

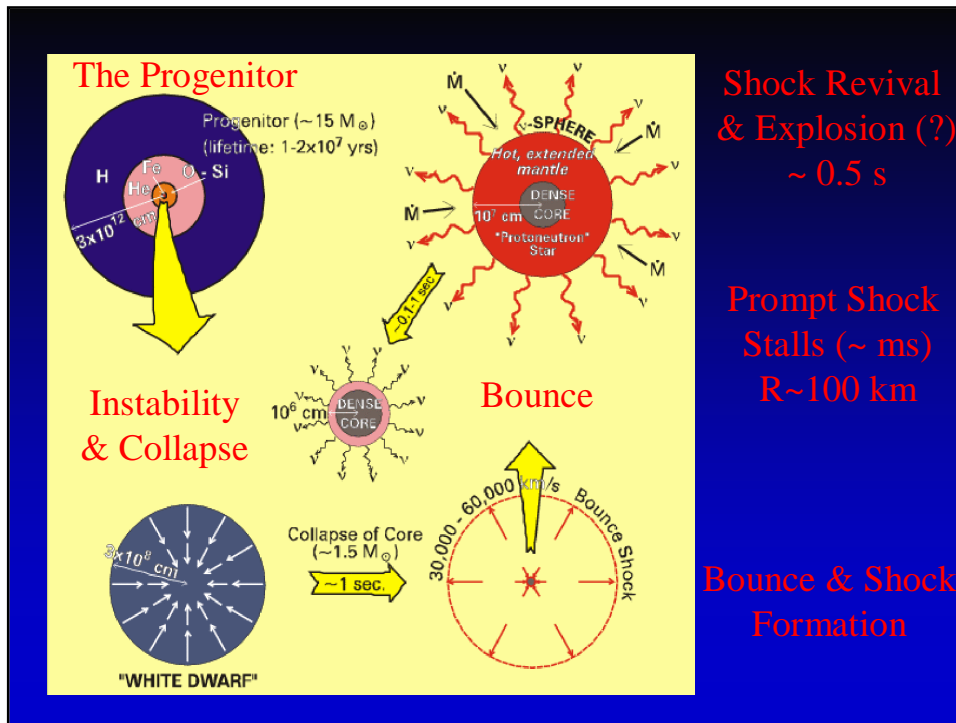
Supernovae, Neutrino-Driven Winds, and Nucleosynthesis

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Outline

- **Introduction: The Supernova Problem**
 - Scales, Important Processes, & Outstanding Questions
- **Dynamical Models of Core-Collapse**
 - Breakout, Spectra, & Detector Signatures
- **Protoneutron Star Winds**
 - Neutrino-Driven Outflows & Nucleosynthesis
 - Magnetic Fields
- **Summary, Conclusions, & The Future**



- **What is a Core-Collapse Supernova?**
 - An explosion initiated by a dynamical instability
 - The death of a massive star
 - The birth cry of a neutron star or black hole
 - An agent of chemical and dynamical galactic evolution
- **Characteristics of Supernovae:**
 - Total kinetic energy: $\sim 10^{51}$ erg
 - Luminous energy radiated: $\sim 10^{49}$ erg
 - Temperatures: 0.1 – 50 MeV (10 MeV $\sim 10^{11}$ K)
 - Densities: $10^9 - 6 \times 10^{14}$ g cm $^{-3}$
- **Neutrino Production:**
 - Neutrino energy radiated: $\sim 2-3 \times 10^{53}$ erg
 - Core is opaque to neutrinos of all flavors!

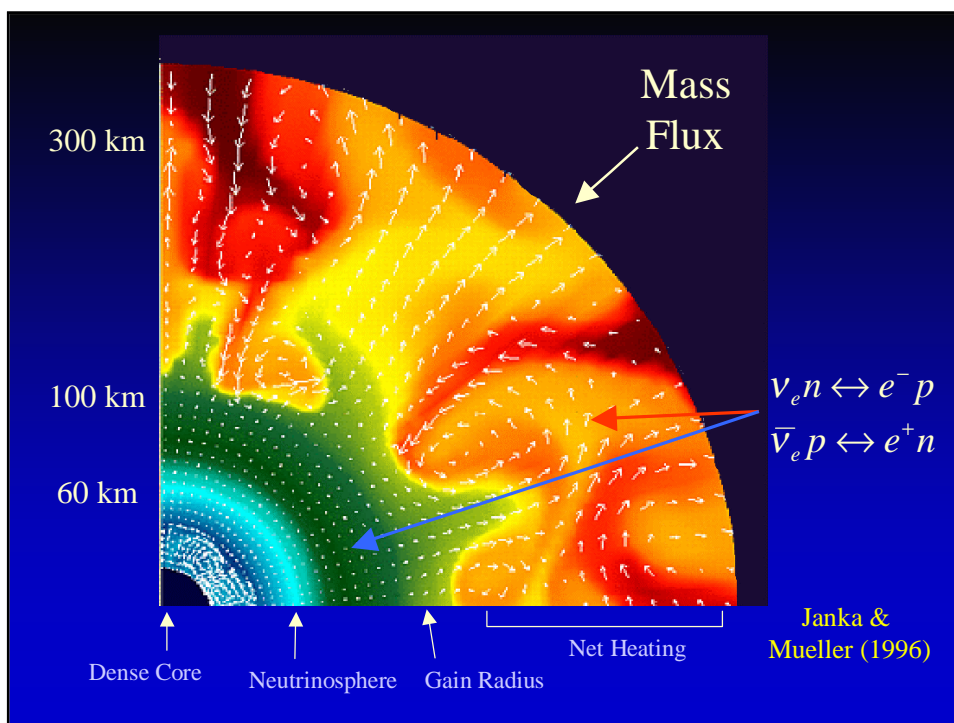
$$\tau_{\text{Diff}} \approx 1 \text{ s}$$

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(Burrows, Hayes, & Fryxell 1995)

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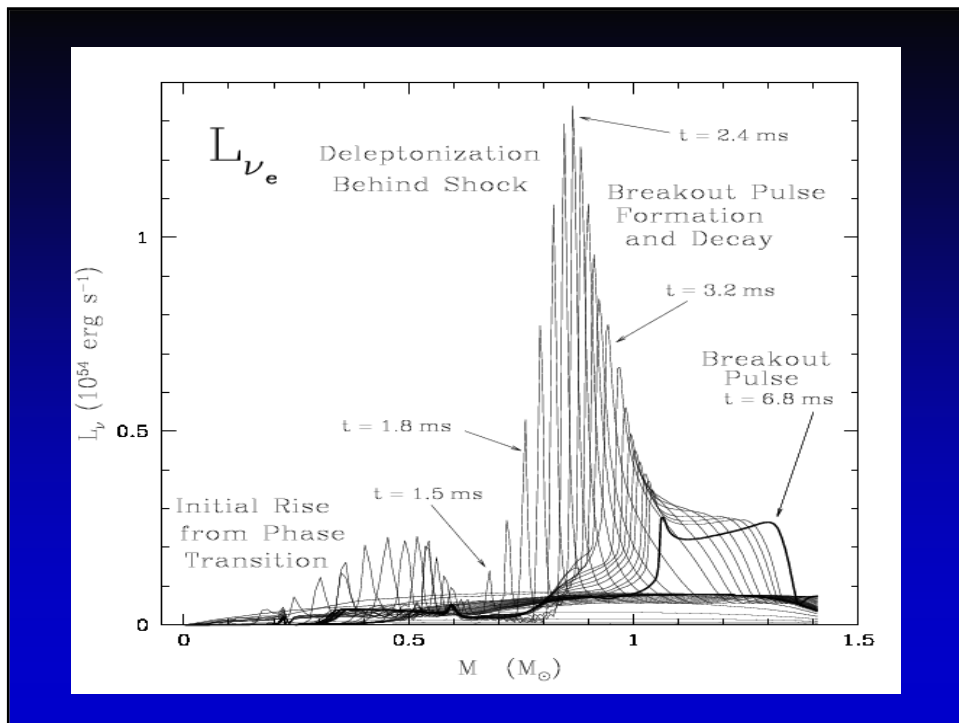
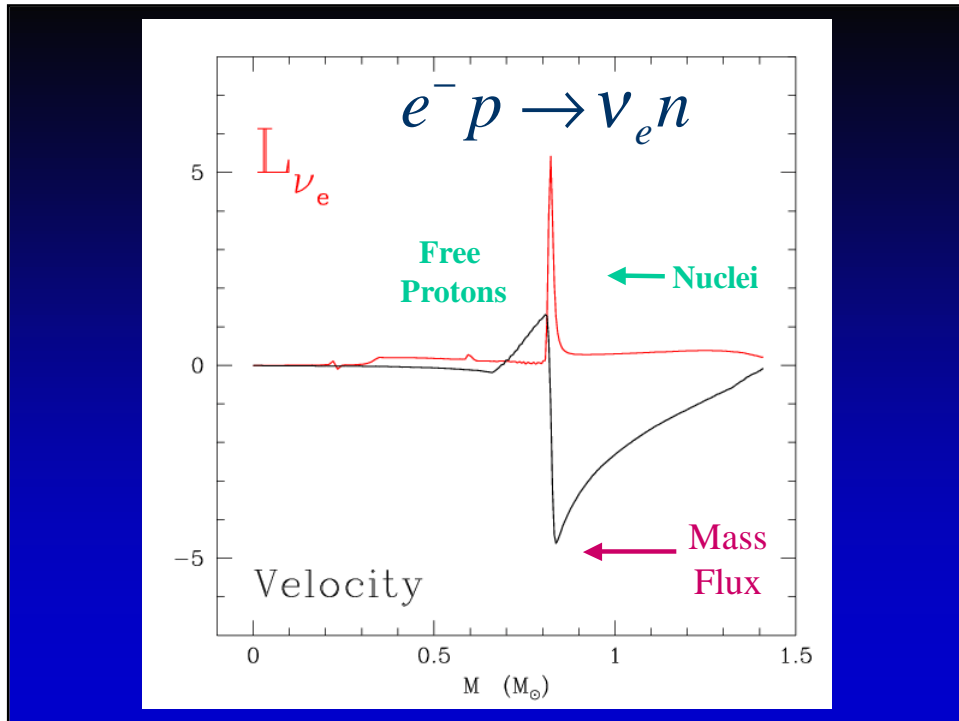
300 km

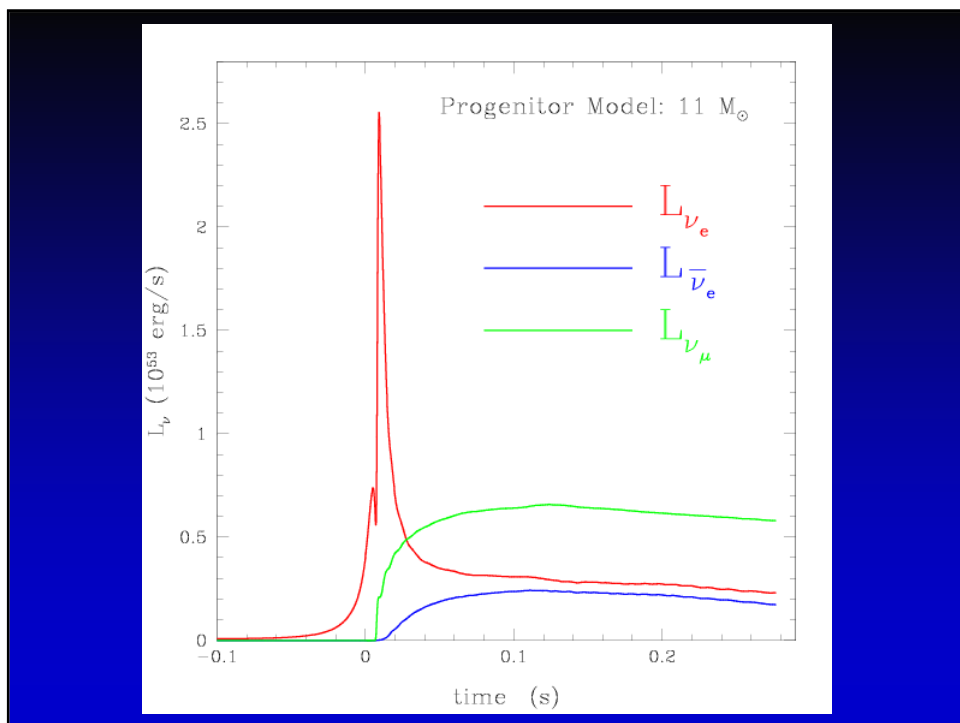
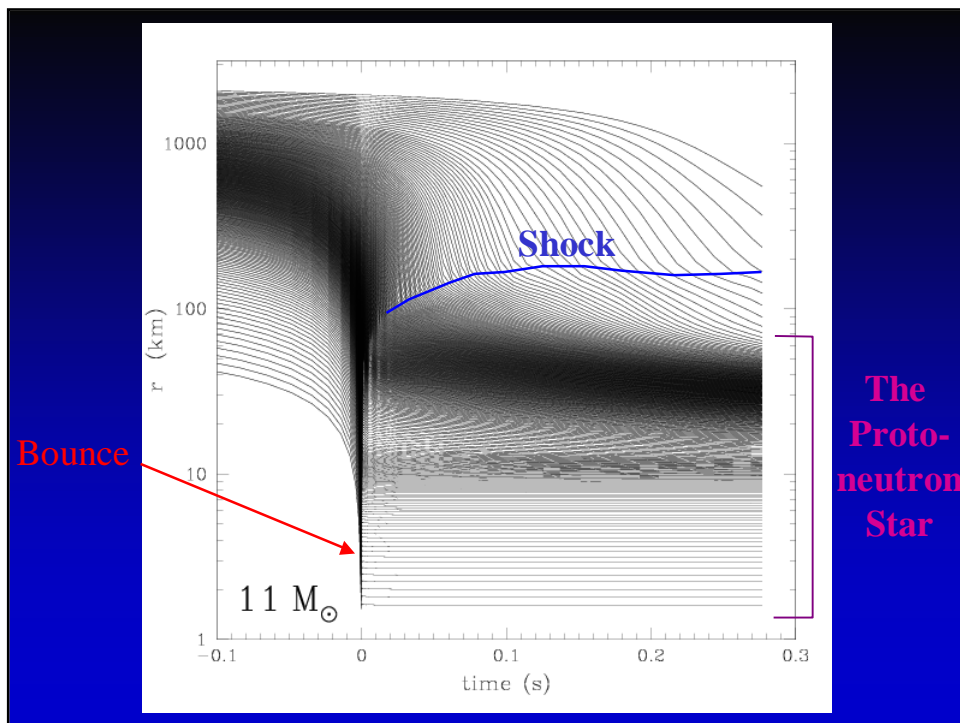


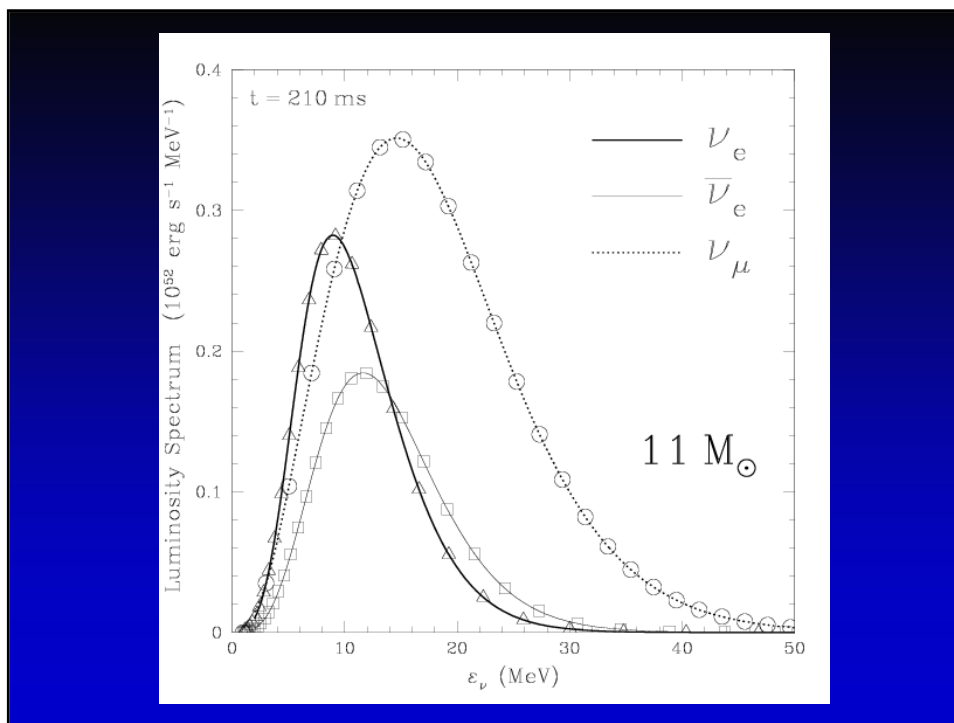
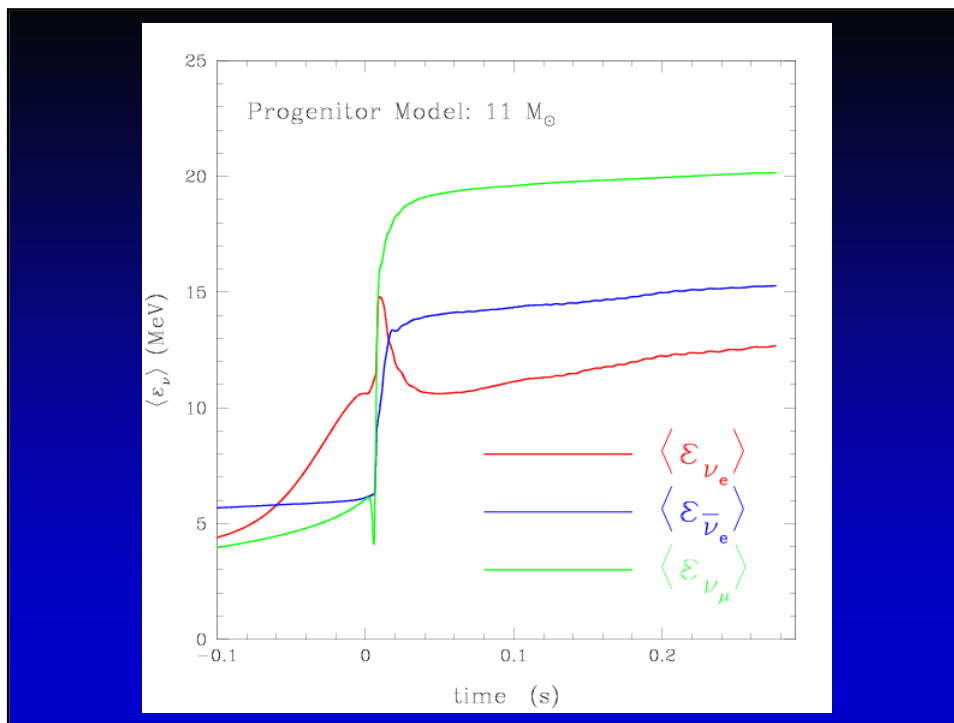
Dynamical Models of Supernovae

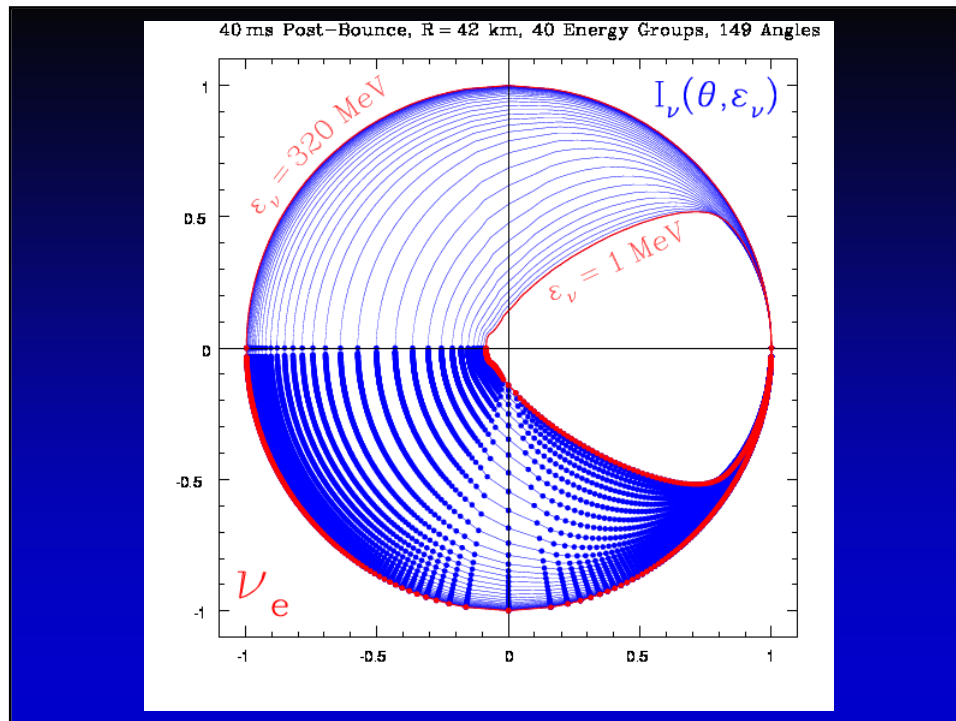
- **Ingredients:** Thompson et al. '02
 - **Hydrodynamics:** Newtonian, explicit, artificial viscosity
 - **High-density EOS:** Lattimer-Swesty, liquid drop model (Lattimer & Swesty 1991)
 - **Neutrino Transport:** Feautrier technique, tangent-ray, ALI (Eastman & Pinto 1993; Burrows et al. 2000)
 - **Microphysics:** Neutrino opacities, emission/absorption, new inelastic scattering algorithm, bremsstrahlung (Rampp & Janka '00, '02; Liebendorfer '00; et al. '01)
- **Tests & Verification**

QuickTime™ and a YUV420 codec decompressor are needed to see this picture.







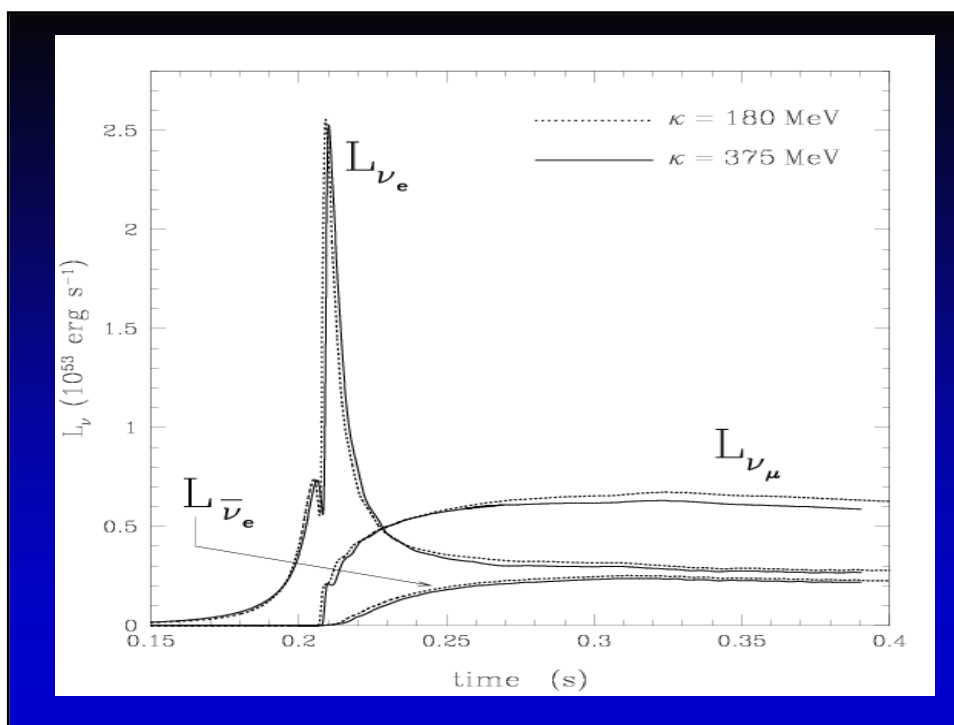
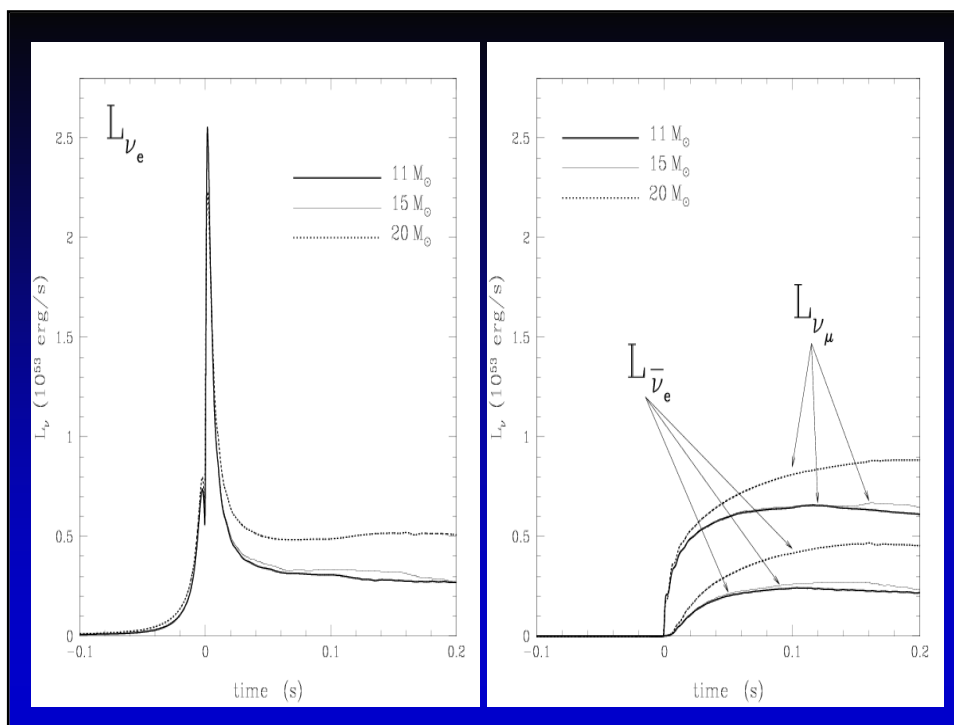


Variations & Detection

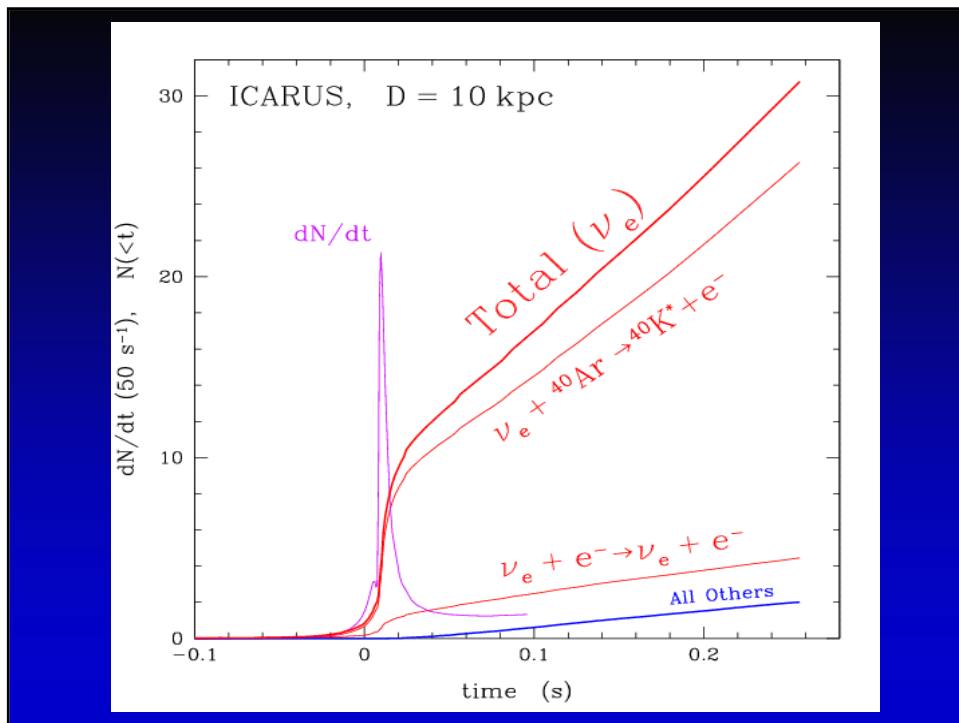
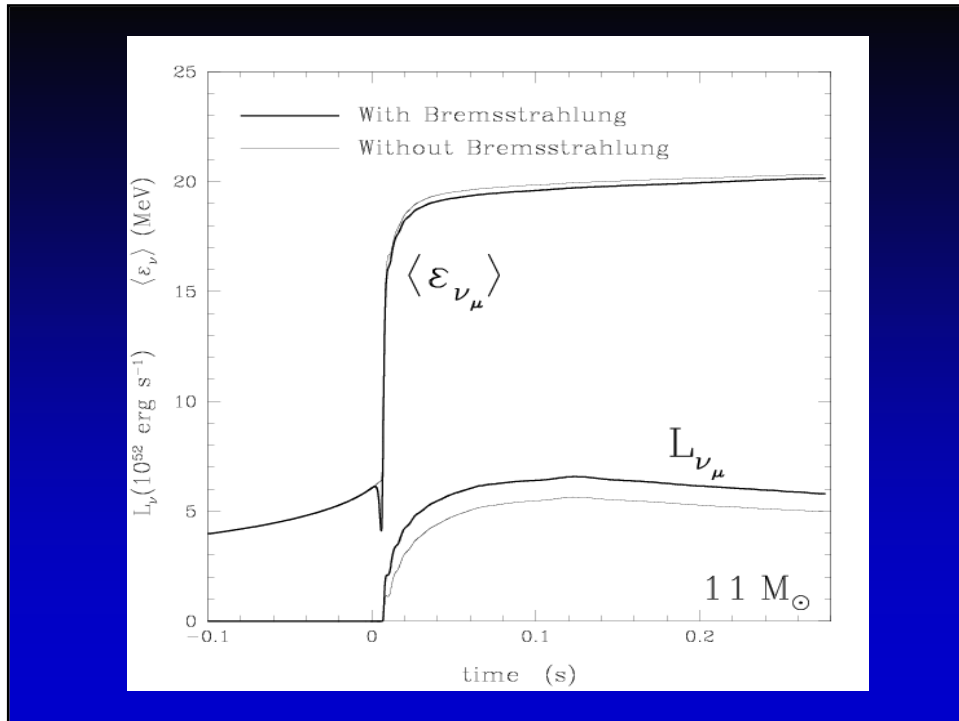
- Progenitor Mass: 11, 15, 20, 30, 40, 75 M_{sun}
- Microphysics:
 - Opacities: Neutral-current scattering/10 (Liebendorfer '00)
 - Nucleon-Nucleon Bremsstrahlung (Thompson et al. '00)
 - Inelastic Neutrino Scattering (Bruenn 1985; Thompson et al. '02)
 - Weak Magnetism/Recoil Correction (Horowitz '97, '02)
- Nuclear Compressibility: 180, 220, 375 MeV Modulus
- SK, SNO, ICARUS

Thompson et al. 2002

Supernovae, Neutrino-Driven Winds, and Nucleosynthesis



Supernovae, Neutrino-Driven Winds, and Nucleosynthesis



Supernova Model Summary

- Dynamical models constructed
- Precision neutrino spectra calculated for several progenitors
(AVAILABLE ONLINE! See me!)
- EOS & Opacity variations assessed
- Detector signals calculated

No Explosions in 1D!

(Burrows, Hayes, & Fryxell 1995)

QuickTime™ and a YUV420 codec decompressor are needed to see this picture.

— 300 km —→

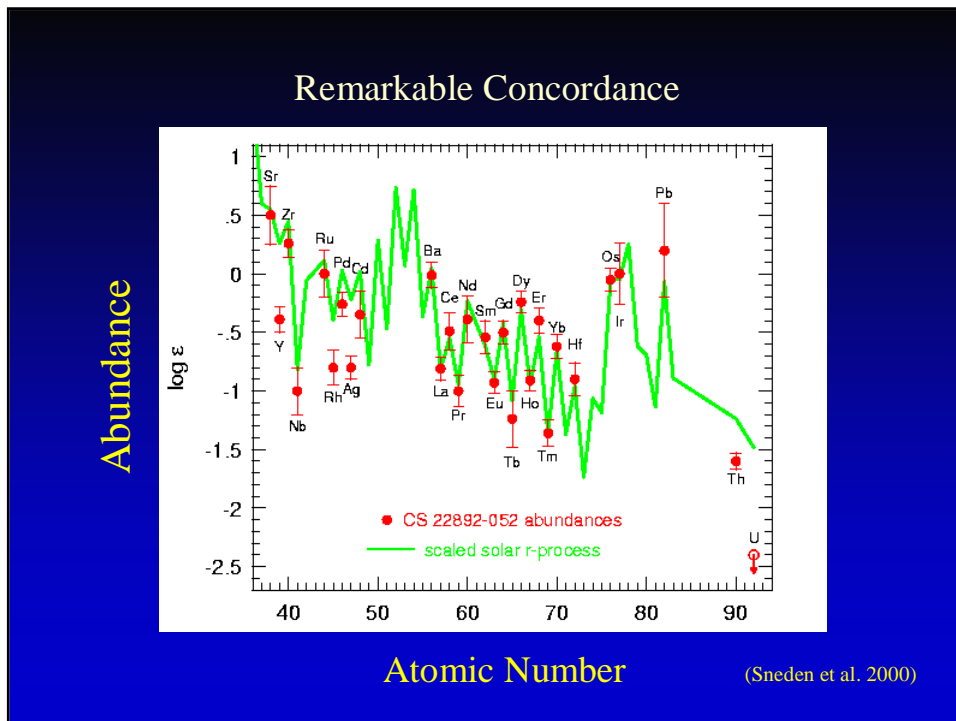
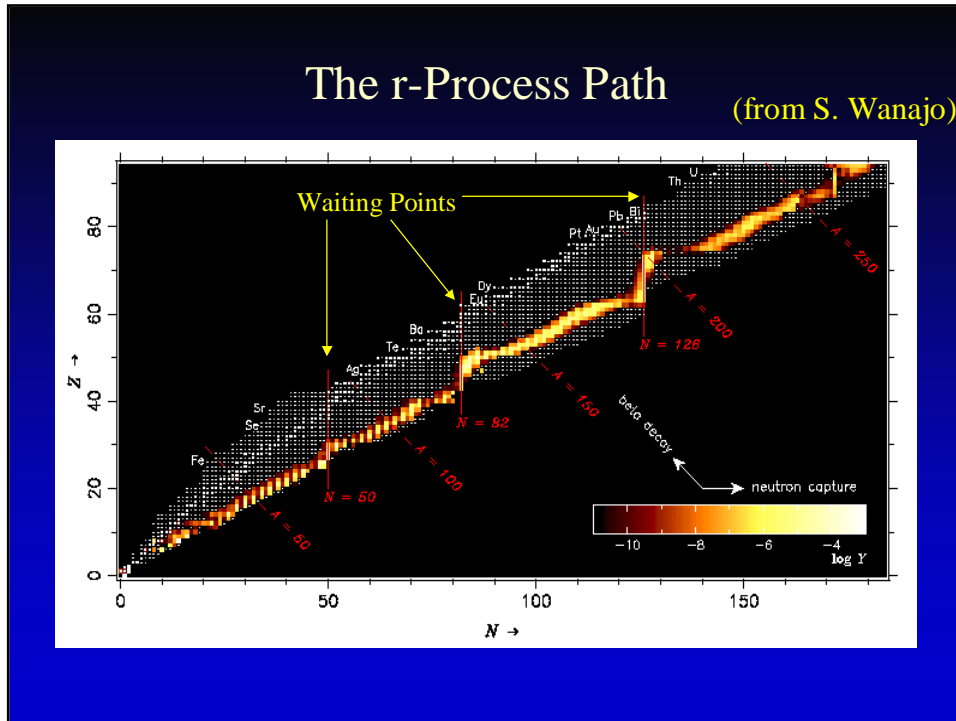
What is the origin of the heavy nuclei we find in nature?

s-Process and r-Process

The Rapid Neutron Capture (r) Process

- Seed nuclei **neutron-capture** on timescales much **shorter** than those for **beta decay**.
- If neutron-to-seed ratio $\sim 100:1 \Rightarrow$ heaviest nuclei.
 - Ratio set by entropy, dynamical timescale, & electron fraction
- The r-process is responsible for roughly **half** of all nuclei above the iron group.

(Burbidge, Burbidge, Fowler, and Hoyle 1957; Cameron 1957)



Motivation:

Suggests a universal r-process site that acts early in the chemical enrichment history of the galaxy.

(Burris et al. 2000; Cayrel et al. 2001)

What is the astrophysical site?

Possible Sites:

- Protoneutron Star Winds
 - Rate ~ 1 every 50 yr
 - Required mass ejected $\sim 10^{-5}$ - $10^{-6} M_{\text{sun}}$
 - Neutron Star-Neutron Star Mergers
 - Rate ~ 1 every 10^5 yr,
 - Required mass ejected $\sim 10^{-1} M_{\text{sun}}$
- Both are neutron rich

What are the physical characteristics of neutrino-driven winds?

Are these conditions suitable for a robust r-Process?

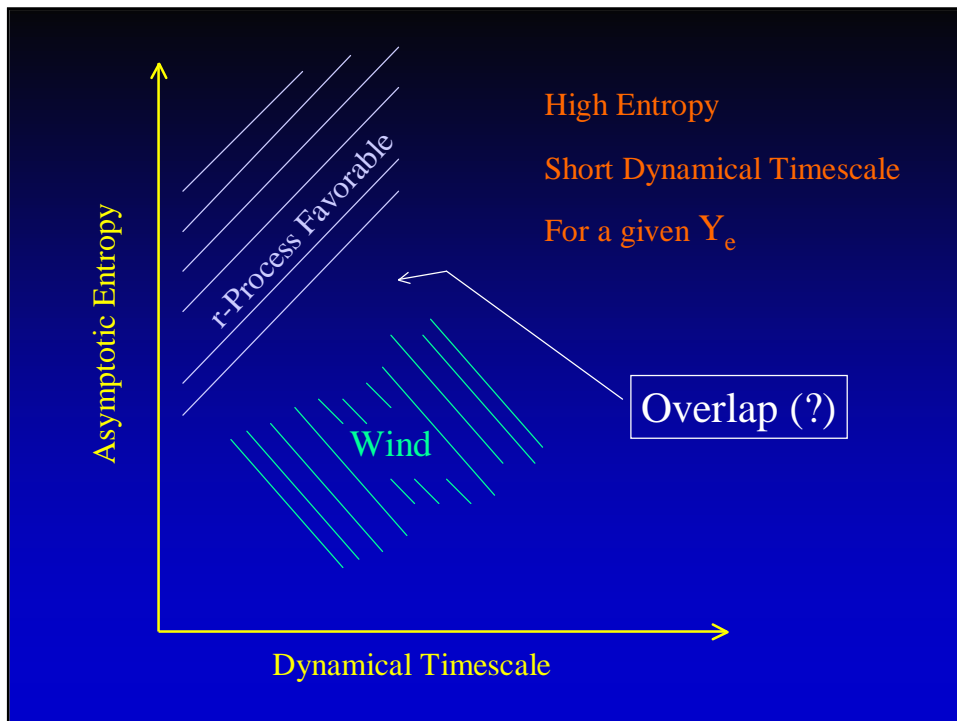
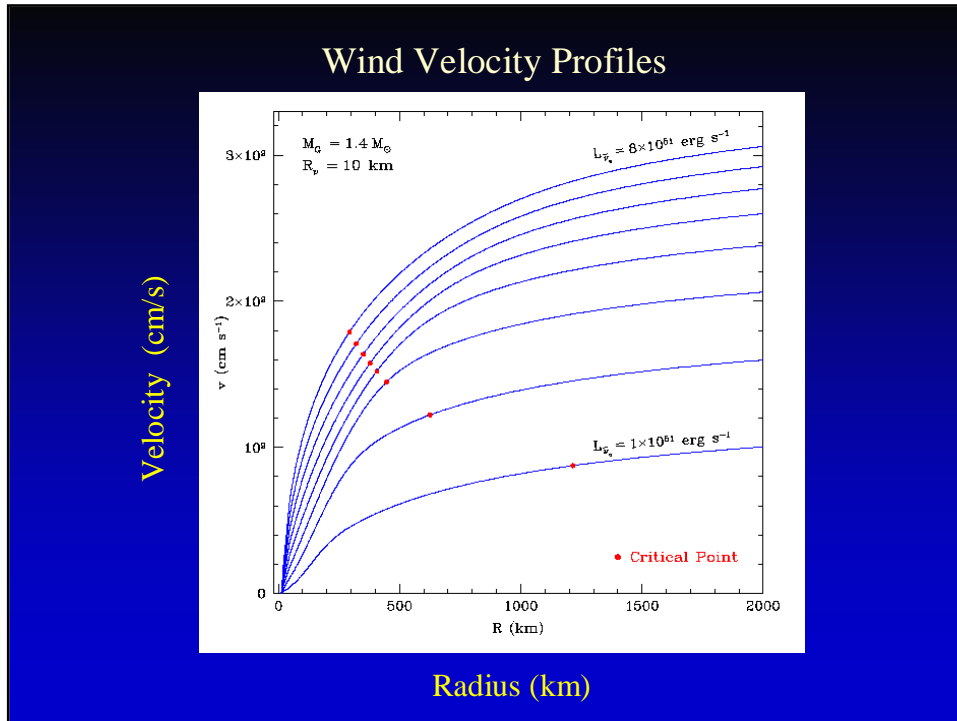
Entropy, dynamical timescale, and electron fraction

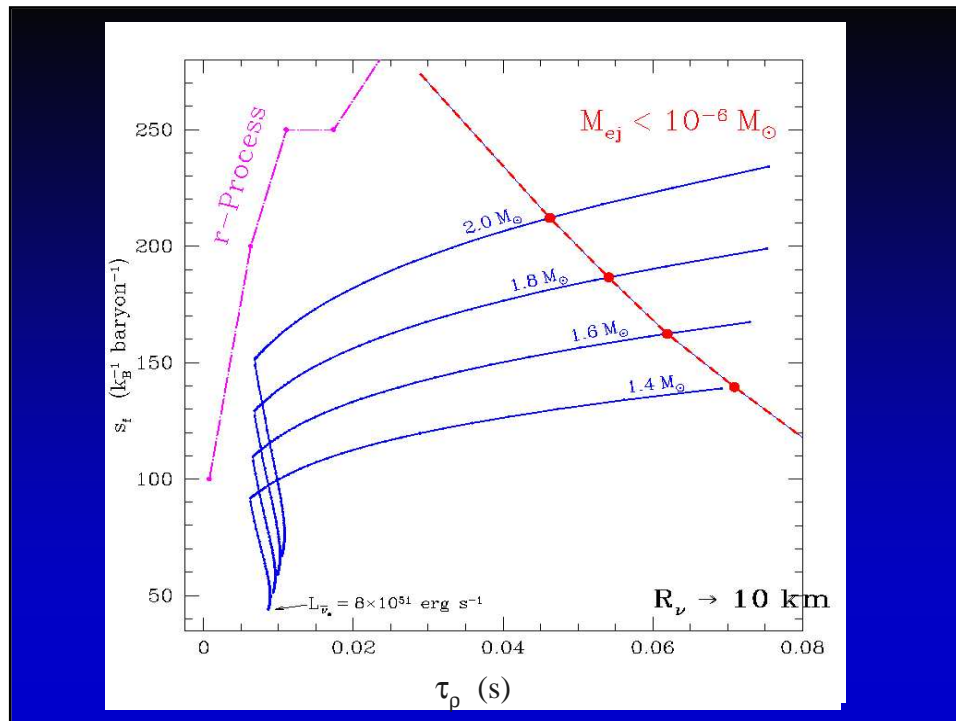
Neutrino-Driven Steady-State Winds

- Physics: eigenvalue problem ($\dot{M} = 4\pi r^2 \rho u$)
 - Three critical, coupled ODEs for ρ , T , and v
 - Neutrino heating and cooling
 - Electron fraction (Y_e) evolution
 - General Relativity, Schwarzschild metric
 - EOS: Arbitrary lepton degeneracy and relativity

Thompson, Burrows, & Meyer 2001

(Paczynski & Prószyński 1986; Duncan et al. 1986)

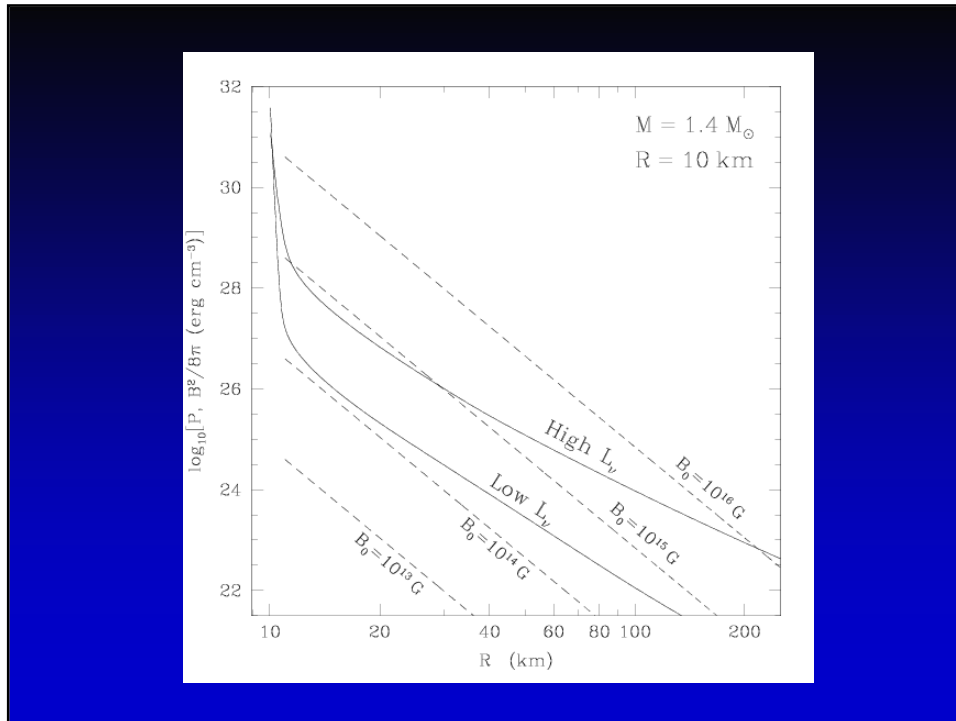




Transonic neutrino-driven winds
from $1.4 M_{\text{sun}}$ $R=10\text{km}$ neutron
stars fail to produce 3rd-peak r-
process nucleosynthesis

Qian & Woosley '94; Otsuki et al. '00;
Thompson et al. '01; Wanajo et al. '01;
Sumiyoshi et al. '00; but, see Terasawa et al. '02

What about magnetic fields
and rotation?



Closed Loops, Trapping, and Entropy Amplification

- Closed loops trap matter, but not permanently

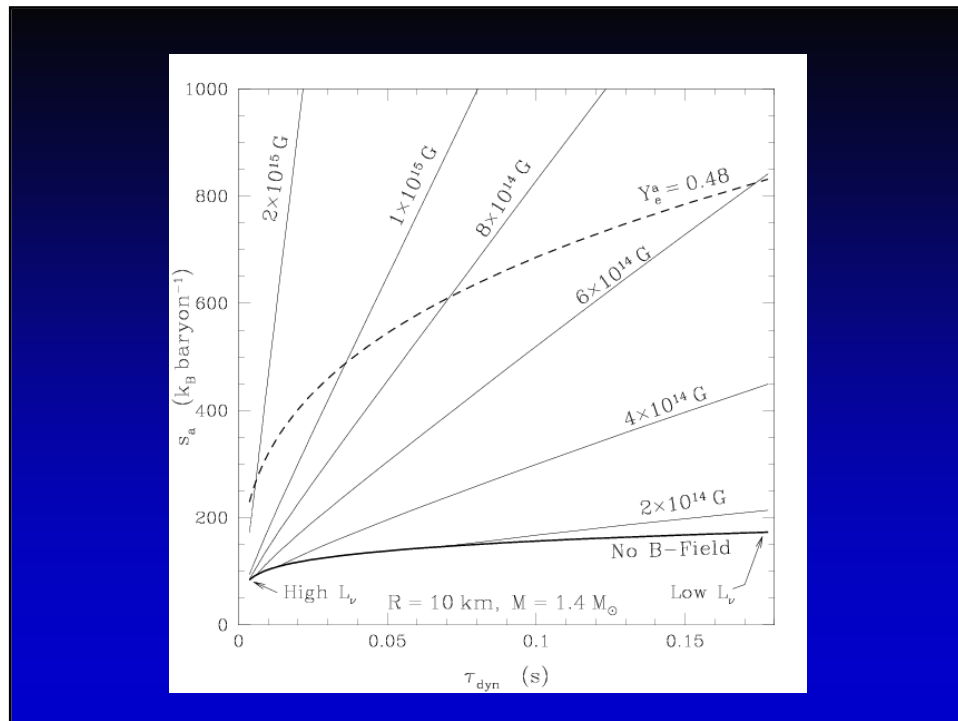
- Trapping Timescale:

$$\tau_{\text{Trap}} \approx [B^2/8\pi - P]/[q\rho]$$

- Entropy Enhancement:

$$\Delta s \approx \tau_{\text{Trap}} q / T$$

Thompson 2003



Magnetic Protoneutron Star Winds

- Neutrino-driven winds from PNSs with magnetar-like field strengths are magnetically dominated.
- In regions where the field dominates P , closed loops will form.
- These loops may trap matter temporarily, leading to higher asymptotic entropy.
- This entropy enhancement may be sufficient for robust r-process nucleosynthesis.

Some Open Questions

- Complex field topologies
- MHD instabilities
- Convection & Rotation - footpoint motion
- Spindown

The Future

- Dynamical models of explosion, wind emergence, and evolution (reverse shocks, fallback, etc)
- 2D & 3D models of neutrino-driven magnetohydrodynamic protoneutron star winds.

The End

The Fundamental Equations

$$\frac{dv}{dr} = \frac{v}{2r} \left[\frac{v_e^2}{y^2} \left(\frac{1 - c_s^2/c^2}{c_s^2 - v^2} \right) - 4c_s^2 \left(\frac{1 - v^2/c^2}{c_s^2 - v^2} \right) \right] + \frac{D}{C_V T} \frac{\dot{q}}{y} \left(\frac{1 - v^2/c^2}{c_s^2 - v^2} \right)$$

$$\frac{d\rho}{dr} = \frac{2\rho}{r} \left(\frac{v^2 - v_e^2/4y^2}{c_s^2 - v^2} \right) + \frac{\rho}{vy} \frac{D}{C_V T} \frac{\dot{q}}{c_s^2 - v^2}$$

$$\frac{dT}{dr} = \frac{2}{r\rho} \frac{D}{C_V} \frac{P + \varepsilon}{c^2} \left(\frac{v^2 - v_e^2/4y^2}{c_s^2 - v^2} \right) + \frac{\dot{q}}{C_V (vy)} \left(\frac{(1 - D/c^2)c_T^2 - v^2}{c_s^2 - v^2} \right)$$

where $D = c^2 \frac{T}{\varepsilon + P} \frac{\partial P}{\partial T} \Big|_{\rho}$ and $y = \gamma (1 - 2GM/rc^2)^{1/2}$

Example: Neutrino-Electron Scattering

$$\dot{q}_{\nu,e} = cn_e n_{\nu_i} \langle \sigma_{\nu_i,e} \omega_{\nu_i,e} \rangle = \int c \frac{dn_e}{d\varepsilon_e} d\varepsilon_e \frac{dn_{\nu_i}}{d\varepsilon_{\nu_i}} d\varepsilon_{\nu_i} \sigma_{\nu_i,e} \omega_{\nu_i,e}$$

$$\omega_{\nu_i,e} \approx \frac{1}{2} (\varepsilon_{\nu_i} - 4T)$$

(Bahcall 1964)

$$\sigma_{\nu_i,e} \approx \kappa_i T \varepsilon_{\nu_i}$$

(Tubbs and Schramm 1975)

$$\dot{q}_{\nu,e} \approx \frac{c}{\rho} \left[\frac{T^3}{(\hbar c)^3} \frac{F_2(\eta_e)}{\pi^2} \right] \frac{L_\nu}{4\pi R^2 c \langle \varepsilon_\nu \rangle} \Phi^4 2\Xi(r) \quad \text{erg g}^{-1} \text{s}^{-1}$$

$$\times \left[\frac{\kappa}{2} \langle \varepsilon_\nu \rangle \frac{F_4(\eta_\nu)}{F_3(\eta_\nu)} T \left(\langle \varepsilon_\nu \rangle \Phi \frac{F_2(\eta_\nu)}{F_3(\eta_\nu)} - 4T \frac{F_3(\eta_\nu)}{F_4(\eta_\nu)} \right) \right]$$

(Also Tubbs 1975; Janka 1999; Salmonson and Wilson 1999)

Neutrino Interactions: Heating and Cooling (\dot{q})

- **Charged Current** : $\nu_e n \leftrightarrow e^- p$ and $\bar{\nu}_e p \leftrightarrow e^+ n$
also affect Y_e evolution
- **Scattering** : $\nu_i e^{\mp,+} \leftrightarrow \nu_i e^{\mp,+}$ and $\nu_i n, p \leftrightarrow \nu_i n, p$
- **Pair Processes** : $e^- e^+ \leftrightarrow \nu_i \bar{\nu}_i$

Tools:

- Algorithm:
 - 2-point boundary value problem
 - Relaxation on adaptive radial mesh
(London and Flannery 1982)

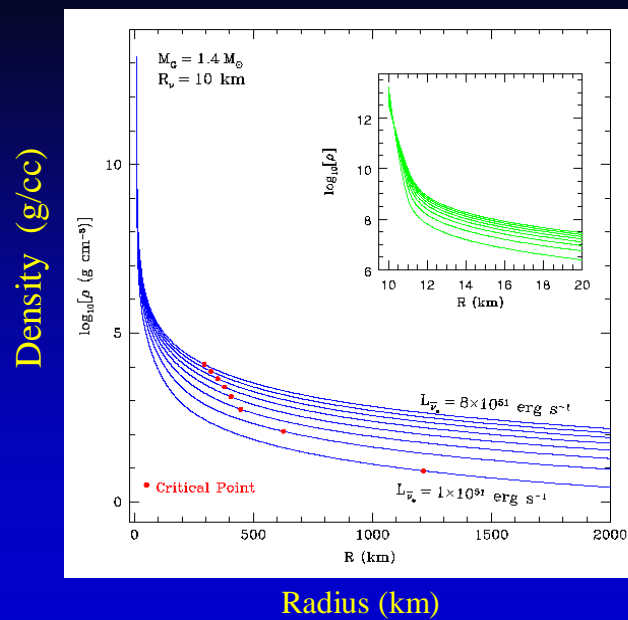
- Physical BCs:

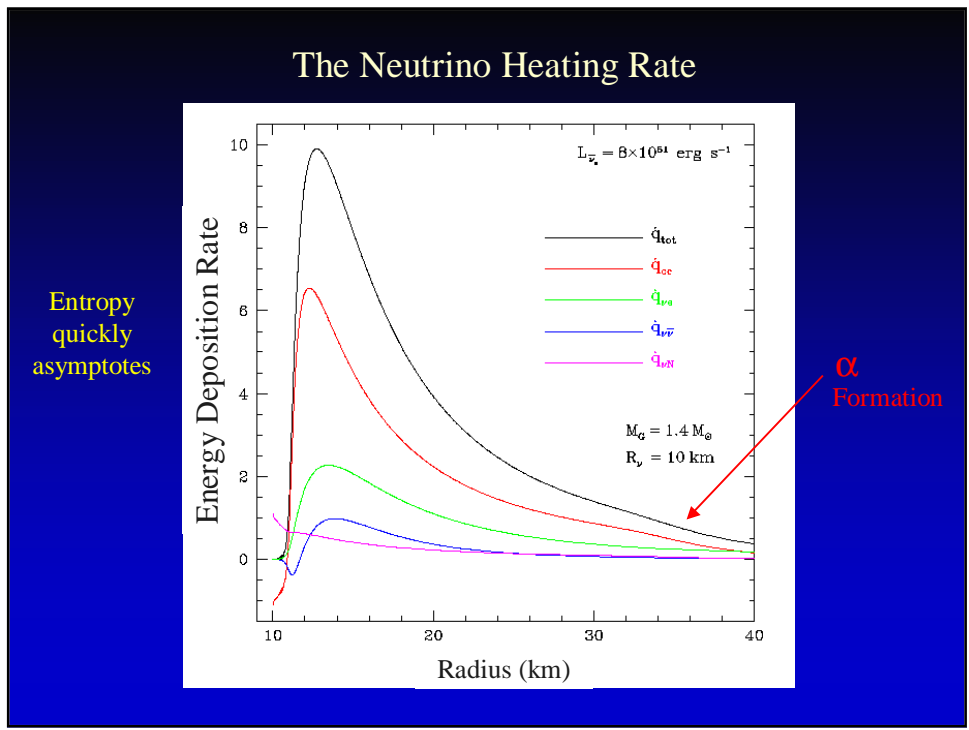
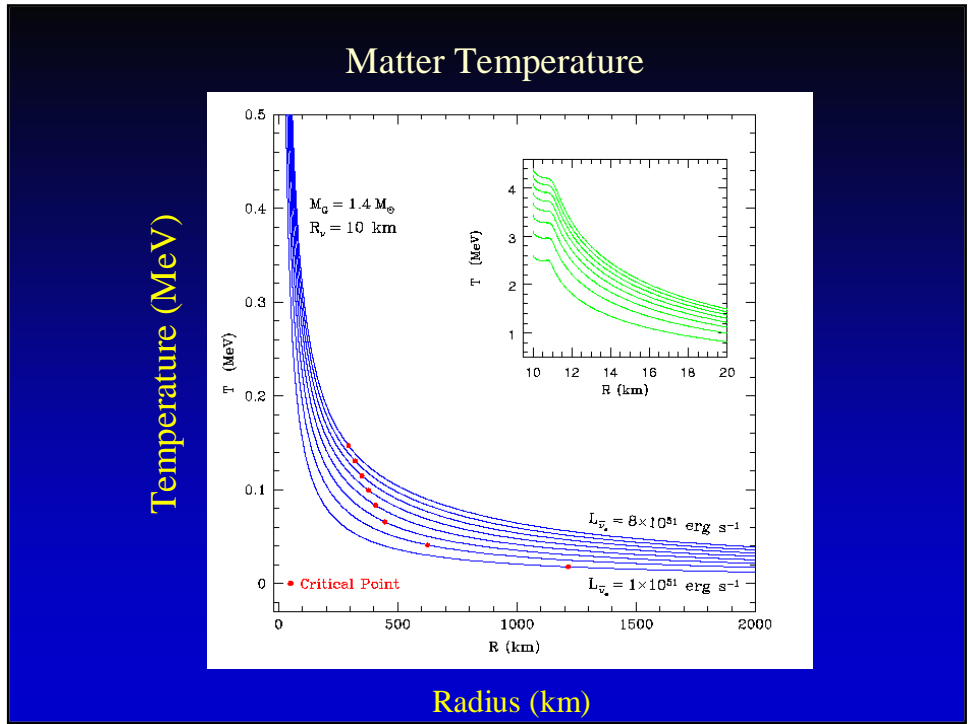
- Optical depth
- Triple Newton-Raphson
- Sonic Point ($v = c_s$)

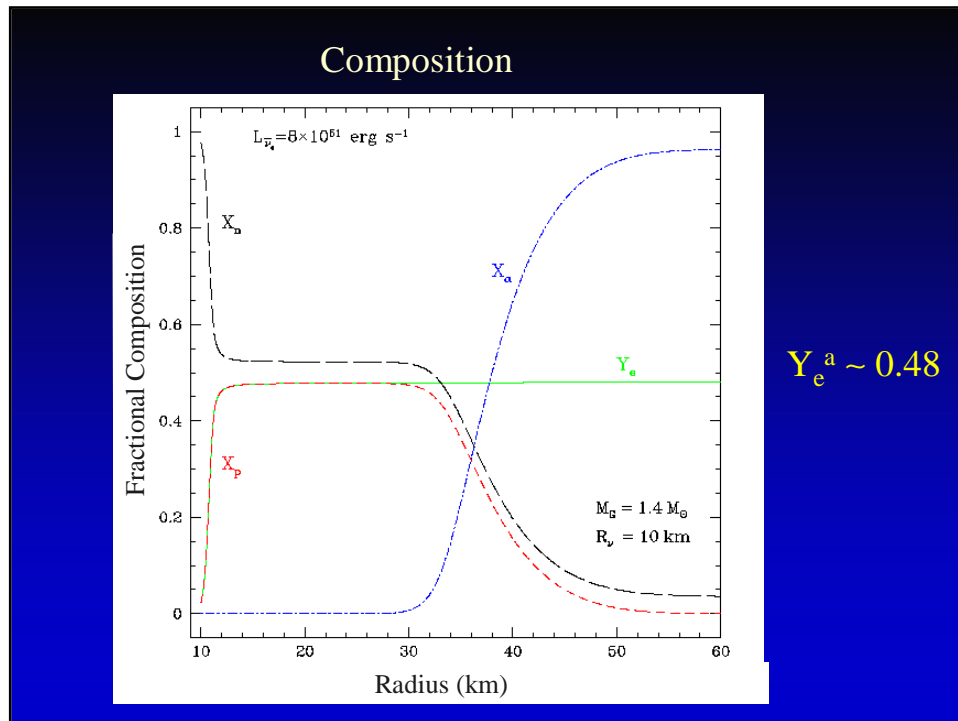
$$\tau_{\nu_e}(R) = \frac{2}{3} = \int_R^{\infty} \kappa \rho dr$$

$$\dot{q}(R)=0; \quad \left. \frac{dY_e}{dr} \right|_R = 0$$

Mass Density Profiles







Dynamical Timescale

- Increased by including GR
 - Redshift
 - Bending of null geodesics
- Not constant, evolves in radius

$$\tau_\rho \propto L_\nu^{-0.75} \langle \epsilon_\nu \rangle^{-2.6}$$

Thompson, Burrows, & Meyer 2001

Asymptotic Y_e

is set by ratio of luminosities and magnitude of energies.

$$Y_e^a \approx \frac{\Gamma_{\nu_e n}}{\Gamma_{\nu_e n} + \Gamma_{\bar{\nu}_e n}} \approx \left(1 + \frac{L_{\bar{\nu}_e} \langle \epsilon_{\bar{\nu}_e} \rangle - 2\Delta + 1.2\Delta^2 / \langle \epsilon_{\bar{\nu}_e} \rangle}{L_{\nu_e} \langle \epsilon_{\nu_e} \rangle + 2\Delta + 1.2\Delta^2 / \langle \epsilon_{\nu_e} \rangle} \right)^{-1}$$

$$Y_e^a > 0.45$$

(Qian et al. 1993)

Asymptotic Entropy

- Increased 20-30 units by GR: more compact
- Important Feedbacks
 - Cannot arbitrarily increase net heating
 - Broader deposition profile increases s .

$$s_a \propto L_\nu^{-0.15} \langle \epsilon_\nu \rangle^{-0.4}$$

Modeling Evolution with Steady-State Winds

- Posit $L_\nu(t)$, $R(t)$, and $\epsilon_\nu(t)$
- Calculate mass ejected

$$M_{\text{ejected}}(t) = \int_0^t \dot{M}(t') dt'$$

- Timescales:
 - Sound crossing time
 - Escape time
 - Luminosity decay time

$$\tau_{\text{sound}}, \tau_{\text{escape}} \ll \tau_{\text{decay}}$$

Assessing the Wind's Potential for r-Process Nucleosynthesis

What is the neutron-to-seed ratio in the wind?

- Entropy
 - asymptotes in ~100 km
- Y_e
 - asymptotes in ~10 km
- Dynamical Timescale
 - sets slope of density profile

$$\tau_\rho = \frac{1}{(v\gamma)} \left| \frac{1}{\rho} \frac{\partial \rho}{\partial r} \right|^{-1}_{T=0.5\text{MeV}}$$