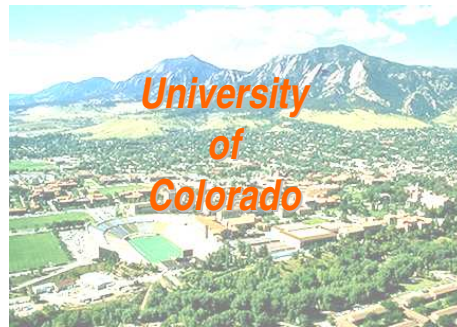



Cosmological Radiative Transfer

Nick Gnedin



Outline


- Reionization inside a computer
- SF & RT: two stumbling blocks
- Is OTVET the answer?
- Adaptive Mesh Refinement vs Ray Tracing
- A call for help.



The ^{very} Brief History of Time


The universe began **HOT** - ended **COLD**

- $T > 4000\text{K}$: hot and photo-*ionized*
- $T \sim 4000\text{K}$: recombination
- $T < 4000\text{K}$: neutral and *cooling down*...
- Today: hot ($\sim 8000\text{K}$) and *ionized*!




Something must have happened in between!

RE-IONIZATION



Imagine If You Were a Hydrogen Atom...



Reionization Inside a Computer

Easy physics:
fast, accurate

Hard physics:

Extra physics:

- dark matter dynamics
- gas dynamics
- ionization/recombination/heating

- star formation ← *phenomenology*
- radiative transfer ← *hard numerics*

- magnetic fields
- dust
- cosmic rays

Perhaps, not needed ?

Star Formation

Schmidt Law:

$$\frac{d\rho_*}{dt} = \epsilon_{SF} \frac{\rho_{gas}}{\tau_{dyn}}$$

• Must be scale-dependent

Although, observed in large and small galaxies



Radiative Transfer

$$\frac{\partial J_\nu}{\partial t} + c\vec{n} \frac{\partial J_\nu}{\partial \vec{x}} = -k_\nu J_\nu + S_\nu$$

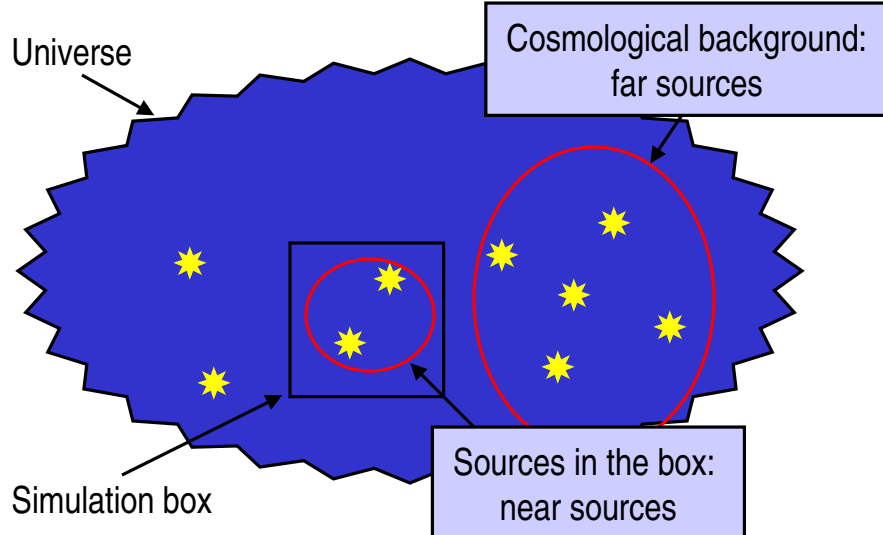
Just another Boltzmann equation, but...

$$\frac{\partial f}{\partial t} + \vec{v} \frac{\partial f}{\partial \vec{x}} - \nabla \phi \frac{\partial f}{\partial \vec{v}} = 0$$

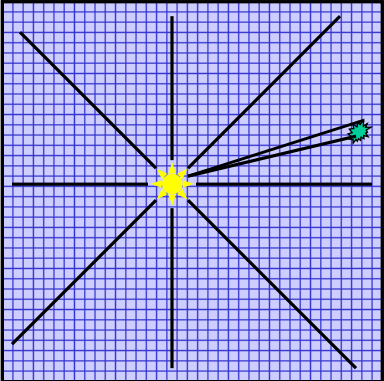
Yet another Boltzmann equation.



Sources: Near vs Far



Ray Tracing



If we want to preserve the resolution of gas solver:
 $N(\text{rays}) \sim N(\text{surface cells})$

↓

Far sources: $\sim M^{5/3} N_v$
 Near sources: $\sim N \times N_s \times N_v$


No way!

Moments of Boltzmann Equation

<p>Massive particles:</p> <ul style="list-style-type: none"> • Density • Momentum • Pressure tensor • ... 	<p>Photons:</p> <ul style="list-style-type: none"> • Energy density • Flux • Radiation Pressure tensor • ...
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The hierarchy needs to be closed.

But: photons move with the speed of light!




Eddington Tensor

$$P_R^{ij} = \frac{4\pi}{c} \langle J_\nu n^i n^j \rangle = E_\nu h^{ij}$$

Eddington tensor \rightarrow h^{ij}

$$\text{Trace}(h^{ij}) = 1$$

5 independent components, not 6!



OTVET Approximation

Simple idea:

- use the Eddington tensor from the optically thin regime.

Optically thin regime:

- collect $1/r^2$ contribution from all sources – just like gravity!

Result:

- **O**ptically **T**hin **V**ariable **E**ddington **T**ensor approximation



Scaling

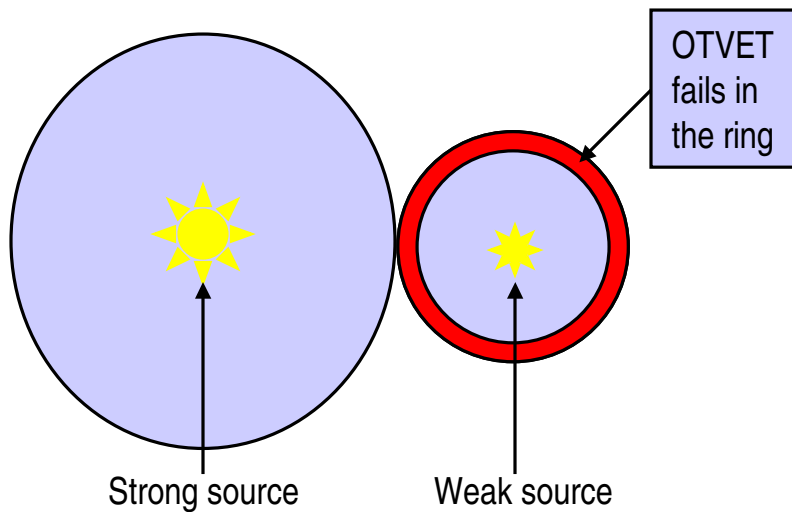
- Moments equations (like hydro) $\sim N \times N_v$
- Eddington tensor (like gravity) $\sim N \times \log(N)$


Advantages:

- **fast**
- controlled: the error of the approximation can be measured posteriori
- no scaling with number of sources: works for any N_S
- energy density and flux are conserved **exactly**



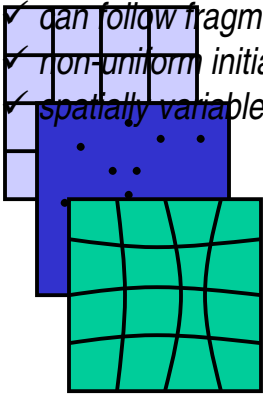
Where It Fails



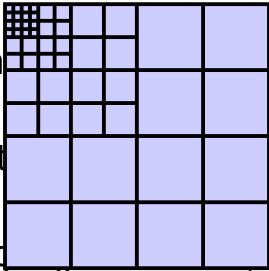


Do We Have a Choice?


XX CenturyXXI Century



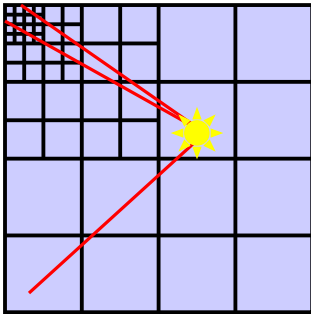
✓ can follow fragmentation
✓ non-uniform initial conditions
✓ spatially variable resolution



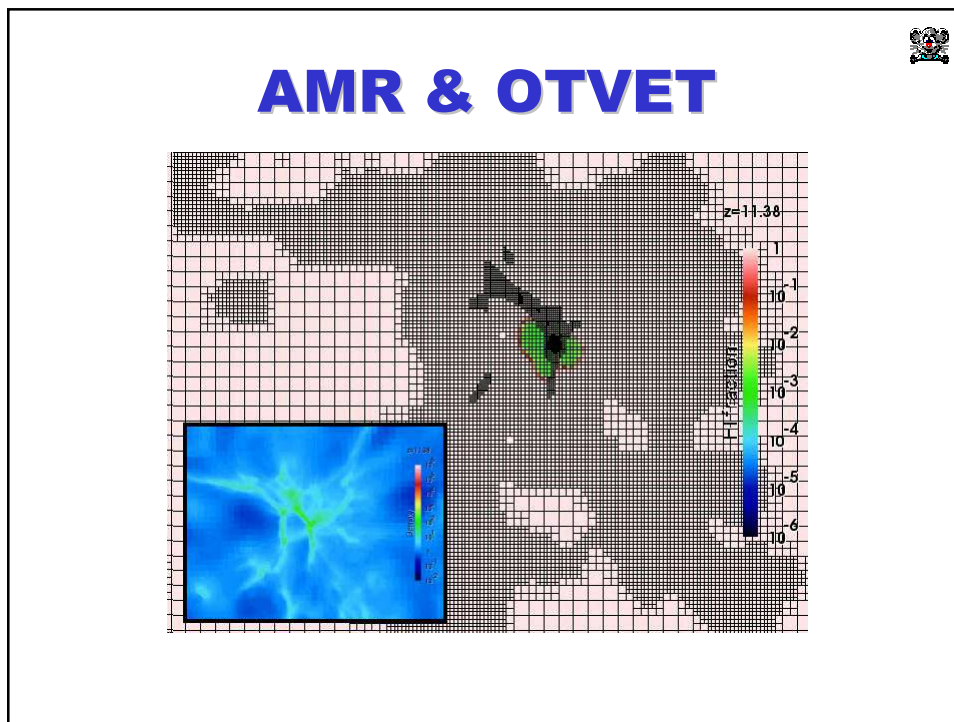
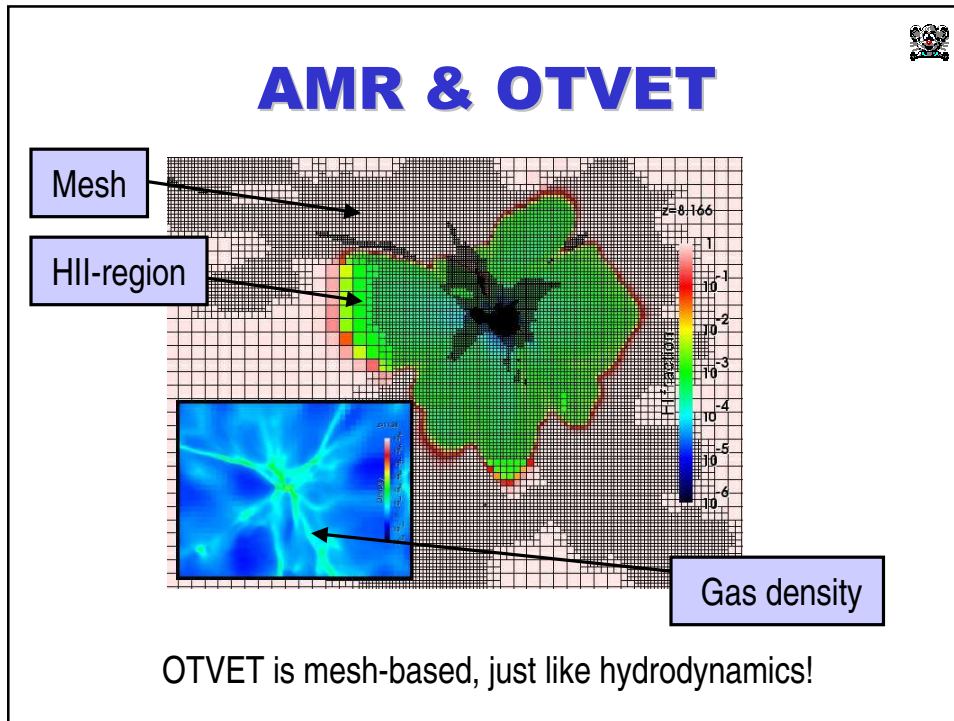
- Eulerian schemes
- Smooth Particle Hydrodynamics (SPH)
- Arbitrary Lagrangian Eulerian (ALE)
- Adaptive Mesh Refinement (AMR)




AMR And Ray Tracing




Ray Tracing on the AMR mesh would require unrealistically many rays – as if the whole mesh was refined.





RT vs N-body

Radiative Transfer		N-body
Ray tracing	↔	Direct summation
OTVET	↔	Hydrodynamics
?	↔	PM
???	↔	P ³ M, Tree, AMR

- 
- ### Conclusions
- Reionization is a complex process: its complete understanding will crucially depend on our ability to simulate radiative transfer with a fast and accurate solver.
 - Large computational boxes will be required, because HII regions are biased and non-trivially shaped.
 - As AMR takes over, we are left with few options for doing RT on adaptive meshes.
 - Graduate students: ***YOU ARE THE ONLY HOPE!***

