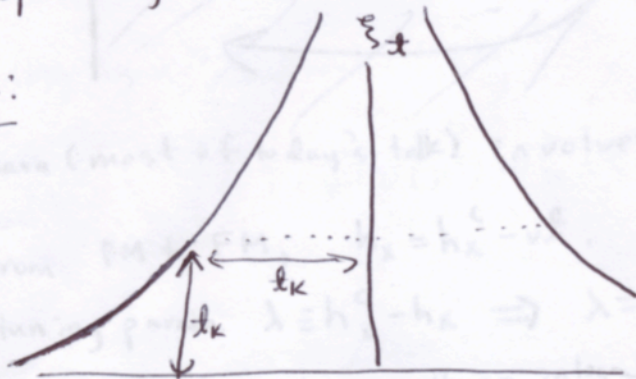


Title: Universality and Dephasing within the Kibble-Zurek Mechanism

Talk has three major goals:

- 1) Provide numerical evidence that standard scaling arguments (including universality) apply to certain non-equilibrium protocols, such as Kibble-Zurek (KZ)
- 2) Discuss certain phenomena which appear only within the KZ scaling theory, i.e., the existence/non-existence of dephasing as a fn. of protocol.
- 3) Describe how KZ and its (natural) extensions ~~can~~ ^{may} be useful to probe universal aspects of dephasing and thermalization.

Basic ideas of KZ:



Fallout of eq. on timescale where system does not have time to respond

"Tuning parameter"
 $\lambda = v t$
 Ramp rate

$$t_k = \sum_t (t_k) \xrightarrow{\text{Algebra}} \boxed{t_k = v \frac{-v\tau}{1+v\tau}} \Rightarrow l_k = t_k = v \frac{-v\tau}{1+v\tau}$$

Original formulation (Kibble 1976, Zurek 1985) - l_k sets density of defects after ramp

Modern perspective (Deng et al. 2008, Chandran et al. 2012, ...) - t_k & l_k set scales in well-defined neq. scaling limit

$$\xi_t^{eq} \rightarrow t_k$$

$$\xi_{eq} \rightarrow l_k$$

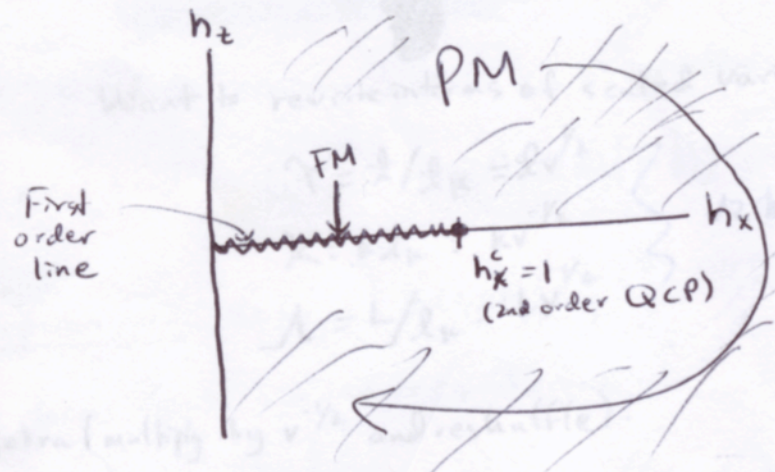
Also exists similar story for non-linear ramps (Barentine & Polk. 2008, Sen et al. 2008)

If $\lambda = v |t|^\nu \text{sgn}(t)$, then $t_k \rightarrow v \frac{-v\tau}{1+v\tau}$

Special case $r=0 \Leftrightarrow$ symmetric quench across CP (v plays the role of quench depth)

The model: (Extended) transverse-field Ising ^{chain} ~~model~~

$$H = -\sum_j [\sigma_j^z \sigma_{j+1}^z + h_x \sigma_j^x + h_z \sigma_j^z]$$



Simplest case (most of today's talk) involves just TFI chain, $h_z = 0$.

Ramp from PM to FM, $h_x = h_x^c - vt$.

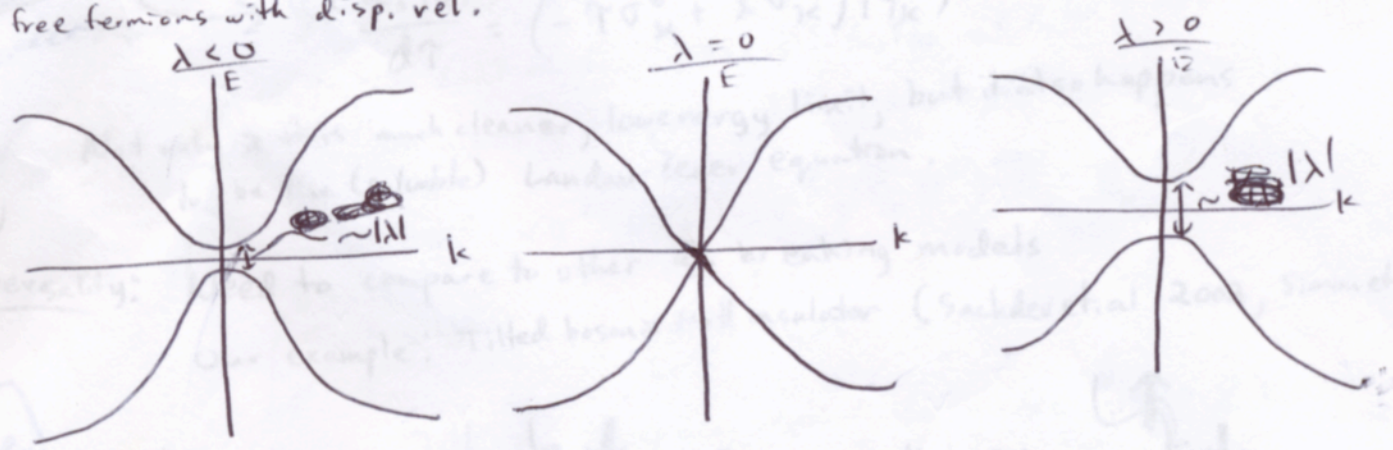
Define tuning param. $\lambda \equiv h_x^c - h_x \Rightarrow \lambda = vt$

This is a ramp for an integrable theory with an integrable critical point

Critical exponents: $\nu = z = 1 \Rightarrow t_k = l_k = \nu^{-1/2}$ (note that $t_k, l_k \rightarrow \infty$ as $\nu \rightarrow 0$)

Solution is via Jordan-Wigner (Sachdev 1999) cf.

Gives free fermions with disp. rel:



$$H_k = [-\lambda + \sigma(k^2)] \sigma_k^z + [k + \sigma(k^2)] \sigma_k^x$$

$$E_k = \pm \sqrt{\lambda^2 + k^2}$$

Dynamic scaling can be motivated from Schrödinger eqn:

$$i \frac{d|\Psi_k\rangle}{dt} = \left[(-v\tau + \mathcal{O}(k^2)) \sigma_k^z + (k + \mathcal{O}(k^3)) \sigma_k^x \right] |\Psi_k\rangle$$



Want to rewrite in terms of scaled variables:

$$\Upsilon = t/l_k = \tau v^{1/2}$$

$$\mathcal{X} = k l_k = k v^{-1/2}$$

$$\mathcal{L} = L/l_k = L v^{1/2}$$

√2.11e (LEFT)

Via some algebra (multiply by $v^{-1/2}$ and reshuffle):

$$i \frac{d|\Psi_x\rangle}{d\Upsilon} = \left[(-\Upsilon + v^{1/2} \mathcal{O}(\mathcal{X}^2)) \sigma_x^z + (\mathcal{X} + v \mathcal{O}(\mathcal{X}^3)) \sigma_x^x \right] |\Psi_x\rangle$$

Scaling limit:

$$\lim_{l_k \rightarrow \infty (v \rightarrow 0)} \left[\text{Schr. eqn} \right]$$

$t/l_k, k l_k, L/l_k \text{ fixed}$

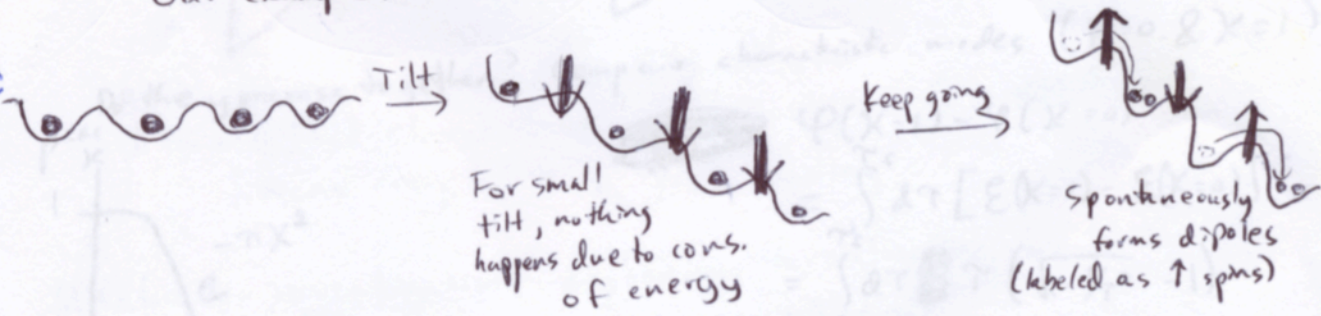
$$i \frac{d|\Psi_x\rangle}{d\Upsilon} = (-\Upsilon \sigma_x^z + \mathcal{X} \sigma_x^x) |\Psi_x\rangle$$

Not only is this much cleaner, low energy limit, but it also happens to be the (solvable) Landau-Zener equation.

Universality: Need to compare to other \mathbb{Z}_2 breaking models

Our example: Tilted bosonic Mott insulator (Sachdev et al. 2002, Simon et al. 2002)

Visible
(HEG/AF)



For small tilt, nothing happens due to cons. of energy

Spontaneously forms dipoles (labeled as ↑ spins)

\mathbb{Z}_2 symmetry \leftrightarrow Dipoles on even/odd bonds

In practice map to spin model (dipole = \uparrow).
 Now QPT from $\downarrow\downarrow\downarrow$ PM to $\uparrow\downarrow\downarrow$ AFM.
 Simulate this model with time-dep DMRG.

Variable (LEFT)

Observables: Need to compare models via observables. Will use two:

Excess heat: $Q(t) \equiv \langle \psi(t) | H(t) | \psi(t) \rangle - \langle \psi_{gs}(t) | H(t) | \psi_{gs}(t) \rangle$

Equal-time Order param. corr.: $G(x, t) \equiv \langle \psi(t) | S_j^z S_{j+x}^z | \psi(t) \rangle$

Scaling forms look similar to equilibrium

$$\frac{Q}{L} \rightarrow \frac{Q t_k}{L t_k} = \frac{Q}{vL} = q(\gamma, \nu)$$

$$G \rightarrow G = x^{-1/4} g(\gamma, \nu, X = x/t_k)$$

(show data for universality on slides)

Part 2: Dephasing - Know that $|\psi_i\rangle \xrightarrow{t} \text{Dephased ensemble (GGE/Therm.)}$

Does dephasing occur within KZ scaling limit?

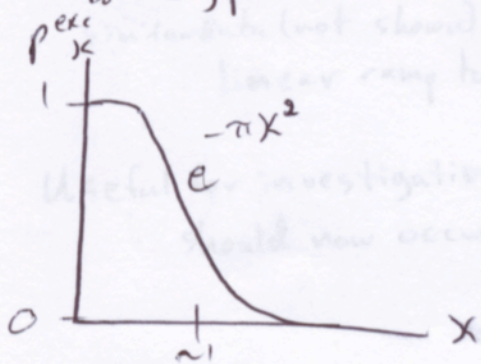
Naive guess is Yes: because gap increases as a function of time.

It is true that individual modes "precess" faster at later times
 But what matters for observables is the phase difference between modes

$$H t_k \rightarrow -\gamma \sigma^z + X \sigma^x = \gamma \left(-\sigma^z + \frac{X}{\gamma} \sigma^x \right) \Rightarrow \vec{B}_x \cdot \vec{\sigma}$$

$\vec{B}_x = \gamma \hat{z} + X \hat{x}$
 $\sigma_x = \frac{X}{\gamma}$

Do they precess together? Compare characteristic modes ($X=0$ & $X=1$)



$$\begin{aligned} \Delta \Phi &= \Phi(X=1) - \Phi(X=0) \\ &= \int_0^{\tau_f} d\tau [E(X=1) - E(X=0)] \\ &= \int_0^{\tau_f} d\tau \gamma (\sqrt{1 - 1/\tau^2} - 1) \\ &\approx \int_0^{\tau_f} d\tau \frac{1}{2\tau} = \frac{1}{2} \ln \tau_f \end{aligned}$$

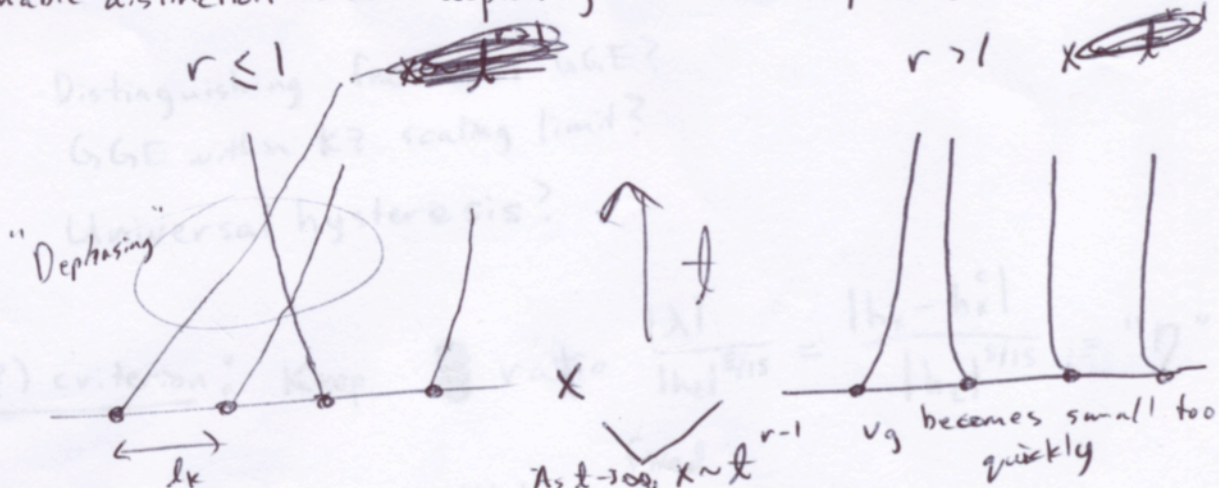
For linear ramp, $\Delta \Phi \rightarrow \infty$ as $\tau_f \rightarrow \infty \Rightarrow$ Dephases

For ramp with arbitrary r , $\Delta\psi \approx \lim_{T_f \rightarrow \infty} \int_0^{T_f} d\tau \frac{1}{2\tau^r}$

$r \leq 1 \Rightarrow \infty \Rightarrow GGE$

$r > 1 \Rightarrow \text{const.} \Rightarrow \text{Modes become phase locked (not GGE)}$

Protocol-tunable distinction between dephasing to GGE & not dephasing



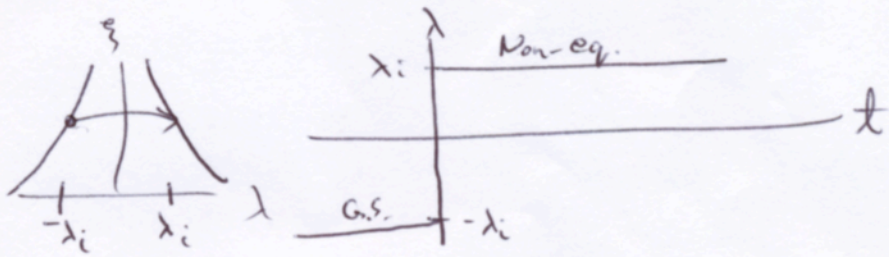
$r = 1$ is marginal case where $\chi \sim \ln t$ ("z = ∞ "?)

Part. 3: Extensions

There exist essentially arbitrary protocols to extend these ideas of KZ scaling theory (at least in theorist fantasy land)

One useful scaling direction: Time after quench/ramp

Simplest case is $r = 0$:



$l_k \sim \xi_{eq}(\lambda_i) \sim |\lambda_i|^{-1}$

Data shows scaling collapse
 similar data (not shown) shows scaling collapse for ~~linear ramp~~
 linear ramp to $r = 1$, then waiting.

Useful for investigating dephasing/thermalization, as it should now occur in finite (scaled) time.

~~Structure~~
Within the Ising model, can investigate other, less trivial scaling directions.

Longitudinal field h_z : $\xi_{\text{eq}} \sim |h_z|^{-8/15}$ ($\nu = 8/15$)

Integrable (but non-trivial E_8 field theory, Zamolodchikov 1989, Coldea et al. 2010)

Distinguishing features of GGE?

GGE within KZ scaling limit?

Universal hysteresis?

Thermalizing (?) criterion : Keep $\frac{|\lambda|}{|h_z|^{8/15}} = \frac{|h_x - h_x^c|}{|h_z|^{8/15}} \equiv \eta$ fixed.

Each value of $\eta \neq 0, \infty$ defines a ^{likely} non-integrable field theory.

Does thermalization occur in the KZ scaling limit?

Time scales for GGE/pretherm/therm. for ramp/quenches near the integrable QCP?