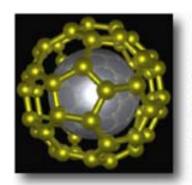
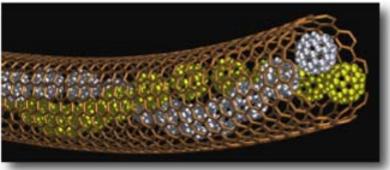


Outstanding problems in using spin states for practical quantum information science



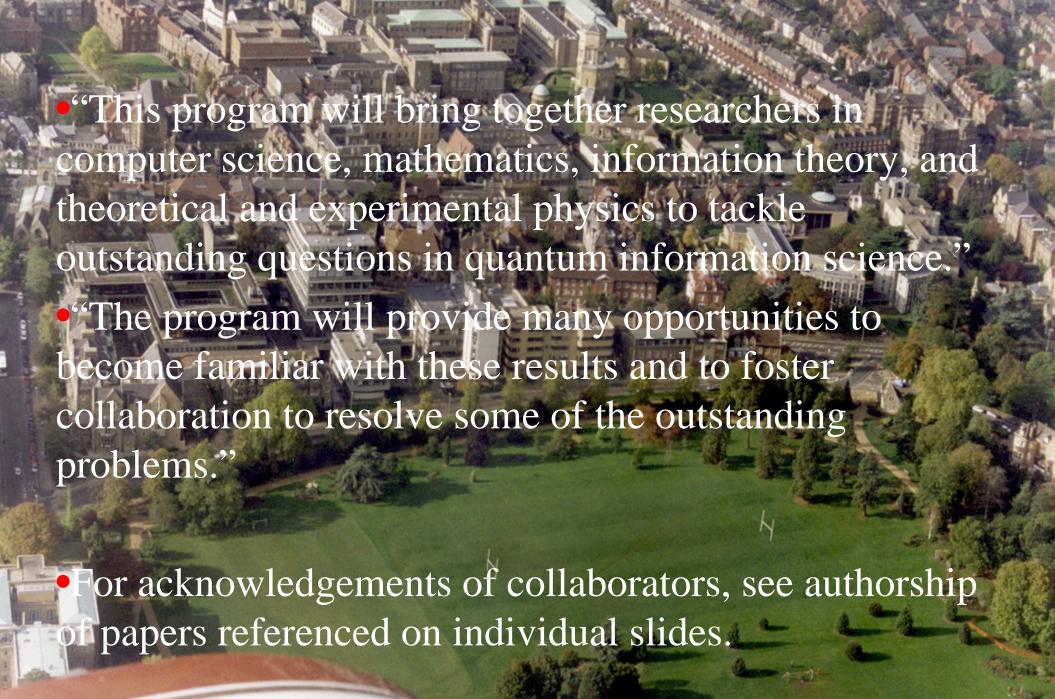


KITP QIS

www.qipirc.org



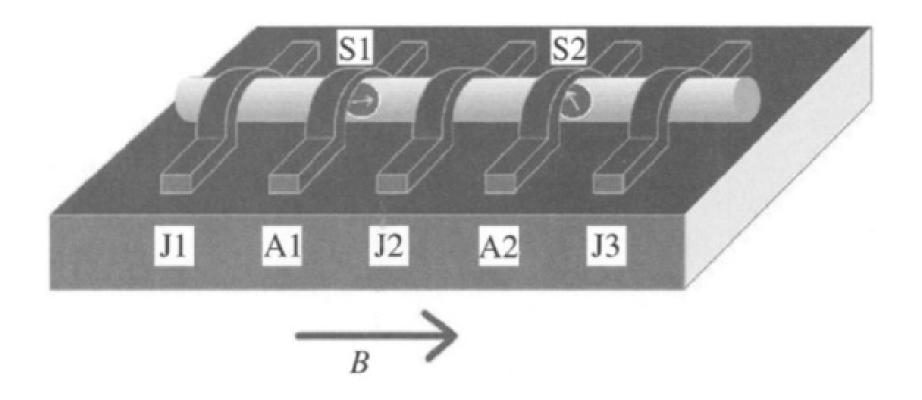








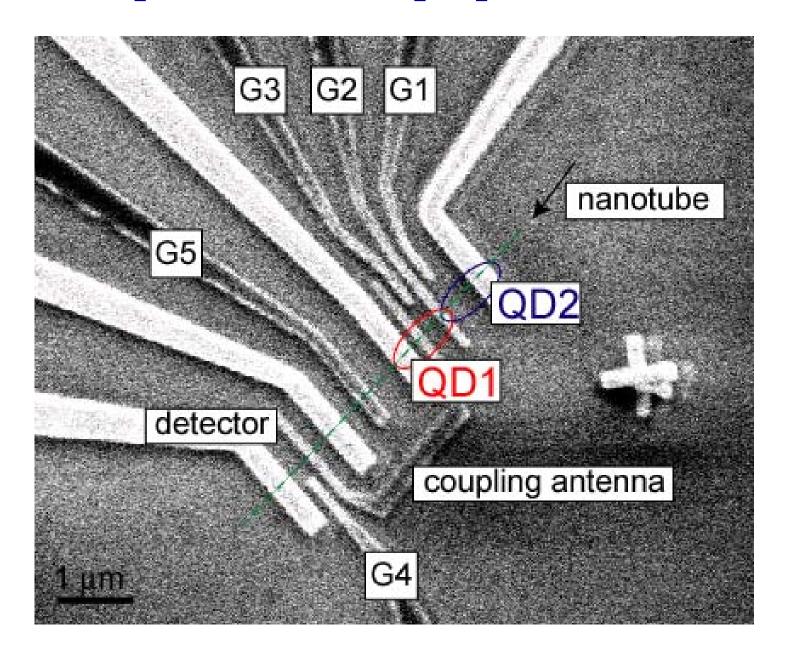
Nanoscale solid-state quantum computing



A. Ardavan *et al.*, *Phil. Trans. R. Soc. Lond.* A **361**, 1473-1485 (2003)



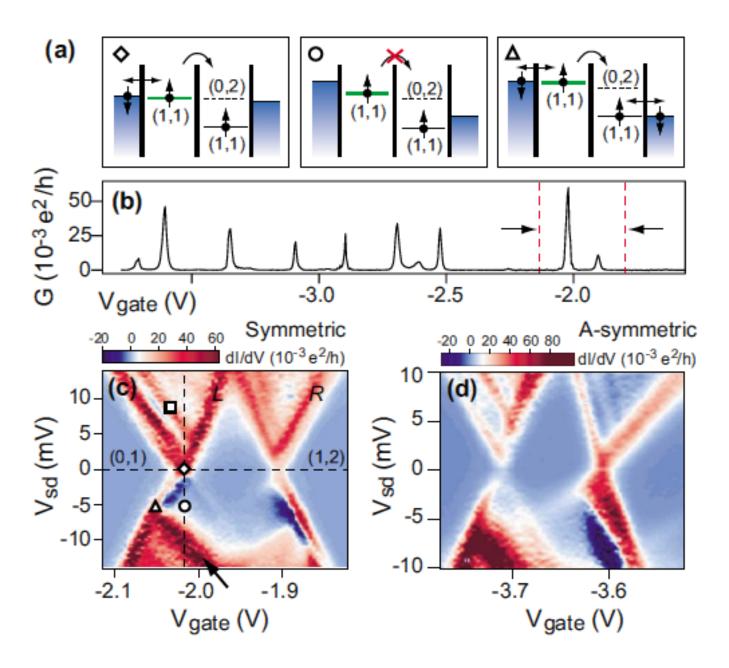
Spin blockade in peapod devices



M.R. Buitelaar et al., Phys. Rev. B 77, 245439 (2008)



Spin blockade in peapod devices



M.R. Buitelaar et al., Phys. Rev. B 77, 245439 (2008)

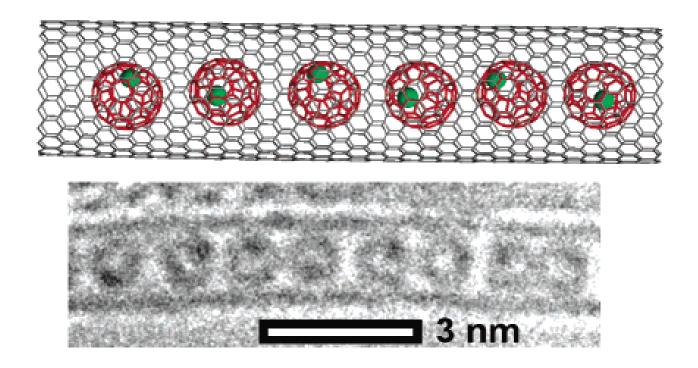


Outstanding problems in using spin states

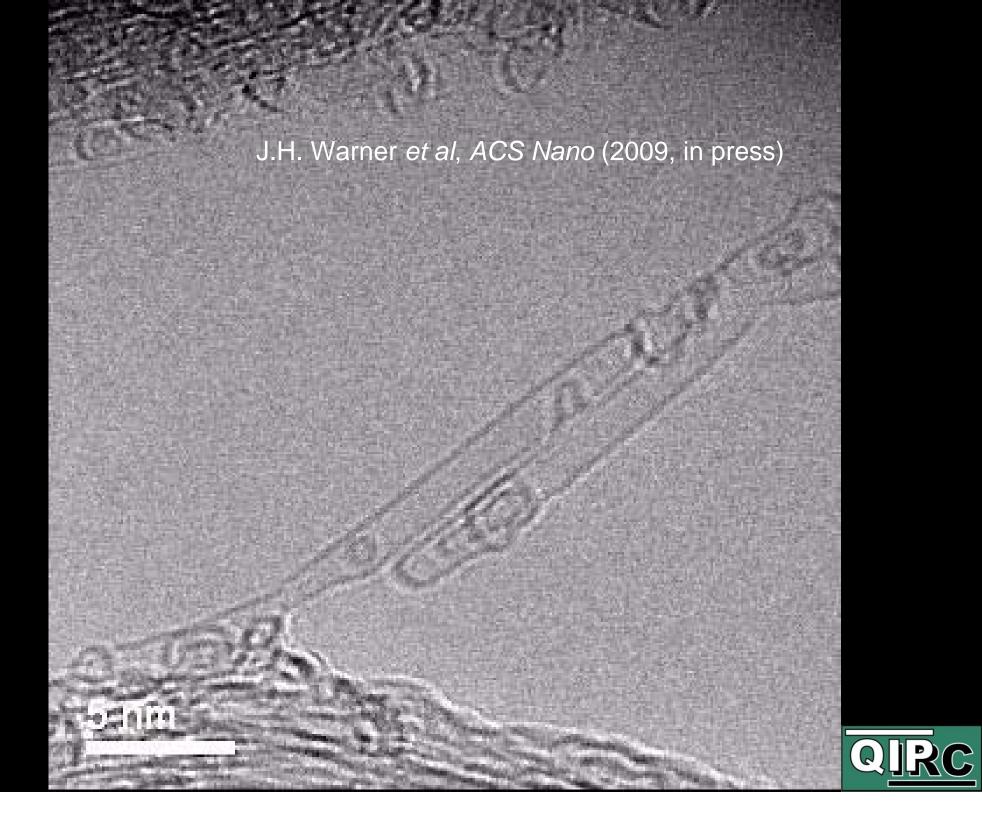
- In peapods:
 - how can you measure single spins?



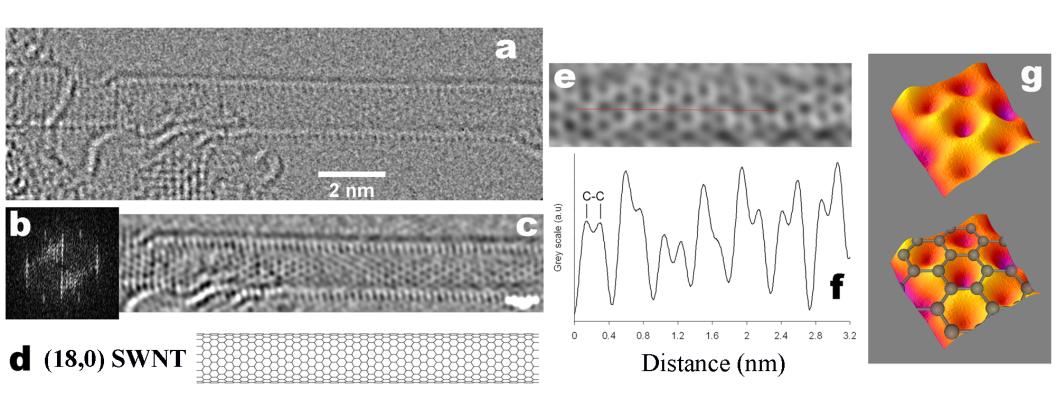
Ce@C₈₂ endohedral fullerenes in nanotubes



A.N. Khlobystov *et al.*, *Accounts of Chemical Research* **38**, 901-909 (2005)



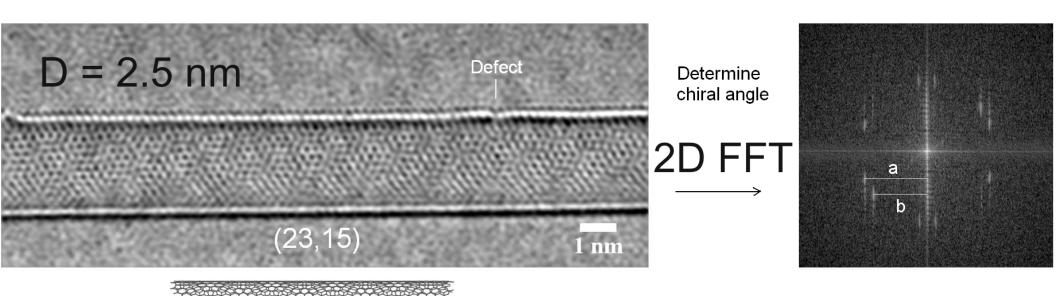
Aberration corrected HRTEM of SWNT



J.H. Warner *et al*, *ACS Nano* **3**, 1557-1563 (2009)



Aberration corrected HRTEM of SWNT



Atomic structural model (23, 15)

J.H. Warner et al, ACS Nano 3, 1557-1563 (2009)

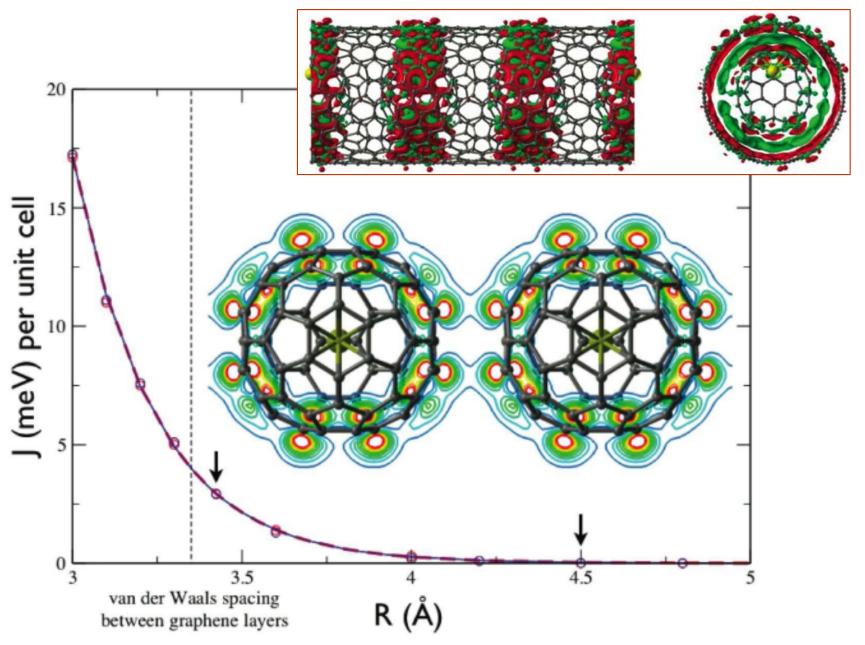


Outstanding problems in using spin states

- In peapods:
 - how can you measure single spins?
 - what will be the EDMR mechanisms?



Modelling charge transfer in Sc@C₈₂@SWNT



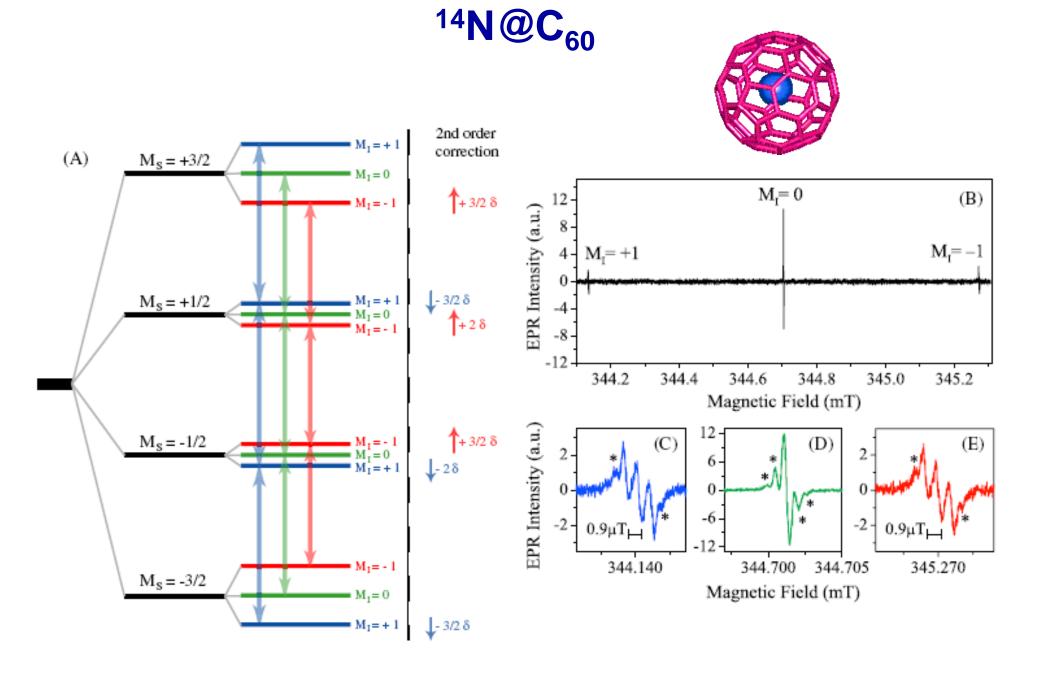
L. Ge et al., Phys. Rev. B 77, 235416 (2008)



Outstanding problems in using spin states

- In peapods:
 - how can you measure single spins?
 - what will be the EDMR mechanisms?
 - can you use RKKY for controlled coupling?

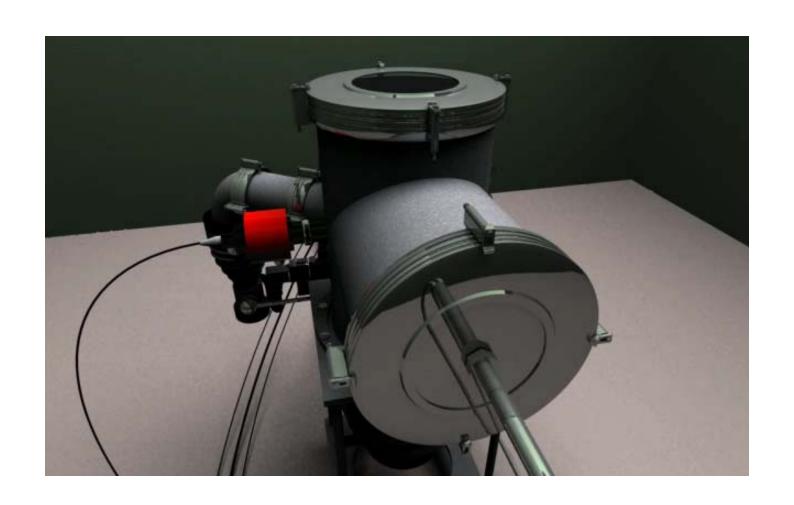




S.C. Benjamin et al., J. Phys.: Condens. Matter 18, S867–S883 (2006)



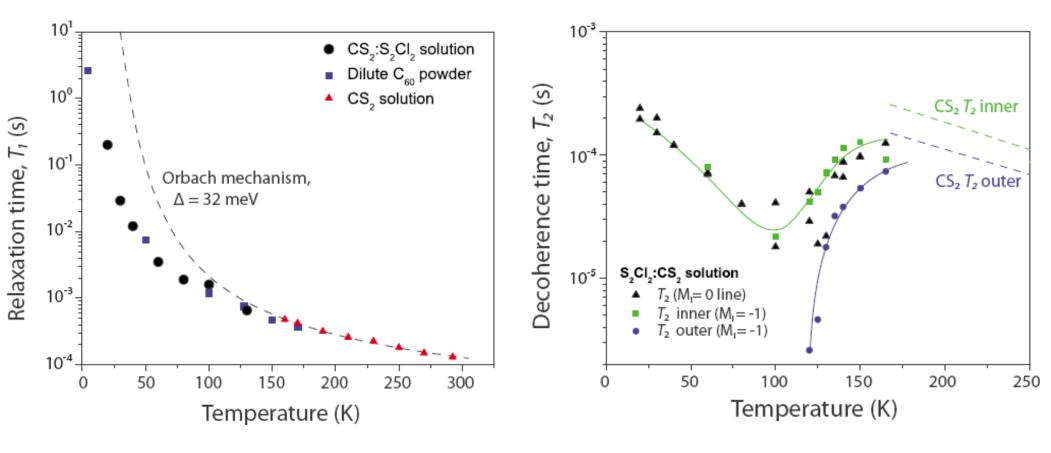








Spin relaxation and decoherence times of ¹⁴N@C₆₀

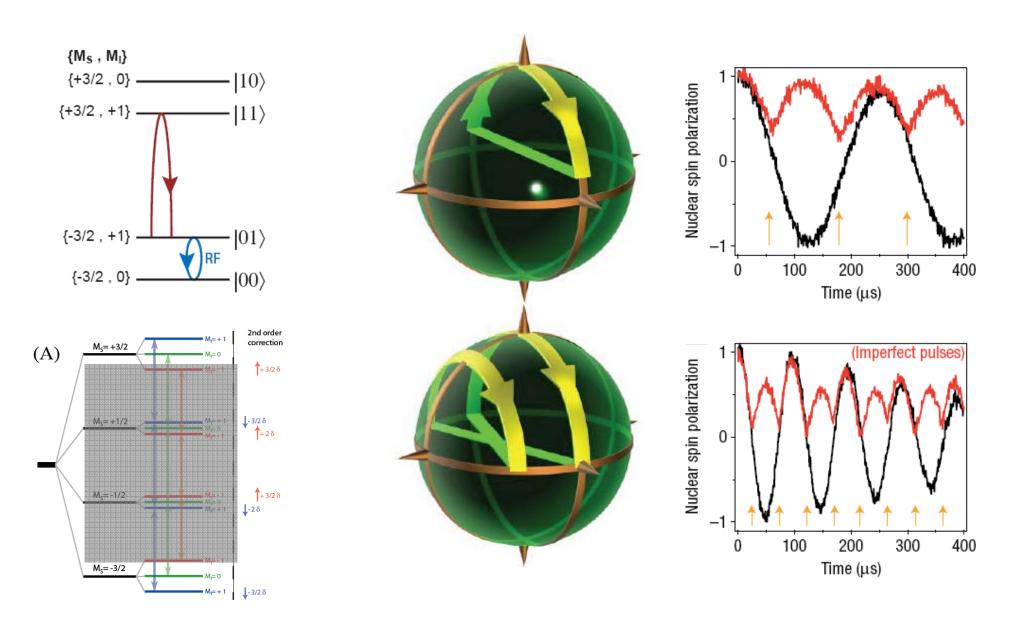


- T₁, above 150 K, follows Orbach mechanism
 - Additional relaxation mechanisms at lower temperature
- T_{2inner} increases below 100 K, and is unaffected by the ¹⁴N spin
 - Fall in T_{2outer} is due to dephasing as a result of zero field splitting

J.J.L. Morton et al., Phys. Rev. B 76, 085418 (2007)



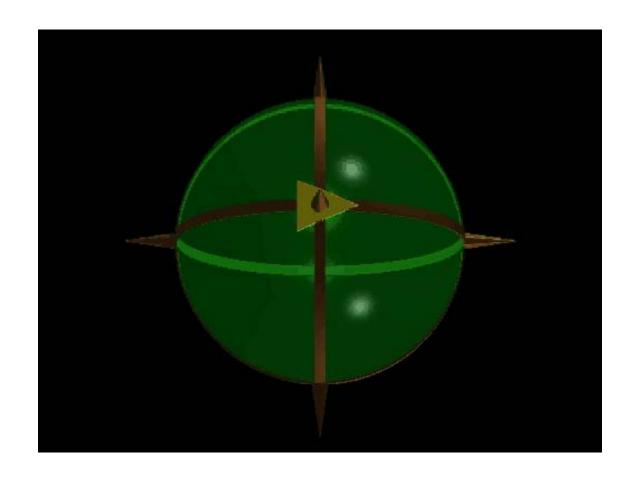
Bang-bang control



J.J.L. Morton *et al.*, *Nature Physics* **2**, 40-43 (2006)



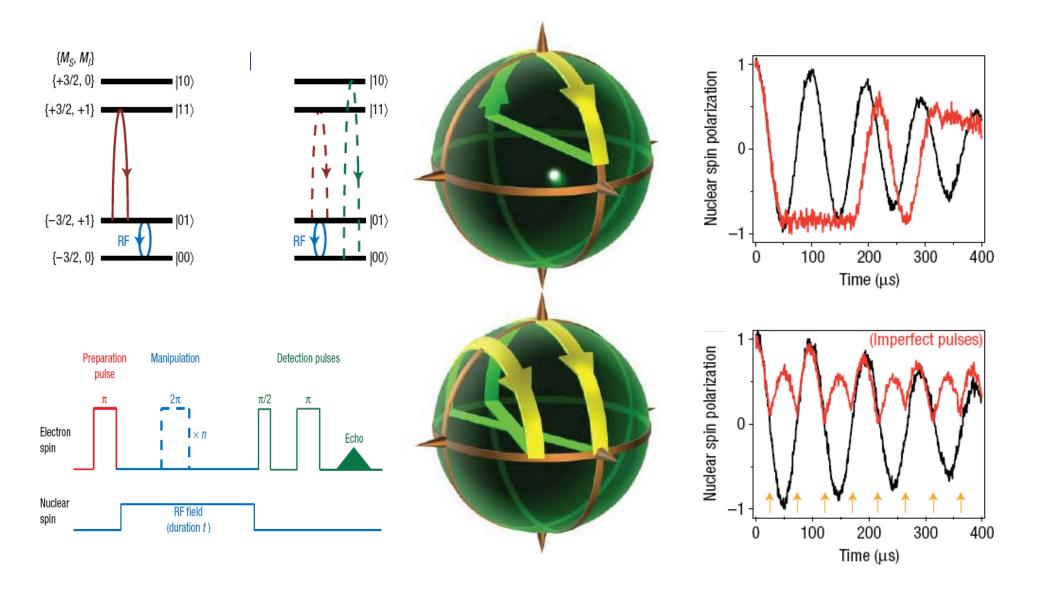
Bang-bang control



J.J.L. Morton *et al.*, *Nature Physics* **2**, 40-43 (2006)



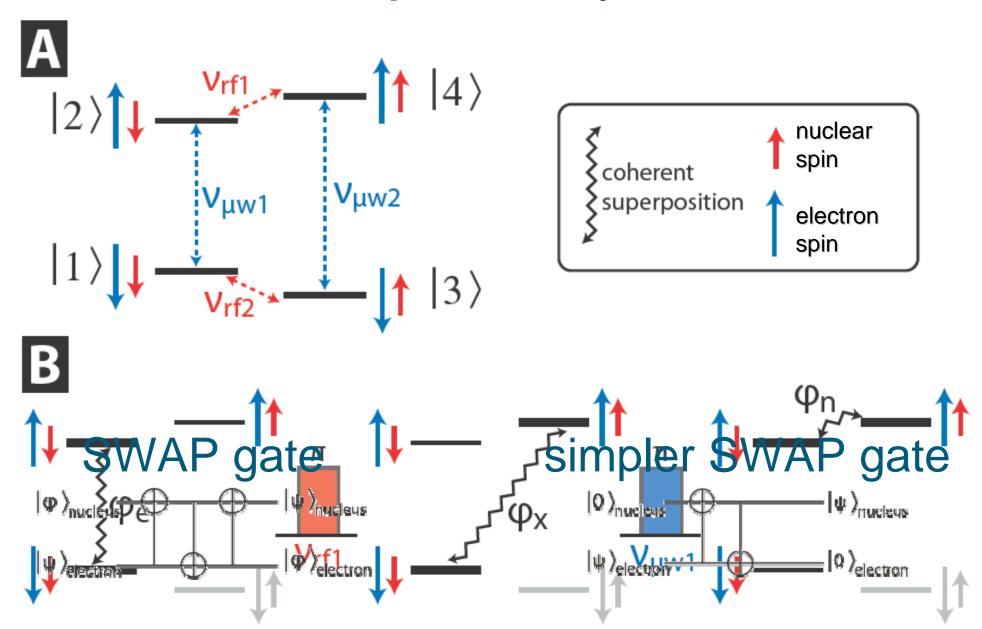
Bang-bang control



J.J.L. Morton *et al.*, *Nature Physics* **2**, 40-43 (2006)



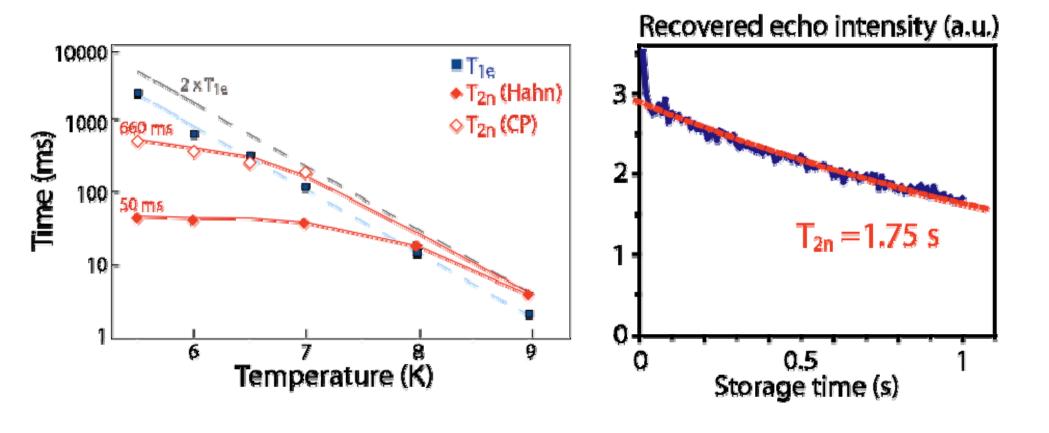
Spin memory



J.J.L. Morton et al., Nature 455, 1085-1088 (2008)



31P in 28Si

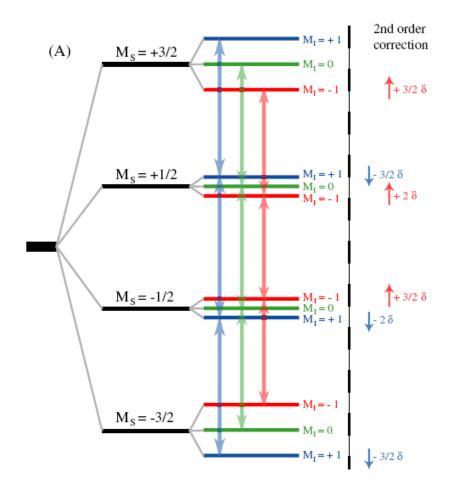


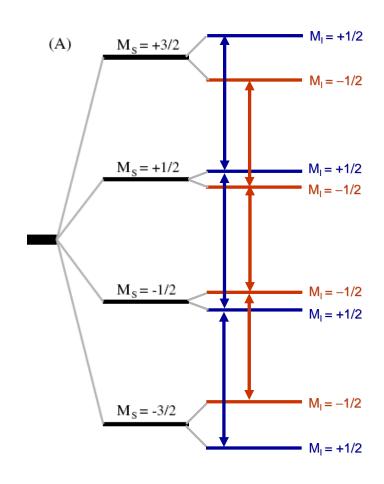
J.J.L. Morton et al., Nature 455, 1085-1088 (2008)

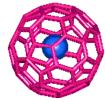


15N@**D**@C₆₀

$^{15}N@C_{60}$

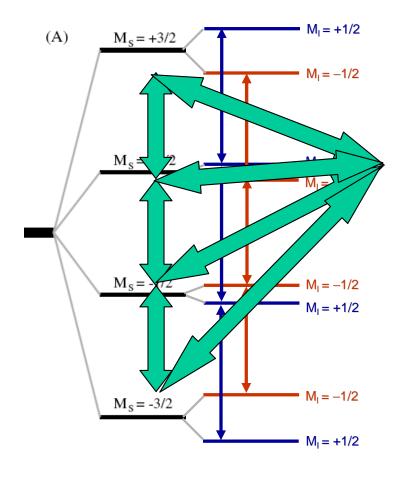




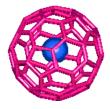




¹⁵N@C₆₀



- 1. $\pi/2_{MW}$ to create coherence
- 2. π_{RF} to flip nuclear spin
- 3. π_{MW} to flip electron spin
- 4. Outer coherences rapidly decay
- 5. Wait several T_{2e}
- 6. π_{MW} to flip electron spin
- 7. π_{RF} to flip nuclear spin
- 8. Original inner coherence of electron spin is now restored



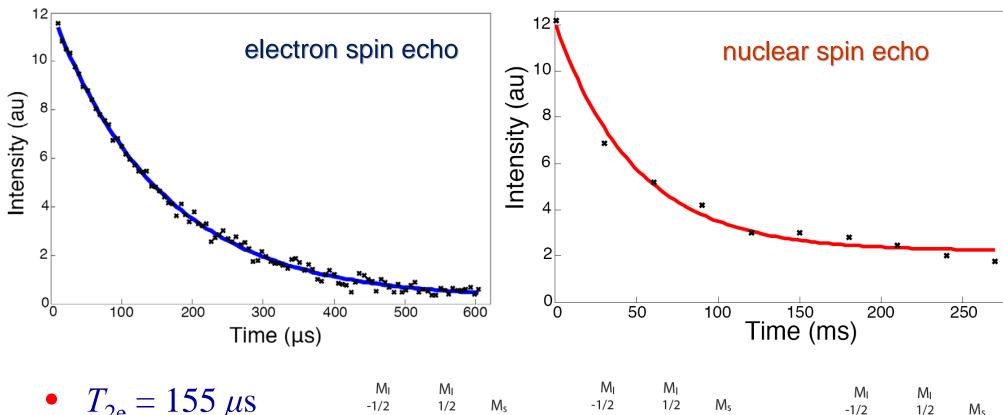


Stored and recovered signal 3500 Initial echo 3000 2500 Intensity (au 2000 1500 Recovered echo 1000 500 π π π 0 refocus transfer to N transfer to E -500 10 10 20 Time (µs) Time (ms) Time (µs)

- Phase cycled m.w. $\pm -x$ and r.f. 0/90
 - Evidence that information is stored and retrieved from nuclear spin
- Two-way transfer fidelity $f^2 = 0.44$
 - Known limitations to fidelity are entirely instrumental



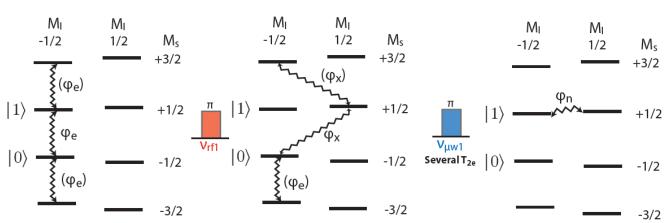
Storage time @ 10 K



•
$$T_{2e} = 155 \,\mu s$$

•
$$T_{2n} = 54 \text{ ms}$$





Richard Brown



Quantum sensing

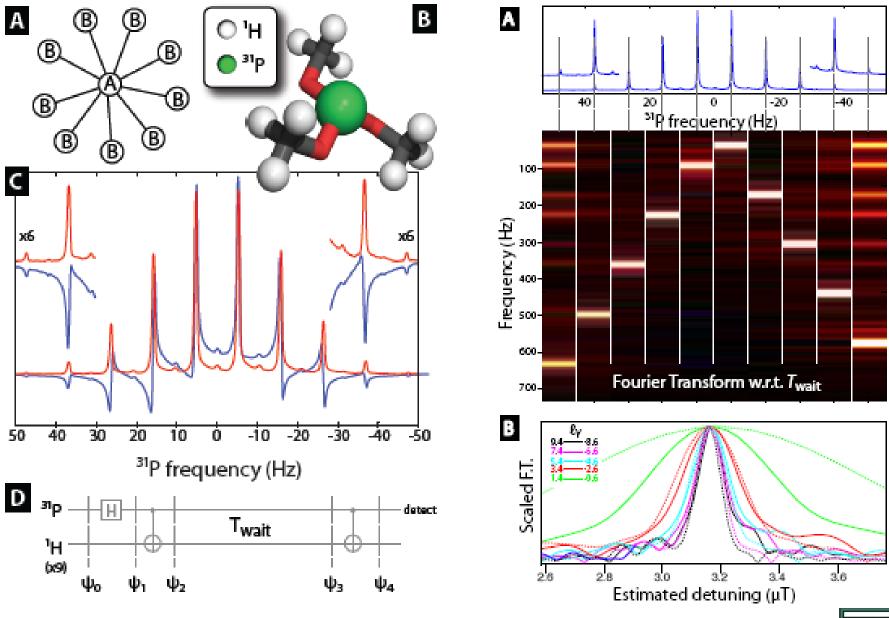
$$|\Psi\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) \qquad |\Psi\rangle = \frac{1}{\sqrt{2}} (|0000\rangle + |1111\rangle)$$

$$|\Psi\rangle = \frac{1}{\sqrt{2}} (|0\rangle + e^{i\phi} |1\rangle) \qquad |\Psi\rangle = \frac{1}{\sqrt{2}} (|0000\rangle + e^{iN\phi} |1111\rangle)$$

- Already demonstrated in optics
- Spin entanglement fully deterministic
- Enhancement by \sqrt{N} in sensitivity to field
- New generation of sensors



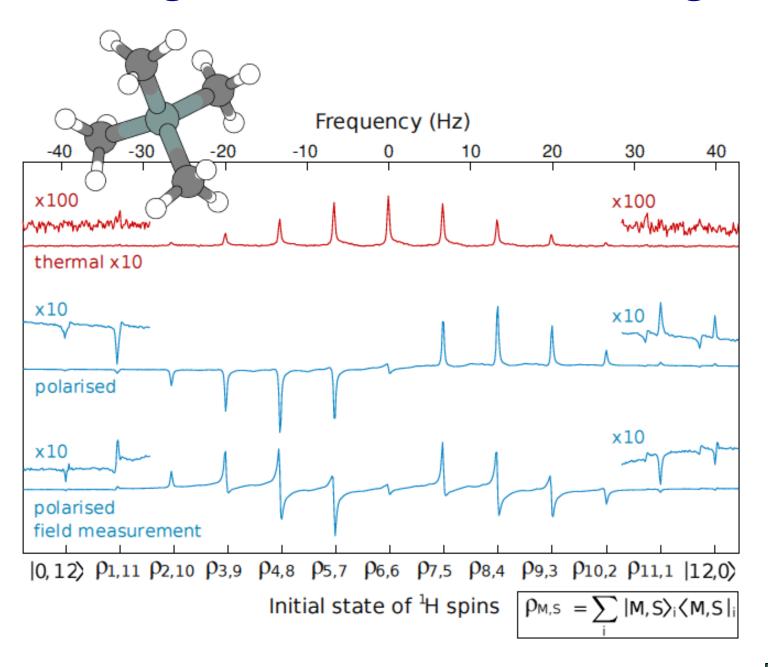
Entanglement enhanced field sensing



J.A. Jones et al., Science **324**, 1166-1168 (2009)



Entanglement enhanced field sensing

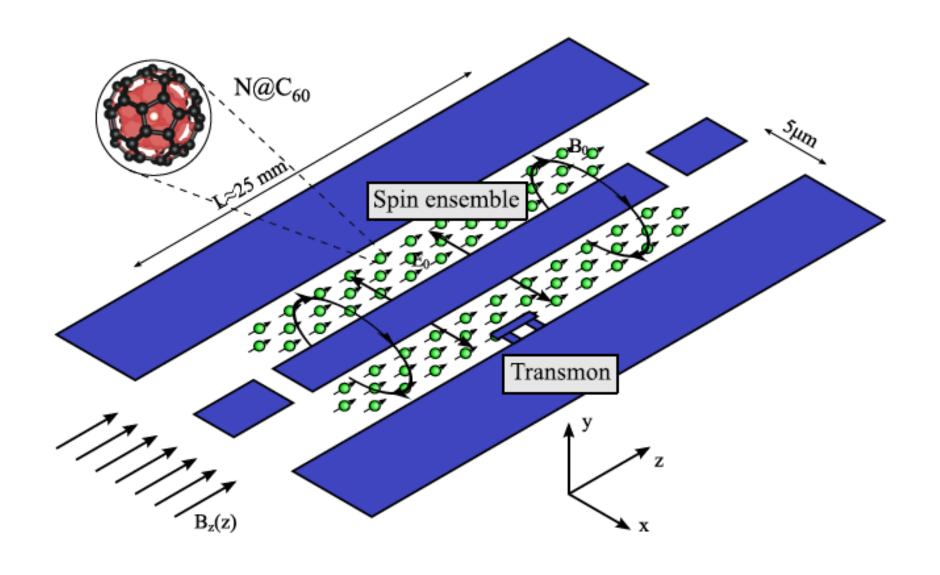




Outstanding problems in using spin states

- In peapods:
 - how can you measure single spins?
 - what will be the EDMR mechanisms?
 - can you use RKKY for controlled coupling?
- In entanglement enhanced sensing:
 - can you achieve a pure state?





J.H. Wesenberg et al., Phys. Rev. Lett. 103, 070502 (2009)



JOURNAL OF APPLIED PHYSICS

VOLUMB 26, NUMBER 11

NOVEMBER, 1955

Spin Echo Serial Storage Memory

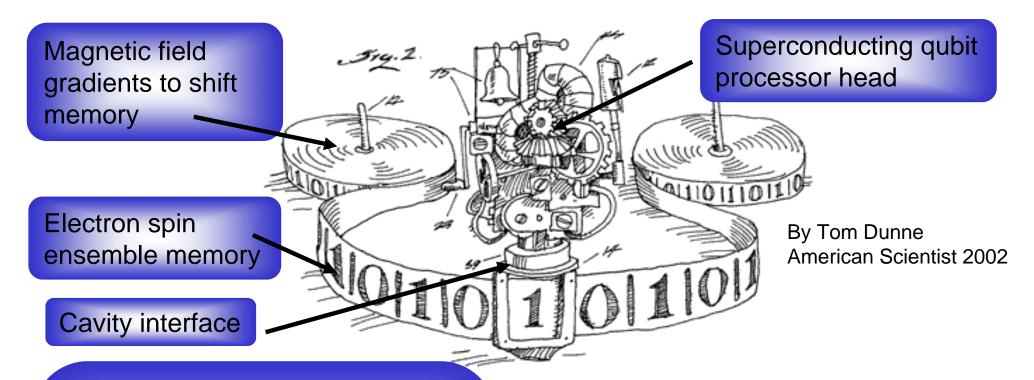
A. G. Anderson, R. L. Garwin, E. L. Hahn, J. W. Horron, G. L. Tucker, † and R. M. Walker.

Watson Scientific Computing Laboratory, Columbia University, New York, New York

(Received July 26, 1955)

By utilizing the method of pulsed nuclear magnetic resonance, radio-frequency energy in the form of pulses can be stored serially in a sample of nuclear spins and recalled at an arbitrary later time within the memory or relaxation time of the spin sample. Weak pulses of radio-frequency energy condition the nuclear suins to start precessing in phase. After they become completely out of phase, a strong recollection pulse brings about a phase reversal of precession and produces a series of spin echoes in a sequence corresponding to direct or reverse order of input pulses. The echo amplitudes in such a series are given as a function of the number and strength of the input pulses and the conditions for maximum storage capacity in a spin ensemble are determined. The maximum specific storage capacity in liquids is expressed in terms of the thermal noise of the detecting apparatus, the effect of self-diffusion of the molecules, and the relaxation times. The origin of undesired spin echoes arising from the interaction of input pulses is discussed, and means for climinating these echoes by frequency and magnetic field modulation are discussed and applied. Extensive use is made of a magnetic field modulation technique to destroy undesired echoes, and to permit novel types of recall of serially stored groups of pulses. Whereas Fernbach and Proctor [J. Appl. Phys. 26, 170 (1955)] have demonstrated multiple pulse storage under conditions which reproduce the input pulse shape, the present investigation is concerned with the storage of a maximum number of pulses whose shape is ideally determined by the nuclear spin band width. In practice, the order of 1000 rf pulses can be stored and recalled by this method in a proton sample several or in volume within a memory time of 10 to 50 milliseconds. Large specific storage capacities expected for existing long relaxation time liquids are not realized because of excessive self-diffusion.





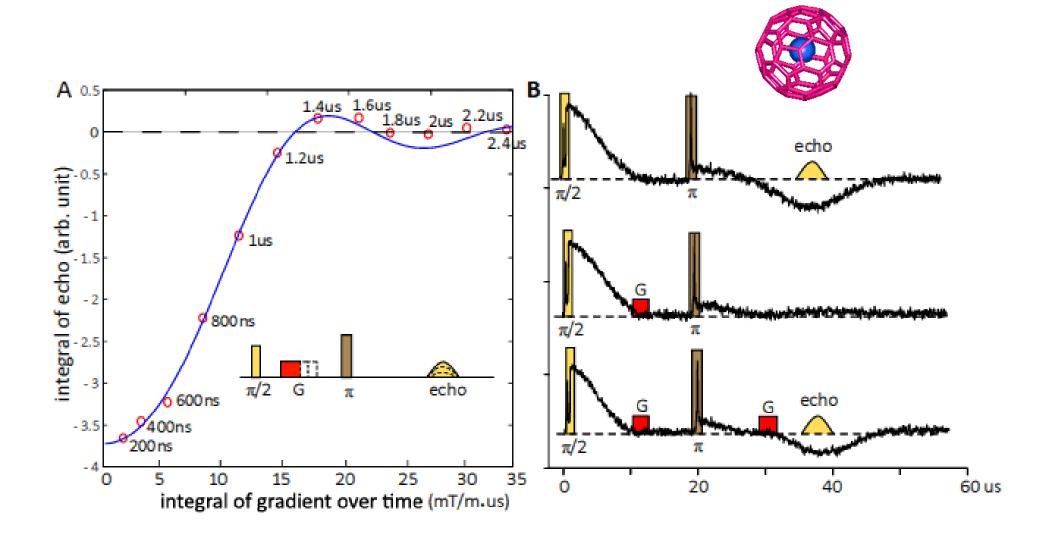
Key properties

- A processor head
- → Processor/Memory interface
- A homogenous storage medium
- ♦ Means of winding the tape

Optimize 4 elements using best suited quantum systems and immediately get many qubits!

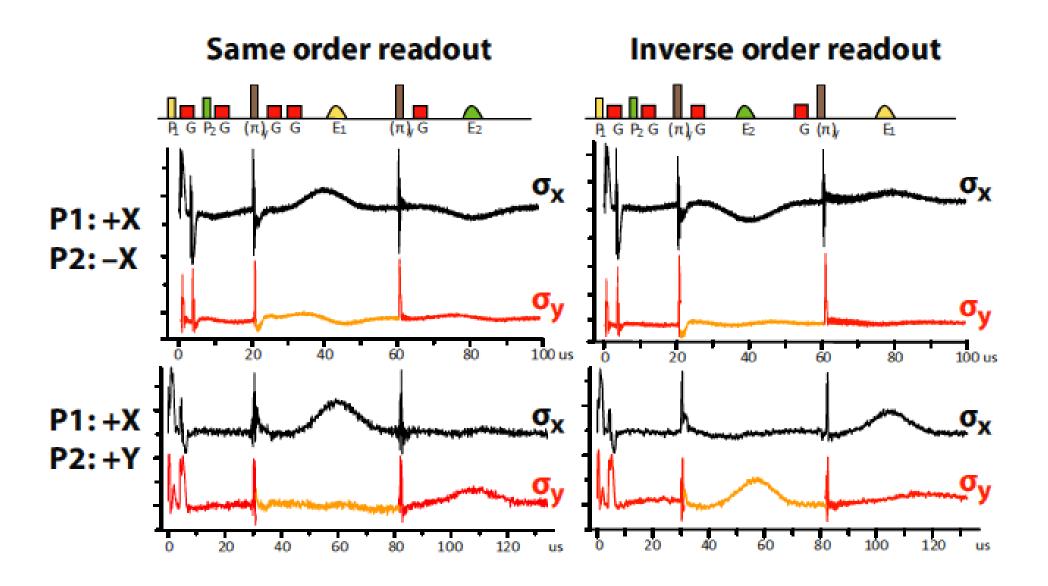
Dave Schuster





H. Wu et al., arXiv:0908.0101

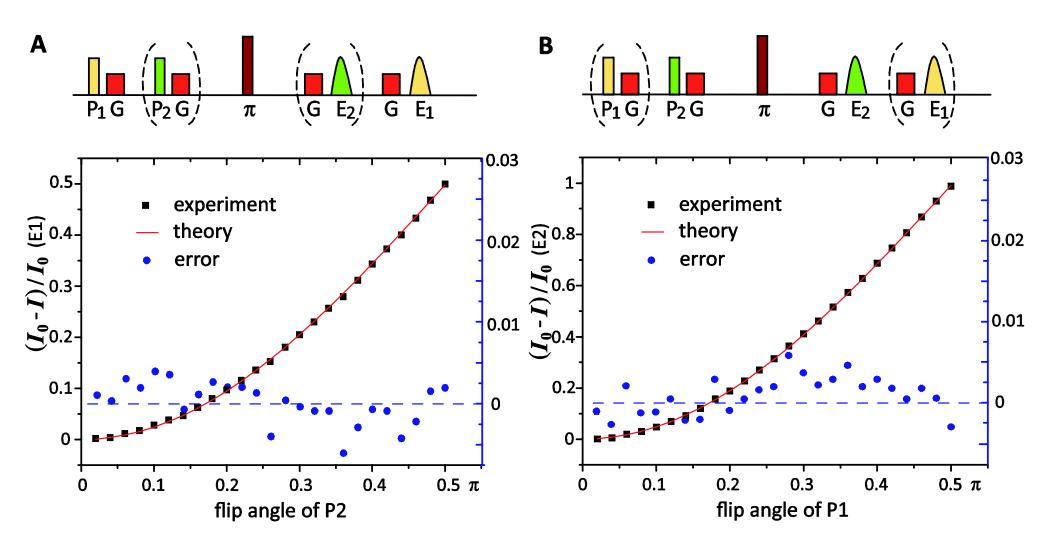




H. Wu et al., arXiv:0908.0101

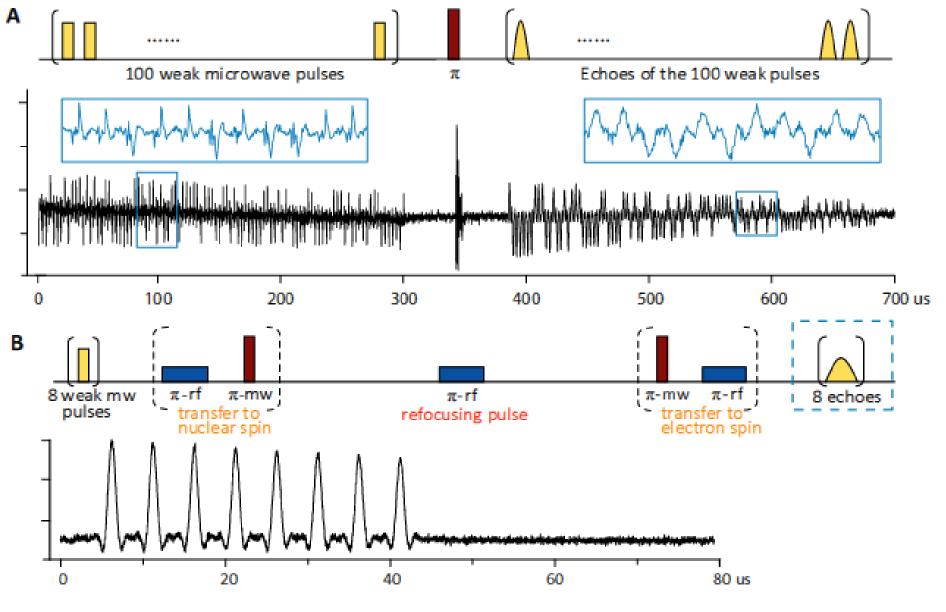


Crosstalk between nominally orthogonal modes



H. Wu et al., arXiv:0908.0101









Outstanding problems in using spin states

- In peapods:
 - how can you measure single spins?
 - what will be the EDMR mechanisms?
 - can you use RKKY for controlled coupling?
- In entanglement enhanced sensing:
 - can you achieve a pure state?
- In collective spin states:
 - how could you avoid using refocusing pulses?
 - can you avoid the Josephson junction bottleneck?



Outstanding problems in using spin states

