

# Optimal control for quantum computing

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# Outline

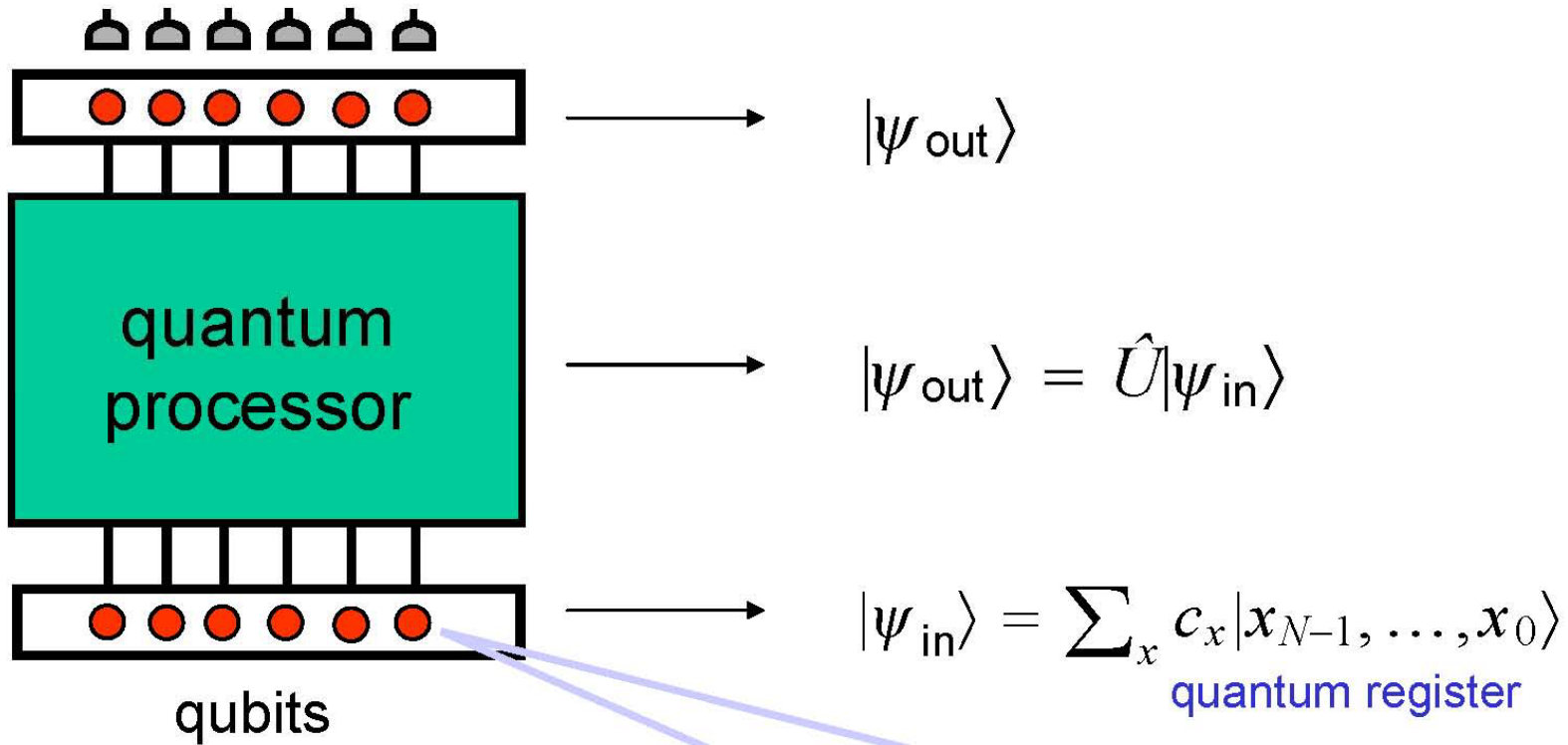
- Scalable quantum information systems
  - what do we need?
- Quantum control
  - what can we do?
- Entanglement generation
  - using different interactions
- What can go wrong
  - and how to fix it
- How much can we push?
  - the Quantum Speed Limit

# Scalability

What do we need?

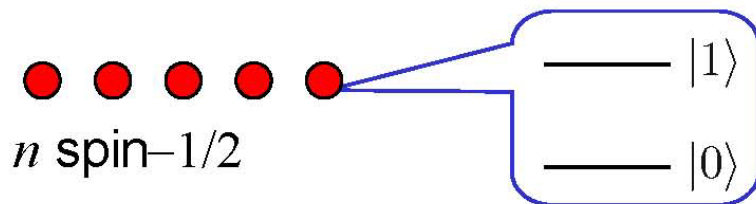
[TC, Grangier, Walraff, Zoller, Nature Phys. '08]

# quantum computing



- quantum memory: qubits

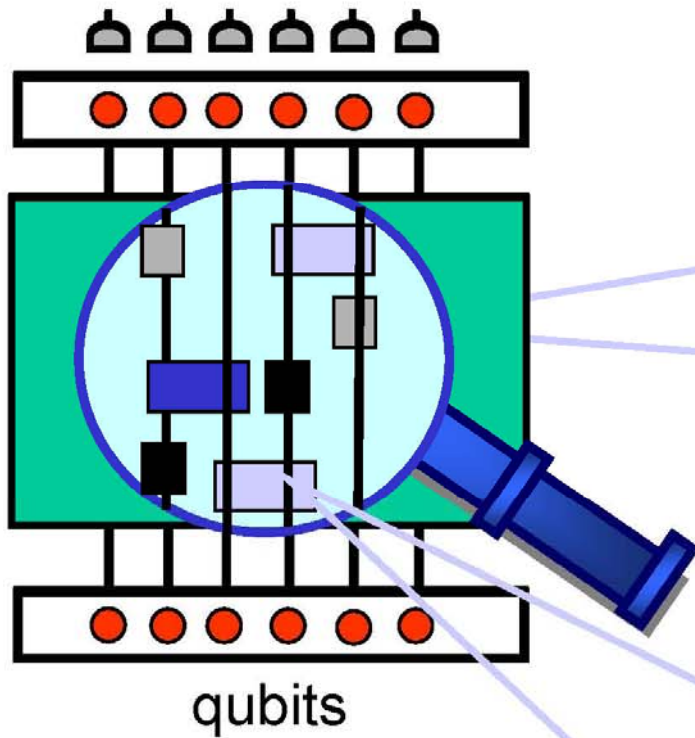
$n$  spin-1/2



example: two qubit entangled state

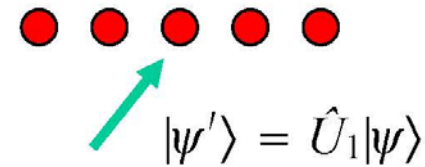
$$|\psi\rangle = c_{00}|00\rangle + c_{01}|01\rangle + c_{10}|10\rangle + c_{11}|11\rangle$$

# quantum computing

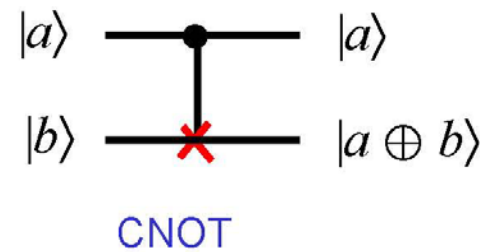
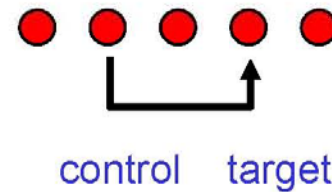


# quantum gates

- single qubit gate



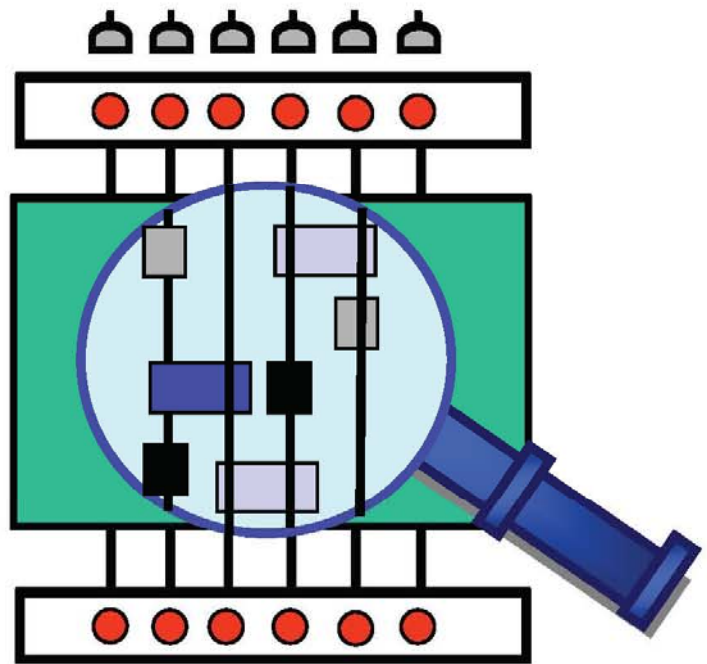
- two qubit gate: entanglement



## truth table

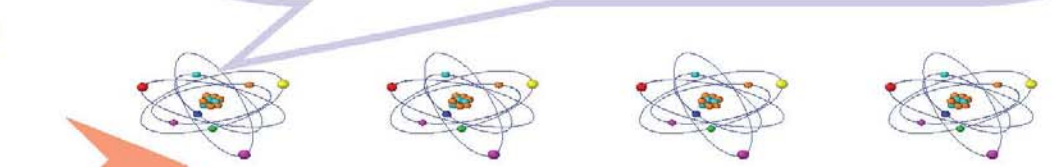
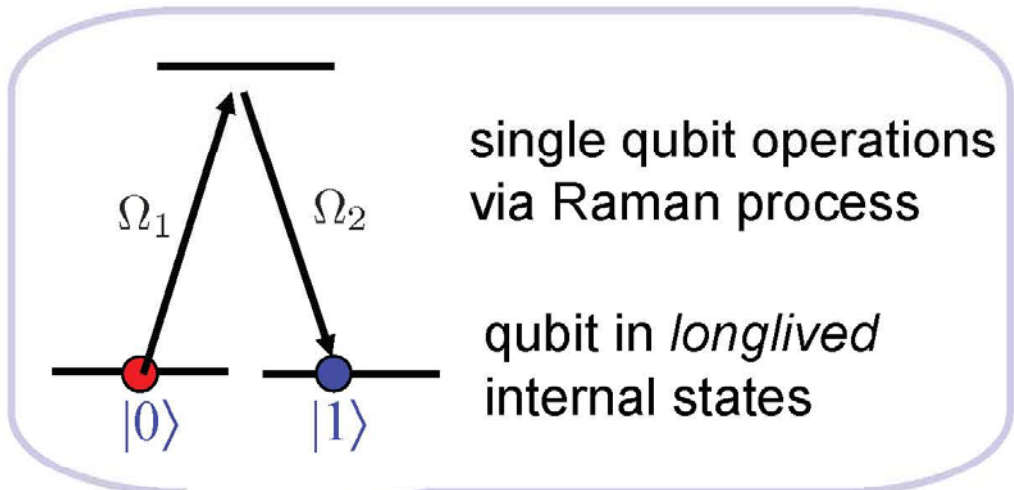
$ 0\rangle 0\rangle$	$\rightarrow$	$ 0\rangle 0\rangle$
$ 0\rangle 1\rangle$	$\rightarrow$	$ 0\rangle 1\rangle$
$ 1\rangle 0\rangle$	$\rightarrow$	$ 1\rangle 1\rangle$
$ 1\rangle 1\rangle$	$\rightarrow$	$ 1\rangle 0\rangle$

# quantum computing



qubits

# physical realization

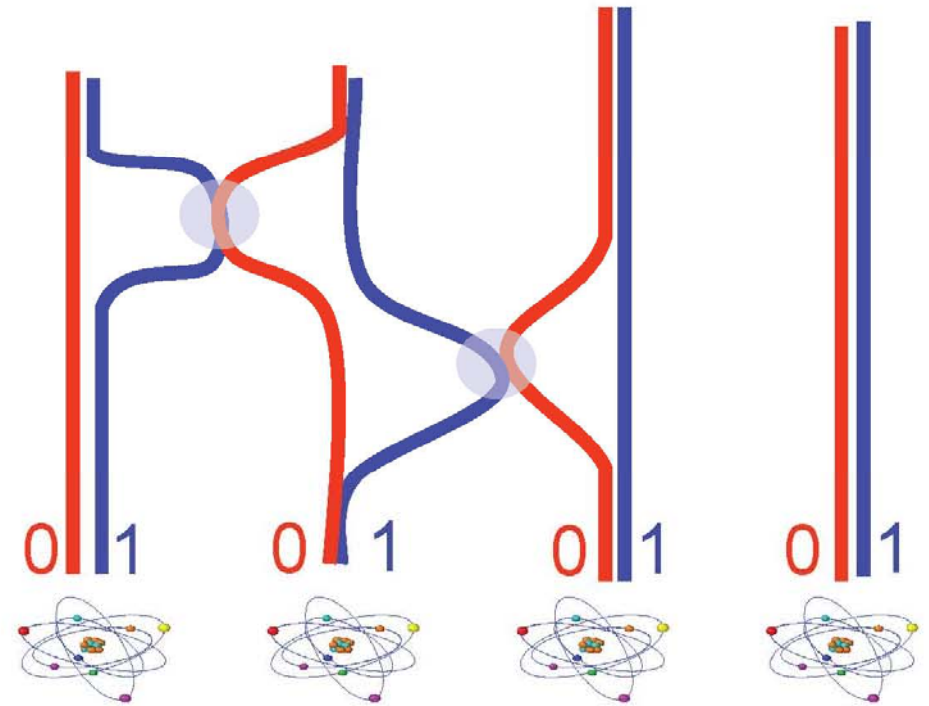
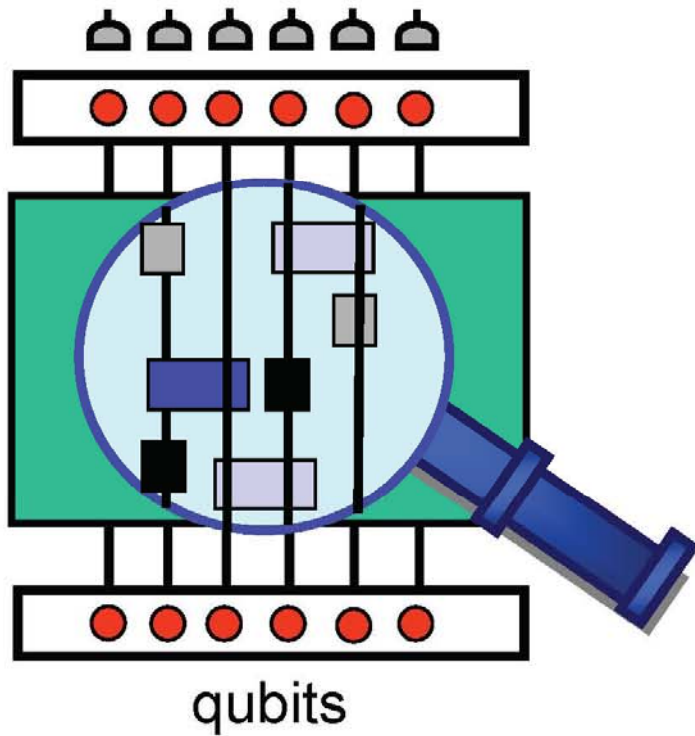


atoms as qubits

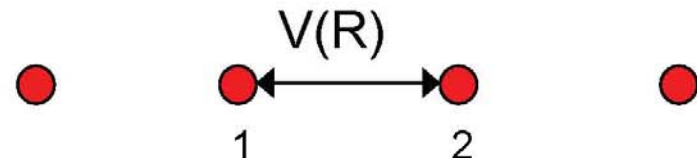
laser

Requirements:  
addressing  
single qubit

# entangling qubits

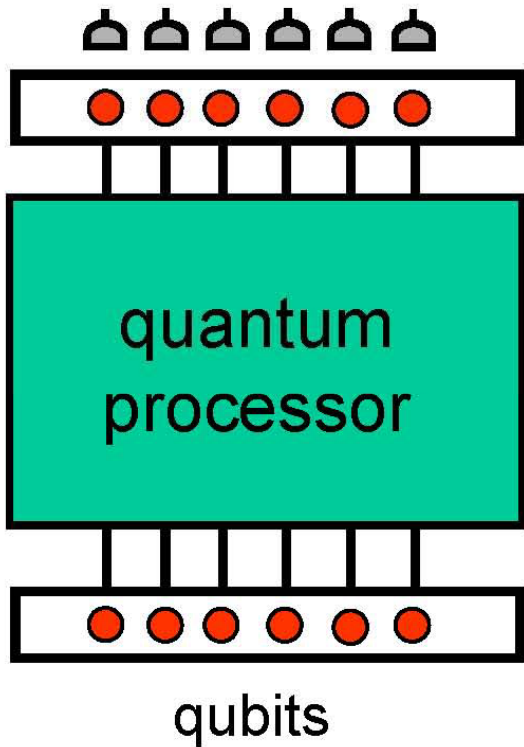


**controllable two body interactions:**  
controlled collisions, ...



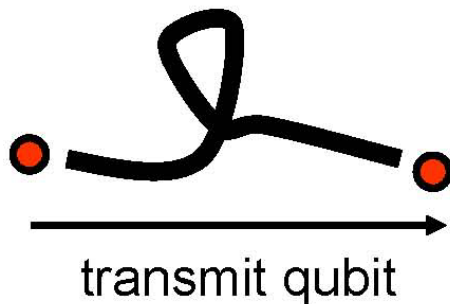
Hamiltonian  $H = \Delta E(t)|1\rangle_1\langle 1| \otimes |1\rangle_2\langle 1|$

so that  $|1\rangle_1|1\rangle_2 \rightarrow e^{i\phi}|1\rangle_1|1\rangle_2$



## DiVincenzo Criteria

1. **scalable** system of well-characterized qubits
2. initialize qubits
3. long decoherence times
4. universal set of quantum gates
5. qubit readout
  
6. interconvert stationary and flying qubits
7. faithful transmission of qubits between specified locations

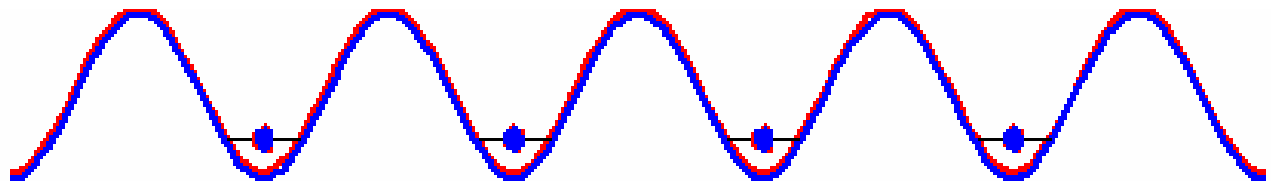


GOAL: satisfy requirements of fault tolerant quantum computing



# Scalability desiderata

- Memory:
  - Quantum register with many qubits
  - Low decoherence rates
- Gates:
  - Fast operation
  - High fidelity
- ...implementation with ultracold systems:
  - Good isolation from environment
  - Individual control
  - Periodic potentials



# Scalability in practice

- Nobody really knows how to build a working QC
  - need to have quite fast AND ultra-accurate gates
- Specific practical problems – for instance
  - transport of particles in traps
  - strong coupling to control vs weak coupling to environment
- Quantum Optimal Control Theory
  - tailored answers to specific problems
  - ...robustness to noise?

# Control

What can we do?

# Example: transport in traps

Initial state

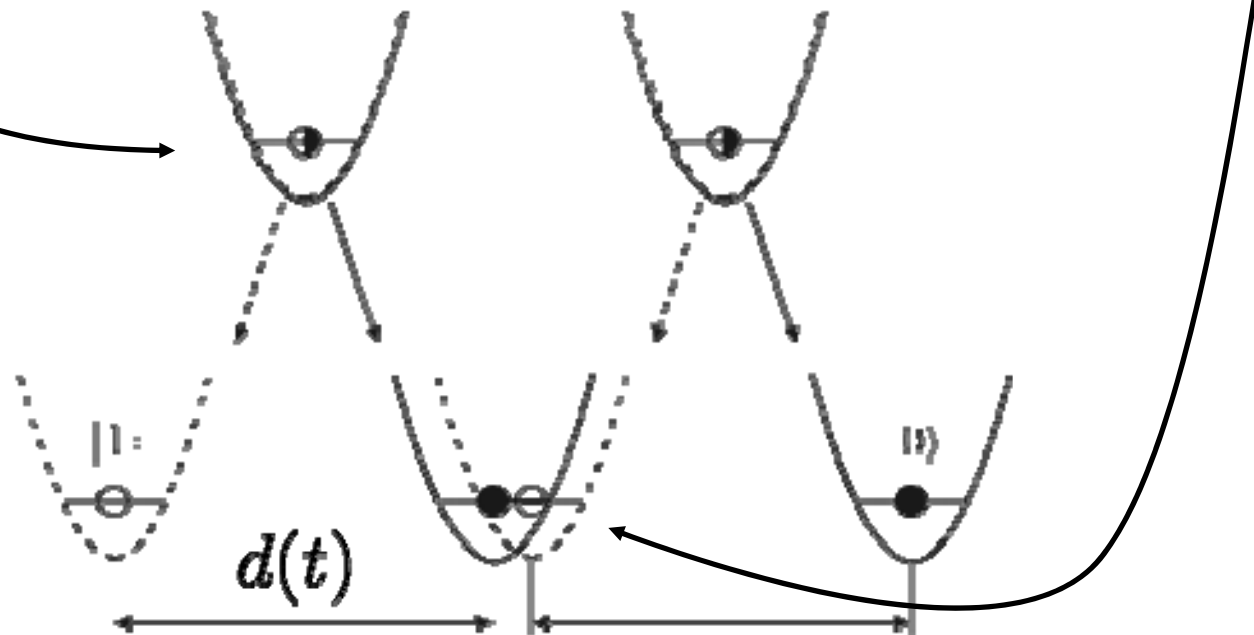
Hamiltonian

Evolved state

$|\psi_0\rangle$

$\mathcal{H}(d(t), t)$

$|\psi(T)\rangle$



Task: minimize error

$$1 - \mathcal{F} = 1 - |\langle \psi_{\text{goal}} | \psi(T) \rangle|^2$$

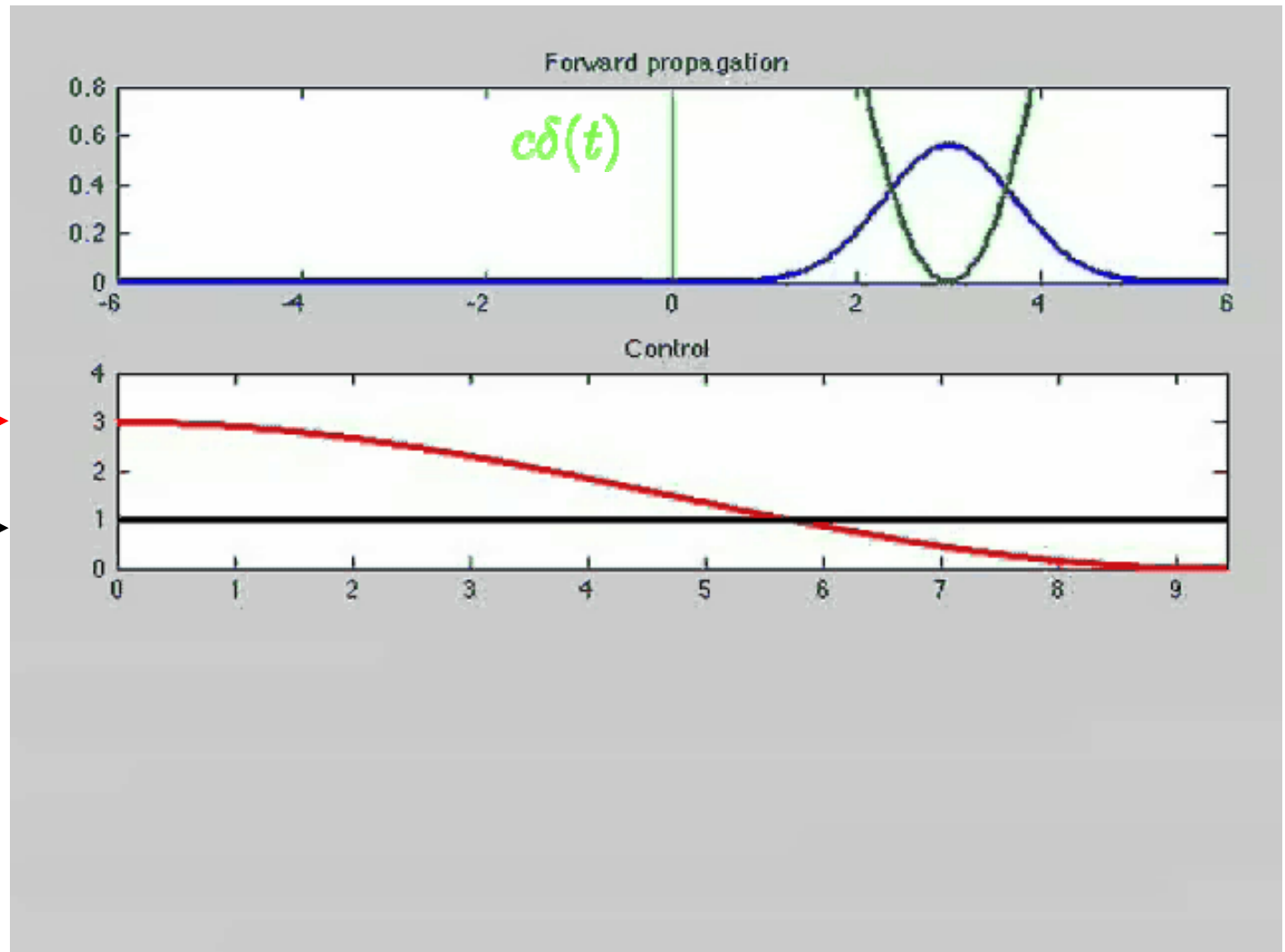
# Optimization via Krotov

$$\frac{m\omega^2(t)}{2} [x - d(t)]^2$$

$d(t)$  →

$\omega(t)$  →

$\mathcal{F}$

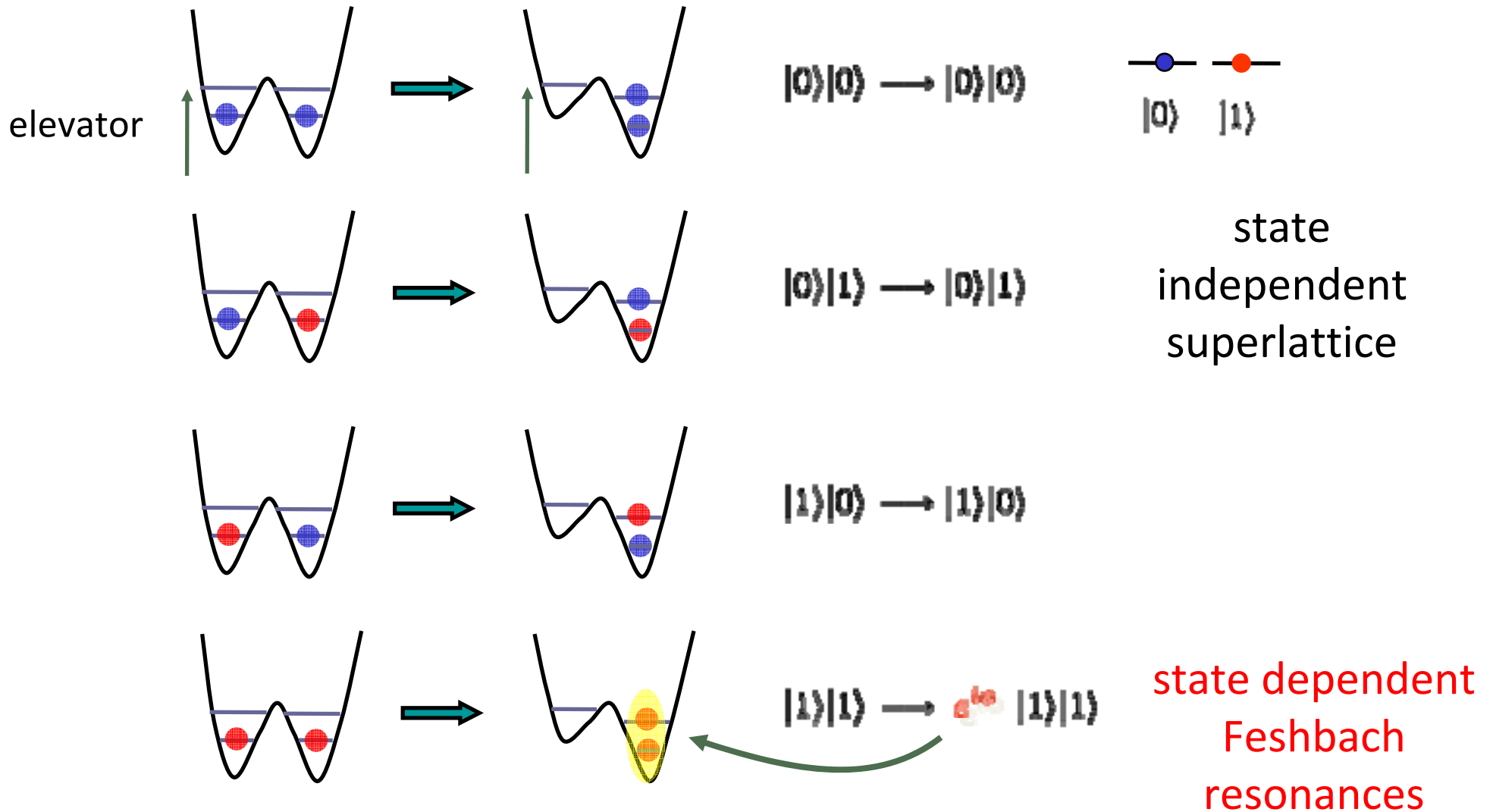


# Entanglement

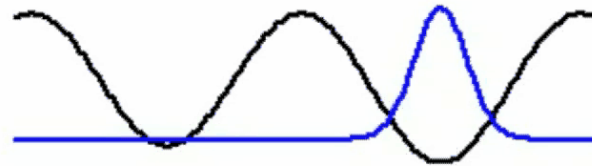
using atomic interactions

# Entangling Feshbach resonances

with P. Julienne, P. Zoller '04

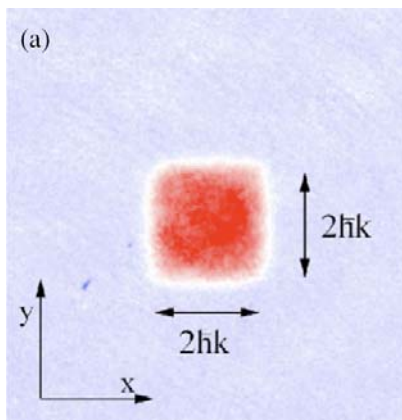


# Transport in dipole traps

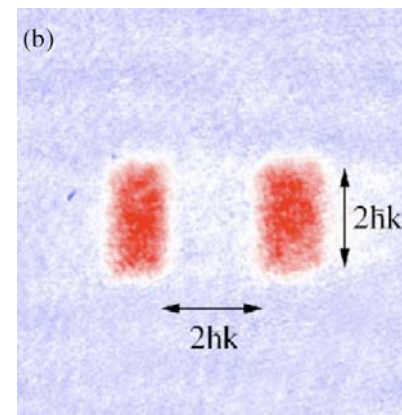


© T. Porto, W. Phillips 2005

Realization of (not time-optimized) transport in an optical lattice



few ms transfer time

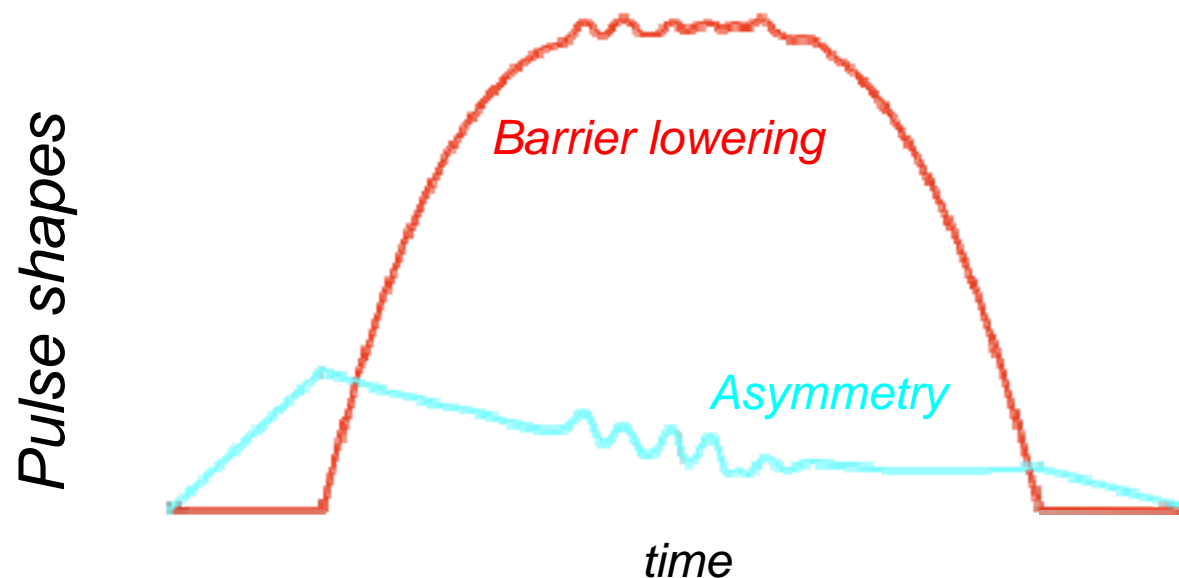


...two-qubit gate: W. Phillips, Nature 2007



# Optimized pulses

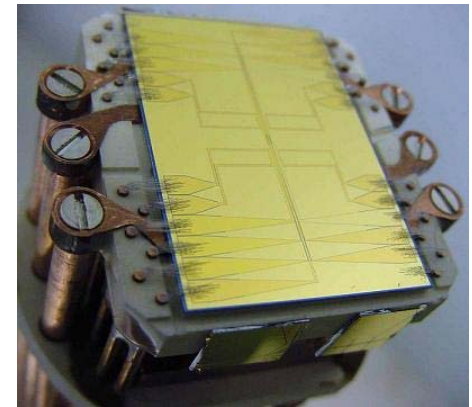
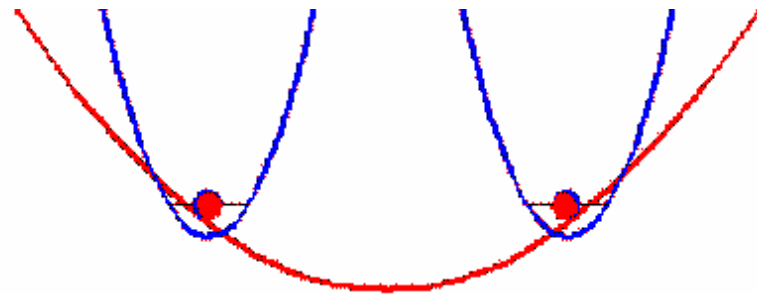
- Optimization algorithm introduces wiggles in pulse shapes
- “Shaking” helps exciting-deexciting
- Frequency higher than gate operation rate



# Switching gates on atom chips

with R. Folman, J. Schmiedmayer '00

- Modulated magnetic field yields state-dependent potential
- State-independent electrostatic attraction switches off the barrier
- Logical phase accumulated at each collision



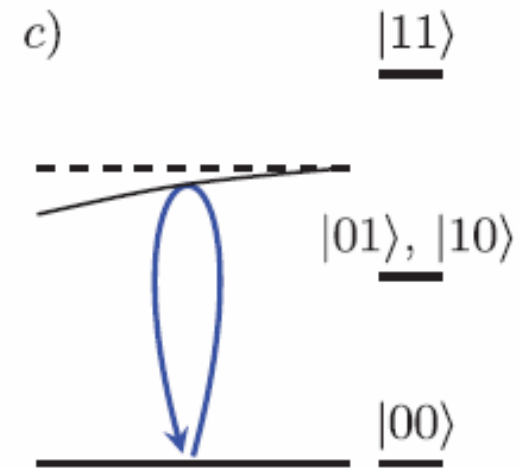
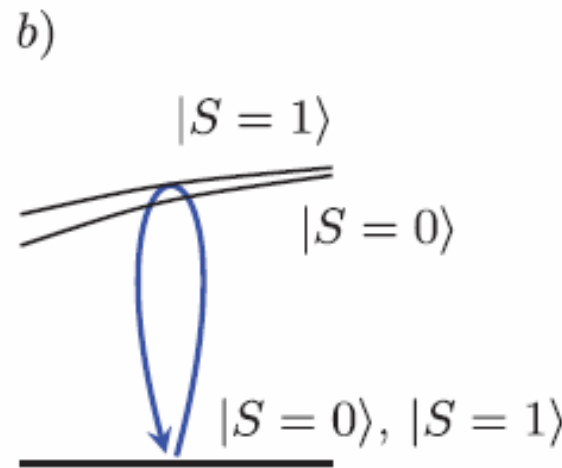
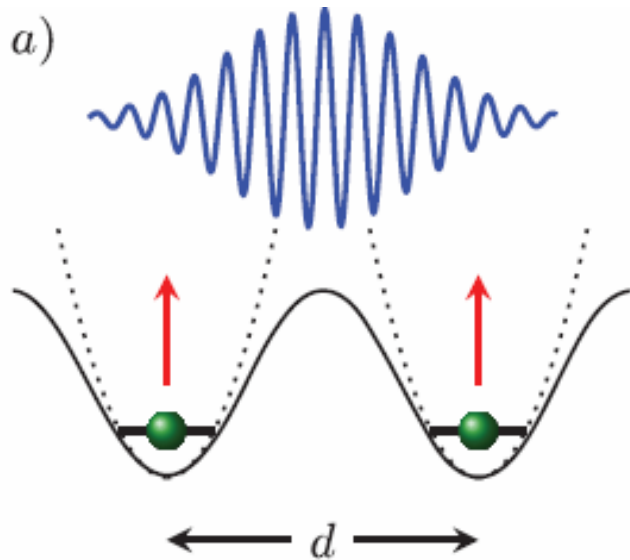
© R. Folman, J. Schmiedmayer



© J. Reichel

# Ultrafast optical gate

ongoing collaboration with C. Koch



$$\begin{aligned}
 |\uparrow\rangle_1 |\uparrow\rangle_2 &= |S=1, m=1\rangle \\
 |\uparrow\rangle_1 |\downarrow\rangle_2 &= \frac{|S=1, m=0\rangle + |S=0, m=0\rangle}{\sqrt{2}} \\
 |\downarrow\rangle_1 |\uparrow\rangle_2 &= \frac{|S=1, m=0\rangle - |S=0, m=0\rangle}{\sqrt{2}} \\
 |\downarrow\rangle_1 |\downarrow\rangle_2 &= |S=1, m=-1\rangle
 \end{aligned}$$

$$\frac{1}{2} \begin{pmatrix} 2 & 0 & 0 & 0 \\ 0 & 1 + e^{i\phi} & 1 - e^{i\phi} & 0 \\ 0 & 1 - e^{i\phi} & 1 + e^{i\phi} & 0 \\ 0 & 0 & 0 & 2 \end{pmatrix}$$

$$\phi = \frac{\pi}{2} \longrightarrow \sqrt{SWAP}$$

# What can go wrong

and how to fix it

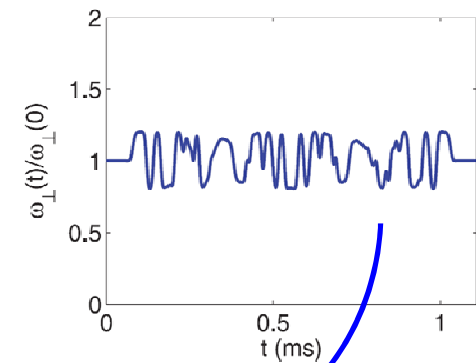
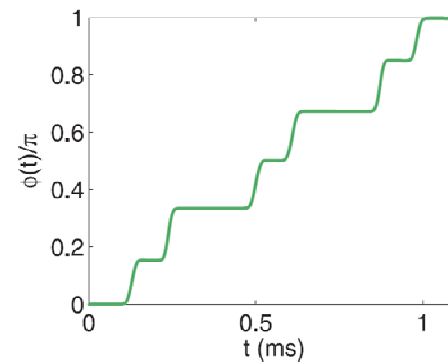
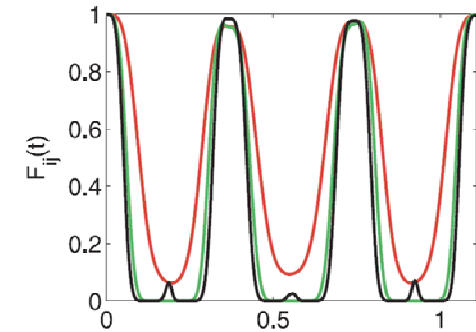
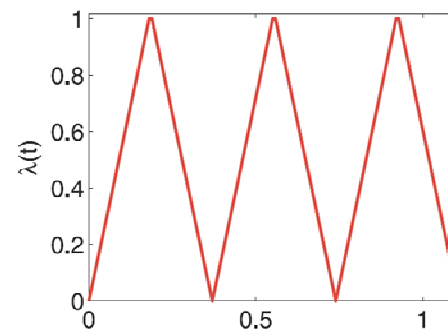
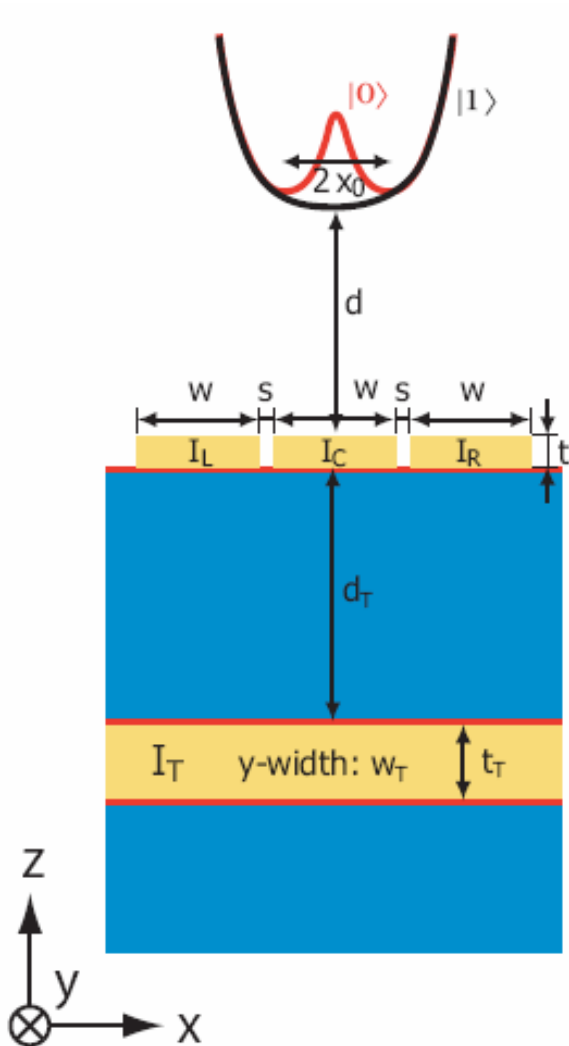
# What can go wrong?

- Anharmonicity in the trapping potentials
- Noise in the control parameters
- Limited bandwidth
- Imperfect pulse calibration
- Leakage
- Finite temperature
- Inhomogeneous broadening
- ...decoherence

Anharmonicity  
in the trapping potential

# Microwave pulse shaping on atom chips

with J. Reichel, T. Hänsch '06



Transverse frequency modulation compensates for interaction-induced wavepacket distortion

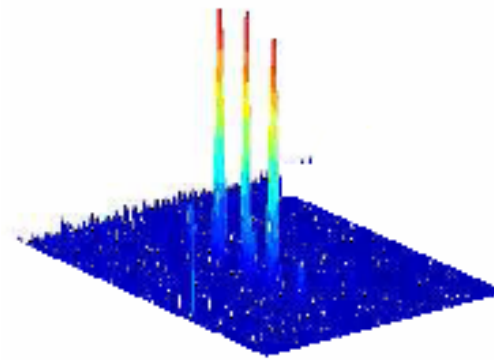
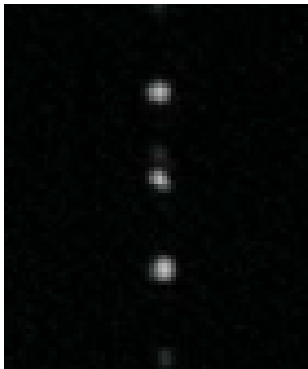
Noise in the control parameters



# Noise in dipole traps

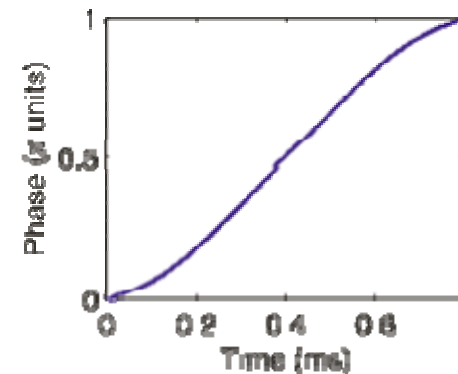
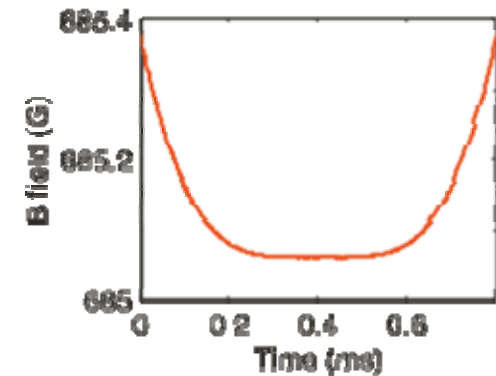
with P. Grangier '05

Holographic tweezer setup



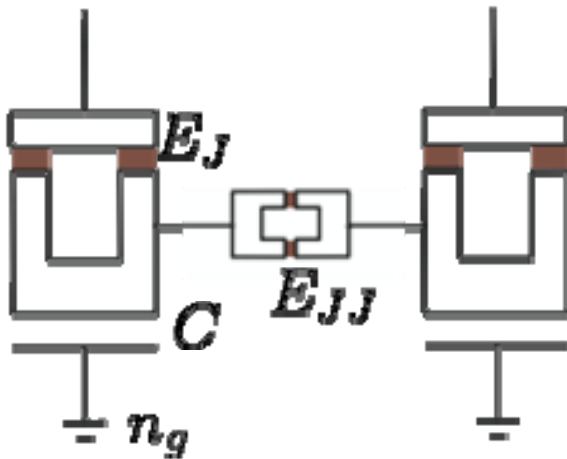
Typical noise for laser intensity and position around  $10^{-2}$

Fidelity decrease  $\sim 10^{-3}$  due to timescale separation between noise and control



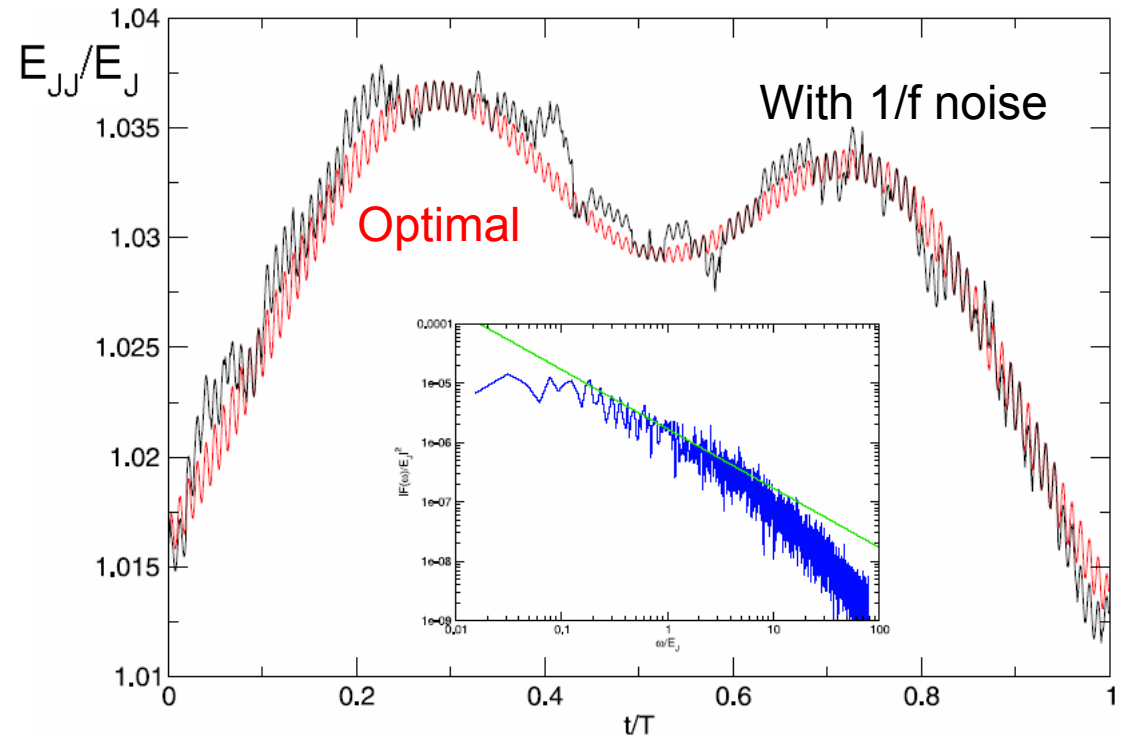
# Noise in Josephson charge qubits

with R. Fazio '07



$$G_{JJ} = \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & \pm i & 0 & 0 \\ 0 & 0 & \pm i & 0 \\ 1 & 0 & 0 & 0 \end{pmatrix}$$

- Qubit: 0 or 1 excess Cooper pair
- Control parameter: Josephson energy  $E_{JJ}$



# Error with/without control

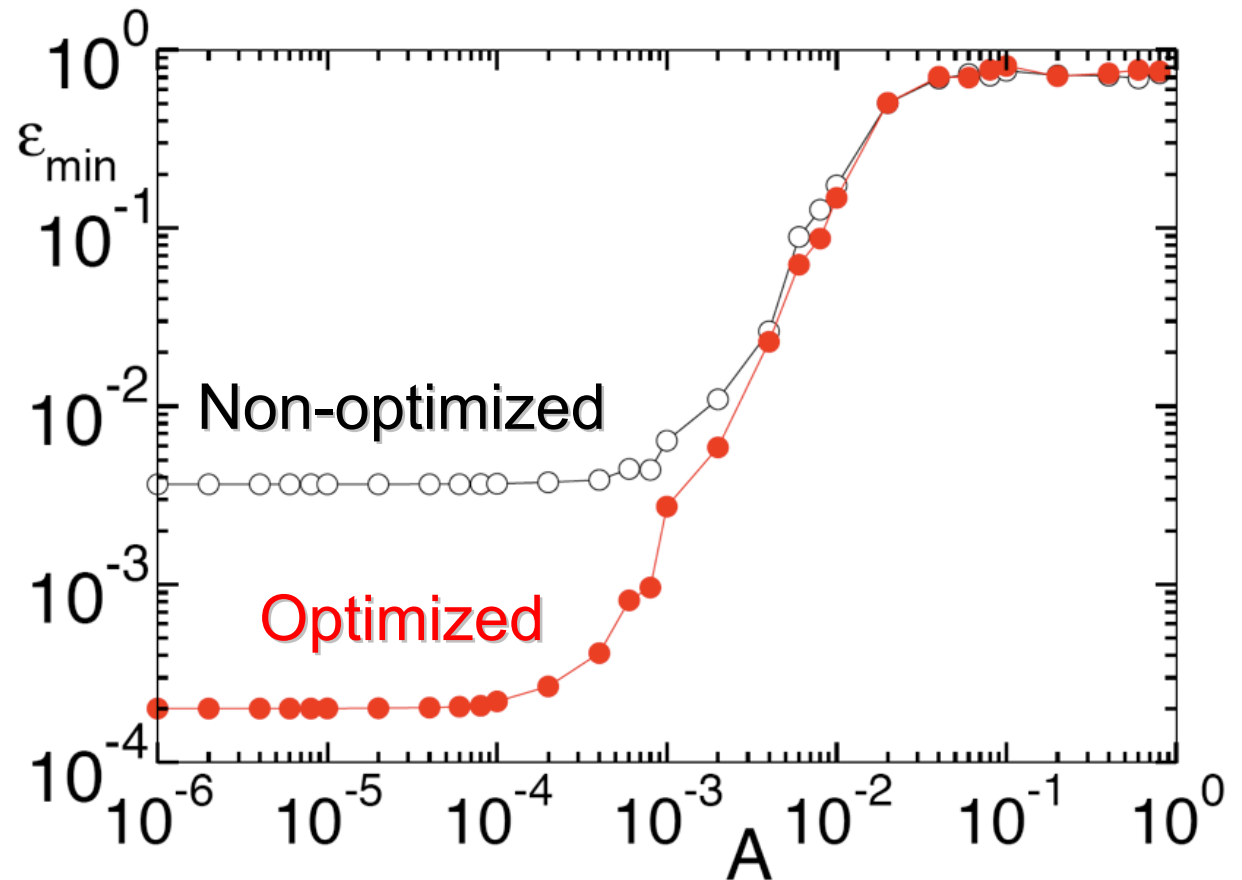
1/f noise

$$S(\omega) \propto A/\omega$$

Typical exp. values

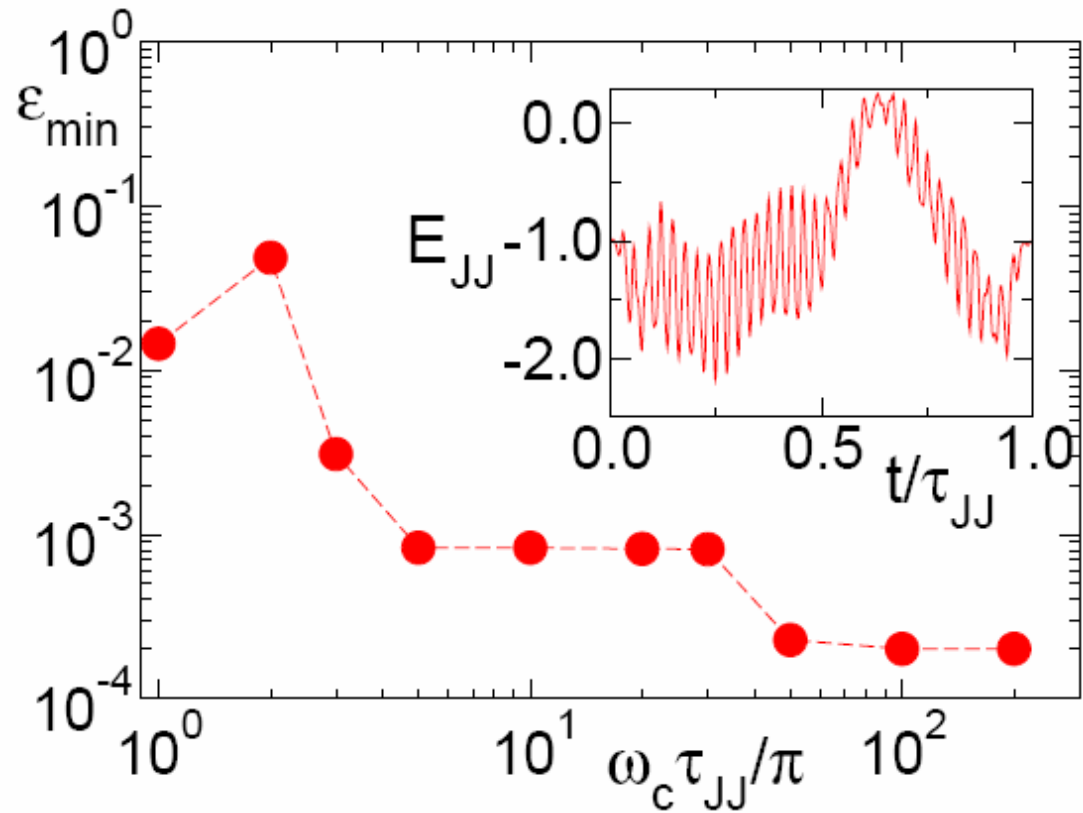
$$A \sim 10^{-5}$$

Fault tolerance  
with realistic  
noise?



Limited bandwidth

# Limited bandwidth in JJ gates

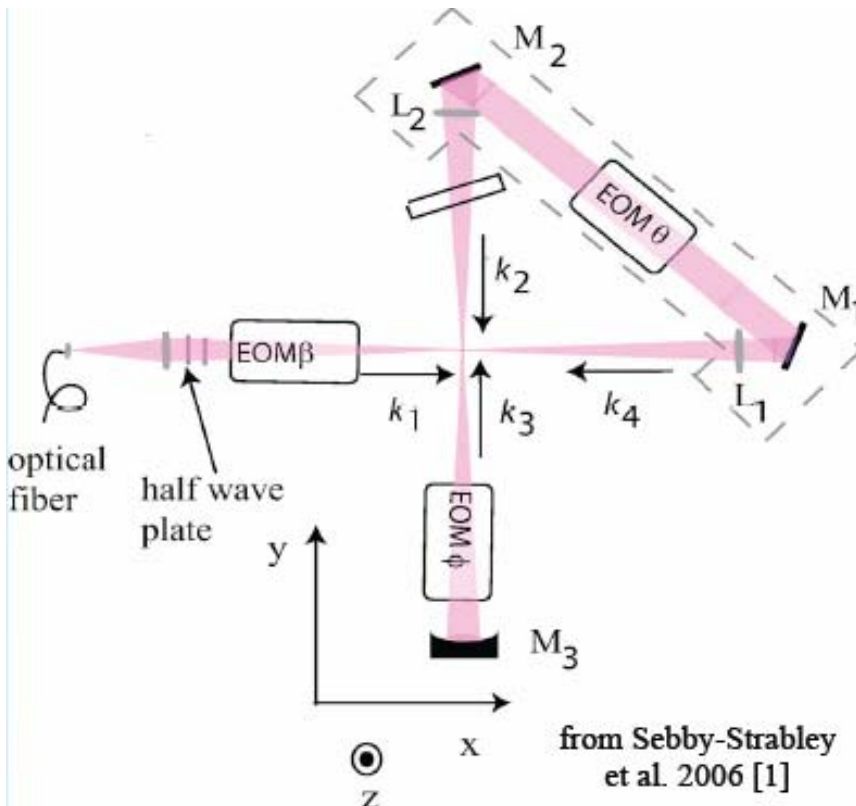
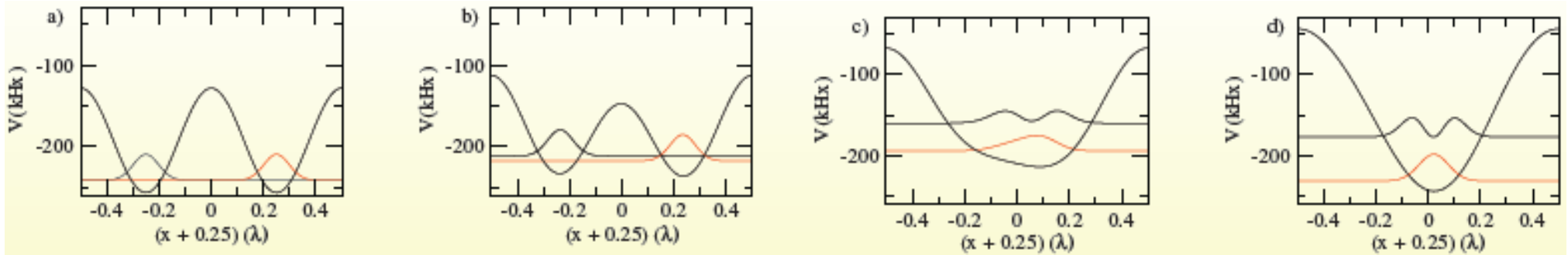


Pulse shape constrained in Fourier space

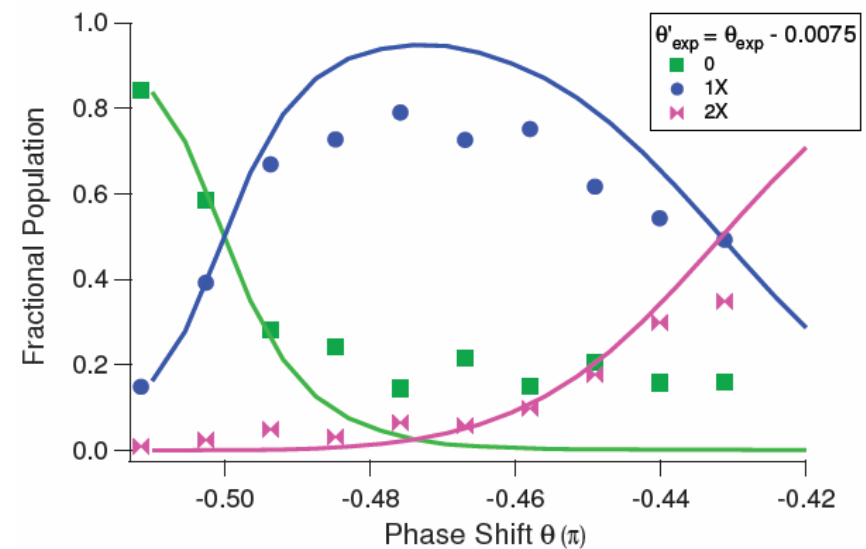
Imperfect pulse calibration

# Transport in a real lattice

with T. Porto, W. Phillips '08



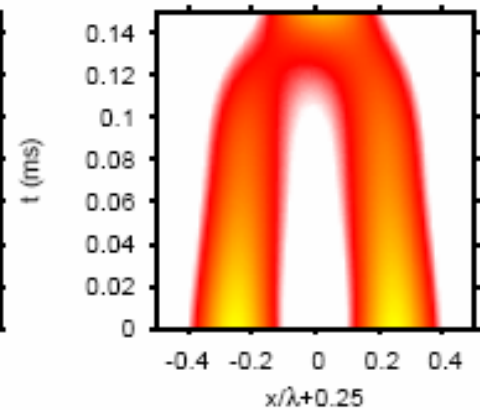
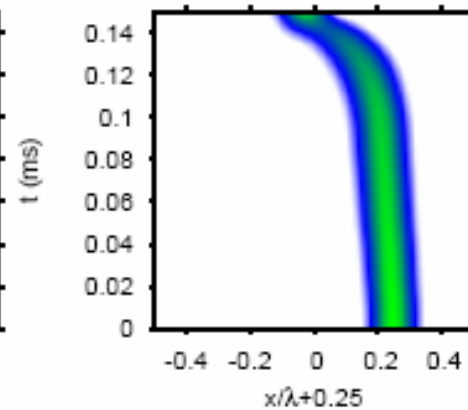
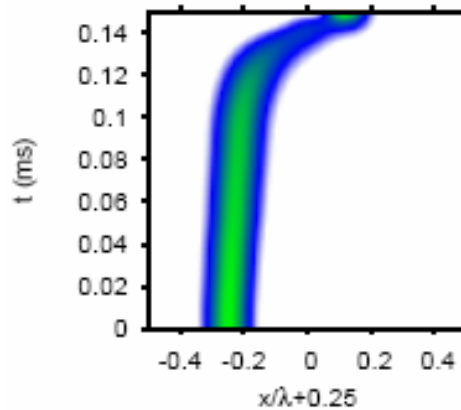
## Calibration of the control simulation



# Optimization results

Without  
optimization:  
 $F=0.22$

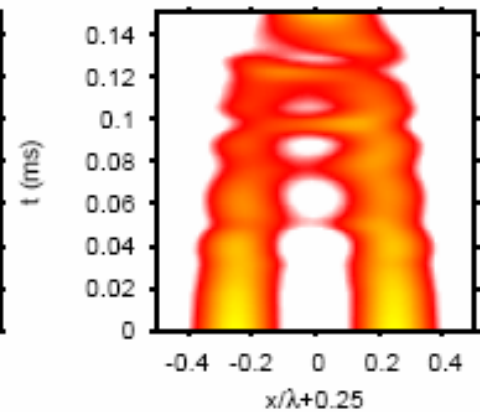
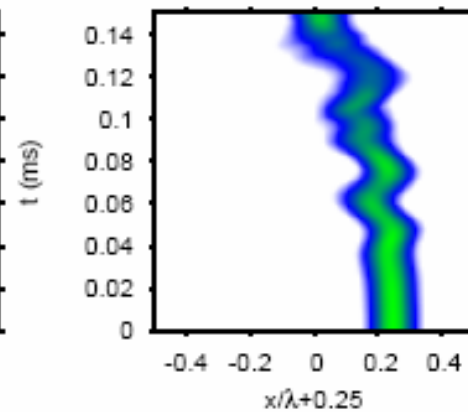
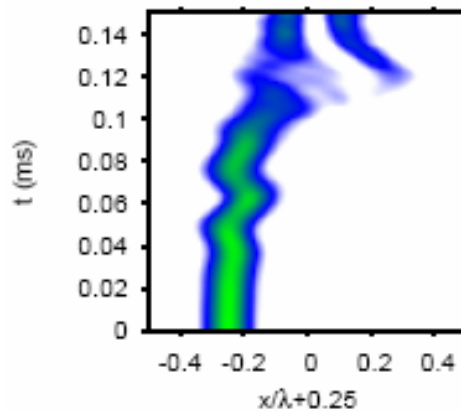
$T=150 \mu\text{s}$



Wavefunctions

Potential

With  
optimization:  
 $F=0.97$

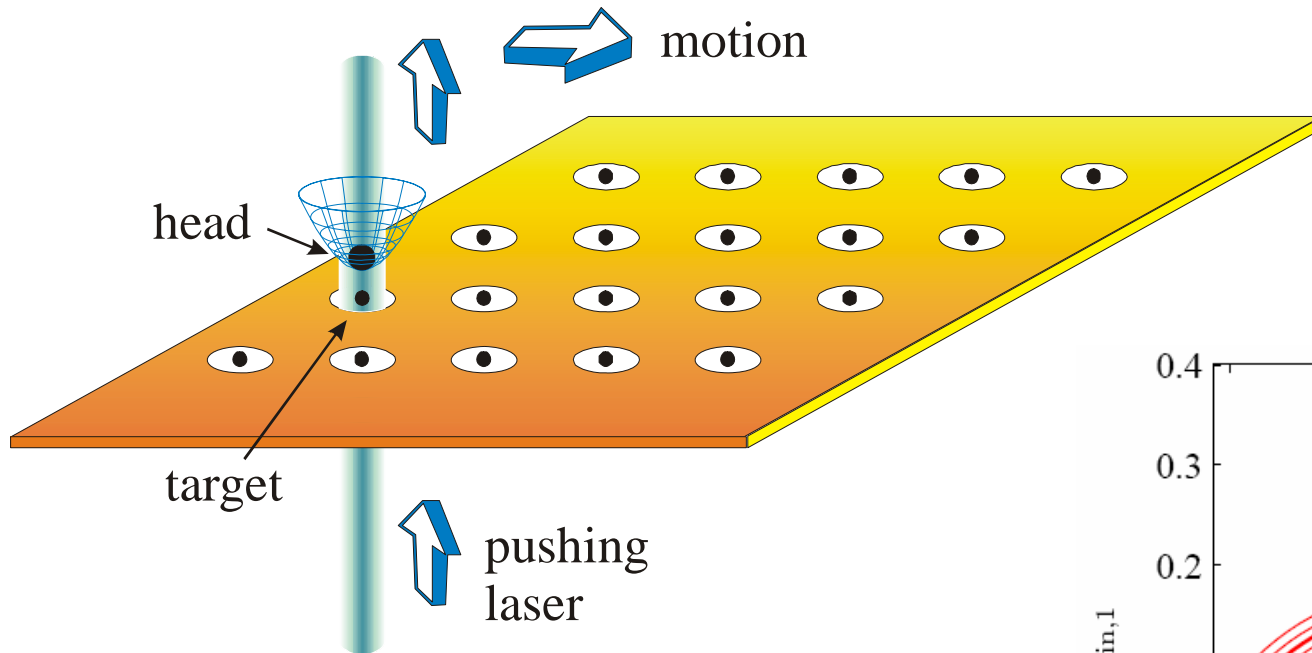




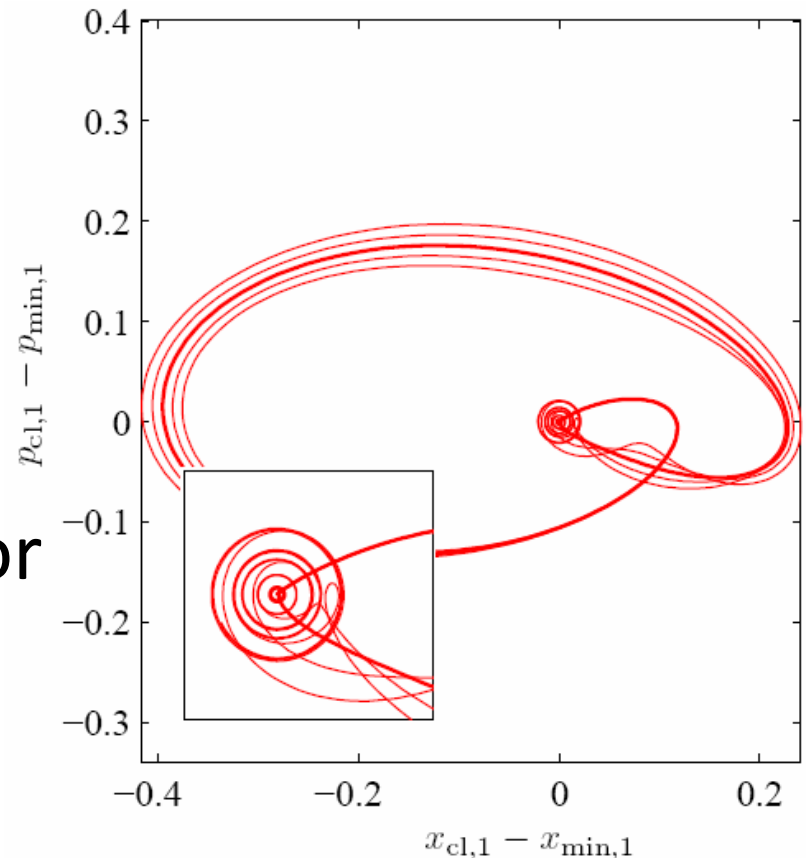
Leakage

# Leakage in ion-pushing gates

with D. Tannor '09

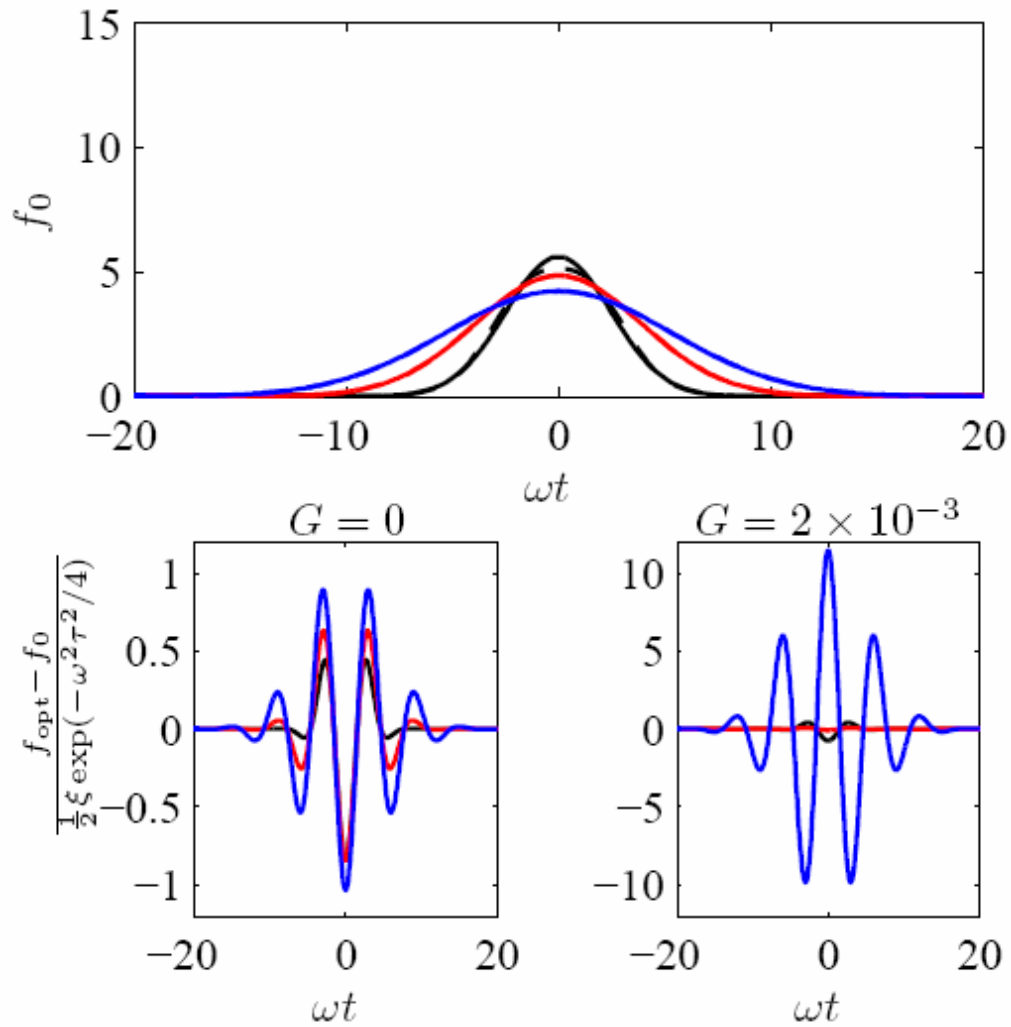


Phase-space ion trajectories for inhomogeneous pushing force

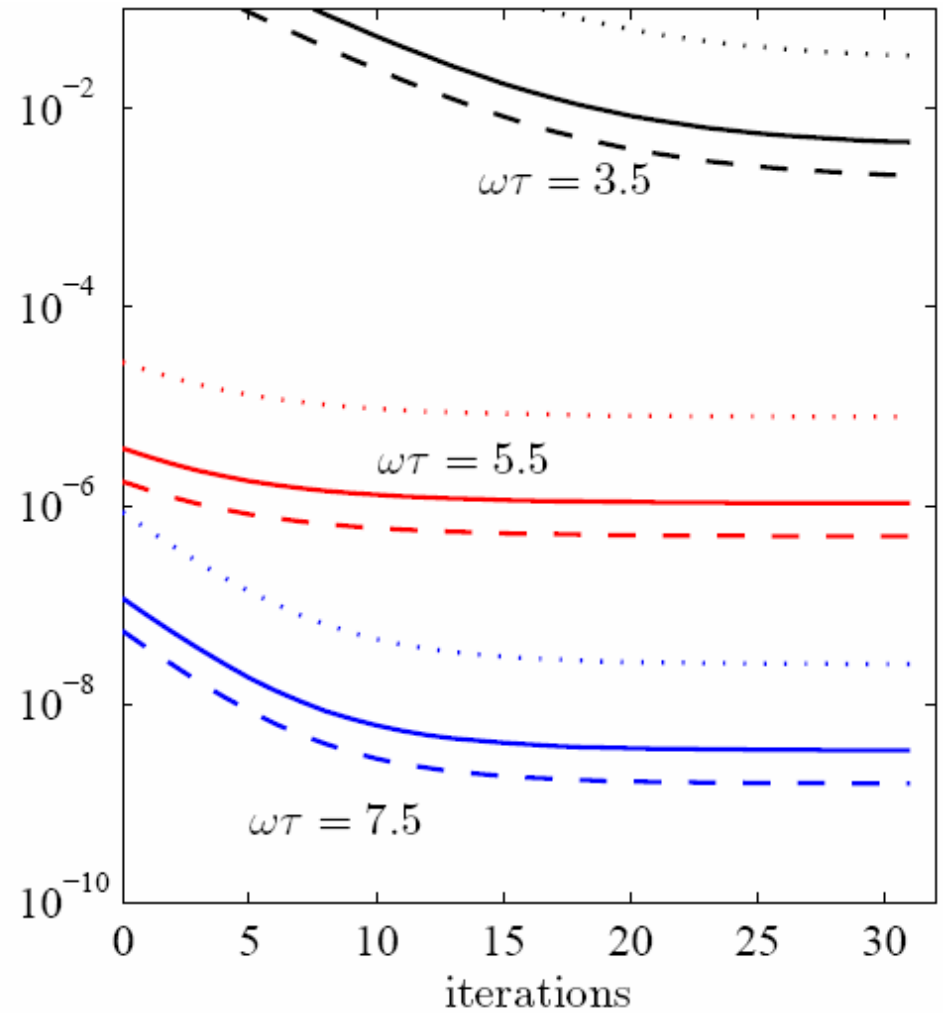


# Krotov optimization

## Pulses



## Errors

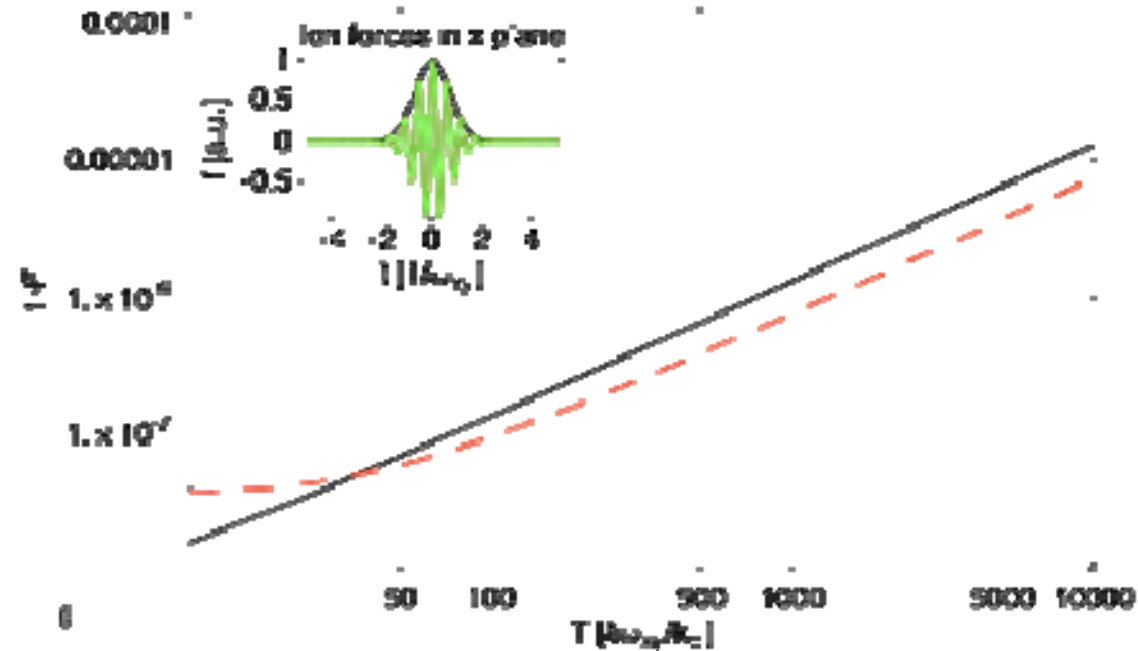
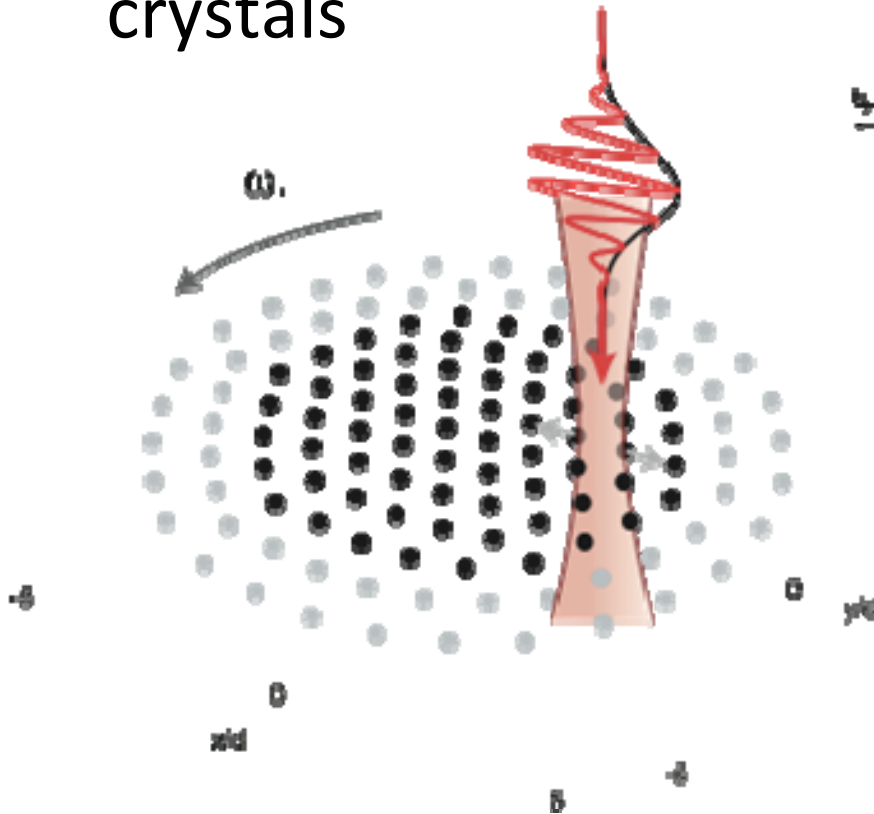


Finite temperature

# Thermal noise in ion Wigner crystals

with J. Taylor '07

Quantum gates  
among ions in 2D  
crystals



Fast-carrier-  
modulation gate:  
decoupling from soft  
phonon modes

# Inhomogeneous broadening

# A broadening model

with N. Khaneja '08

Simple HO transport problem

$$V = \frac{m\omega^2}{2} [x - d(t)]^2$$

... when you move it, you make mistakes

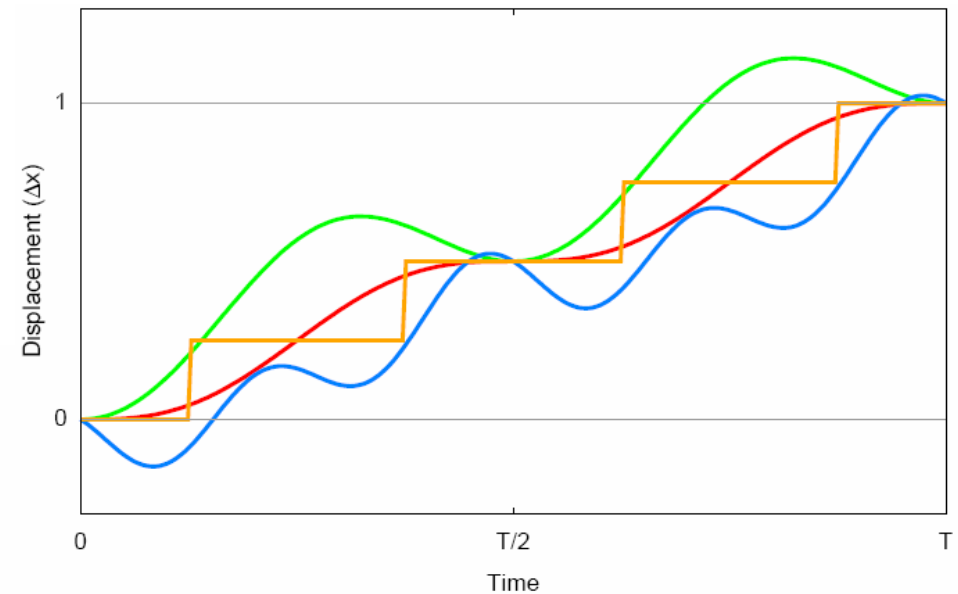
$$d(t)$$

— Optimal path  $d(t)$

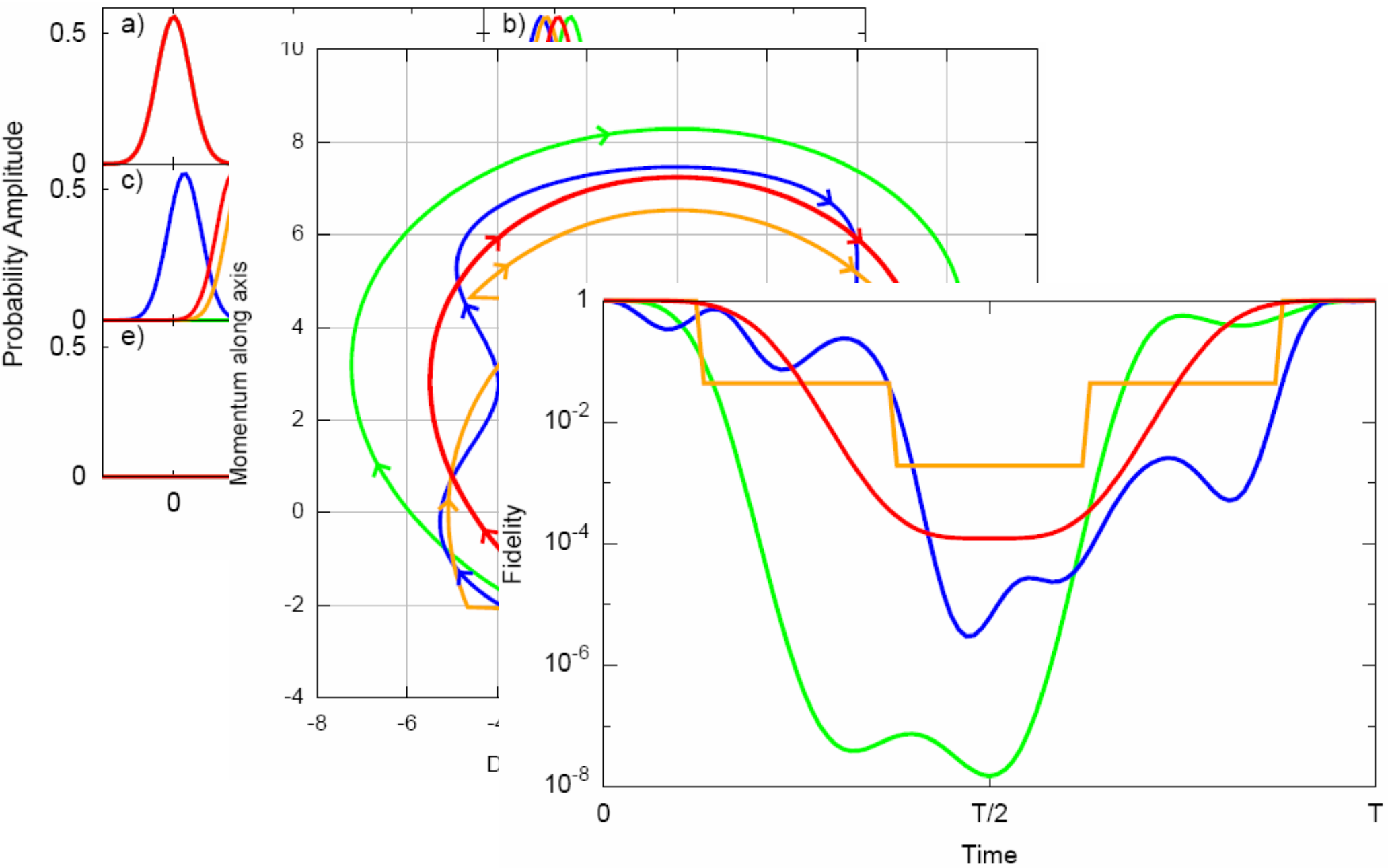
—  $d(t) + \dot{d}(t)$

—  $d(t) + \alpha_0 + \sum_{n=1}^2 (\alpha_1 \cos [2n\pi \frac{t}{T}] + \alpha_2 \sin [2n\pi \frac{t}{T}])$

—  $d(t)$  piecewise



# Distorted transport results





# How much can we push?

The Quantum Speed Limit

# Quantum Speed Limit

The speed at which a quantum state evolves is linked to the dynamics of the Hamiltonian.

$$E = \langle \Psi | H | \Psi \rangle$$

Initial energy                      Initial state

$$\Delta E = \sqrt{\langle \Psi | (H - E)^2 | \Psi \rangle}$$

Energy variance

Minimum time required for a quantum state to evolve to an orthogonal state

$$T_{\min}(E, \Delta E) \equiv \max \left( \frac{\pi \hbar}{2E}, \frac{\pi \hbar}{2\Delta E} \right)$$

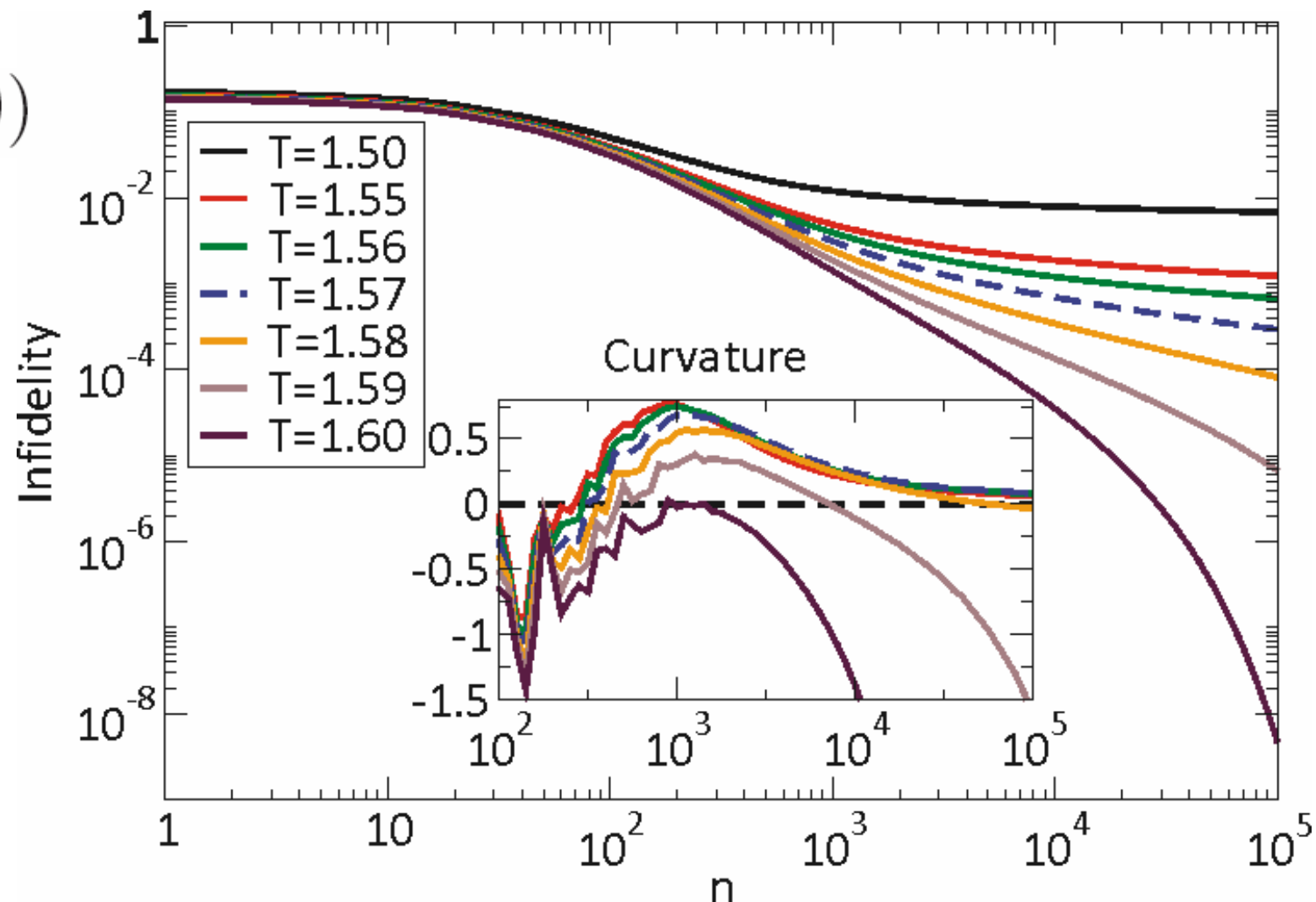
# Toy model: Landau-Zener crossing

$$H[\Gamma(t)] = \begin{pmatrix} \Gamma(t) & \omega \\ \omega & -\Gamma(t) \end{pmatrix}$$

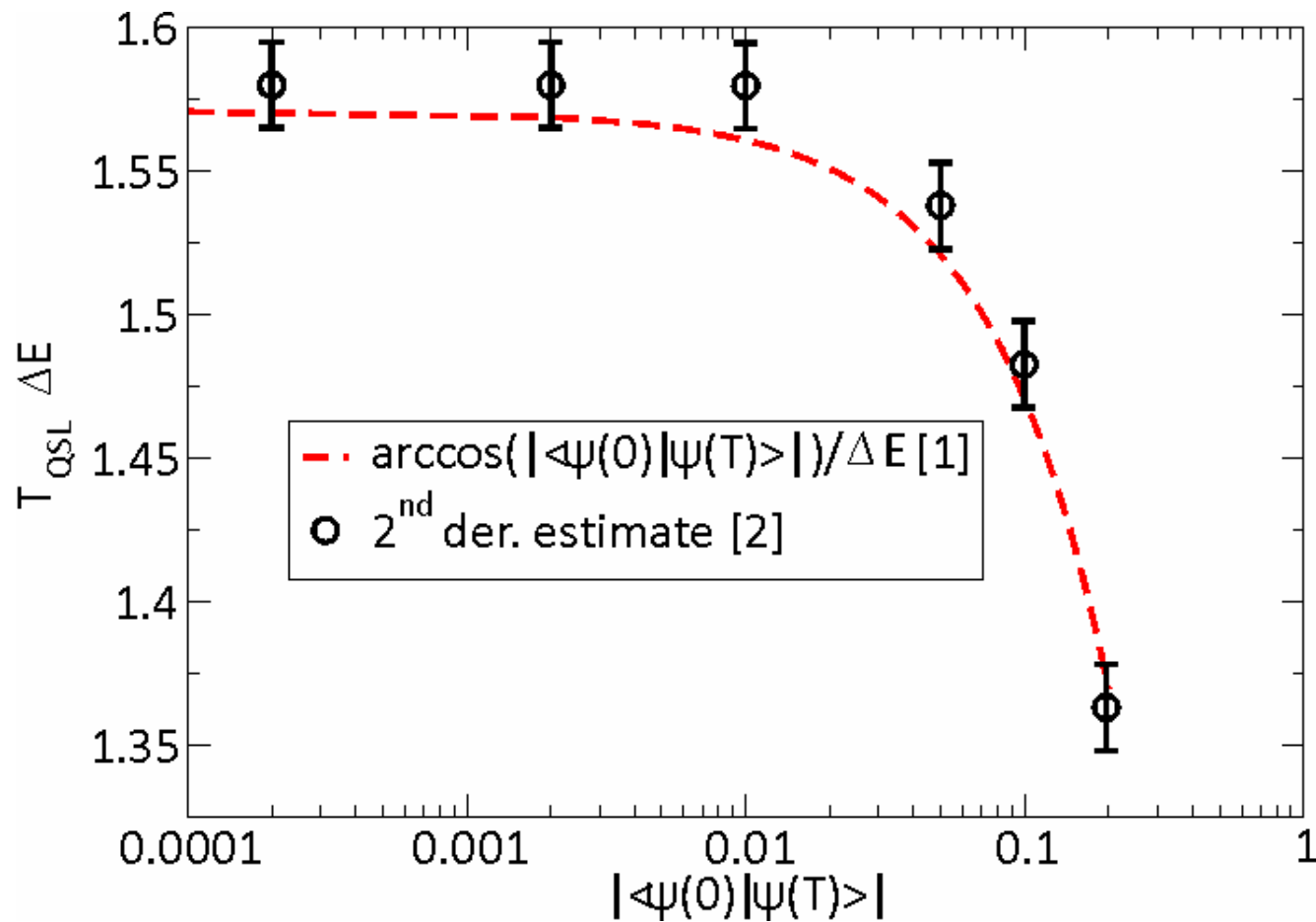
with R. Fazio '09

$$\Gamma(T) = -\Gamma(0)$$

Goal: back to ground state after crossing



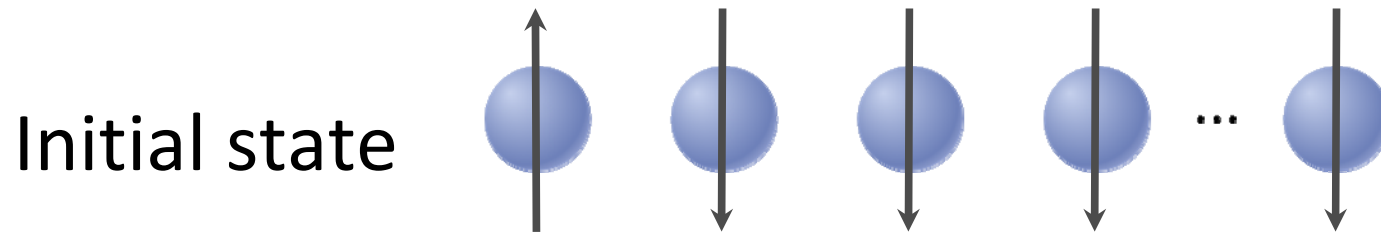
# Heuristic vs. analytic QSL



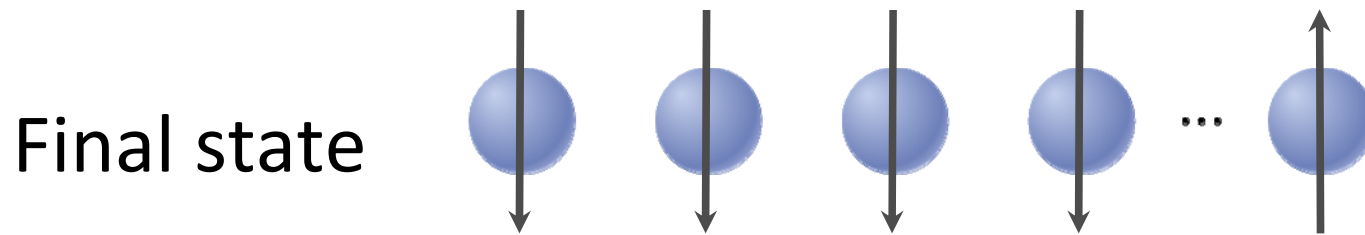
[1] K. Bhattacharyya, J. Phys. A: Math. Gen. 16, 2993 (1983).

[2] T. Caneva, M. Murphy, TC, R. Fazio, S. Montangero, V. Giovannetti, G. E. Santoro, arXiv:0902.4193

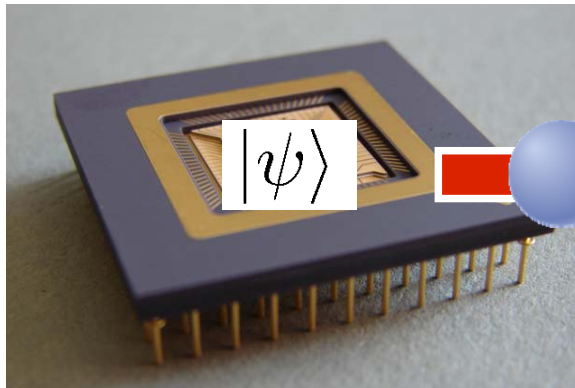
# Spin chain transport



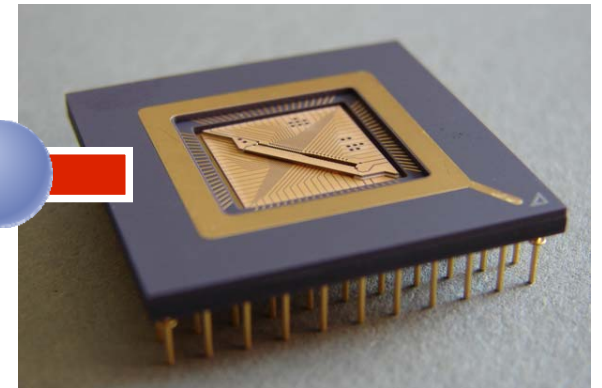
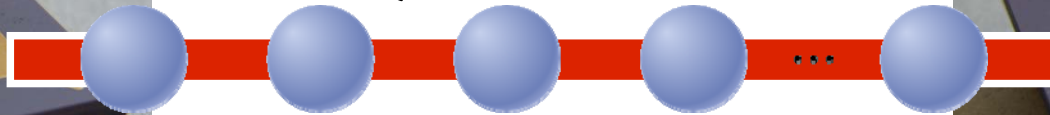
$$|\psi(0)\rangle = (\alpha|0\rangle + \beta|1\rangle) \otimes |00\dots\rangle$$



$$|\psi(0)\rangle = |00\dots\rangle \otimes (\alpha|0\rangle + \beta|1\rangle)$$



Quantum wire

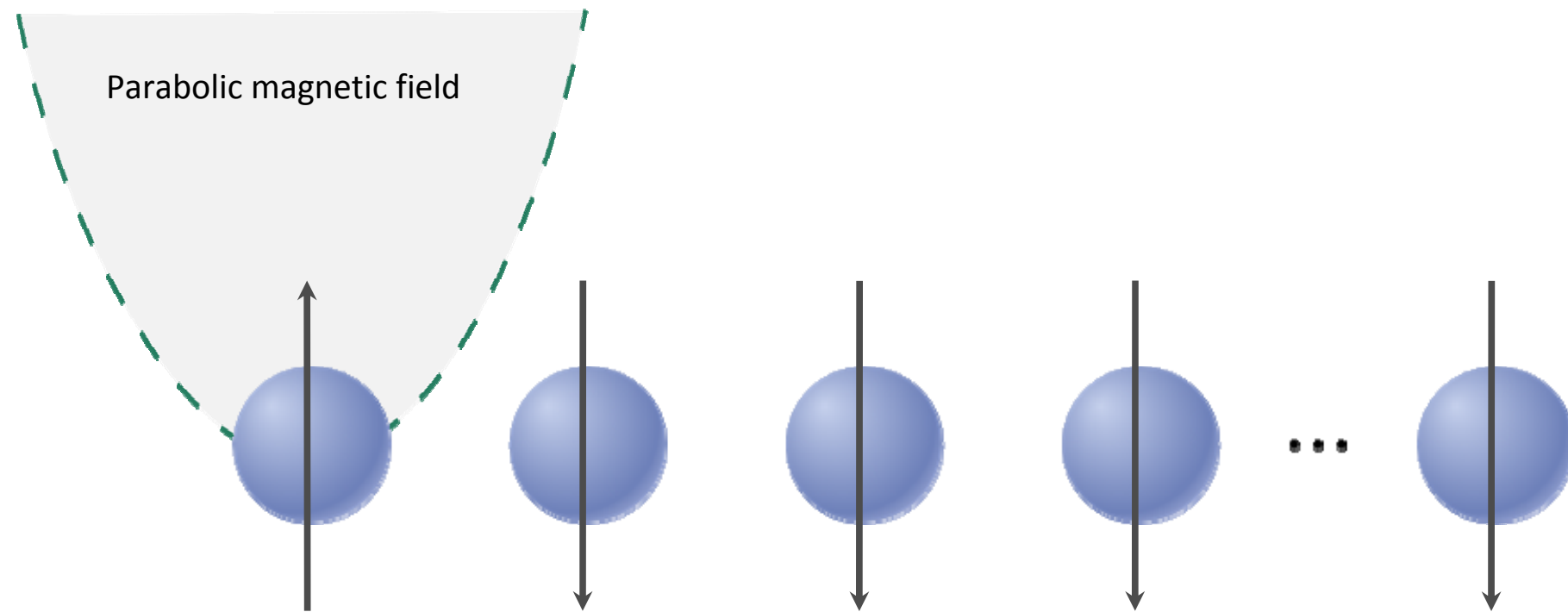


S. Bose, PRL **91** 207901 (2003)

# The transport mechanism

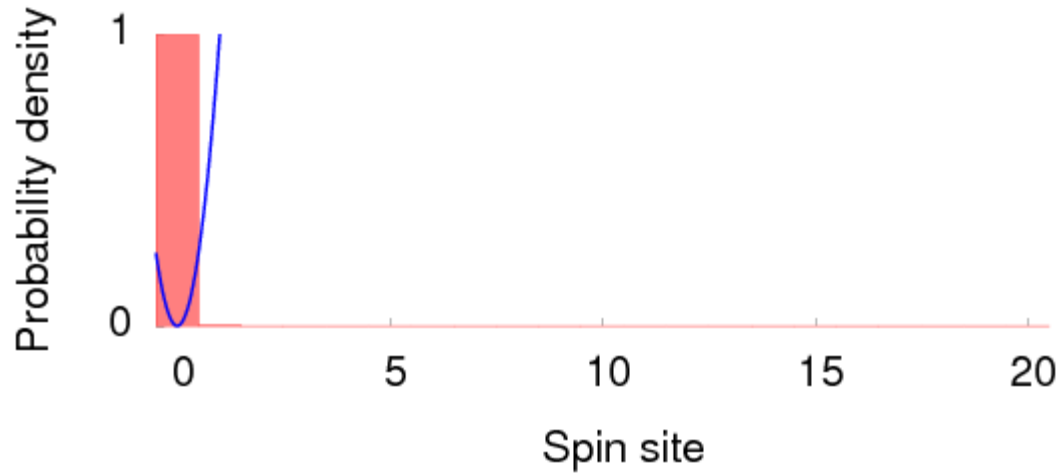
$$H = -\frac{J}{2} \sum_{n=0}^{N-2} \vec{\sigma}_n \cdot \vec{\sigma}_{n+1}$$

spin – spin coupling  
(nearest neighbour)



# The transport mechanism

No problem adiabatically,



But if we try naively to go faster...

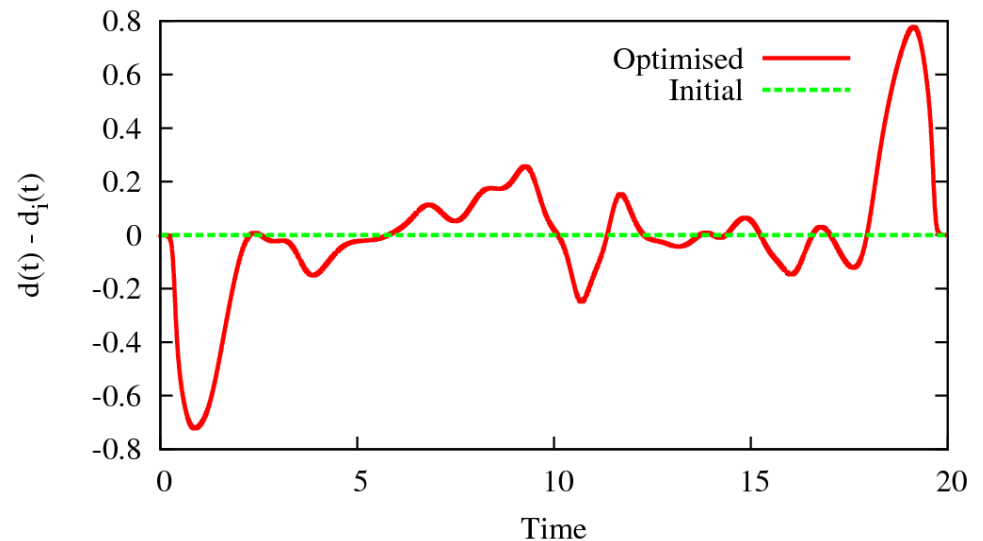
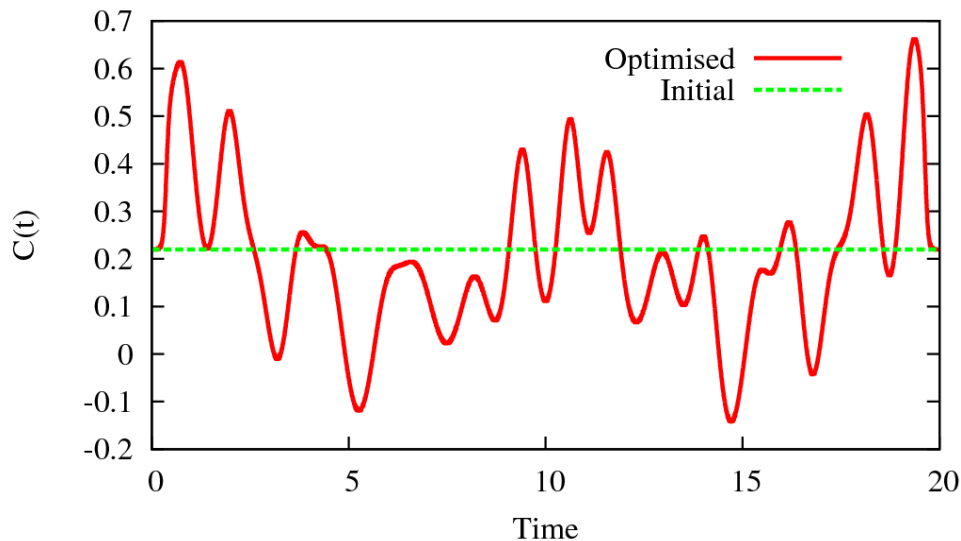
# Optimal control formulation

The Heisenberg Hamiltonian:

$$H = -\frac{J}{2} \sum_{n=0}^{N-2} \vec{\sigma}_n \cdot \vec{\sigma}_{n+1} + \sum_{n=0}^{N-1} \frac{C(t)}{2} (n - d(t))^2 \sigma_n^z,$$

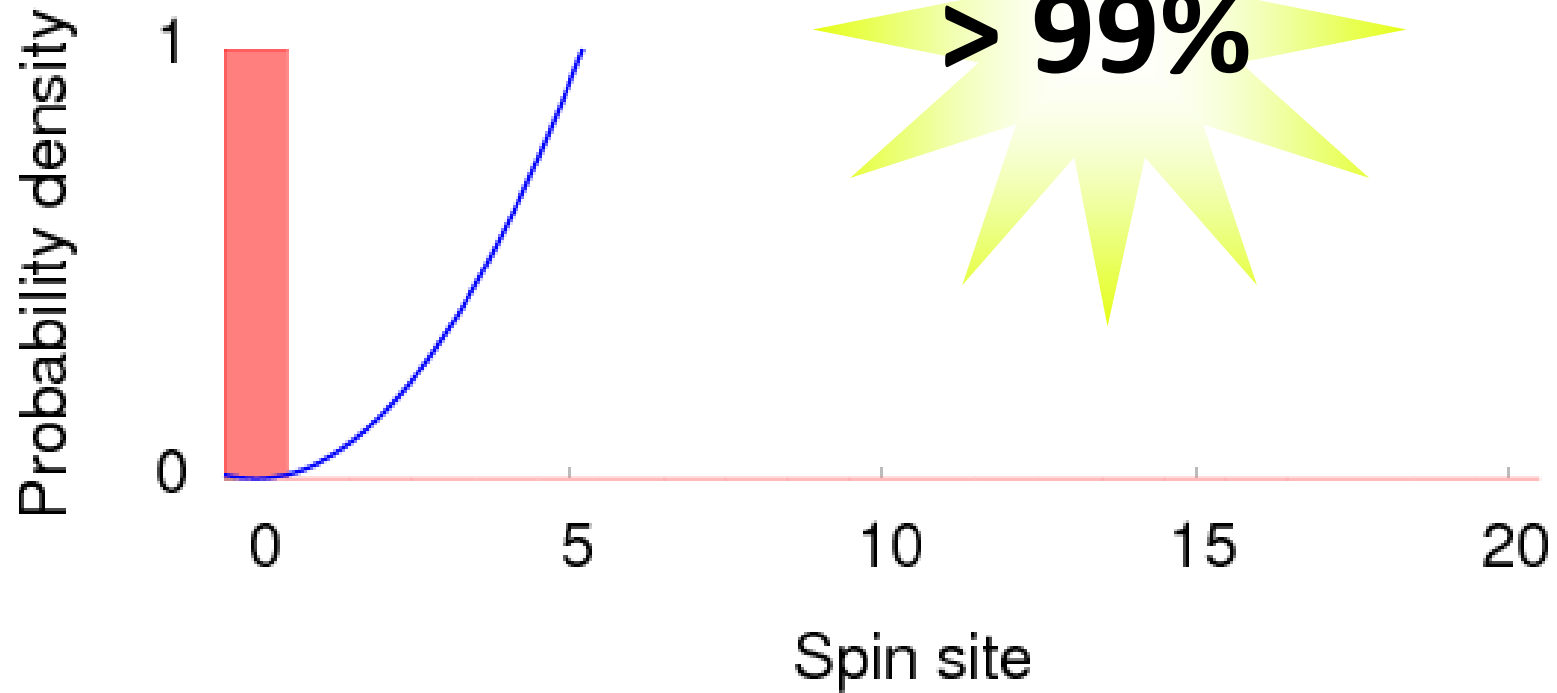
control parameters

Use the Krotov optimisation algorithm to increase the transfer fidelity



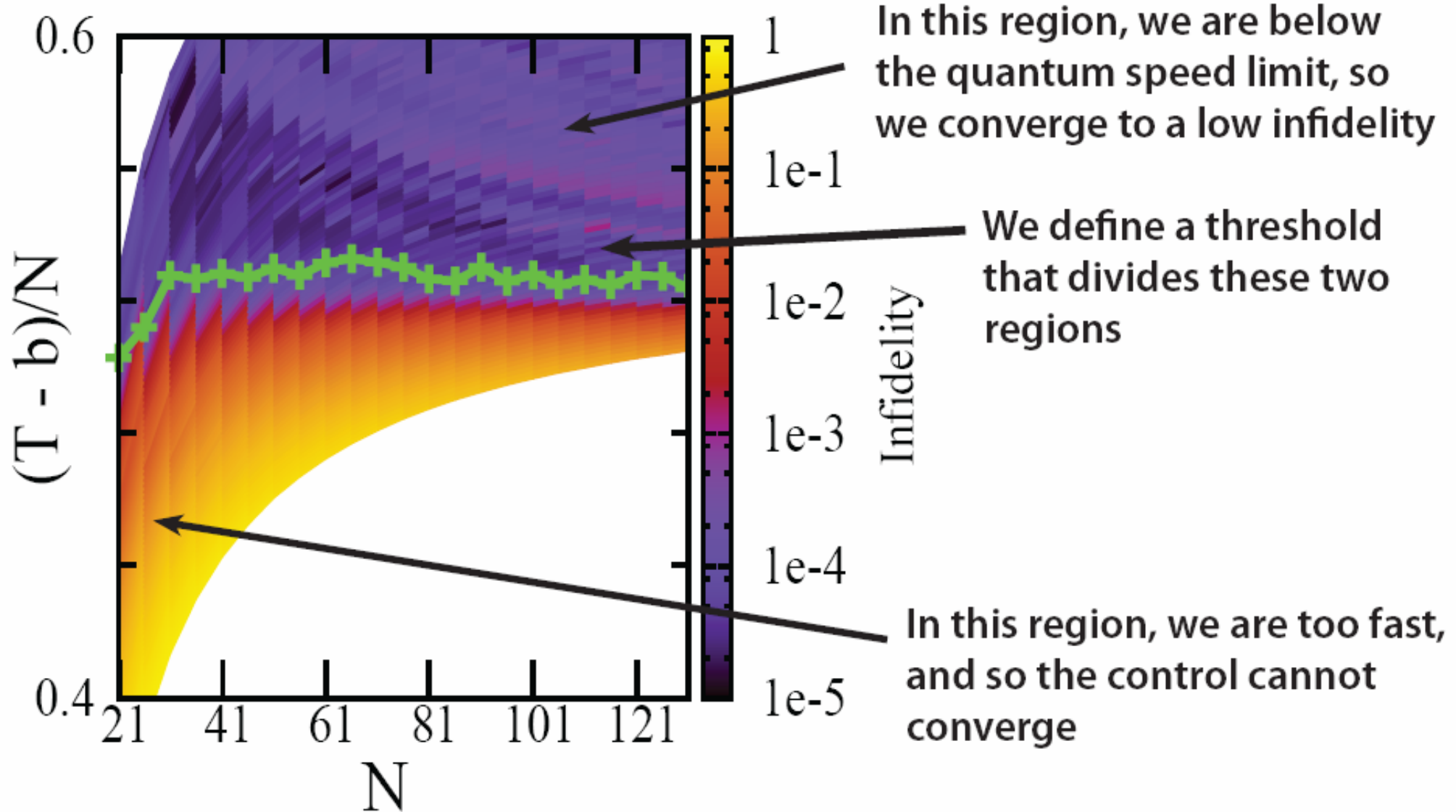


# Optimised Transport

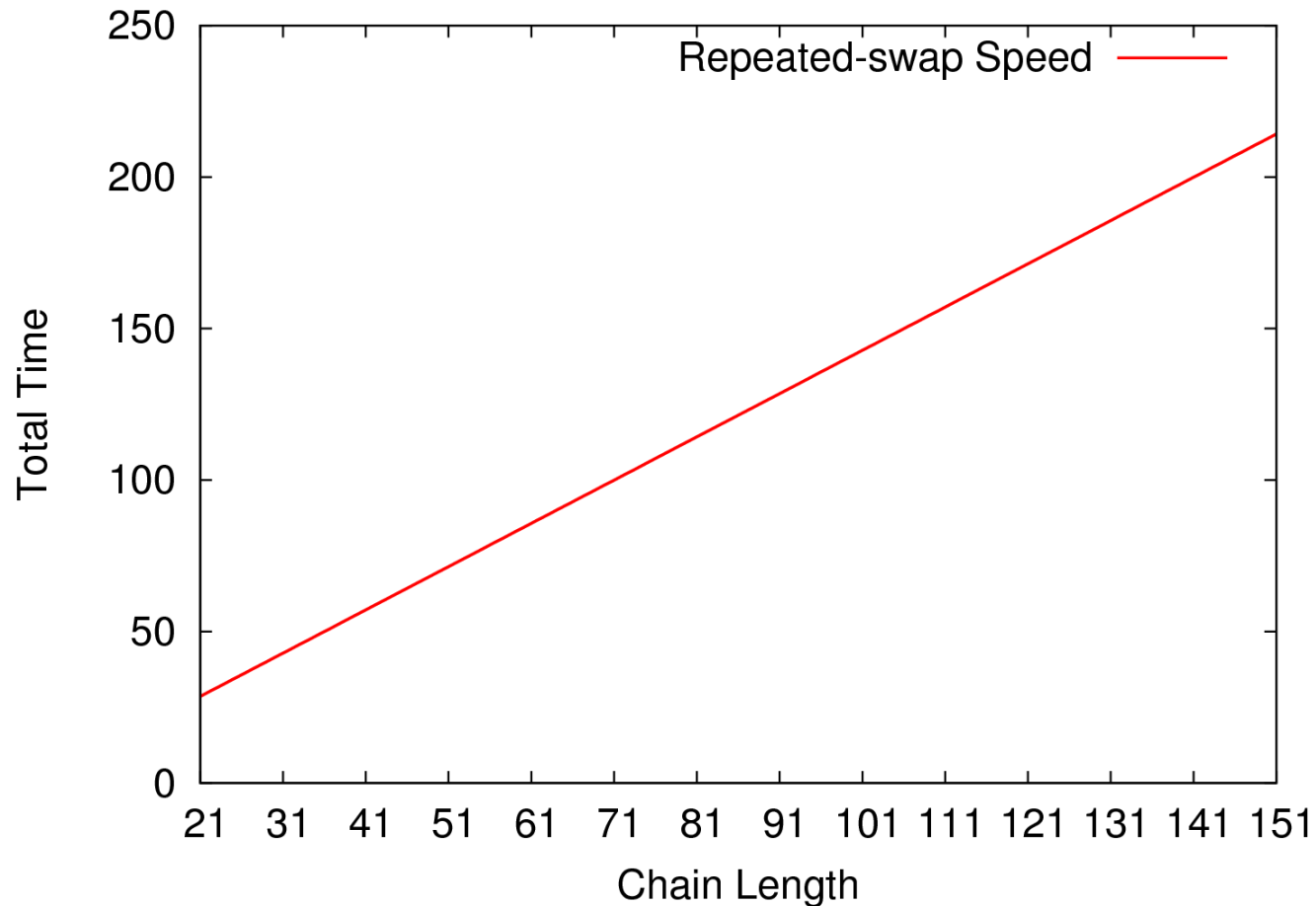


This is  $\sim 200$  times faster

# Final transport infidelity



# Phenomenological Model

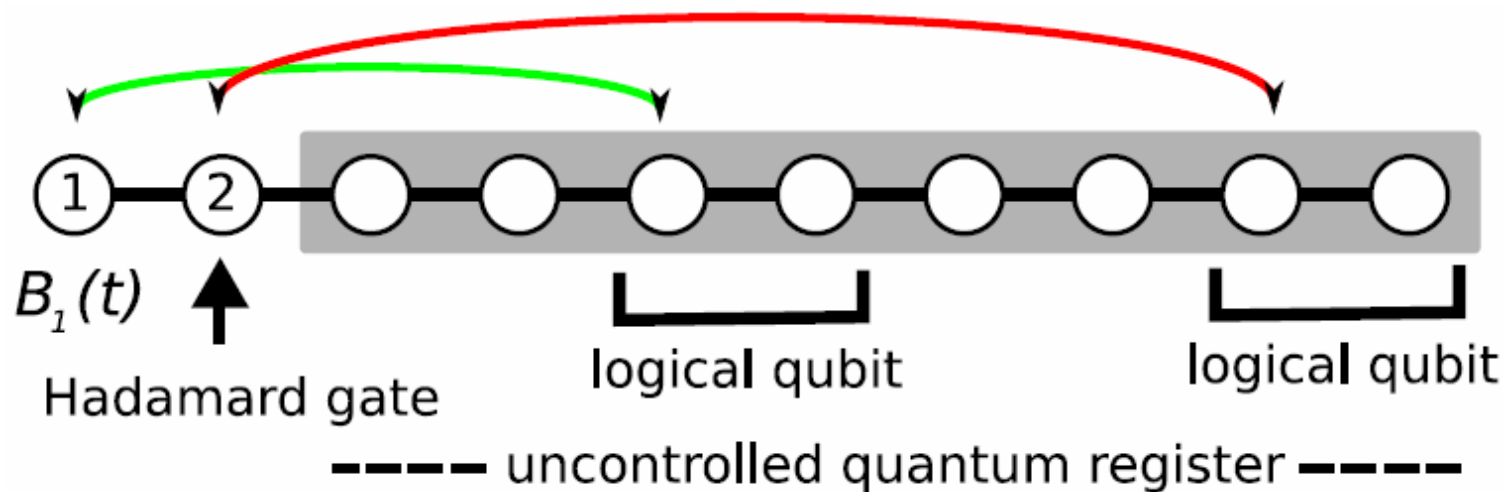


Theory: 
$$\tau_{\text{QSL}} \equiv \max \left\{ \frac{\pi \hbar}{2J}, \frac{\pi \hbar}{2\Delta\mathcal{E}} \right\} \quad \Delta\mathcal{E} = \frac{1}{T} \int_0^T \Delta E_i(t) dt$$

Actual optimization speed limits is **lower**

# Can this be used for computing?

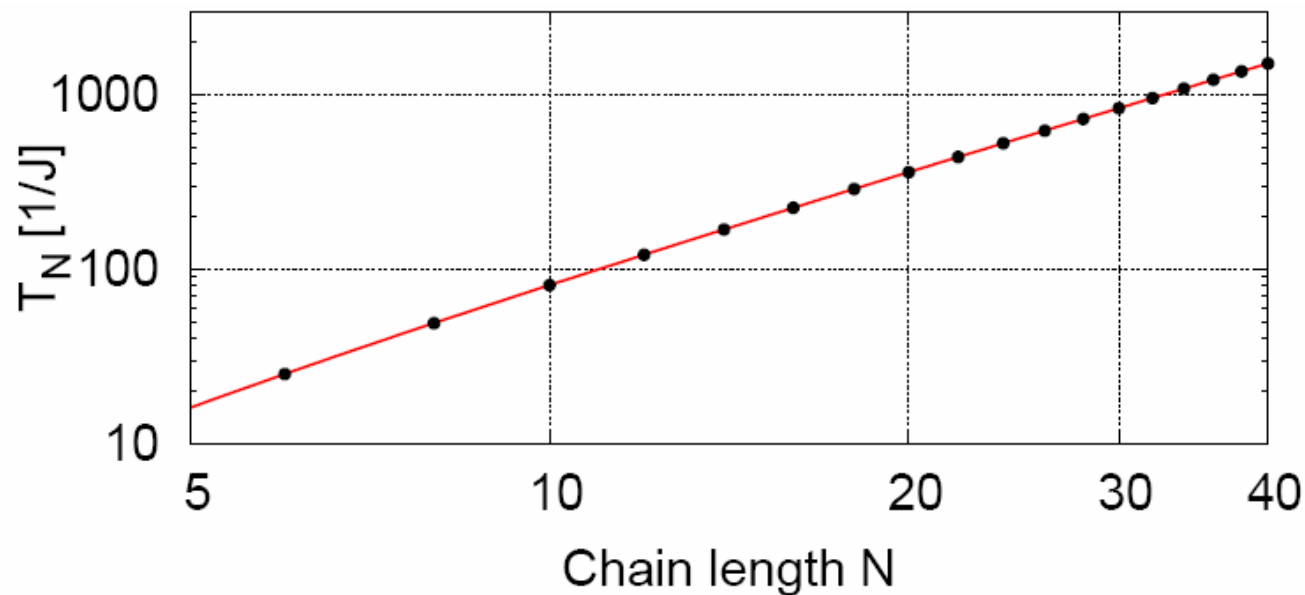
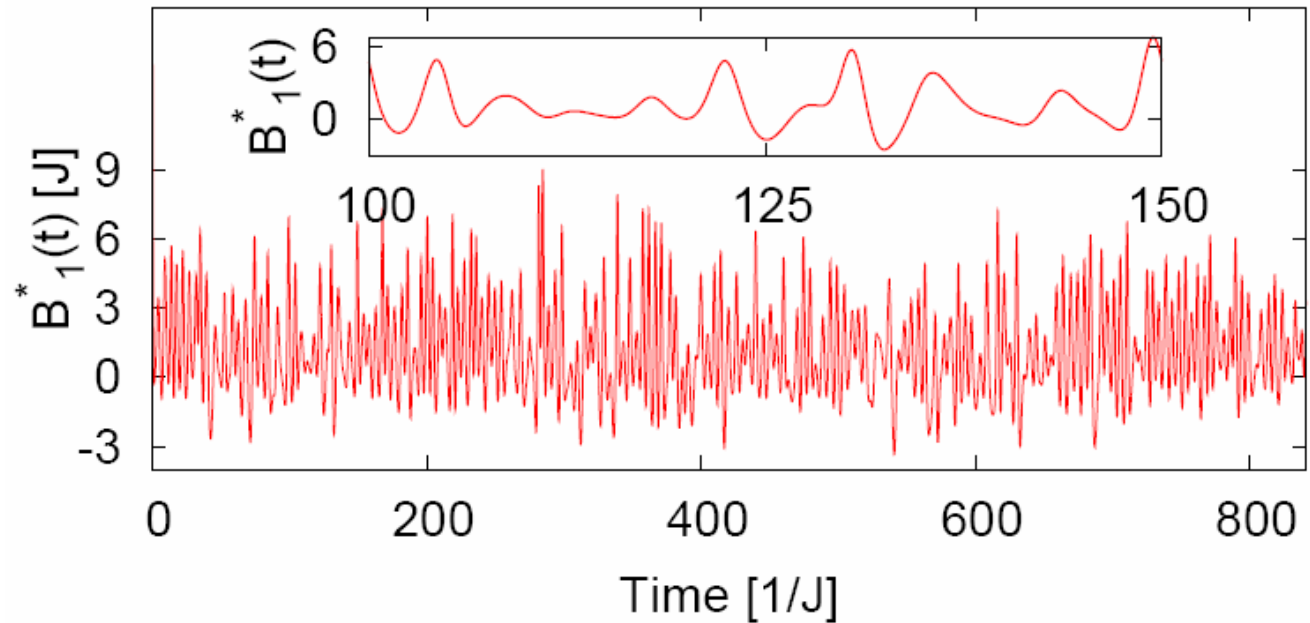
Scalable quantum computation via local control of only two qubits



D. Burgarth, K. Maruyama, M. Murphy, S. Montangero,  
T. Calarco, F. Nori, M. Plenio, arXiv:0905.3373

# Scaling of the operation time

Sample  
control  
pulse



$$T_N = (N - 1)^2$$

# Conclusions

- Quantum optimal control does work for quantum information processing
- It allows fixing a range of real issues
- Its limits deserve further exploration

# Work done with...

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- C. Koch (Berlin)
- R. Fazio (Pisa)
- P. Grangier (Orsay)
- T. Hänsch (Munich)
- P. Julienne, W. Phillips (Gaithersburg)
- M. Lukin (Harvard)



- D. Tannor (Weizmann)...and with no military funding.
- P. Zoller (Innsbruck)