TeV Scale SUSY – What Now?

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L. Roszkowski, KITP, UCSB, 9 July 2013

BayesFITS Group in Warsaw

- New research group at National Centre for Nuclear Research, formed in mid-2011
- > Funded by a 4.5-yr grant (~1.5MEuro) from Foundation for Polish Science
- Currently four postdocs, 1 PhD student, plus several local and external collaborators



LR

Shoaib Munir

Enrico Sessolo



Kamila Kowalska

Sebastian Trojanowski

- So far 9 papers out, 6 published or accepted (PRD, JHEP), more in pipeline
- **Research area:**
 - ``new physics'' (SUSY) and astroparticle physics (dark matter) in the LHC era
 - Early Universe, relics, etc
 - Flavor physics, ...

Outline

- \diamond Why TeV-scale for SUSY
- \diamond How to compare theory with data
- \diamond Implications of mh~126 GeV for favored SUSY mass scale
- \diamond Impact of DM relic abundance and searches
- \diamond CMSSM and beyond (CNMSSM, MSSM)
- ♦ Summary

Based on:

- Two ultimate tests of constrained SUSY, <u>1302.5956</u>
- The Constrained NMSSM with a 125 GeV Higgs boson -- A global analysis, <u>1211.1693</u>
- Di-photon rate enhancement in the NMSSM with nearly degenerate scalar and pseudoscalar Higgs bosons,

1305.0591

• Constrained MSSM favoring new territories: The impact of new LHC limits and a 125 GeV Higgs boson, 1206.0264 ...with updates



Where is SUSY?

After LHC(7/8TeV):

We know better now where SUSY is not.

Hints where SUSY may actually be.



BSM: hints from the LHC....



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Main news from the LHC so far...



➢ No (convincing) deviations
from the SM persymmetry or S
BR($\overline{B}_s \rightarrow \mu^+\mu^-$) = $(3.2^{+1.5}_{-1.2}) \times 10^{-9}$



Stringent lower limits on superparner masses

SUSY masses reaching 1 TeV scale+...



...and from the media...

Is Supersymmetry Dead?

The grand scheme, a stepping-stone to string theory, is still high on physicists' wish lists. But if no solid evidence surfaces soon, it could begin to have a serious PR problem

SCIENTIFIC AMERICAN[™]

April 2012

Nothing new...



The negative result illustrates the risks of Big Science, and its often sparse pickings.

By MALCOLM W. BROWNE

Three hundred and fifteen physicists worked on the experiment.

Their apparatus included the Tevatron, the world's most powerful particle accelerator, as well as a \$65 million detector weighing as much as a warship, an advanced new computing system and a host of other innovative gadgets.

But despite this arsenal of brains and technological brawn assembled at the Fermilab accelerator laboratory, the participants have failed to find their quarry, a disagreeable reminder that as science gets harder, even Herculean efforts do not guarantee success.

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27 August 2011 Last updated at 06:41 GMT

7.4K 🗲 Sharo 📑 💟 🗠 🖨

LHC results put supersymmetry theory 'on the spot'



By Pallab Ghosh Science correspondent, BBC News

Results from the Large Hadron Collider (LHC) have all but killed the simplest version of an enticing theory of sub-atomic physics.

Researchers failed to find evidence of so-called "supersymmetric" particles, which many physicists had hoped would plug holes in the current theory.

Theorists working in the field have told BBC News that they may have to come up with a completely new idea. Supersymmetry predicts the existence of mysterious super particles.

Data were presented at the Lepton Photon science meeting in Mumbai.

Related Stories

Energy, luminosity and the number of physicist failing to find SUSY have all increased by factor of 10...

Constrained SUSY – still alive?

The constrained MSSM (CMSSM) paradigm is "hardly tenable"

At Open Symposium of the European Strategy Preparatory Group, Krakow, Poland, 10-12 Sept. 2012

Constrained SUSY is in coma

A. Masiero, PLANCK-13

Really?



Constrained SUSY – still alive?

The constrained MSSM (CMSSM) paradigm is "hardly tenable"

At Open Symposium of the European Strategy Preparatory Group, Krakow, Poland, 10-12 Sept. 2012

Conventional susy models (CMSSM, NMSSM,) do not work as such and should finally rest in peace

F. Zwirner, Moriond EW (2013) summary talk

Constrained SUSY is in coma

A. Masiero, PLANCK-13







My conjecture:

(Coined before LHC era.)

SUSY cannot be experimentally ruled out.

It can only be discovered.

Or else abandoned.



SUSY: Constrained or Not?

• Constrained:

Low-energy SUSY models with unification relations among gauge couplings and (soft) SUSY mass parameters





figure from hep-ph/9709356

Virtues:

- Well-motivated
- Predictive (few parameters)
- Realistic

Many models:

- CMSSM (Constrained MSSM): 4+1 parameters
- NUHM (Non-Universal Higgs Model): 6+1
- CNMSSM (Constrained Next-to-MSSM) 5+1
- CNMSSM-NUHM: 7+1
- String-inspired, split, ``natural'', etc^{szkowski, KITP, UCSB, 9 July 2013}

Phenomenological:

Supersymmetrized SM...

Features:

- Many free parameters
- Broader than constrained SUSY



Many models:

- general MSSM over 120 params
- MSSM + simplifying assumptions
- pMSSM: MSSM with 19 params
- p9MSSM, p12MSSM, pnMSSM, ...

The 126 GeV SM-Like Higgs Boson

A blessing or a curse for SUSY?

The 126 GeV Higgs Boson and SUSY

A blessing...

Fundamental scalar -> SUSY
 Light and SM-like -> SUSY



L. Roszkowski, KITP, UCSB, 9 July 2013

The 126 GeV Higgs Boson and SUSY

A blessing...

Fundamental scalar -> SUSY
 Light and SM-like -> SUSY

Low energy SUSY prediction: Higgs with mass up to ~135 GeV

Constrained SUSY prediction: SM-like Higgs with mass up to ~130 GeV

The 126 GeV Higgs Boson and SUSY



L. Roszkowski, KITP, UCSB, 9 July 2013

 $M_{\rm SUSY} \equiv \sqrt{m_{\tilde{t}_1} m_{\tilde{t}_2}}$ 16

How to compare theory with experiment

Rigid step-function application of limits/allowed ranges (e.g. DM relic abundance, etc) Mahmoudi et al, Hewett et al, ...
 Frequentist (chi^2-based) MasterCode, Fittino, ...
 Bayesian BayesFITS, Allanach, SuperBayes, Balazs, Kraml...

Frequentist: "probability is the number of times the event occurs over the total number of trials, in the limit of an infinite series of equiprobable repetitions"

Bayesian: "probability is a measure of the degree of belief about a proposition"

Both F and B are based on the likelihood function.



The Likelihood function

Positive measurements: Take a single observable $\xi(m)$ that has been measured $(e.g., M_W)$ c – central value, σ – standard exptal error define $\chi^2 = \frac{[\xi(m) - c]^2}{c^2}$ assuming Gaussian distribution $(d \rightarrow (c, \sigma))$: $\mathcal{L} = p(\sigma, c | \boldsymbol{\xi}(m)) = rac{1}{\sqrt{2\pi}\sigma} \exp\left[-rac{\chi^2}{2} ight]$ when include theoretical error estimate τ (assumed Gaussian): $\sigma \rightarrow s = \sqrt{\sigma^2 + \tau^2}$ TH error "smears out" the EXPTAL range for several uncorrelated observables (assumed Gaussian): $\mathcal{L} = \exp\left[-\sum_{i} \frac{\chi_{i}^{2}}{2}\right]$

Bayes

L. Roszkowski, KITP, UCSB, 9 July 2013

Limits: **Helalive Likelhood** 0.6 0.4 1 1.1 1.2 1.3 0.7 0.8 0.9 Smear out bounds. Add theory error.

Central object: Likelihood function

LHC direct limits:

Need careful treatment. Typically use Poisson.

Bayesian statistics



Bayes theorem:

$$p(m|d) = rac{p(d|m) \, \pi(m)}{p(d)}$$

Prior $\pi(m)$ – what we know about the model m before seeing the data dLikelihood p(d|m) – the probability of obtaining data d if model m is true Posterior p(m|d) – the probability about m after seeing d. Evidence p(d) – normalization factor, important for model comparison





Minimum chi2 approach: find best-fit and draw confidence regions about it L. Roszkowski, KITP, UCSB, 9 July 2013

Constrained Minimal Supersymmetric Standard Model (CMSSM)





figure from hep-ph/9709356

At $M_{\rm GUT} \simeq 2 \times 10^{16} \, {\rm GeV}$:

- ${}^{ {}_{ { \hspace{-.1em} I} }}$ gauginos $M_1=M_2=m_{\widetilde{g}}=m_{1/2}$
 - scalars $m_{\widetilde{q}_i}^2 = m_{\widetilde{l}_i}^2 = m_{H_b}^2 = m_{H_t}^2 = m_0^2$

9 3-linear soft terms
$$A_b = A_t = A_0$$

- $\begin{array}{l} \bullet \quad \text{radiative EWSB} \\ \mu^2 = \frac{m_{H_b}^2 m_{H_t}^2 \tan^2 \beta}{\tan^2 \beta 1} \frac{m_Z^2}{2} \end{array}$
- five independent parameters: $m_{1/2}, m_0, A_0, \tan\beta, \operatorname{sgn}(\mu)$
- well developed machinery to compute masses and couplings





CMSSM: numerical scans

- Perform random scan over 4 CMSSM +4 SM (nuisance) parameters <u>simultaneously</u>
- Very wide ranges: $100 ext{ GeV} \leq m_0 \leq 20 ext{ TeV}$ $100 ext{ GeV} \leq m_{1/2} \leq 10 ext{ TeV}$ $-20 ext{ TeV} \leq A_0 \leq 20 ext{ TeV}$ $3 \leq ext{ tan } eta \leq 62$

 Use Nested Sampling algorithm to evaluate posterior

Nuisance	Description	Central value \pm std. dev.	Prior Distribution
M_t	Top quark pole mass	$173.5 \pm 1.0 \mathrm{GeV}$	Gaussian
$m_b(m_b)_{ m SM}^{\overline{MS}}$	Bottom quark mass	$4.18\pm0.03{\rm GeV}$	Gaussian
$\alpha_s(M_Z)^{\overline{MS}}$	Strong coupling	0.1184 ± 0.0007	Gaussian
$1/\alpha_{\rm em}(M_Z)^{\overline{MS}}$	Inverse of em coupling	127.916 ± 0.015	Gaussian

Use 4 000 live points

Use Bayesian approach (posterior)



Hide and seek with SUSY

The experimental measurements that we apply to constrain the CMSSM's parameters. Masses are in GeV.

	Measurement	Mean or Range	Error: (Exp., Th.)	Distribution
?→	Combination of:			
	CMS razor 4.4/fb , $\sqrt{s}=7{\rm TeV}$	See text	See text	Poisson
	CMS $\alpha_T \ 11.7/\text{fb}$, $\sqrt{s} = 8 \text{ TeV}$	See text	See text	Poisson
	m_h by CMS	$125.8{ m GeV}$	$0.6{ m GeV}, 3{ m GeV}$	Gaussian
	$\Omega_\chi h^2$	0.1120	0.0056,10%	Gaussian
	$\delta \left(g-2 ight)^{ m SUSY}_{\mu} imes 10^{10}$	28.7	8.0, 1.0	Gaussian
	$\mathrm{BR}\left(\overline{B} \to X_s \gamma\right) \times 10^4$	3.43	0.22, 0.21	Gaussian
	$BR(B_u \to \tau \nu) \times 10^4$	1.66	0.33, 0.38	Gaussian
	ΔM_{B_s}	$17.719{\rm ps}^{-1}$	$0.043 \mathrm{ps^{-1}},\ 2.400 \mathrm{ps^{-1}}$	Gaussian
	$\sin^2 heta_{ m eff}$	0.23116	0.00012, 0.00015	Gaussian
	M_W	80.385	0.015, 0.015	Gaussian
	$BR(B_s \to \mu^+ \mu^-)_{current} \times 10^9$	3.2	+1.5 - 1.2, 10% (0.32)	Gaussian
	BR $(B_s \to \mu^+ \mu^-)^{\text{currents}}_{\text{proj}} \times 10^9$	$3.5(3.2^*)$	$0.18~(0.16^*),~5\%~[0.18~(0.16^*)]$	Gaussian

SM value: $\simeq 3.5 \times 10^{-9}$

10 dof



SUSY - most important constra af the LH

Higgs mass

CMS: $m_h \sim 125.8 \text{ GeV}$ (in ZZ); $m_h = 124$. ATLAS: $m_h = 124.3 \text{ GeV}$ (in ZZ); $m_h =$

Direct search limits

B s -> mu mu

$$\mu = 1.7^{+0.5}_{-0.4} \text{ ATLAS} \quad \mu = 0.91^{+0.3}_{-0.24} \text{ CMS}$$

limit...

> Dark matter density $\mu = 1.5 \pm 0.6$ ATLAS $\mu = 0.76 \pm 0.21$ CMS Positive measurement, inconsiste

Thursday, March 7, 2013

Lower

LHCb (Nov 2012)

The

> Other flavor (b to s gamma, etc)

EW observables (M_W,...)



V,N

g fusior



~126 GeV Higgs in SUSY

- In SUSY m_h is a calculated quantity.
- 1-loop corr: positive, up to ~45 GeV

$$\Delta m_h^2 = \frac{3m_t^4}{4\pi^2 v^2} \left[\ln\left(\frac{M_{\rm SUSY}^2}{m_t^2}\right) + \frac{X_t^2}{M_{\rm SUSY}^2} \left(1 - \frac{X_t^2}{12M_{\rm SUSY}^2}\right) \right]$$

• 2-loop corr: negative, ~3 GeV

two most complete calculations differ by a 2-5 GeV (DR-bar (Slavich,...) used in SoftSusy, Spheno, Suspect, and on-shell (Hollik,...) in FeynHiggs

```
Substantial theory error!
```



$$M_{\rm SUSY} \equiv \sqrt{m_{\tilde{t}_1} m_{\tilde{t}_2}}$$
$$X_t = A_t - \mu \cot \beta$$

0 X (TeV)

Diouadi. arXiv:hep-ph/0503173



• (3-loop corr: positive, <~2-3 GeV)

P. Kant J. Feng, et al

Two ways to obtain m_h~126 GeV: 1. increase M_SUSY -> heavy superparners! or

2. take large |X_t|~|A_tststop__12 atv213TeV

Applies to SUSY generally, not just constrained models.

~126 GeV Higgs in the CMSSM

0.2

0.0

Include only m_h~126 GeV ۲

and lower limits from direct $m_{h_1} \simeq m_{h_2} \simeq 125.3$ SUSY searches

$$\mathcal{L} \sim e^{rac{(m_h-125.8\,{
m GeV})^2}{\sigma^2+ au^2}}$$

 $\sigma = 0.6 \text{ GeV}, \tau = 2 \text{ GeV}$

We use DR-bar approach (SoftSusy). It gives larger m h.

~126 GeV Higgs mass implies multi-TeV scale for SUSY

Consistent with:

- SUSY direct search lower limits at LHC ٠
- constraints from flavor ٠

A curse...

 $m_{h_2} \simeq 125.3$





~126 GeV Higgs in the CMSSM



~126 GeV Higgs mass implies multi-TeV SUSY masses

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If m_h were, e.g., 116 GeV...



...significant tension with LHC bounds





(Abreu, ...)

Generically ~126 GeV Higgs mass implies multi-TeV scale for SUSY

A weak upper bound on M_SUSY

Even weaker for small tanbeta





Dark matter density

 Unified SUSY: neutralino relic density is typically 1-2 orders of magnitude too large



BayesFITS (2013)

 $[cv - 2\sigma, cv + 2\sigma]$

2 sigma range

15000

20000









5000

10000

m_o (GeV)



10000

8000

m^{1/2} (GeV)

2000

0

Relic density

CMSSM. u >0

CMSSM: these are the only DM-favored regions

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~126 GeV Higgs vs stop mass



Best fit to ~126 GeV Higgs for M_SUSY~1 TeV or >> 1 TeV

best-fit point $\chi^2_{\rm min}/{
m dof} = 18.26/10$

 $[\chi^2_{
m min}/{
m dof}\simeq 4/9$ when drop $(g-2)_\mu]$

- Dark matter density selects specific regions
- CMSSM: DM regions (almost) disconnected
 other models: they overlap

1302.5956

Can such multi-TeV ranges of SUSY parameters be experimentally tested?



Are we done with the LHC?



There is more out there than meets the eye



Direct Detection AD 2011 - Before LHC





CMSSM and 1-tonne DM detectors



1-tonne DM detectors to cover most of CMSSM predictions

... over ALL multi-TeV ranges of mass parameters

(Except for some cases at mu<0)

LUX (2014) to improve sensitivity by ~1 decade

Generic prediction of multi-TeV SUSY: ~1TeV LSP (higgsino)



- Other flavor (b to s gamma, etc)
- > EW observables (M_W,...)







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= 45

, $tan\beta$

GeV

1.4

(TeV)

7

40

SUSY

Effect of precise $BR(\overline{B}_s \to \mu^+ \mu^-)$



If $BR(\overline{B}_s \to \mu^+ \mu^-) \simeq SM$ value with 5-10% precision (both TH and EXPT)

 \Rightarrow A funnel region gone





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Effect of precise $BR(\overline{B}_s \to \mu^+ \mu^-)$



If $BR(\overline{B}_s \to \mu^+ \mu^-) \simeq SM$ value with 5-10% precision \Rightarrow A funnel region gone

Ways to rule out the CMSSM:

- No DM signal in 1-tonne detectors
- DM signal at ~500 to 750

GeV

SC: for $\mu < 0 \ \sigma_p^{\text{SI}}$ lower (cancellations)

NUHM, CNMSSM: similar ranges of sigma_p but DM-favored regions overlap

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Constrained SUSY is alive and well...

 Even the simplest unified SUSY model (CMSSM) is <u>consistent</u> with all data (Higgs mass, DM relic density, direct limits, flavor-violating processes, ...)

...except for g-2, R(gamma gamma)

- M_SUSY >~ (or even >>) 1 TeV favored by ~126 GeV Higgs
- In less unified models somewhat lower SUSY masses are allowed (but not by much)

...except for very fine tuned corners

CMSSM and beyond



Generally: ~126 GeV Higgs: need ~>1TeV, typically multi-TeV M_SUSY scale

Warning: different models scanned with different ranges, precision, completeness, etc

CNMSSM...

Constrained Non-Minimal SSM

- The MSSM suffers from the mu-problem: $\mu \sim \mathcal{O}(M_{\mathrm{SUSY}})$
- The NMSSM extends the MSSM by addition of a gauge singlet superfield S

$$W = \lambda S H_u H_d + \frac{\kappa}{3} S^3 + (\text{Yukawa couplings})$$

S gets a vev $\,s=\langle S
angle\,\,$ ==> Superpotential develops an effective term $\mu_{
m eff}=\lambda s$

• Extra terms in the soft SUSY-breaking Lagrangian for the Higgs Sector

$$V_{
m soft} = m_{H_u}^2 |H_u|^2 + m_{H_d}^2 |H_d|^2 + m_S^2 |S|^2 + \left(\lambda A_\lambda S H_u H_d + \frac{1}{3} \kappa A_\kappa S^3 + \text{h. c.}
ight)$$

Higgs and neutralino sectors extended wrt MSSM:+1 CP-even scalar
+1 CP-odd pseudoscalar
+1 neutralinoLeads to interesting phenomenology+1 neutralino

Richer Higgs phenomenology than MSSM:

 $\begin{array}{l} \rightarrow \ m_{h_1} \simeq 125 \quad {\rm GeV}, \ h_2 \ undetected \\ \rightarrow \ m_{h_2} \simeq 125 \quad {\rm GeV}, \ h_1 \ undetected \\ \rightarrow \ m_{h_1} \simeq \ m_{h_2} \simeq 125 \ {\rm GeV} \\ \rightarrow \ m_{h_1} \simeq \ m_{a_1} \simeq 125 \ {\rm GeV} \end{array} \xrightarrow{\bullet} particularly interesting case \\ {\rm L. Roszkowski, \ KITP, \ UCSB, 9 \ July \ 2013 } \end{array}$

Global Bayesian analysis of the CNMSSM

CNMSSM parameters:

arXiv:1211.1693

 $m_0, m_{1/2}, A_0, \tan\beta, \lambda, \operatorname{sgn}(\mu_{\text{eff}})$

Measurement	Mean or range	Error (Exp., Th.)	Distribution
CMS razor 4.4/fb	Likelihood map		Poisson
$m_{h_{\rm sig}}$ (GeV)	125.8	0.6, 3	Gaussian
$R_{h_{\rm sig}}(\gamma\gamma)$	1.6	0.4, 15%	Gaussian
$R_{h_{\rm sig}}(ZZ)$	0.80	+0.35-0.28,15%	Gaussian
$m_{h_{\rm hid}}$ (GeV)	< 122.7, > 128.9	0, 3	Error Fn
$R_{h_{\rm hid}}(X)$	$\mu_{95}(X)$ from CMS	0, 15%	Error Fn
$\Omega_{\chi} h^2$	0.1120	0.0056, 10%	Gaussian
$\delta \left(g-2 ight)^{ m SUSY}_{\mu} imes 10^{10}$	28.7	8.0, 1.0	Gaussian
$\mathrm{BR}\left(\overline{\mathrm{B}} \to \mathrm{X_s}\gamma\right) \times 10^4$	3.43	0.22, 0.21	Gaussian
$BR(B_u \rightarrow \tau \nu) \times 10^4$	1.66	0.66, 0.38	Gaussian
$\Delta M_{B_s} (\mathrm{ps}^{-1})$	17.719	0.043, 2.400	Gaussian
$BR(B_s \to \mu^+\mu^-)$	3.2×10^{-9}	+1.5-1.2, 10%	Gaussian



LHC constraints:

CMS Razor limit on (m_0,m_1/2)

• For a `signal' Higgs: $\mathcal{L}_{\psi_i}(h_{\text{sig}}) = \exp\left[(\psi_i(\text{obs}) - \psi_i(h_{\text{sig}}))^2/2(\tau_{\psi_i}^2 + \sigma_{\psi_i}^2)\right]$

$$\psi_i = m(h_{sig}), R_X(h_{sig}); R_X(h_{sig}) = \frac{\sigma(pp \to h_{sig})}{\sigma(pp \to h_{SM})} \times \frac{BR(h_{sig} \to X)}{BR(h_{SM} \to X)}$$

 $R_X(obs) \simeq \frac{\sigma_X}{\sigma_{SM}}$ (CMS) for $X = \gamma\gamma, ZZ$

• For the unobserved Higgs: an exclusion likelihood for mass and R_X $X = \gamma \gamma, ZZ, \tau \tau, WW$

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Constrained NMSSM



arXiv:1211.1693

Constrained NMSSM



...excluded by Xenon100 limit on DM

arXiv:1211.1693

Constrained NMSSM

 $\mathrm{Case}\; m_{h_1} \simeq m_{a_1} \simeq 126\,\mathrm{GeV}$

arXiv:1305.0591

- requires relaxing strict CNMSSM conditions $m_{H_u}, m_{H_d}, m_S \neq m_0$ (CNMSSM-NUHM) $A_{\lambda} = A_{\kappa} \neq A_0$
- Goal: boost $\gamma\gamma$ mode but keep ZZ/WW SM-like





[GeV

Ŕ

¥ 0.35

0.3

0.25

0.2 0.4

127

126 29 125

f⁼124

122

121

120

122

123100124



Up to 60% enhancement in di-photon mode!

Bonus: significant enhancement also in bb and tau-tau decay modes (all the cases)



Three solutions:

4000

non100

LUX

Kenon1T

122087 K 2n2 5 01 37 130 GeV

1 45 4 Bb < Rbb

5009

1600

700

127400128

6000

<u>69990 -4000 0.52000 90.55 2000</u>

LSP higgsino[®]like

 $1.15 < R^{bb} < 1.3$

1360

ma, [GeVnhy0 [GeV]

2005

CNMSSM NUHM

Singlino-higgsinox

0.2

0.1



 $1.45 < R^{bb}$

LSP bino-higgsino like (Focus Point)





CNMSSM-NUHM $\mu < 0$ $122 \text{ GeV} < m_{h_1}, m_{a_1} < 130 \text{ GeV}$ FP region $5 < R_{rr,bb}^{bb} < 10$ $1 < R_{rr,bb}^{bb} < 2$ $2 < R_{\tau\tau,bb}^{bb} < 5$ • $10 < R_{\tau\tau,bb}^{bb}$ [GeV] -100 $+ \kappa s$ -200 A^{SUSY} + -300 -400 -500 0.1 0.6 0 0.2 0.3 0.4 0.5 к

50

Need associated b-bbar-h production

All this can be realized in associated bbar H production mode, and not gluon-fusion mode!

It is tiny in the SM, but tanb^2-enhanced in the MSSM.

In all 3 cases, tan can be large, e.g.:



(V. Sharma)





~126 GeV Higgs in general MSSM

More parameters, more freedom

...here 9 parameters: p9MSSM



...except for very fine tuned corners which allow much lighter staus, stops,

p9MSSM: Constraints

Measurement	Mean or range	Error: exp., th.	Distribution	
CMS $\alpha_T \ 11.7/\text{fb}$, $\sqrt{s} = 8 \text{ TeV}$	See text.	See text.	Poisson	
m_h (by CMS)	$125.8\mathrm{GeV}$	$0.6{ m GeV}, 3{ m GeV}$	Gaussian	
$\mid \Omega_{\chi} h^2$	0.1199	0.0027,10%	Gaussian	
$BR(\overline{B} \to X_s \gamma) \times 10^4$	3.43	0.22, 0.21	Gaussian	
$BR(B_u \to \tau \nu) \times 10^4$	1.66	0.33, 0.38	Gaussian	
ΔM_{B_s}	$17.719\mathrm{ps}^{-1}$	$0.043 \mathrm{ps}^{-1}, \ 2.400 \mathrm{ps}^{-1}$	Gaussian	
$\sin^2 heta_{ m eff}$	0.23146	0.00012, 0.00015	Gaussian	A
M_W	$80.399{ m GeV}$	$0.023{ m GeV},0.015{ m GeV}$	Gaussian	
$ BR (B_s \to \mu^+ \mu^-) \times 10^9$	3.2	+1.5, -1.2, 10%	Gaussian	
$m_b(m_b)^{\overline{MS}}$	$4.18\mathrm{GeV}$	$0.03\mathrm{GeV},0$	Gaussian	
M_t	$173.5\mathrm{GeV}$	$1.0\mathrm{GeV},0$	Gaussian	
$\delta (g-2)^{\text{SUSY}}_{\mu} \times 10^{10}$	28.7	8.0, 1.0	Gaussian	
XENON100 (2012)	See text.	See text.	Poisson	
$CMS \ 3l + E_T^{miss} \ 9.2/fb, \sqrt{s} = 8 \text{ TeV}$	See text.	See text.	Poisson]

Table 2: The experimental constraints that we include in our likelihood functions to constrain our p9MSSM model. We denote the first block of constraints as **basic**.

Now include (optionally) DM direct detection limit in the likelihood function

Plus LEP, Tevatron:

$$\begin{array}{lll} m_{\chi} &> \ 46\,{\rm GeV}, \\ m_{\tilde{e}} &> \ 107\,{\rm GeV}, \\ m_{\tilde{g}} &> \ 500\,{\rm GeV}, \\ m_{\chi_{1}^{\pm}} &> \ 94\,{\rm GeV} \ {\rm if} \ {\rm m}_{\chi_{1}^{\pm}} - {\rm m}_{\chi} > 3\,{\rm GeV} \ {\rm and} \ {\rm tan} \ \beta < 40 \\ m_{\tilde{\mu}} &> \ 94\,{\rm GeV} \ {\rm if} \ {\rm m}_{\tilde{\mu}} - {\rm m}_{\chi} > 10\,{\rm GeV} \ {\rm and} \ {\rm tan} \ \beta < 40 \\ m_{\tilde{\tau}} &> \ 81.9\,{\rm GeV} \ {\rm if} \ {\rm m}_{\tilde{\tau}_{\rm R}} - {\rm m}_{\chi} > 15\,{\rm GeV}, \\ m_{\tilde{b}_{1}} &> \ 89\,{\rm GeV} \ {\rm if} \ {\rm m}_{\tilde{b}_{1}} - {\rm m}_{\chi} > 8\,{\rm GeV}, \\ m_{\tilde{t}_{1}} &> \ 95.7\,{\rm GeV} \ {\rm if} \ {\rm m}_{\tilde{t}_{1}} - {\rm m}_{\chi} > 10\,{\rm GeV}. \end{array}$$

Direct Detection of DM in MSSM



Figure 6: p9MSSM points that are allowed at 2σ by the **basic** constraints on the $(m_{\chi}, \sigma_p^{\text{SI}})$ plane. The points consistent at 2σ with the **basic** and XENON100 constraints are divided by the composition of the neutralino: gaugino-like (green squares), mixed (blue circles), or higgsino-like (red stars). Points excluded at the 95% C.L. by **basic**+XENON100 are shown as gray crosses. (a) $\Sigma_{\pi N} \simeq 43 \pm 12 \text{ MeV}$, (b) $\Sigma_{\pi N} \simeq 66 \pm 6 \text{ MeV}$.

Biggest uncertainty in DD c.s.

MSSM: signal likely but not guaranteed

L. Roszkowski, KITP, UCSB, 9 July 2013



(g-2) anomaly:

- Inconsistent with SUSY with slepton-squark unfication
- Implies:
 - m_chi~<500 GeV
 - smuon, sneutrino~<600 GeV</p>

(2sigma, p9MSSM)



Window of hope for LHC14

... a question on many people's mind...

But what about fine-tuning/naturalness?!

- I prefer to follow what the data implies, rather than theoretical prejudice
- Stabilizing mass hierarchy: initial motivation for SUSY but why should we treat it as a sacred cow
- Naturalness: fundamental Higgs -> SUSY
- 126 GeV -> generically 1TeV <~ M_SUSY tens of TeV
- Initial motivation for cosmic inflation was to rid the Universe of unwanted relics like monopoles. Now: primordial density perturbation
- Fine-tuning is needed at any scale above the EW scale!

1 TeV is not a magic number

- If SUSY is discovered, large FT issue will have to be understood/accepted
- If SUSY is not discovered, the issue will become irrelevant
- Naturalness argument gone astray:

$$rac{m_t}{m_b} \sim rac{m_c}{m_s} \simeq 14 \; \Rightarrow \; m_t \simeq 60 \, {
m GeV}$$



To take home:

Even the simplest constrained SUSY model CMSSM is consistent with all experimental constraints.

except (q-2) muon, R(gamma gamma)

- (Other simple constrained SUSY models: more freedom or similar story.)
- > Higgs of 126 GeV --> typically M SUSY at multi-TeV scale.

Plus a window of light stop 1 (~1TeV) – best fit region (stau coann.)

- > 1-tonne DM detectors to probe most CMSSM parameters. Far beyond direct LHC reach. Big bite by LUX in 2014. Other simple constrained SUSY models: similar story. MSSM: wide ranges within/outside 1-tonne detector's reach.
- TTeV (higgsino) LSP DM generic prediction of constrained SUSY models (and also MSSM – but inconsistent with g-2!)
- MSSM: (g-2)_muon: some EWinos within ~500 GeV



CNMSSM+NUHM: h1+a1 degeneracy: simultaneous enhancement of Higgs to 2-photon, tau-tau and bbar, but not ZZ, WW signal in (SM-subdominant)³ bbbar H mode

The real message:

SUSY may be too heavy for the LHC

Dark matter searches are likely to come to the rescue

BACKUP

High FT: problem or a hint?

m_h~126 GeV -> M_SUSY ~> 1 TeV -> high FT is basically ``an experimental fact"
...despite various ``islands" of smaller M_SUSY here and there

► EWSB:
$$\mu^2 = -\frac{1}{2}M_Z^2 + \frac{m_{H_d}^2(M_{\text{SUSY}}) - \tan^2\beta m_{H_u}^2(M_{\text{SUSY}})}{\tan^2\beta - 1} \qquad m_{H_u,d}^2: \text{ tree} + 1\text{L corrs}$$

→ FT argument: $m_{H_u}^2$, $m_{H_u}^2$ and μ^2 need to be all fine-tuned to give M_Z^2

- Standard approach: look for ways to reduce it
- > Another approach: accept it as an anthropic ``accident'' (Ibanez)
- Our way: Do the regions favored by m_h~126 GeV and DM density map out certain relations at the GUT scale?

Is nature telling us something?