Entanglement hydrodynamics

and comments on tensor networks



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In collaboration with: Agón, Bao, Casini, Cotler, Hertzberg, Liu, Mueller, Stanford, van der Schee, Virrueta

Geometry from the Quantum, KITP, 01/13/2020



Transport



Thermalization



Butterfly effect

Phenomena associated with chaotic dynamics:







Thermalization



Butterfly effect

Ultimate goal: understand these phenomena and their relation in quantum systems, relate them to gravity through AdS/CFT

Goal of talk:

- Develop an effective theory of entanglement dynamics in the hydrodynamic limit
- Study its interplay with other chaotic phenomena
- Comment on relations to tensor network approaches to AdS/CFT

Setup:

- Degrees of freedom interacting strongly through local **chaotic** Hamiltonian.
- In highly excited state, out of equilibrium at t = 0, in equilibrium for $t \to \infty$.
- Foundational question in statistical physics. Subject of intense current activity in HEP, CMT, QI, and AMO experiments.



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Study the setup using holographic duality:

- A QFT settling to thermal equilibrium is dual to a collapsing black hole.
- No small parameters, holography is indispensible in understanding real time quantum dynamics.
- Entanglement plays a crucial role in thermalization.



We have an effective theory for describing conserved densities.

- Hydrodynamics applies universally for all chaotic systems. Generalized hydrodynamics for integrable systems.
- Navier-Stokes equations: $\partial_t v + (v \cdot \nabla)v \nu \nabla^2 v = -\nabla p$



Quantum system with many interacting degrees of freedom



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Quantum system with many interacting degrees of freedom

 Relativistic hydro from hep-th POV is an EFT based on systematic long distance, late time expansion. Fluid variables:

 $T_{ab} = (\rho + p)u_a u_b + p \eta_{ab} + \Pi_{ab}$

- Hydrodynamics follows from the conservation of T_{ab} . Solution determines $\langle T_{ab}\rangle$ out of equilibrium.

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• Fluid/gravity constructs black holes with bumpy horizons from fluid flows. [Bhattacharyya et al.]

$$ds^{2} = \frac{1}{z^{2}} \left[2u_{a}(x)dx^{a}dz + \left[\eta_{ab} + \left(1 - a\left(\frac{dz}{4\pi T(x)}\right) \right) u_{a}(x)u_{b}(x) \right] dx^{a}dx^{b} \right]$$

+ (gradients)

d

• Alternative history: String theorists discover hydrodynamics by studying AdS black holes.

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- Alternative history: String theorists discover hydrodynamics by studying AdS black holes.
- Interested in more data than $\langle T_{ab}\rangle$: entanglement entropy, butterfly effect, etc.
- I want to follow the "alternative history" path to discover a hydrodynamic effective theory of entanglement dynamics (and operator growth).
- Hydrodynamics is universal, there is evidence for the universality of the effective theory of entanglement hydrodynamics.

Quantum system with many interacting degrees of freedom



Outline





Transport

- Hydro as an EFT
- Holography for real time dynamics

Thermalization

- Entanglement entropy as a probe
- Membrane theory is the EFT
- Interplay with hydro



Comments on tensor networks

- Membrane theory from random circuits
- Interplay with operator growth

Conclusions and open questions

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Quantum thermalization and subsystems

Quantum thermalization

- Pure state with nonzero energy density: $|\psi(0)\rangle$ Unitary time evolution: $|\psi(t)\rangle = e^{-iHt}|\psi(0)\rangle$
- $\rho(t) \equiv |\psi(t)\rangle\langle\psi(t)| \not\rightarrow \frac{e^{-\beta H}}{Z}$ cannot mean thermalization.

 $\rho(t)$ encodes all the information in $|\psi(0)\rangle$, but at late times in a very nonlocal way.

Quantum system with many interacting degrees of freedom

Fully isolated from environment

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- Consider subsystems, reduced density matrix: $ho_A = {
 m Tr}_{ar A} \, |\psi
 angle \, \langle\psi|$
- Thermalization: $\rho_A(t) \rightarrow \rho_A^{(eq)}(\beta) = \text{Tr}_{\bar{A}} \frac{e^{-\beta H}}{Z}$

For $t \to \infty$, in the thermodynamic limit $\bar{A} \to \infty$, with β determined by the energy density. Entanglement is crucial in making this possible.



Entanglement entropy

Entanglement entropy is a good diagnostic of thermalization, we **focus on this quantity**.

• In ground states of local Hamiltonians the entropy scales with the area:

 $S_A = \# \frac{\operatorname{area}(\Sigma)}{\delta^{d-2}} + \dots$

• A generic state in the Hilbert space shows volume scaling: $S_A = s_{\rm th} \operatorname{vol}(A) + \dots$



S Typical point inside unentangled with outside



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 Purest setup is a quench: start with ground state of a local Hamiltonian, change the Hamiltonian suddenly, and let the system evolve. (No transport.)









Can reformulate holographic surface extremization in d+1 dimensions as membrane minimization in d dimensions in the limit $R, t \gg t_{\rm loc}$. [MM₂] A

- Detailed understanding of HRT surfaces. The surface has three parts: [MM₁]
 - 1. Outside the horizon part gives (divergent) area law.
 - 2. Behind the horizon region.
 - 3. Behind the shell part gives entropy in the vacuum.



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- Push HRT surface to the boundary along constant infalling time.
- Scaling limit: $x^{\mu} \rightarrow R x^{\mu}$, $z \rightarrow z$ Area functional independent of the derivatives of z. Solve algebraic EoM, plug back into action to derive membrane theory.

$$S[A] = s_{\rm th} \int d^{d-1}\xi \sqrt{\gamma} \frac{\mathcal{E}(v)}{\sqrt{1-v^2}}$$

Horizon ~ boundary
$$t \int v$$

t for the second second





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• Using the NEC, can prove the following properties of $\mathcal{E}(v)$. $\mathcal{E}(v)$ can be thought of as a transport coefficient.





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• Simple bound on saturation time from operator growth: [MM, Stanford] $t_S \ge R/v_B$ For elongated shapes in 4D we find: $t_S = R/v_B$ Black holes often saturate entanglement entropy the fastest.

Interplay with hydro

The membrane theory is robust, can be generalized away from global quenches. [MM, Virrueta]

• Fluid/gravity black brane dual to an inhomogenous state in local thermal equilibrium. To subleading order, we get the membrane coupled to hydrodynamics:

$$S = \int d^{d-1}\xi \,\sqrt{\gamma} \,s_{\rm th}(x) \frac{\mathcal{E}(v)}{\sqrt{1-v^2}} + \dots, \quad v(x) \equiv \frac{(n \cdot u(x))}{\sqrt{1+(n \cdot u(x))^2}}$$

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- Adaptable to other inhomogenous setups, can incorporate β/R and $1/\lambda$ corrections without change in the structure of the membrane theory. 1/N corrections would be most interesting.
- Membrane theory is versatile, has connections to operator growth and hydrodynamics, and has all the features to be a universal theory.

Features of the thermalization:

- Conserved densities described by hydro.
- State of the entire system cannot become thermal. Small subsystem thermalize by becoming entangled with the rest of the system.

 $S_A(t) \rightarrow S_A^{(eq)}(\beta) = s_{th}(\beta) \operatorname{vol}(A)$ Captures the essence of thermalization.

Goal: Find effective theory (akin to hydro) of entanglement dynamics.

- Alternative history method: Discovered membrane theory by studying AdS black holes, has structure applicable to all chaotic theories.
- In the following conduct further tests. Elucidate connections to other manifestations of chaotic dynamics.





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Conclusions and open questions

The same description of entanglement dynamics arises in CMT.

• Random quantum circuit model for the evolving wave function.



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Minimal cut computes the entropy. [Nahum, Ruhman, Vijay, Haah]



Minimal membrane phenomenology of entropy dynamics. [Jonay, Huse, Nahum]

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- Analytic arguments in Floquet systems. [Nahum, Zhou] Evidence in chaotic spin chains. [Jonay, Huse, Nahum]
- Remarkable unification of CMT and HEP approaches: Membrane description of EE growth in quenches.

Tensor networks and holography

The analogy between minimal cuts and the RT surface computing entanglement entropy has inspired toy models of holography. u = -3

 AdS/MERA analogy, [Swingle] perfect and random tensor networks [Pastawski et al; Hayden et al.]



 Suggestive results for maximal volume slice. [Hartman, Maldacena; Roberts, Stanford, Susskind] But HRT surfaces for different shapes do not lie on same Cauchy slice.



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 Entanglement of local operator with growing footprint is computed by membrane in time fold geometry. [Roberts, Stanford, Susskind; Jonay, Huse, Nahum; MM, Virrueta]



• Quantitative connection to TNs through EoM, bulk geometry encoded in $\mathcal{E}(v)$.

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Universality classes of entropy dynamics

I propose that there are two universality classes of entropy dynamics at long distances and late times (in translationally invariant systems).

- 2d integrable models, RCFTs, d>2 free theories are described by the **quasiparticle theory**.
- The holographic results can be reformulated in terms of a **membrane theory**, which then can be adopted to any chaotic system. Applies to holographic theories, random circuits, evidence for chaotic spin chains. [Jonay et al., MM₂]
- Is there something in between?
- Analogous to the dichotomy between generalized hydrodynamics applicable to integrable systems (giving ballistic transport) and hydrodynamics (describing diffusive transport).

- Qualitative picture of entanglement entropy at time t of a region of characteristic size R, R, t ≫ t_{loc}.
 [Cardy, Calabrese; Hartman, Maldacena; Liu, Suh]
- EE in free scalar theory for a disk, dots are data points, line is quasiparticle theory [Cotler, Hertzberg, MM, Mueller]

 EE in holographic theories for a disk, data collapse, solid line is membrane theory, deviation is controlled by 1/R [MM₁]



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- Universality classes of thermalization:
 Quasiparticle theory vs Membrane theory
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 - Membrane couples to hydrodynamics
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 - Membrane is a cut through TN, TN is obtained after solving bulk EoMs



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 - Membrane is a cut through TN, TN is obtained after solving bulk EoMs
- Rich applications
 - Entropy cone inequalities generalized to time dependent settings. [Hayden, Headrick, Maloney; Bao et al.; Bao, MM]
 Bit threads reformulation. [Freedman, Headrick; Agon, MM]
- Membrane theory has all the features to be a universal theory.

Open questions and outlook for gravity

Open questions and some hints

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- Are new quantum extremal surfaces, islands be captured by the membrane theory? Hint: It looks plausible that 1/N corrections can be captured by the membrane theory. It may be that we get multiple minimal membranes for evaporating BH.
- Is the membrane theory a good starting point to getting gravitational dynamics out of entanglement?
 Hint: Slogan: "Gravity is the hydrodynamics of entanglement." May have to go to shorter times and distances in CFT to see dynamical geometry.

Backup slides

Calabrese-Cardy model: energy injection from quench creates a finite density of EPR pairs, subsequently travel freely at the speed of light isotropically.

- Leads to linear growth with $v_E = 1$ in 2d.
- Higher dimensions: entanglement spreading depends on entanglement pattern on the light cone μ[L_Σ].
 Contribution from each light cone has to be added.
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Bound on the entanglement speed from SSA:

$$v_E \le v_E^{(\text{EPR})} = \frac{\Gamma(\frac{d-1}{2})}{\sqrt{\pi}\Gamma(\frac{d}{2})} < v_E^{(\text{SBH})}$$

Slower than holography.



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Slower than holography.

- In strongly coupled systems, entanglement grows faster than what's possible for free particles streaming at the speed of light!
- Consider the effect of interactions: tensor network picture emerging from scattering particles is natural. [Hartman, Maldacena; Casini, Liu, MM]



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• Time evolution of a Gaussian initial state is Gaussian (with time dependent complex kernel).

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$$\chi_I = \begin{pmatrix} \phi_i \\ \pi_i \end{pmatrix}, \qquad [\chi_I, \chi_J] = i J_{IJ}$$
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 The symplectic eigenvalues of the correlation matrix give the eigenvalues of the reduced density matrix:

$$\begin{split} \tilde{\chi} &= S\chi \,, \qquad SJS^T = J \,, \\ \tilde{\Gamma} &= S\Gamma S^T = \begin{pmatrix} \operatorname{diag}\left(\gamma_k\right) & 0 \\ 0 & \operatorname{diag}\left(\gamma_k\right) \end{pmatrix} \end{split}$$

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 Numerical results for 3d boundary state quench for scalar field. [Cotler, Hertzberg, MM, Mueller]





Entropy cone

Entanglement entropy in static holographic states obeys inequalities, that are not true in general in QM.

• The best known one is the monogamy of mutual information. [Hayden, Headrick, Maloney] It can be proven using the same steps as in the proof of SSA.

 $S(AB) + S(BC) + S(AC) \geq S(A) + S(B) + S(C) + S(ABC)$



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 [Bao et al.] Holography is not essential, only need that the entropy is proportional to a partionable geometric minimization problem.
- HRT is an extremization of codimension-2 surface, no proof (or counterexample) is known for many-party inequalities. Inclusion-exclusion applies to the membrane theory, hence proof for time dependent states (large regions, late times). [Bao, MM]

Bit threads

The Ryu-Takayanagi prescription can be reformulated in the language of bit threads. [Freedman, Headrick]

[Freedman, Headrick] • Maximize $\int_A \sqrt{h} n_\mu w^\mu$

Constraints: $\nabla_{\mu}w^{\mu} = 0$, $1 - |w^{\mu}| \ge 0$

• Covariant generalization to HRT [Headrick, Hubeny]

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- Covariant generalization to HRT [Headrick, Hubeny]
- Membrane theory can also be similarly reformulated without reference to holography. Only one constraint changes [Agon, MM]

 $\nabla_{\mu}w^{\mu} = 0, \quad H(w_t) - |\vec{w}| \ge 0$

 $H(w_t)$ is the Legendre transform of $\mathcal{E}(v)$:

 $H(w_t) \equiv \mathcal{E}(v) - v \,\mathcal{E}'(v), \quad w_t \equiv -\mathcal{E}'(v)$

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The Ryu-Takayanagi prescription can be reformulated in the language of bit threads. [Freedman, Headrick]

• Maximize $\int_A \sqrt{h} n_\mu w^\mu$

Constraints: $\nabla_{\mu}w^{\mu} = 0$, $1 - |w^{\mu}| \ge 0$

- Covariant generalization to HRT [Headrick, Hubeny]
- Membrane theory can also be similarly reformulated without reference to holography. Only one constraint changes [Agon, MM]

$$abla_{\mu} w^{\mu} = 0, \quad H(w_t) - |\vec{w}| \ge 0$$

 $H(w_t)$ is the Legendre transform of $\mathcal{E}(v)$:

 $H(w_t) \equiv \mathcal{E}(v) - v \,\mathcal{E}'(v), \quad w_t \equiv -\mathcal{E}'(v)$

- The map that reconstructs the HRT surface from the minimal membrane can be used to push the membrane theory bit thread into the bulk.
- Membrane theory is versatile, has connections to operator growth and hydrodynamics, and has all the features to be a universal theory.

Entanglement entropy obeys inequalities, natural to consider bounds in the quench setup.

• $v_E \leq 1$ can be proven using Lorentz invariance and the SSA inequality, [Casini, Liu, MM] or the monotonicity of relative entropy. [Afkhami-Jeddi, Hartman]

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- Monotonicity of (thermal) relative entropy for subsystems combined with emergent v_B light cones at finite temperature in chaotic systems:

 $S[A(t)] \le S[A'(t')] + s_{\rm th} \left(V[A(t)] - V[A'(t')] \right)$

Gives bound for all times. Can be combined with another proposed inequality. [MM, Stanford]



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$$v_E \le v_B \,, \quad t_S \ge \frac{R_{\text{insc}}}{v_B}$$

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• Membrane theory proof: there exists a maximal membrane tension compatible with the general properties discussed before.

$$\mathcal{E}_{\max}(v) = v_E + \left(1 - \frac{v_E}{v_B}\right)|v| \qquad (|v| \le v_B)$$

The resulting minimal membrane is a combination of a cylinder and the cone saturating the combined inequalities. $[MM_2]$

