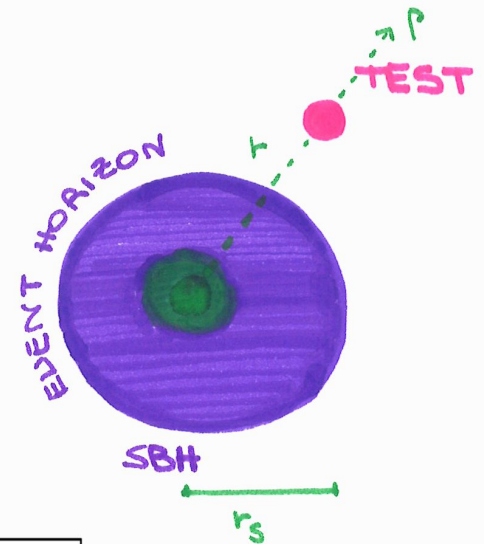
$$h_{\text{SBH}}^{\text{TEST}}(p, r) = \frac{p^2}{2} + \frac{1}{2} \underbrace{\left(\frac{L^2}{r^2} + \mu^2 \right)}_{V(r)} \left(1 - \frac{r_s}{r} \right)$$

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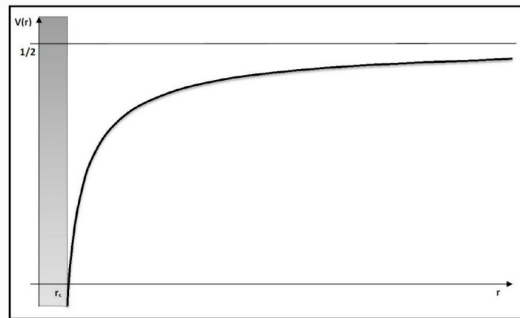
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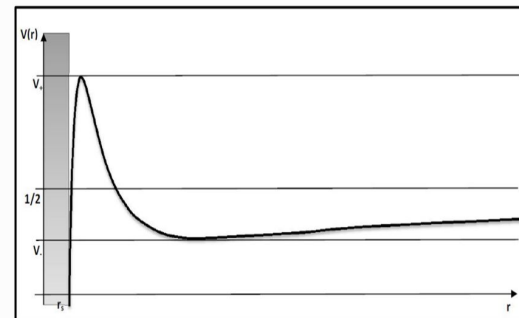


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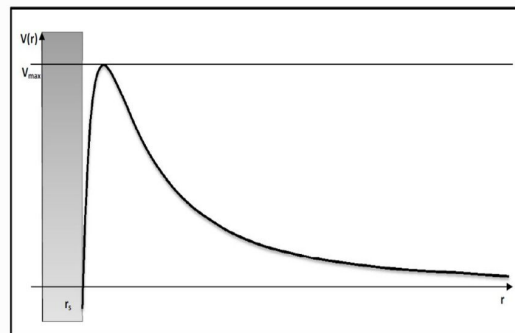


SBH potential for matter if $L^2 < 3r_s^2$



SBH potential for matter if $L^2 > 3r_s^2$

$\mu^2 = 0$



SBH potential for light

QUANTUM MECHANICS vs GENERAL RELATIVITY



QUANTUM MECHANICS vs GENERAL RELATIVITY



ALGEBRA

vs

GEOMETRY

MICRO

vs

MACRO



EVOLUTION IN TIME

vs

SYMMETRY TRANSFORMATION

QUANTUM MECHANICS vs GENERAL RELATIVITY



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generalized coherent states

GCS



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GCS

large- N quantum (field) theory

large- N Q(F)T



QUANTUM MECHANICS vs GENERAL RELATIVITY



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generalized coherent states

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Page & Wootters mechanism

PalW



QUANTUM MECHANICS vs GENERAL RELATIVITY

ALGEBRA vs GEOMETRY



generalized coherent states

GENERALIZED COHERENT STATES

GCS

group-theoretic construction

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GCS

group-theoretic construction

$\mathcal{Q} \rightarrow \mathfrak{g} \& \mathfrak{h} \rightarrow \mathfrak{g}$ group (Lie)

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GCS

group-theoretic construction

$\mathcal{Q} \rightarrow \mathfrak{g}_1 \& \mathfrak{H} \rightarrow \mathfrak{g}$ "dynamical" group (Lie)

propagators of \mathcal{Q} are elements of a unitary irreducible representation of \mathfrak{g} obtained from \mathfrak{g}_1 via a Lie exponential map

GENERALIZED COHERENT STATES

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$|R\rangle \in \mathfrak{h}$ reference state

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$\mathfrak{g}/\mathfrak{F}$

coset

$\hat{\Omega} \in \mathfrak{g}/\mathfrak{F}$



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\mathcal{M}



manifold

GCS

$$\hat{H} \in \mathcal{U}/\mathcal{R}$$

$$\hat{H} = e^{\sum_{\mathbf{p}} (\Omega_{\mathbf{p}} \hat{E}_{\mathbf{p}} - \Omega_{\mathbf{p}}^* \hat{E}_{-\mathbf{p}})}$$

$$|R\rangle = e^{\sum_{\mathbf{p}} (\Omega_{\mathbf{p}} \hat{E}_{\mathbf{p}} - \Omega_{\mathbf{p}}^* \hat{E}_{-\mathbf{p}})} |R\rangle \in \mathcal{H}$$

GCS

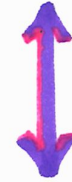
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$$\Omega := (\Omega_1, \Omega_2, \dots)$$



$$\in \mathcal{M}$$

manifold with a symplectic structure

GCS

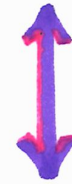
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$$\Omega \rightarrow \zeta \rightarrow \tilde{\mathcal{Z}}$$

complex projective coordinates

$$\{f, g\} = i \sum_{\beta} \left(\frac{\partial f \partial g}{\partial \tau_{\alpha} \partial \tau_{\beta}^*} - \frac{\partial f \partial g}{\partial \tau_{\beta}^* \partial \tau_{\alpha}} \right)$$

GCS

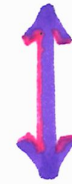
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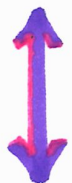
$$\zeta = \frac{1}{\sqrt{2}} (w - iw^*)$$

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$$d\mu(\Omega) : \int_{\mathcal{M}} d\mu(\Omega) |\Omega\rangle\langle\Omega| = \hat{\mathbb{1}}_{\mathfrak{H}}$$

invariant measure

PSEUDO-SPIN COHERENT STATES $SU(1,1)$

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$SU(1,1)$ algebra : $\text{span} \{ \hat{K}_0, \hat{K}_1, \hat{K}_2 \}$ $[\hat{K}_\alpha, \hat{K}_\beta] = i\epsilon_{\alpha\beta\gamma} \hat{K}_\gamma$

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\updownarrow
Bergmann index

$$\hat{K}^2 |k, m\rangle = k(k-1) |k, m\rangle$$

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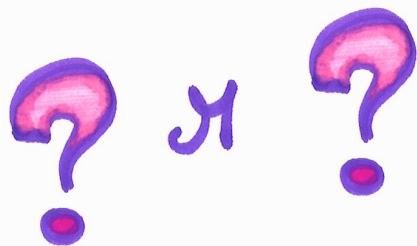
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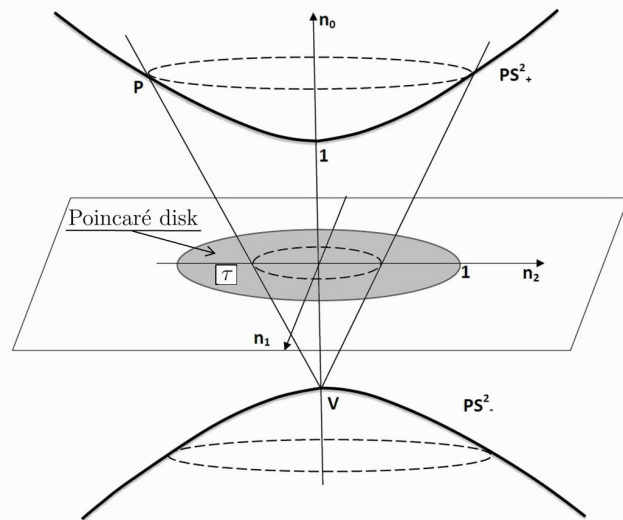
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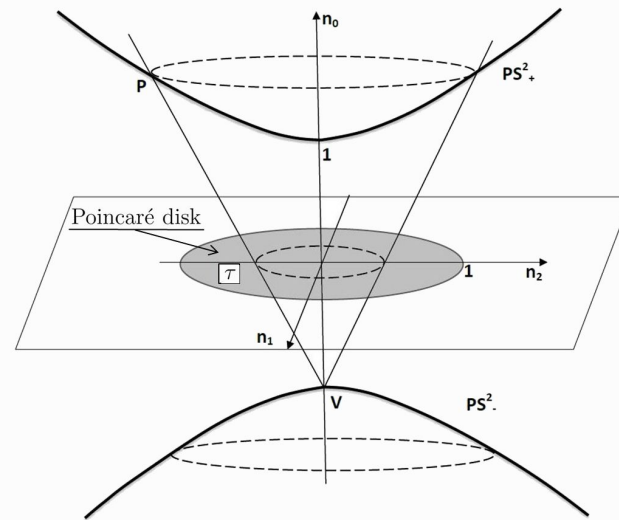
THE MANIFOLD PS^2 (PSEUDO - SPHERE)



two-sheets hyperboloid

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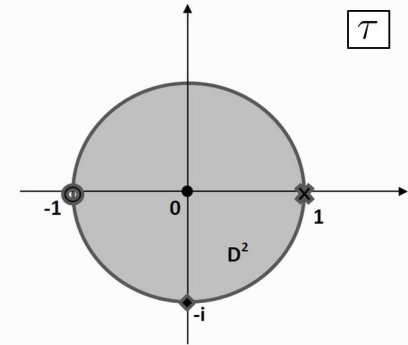
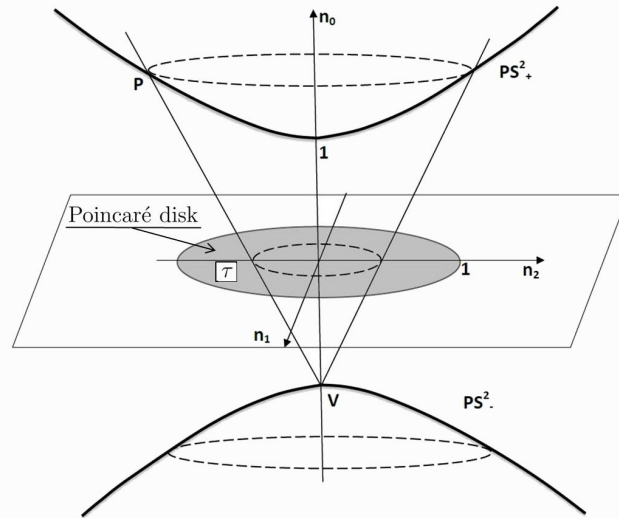
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Poincaré disk D^2

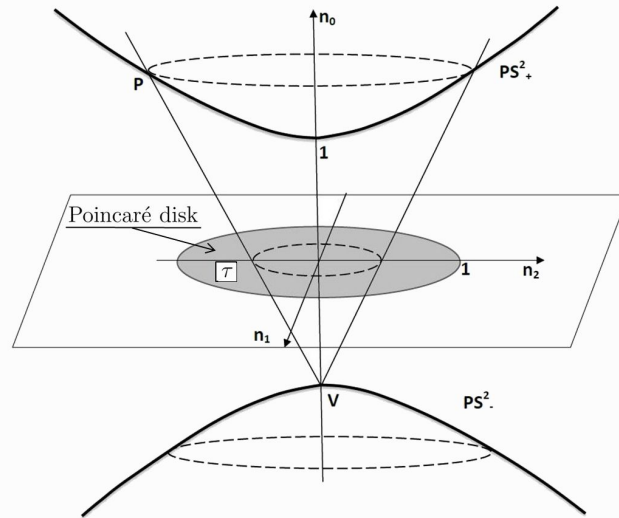
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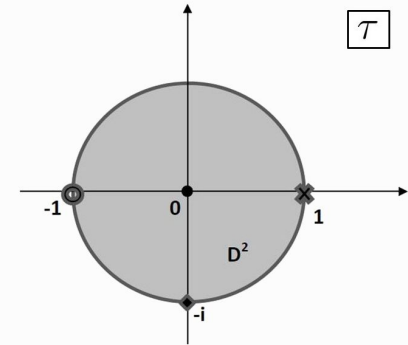
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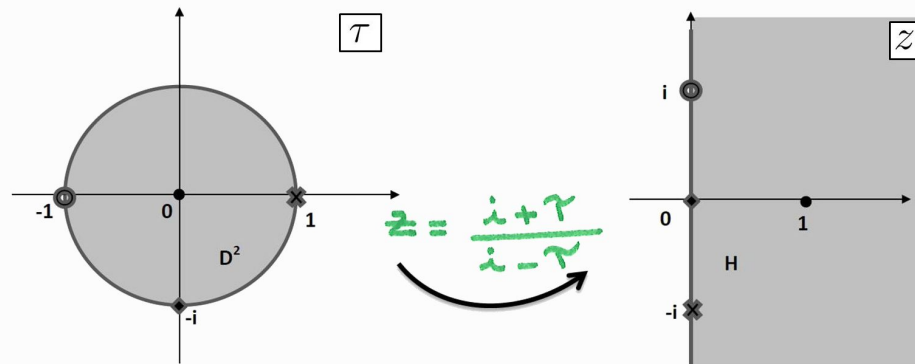
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Poincaré half-plane H

QUANTUM MECHANICS vs GENERAL RELATIVITY

MICRO vs MACRO



large-N quantum (field) theory

QUANTUM theory Q

Hilbert space \mathcal{H} •

Lie algebra \mathfrak{g} •

Hamiltonian operator $\hat{H} \in \mathfrak{g}$ •

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CLASSICAL theory \mathcal{C}

* manifold \mathcal{M}

* symplectic form on \mathcal{M} that defines
Poisson Brackets $\{\cdot\}_{PB}$

* Hamiltonian function $h_{cl}: \mathcal{M} \rightarrow \mathbb{R}$

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consider $\mathcal{Q}_{k \gg 0} \longrightarrow$

construct $|\Omega\rangle_k \longrightarrow$

find \mathcal{M}_k

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$\lim_{k \rightarrow 0}$

WHEN

$$\lim_{k \rightarrow 0} Q_k = \mathbb{C}$$



WHEN $\lim_{k \rightarrow 0} Q_k = \mathbb{C}$?

four conditions must hold

1 irreducibility of $G_k \rightarrow c_k \int_{\mathcal{H}} d\mu(\Omega) |\Omega\rangle\langle\Omega| = \mathbb{1}$

2 uniqueness of the "zero" operator

3 exponentially decrease of different coherent states overlap

4 classical limit of the Hamiltonian $\lim_{k \rightarrow 0} k \frac{\langle \Omega | \hat{H} | \Omega \rangle_k}{\langle \Omega | \Omega \rangle_k} < \infty$

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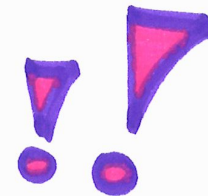
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 $\underbrace{\hspace{10em}}_{\#_k(\Omega)}$

IF 1-4 HOLD $\exists \mathcal{M}$ ON WHICH A CLASSICAL DYNAMICS CAN BE DEFINED



$$\lim_{k \rightarrow 0} k \#_k(\Omega) = h_{cl}(v, w)$$



LARGE- N LIMIT OF VECTOR MODELS

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Q_N describes a system of N spinless particles $\dot{+} [\hat{q}_i, \hat{p}_i] = \frac{1}{N} \delta_{ij} \hat{1}$
with a global symmetry whose invariants are

$$\hat{A} = \frac{1}{2} \sum_i \hat{q}_i^2$$

$$\hat{B} = \frac{1}{2} \sum_i (\hat{q}_i \hat{p}_i - \hat{p}_i \hat{q}_i)$$

$$\hat{C} = \frac{1}{2} \sum_i \hat{p}_i^2$$

with Hamiltonian

$$\hat{H}_N = N h [\hat{A}, \hat{B}, \hat{C}]$$

h an arbitrary polynomial

LARGE-N LIMIT OF VECTOR MODELS

Q_N describes a system of N spinless particles \dot{q}_i, \dot{p}_i with $[\hat{q}_i, \hat{p}_i] = \frac{1}{N} \delta_{ij} \mathbb{1}$
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identify \mathfrak{g}

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h an arbitrary polynomial

set $k = \frac{1}{N}$

identify \mathfrak{g} and the related coherent states

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recognize pseudo-spin CS

\mathcal{M} pseudo sphere PS^2

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$$\langle \Omega | \Omega' \rangle = \frac{(1-|\tau'|^2)^{NK} (1-|\tau|^2)^{NK}}{(1-\tau' \tau^*)^{2NK}}$$

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one can define a classical dynamics on PS^2

$$z = \frac{i+\tau}{i-\tau} = \rho - i\nu = \frac{k}{w} - i\nu \quad \{f, g\}_{PB} = \frac{\partial f}{\partial \nu} \frac{\partial g}{\partial w} - \frac{\partial g}{\partial \nu} \frac{\partial f}{\partial w}$$

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$$A(\nu, w) = w \quad B(\nu, w) = 2\nu w \quad C(\nu, w) = w \left(\frac{k^2}{w^2} + \nu^2 \right)$$

notice that $K^2(\nu, w) = k^2$ constant in \mathcal{M}

conserved angular momentum of the classical theory



EXAMPLE : FREE PARTICLES

$$\hat{H}_N = \frac{N}{2} \sum_i \hat{p}_i^2 = N\hat{C}$$

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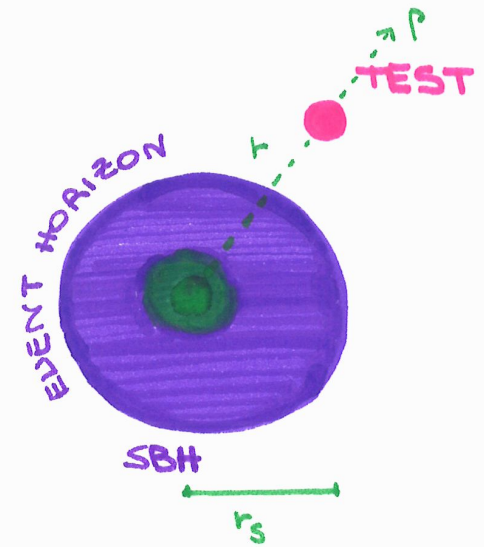
with $4k^2 = L^2$ angular momentum

CONSERVED



THE SCHWARZSCHILD BLACK HOLE

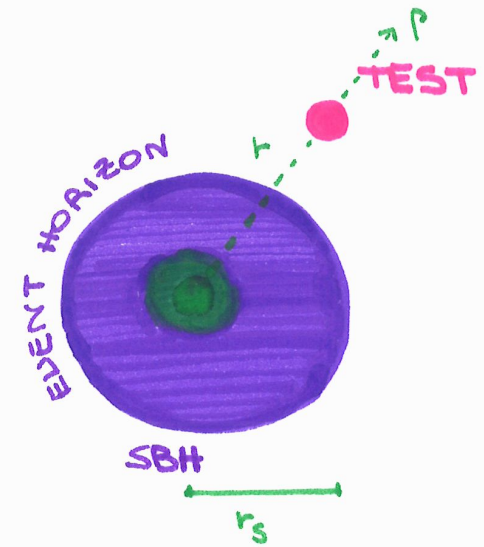
$$h_{\text{SBH}}^{\text{TEST}}(p, r) = \frac{p^2}{2} + \frac{1}{2} \left(\frac{L^2}{r^2} + \mu^2 \right) \left(1 - \frac{r_s}{r} \right)$$



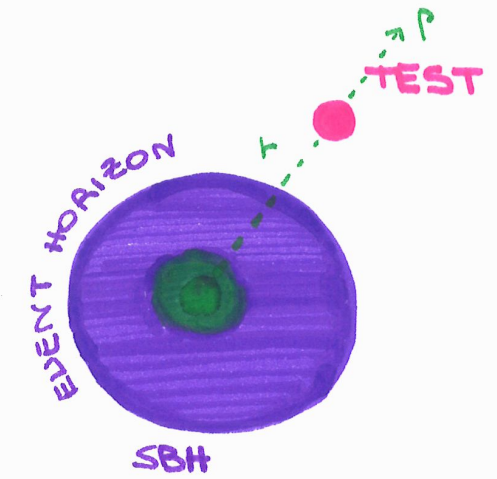
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$$h_{\text{SBH}}^{\text{TEST}}(p, r) = \frac{p^2}{2} + \frac{1}{2} \left(\frac{L^2}{r^2} + M^2 \right) \left(1 - \frac{r_s}{r} \right)$$

WHAT IS WHAT

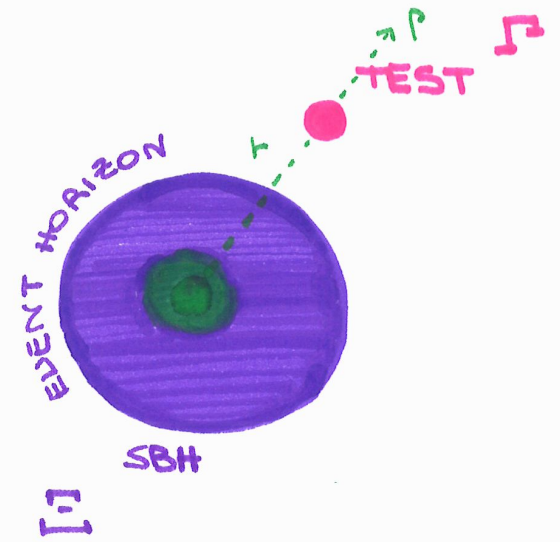


COMPOSITE SYSTEM



COMPOSITE SYSTEM

$$\Psi = \text{[musical note]} + \text{[E]}$$



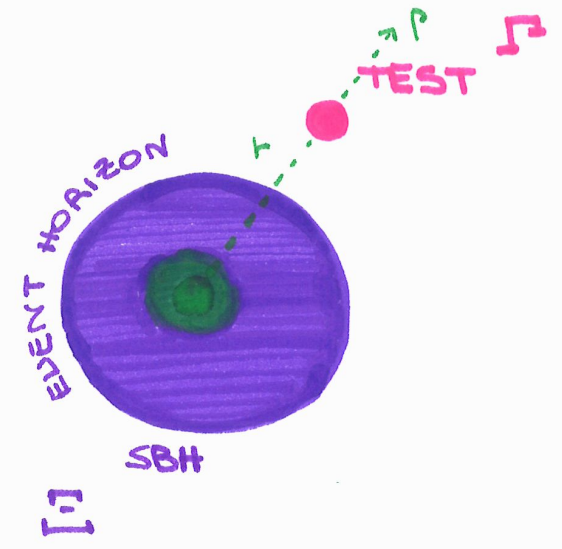
COMPOSITE SYSTEM



$$\Psi = \text{[musical note]} + \text{[E]}$$

$h_{\text{SBH}}(\rho, r)$

A red arrow points from the musical note to the text below, and a purple arrow points from the text below to the 'E' in the equation.

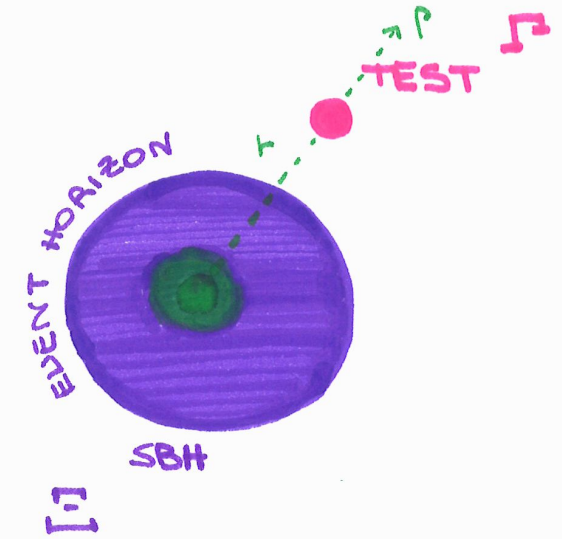


COMPOSITE SYSTEM



$$\Psi = \text{[musical note]} + \text{[E]}$$

$\hbar_{\text{SBH}}(\rho, r)$



PARAMETRIC REPRESENTATION (WITH COHERENT STATES FOR [E])

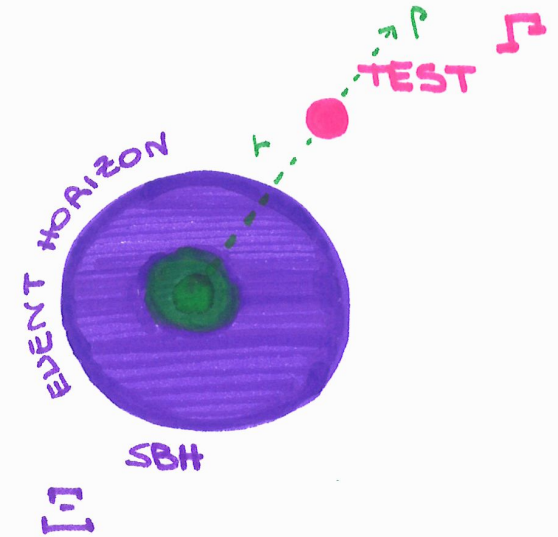
ECS

COMPOSITE SYSTEM



$$\Psi = \text{[red symbol]} + \text{[blue symbol]}$$

TEST
SBH (p,r)



PARAMETRIC REPRESENTATION (WITH COHERENT STATES FOR [blue symbol])
ECS

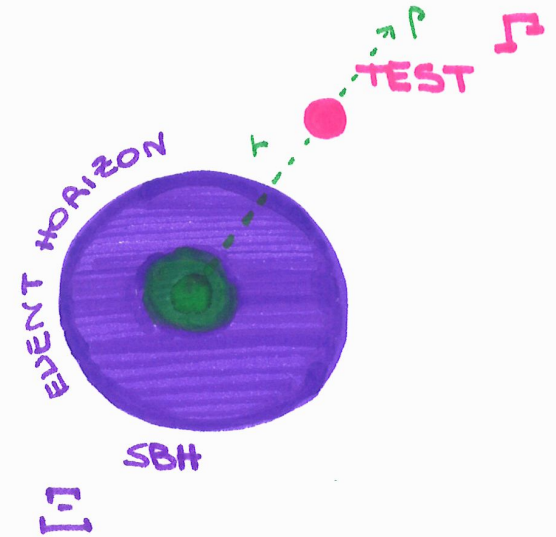
$$|\Psi\rangle = \sum_{\alpha\beta} c_{\alpha\beta} |\xi\rangle \otimes |\chi\rangle$$

COMPOSITE SYSTEM



$$\Psi = \text{[red symbol]} + \text{[blue symbol]}$$

TEST
SBH (p, r)



PARAMETRIC REPRESENTATION (WITH COHERENT STATES FOR [])
ECS

$$|\Psi\rangle = \sum_{\alpha\beta} c_{\alpha\beta} |\xi\rangle \otimes |\chi\rangle = \int d\mu(\alpha) \alpha(\alpha) |\alpha\rangle \otimes |\phi(\alpha)\rangle$$

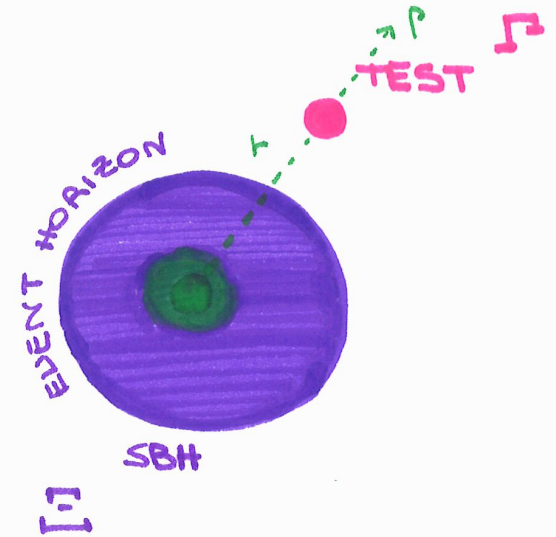
$\int d\mu(\alpha) |\alpha\rangle \otimes |\alpha\rangle$

COMPOSITE SYSTEM



$$\Psi = \mathcal{H} + \Xi$$

\mathcal{H} (red musical note) + Ξ (blue musical note)
 $h_{\text{SBH}}(\rho, r)$ (green text)
 TEST (red text)



PARAMETRIC REPRESENTATION (WITH COHERENT STATES FOR Ξ)
 ECS

$$|\Psi\rangle = \sum_{\mathcal{H}, \Xi} c_{\mathcal{H}, \Xi} |\Xi\rangle \otimes |\mathcal{H}\rangle = \int_{\mathcal{H}} d\mu(\omega) \alpha(\omega) |\omega\rangle \otimes |\phi(\omega)\rangle$$

$\int_{\mathcal{H}} d\mu(\omega) |\omega\rangle \langle \omega|$ (under the sum)

$$\int_{\mathcal{H}} d\mu(\omega) \alpha^2(\omega) = 1$$

$$\langle \phi(\omega) | \phi(\omega) \rangle = 1 \quad \forall \omega \in \mathcal{H}$$

HOW CAN AN ENVIRONMENT
SHAPE THE SPECTRUM OF
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HOW CAN AN ENVIRONMENT SHAPE THE SPECTRUM OF A QUANTUM PARTICLE



- TIMELESSNESS $\hat{H}_\psi |\psi\rangle = 0$

- NON-INTERACTING PARTITION $\hat{H}_\psi = \hat{H}_\equiv \otimes \hat{1}_\Omega + \hat{1}_\equiv \otimes \hat{H}_\Omega$

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THESE ASSUMPTIONS ARE THE SAME UPON WHICH THE

PaW MECHANISM

and recent updates are based

with \equiv the clock for Ω



● TIMELESSNESS

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$$\langle \bar{\Omega} | \hat{H}_\psi | \psi \rangle = 0$$

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LARGE-N \equiv

$$N \rightarrow \infty$$

$$\int_{\mathcal{H}} d\mu(\Omega) \chi(\Omega) \delta(\Omega - \bar{\Omega}) \left[\hat{H}_\equiv(\Omega, \bar{\Omega}) |\phi(\Omega)\rangle - \hat{H}_\Omega |\phi(\Omega)\rangle \right] = 0$$

$$\hat{H}_\Omega |\phi(\Omega)\rangle = \hat{H}_\equiv(\Omega) |\phi(\Omega)\rangle$$

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$$\hat{H}_\Omega |\phi(\Omega)\rangle = \hat{H}_\equiv(\Omega) |\phi(\Omega)\rangle$$

COULD THIS BE

$$w_{\equiv}^{\Omega}(\Omega) = w_{\text{SBH}}^{\text{TEST}}(p, r)$$



SBH AS MACROSCOPIC QUANTUM SYSTEMS

SBH AS MACROSCOPIC QUANTUM SYSTEMS



SBH AS MACROSCOPIC QUANTUM SYSTEMS



$$h_{ce}(v, w) = h_{SBH}^{TEST}(p, r)$$

SBH AS MACROSCOPIC QUANTUM SYSTEMS

? Q_N ?

$$h_{cl}(r, w) = h_{SBH}^{TEST}(p, r)$$

$$\lim_{N \rightarrow \infty} \frac{1}{N} \frac{\langle \Omega | \hat{H}_N | \Omega \rangle}{\langle \Omega | \Omega \rangle} = \frac{p^2}{2} + \frac{1}{2} \frac{L^2}{r^2} - \mu^2 \frac{r_s}{2r} - \frac{L^2 r_s}{2r^3} + \frac{\mu^2}{2}$$

SBH AS MACROSCOPIC QUANTUM SYSTEMS



$$h_{cl}(r, w) = h_{SBH}^{TEST}(p, r)$$

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$$g = SU(1, 1)$$

SBH AS MACROSCOPIC QUANTUM SYSTEMS

? Q_N ?

$$h_{cl}(r, w) = h_{SBH}^{TEST}(p, r)$$

$$\lim_{N \rightarrow \infty} \frac{1}{N} \frac{\langle \Omega | \hat{H}_\pm | \Omega \rangle}{\langle \Omega | \Omega \rangle} = \frac{p^2}{2} + \frac{1}{2} \frac{L^2}{r^2} - \mu^2 \frac{r_s}{2r} - \frac{L^2 r_s}{2r^3} + \frac{\mu^2}{2}$$



$\mathfrak{g} = SU(1,1) + \text{HAWKING RADIATION}$

TWO-MODE REALIZATION OF $SU(1,1)$ ALGEBRA



TWO-MODE REALIZATION OF SU(1,1) ALGEBRA

$$\hat{a} = \sum_i \hat{a}_i \quad \hat{b} = \sum_i \hat{b}_i$$

$$[\hat{a}_i, \hat{a}_j^+] = \frac{1}{N^2} \delta_{ij} \hat{1} = [\hat{b}_i, \hat{b}_j^+] \quad [\hat{a}_i, \hat{b}_j^{(+)}] = 0 \quad i, j : 1, 2, \dots, N$$

$$[\hat{a}, \hat{a}^+] = [\hat{b}, \hat{b}^+] = \frac{1}{N} \hat{1} \quad [\hat{a}, \hat{b}^{(+)}] = 0$$

$$\text{SU}(1,1) \quad \left\{ \begin{array}{l} \hat{K}_+ = \hat{a}^+ \hat{b}^+ \\ \hat{K}_- = \hat{a} \hat{b} \\ \hat{K}_0 = \frac{1}{2} (\hat{a}^+ \hat{a} + \hat{b} \hat{b}^+) \end{array} \right. \quad \hat{K}_{\pm} = \hat{K}_1 \pm i \hat{K}_2$$

$$2\hat{A} = \hat{N}_a + \hat{N}_b - i(\hat{a}^+ \hat{b}^+ - \hat{a} \hat{b}) \quad 2\hat{C} = \hat{N}_a + \hat{N}_b + i(\hat{a}^+ \hat{b}^+ - \hat{a} \hat{b})$$


$$\hat{B} = \hat{a}^+ \hat{b}^+ + \hat{a} \hat{b}$$

$$\hat{H}_{\equiv} = N [\hat{C} + \alpha \hat{A} + \beta \hat{B} + \epsilon \hat{1}]$$

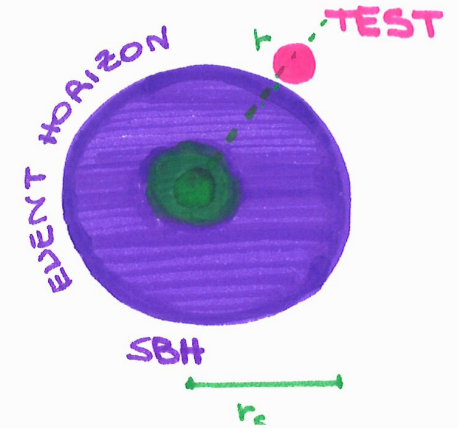


$$\hat{H}_{\text{eff}} = N [\hat{C} + \alpha \hat{A} + \beta \hat{B} + \epsilon \hat{1}]$$




$$h_{\text{cl}}(p, r) = \frac{p^2}{2} + \frac{L^2}{2r^2} + \frac{\alpha}{2} r^2 + \beta p r + \epsilon$$

$$\hat{H}_{\equiv} = N [\hat{C} + \alpha \hat{A} + \beta \hat{B} + \epsilon \hat{1}]$$



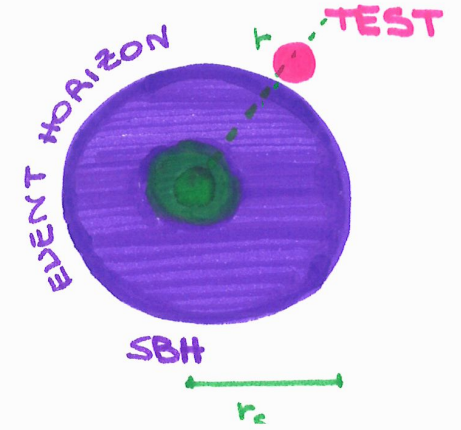
$$\longrightarrow h_{cl}(p, r) = \frac{p^2}{2} + \frac{L^2}{2r^2} + \frac{\alpha}{2} r^2 + \beta p r + \epsilon$$

$$= h_{SBH}^{TEST}(p, r - r_s \ll r_s) \sim \frac{p^2}{2} + \frac{1}{2r_s^2} \left(\frac{L^2}{r_s^2} + p^2 \right) \delta + o\left(\frac{\delta^2}{r_s^2}\right)$$

NEAR-HORIZON $\delta := r - r_s \ll r_s$

NEAR HORIZON DYNAMICS

$$r - r_s \ll r_s$$



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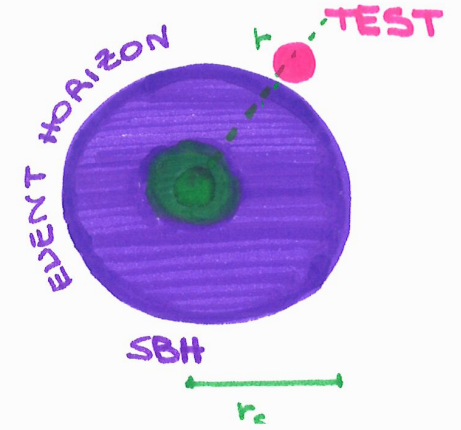
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NEAR HORIZON DYNAMICS

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free quadratic term

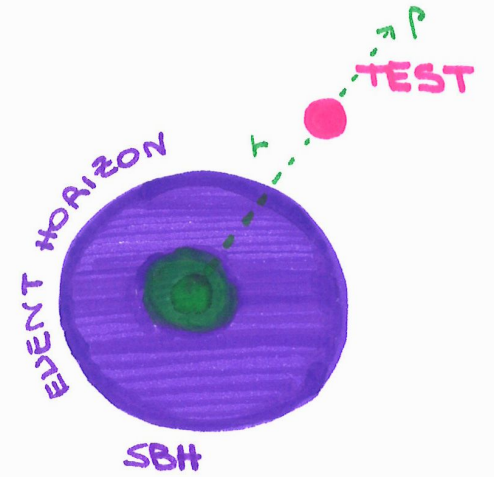


pairs creation-annihilation term
(Hawking?)

$$\hat{H}_\Omega |\phi(\Omega)\rangle = H_\Omega(\Omega) |\phi(\Omega)\rangle$$

COULD THIS BE

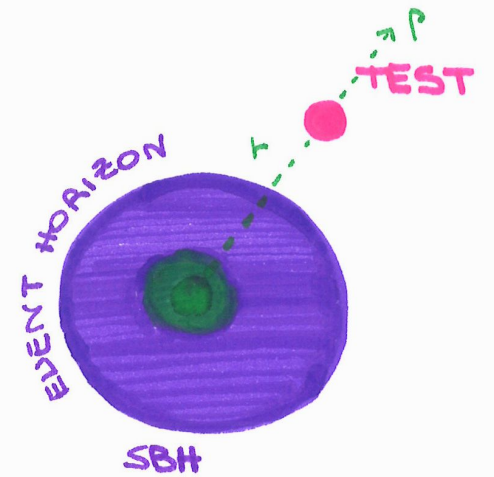
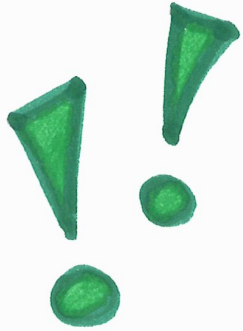
$$h_{\Omega}^{\Omega} = h_{\text{SBH}}^{\text{TEST}}(p, r)$$



$$\hat{H}_\Omega |\phi(\Omega)\rangle = H_\Xi(\Omega) |\phi(\Omega)\rangle$$

THE MACROSCOPIC Ξ
IS THE SBH FOR THE

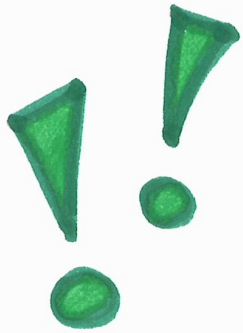
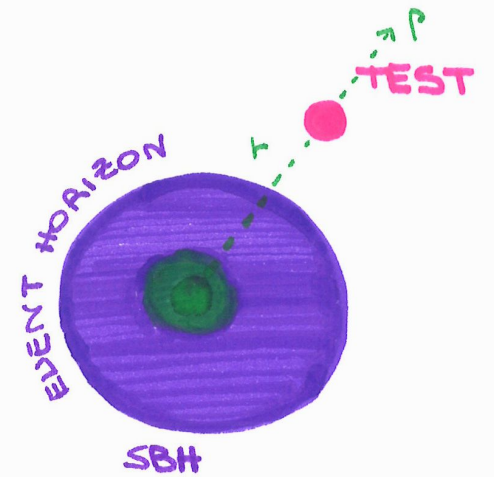
TEST PARTICLE μ



$$\hat{H}_\Omega |\phi(\Omega)\rangle = H_{\Xi}(\Omega) |\phi(\Omega)\rangle$$

THE MACROSCOPIC Ξ
IS THE SBH FOR THE

TEST PARTICLE μ



Its presence is reflected into the behaviour
test particle μ via the parametric dependence of its spectrum
on the couple (p,r) provided by Ξ

(remember that $(p,r) \rightarrow \Omega \in \mathcal{M} \rightarrow$ coherent states for Ξ)

QUANTUM MECHANICS vs GENERAL RELATIVITY

EVOLUTION IN TIME vs SYMMETRY TRANSFORMATION

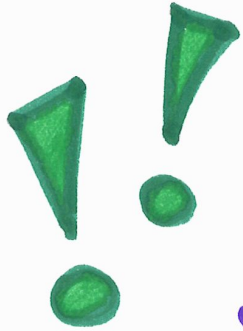
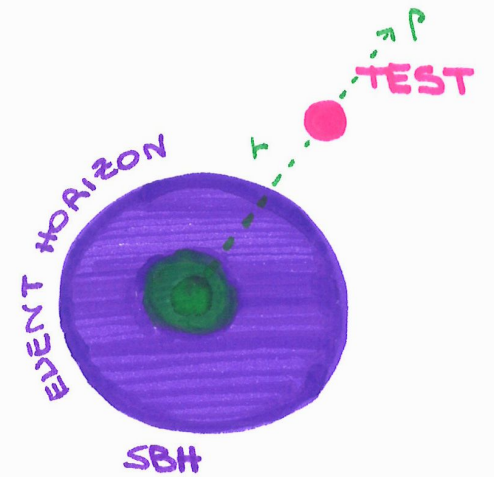


Page & Wootters mechanism

$$\hat{H}_\Omega |\phi(\Omega)\rangle = H_\Xi(\Omega) |\phi(\Omega)\rangle$$

THE MACROSCOPIC Ξ
IS THE SBH FOR THE

TEST PARTICLE μ

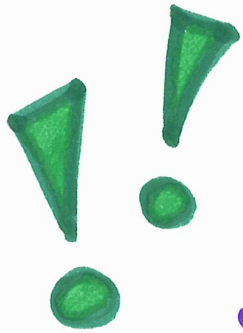
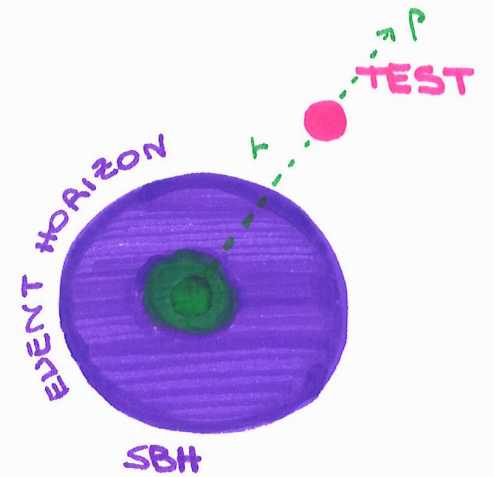


- TIMELESSNESS
- NON-INTERACTING PARTITION
- ENTANGLEMENT

$$\hat{H}_\Omega |\phi(\Omega)\rangle = H_\Xi(\Omega) |\phi(\Omega)\rangle$$

THE MACROSCOPIC Ξ
IS THE SBH FOR THE

TEST PARTICLE Ω



- TIMELESSNESS
- NON-INTERACTING PARTITION
- ENTANGLEMENT

THESE ASSUMPTIONS ARE THE SAME UPON WHICH THE

PaW MECHANISM

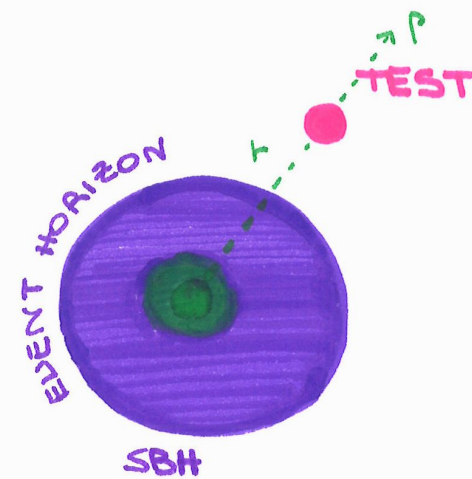
and recent updates are based

with Ξ the clock for Ω



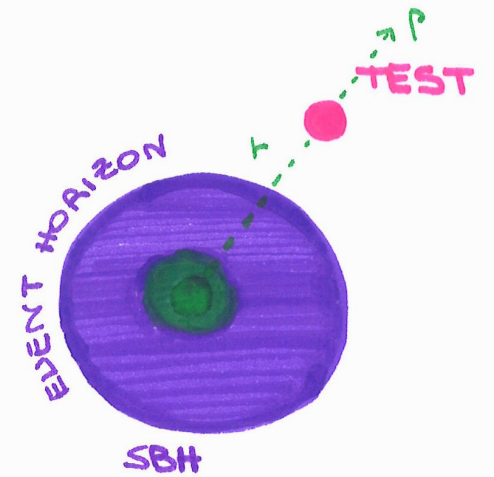


IS THE SBH





Ξ IS THE SBH

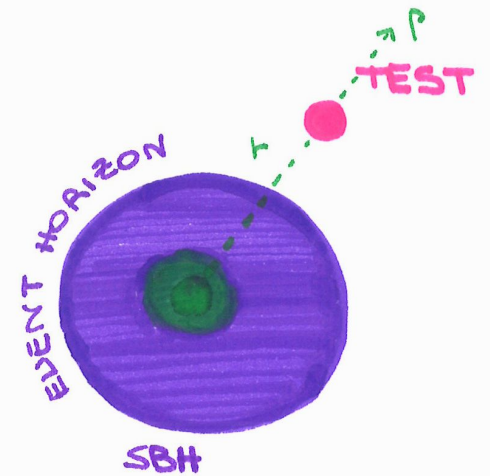


According to the PAW mechanism

Ξ IS A CLOCK FOR Ω

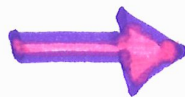


Ξ IS THE SBH



According to the PAW mechanism

Ξ IS A CLOCK FOR Ω



THE SBH IS THE
CLOCK

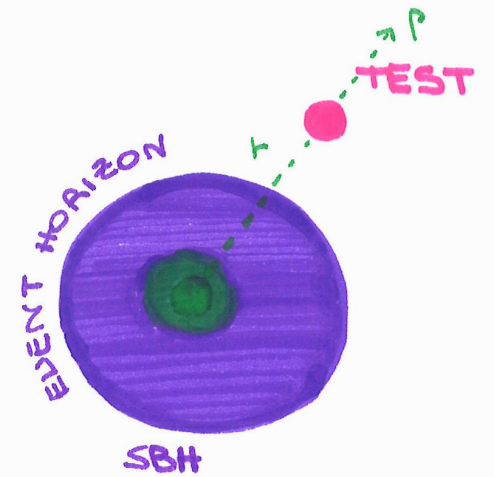
FOR WHAT GOES AROUND IT



THANKS

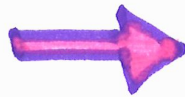


Ξ IS THE SBH



According to the PAW mechanism

Ξ IS A CLOCK FOR Ω



THE SBH IS THE

CLOCK

FOR WHAT GOES AROUND IT

