Neutrino Bounds on Astrophysical Sources and New Physics

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March 3, 2003 KITP Neutrino Conference

- Introduction: neutrino and hadron shower Characteristics
- Neutrino acceptances and exposures: hadronic and electromagnetic
- Bounds on the high energy neutrino flux
- Model-independent bounds on new physics interactions
- Implications for TeV-scale gravity
- Conclusions

Work done with: Luis Anchordoqui, Jonathan Feng and Alfred Shapere Phys. Rev. D **66**, 103002 (2002) [hep-ph/0207139].

Shower characteristics

Neutrino showers

- Long interaction length > any atmospheric depth
 ⇒ showers above detector
- Large EM component
- Curved front
- Signal spread in time (μs)

Hadron showers

- Short Interaction length ~ 40 g/cm² ⇒ shower maximum high in the atmosphere (X_{max} ~ 800 - 900 g/cm²)
- EM component damped in ~ 40 − 60 g/cm² ⇒ only muons survive at ground
- Flat shower front (> 100 km)
- Short signal (ns)

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$$(A\Omega)_{\text{eff}}(E_{\text{sh}},t) \equiv \int_{\theta_{\min}}^{\theta_{\max}} A(t) \mathcal{P}(E_{\text{sh}},\theta,t) 2\pi \sin\theta \, d\theta$$

Effective aperture P. Billoir, in Venice 1999, Neutrino telescopes, Vol. 2, p. 111

$$\int \theta^{max} \int (\partial \phi - \partial \phi) d\phi = 0$$

 $h_{
m max}$

= 15 km, $H \approx 8$ km

= total observation time of the detector

$$(4\Omega)_{\text{eff}}(E_{\text{sh}},t) \equiv \int_{0}^{\theta_{\text{max}}} A(t) \mathcal{P}(E_{\text{sh}},\theta,t) 2\pi \sin\theta \, d\theta$$

Exposure

 $\mathcal{E}(E_{\rm sh}) \approx \int_0^T dt \int_0^{h_{\rm max}} (A\Omega)_{\rm eff}(E_{\rm sh}, t) \frac{\rho_{\rm atm}(0)}{\rho} e^{-h/H} dh$

 $ho_{
m water}$

Neutrino exposure: ground arrays



Neutrino exposure: fluorescence detectors (FD)

Hadronic and EM similar

- FD's sensitive to total EM activity along shower axis
- Because π^{\pm} mostly interact before decay, 80-90% of energy in hadronic showers is EM
- Adopt total FE exposure for both hadronic and EM showers from a total of five running periods ("epochs") (1983-1992) as reported in

R. M. Baltrusaitis *et al.*, Nucl. Instrum. Meth. A **240** (1985) 410; R. M. Baltrusaitis *et al.*, Nucl. Instrum. Meth. A **264** (1988) 87. D. J. Bird *et al.* [HIRES Collaboration], Astrophys. J. **424** (1994) 491.

• The additional periods enhance the first-epoch FE exposure by a factor of 3.



Results of searches for deeply penetrating QH showers

AGASA Collaboration

- Searched for QH showers with $X_{\rm max} \ge 2500$ g/cm²
- X_{\max} determined by
 - fit to lateral distribution of charged particles at ground level
 - fit to curvature of shower front
- expected hadronic backg'd 1.72 events
- only one event with $X_{\rm max}$ clearly > 2500 g/cm²

Fly's Eye

- X_{\max} determined by 3-parameter fit to charged particle density
- 5000 events, 11 years, no neutrino candidates
- combined data imply upper bound of 3.5 neutrino-induced events at 95% CL

Working equations for bounds

Event rate

$$N = \sum_{i,X} \int dE_i \, N_A \, \frac{d\Phi_i}{dE_i} \, \sigma_{iN \to X}(E_i) \, \mathcal{E}_{iX}(E_i)$$

- $i = \nu_e, \bar{\nu}_e, \nu_\mu, \bar{\nu}_\mu, \nu_\tau, \bar{\nu}_\tau$
- $d\Phi_i/dE_i$ = source flux of neutrino species *i*.
- $\mathcal{E}_{iX}(E_i)$ = appropriate exposure measured in cm³ we sr \cdot time.

Exposures for SM processes

- 20% of energy \rightarrow hadronic recoil
- Examples for AGASA (charged current)

 $\mathcal{E}_{\nu_e X}(E_{\nu_e}) = \min\{\mathcal{E}_{had}(0.2E_{\nu_e}) + \mathcal{E}_{EM}(0.8E_{\nu_e}), \mathcal{E}_{sat}\}$ $\mathcal{E}_{\nu_{\mu}X}(E_{\nu_{\mu}}) = \mathcal{E}_{\text{had}}(0.2E_{\nu_{\mu}})$

Examples for Fly's Eye (charged current):

 $\mathcal{E}_{\nu_e X}(E_{\nu_e}) = \mathcal{E}(E_{\nu_e})$ $\mathcal{E}_{\nu_{\mu}X}(E_{\nu_{\mu}}) = \mathcal{E}(0.2E_{\nu_{\mu}})$

• Neutral current RHS = $\mathcal{E}_{had}(0.2E_{\nu_i})$ (AGASA), $\mathcal{E}(0.2E_{\nu_i})$ (FE)



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Bounds on astrophysical neutrino fluxes (cont'd)

Model-independent local bounds (continued)

• Take $\Delta=1,$ and total mixing

$$\langle E_i d\Phi_i / dE_i \rangle = \frac{1}{6} \langle E_\nu d\Phi_\nu / dE_\nu \rangle$$

and use SM cross sections to obtain model-independent local upper bound on total ν flux at 95% CL.

Global bounds assuming particular flux behavior

• Illustrative choices

$$\frac{d\Phi_{\nu}}{dE_{\nu}} = J_0 \left(\frac{E_{\nu}}{E_0}\right)^{-\gamma} , \gamma = 1.5 \text{ or } 2.0$$

• Can now integrate and obtain bounds over entire energy range for each γ .



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Haim Goldberg Northeastern University, Boston Radio Ice Čerenkov Experiment:

searches for radio pulses from EM showers created by electron neutrino collisions in ice.

Goldstone Lunar Ultra-high energy neutrino Experiment:

searches for microwave Čerenkov pulses from EM showers induced by neutrinos in the Moon's rim.

Bounds on neutrino cross sections

Assume Protheroe-Johnson flux

- Input nucleon spectrum E^{-2} , cutoff energy $10^{12.5}$ GeV, source evolution $(1+z)^4$ as a minimum
- \bullet Expected event rate at AGASA and FE with SM cross section \sim 0.02/yr
- Keep new physics flavor-blind, negligible EM component
- Illustrate with two cases for new physics

(1) $y = E_{\rm sh}/E_{\nu} = 1$: all energy into shower. Example: TeV scale black hole production. D'Eath,

Payne, Phys. Rev. D **46**, 694 (1992).

(2) y = 0.1: leading particle effect. Example: KK graviton exchange in NC interaction Kachelriess,

Plumacher, hep-ph/0109184.

 Model-independent condition on cross section is now

 $N_A \left\langle \sigma_{\nu N \to X}(E_{\nu}) \right\rangle \left\langle \mathcal{E}(y E_{\nu}) \right\rangle \left\langle E_{\nu} d\Phi_{\nu} / dE_{\nu} \right\rangle < 3.5$

averaged over an energy interval of 1 e-folding.



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νN cross section bounds (cont'd)

Comparison to previous bounds C. Tyler, A. V. Olinto and G. Sigl, Phys. Rev. D **63**, 055001 (2001) [hep-ph/0002257] (TOS)

- Updated exposure \rightarrow gain factor of 9
- Cosmogenic flux source cutoff energy of $10^{12.5}$ vs. $10^{11.5}$ in TOS (just at cutoff CR energy) \rightarrow gain factor of 4
- 95% CL limits \rightarrow lose factor of 4
- → net improvement on bounds by about order of magnitude

Implications for TeV-scale gravity

Extra dimensions

- General idea: our 4-dimensional universe is embedded in a larger geometry with *n* extra spatial dimensions
- Compactify on n-torus, common radius R
- Regaining Newton's law at distances large compared to *R* implies

$$M_{\rm Pl}^2 = 8\pi \ M_D^{2+n} \ R^n$$

with M_D related to the (4 + n)-dimensional Planck mass

• Exciting possibility: R is large enough so that $M_D \sim 1 \text{ TeV}$ I. Antoniadis, Phys. Lett. B 246, 377 (1990); J. D. Lykken, Phys. Rev. D 54, 3693 (1996) [hep-th/9603133]; N. Arkani-Hamed, S. Dimopoulos and G. R. Dvali, Phys. Lett. B 429, 263

(1998) [hep-ph/9803315]

• Striking implication: existence of tower of massive gravitons $(K\bar{K})$

Tabletop Gravity

Newtonian gravity good down to 0.2 mm Hoyle et al., PRL86 (2001) 1418 \rightarrow for $n = 1 M_D > 5000$ TeV: uninteresting

Astrophysics

For n = 2(3) limits on rate of supernova cooling through KK graviton emission require

 $M_D > 10(600) {
m ~TeV}$ Cullen-Perelstein, PRL83 (1999) 268; Hannestad-Raffelt, PRL87

(2001) 051301; hep-ph/0110067

Accelerator experiments

• LEP – direct graviton emission (single photons, Z's): for $n = 4(6), \ M_D > 870(610)$ GeV L3 Collaboration,

PLB470 (1999) 281

• Tevatron – virtual KK graviton exchange in $e^+e^$ and $\gamma\gamma$ production: $M_D > 1.0 - 1.2$ TeV, depending on brane tension cutoff

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Proposal that collapse to TeV-scale BHs occurs in high energy particle collisions Banks, Fischler, hep-th/9906038; Emparan, Horowitz, Myers, Phys. Rev. Lett.

Black holes

85, 499 (2000) [hep-th/0003118]; Giddings, Katz, J. Math. Phys. 42, 3082 (2001) [hep-th/0009176]; Giddings, Thomas, hep-ph/0106219;

Dimopoulos,Landsberg, Phys. Rev. Lett. 87, 161602 (2001) [hep-ph/0106295].

- These evaporate primarily to SM particles
- Parton-parton cross section is \sim geometric $\hat{\sigma} \simeq \pi r_s^2$, where Schwarzschild radius

$$r_s(M_{\rm BH}) = \frac{1}{M_D} \left[\frac{M_{\rm BH}}{M_D} \right]^{\frac{1}{1+n}} \left[\frac{2^n \pi^{(n-3)/2} \Gamma(\frac{n+3}{2})}{n+2} \right]^{\frac{1}{1+n}}$$

Black hole production by neutrinos in cosmic rays



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Bounds on TeV-scale gravity

- Total exposure + cosmogenic flux + σ \longrightarrow expected event rates for different value of x_{\min} and M_D
- Requiring that event rate satisfy the 95% CL bound \longrightarrow determine lower bound on M_D for each x_{\min}



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Bounds on TeV-scale gravity (cont'd)

Can now generate 95% exclusion plot in $M_D - x_{\min}$ space, $n = 1 \dots 7$ from bottom.



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Comparison with D0 Collaboration bounds

• These are obtained from 95% CL upper bounds on anomalous 4-point functions generating $\gamma\gamma$ or

 e^+e^- pairs. B. Abbott *et al.* [D0 Collaboration], Phys. Rev. Lett. **86**, 1156 (2001) [arXiv:hep-ex/0008065].

• The 4-pt function arises through virtual KK graviton *s*-channel intermediate states. Giudice, Rattazzi,

Wells, Nucl. Phys. B 544, 3 (1999); Han, Lykken and Zhang, Phys. Rev. D 59, 105006 (1999); Hewett, Phys.

Rev. Lett. **82**, 4765 (1999).

• Sum diverges, so need cutoff introduced as brane-softening factor Λ - expect $\Lambda \leq M_D$.

M. Bando *et al*, Phys. Rev. Lett. **83**, 3601 (1999); M. Bando, T. Noguchi, arXiv:hep-ph/0011374; H. Murayama, J. D. Wells, arXiv:hep-ph/0109004.

Table 1: Lower limits on M_D at 95% CL in TeV.

Λ/M_D	$M_{D,\min}$ (TeV)		
	n = 4	n = 6	n=7
0.5	0.80	0.63	0.58
0.6	0.88	0.76	0.73
0.7	0.95	0.89	0.88
0.8	1.01	1.01	1.04
0.9	1.07	1.14	1.21
1.0	1.13	1.26	1.38

• For $\Lambda < M_D$, present bounds 1.3 - 1.4 TeV exceed Tevatron bounds ~ 1 TeV.

Summary of results and conclusions

- New bounds obtained on the cosmic neutrino flux from existing limits on quasi-horizontal deeply developing showers, taking into account the combined exposures of the AGASA and Fly's Eye experiments. Results significantly strengthen existing limits.
- They also present severe constraints on top-down models where the cascade decay of exotic elementary X particles or topological defects are responsible for the events detected with energies ≥ 10¹¹ GeV. This is because neutrinos are typically a significant component in X decays, and have a hard spectrum extending up to M_{GUT} ~ 10¹⁶ GeV, and our bounds are typically exceeded when the proton flux from top-down models is normalized to the observed spectrum.
- Complete neutrino exposure was combined with the flux of cosmogenic neutrinos, to derive model-independent upper bounds on the neutrino-nucleon cross section. These bounds strengthen existing limits by roughly one order of magnitude.

Summary of results and conclusions (cont'd)

• Considered TeV-scale gravity models to study BH production. Upper bounds on the neutrino-nucleon cross section implied lower limits on the fundamental Planck scale, which represent the best existing limits on TeV-scale gravity for $n \ge 5$ extra spatial dimensions.