

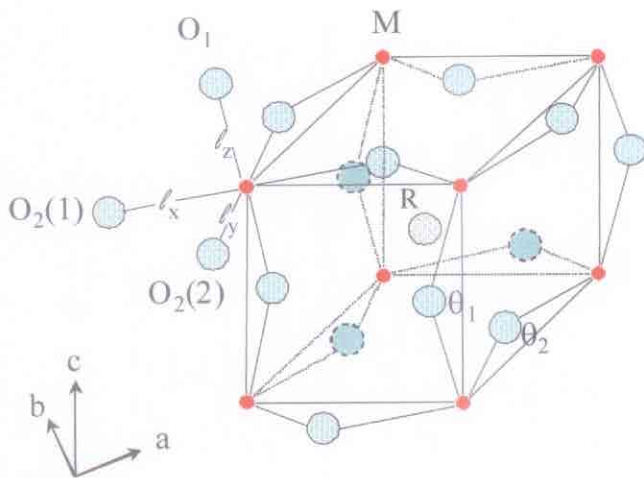
Magnetic properties of perovskites with d-orbital degeneracies

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1. Universal octahedral-site distortion in orthorhombic RMO_3 perovskites.
2. RTiO_3
3. RVO_3
4. RMnO_3
5. $\text{A}^{2+}\text{Ru}^{4+}\text{O}_3$

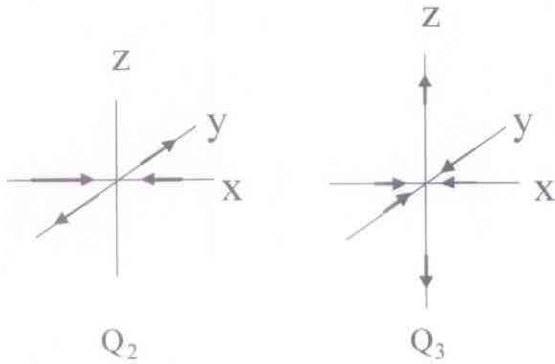
The crystal structure of orthorhombic RMO_3 perovskites



The basic perovskite unit cell with the orthorhombic distortion with Pbnm space group.

Coop. rotation about $[110]$, $[001]$
 $\longrightarrow a \leq c/\sqrt{2} < b$

Problem: Need $\text{MO}_{6/2}$ site distortion to accommodate fixed rotation axis.
O'Keeffe & Hyde (1977)



$$Q_2 = l_x - l_y, \quad Q_3 = (2l_z - l_x - l_y)/\sqrt{3},$$

$$\rho_o = (Q_2^2 + Q_3^2)^{1/2} \text{ and } \phi = \tan^{-1}(Q_3/Q_2)$$

IR = R^{3+} -ion radius (9-fold coordination)

Calculated vs Experimental ω and α

SPuDS program is used for calculating ω vs IR

Lufaso & Woodward, Acta Cryst. B 57, 725 (2001)

Assumes:

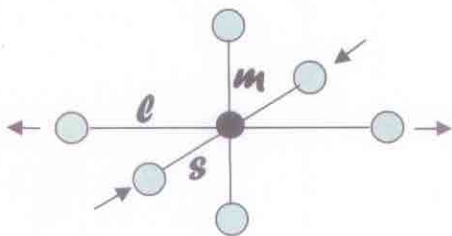
- (1) Rigid $\text{MO}_{6/2}$ rotations
- (2) (M-O) from bond valance sum rule.

Experiments:

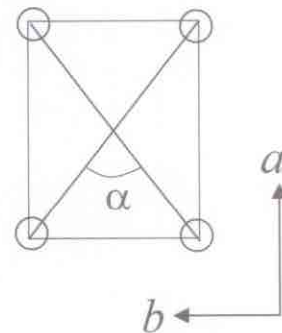
RFeO_3 (Pbnm for all R), *Marezio et al, Acta Cryst. B 26, 2008 (1970)*

Mat. Res. Bull. 6, 23 (1971).

Apparent octahedral-site distortions:



Max. distortion @ IR $\approx 1.11 \text{ \AA}$



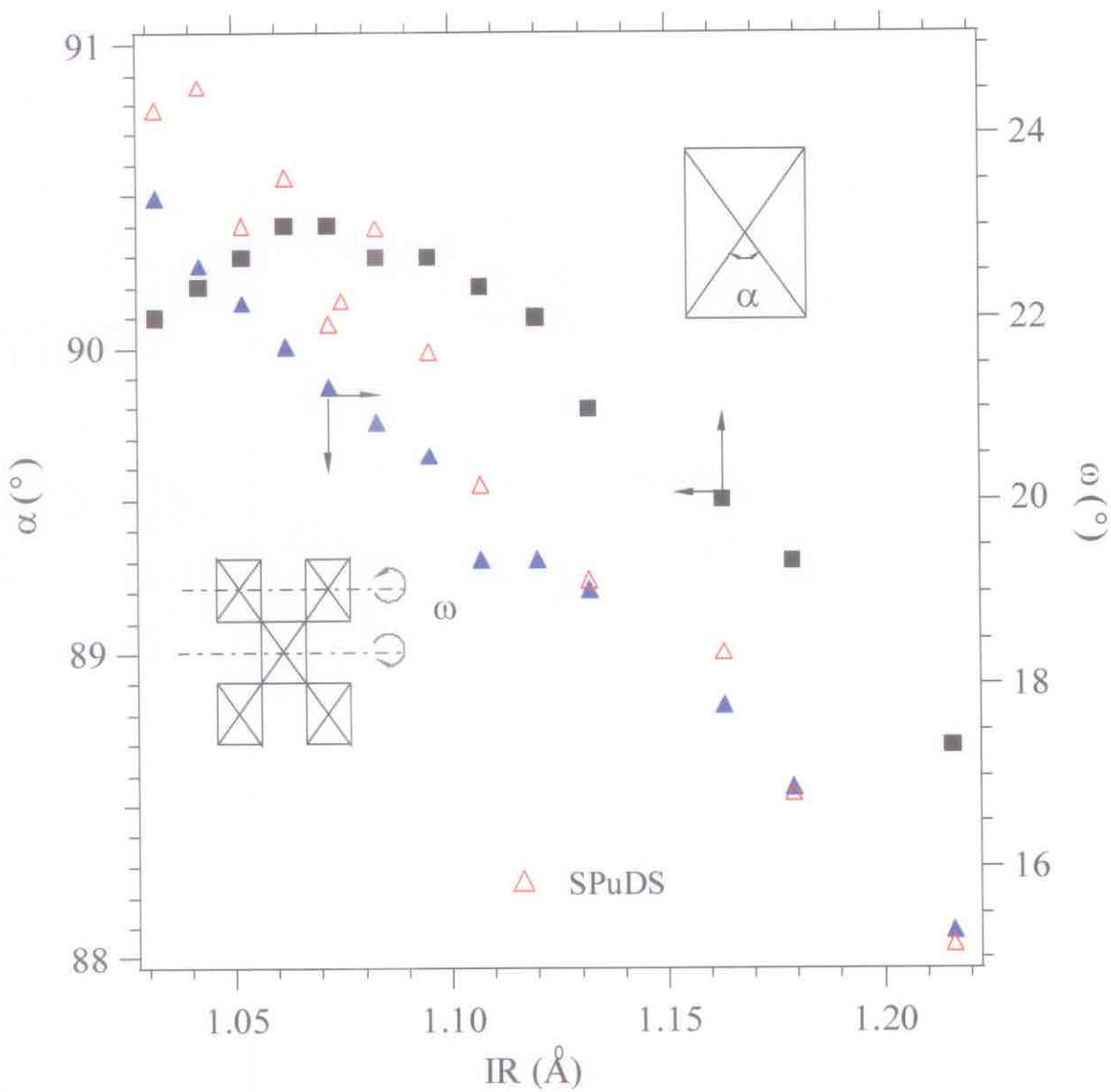
$\alpha < 90^\circ$ for IR $> 1.11 \text{ \AA}$

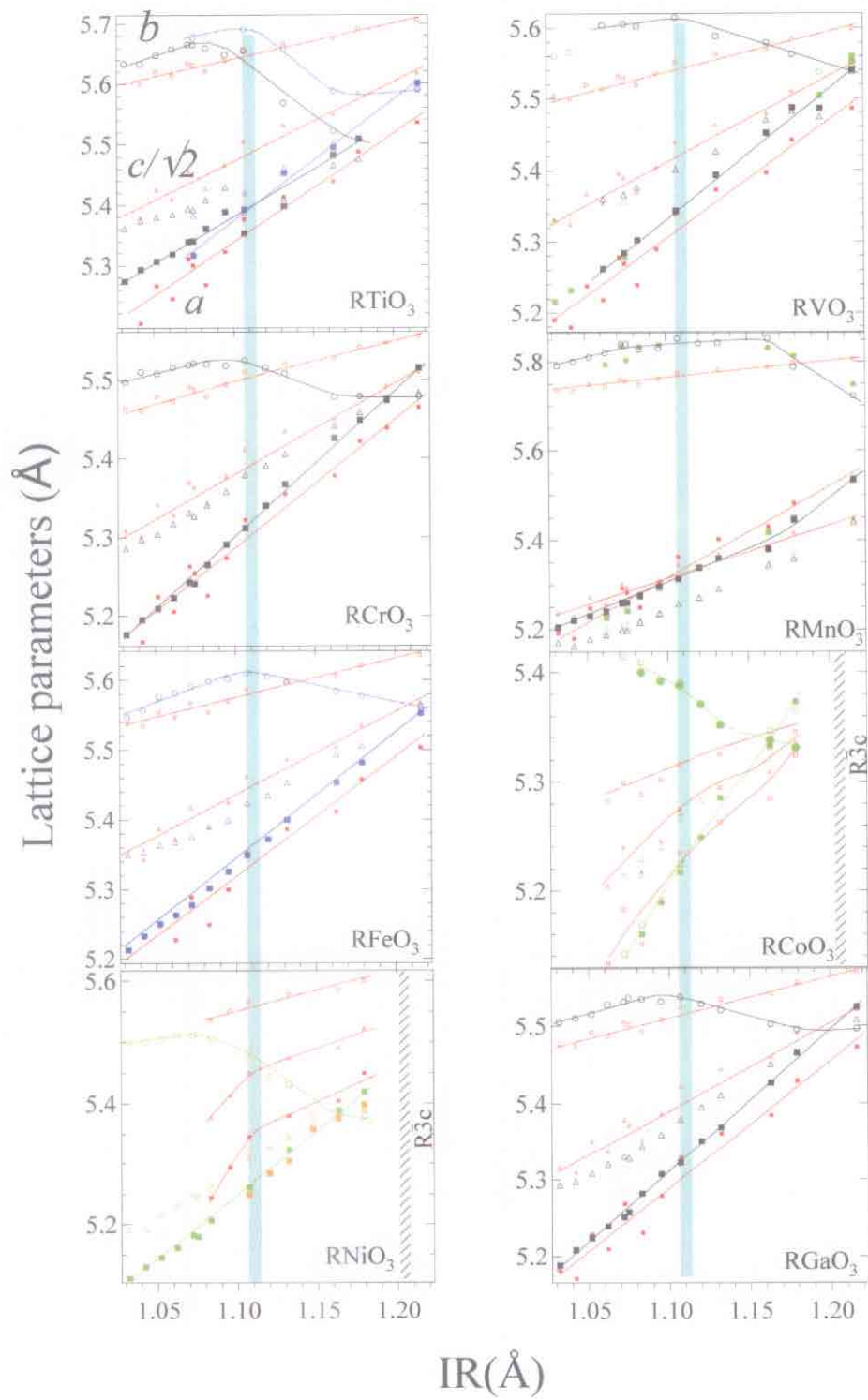
Note: With increasing IR,

Ortho (Pbnm) \longrightarrow Rhomb ($\text{R}\bar{3}\text{c}$) \longrightarrow Cubic (Pm3m)

The rhombohedral structure accommodates larger IR.

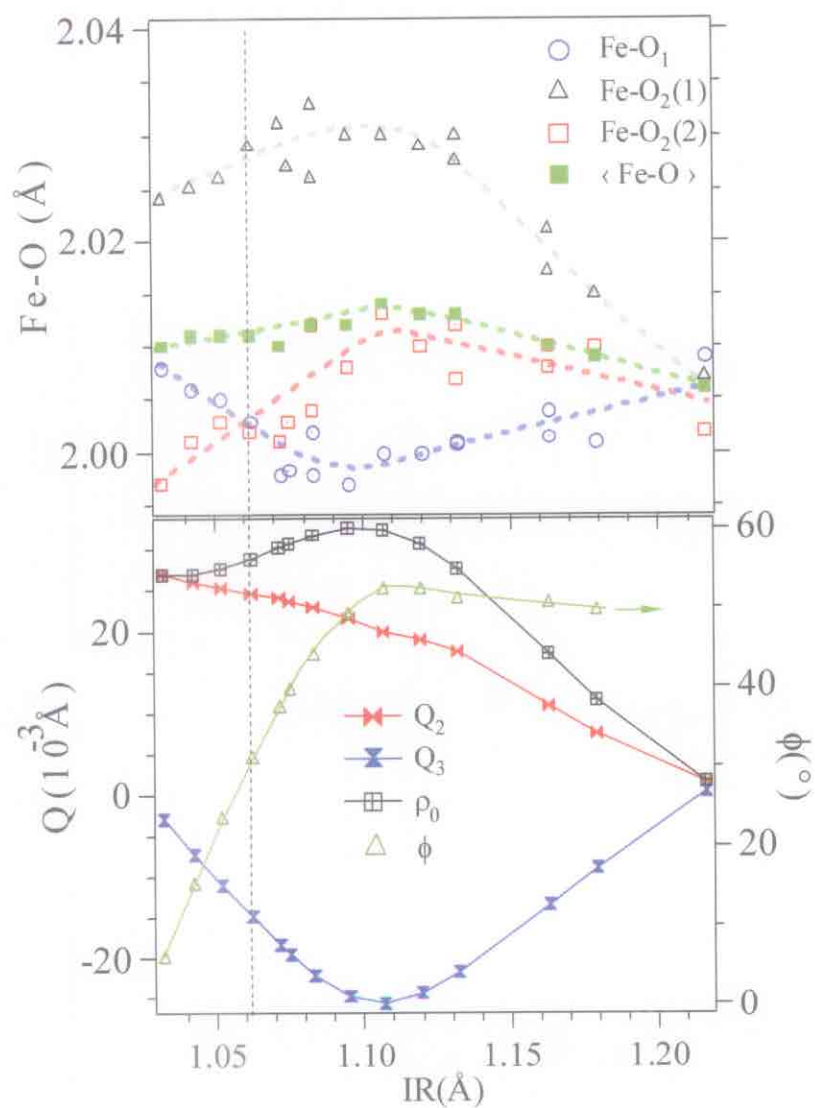
Angles ω and α vs. IR for $RFeO_3$ (Data from Marezio et al)



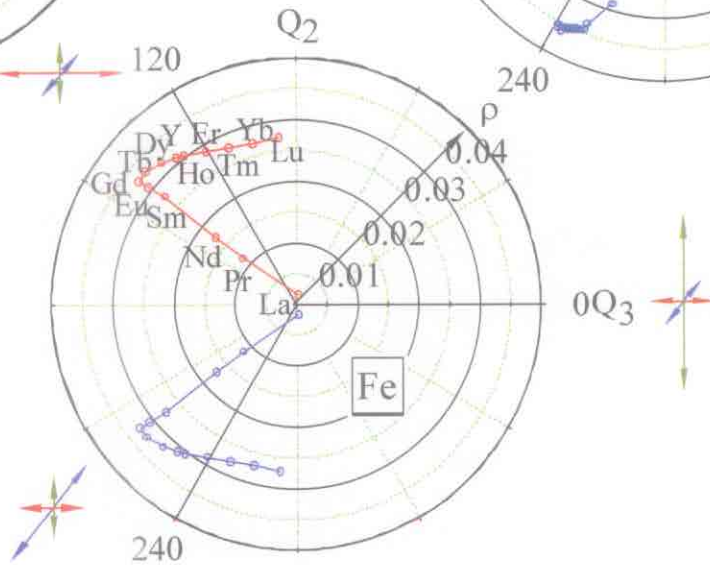
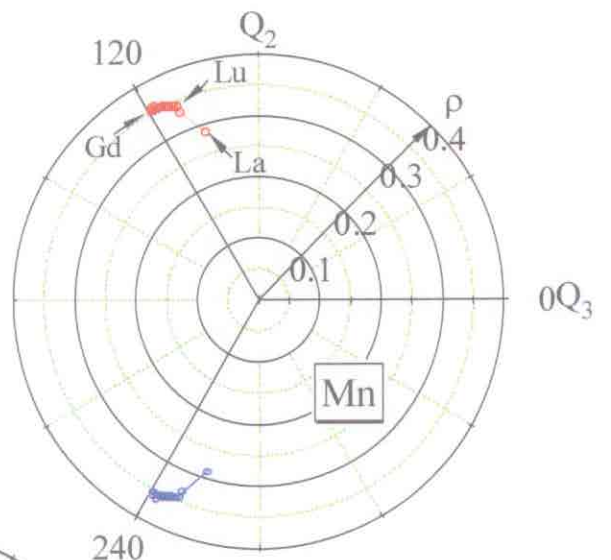
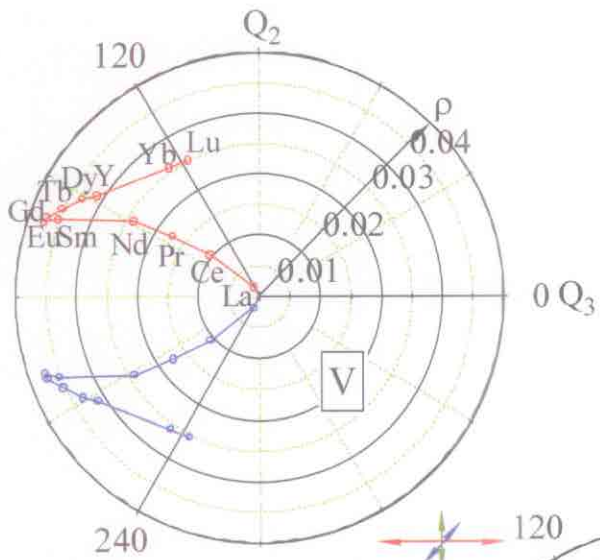


Lattice parameters with red symbols are calculated by SPuDS

Fe-O bond lengths in $R\text{FeO}_3$



Note: $\phi = \tan^{-1}(Q_3/Q_2) = 30^\circ$
 for the octahedral-site distortion $c/a > 1$



$$Q_2 = \mathbf{e}_x - \mathbf{e}_y, \quad Q_3 = (2\mathbf{e}_z - \mathbf{e}_x - \mathbf{e}_y) / \sqrt{3},$$

$$\rho_o = (Q_2^2 + Q_3^2)^{1/2} \text{ and } \phi = \tan^{-1}(Q_3/Q_2)$$

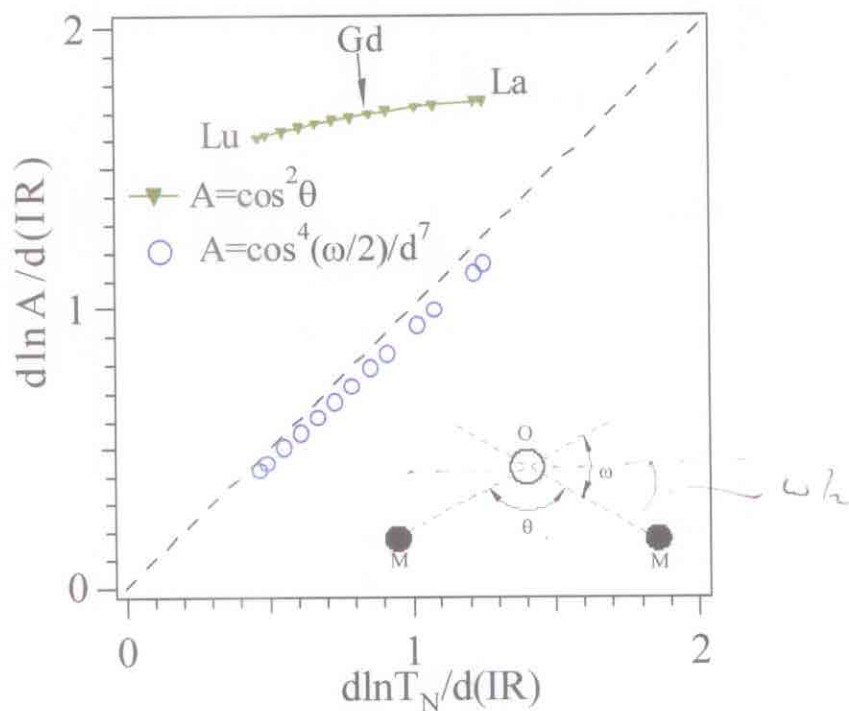
T_N observed vs the calculated value from superexchange interaction for $RFeO_3$

Superexchange: $T_N \sim b^2/U \sim A(\theta, d)$;

θ : Fe-O-Fe bond angle; d : Fe-O bond length.

For fixed U and $A \sim b^2$

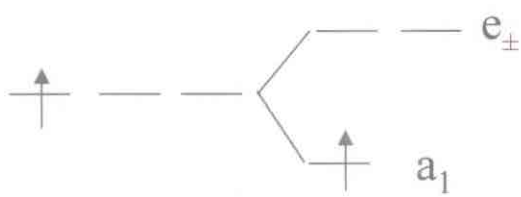
→ $d \ln A / d(\text{IR})$ vs $d \ln T_N / d(\text{IR})$ is straight line at 45° through origin.
 $A(\text{IR})$ from the structural data, $T_N(\text{IR})$ from Mössbauer data.



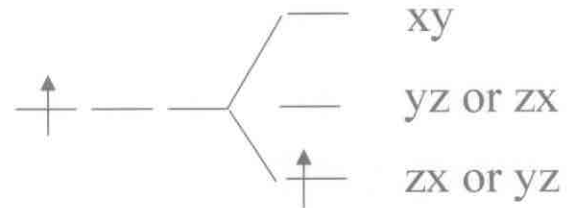
The RTiO₃ Perovskites

Ti³⁺:t¹σ*⁰

Possible site orbital order

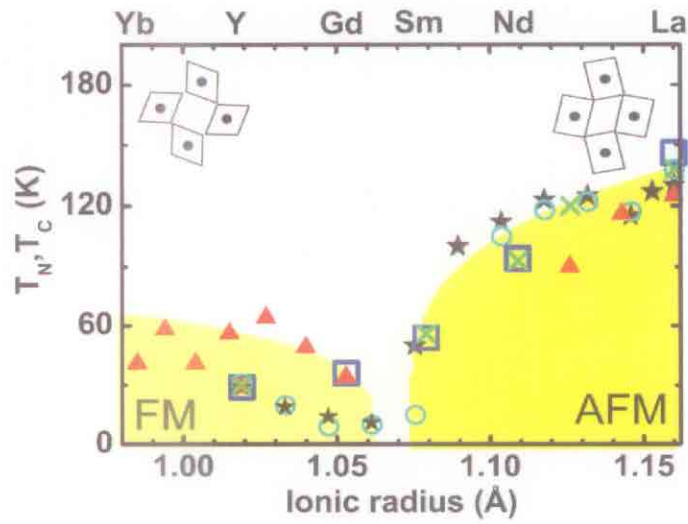
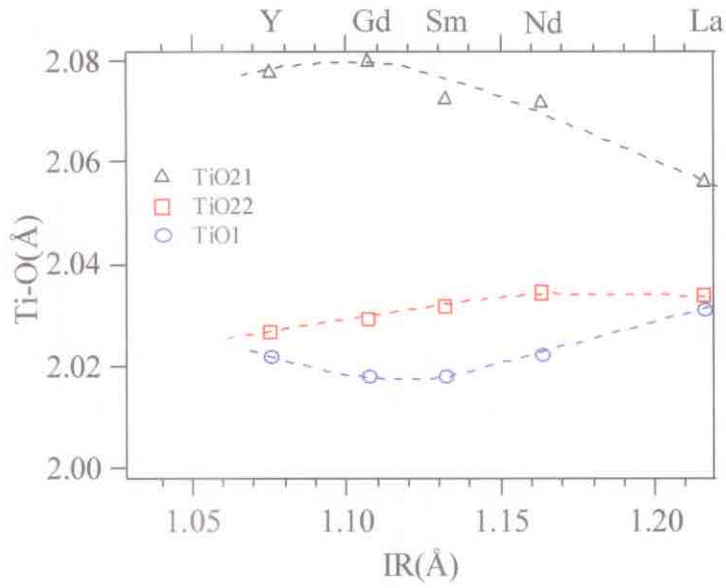


$\alpha < 90^\circ$
AF



$\alpha = 90^\circ$,
orthorhombic, F

RTiO₃

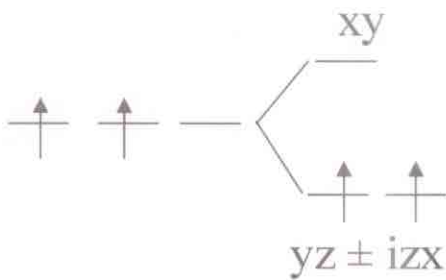


A.C. Komarek *et al* PRB 75, 224402(2007)

RVO₃ perovskites

V³⁺:t²e⁰: localized spins

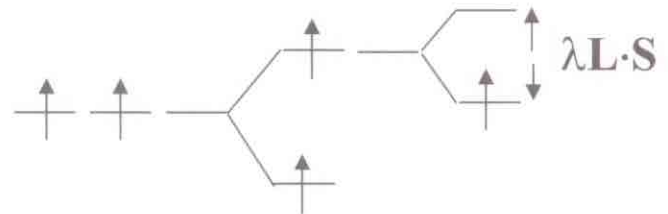
Possible site orbital order



Tetr $c/a > 1$

$T_{OO} > T_N$

Cooperative antiferroic



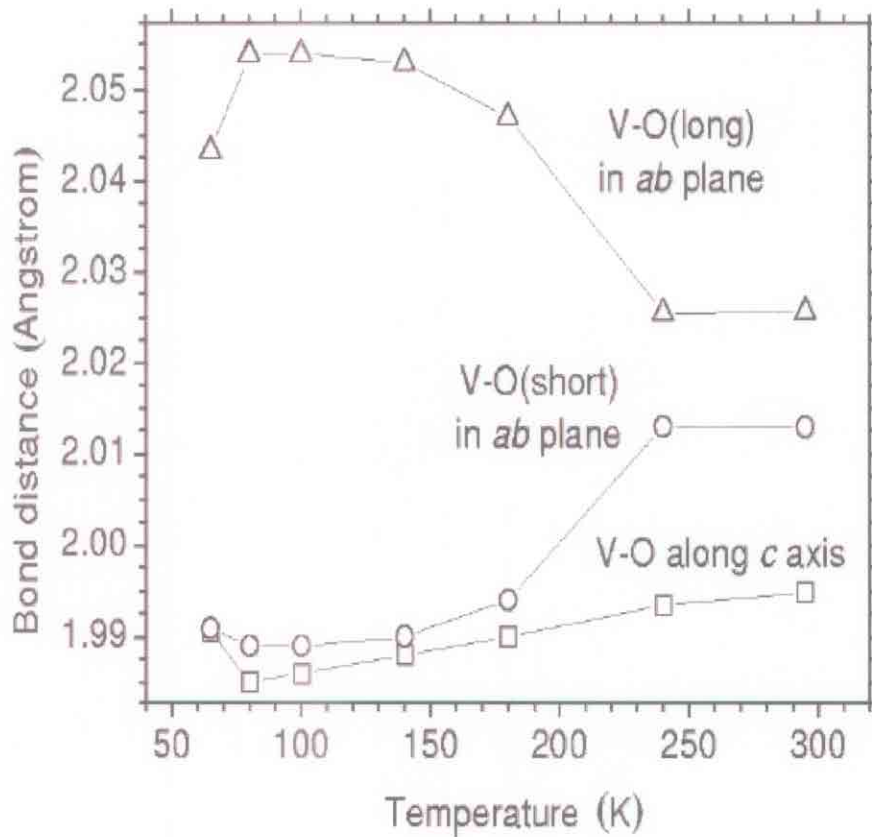
Pseudocubic ($c/a \approx 1$)

$T_{OO} \leq T_N$?

Cooperative ferroic

YVO₃ site distortion

Note: Distortion at $T > T_{OO}$ largest for IR $\approx 1.11 \text{ \AA}$



Temperature dependence of V-O bond lengths determined by using neutron powder diffraction.

Blake et al *Phys. Rev. B* 65, 174112 (2002)

Orbital & Spin ordering in YVO_3

Orbital ordering and spin-spin interactions:

$$G_{OO}: T_{CG} < T < T_{OO}$$

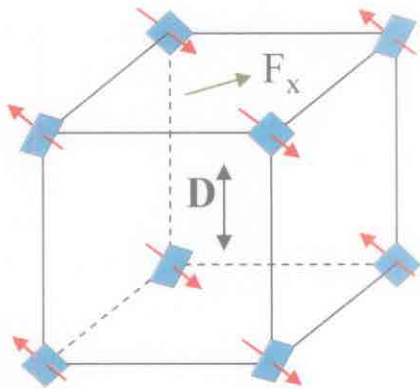
$$C_{OO}: T < T_{CG}$$

$$AF: xy^1 - xy^1, yz^1 - yz^1, zx^1 - zx^1$$

$$F: yz^1 - yz^0, zx^1 - zx^0$$

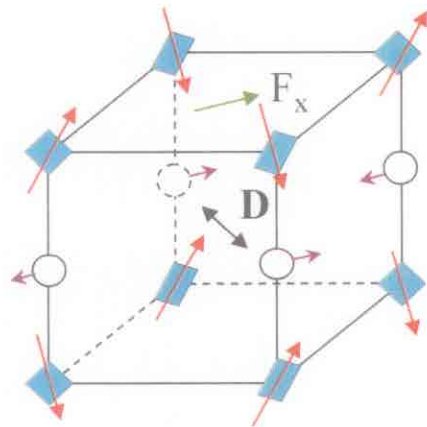
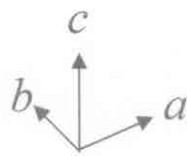
$$G_{OO} C_{SO}: T_{CG} < T < T_N$$

$$C_{OO} G_{SO}: T < T_{CG}$$



$$T_{CG} < T < T_N$$

$$C_y F_x$$



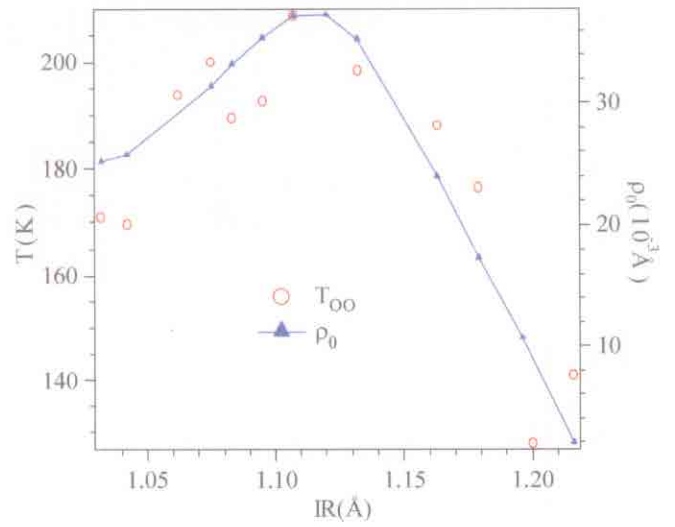
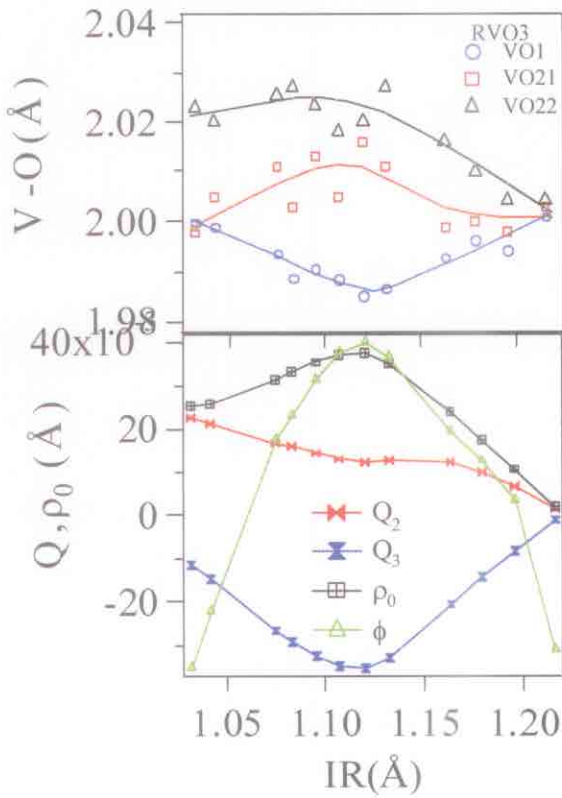
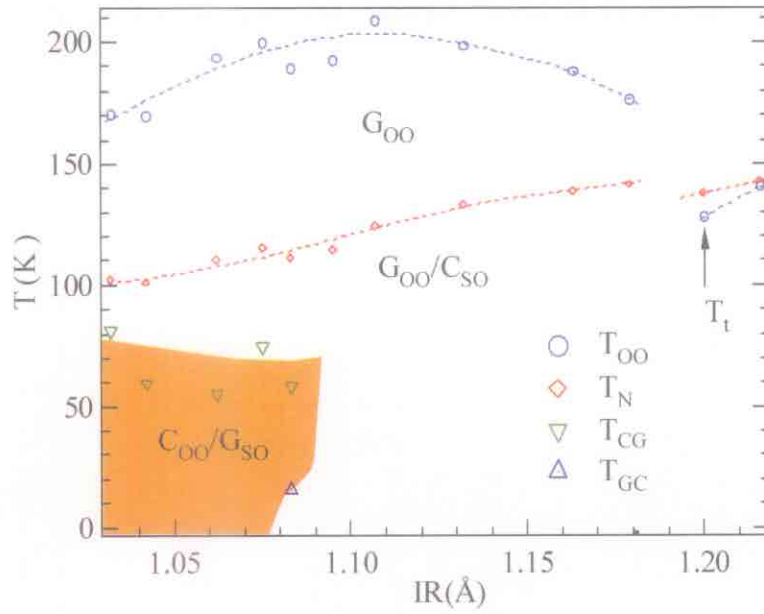
$$T < T_{CG}$$

$$G_z F_x$$

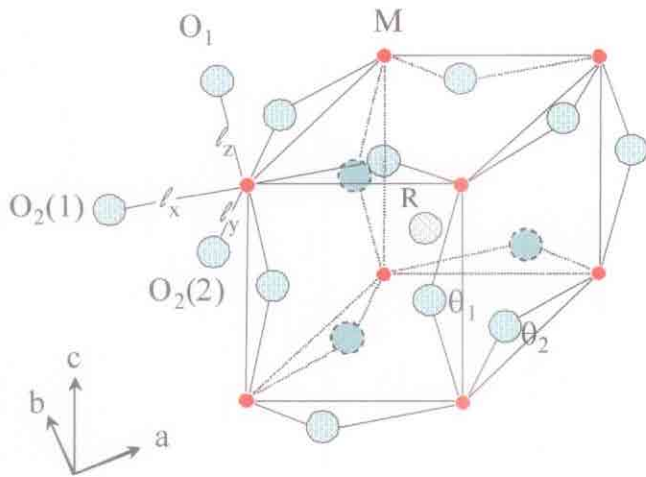
$$\mathbf{D}_{ij} \cdot \mathbf{S}_i \times \mathbf{S}_j \text{ gives } F_x$$

$$\text{VO}_{6/2} \text{ rotations also give } F_x \text{ for } T < T_{CG}$$

RVO₃



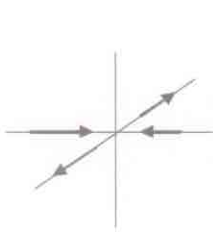
Jahn-Teller deformation in RMnO_3



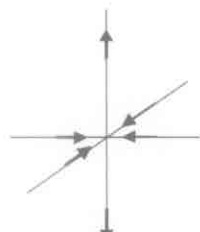
$$l_x = \mathbf{e}$$

$$l_z = \mathbf{m}$$

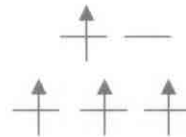
$$l_y = \mathbf{s}$$



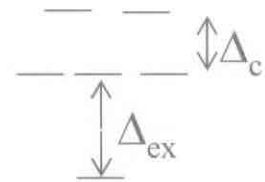
Q_2



Q_3

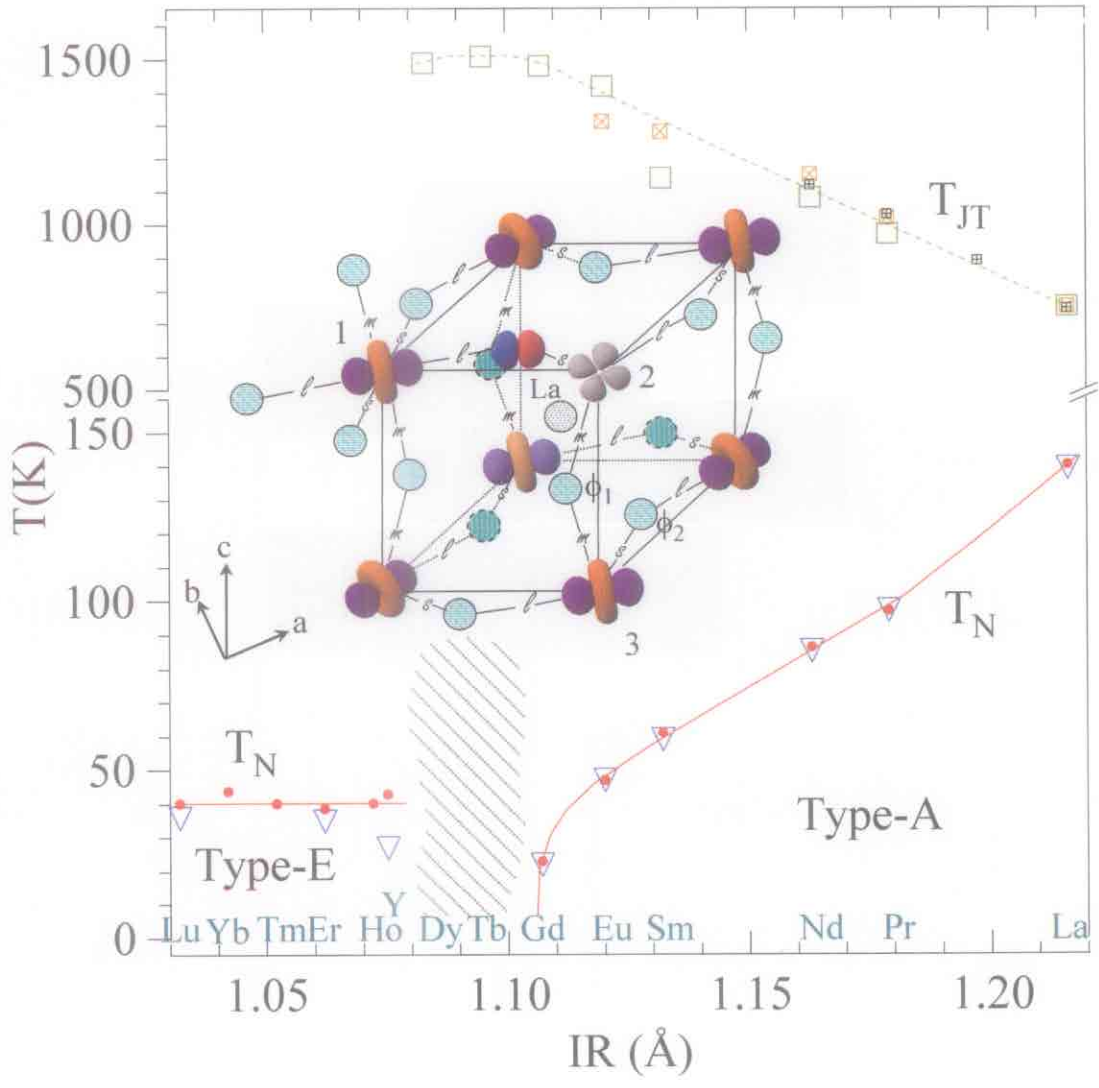


α -spin

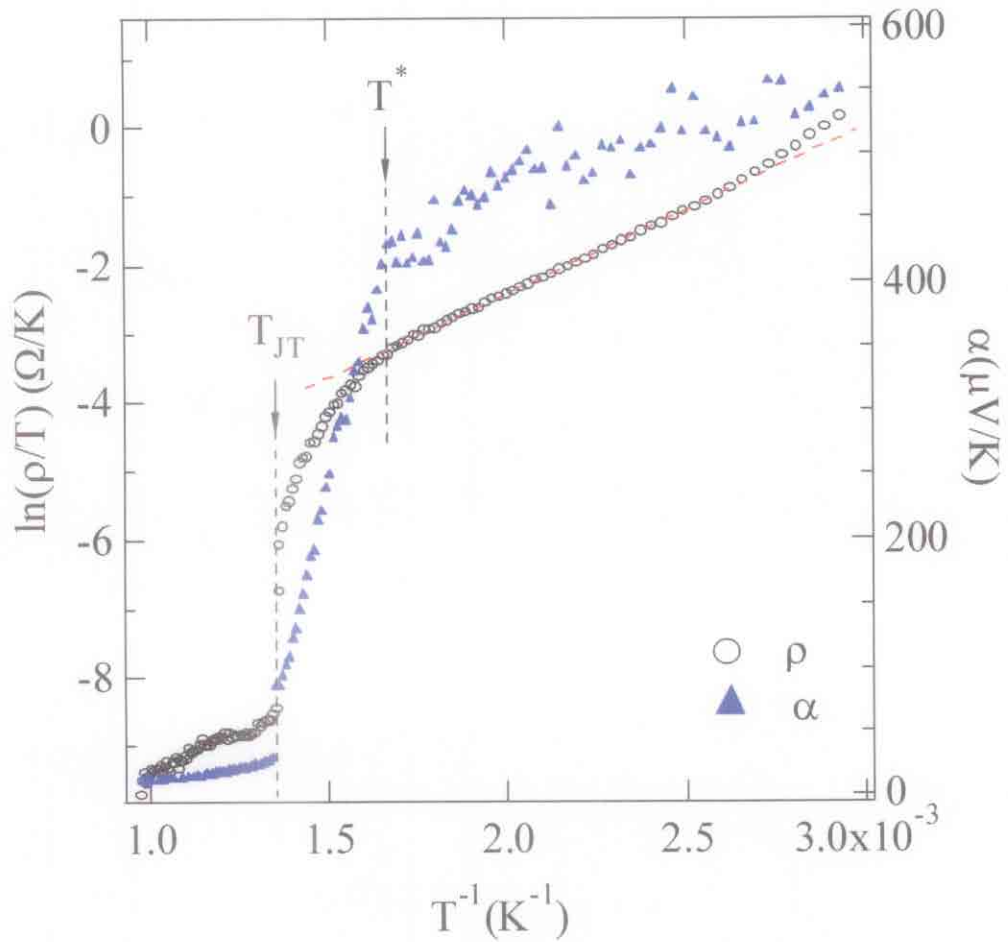


β -spin

RMnO₃



LaMnO₃



(Measured in vacuum)

$T > T_{JT}$: Orbital disorder
Partial $2Mn^{3+} \rightarrow Mn^{2+} + Mn^{4+}$

$T^* < T < T_{JT}$: short-range orbital disorder

Mn-O-Mn Superexchange

c axis: $e^0\text{-O-}e^0$ (AF) & $t^3\text{-O-}t^3$ (AF)

a-b plane: $t^3\text{-O-}t^3$ (AF) & $e^1\text{-O-}e^0$ (F)

For type-A magnetic order

$$T_N \sim b_\sigma^2 / \Delta_\sigma$$

Where

$$b_\sigma^2 \sim \cos^4(\gamma/2) \cos^4(\omega/2) |(3x^2 - r^2) |H'| (x^2 - z^2) |^2$$

$$\Delta_\sigma \approx \Delta_{JT} \sim kT_{JT}$$

$$\gamma = 30^\circ - \tan^{-1}(Q_3/Q_2)$$

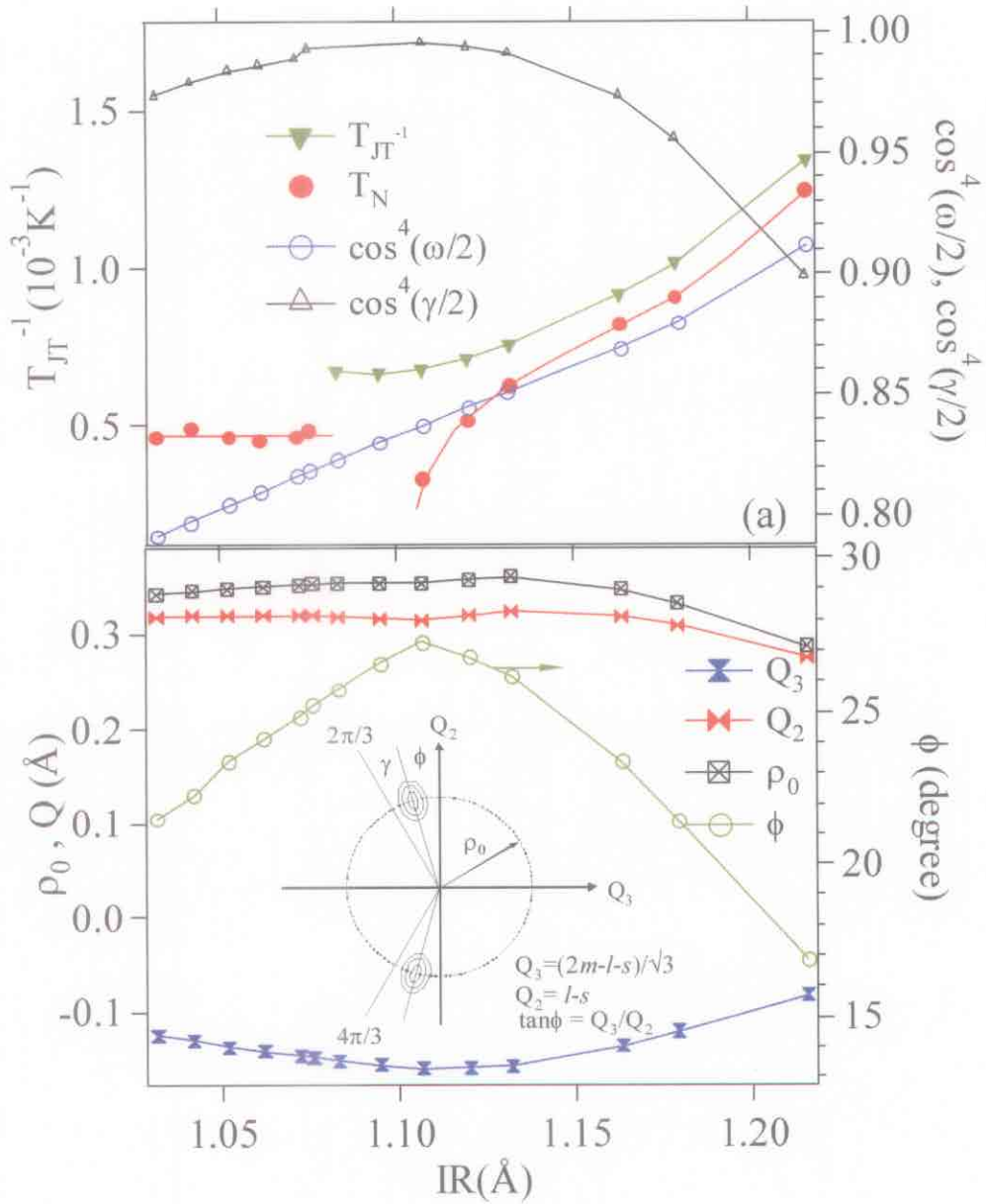
Note:

(1) $d[\ln \cos^4(\omega/2)]/d(\text{IR}) : d \ln T_N / d(\text{IR}) = 0.8 : 9$

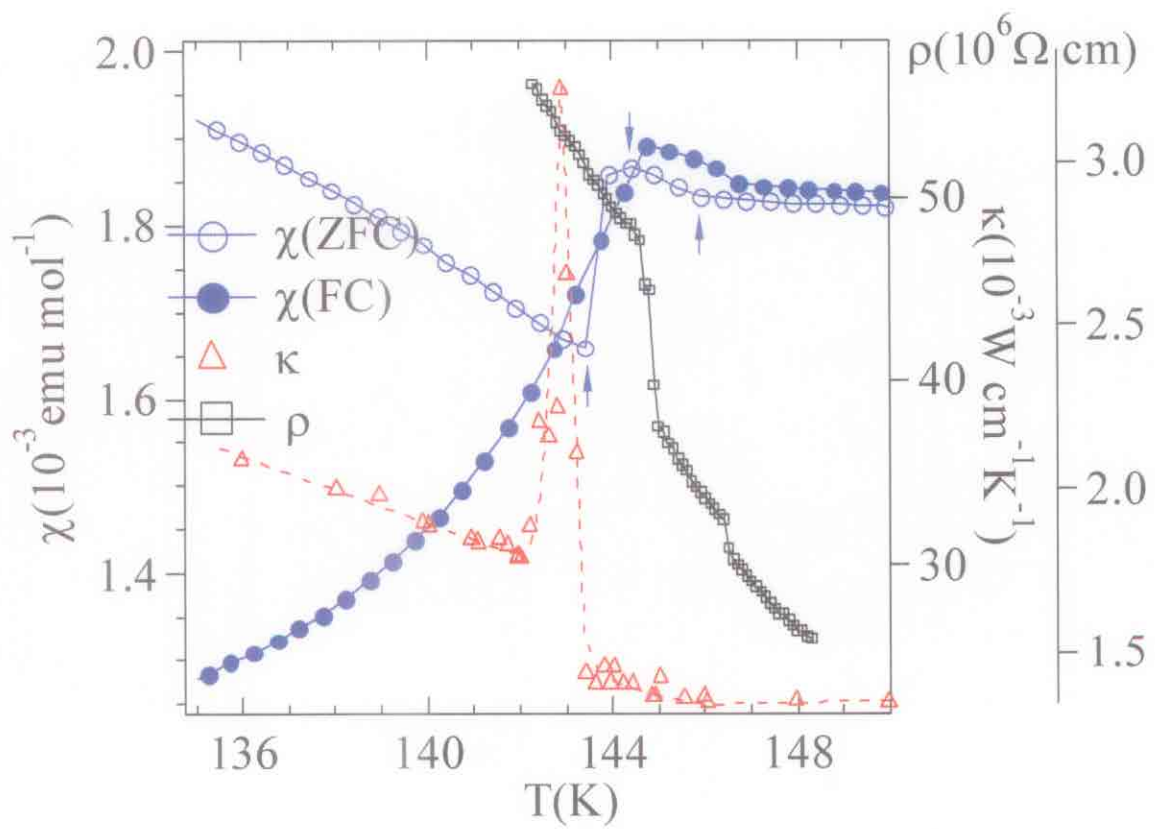
(2) $d[\ln \cos^4(\gamma/2)]/d(\text{IR})$ is negative, whereas $d \ln T_N / d(\text{IR})$ is positive.

(3) $d[\ln(1/T_{JT})]/d(\text{IR}) : d \ln T_N / d(\text{IR}) = 7 : 9$

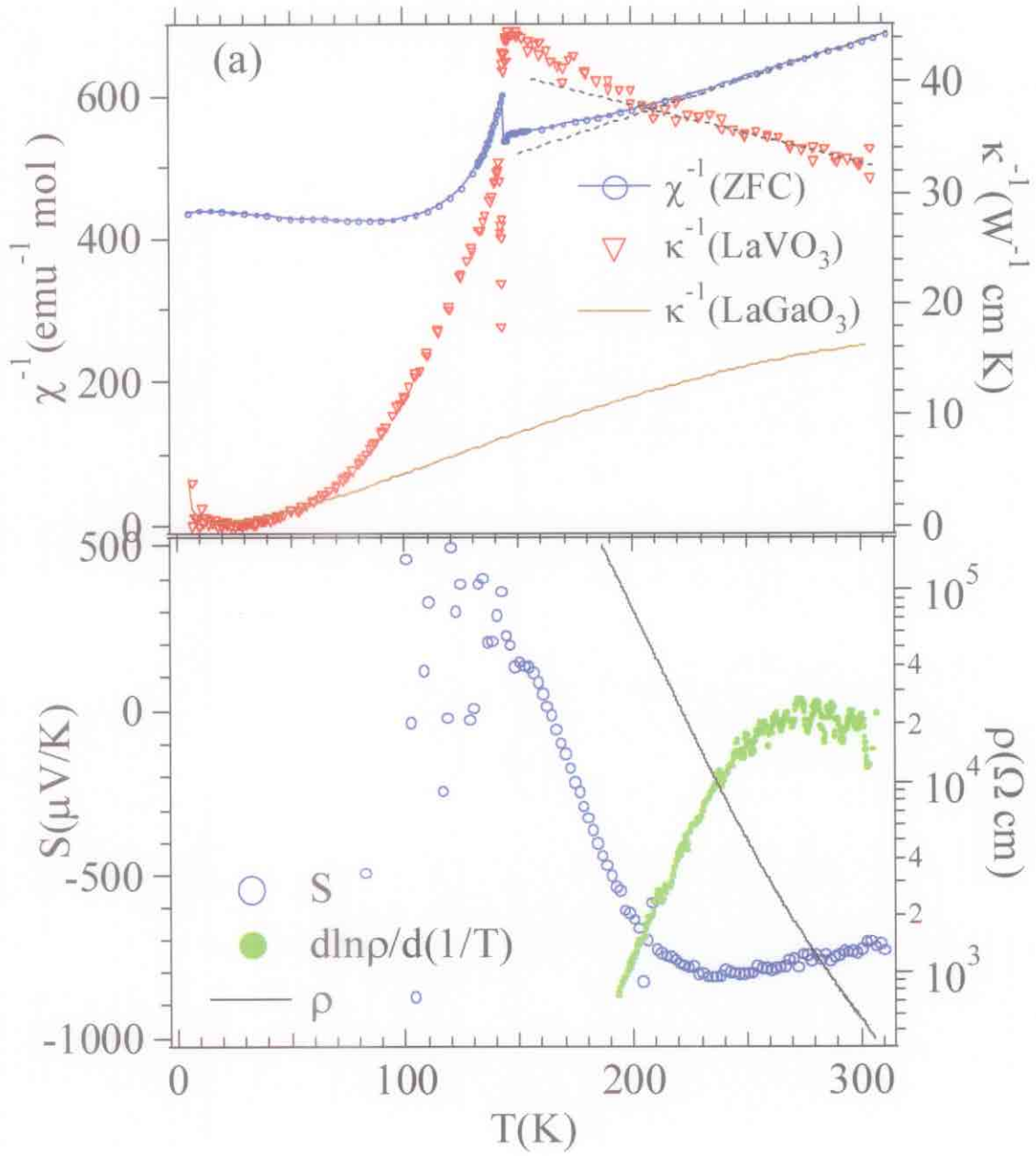
RMnO₃



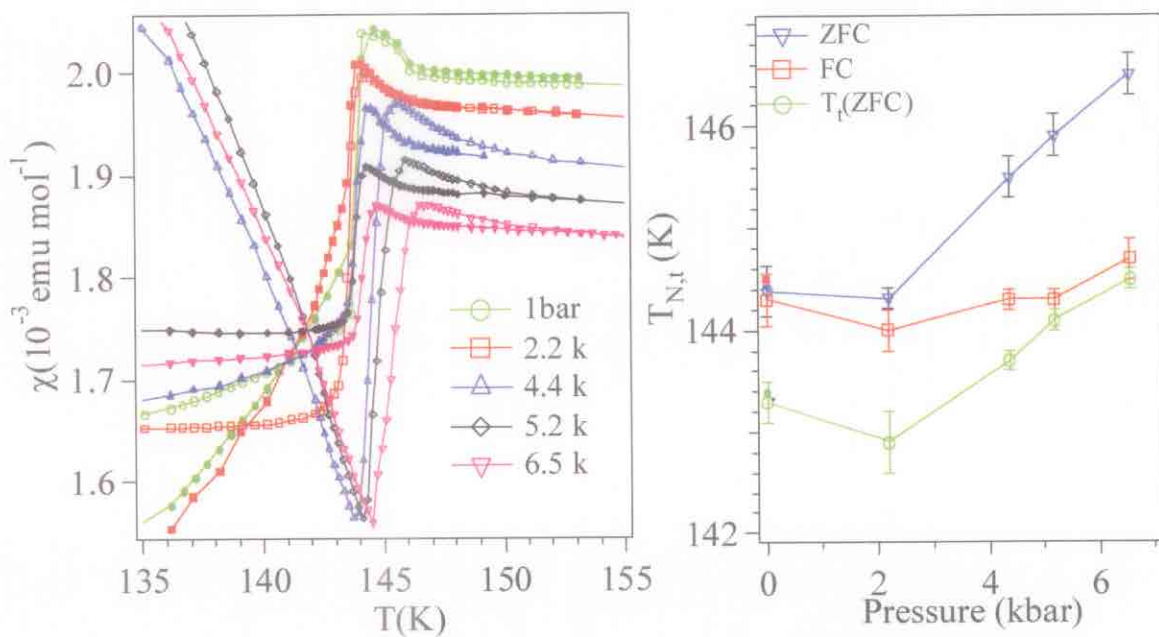
LaVO₃



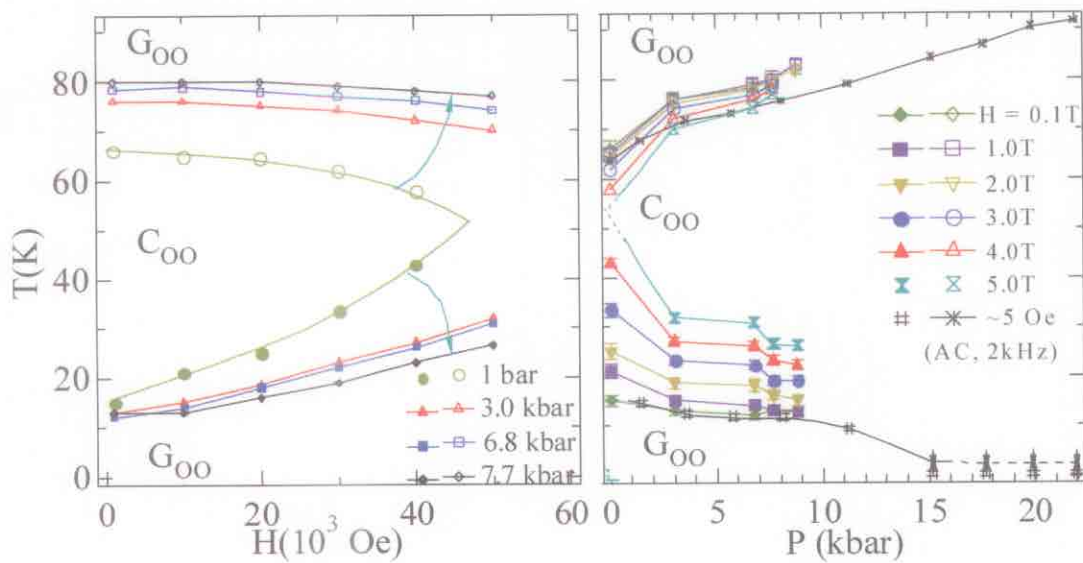
LaVO₃



LaVO₃



DyVO₃

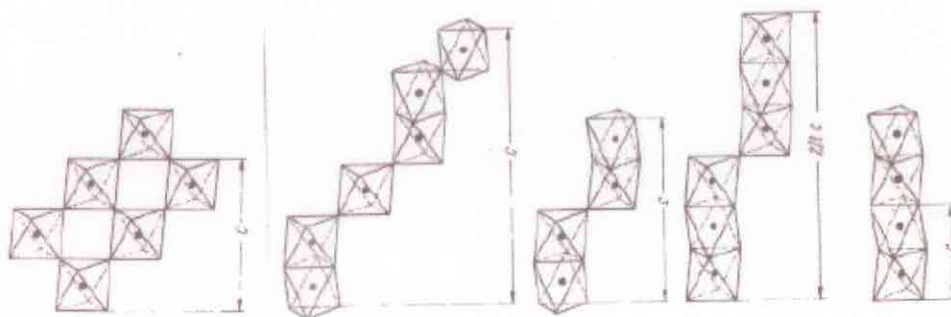
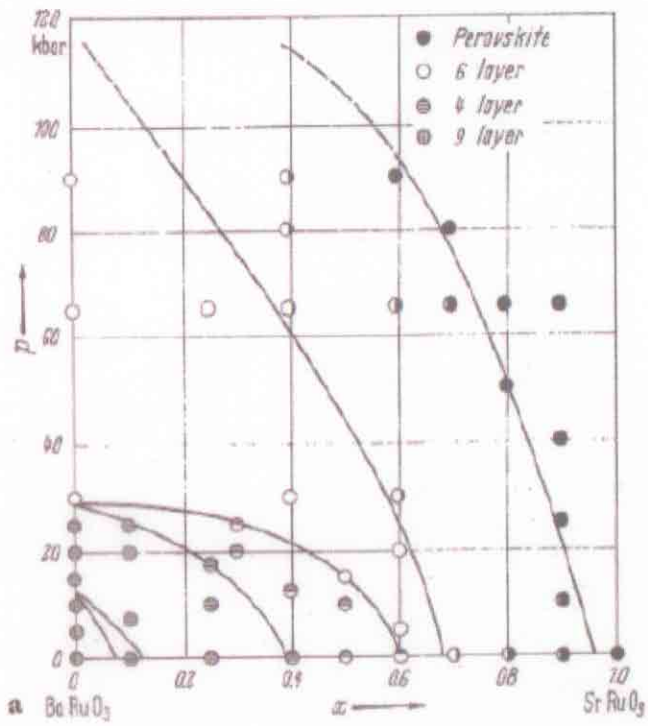


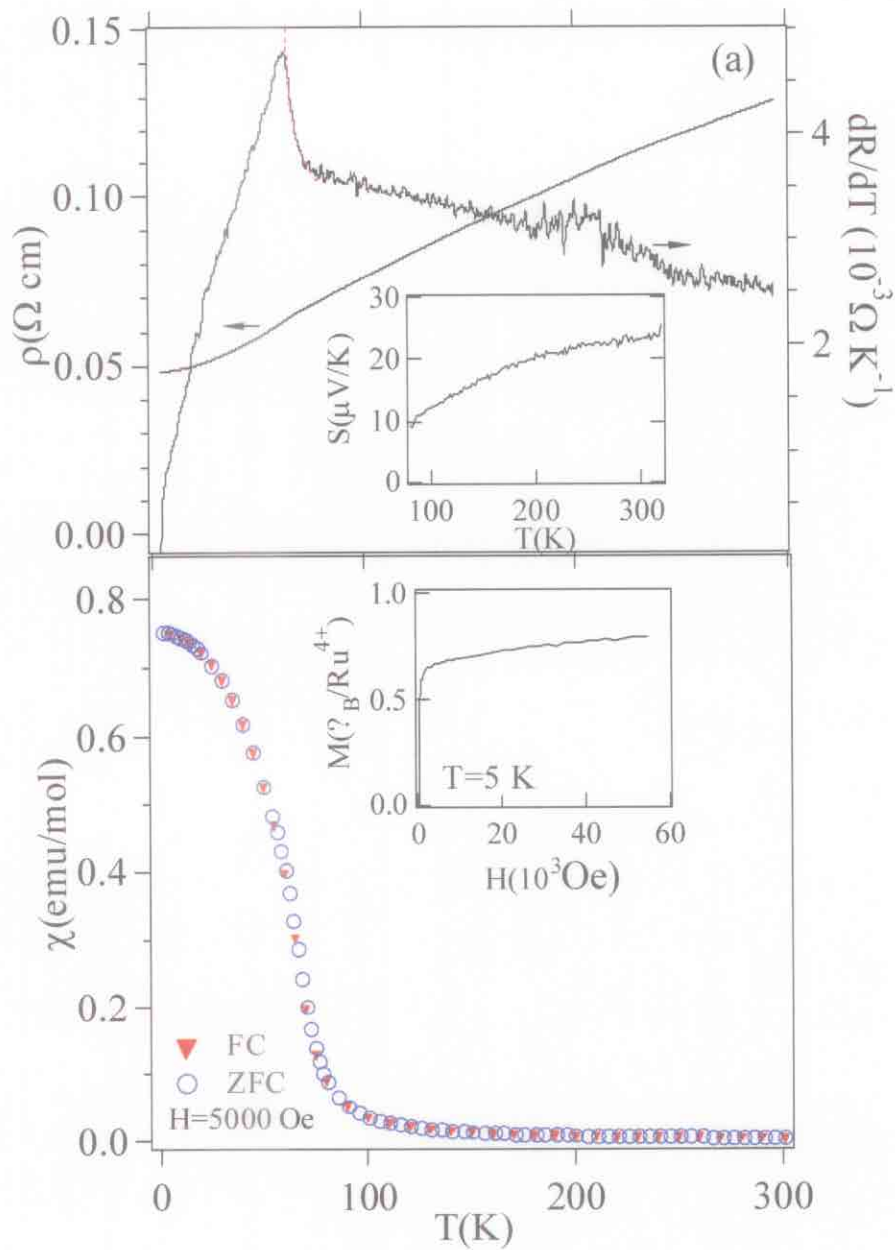
$\text{Sr}_{1-x}\text{Ca}_x\text{RuO}_3$

SrRuO₃: Itinerant-electron ferromagnet
Low-spin $\text{Ru}^{4+}:\pi^*4\sigma^*0$
 $T_c = 164 \text{ K}$, $\mu_{\text{Ru}} = 1.4 \mu_B$ @ 5 K
 $\mu_{\text{eff}} = 2.6 \mu_B/\text{Ru}$
($S = 1$ μ_{eff} (spin-only) = $2.8 \mu_B/\text{Ru}$)

Sr_{1-x}Ca_xRuO₃: Weiss $\theta > T_c$ for $x = 0$ decreases faster than T_c with increasing x . $T_c \rightarrow 0$ and θ changes sign at $x \approx 0.8$

CaRuO₃: No long-range magnetic order.
 Ti^{4+} substitution shows CaRuO_3 on verge of ferromagnetic order
[Cava & He, *Phys. Rev. B* **63**, 172403 (2001)]

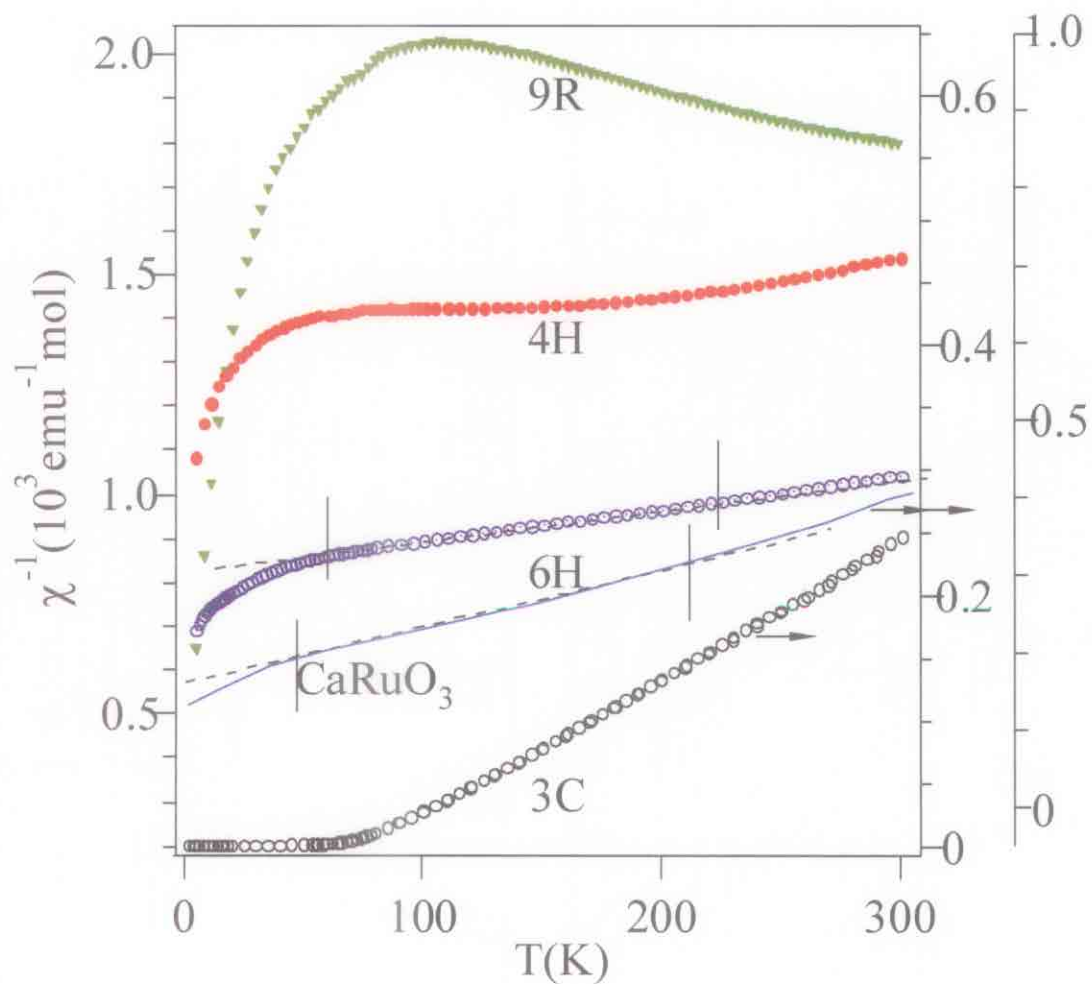


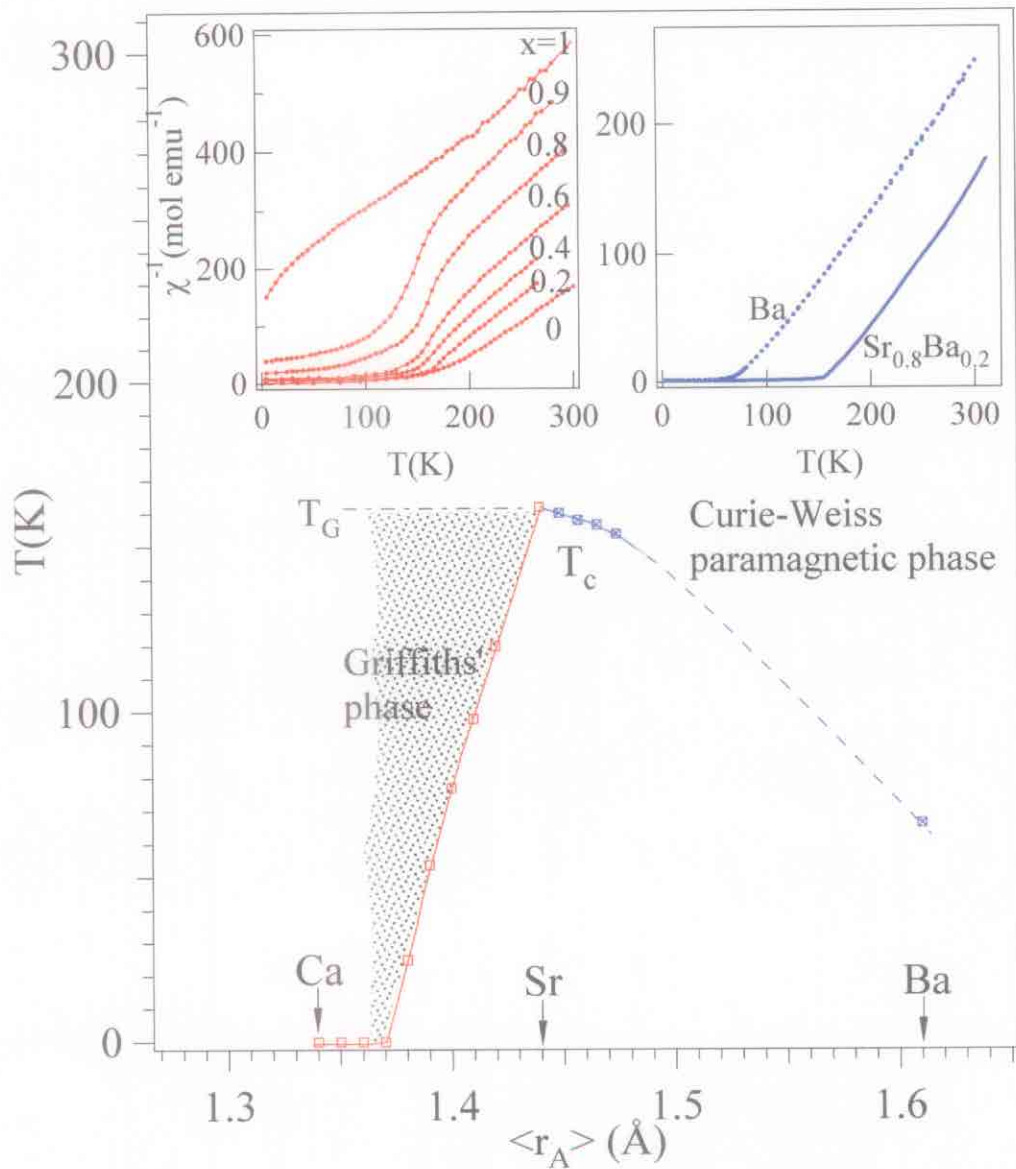


Cubic BaRuO₃: 1100°C @ 18 GPa KAWAI
 mutianvil press C.-Q. Jin *et al* (unpublished)

9R: Drillon *et al*, J. Chem. Soc. Faraday Transactions
75, 1193 (1979)

CaRuO₃: Yashimura *et al*, Phys. Rev. Lett. **83**, 4397
(1999)





C.-Q. Jin *et al* (unpublished)

Conclusions:

- Acidity of A-O bond: Ca > Sr > Ba
- Intraatomic $\lambda L \cdot S$ on a Ru⁴⁺ ion suppresses the interatomic spin-spin interaction in neighborhood of Ca.