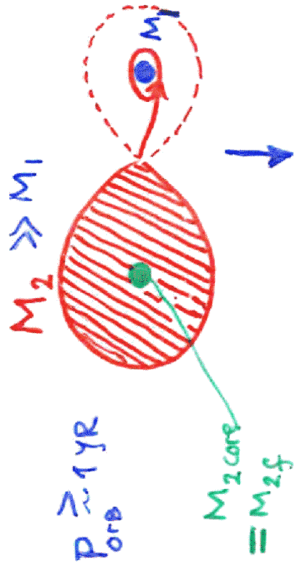


COMMON-ENVELOPE EVOLUTION

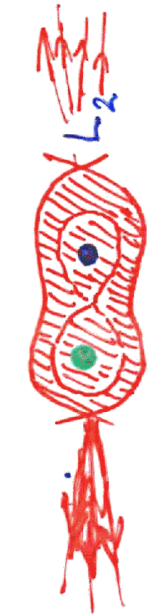
(Paczynski 1975
Ostriker 1975
Webbink 1975)

- TAAM, Bodenheimer,
Ostriker (1978)
(HMXB)
- Meijer & Mejer-Hof
meister (1979)
(Formation of CV)

RED (SUPER-) GIANT



LARGE FRICTIONAL DRAG & EJECTION OF COMMON ENVELOPE



COMMON ENVELOPE

$P_{orb} \leq 1 \text{ day}$



REASONS FOR COMMON-ENVELOPE FORMATION

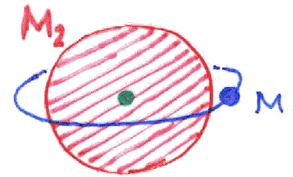
LACK OF CO-ROTATION (Paczynski 1975
Ostriker 1975
Webbink 1975)

1. TIDAL INSTABILITY (G. Darwin 1895):

MOMENT OF INERTIA of GIANT >

$$\frac{1}{3} (\text{MOM. OF Inertia BINARY})$$

[Requires $M_2 \geq (5-10) M_1$]



2. UNSTABLE MASS TRANSFER:

$$\tau_{transfer} \ll \tau_{thermal} (\text{accretor})$$

THIS CAN BE DUE TO:

A. Mass ratio $q = M_1/M_2 \ll 1$

($\tau_{th} \propto 1/M^2$) \Rightarrow if $q < 0.5$: PROBLEM

OR B. CONVECTIVE ENVELOPE of DONOR EXPANDS UPON MASS LOSS

$$P \propto \rho^{5/3} \rightarrow \text{Polytrope } n = 3/2$$

MASS LOSS ON DYNAMICAL TIMESCALE

$$\tau_M = 50 \text{ min} \left(\frac{P_0}{P_*} \right)^{1/2}$$

ROUGH ESTIMATE OF FINAL P_{orb}
(ENERGY ARGUMENT) webbink '79



ASSUME: $E_{bind}(\text{ENVELOPE}) \approx \text{LOSS OF ORBITAL ENERGY}$

$$\frac{G(M_{2f} + M_{2e})M_{2e}}{\lambda a_1 r_{L2}} = \alpha_{CE} \left[\frac{G M_1 M_{2f}}{2a_2} - \frac{G(M_{2f} + M_{2e})M_1}{2a_1} \right]$$

weight factor < 1

$$\therefore a_2 = a_1 \frac{M_1 M_{2f}}{M_{2f} + M_{2e}} \left/ \left(M_1 + \frac{2M_{2e}}{\alpha_{CE} \lambda r_{L2}} \right) \right.$$

EXAMPLE: Meyer & Meyer-Hofmeister (1979):

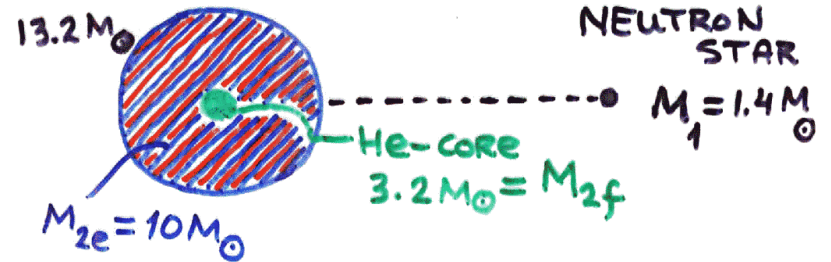
$\alpha_{CE} = 1$, $M_1 = M_{2f} = 1$, $M_{2e} = 4$, $r_{L2} = 0.5$, $\lambda = 0.5$

$$\therefore \frac{a_2}{a_1} = \frac{1}{5} / \left(1 + \frac{8}{0.25} \right) = \frac{1}{5 \times 33} = \frac{1}{165} \quad \text{for } \alpha = 1$$

$$P_2/P_1 = (a_2/a_1)^{3/2} \left(\frac{M_0}{M_f} \right)^{1/2} = \frac{1}{1300}$$

if $P_1 \approx 650^d \Rightarrow P_2 = 0.5^d$

EXAMPLE



$\alpha = 1 \Rightarrow a_2/a_1 = 1/250$

$\therefore \text{FOR } \left\{ \begin{array}{l} a_1 = 2.5 \text{ AU} \\ P_1 = 1.03 \text{ yr} \end{array} \right\} \Rightarrow \left\{ \begin{array}{l} a_2 = 0.01 \text{ AU} \\ = 2.16 R_{\odot} \\ P_2 = 4.09 \text{ HR} \end{array} \right.$



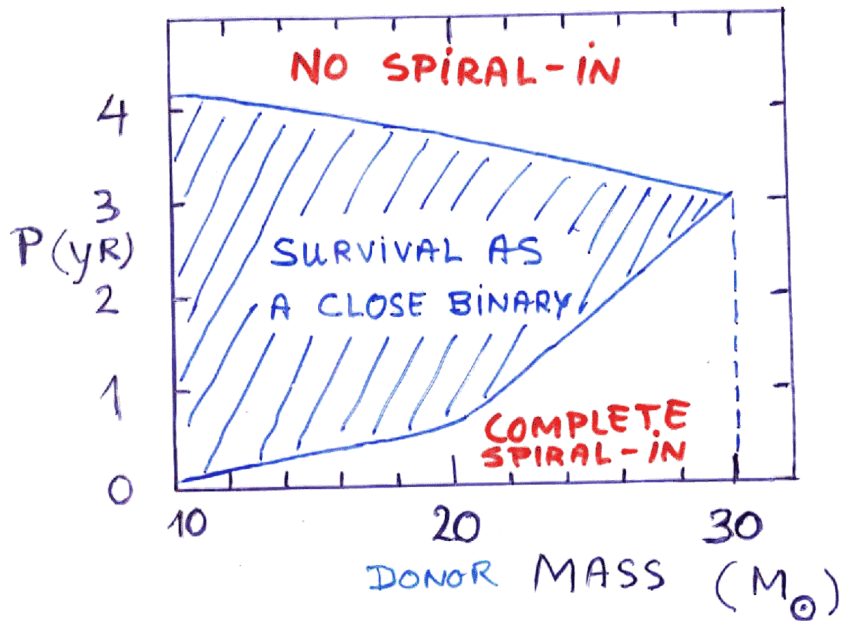
$R_{Roche}(M_{2f}) = 0.98 R_{\odot}$
 $R_{M_{2f}}(\text{He-star}) < 0.5 R_{\odot}$

$a_2 = 2.2 R_{\odot}$
 $P_2 = 4.09 \text{ HR}$

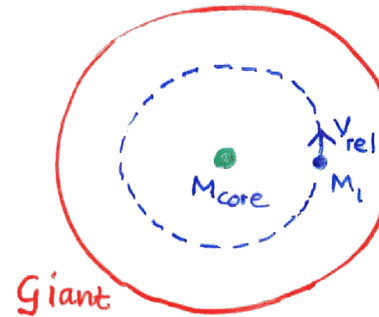
He-star will finally explode

\Downarrow
SHORT-PERIOD eccentric

OUTCOME OF COMMON-ENVELOPE
EVOLUTION OF MASSIVE X-BINARIES
(TAAM & BODENHEIMER
1992)



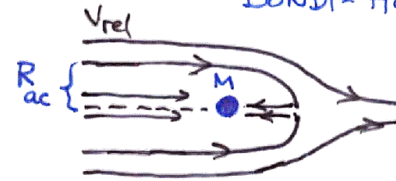
INCLUDES EFFECT
OF STRONG WIND MASS LOSS
OF SUPERGIANTS



INGREDIENTS OF
NUMERICAL MODEL
CALCULATIONS

(e.g.: Taam et al. 1995;
Terman et al. 1993
Podsiadlowski & co-
workers, 2002 \Rightarrow).

a) Energy dissipation: FRICTIONAL DRAG
BONDI-HOYLE APPROXIMATION



$$L_{drag} \sim \dot{M} V_{rel}^2$$

where:

$$\dot{M} = \pi R_{ac}^2 \rho V_{rel}$$

$$R_{ac} = \frac{2GM_1}{V_{rel}^2 + c_s^2}$$

$$V_{rel} = V_{orb} - V_{env}$$

b) ORBITAL EVOLUTION: LOSS OF GRAV. BINDING ENERGY

$$\frac{d}{dt} \left(\frac{GM_{core} \cdot M_1}{2a} \right) = -L_{drag}$$

N.B.: R_{ac} is mostly \gg Radius of M_1

$$R_{ac} = \frac{M}{M_{\odot}} \left(\frac{620 \text{ km/s}}{V_{rel}} \right)^2 R_{\odot}$$

e.g. if $a = 1 \text{ AU}$, $M_2 = 2M_1$, $M_1 = 1M_{\odot} \Rightarrow V_{rel} = 60 \text{ km/s}$

Terman, Taam & Hernquist 1994, ApJ, 422, 729
 3D - SPH simulation 4.67 M_⊙ RED GIANT + 0.94 M_⊙ DWARF

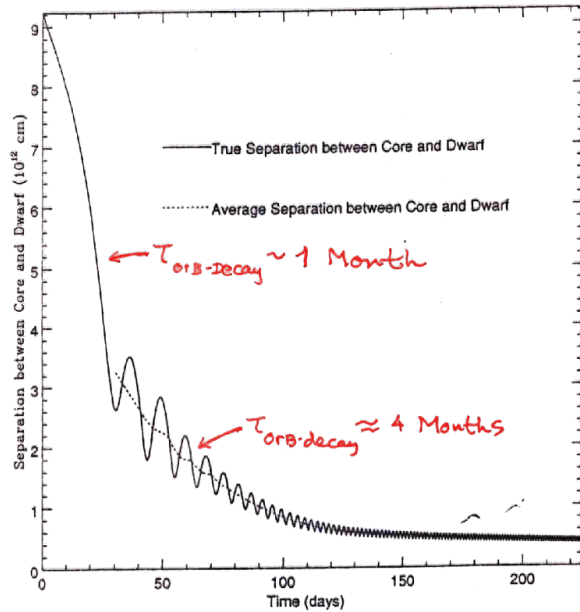


FIG. 1.—The variation of the orbital separation between the core of the red giant and the dwarf companion as a function of time. The solid curve corresponds to the actual separation between components, and the dashed line corresponds to the average separation. The separation is expressed in units of 10¹² cm, and the evolution time is expressed in days.

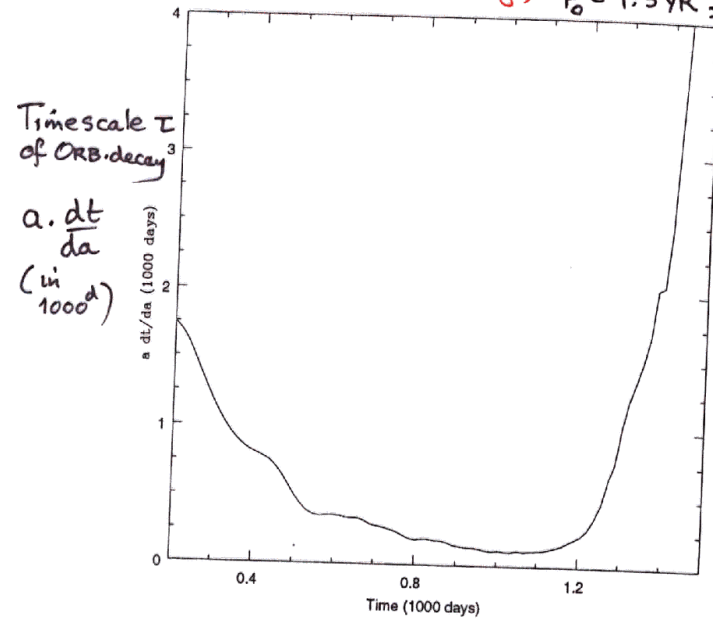
$P_0 = 0.47 \text{ yr } (a_0 = 1.07 \text{ AU}), M_{\text{core}} = 0.62 M_{\odot}$

Initial System in Co-Rotation $R_{\text{core}} = 5.47 R_{\odot}$

80% of mass ejected in ORB. Plane

They find $\alpha_{\text{CE}} \approx 0.15$ (\Rightarrow much of frictional energy release converted into HEAT and ...)

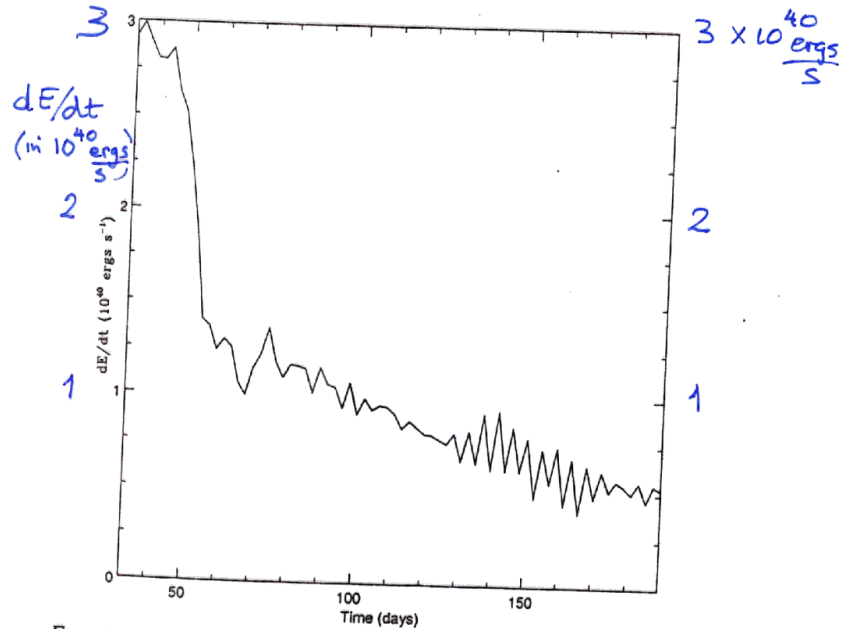
Taam (1995)
 16 M_⊙ Red Supergiant + 1.4 M_⊙ Neutron Star
 (Non-Rotating) $P_0 = 1.3 \text{ yr}, a_0 \approx 3 \text{ AU}$



The time scale of the orbital decay (a/\dot{a}) as a function of time. Both its of 1000 days.

- "Plunge" on timescale $\approx 200^d$
- $\tau_{\text{min}} = 95 \text{ d}$
- After $\approx 200^d$ of heavy frictional energy dissip., core reaches co-rotation

Ierman et al. (1994)
 4.67 M_⊙ Red Giant + 0.94 M_⊙ dwarf



16 M_⊙ Red Giant + 1.4 M_⊙ Neutron Star
 (SPH-simul. Taam 1995)

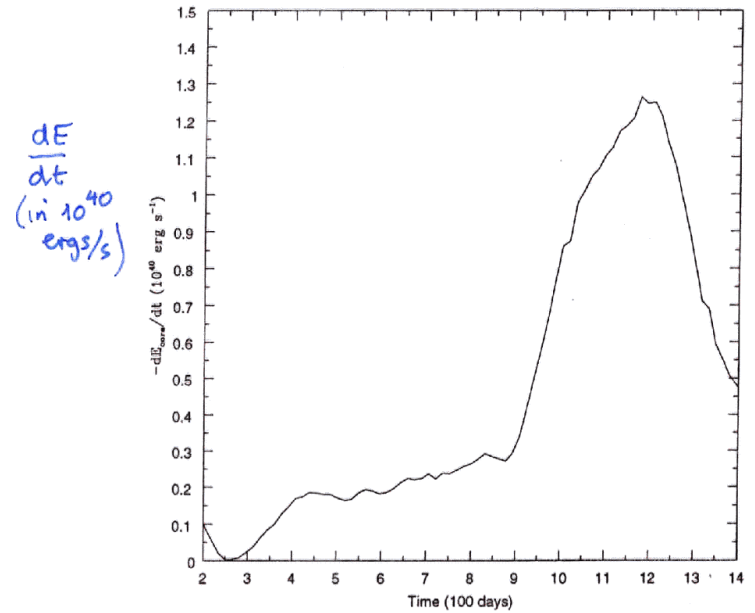
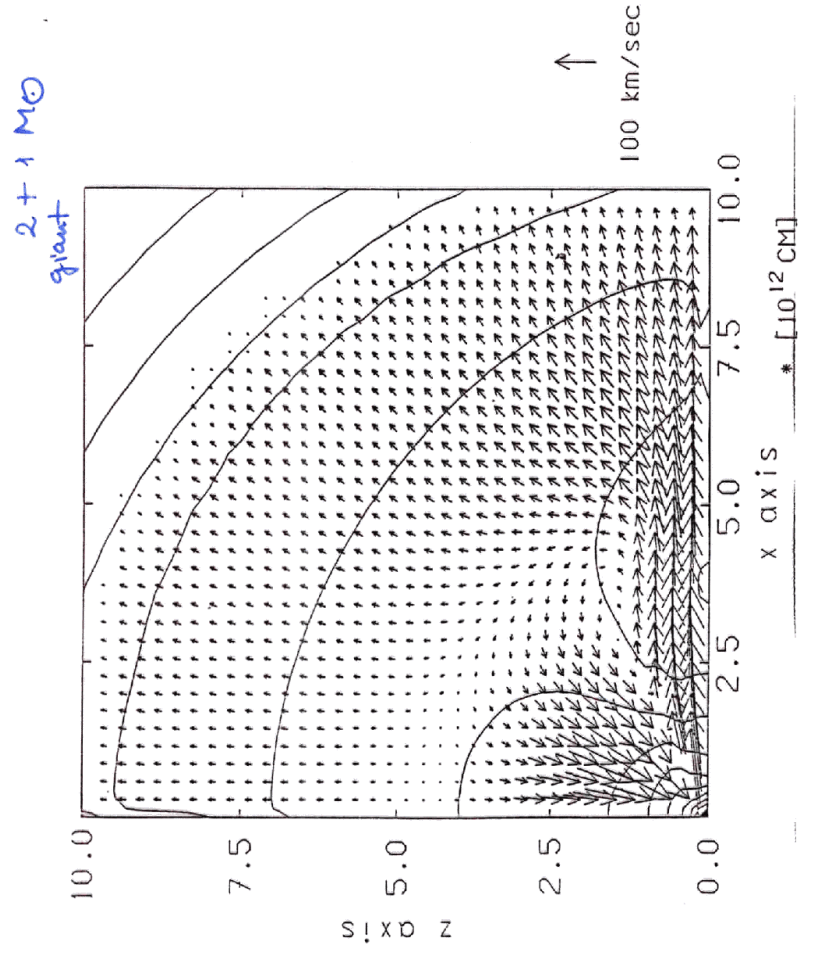
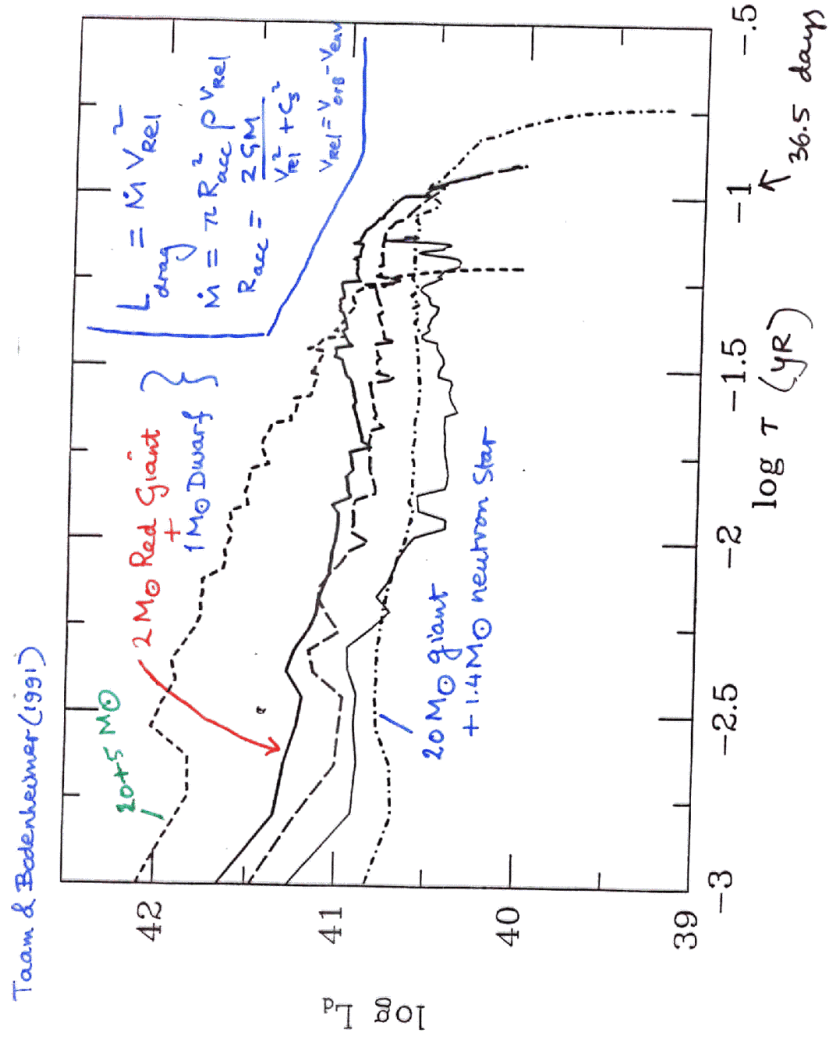


FIG. 4.—The energy dissipation rate deposited in the common envelope by the interaction of the two cores as a function of time. The energy dissipation rate is in units of 10^{40} ergs s^{-1} , and time is measured in days.

$\frac{dE}{dt} = 1.3 \times 10^{40}$ ergs/s for ≈ 50 days
 $> 10^{40}$ ergs/s for ≈ 200 days

"PLUNGE": $L_d \approx (2-3) \times 10^{40}$ ergs/s for $\approx \underline{50}^d$

Timescale of Energy-release is typically: $\tau_{dyn} = \tau_{orbit\ companion}$
 In Early Spiral-in: $\tau \approx 100^d$
 Later: $\tau \approx \tau_{orb}$ (Short)



COMPLETE MERGERS :

1) If star 1 is compact, and final $a_2 < R_{core} \rightarrow$ complete merger
 (e.g.: SHORT-PERIOD HMXBs
 Like Cen X-3 ($P_{orb} = 2.1$ days))

2) If star 1 is a Main-seq. star, but during spiral-in overflows its Roche lobe with respect to M_{core}



↓
 STAR 1 dissolves and merges with core &/or envelope of star 2

- Rapidly Rotating giant
 - Ejection of "Ring-like" structures like in SN1987A (MORRIS & Podsiadlowski 2006)
- 20 M_{\odot} + 2-5 M_{\odot} companion.

density enhancements in the outflow ←

- Low mass giants (Elliptical Galaxies)

Typical Energy Release Rates "Mergers" (SPIRAL-INs)

$$M_2 = 1.2 M_{\odot} \left. \begin{array}{l} \\ M_{core} = 0.5 M_{\odot} \end{array} \right\} A_{AGB} + M_1 = 0.6 M_{\odot}$$

Then: $a_2/a_1 \approx 1/16$ for $\alpha_{CE} = 0.5$

If $a_1 = 0.5 AU \Rightarrow a_2 = 6.7 R_{\odot}$

and $\Delta E_{released} \approx \frac{GM_{core} \cdot M_1}{6.7 R_{\odot}} \approx 2 \cdot 10^{47}$ ergs

$T_{spiral-in} \approx T_{dyn} \approx 2$ months (half of it used for envelope ejection)

$L_{peak, thermal} \approx 2 \cdot 10^{40}$ ergs/s for ~ 60 d

- ! This is, in fact, independent of value of M_1 :

If M_1 twice as small $\Rightarrow a_2$ twice as small $\Rightarrow \Delta E_{released}$ is same; T also same.

- High-mass giants :

$L_{peak} \approx 10^{41} - 10^{42}$ (at most) ergs/s
 on similar timescale

Expected Frequency of Events
for steady Star Formation Rate
and main-sequence compan. mergers.

ASSUME:

- 25 per cent of stars in binaries
with a in range $4 R_{\odot}$ to 5 AU
- $f(\log a) = \text{const.} = 0.1$ for
 $\log a (\text{AU}) = -1.7$ to 0.7
 $\uparrow 4 R_{\odot}$ $\downarrow 5 \text{ AU}$
- $g(q = M_1/M_2) = \text{const.}$ (uncertain!)
- inspiral companions have $q \leq 0.25$
- Birthrate $\frac{dN(M)}{dt} = \psi(M) = A \cdot M^{-2.5}$
- Only stars with convective envelopes
have dyn. timescale spiral-in

RESULTING RATES:

- $M_2 \geq 10 M_{\odot} \Rightarrow R_{\text{mergers}} \approx 0.02 R_{\text{SN core coll.}}$
- $M_2 = 4-10 M_{\odot} : R_{\text{mergers}} = 0.08 R_{\text{SN cc}}$
- $M_2 = 1-4 M_{\odot} : R_{\text{mergers}} \approx R_{\text{SN}}$
(Largest contrib. from close to $M_2 \approx 1$)
(assumed: All stars $\geq 10 M_{\odot}$ go to
core collapse SN).