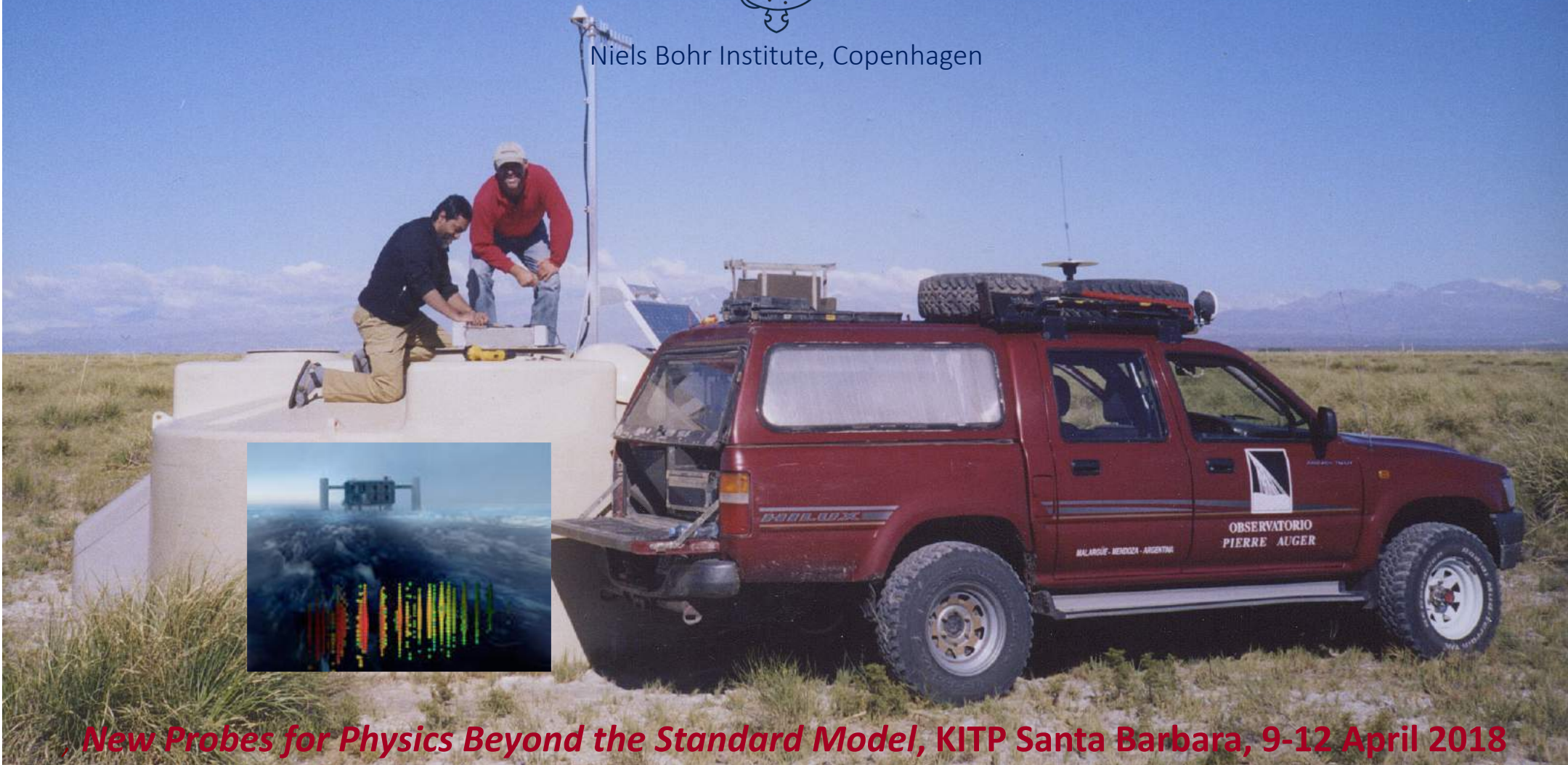


Probing new physics with ultra-high energy cosmic neutrinos

Subir Sarkar

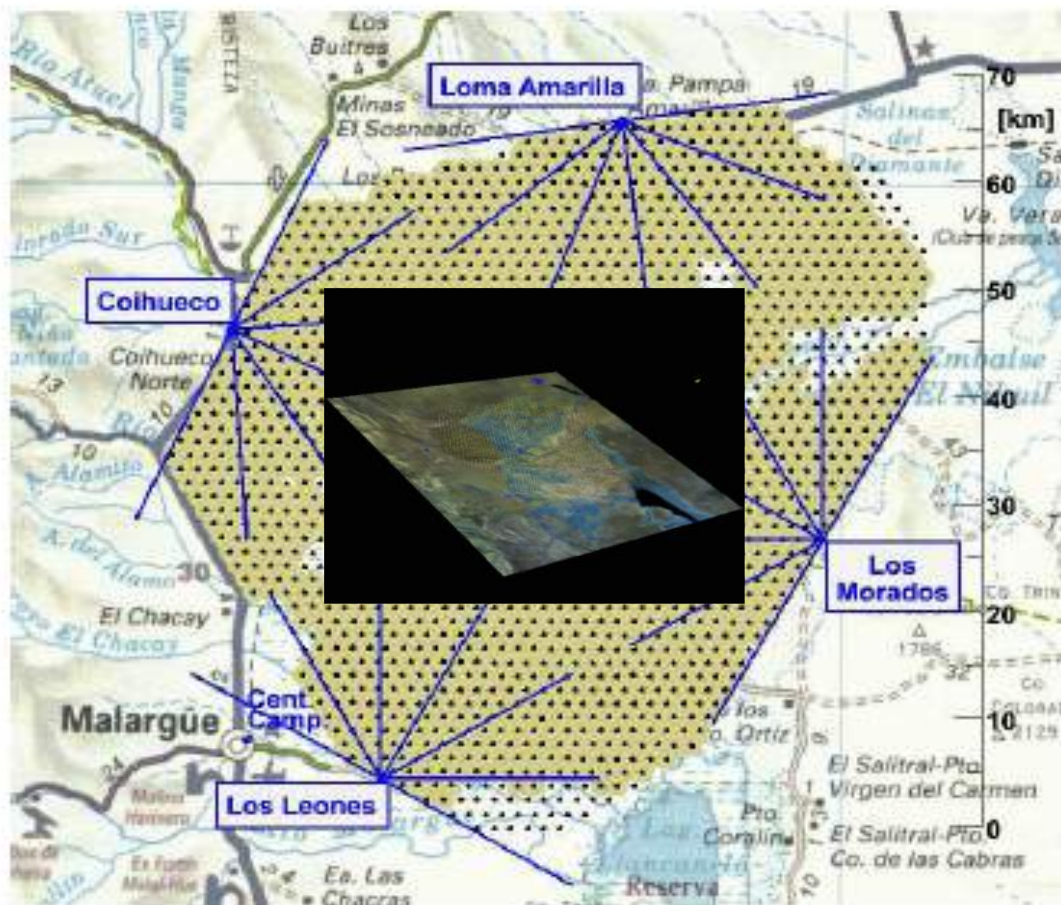


Niels Bohr Institute, Copenhagen

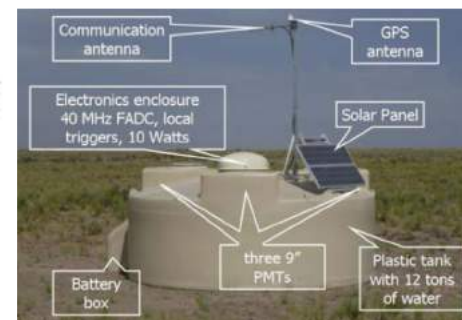


New Probes for Physics Beyond the Standard Model, KITP Santa Barbara, 9-12 April 2018

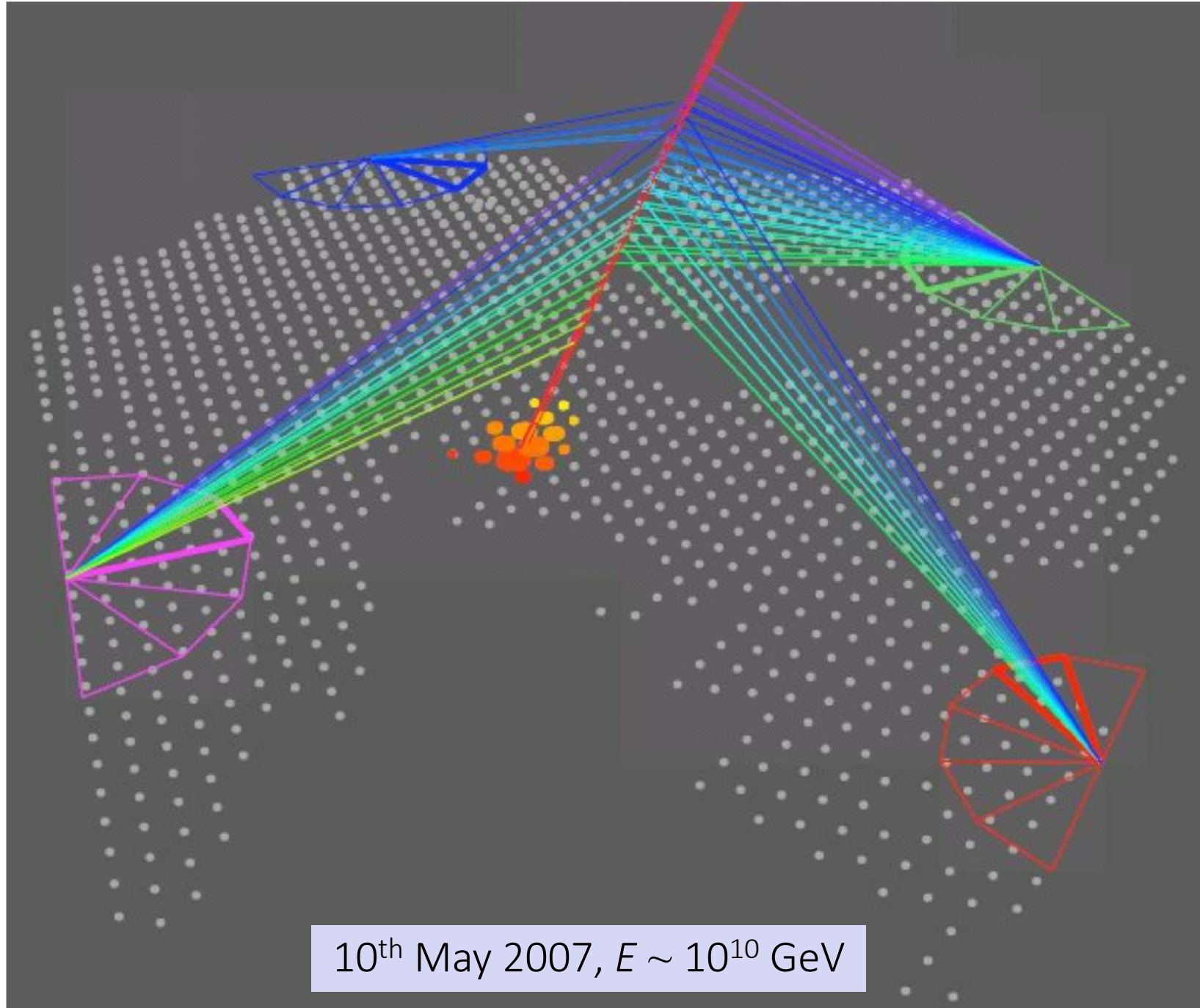
The Pierre Auger Observatory



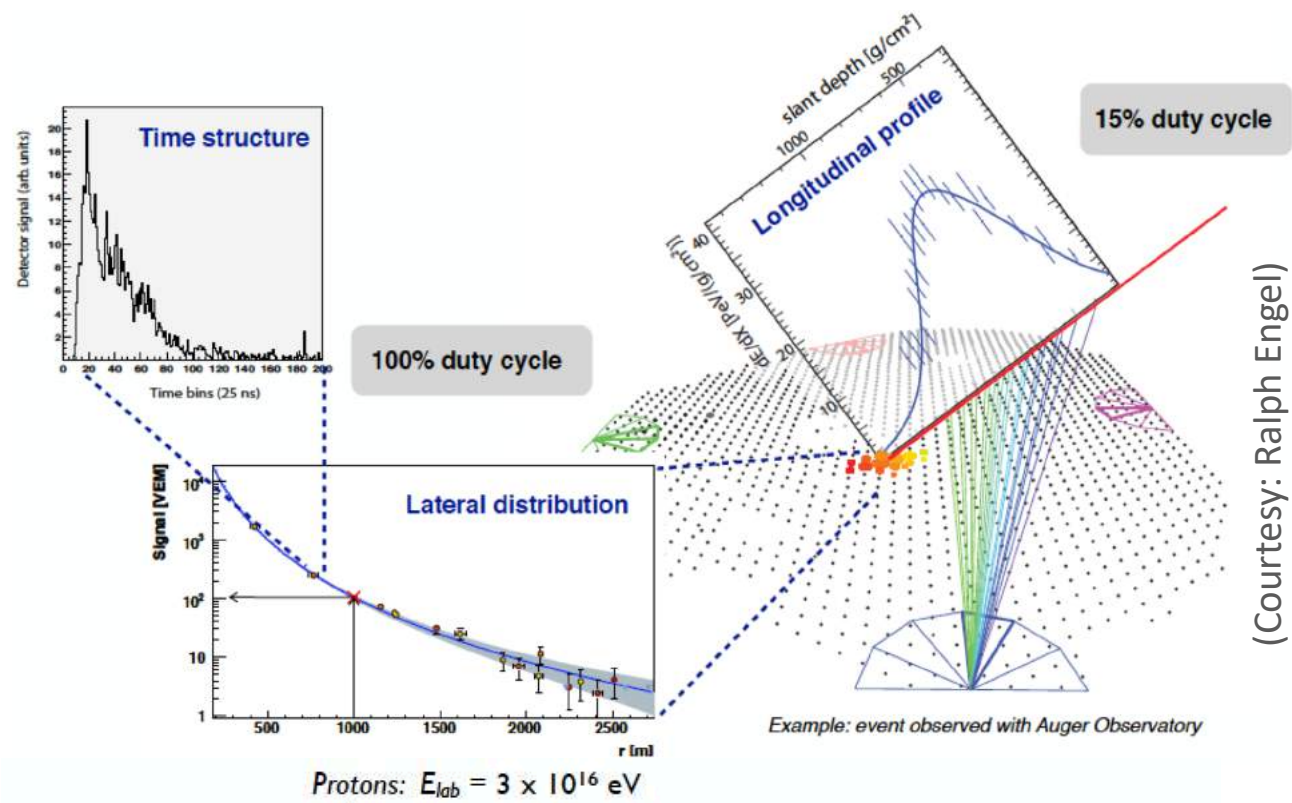
- 1600 water-cherenkov detectors
- Aperture $> 7000 \text{ km}^2 \text{ sr yr}$
- 4×6 telescopes



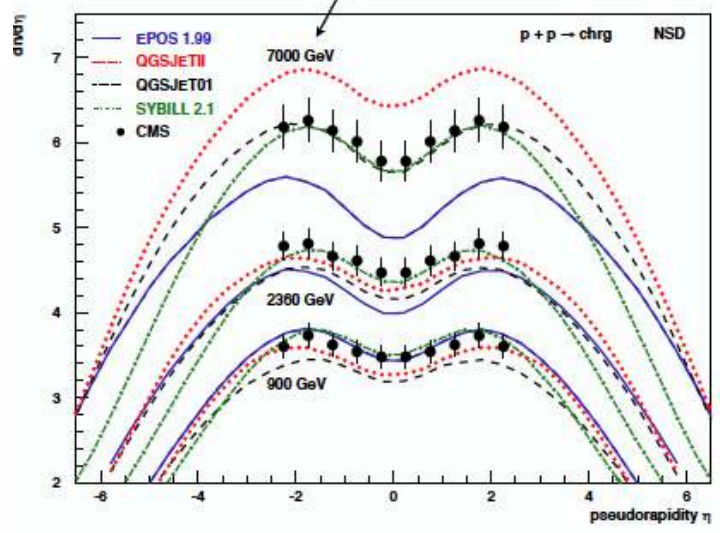
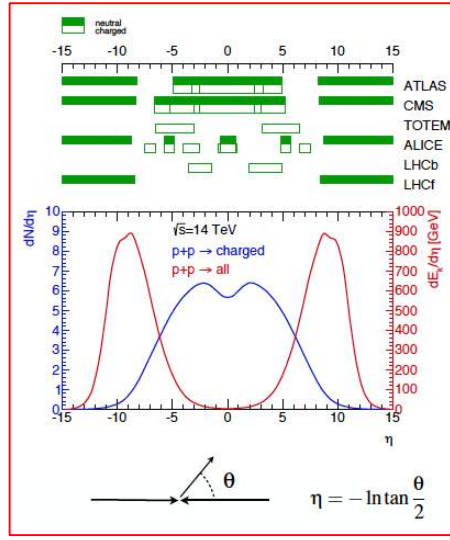
With this detector we see the *highest* energy particles in the universe



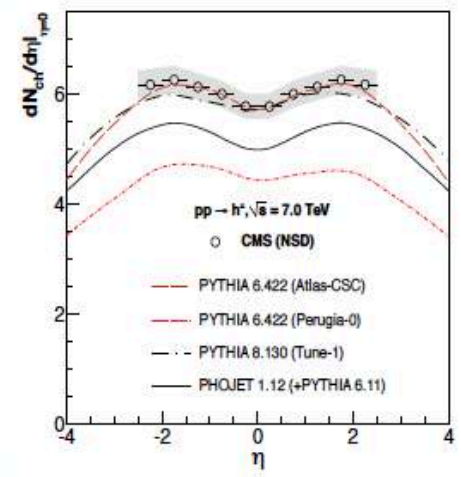
The energy estimate from the lateral distribution of the signal in the **surface detector** array is calibrated *directly* by the (calorimetric) **fluorescence detector** signal ... there is *no* reliance on shower simulations



(Courtesy: Ralph Engel)

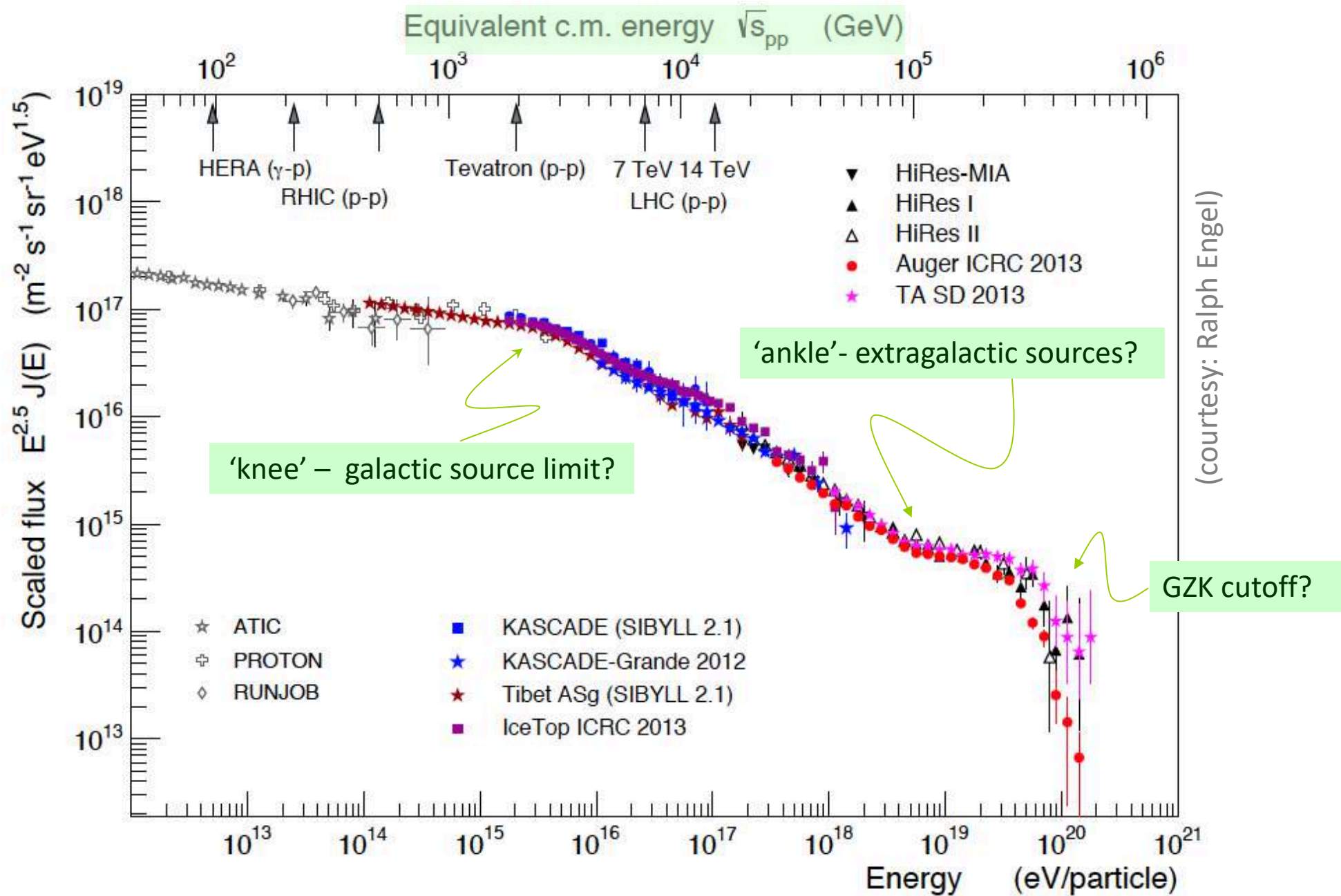


Detailed LHC comparison
(D'Enterria et al., APP 35, 2011)



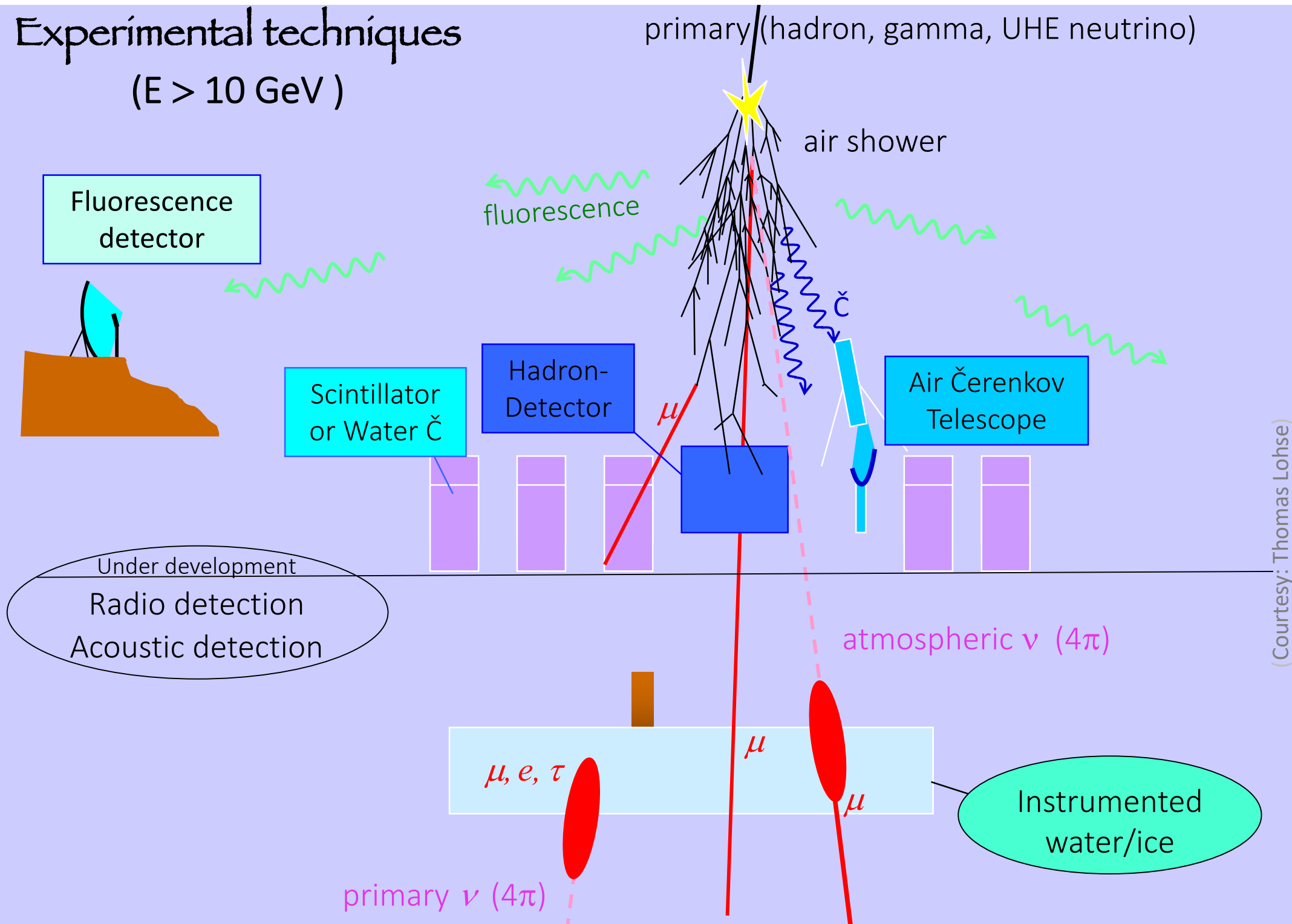
In any case the CR cascade MCs do very well at matching LHC data ... *better* than collider MCs!

By studying cosmic ray interactions we can probe energies well *beyond* the reach of terrestrial accelerators



Experimental techniques

($E > 10 \text{ GeV}$)



Colliders versus Cosmic rays

The LHC has achieved 13 TeV cms ...

But 10 EeV (10^{19} eV) cosmic ray initiating giant air shower

⇒ ~100 TeV cms (... although rate $\lesssim 1/\text{day}$ in 3000 km² array)

New physics would be hard to see in hadron-initiated showers

(BSM cross-section $< \text{TeV}^{-2}$ versus hadronic cross-section $\sim \text{GeV}^{-2}$)

... but may have a dramatic impact on *neutrino* interactions
(since the ν - N cross-section is very small to start with)

→ can probe new physics (both in and) beyond the Standard Model by studying ultra-high energy cosmic neutrinos

Where there are high energy cosmic rays,
there *must* also be neutrinos ...

GZK interactions of extragalactic UHECRs on the CMB

“guaranteed” cosmogenic neutrino flux

... reduced significantly if the primaries are *not* protons but heavy nuclei

UHECR candidate accelerators (AGN, GRBs, ...)

“Waxman-Bahcall limit” ... normalised to observed UHECR flux

... sensitive to ‘cross-over’ energy above which extragalactic component dominates

‘Top down’ sources (superheavy dark matter, topological defects)

motivated by trans-GZK energy events observed by AGASA

... such models now *ruled out* by the limit from Auger on primary photons
(QCD fragmentation in parton shower dominantly creates photons, *not* nucleons)

The sources of cosmic rays *must* also be neutrino sources

Waxman-Bahcall Bound :

- $1/E^2$ injection spectrum (Fermi shock).
- Neutrinos from photo-meson interactions in the source.
- Energy in ν 's related to energy in **CR**'s :

$$[E_\nu^2 \Phi_\nu]_{\text{WB}} \approx (3/8) \xi_Z \epsilon_\pi t_H \frac{c}{4\pi} E_{\text{CR}}^2 \frac{d\dot{N}_{\text{CR}}}{dE_{\text{CR}}}$$

Fraction of CR primary energy converted to neutrinos

From rate of UHE CR's (10^{19} - 10^{21} eV)

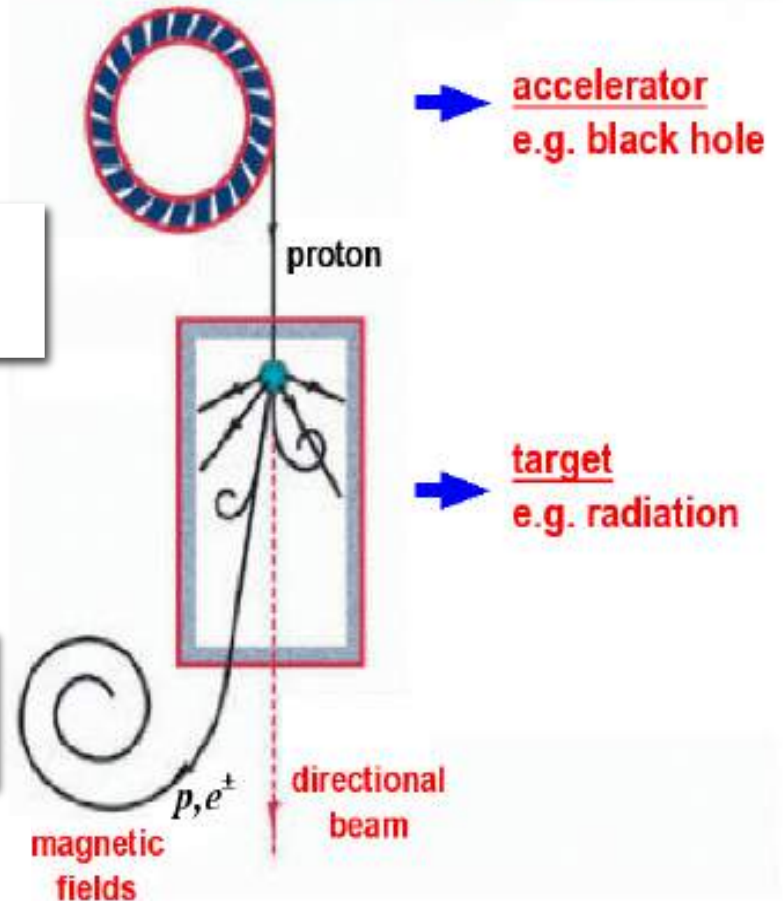
Hubble time

$$\approx 2.3 \times 10^{-8} \epsilon_\pi \xi_Z \text{ GeV cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$$

➔ Making a reasonable estimate for ϵ_π etc allows this to be converted into a **flux expectation**

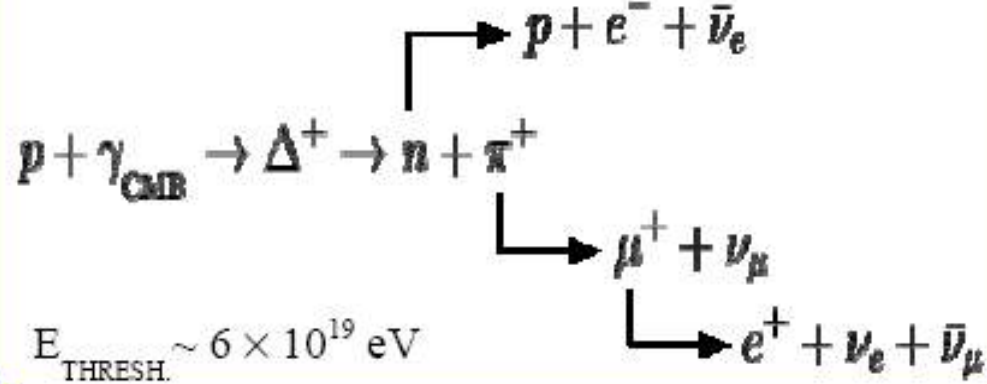
(would be *higher* if extragalactic cosmic rays become dominant at energies below the 'ankle')

COSMIC BEAM DUMP : SCHEMATIC



The “guaranteed” cosmogenic neutrino flux

GZK mechanism :

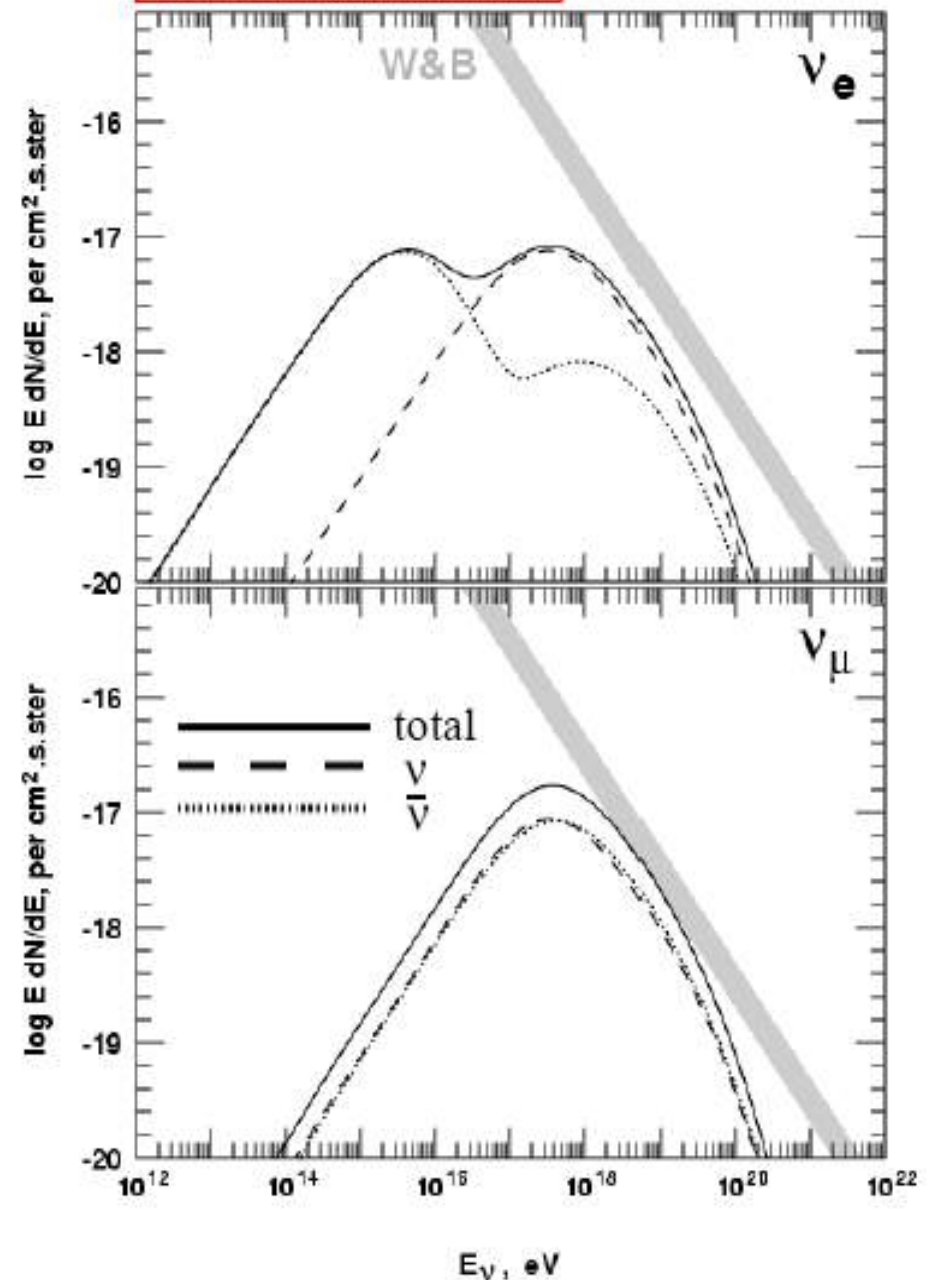


✦ Uncertainties in flux calculations :

- ▶ UHECR luminosity; $\rho_{\text{CR}}(\text{local}) \neq \langle \rho_{\text{CR}} \rangle$
- ▶ injection spectrum
- ▶ cosmological evolution of sources
- ▶ IRB & optical density of sources

factors of ~2 uncertainty each;
factor of ~4 overall (?)

Engel, Seckel, Stanev (2001)



... can pin down by normalising to the γ -ray flux from GZK process (Ahlers *et al*, *Astropart.Phys.* 34:106,2010)

We can work out the interaction rate via ν - N deep inelastic scattering (this is the dominant process above ~ 10 GeV)

$$\frac{\partial^2 \sigma_{\nu, \bar{\nu}}^{CC, NC}}{\partial x \partial y} = \frac{G_F^2 M E}{\pi} \left(\frac{M_i^2}{Q^2 + M_i^2} \right)$$

$Q^2 \uparrow \Rightarrow$ propagator \downarrow

$$\left[\frac{1 + (1 - y)^2}{2} F_2^{CC, NC}(x, Q^2) - \frac{y^2}{2} F_L^{CC, NC}(x, Q^2) \right.$$

$$\left. \pm y \left(1 - \frac{y}{2} \right) x F_3^{CC, NC}(x, Q^2) \right]$$

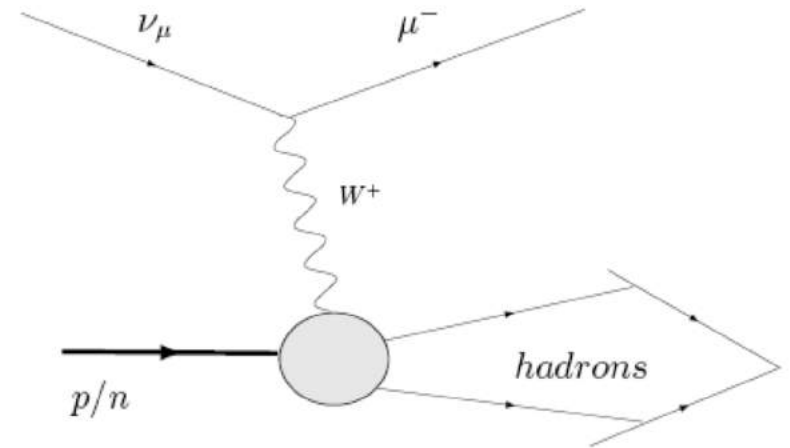
$Q^2 \uparrow \Rightarrow$ parton distribution functions \uparrow

Most of the contribution to #-secn comes from:

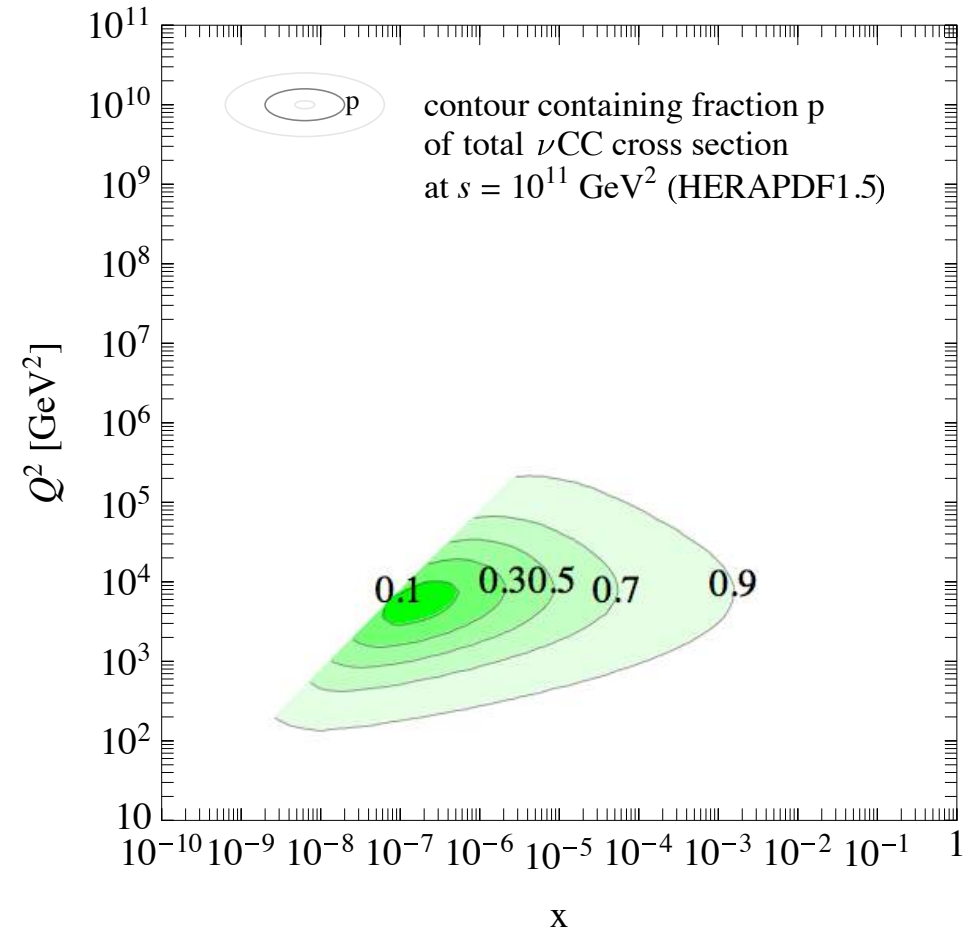
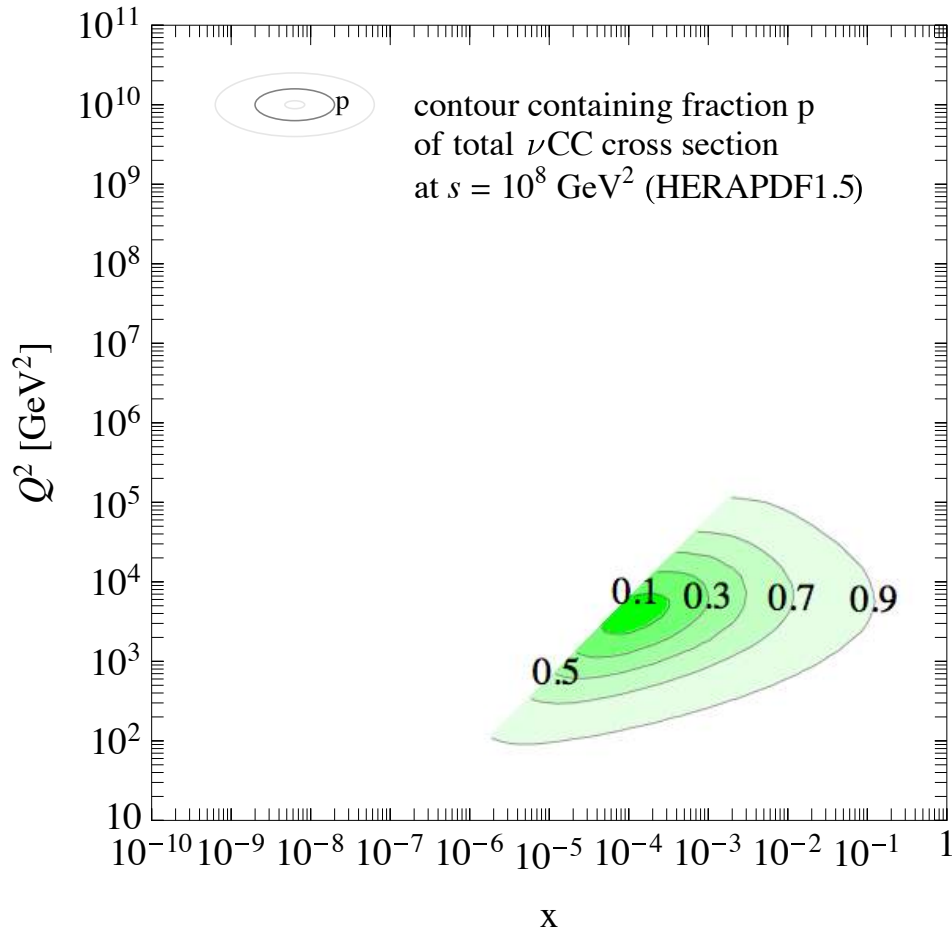
$$Q^2 \sim M_W^2 \text{ and } x \sim \frac{M_W^2}{M_N E_\nu}$$

At leading order (LO) : $F_L = 0$, $F_2 = x(u_\nu + d_\nu + 2s + 2b + \bar{u} + \bar{d} + 2\bar{c})$,
 $x F_3 = x(u_\nu + d_\nu + 2s + 2b - \bar{u} - \bar{d} - 2\bar{c}) = x(u_\nu + d_\nu + 2s + 2b - 2\bar{c})$

Can calculate numerically at Next-to-Leading-Order (NLO) ... *no* significant further change at NNLO



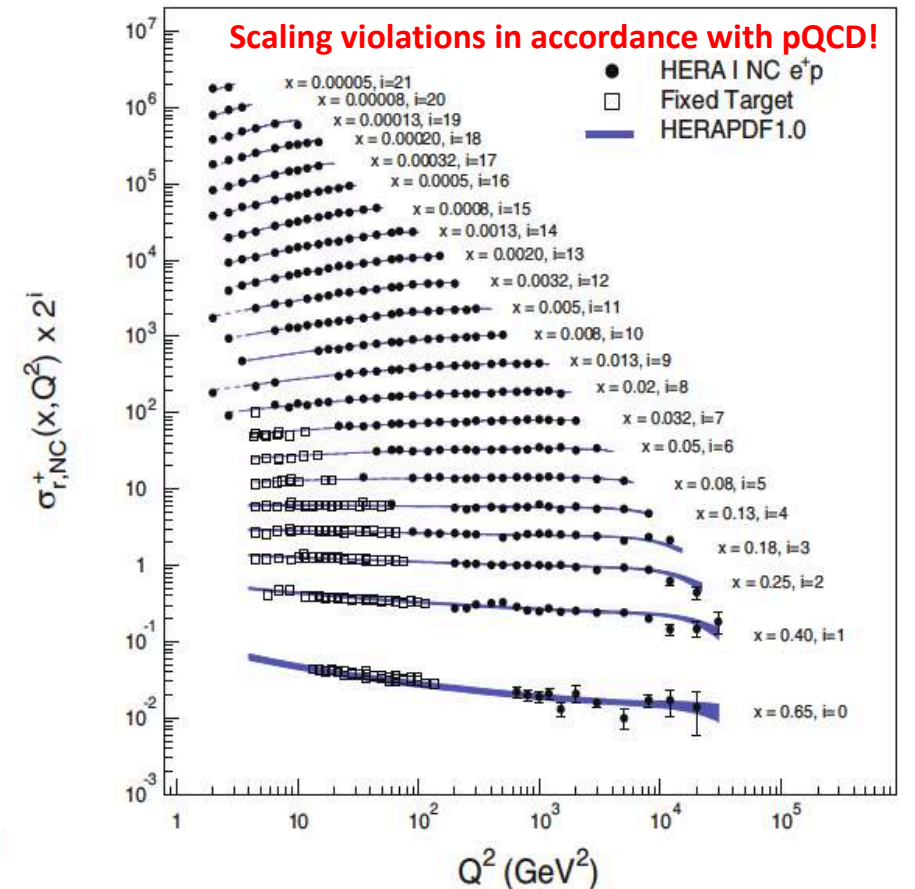
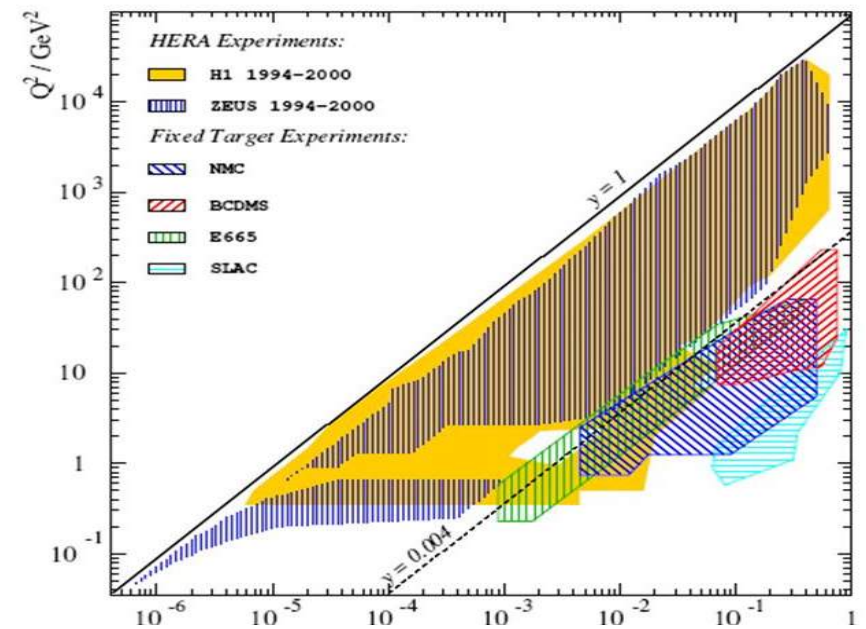
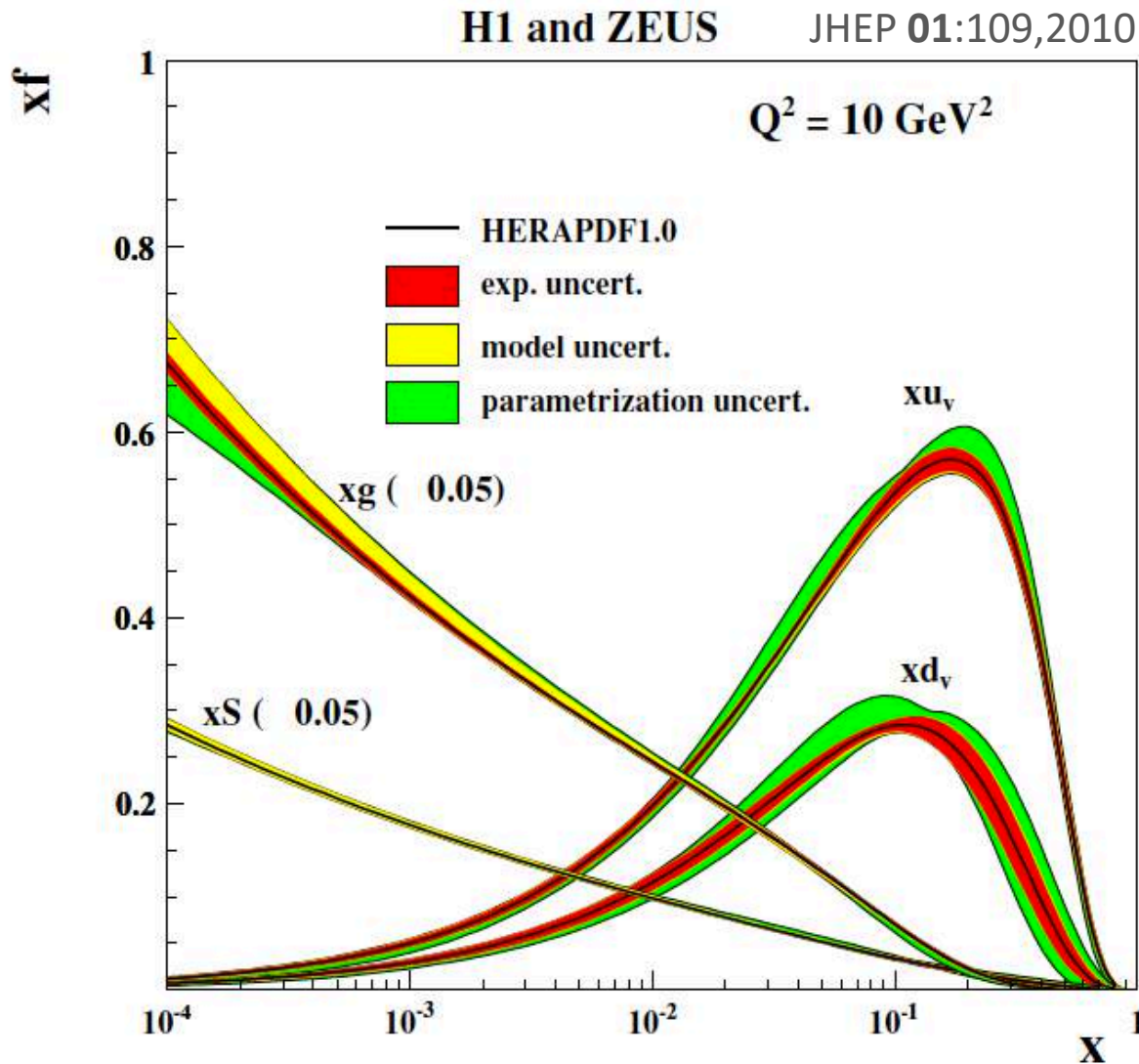
As the neutrino energy increases, *lower* values of Bjorken- x are being probed



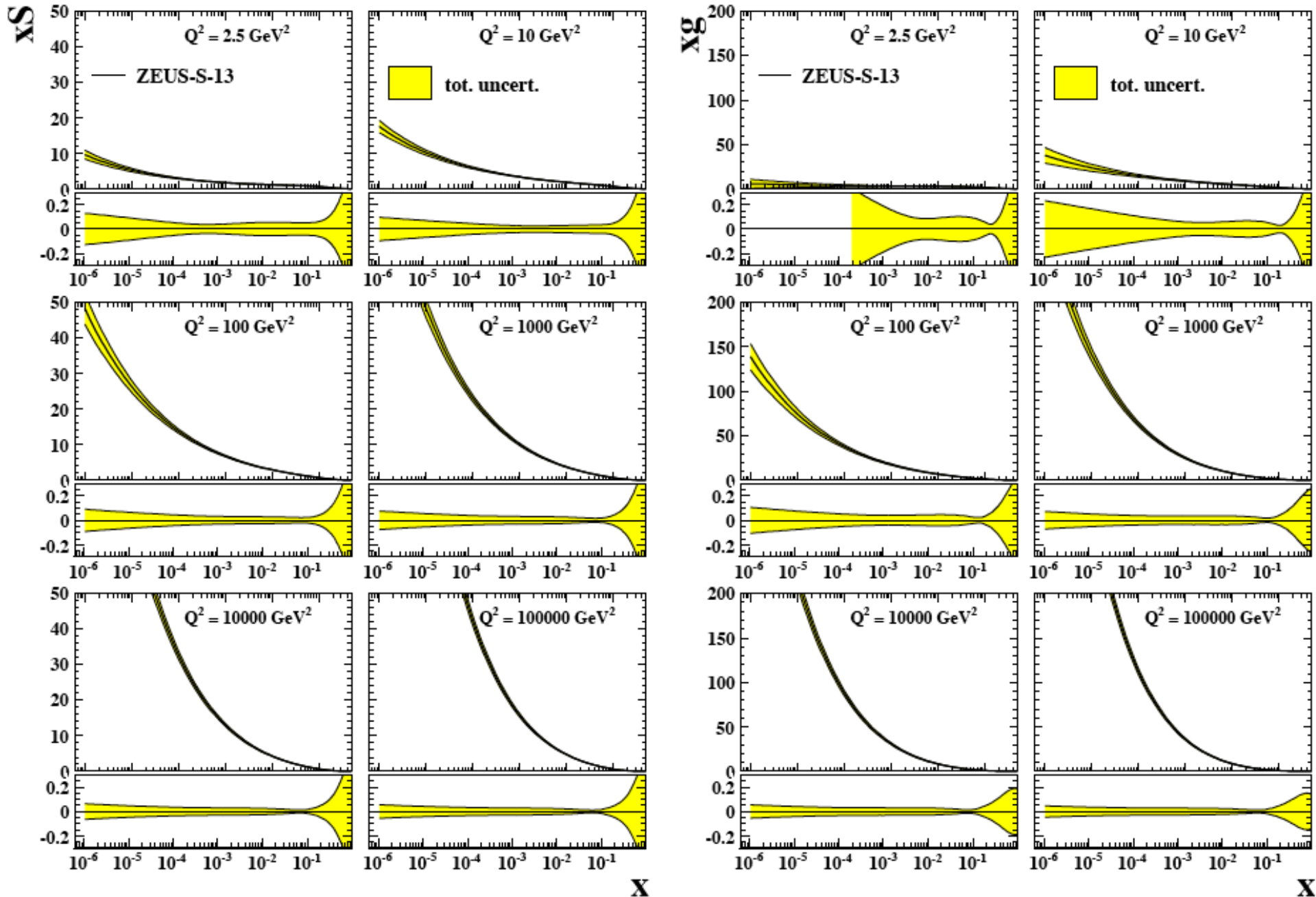
Mertsch, Cooper-Sarkar, Sarkar, JHEP 08:042, 2011

So to determine the DIS cross-section accurately it is essential to have measurements of PDFs down to as low x as is possible ... for energies $>$ PeV we need to evolve these further using the pQCD DGLAP formalism

The H1 & ZEUS experiments at HERA were the first to measure DIS at high Q^2 and low Bjorken- x ... surprising finding was the *steep* rise of the **gluon PDF** at low x



Parton distribution functions from the ZEUS-S global data analysis

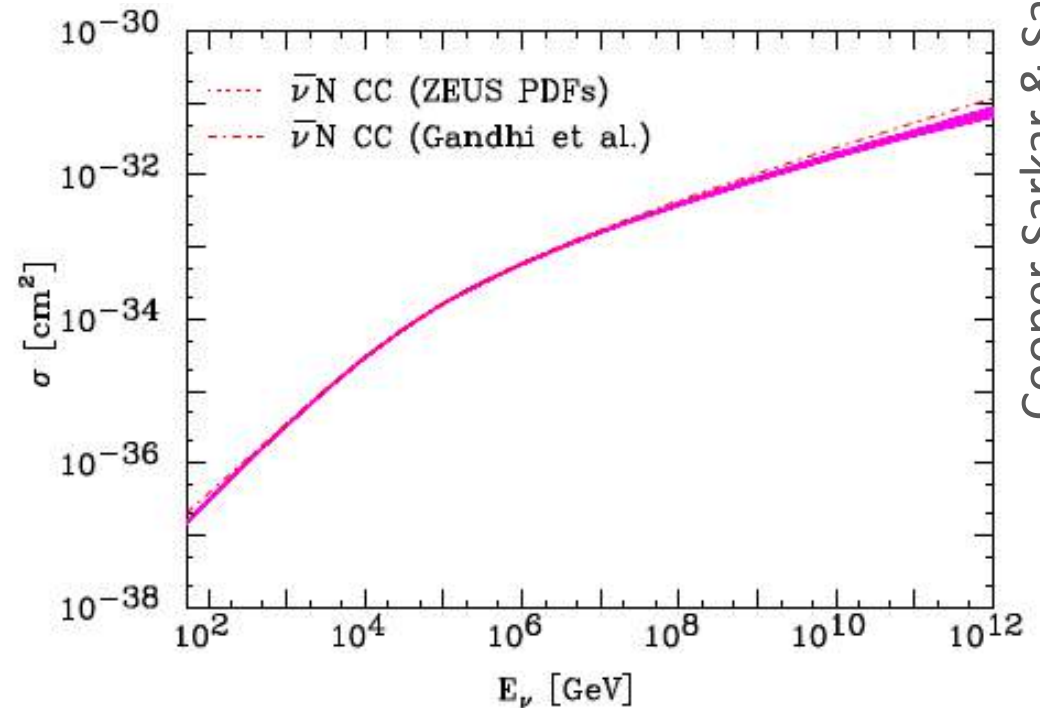
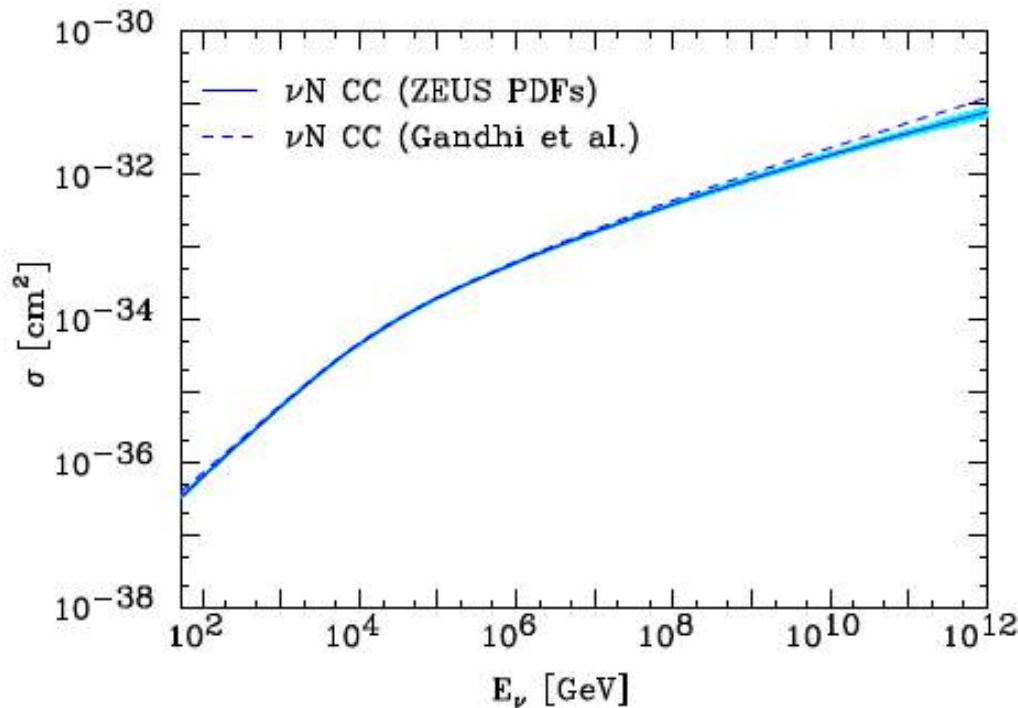
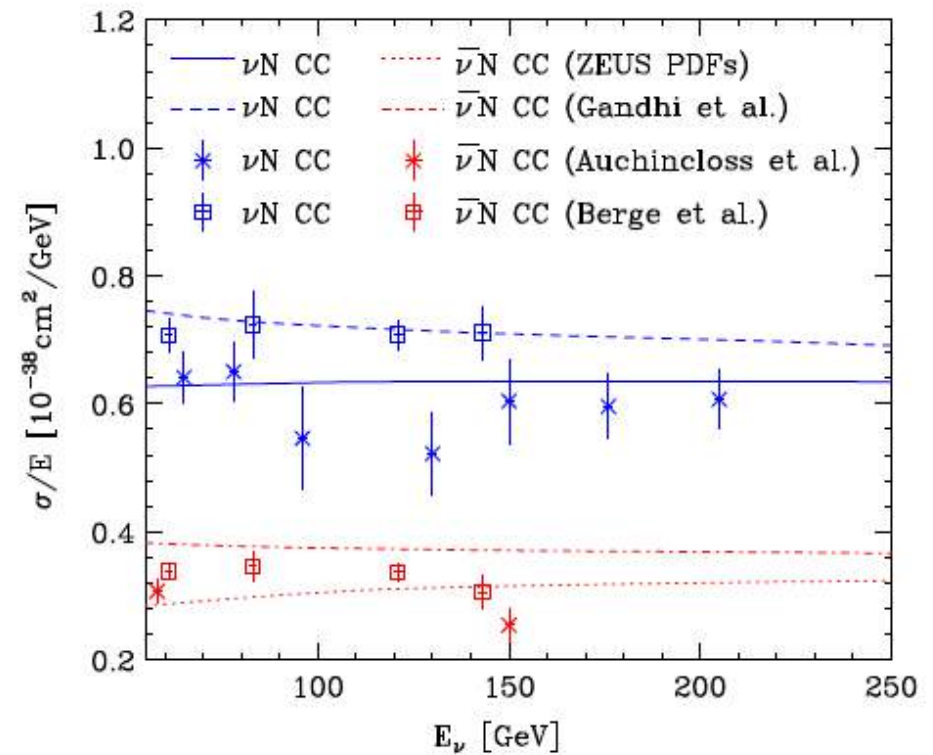


Cooper-Sarkar & Sarkar, JHEP 01:075, 2008

using DGLAP evolution of the PDFs (at NLO, including heavy quark corrections)

The $\#$ -section we found using ZEUS-S PDFs was up to $\sim 40\%$ different from the previous 'standard' calculation (based on CTEQ4) ... More importantly we could quantify the *uncertainty* in the perturbative calculation

At very high energies where very low- x is being probed, recombination/saturation effects may alter the cross-section by a factor of ~ 2 ... however DGLAP evolution appears to fit *all* exptal. data – so *no* imperative for this yet!



IceCube Neutrino Observatory

IceTop: 1 km² surface array (81 'Auger' tanks)

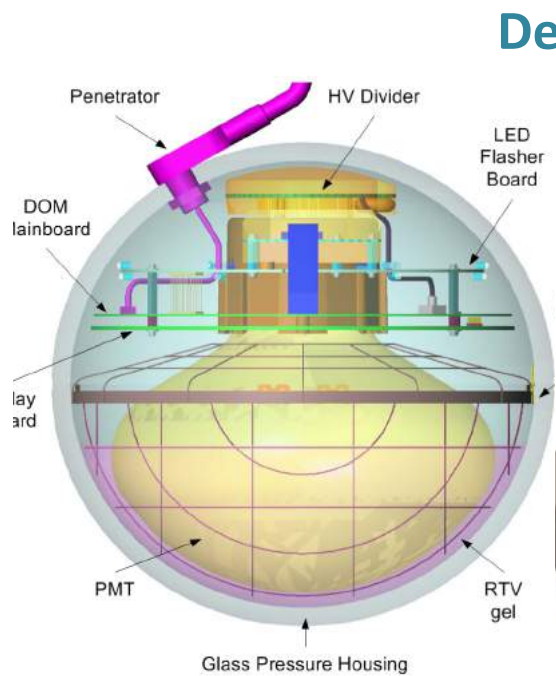
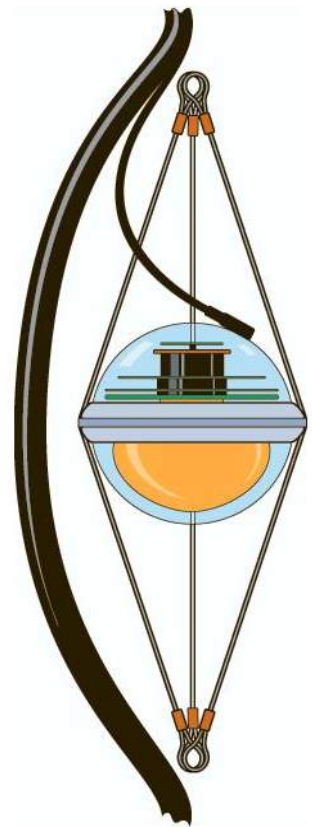
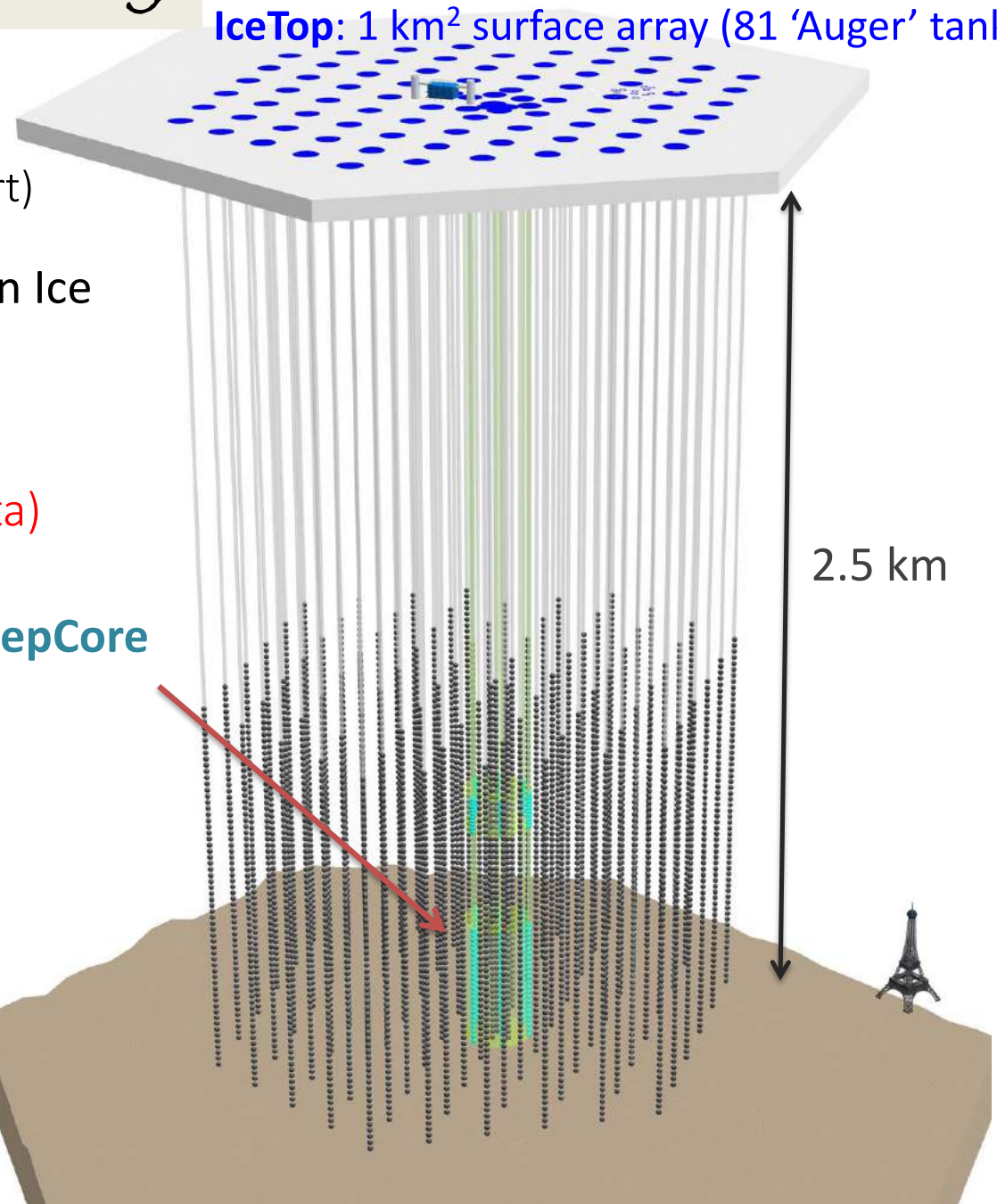
86 strings (125 m between strings)

60 Optical Modules per string (17 m apart)

5160 Digital Optical Modules (DOMs) in Ice

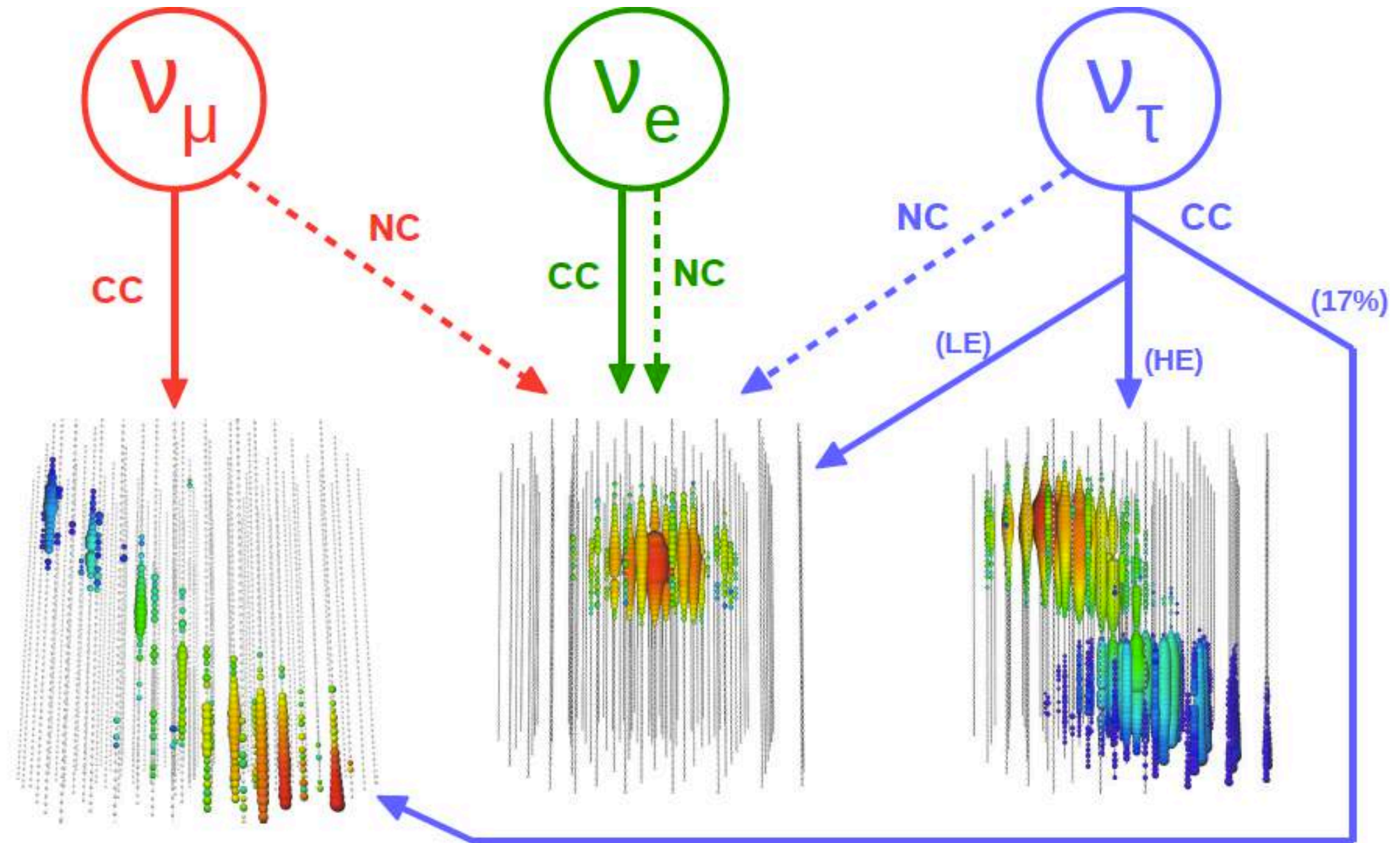
1 km³ ⇒ Gton instrumented volume

Construction: 2004-11 (now 7 yr+ of data)



Cost: 279 M\$ ⇒ <30 cents per ton

Neutrino flavour discrimination in IceCube



Track topology

Good pointing ($\sim 0.2^\circ - 1^\circ$)

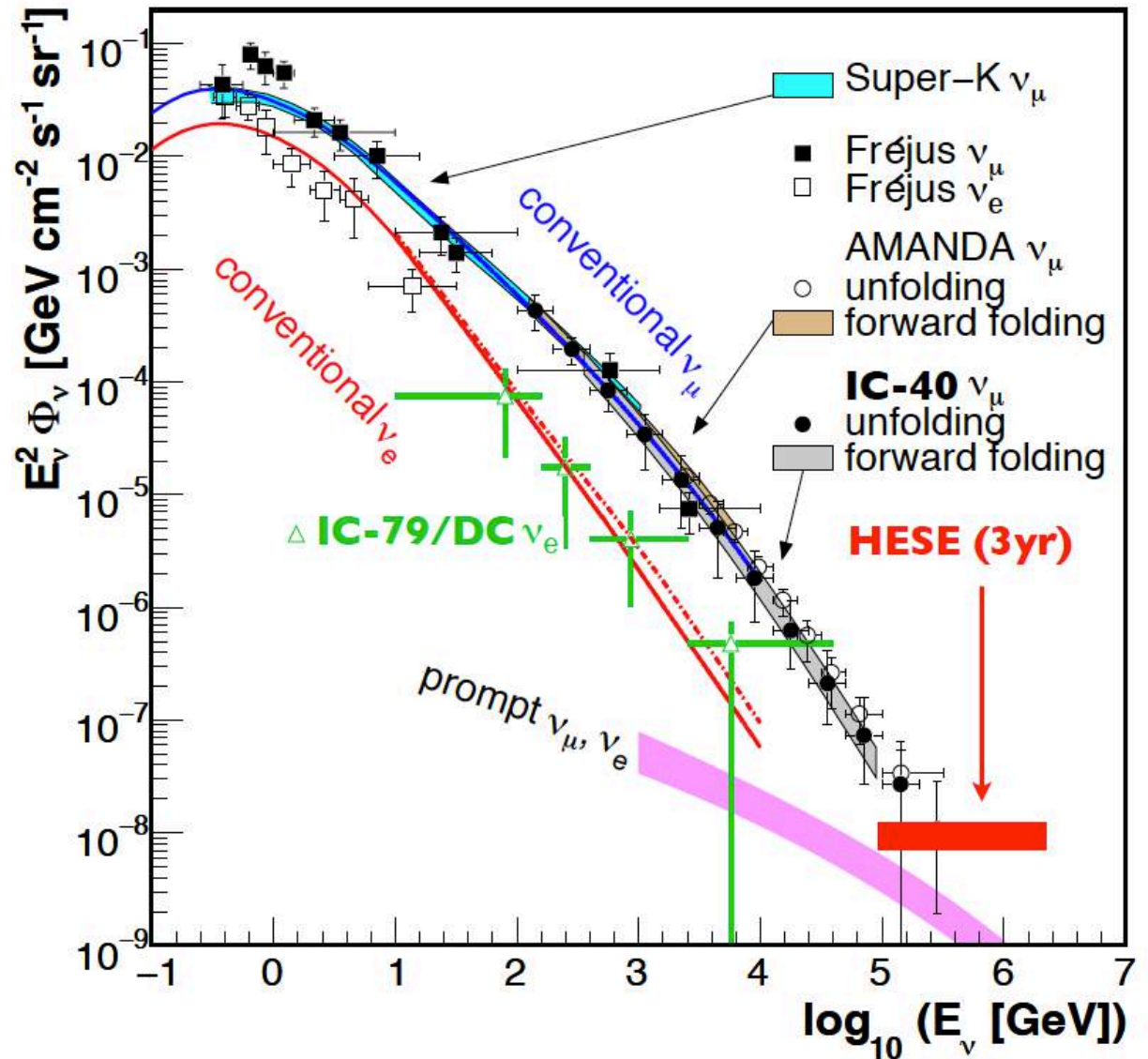
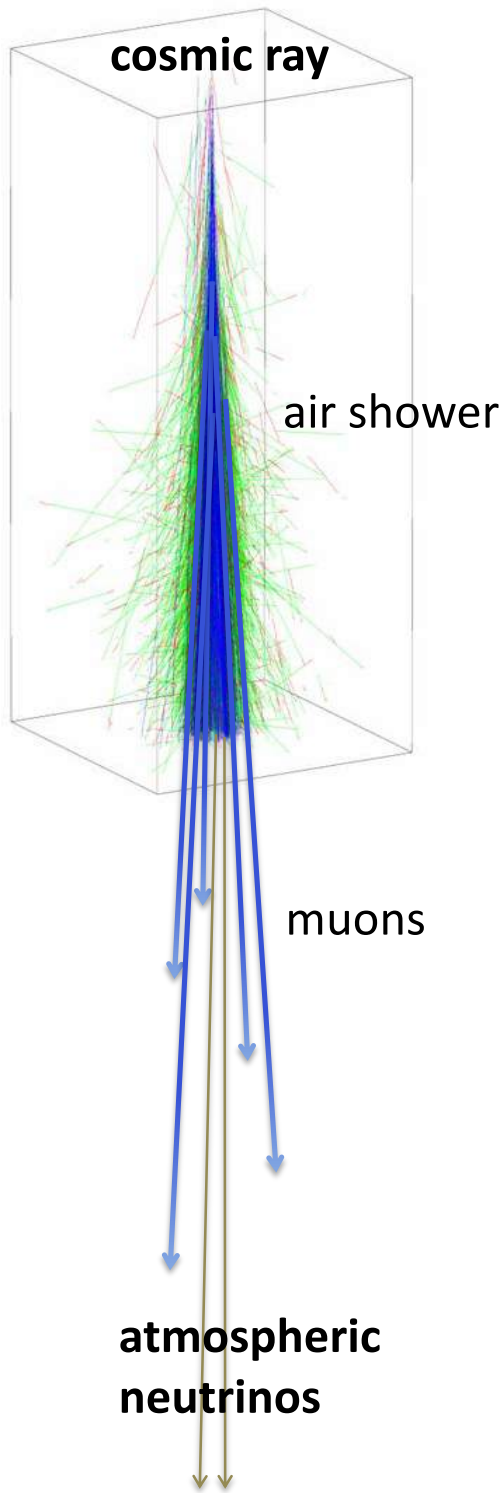
but only lower bound on neutrino energy

Cascade topology

Good energy resolution ($\sim 15\%$)

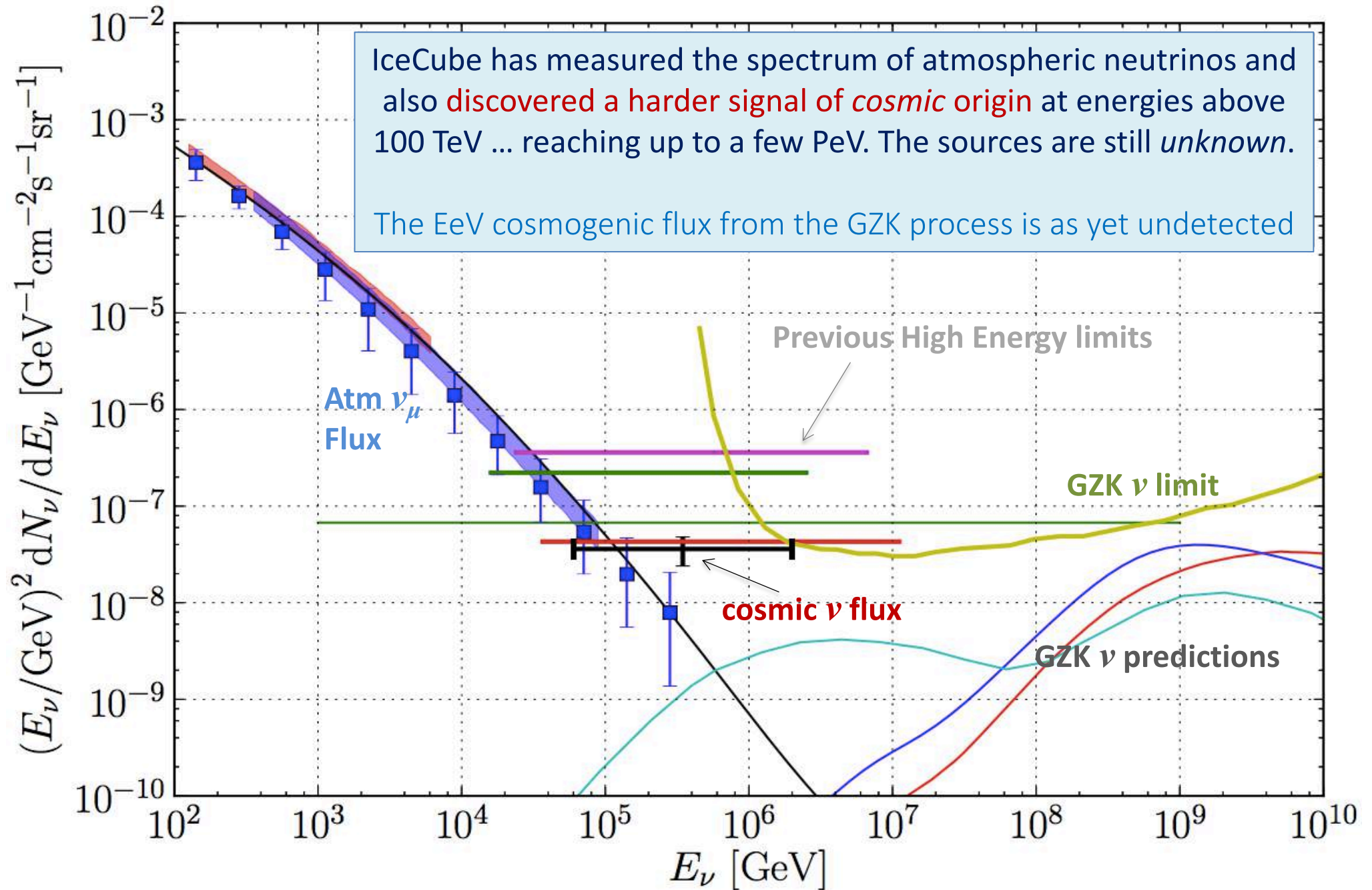
but poor pointing ($\sim 10^\circ - 15^\circ$)

Atmospheric Neutrino Spectrum

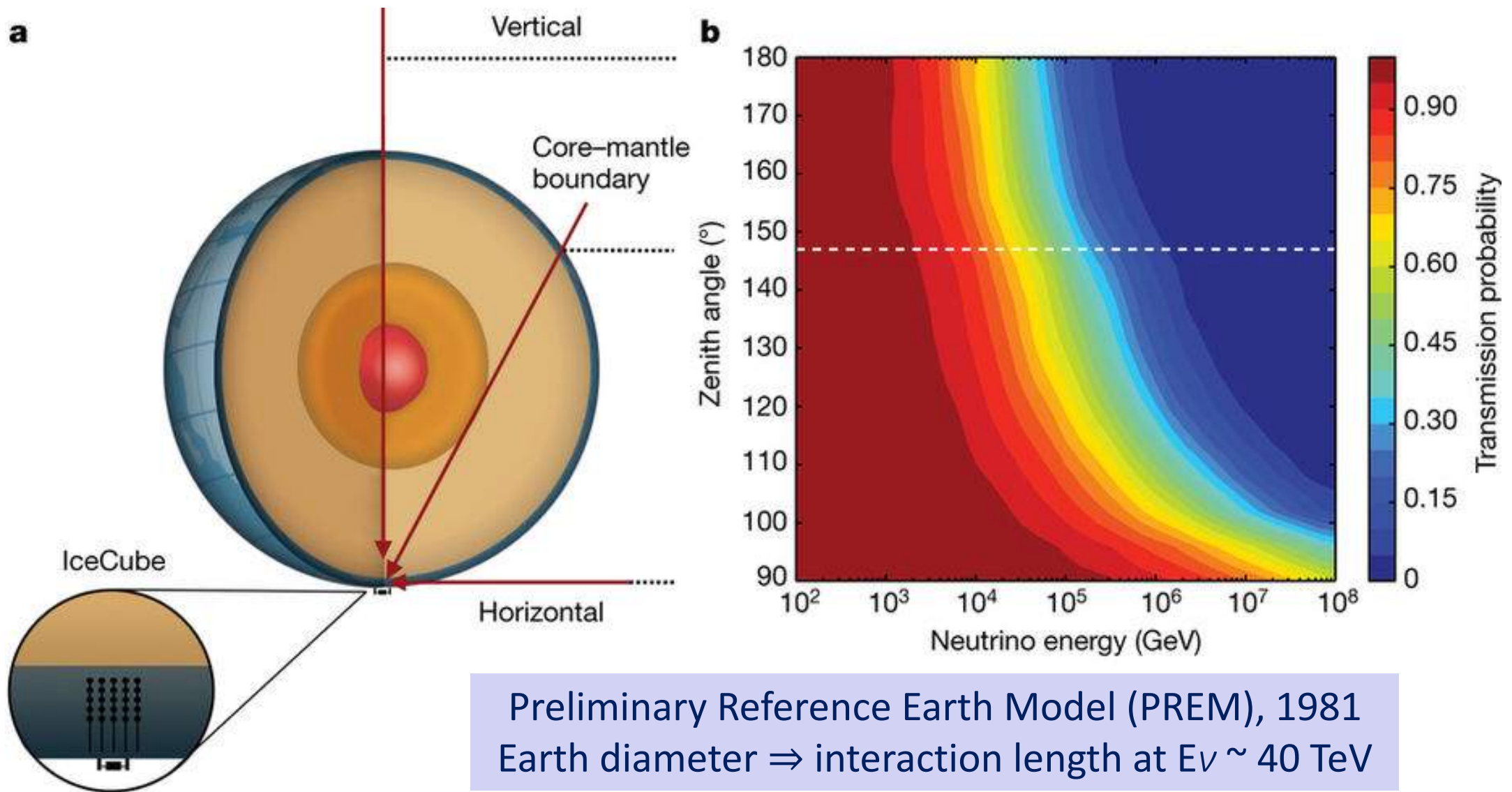


IceCube has measured the atmospheric neutrino background ... in good *agreement* with the number expected from cosmic ray interactions in the atmosphere creating pions and kaons (the 'prompt' flux from charmed meson decays not detected yet)

Current picture of high energy neutrino energy spectrum

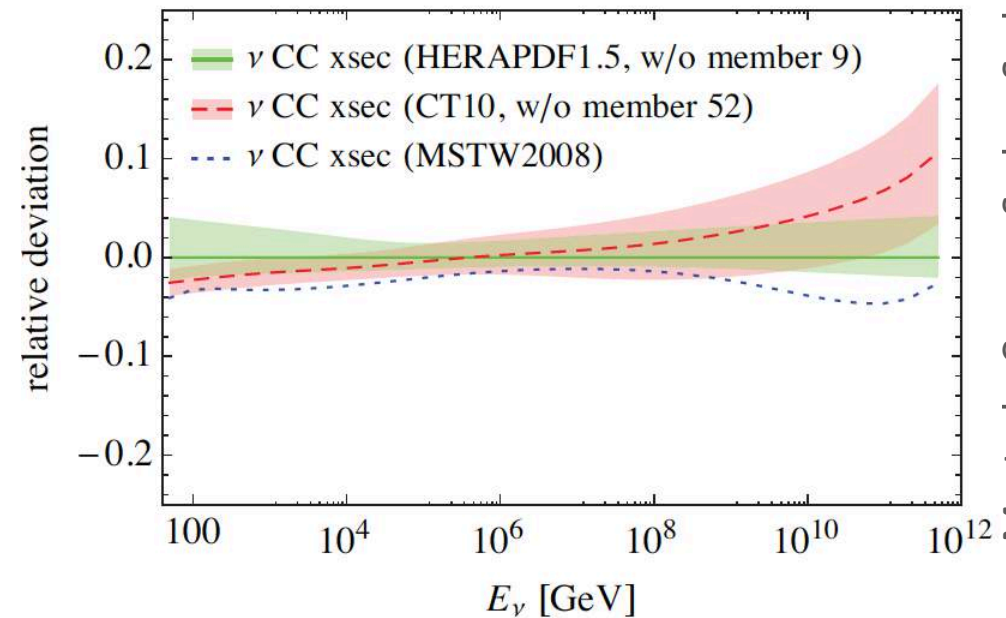
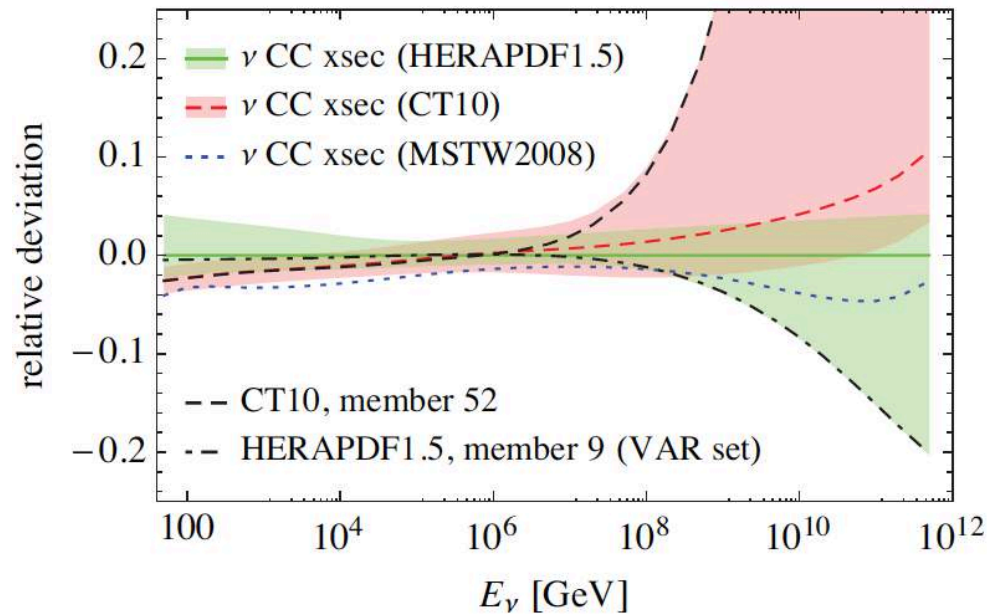
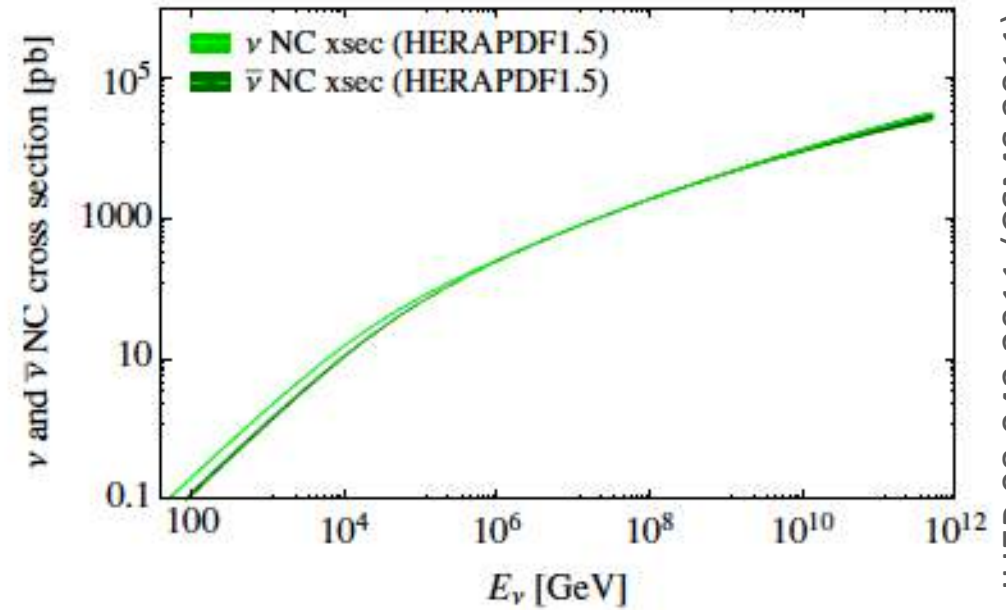
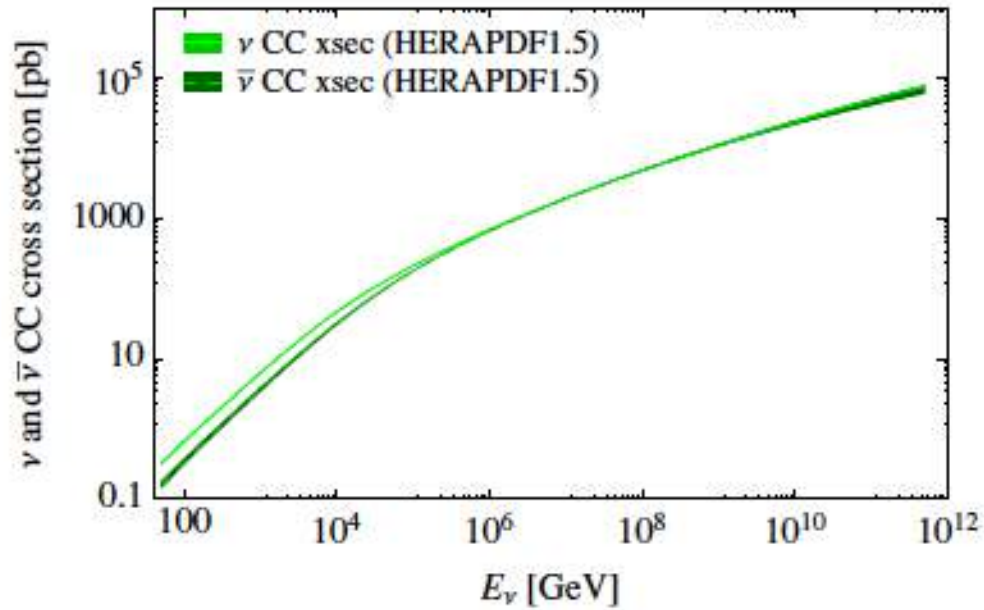


We can now *measure* the ν - N cross-section by examining the zenith angle dependence of the ν flux (assumed* to be isotropic)



*Isotropy certainly true for atmospheric neutrinos (which dominate up to $\sim 10^5$ GeV) and for extragalactic flux too ... galactic component is $< 18\%$

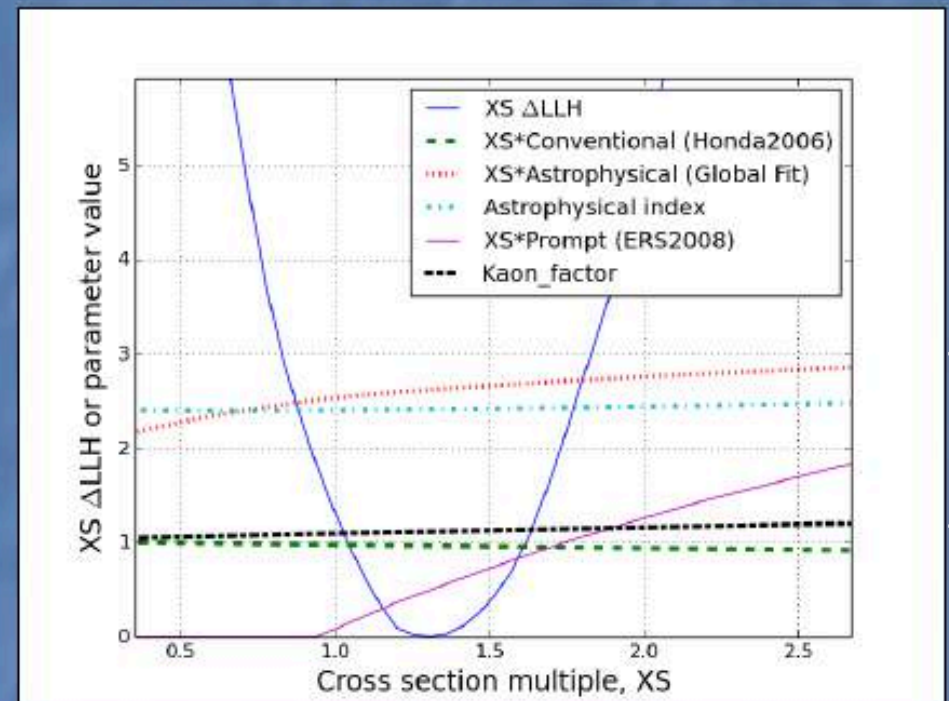
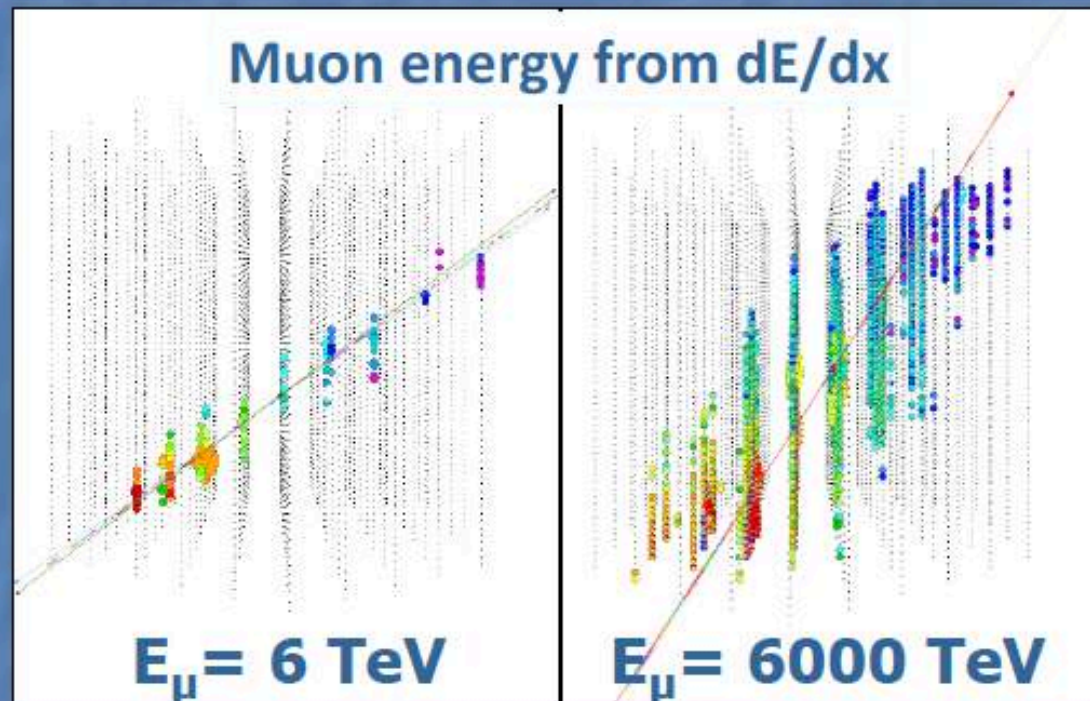
Meanwhile we have recalculated the ν - N #-section @ NLO with \sim few % accuracy using HERAPDF1.5



... finding good agreement between different PDF sets (*after* we reject unphysical members – which would have yielded e.g. a *negative* F_L or too steep rise in #-section)

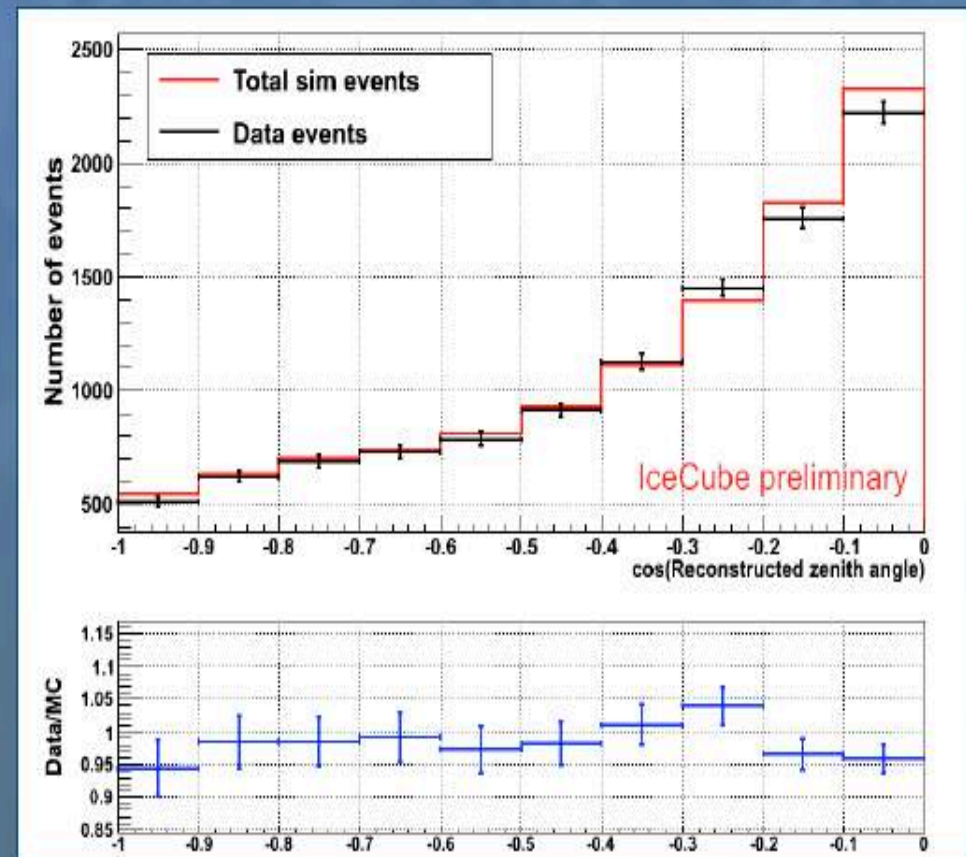
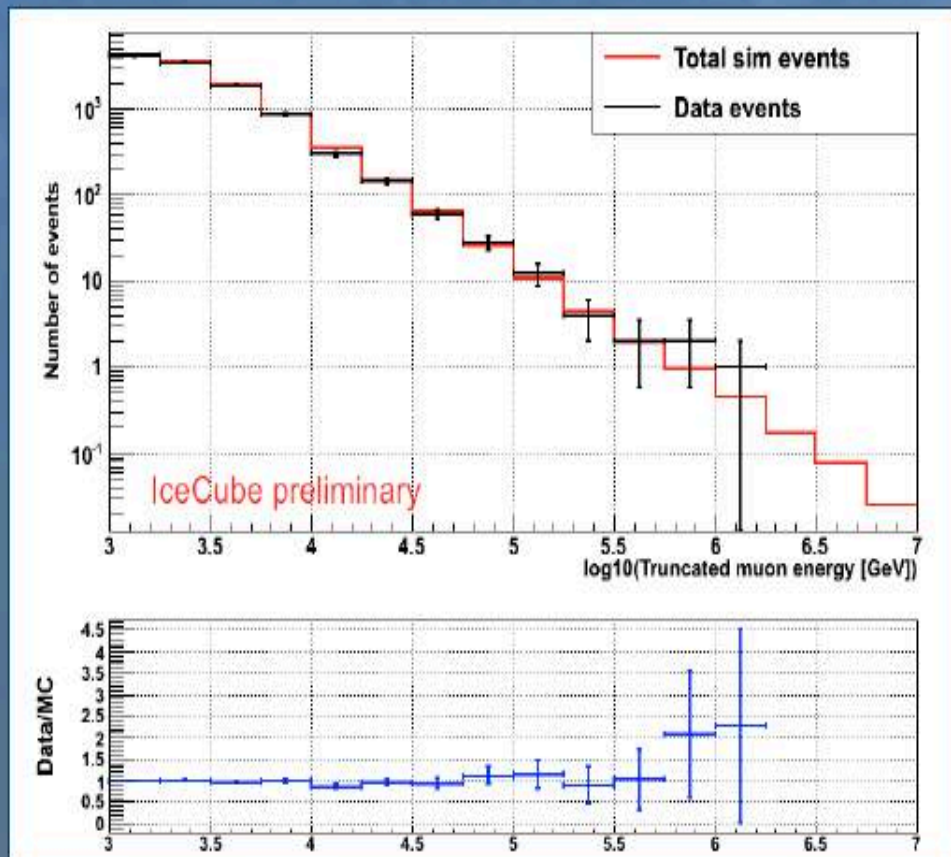
Experimental method

- **Event selection yielded 10,784 muon neutrinos in 2010 data year**
 - Muon energy determined by Truncated Energy method [IceCube 2013]
- **Two-dimensional LLH fit in muon energy and zenith angle**
- **Constrained by priors from other experiments**
 - Astrophysical and prompt fluxes from IceCube [IceCube 2015]
- **Best fit is multiple of Standard Model expectation from CSMS 2011**
 - Fit parameters include fluxes of conventional, astrophysical, prompt, plus $\nu_\mu \bar{\nu}_\mu$ ratio, kaon-pion ratio, DOM efficiency
 - Systematics include ice model, Earth model, atmospheric temperature model, and choice of astrophysical and prompt flux priors

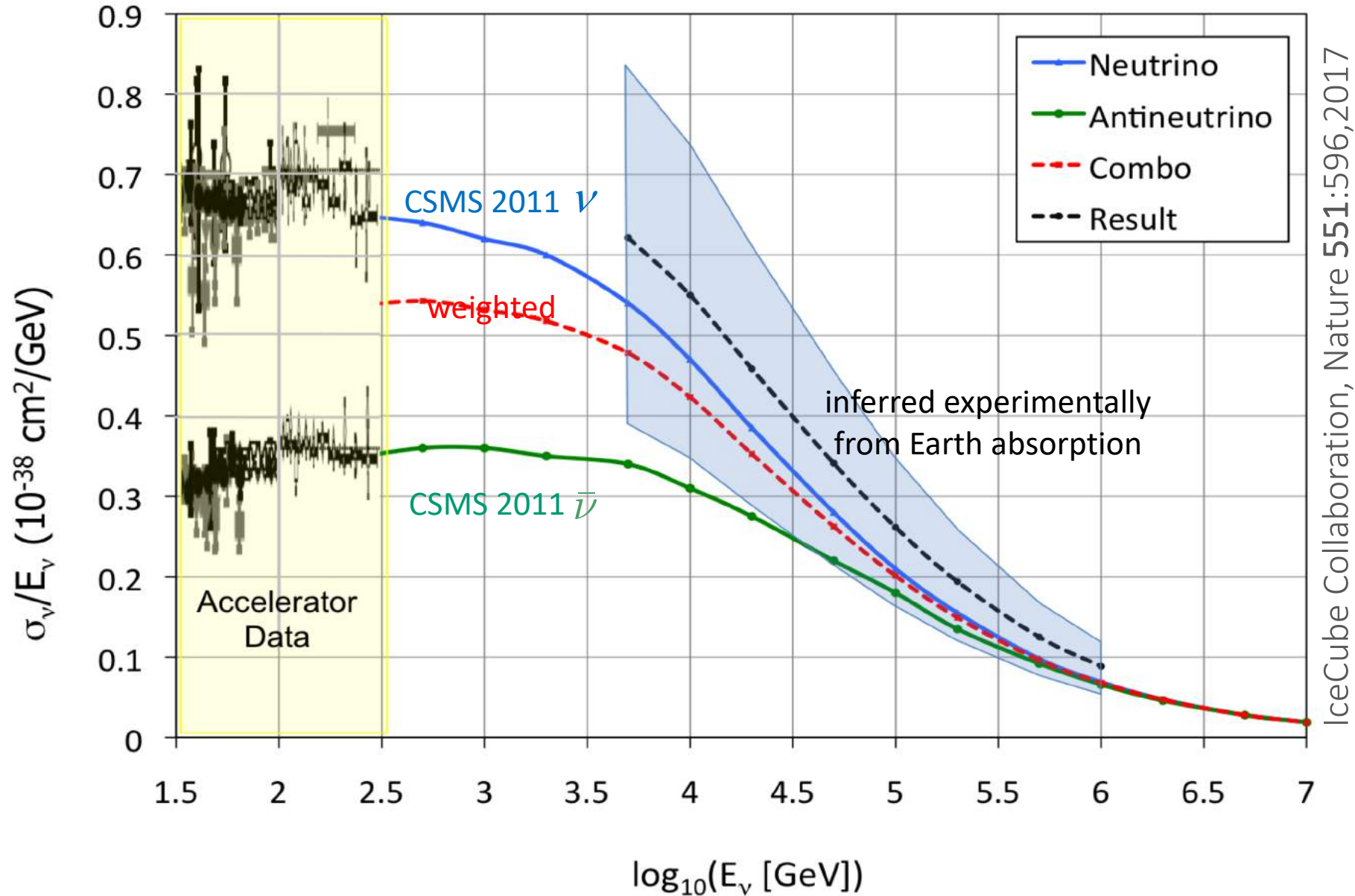


Results

- Total ν_{μ} -nucleon cross section = $1.30^{+0.30}_{-0.26}$ (stat.) $^{+0.32}_{-0.39}$ (syst.) times CSMS 2011 expectation
 - Energy range 5.6 TeV to 620 TeV
- In agreement with the Standard Model cross section at high energy
- Plans for follow-up analysis using 5+ years of data

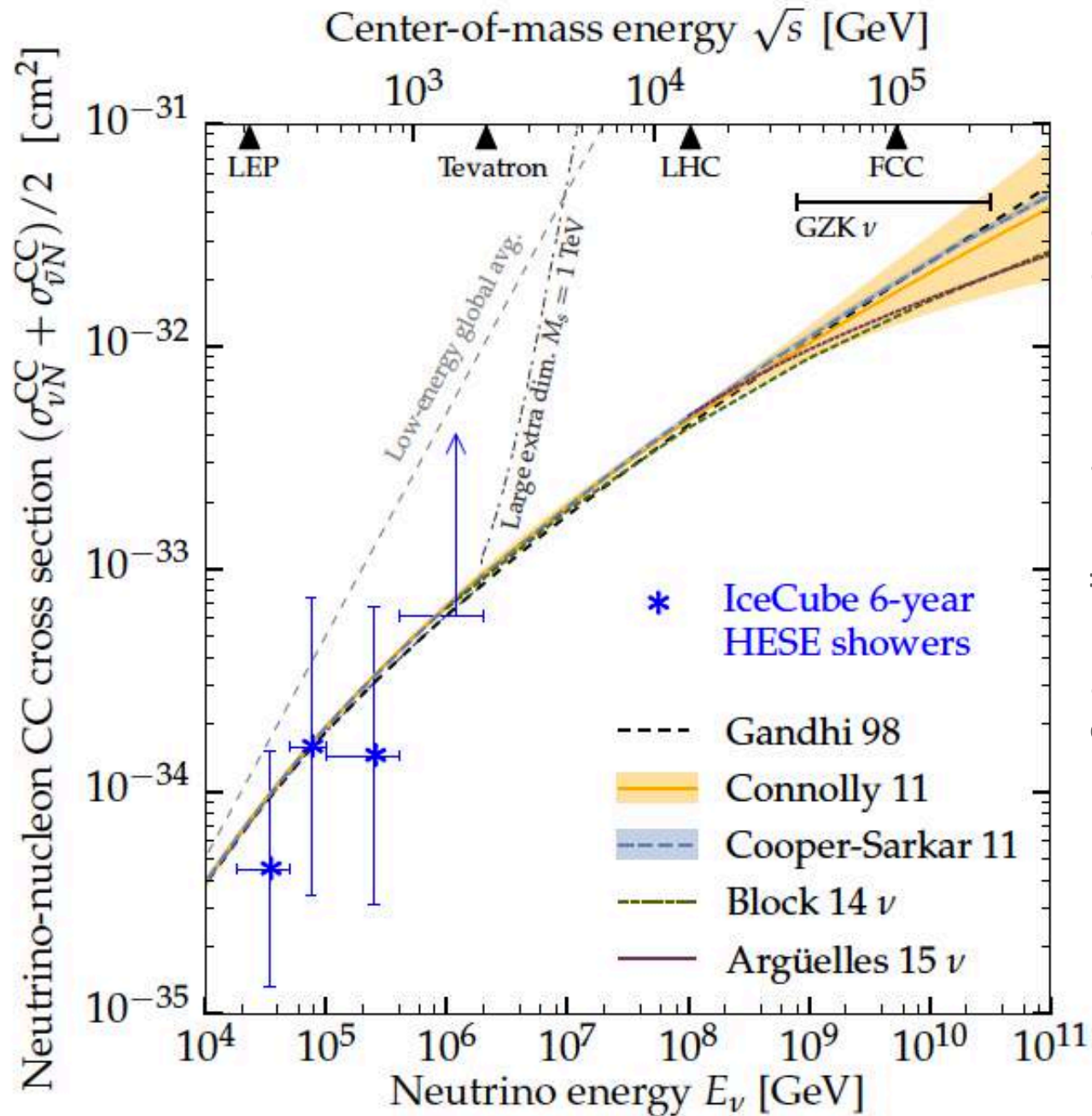


No evidence of deviation (within $\pm 30\%$) from SM up to 620 TeV



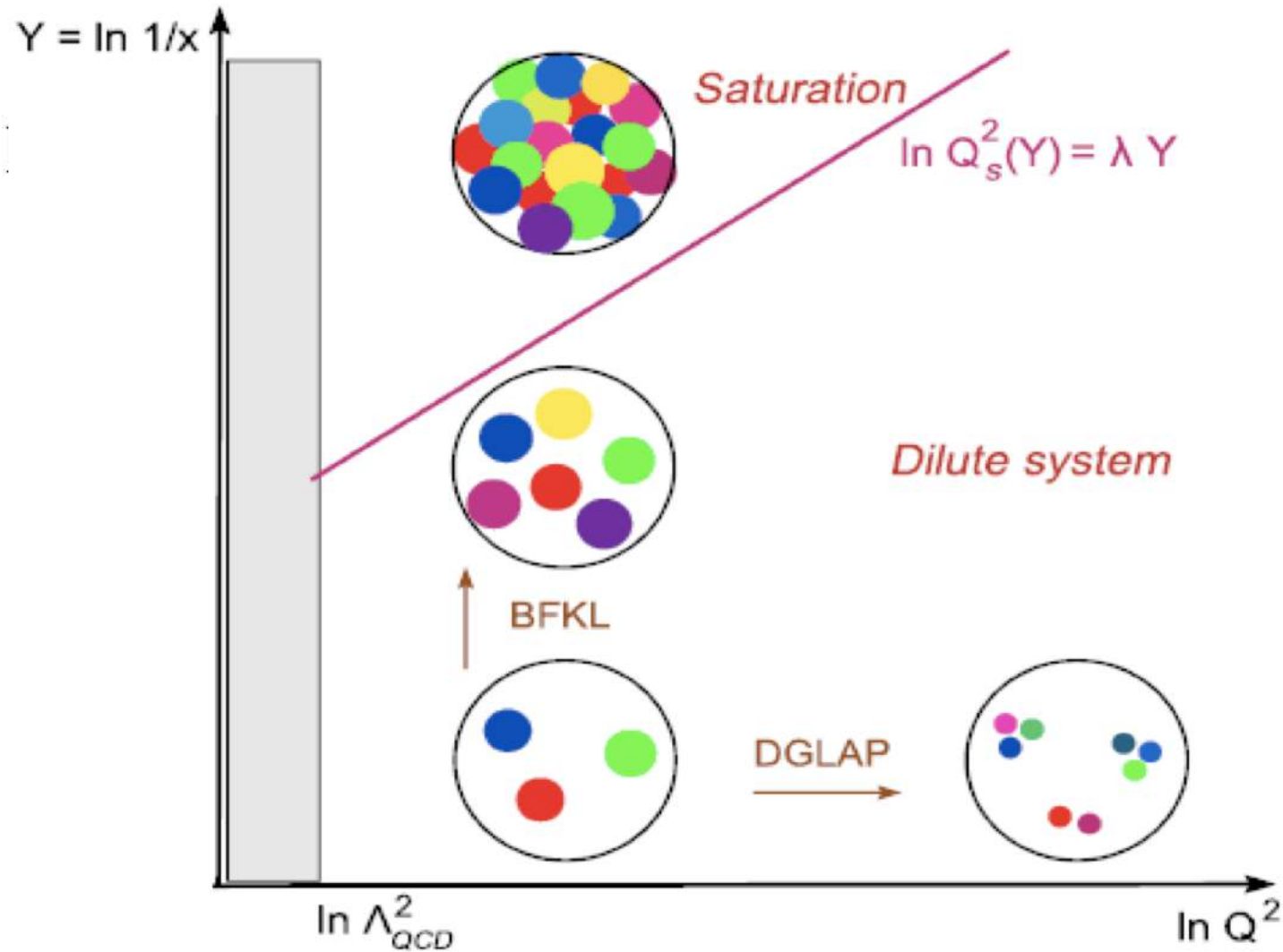
Powerful probe of new physics beyond the SM – from an *astroparticle* experiment ...
should be able to probe up to $\sim 10^9$ GeV using cosmogenic ν ... with **IceCube-Gen2**!

Constrains e.g. large new dimension at low scales (but LHC got there first!)



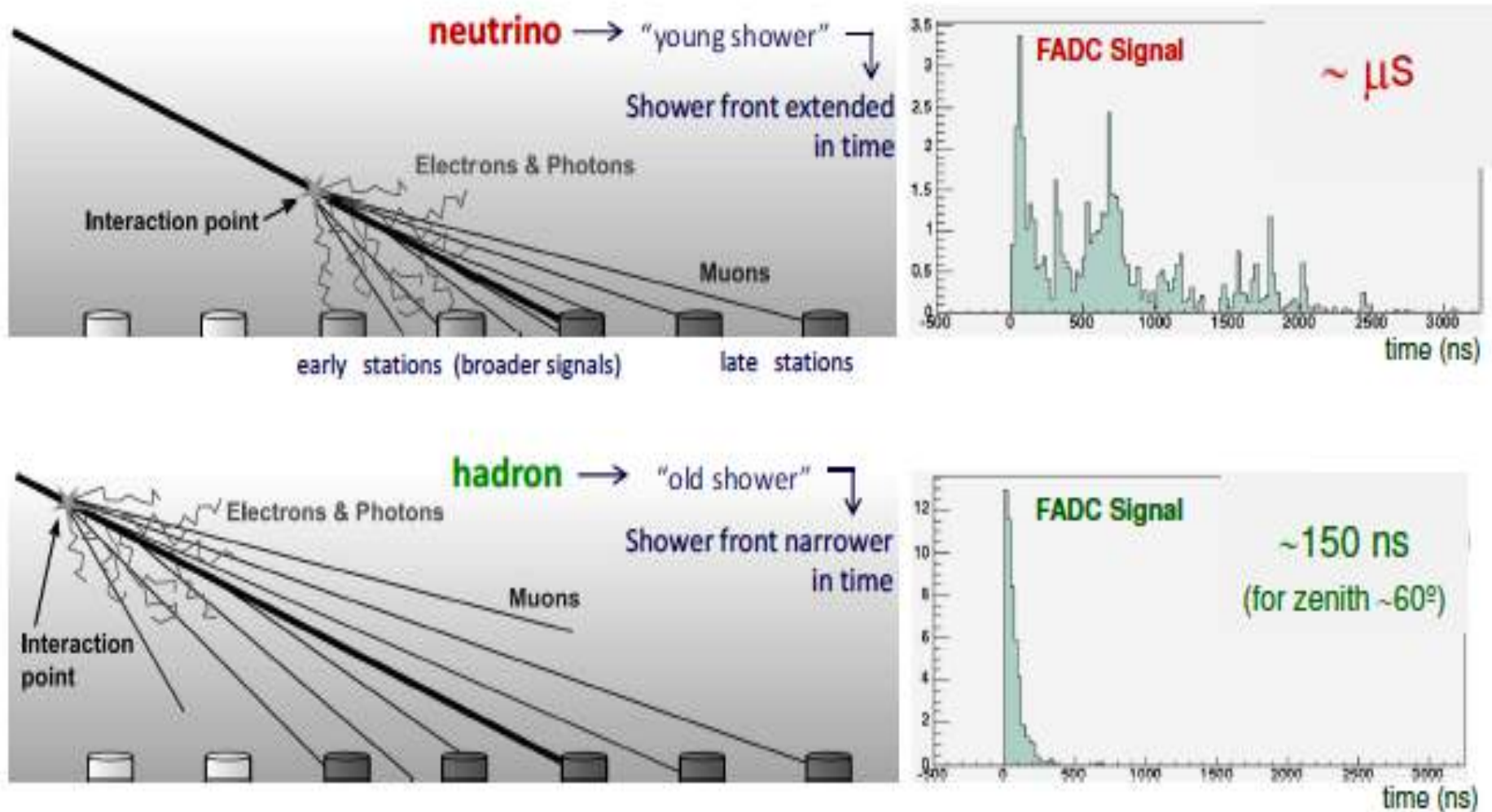
Bustamante & Connolly, arXiv:1711.11043

As the gluon density rises at low x , non-perturbative effects must become important ... a new phase of QCD - **Colour Glass Condensate** - has been postulated to exist (and has some support from RHIC and ALICE data)



This would strongly suppress the ν -N #-secn below its (unscreened) SM value ... can we test this experimentally with UHE cosmic neutrinos?

An unexpected bonus – UHE neutrino detection with air shower arrays



(courtesy:Sergio Navas)

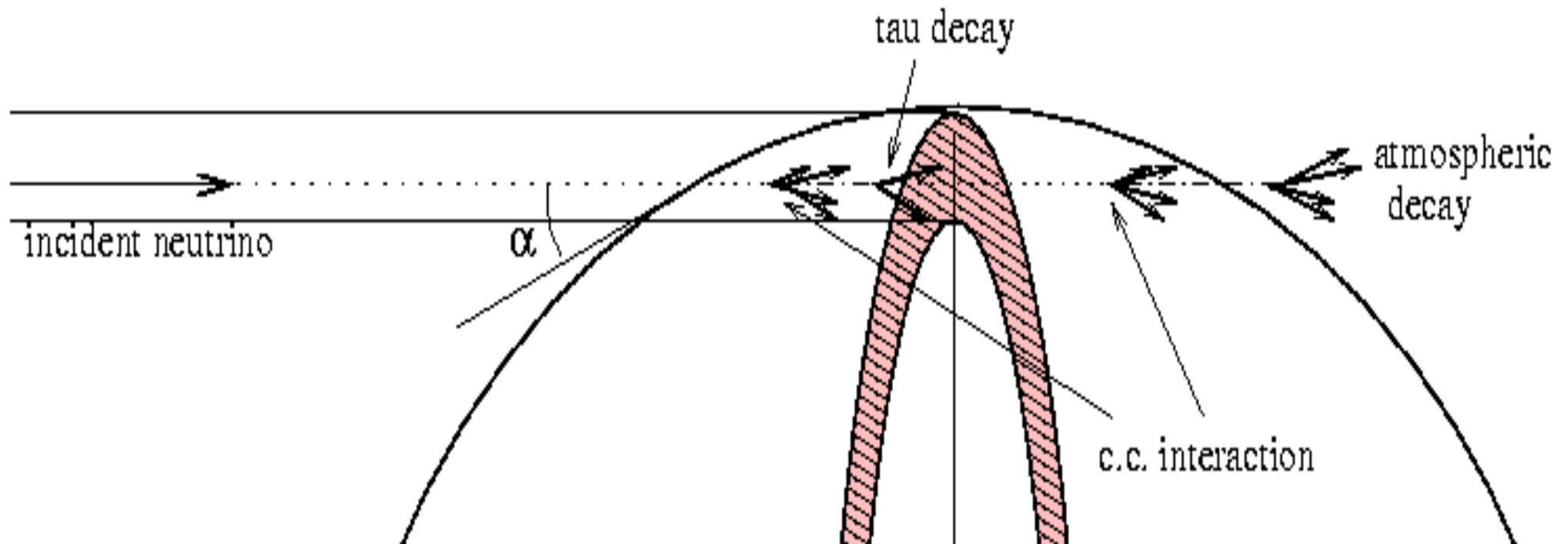
When a cosmic ray (hadron) interacts close to the horizon, the large path length in the atmosphere ensures absorption of charged particles apart from very high energy muons ... However neutrinos can penetrate through the atmosphere and interact close to the array so if we see a *young shower* at a *large zenith angle*, that is a candidate for a UHE neutrino!

Event rate \propto cosmic neutrino flux (all flavours) and ν -N DIS cross-section

An unexpected bonus – UHE neutrino detection with air shower arrays

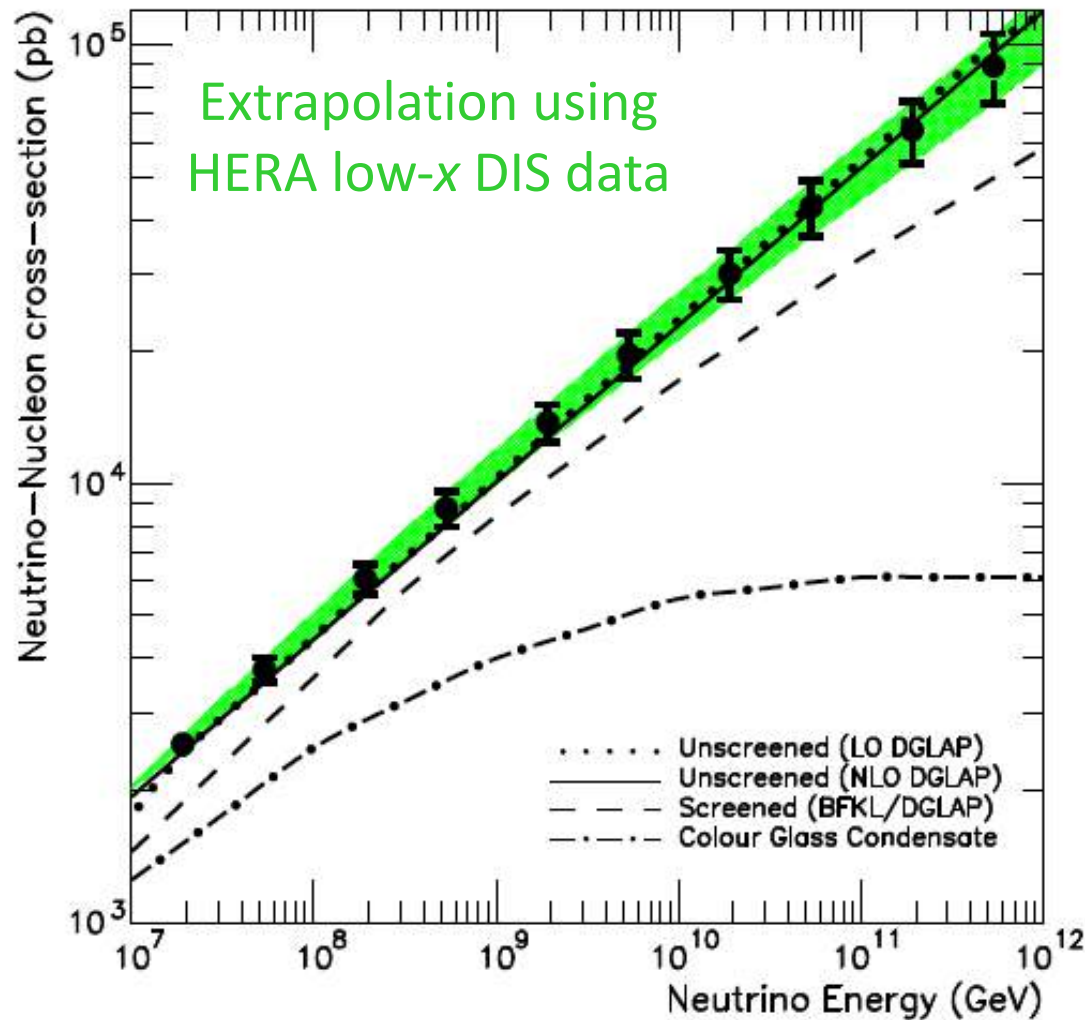
Auger can also see Earth-skimming $\nu_\tau \rightarrow \tau$ which generates *upgoing* hadronic shower (detectable only because the surface detector tanks are raised above the ground)

Neutrino oscillations en-route to Earth should *equilibrate* flavours with $\nu_e:\nu_\mu:\nu_\tau::1:1:1$ so there will be tau neutrinos in the cosmic beam regardless of initial composition

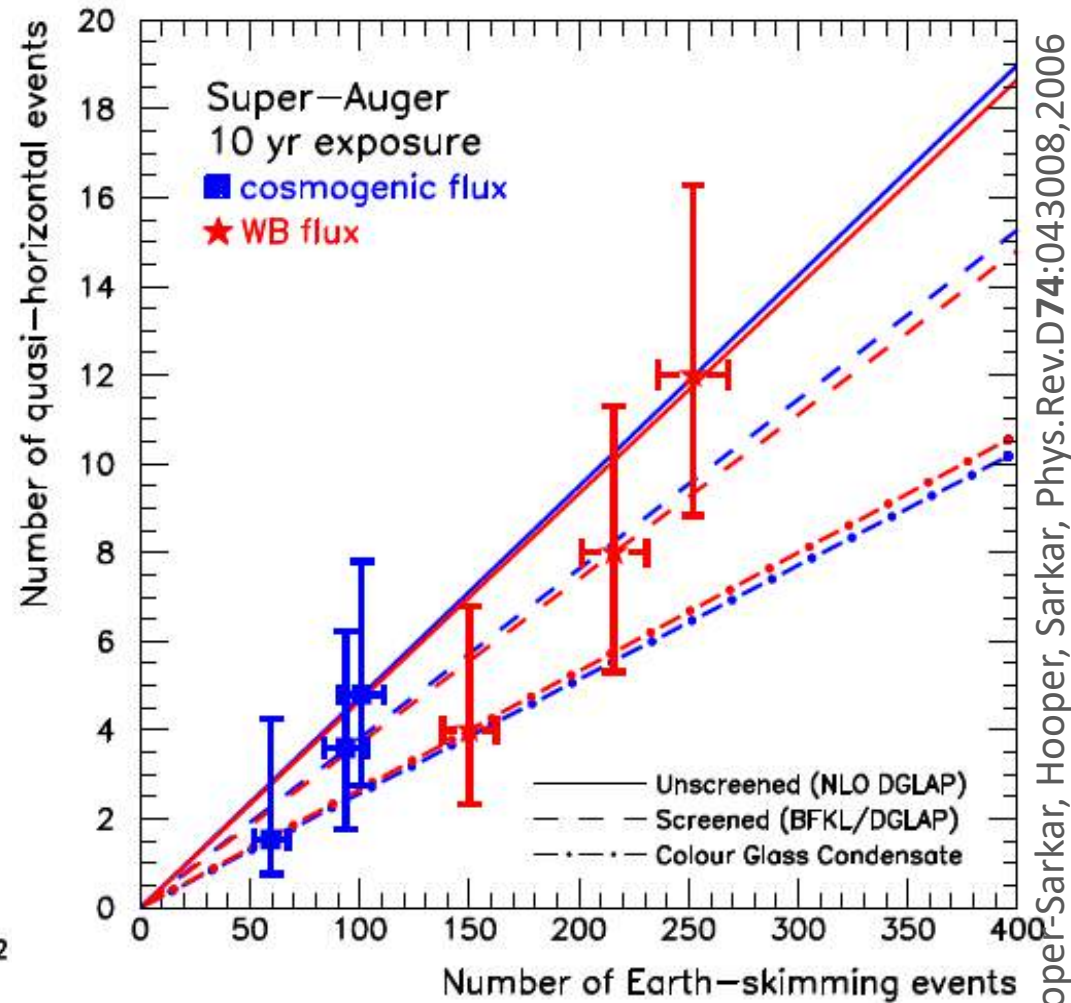


The rate is still \propto the cosmic neutrino flux, but *not* to the ν - N #-section (since higher values also imply stronger *absorption* in the Earth)

Hence low- x QCD *can* be probed with cosmic UHE neutrinos with a *very* large detector... *unfortunately Auger-N was not built*



The steep rise of the gluon density at low- x must saturate (unitarity!) \Rightarrow suppression of the ν - N #-section



The ratio of quasi-horizontal (all flavour) and Earth-skimming (ν_τ) events *measures* the #-section

Meanwhile radio signals spotted from unusual *upward-going* air showers



A neutrino induced cascade produces a coherent radio Cherenkov pulse.

~680km to horizon ->
 $1.5 \times 10^6 \text{ km}^3$ interaction volume

Incident neutrinos
 With energies above ~0.5 EeV

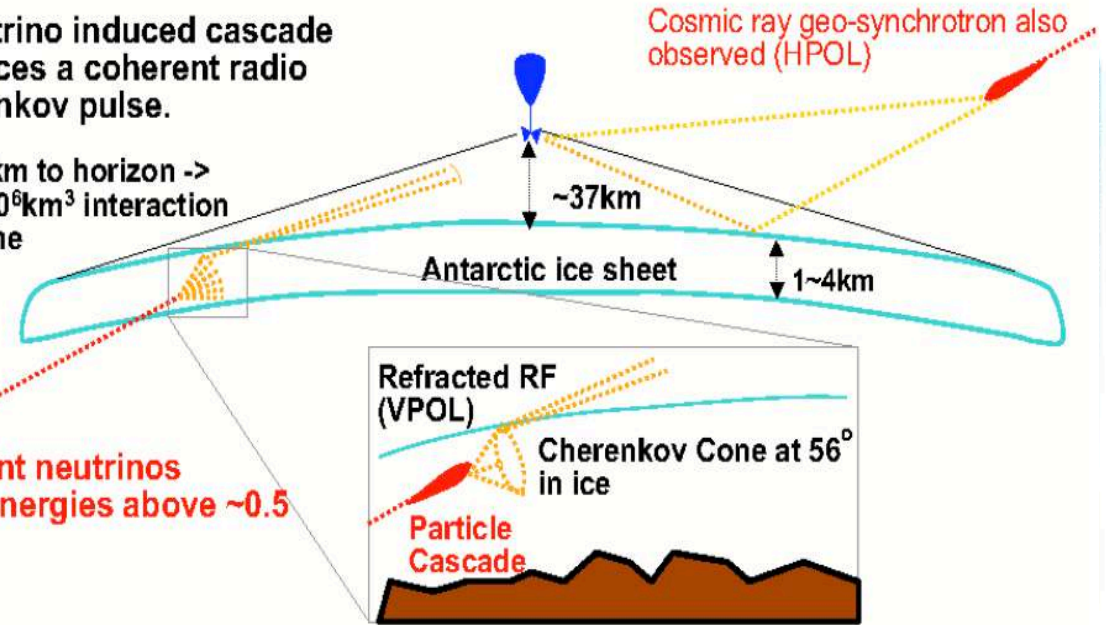


TABLE I: ANITA-I,-III anomalous upward air showers.

event, flight	3985267, ANITA-I	15717147, ANITA-III
date, time	2006-12-28,00:33:20UTC	2014-12-20,08:33:22.5UTC
Lat., Lon. ⁽¹⁾	-82.6559, 17.2842	-81.39856, 129.01626
Altitude	2.56 km	2.75 km
Ice depth	3.53 km	3.22 km
El., Az.	$-27.4 \pm 0.3^\circ$, $159.62 \pm 0.7^\circ$	$-35.0 \pm 0.3^\circ$, $61.41 \pm 0.7^\circ$
RA, Dec ⁽²⁾	282.14064, +20.33043	50.78203, +38.65498
$E_{shower}^{(3)}$	$0.6 \pm 0.4 \text{ EeV}$	$0.56^{+0.3}_{-0.2} \text{ EeV}$

¹ Latitude, Longitude of the estimated ground position of the event.

² Sky coordinates projected from event arrival angles at ANITA.

³ For upward shower initiation at or near ice surface.

Observation of an Unusual Upward-going Cosmic-ray-like Event in the Third Flight of ANITA

We report on an upward traveling, radio-detected cosmic-ray-like impulsive event with characteristics closely matching an extensive air shower. This event, observed in the third flight of the Antarctic Impulsive Transient Antenna (ANITA), a NASA-sponsored long-duration balloon payload, is consistent with a similar event reported in a previous flight. These events may be produced by the atmospheric decay of an upward-propagating τ -lepton produced by a ν_τ interaction, although their relatively steep arrival angles create tension with the standard model (SM) neutrino cross section. Each of the two events have *a posteriori* background estimates of $\lesssim 10^{-2}$ events. If these are generated by τ -lepton decay, then either the charged-current ν_τ cross section is suppressed at EeV energies, or the events arise at moments when the peak flux of a transient neutrino source was much larger than the typical expected cosmogenic background neutrinos. Gorham *et al*, arXiv:1803.05088

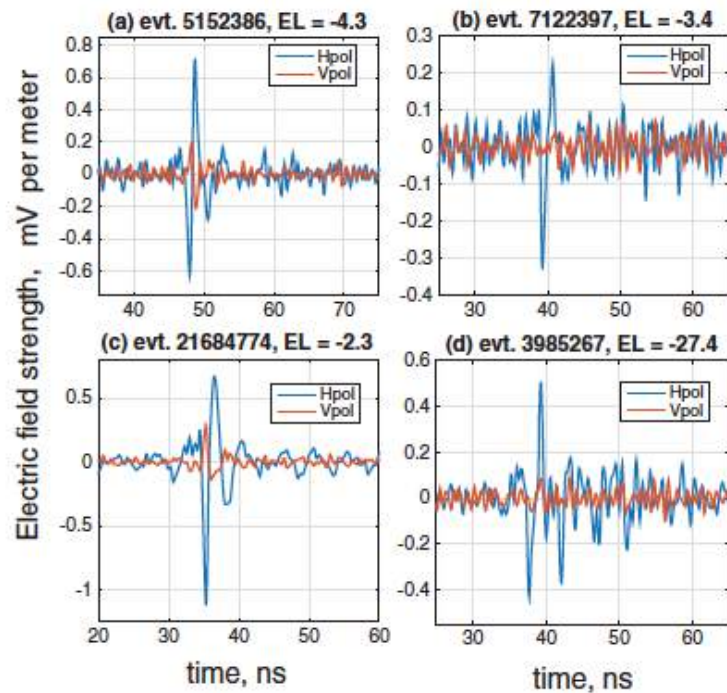
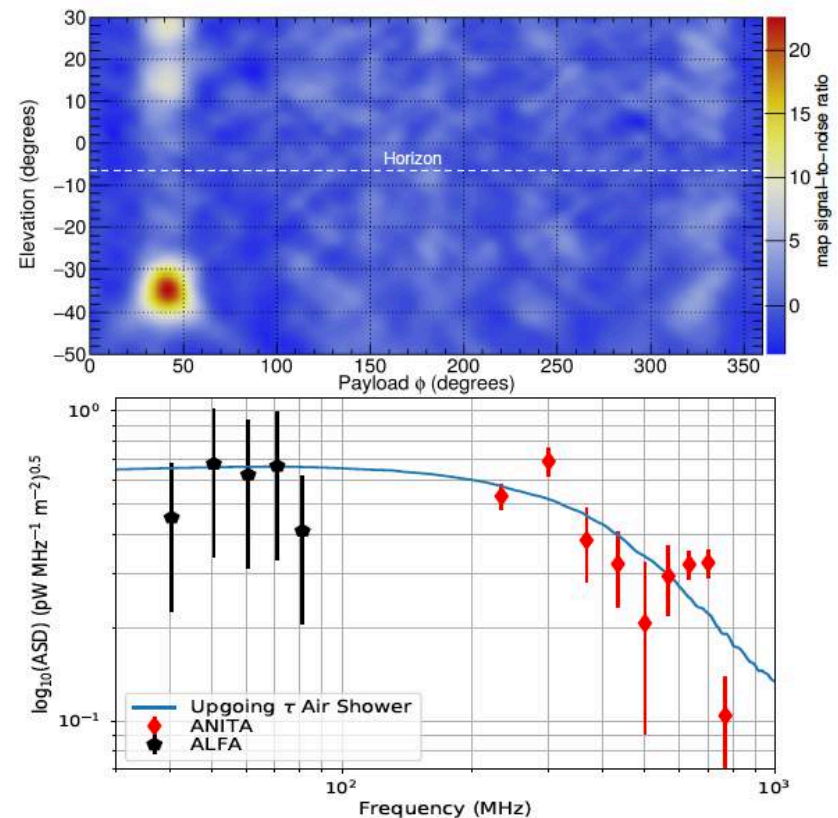


FIG. 1. Waveforms for the four events described here. Events are indexed here and in the text by the letters A, B, C, and D.

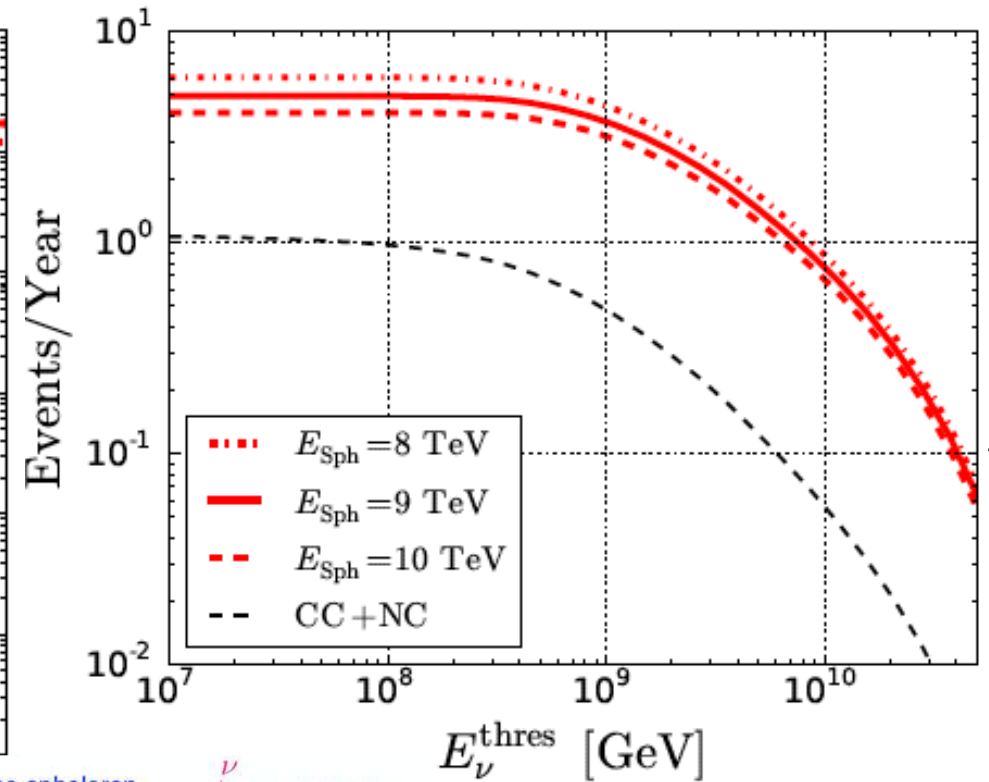
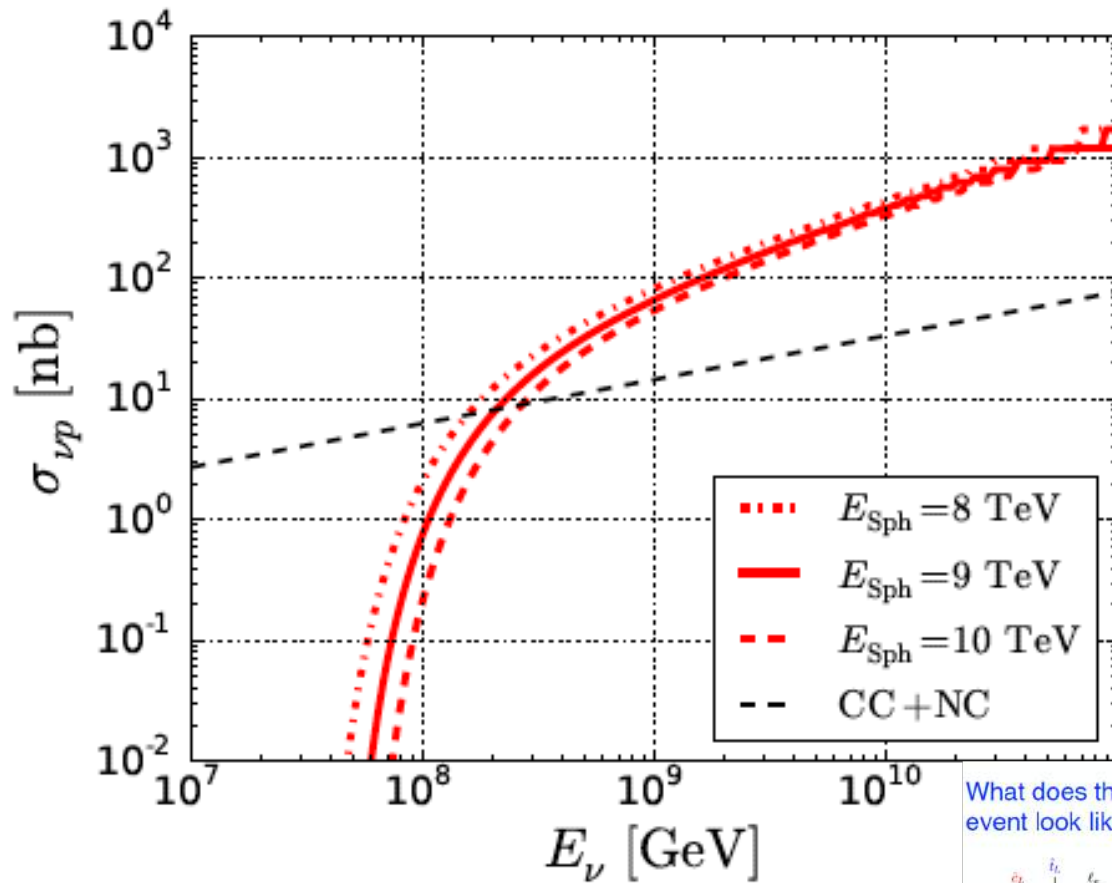


... or the ν - N cross-section may be much *higher* than in the SM

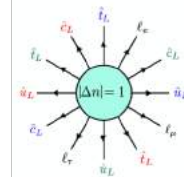
Non-perturbative transitions between degenerate $SU(2)$ vacua (with different $B+L$ #) are exponentially suppressed below the “sphaleron” mass: $\sim M_W/\alpha_W \sim 9$ TeV (update by Tye & Wong, PRD 92:045005,2015) ... *large* cross-sections are predicted for ν - N scattering at higher cms energies

$$E_\nu \geq E_{\text{sph}}^2/2xm_N \simeq 4 \times 10^7/x \sim 10^{9-11} \text{ GeV}$$

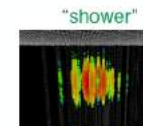
Han & Hooper, PLB 582:21,2004



What does the sphaleron event look like?

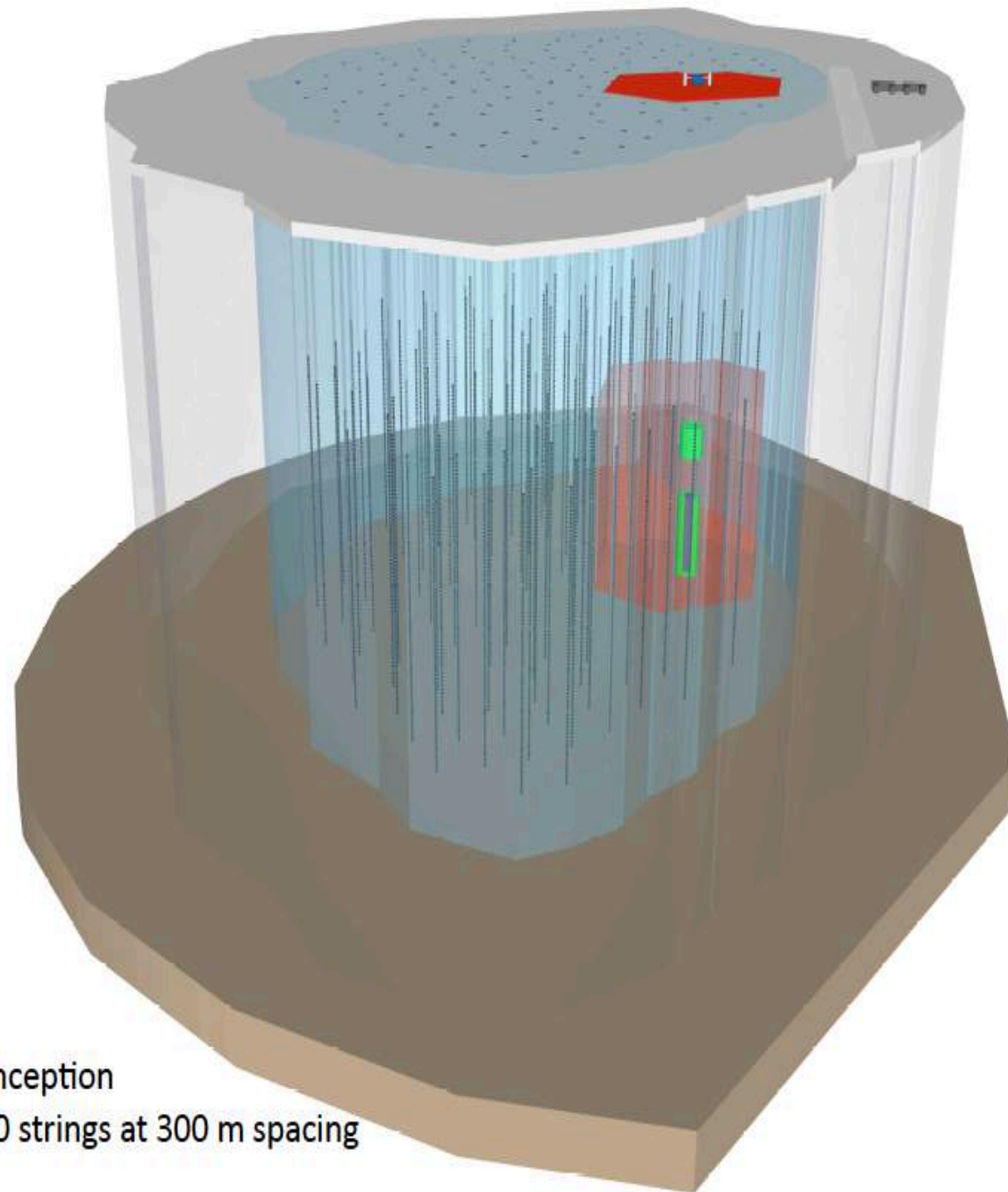


$E_\nu \gtrsim 10^{8-10}$ GeV



IceCube has sensitivity to EW sphalerons comparable to LHC!

To do astronomy *and* particle physics with cosmic neutrinos we must think BIG!



IceCube-Gen2

Artist conception
Here: 120 strings at 300 m spacing