

Quantum integrability and gauge theory

Nikita Nekrasov

w/ S. Shatashvili

Vacua of supersymmetric gauge theories
and
Quantum integrable systems

two dimensional $\mathcal{N}=(2,2)$ super-Poincaré invariance
4 supercharges



(twisted) chiral ring \rightarrow commutative associative algebra
for which the space \mathcal{M} is a representation



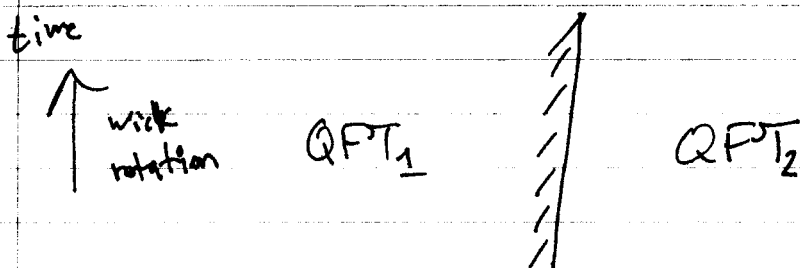
Quantum integrals of motion
of some QIS

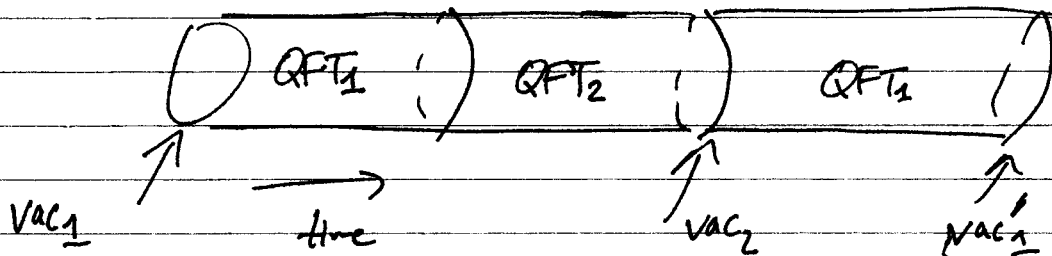


Infinite dimensional symmetry algebra
Spectrum-generating
Yangian or $U_q(\widehat{\mathfrak{sl}}_n)$

- get symmetries acting
- (i) between vacua of a given QFT
 - (ii) map different QFTs to each other

Symmetries realized by domain walls





One ingredient of $N=(2,2)$ supersymmetric gauge theory

twisted superpotential $\tilde{W}(\sigma)$
 σ - \mathbb{C} -scalar in the vector multiplet

vacua are determined by $\frac{\partial \tilde{W}}{\partial \sigma_i} = \sum_j c_{ij} n_j$ $n_j \in \mathbb{Z}$
 $j=1, \dots, r$

at low energy $\sigma \in t_{\mathbb{C}}$

For chiral fields

$$W(x) \rightarrow V = \left\| \frac{\partial W}{\partial x} \right\|^2$$

$$\mathcal{L}^{\text{bos}} = \|F\|^2 + \text{Im} \frac{\partial \tilde{W}}{\partial \sigma_i} F_i + \sum_i \left| \text{Re} \frac{\partial \tilde{W}}{\partial \sigma_i} \right|^2$$

\nearrow curvature of abelian gauge field

$$\frac{1}{2\pi i} \int F_j = n_j \quad \text{flux quantized}$$

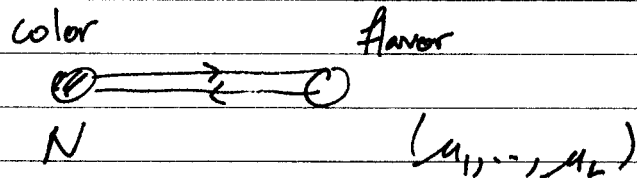
$$e^{i n_j \int F_j} \quad \text{Lagrange multiplier, resum}$$

97' Loseu + SS + NN

2d

$$t = \frac{\theta}{2\pi} + i\tau \quad \uparrow \text{FI term}$$

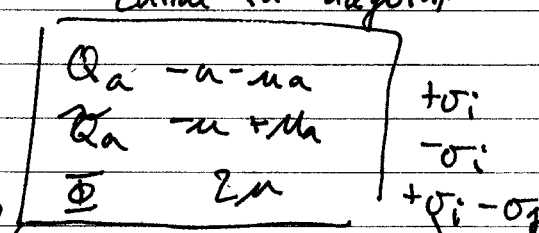
Ex $U(N)$ gauge theory
 $\mathcal{N} = (4,4)$ with L hypermultiplets in fundamental rep



+ twisted masses, consistent w/ $\mathcal{N} = (2,2)$ SUSY
 global sym

$$\mathcal{N} = 4 \text{ vector} = (\sigma, \mu_a) = \begin{matrix} \Phi \\ \uparrow \\ \Phi \end{matrix} = \begin{matrix} \uparrow \\ \text{chiral in adjoint} \end{matrix} \quad \mathcal{N} = 2 \text{ vector}$$

$$W = \overline{Q} \begin{matrix} -1 & +2 & -1 \\ \hline \end{matrix} Q \Phi Q$$



\Rightarrow additional twisted mass u (corresponding to sym)

When (u_1, \dots, u_L, u) are generic

all matter fields are massive \Rightarrow can be integrated out

$$\widetilde{W}^{\text{eff}}(\sigma_1, \dots, \sigma_N; u, u) = \widetilde{w}(\sigma_i)$$

$$= \sum_{i, j} \widetilde{w}(\sigma_i - \sigma_j + 2u)$$

$$+ \sum_i \widetilde{w}(\sigma_i - \mu_a - u) + \widetilde{w}(\mu_a - \sigma_i - u)$$

$$+ \sum_i \sigma_i$$

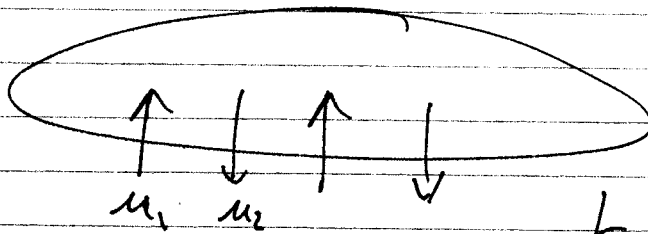
$$\widetilde{w}(x) = x(\log x - 1)$$

Vacua

$$\prod_{a=1}^L \left(\frac{\sigma_i - \mu_a - u}{\sigma_i - \mu_a + u} \right) = e^{2\pi i t} \prod_{t \neq i} \left(\frac{\sigma_i - \sigma_j - 2u}{\sigma_i - \sigma_j + 2u} \right)$$

$$\bar{\omega}_\lambda(x) = x \left(\log \left(\frac{x}{\hbar} \right) - 1 \right)$$

Bethe equations for $XXX_{1/2}$ spin chain
inhomogeneous, twisted



L spins on a circle

local magnetic field $e^{2\pi i t} \sigma_3$
in a sector w/ N spin \uparrow

$$\mathcal{H}_{XXX} = (\mathbb{C}^2)^{\otimes L} = \bigoplus_{N=0}^L \mathcal{H}_N$$

$$\dim \mathcal{H}_N = \binom{L}{N}$$

Susy vacua of $U(N)$ $N=(1,1) + L$ flavors

$$\mathcal{H}_N$$

$\mathcal{Y}(sl_2)$ for which \mathcal{H}_{XXX} is reducible (for generic μ 's)

$$\mathcal{H}_{XXX} = V_{1/2} \left(\frac{\mu_1}{2u} \right) \otimes \dots \otimes V_{1/2} \left(\frac{\mu_L}{2u} \right)$$

μ_i : generic irreducible rep of Yangian

if $\mu_i - \mu_j = 2a$ \mathcal{H}_{xxx} is not irreducible

level crossing between Q and \tilde{Q}

eg $\mu_a - \mu_{a+1}$ meson $O_a = \tilde{Q}_{a+1} Q_a$ becomes massless

Can deform in two ways

$$\langle O_a \rangle \neq 0$$

VEV

$$W \rightarrow W + m O_a$$

deform by mass term

\rightarrow generate couplings $\tilde{Q} \Phi^2 Q \rightarrow$ spin 1 impurities
in the spin chain

$$W^{\text{gen}} = \sum_a \tilde{Q}_a \Phi^{2s_a} Q_a$$

s_a half-integer

$$-2s_a + 2 - 2s_a$$

$$\tilde{Q}_{a+1} \Phi Q_{a+1} + \tilde{Q}_a \Phi Q_a + m \tilde{Q}_{a+1} Q_a$$

integrate out \tilde{Q}_{a+1}

$$\Rightarrow Q_a = m^{-1} \Phi Q_{a+1}$$