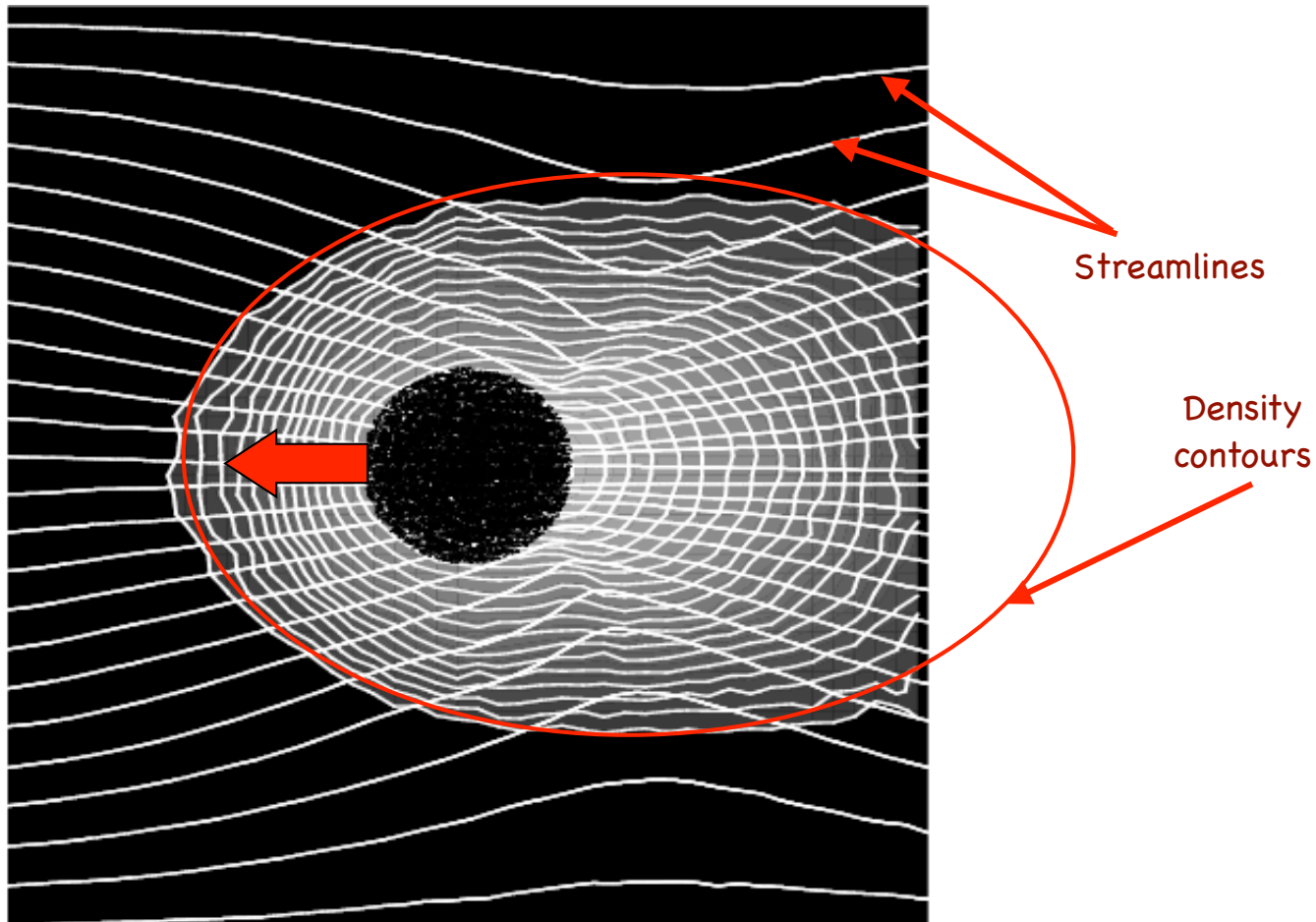


# Gaseous Dynamical Friction in Presence of BH Radiative Feedback

Cole Miller and Tamara Bogdanović

# MBHs transported by dynamical friction

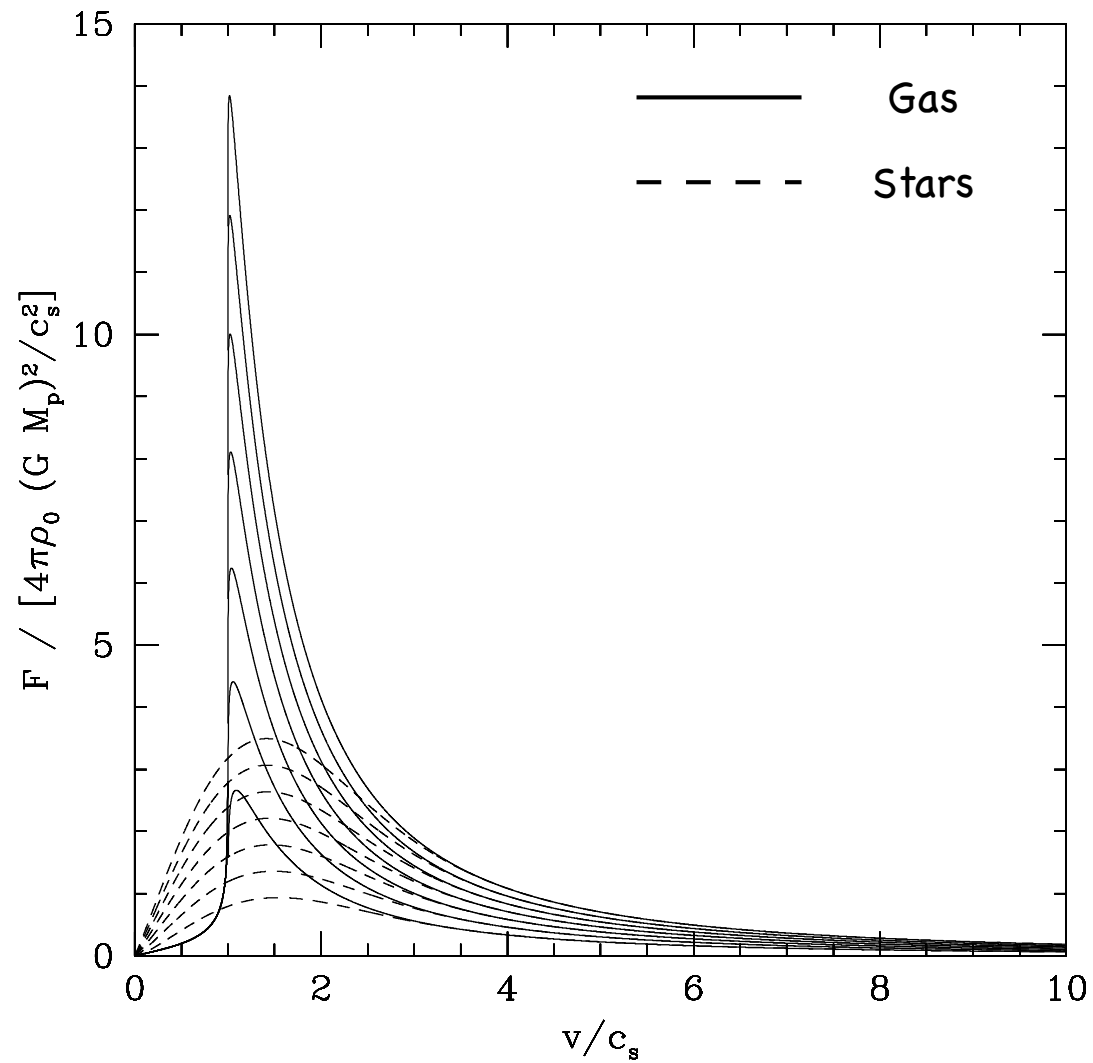
Separation (1000 → 1) pc    Time scale: Million - Billion years



$$F_{\text{DF}} = -I \times 4\pi(GM_p)^2 \rho_0 / V^2$$

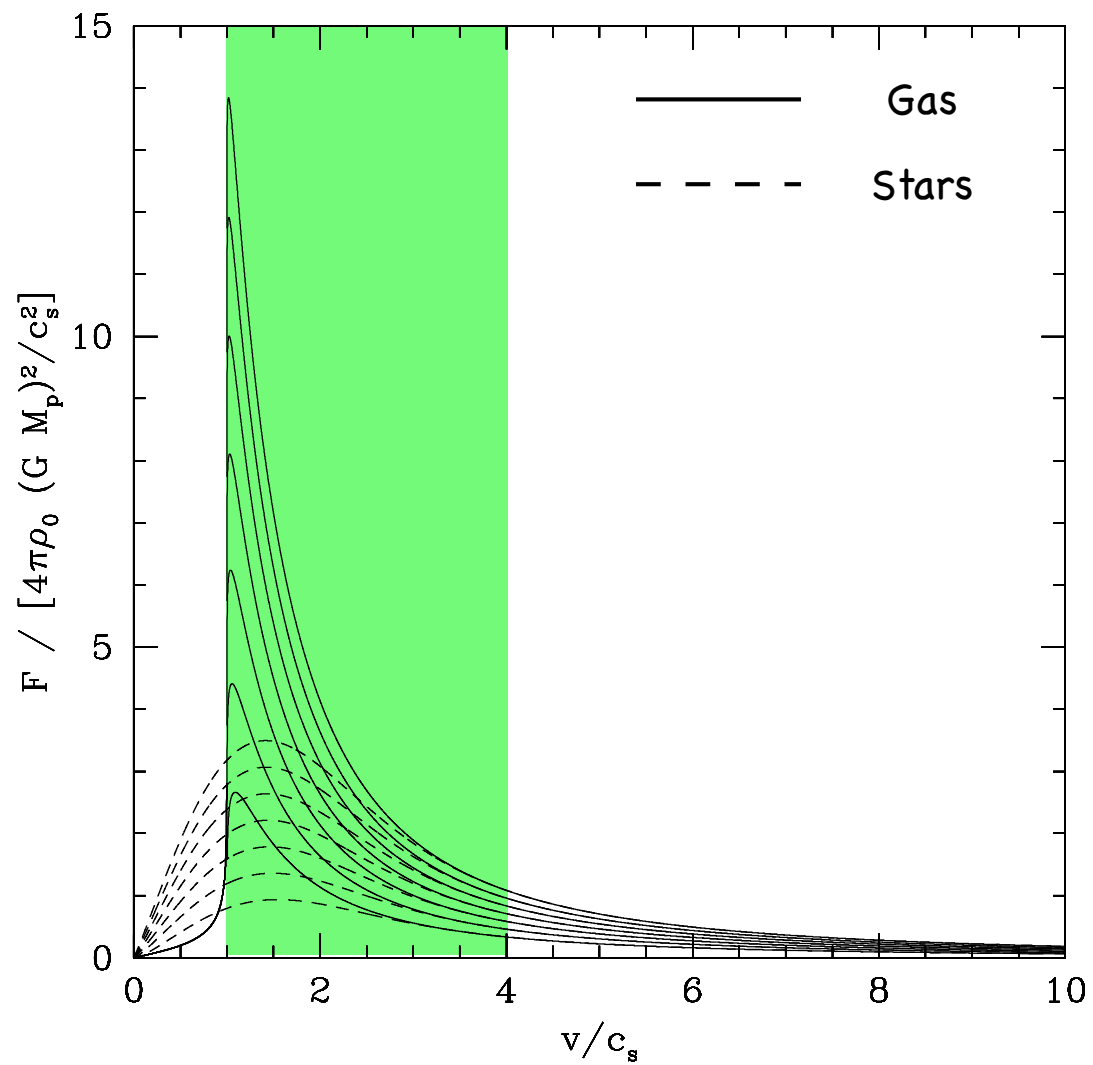
(Chandrasekhar 43, Ostriker 99,  
image from Semelin & Combes 02)

# Stellar vs. gaseous dynamical friction



(Ostriker 99)

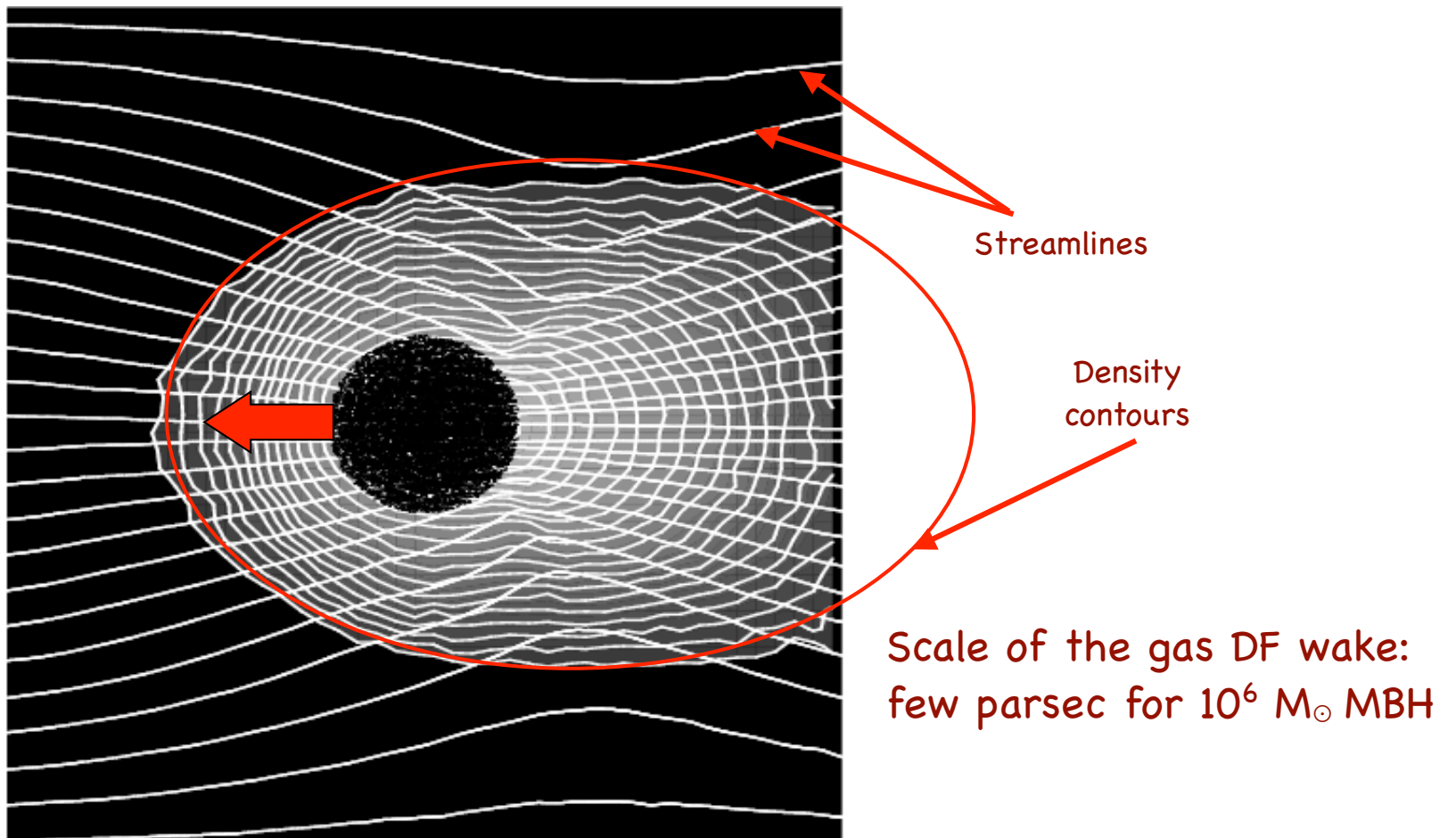
# Stellar vs. gaseous dynamical friction



(Ostriker 99)

# MBHs transported by dynamical friction

Separation (1000 → 1) pc    Time scale: Million - Billion years



$$F_{\text{DF}} = -I \times 4\pi(GM_p)^2 \rho_0 / V^2$$

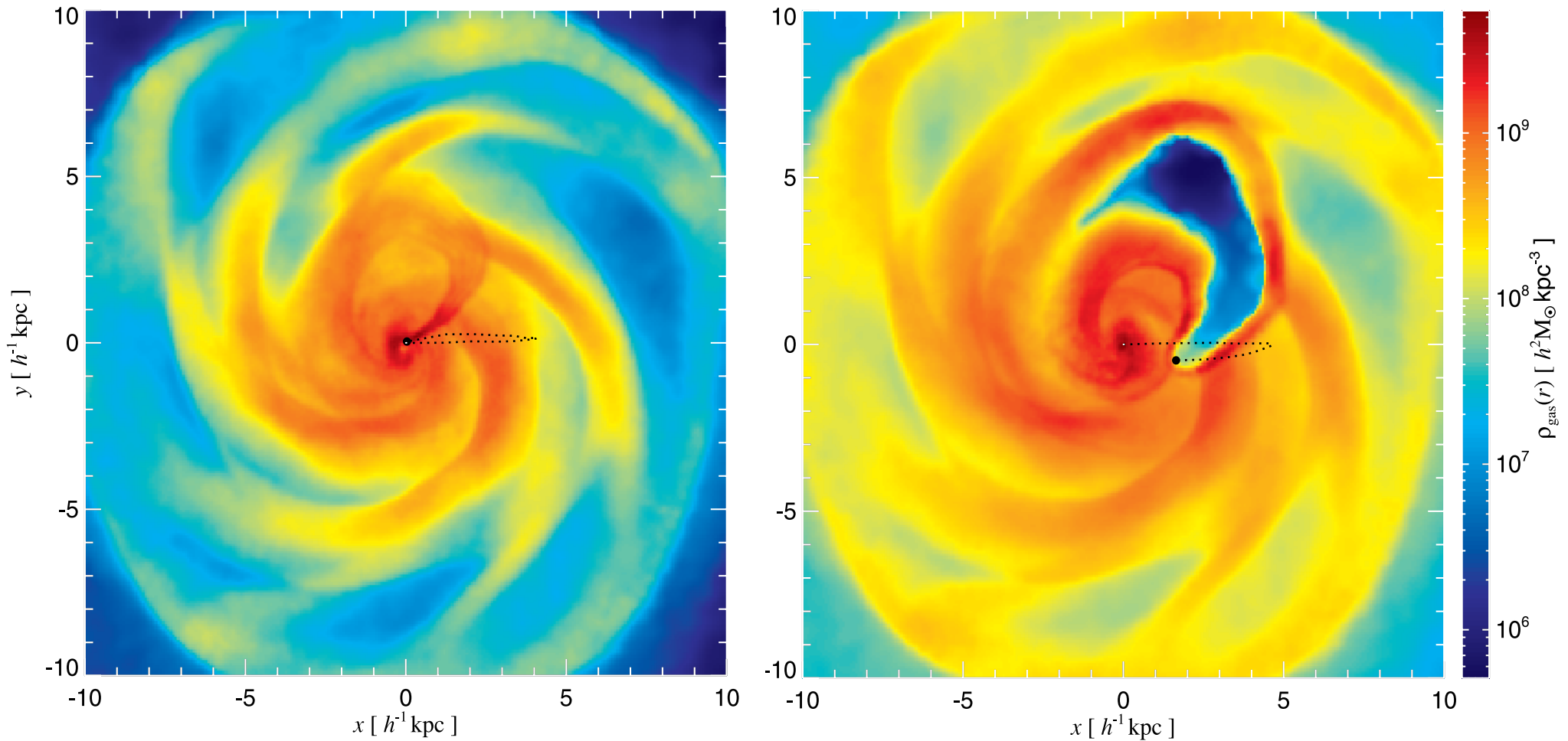
(Chandrasekhar 43, Ostriker 99,  
image from Semelin & Combes 02)

# Wake evacuation by MBH feedback

Without MBH feedback

Mass density

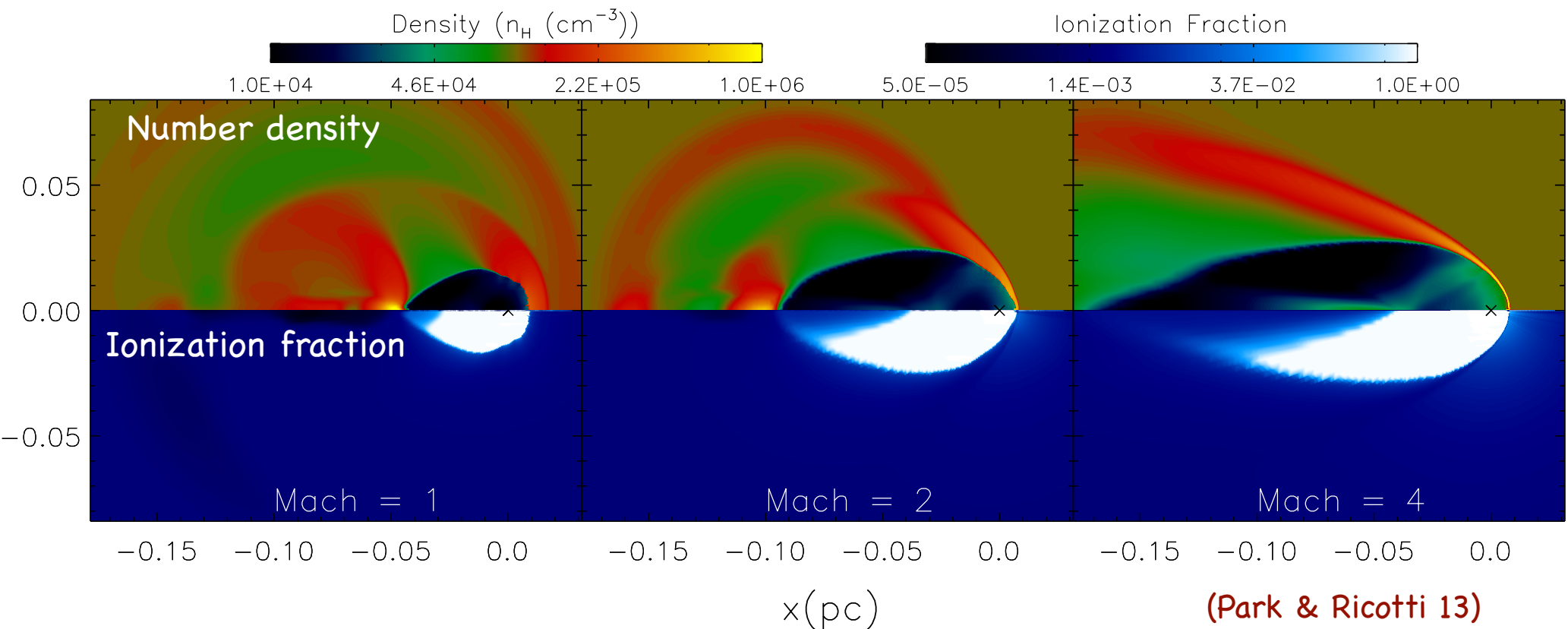
With MBH feedback



(figures from Sijacki+ 11; Souza-Lima+ 17)

# Wake evacuation by accreting BHs (big and small)

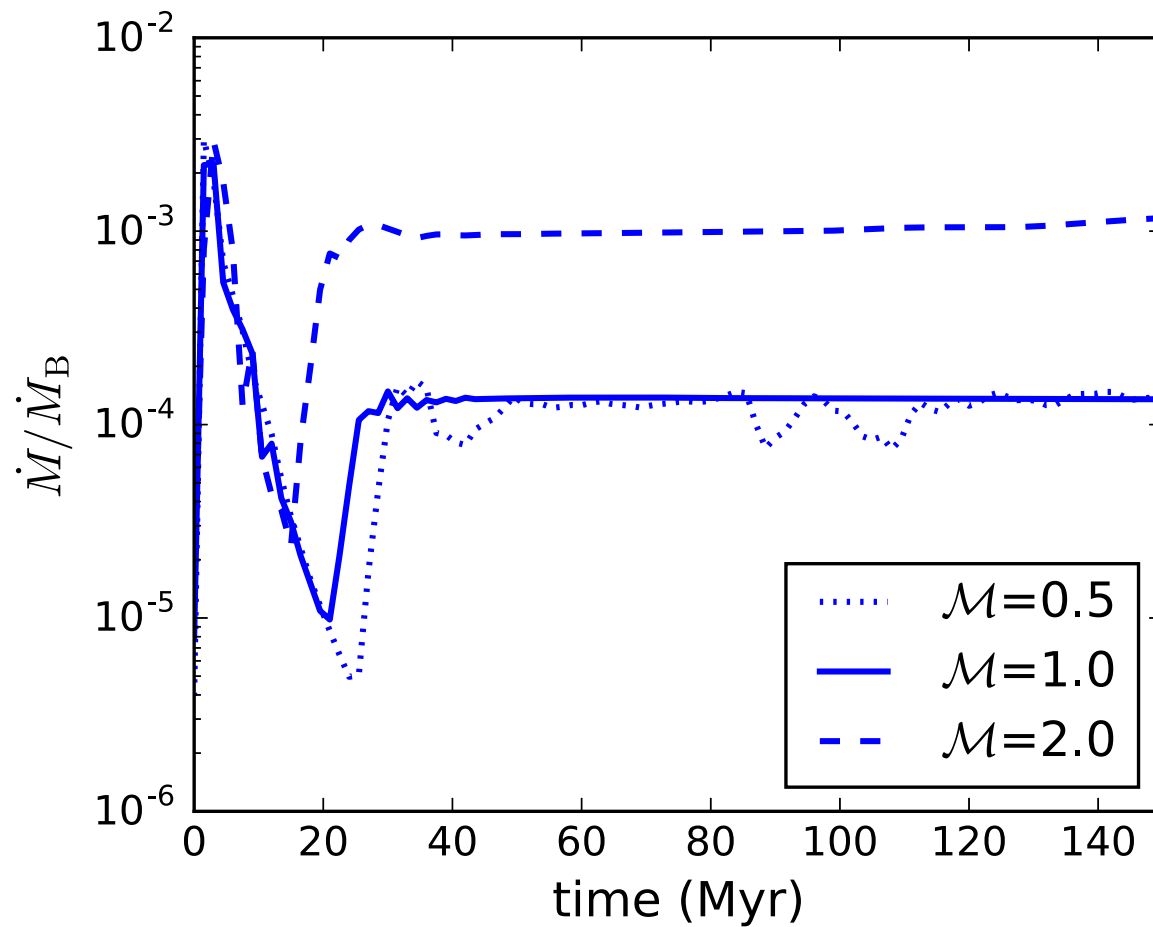
$M_{\text{bh}} = 100M_{\odot}$  BH, immersed in  $T_{\infty}=10^4$  K gas w/  $n_{\infty} = 10^5 \text{ cm}^{-3}$



Density of the shell increases as  $\propto \mathcal{M}^2 \Rightarrow$  accretion rate is suppressed

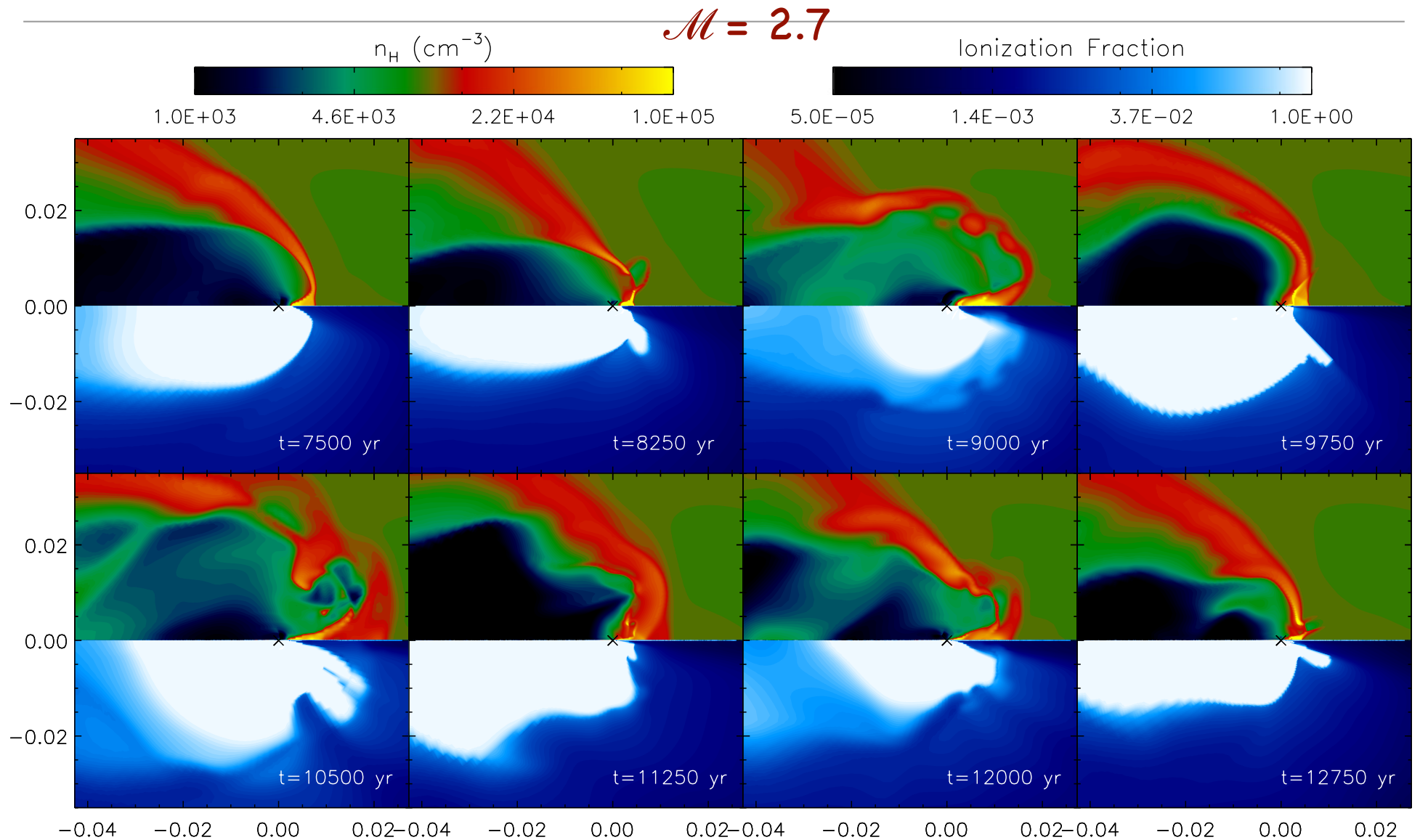
# Radiation mediated accretion ( $1 < \mathcal{M} < 4$ )

$$\dot{M} \approx 1.2 \times 10^{-2} \mathcal{M}^2 \dot{M}_B \left( \frac{T_\infty}{T_{\text{in}}} \right)^{5/2}$$





# A) Wake evacuation is prevented for fast perturbers

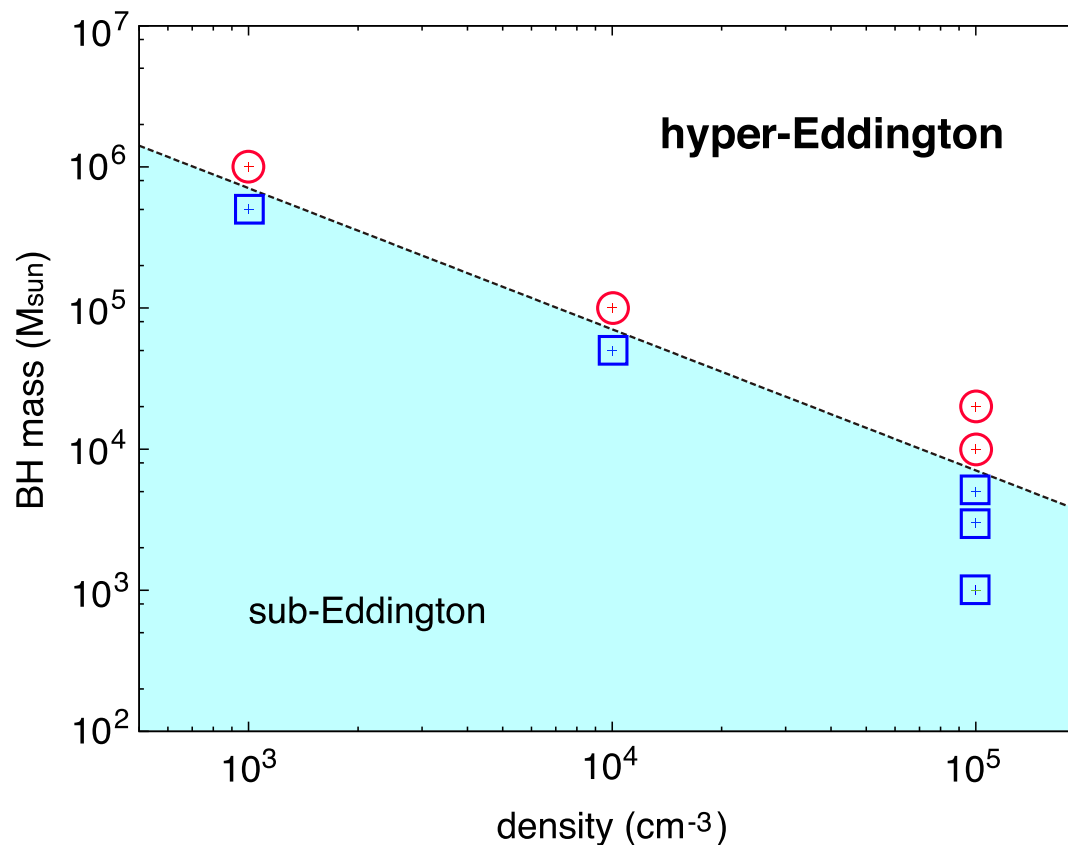


$M_{\text{bh}} = 100M_{\odot}$  ,  $T_{\infty} = 10^4 \text{ K}$  ,  $n_{\infty} = 10^5 \text{ cm}^{-3}$

(Park & Ricotti 13)

## B) Wake evacuation is prevented for super-Eddington accretors

Spherically symmetric accretion flows transition to unimpeded hyper-Eddington accretion when  $R_{\text{HII}} < R_{\text{B}}$  (Park+ 14; figure from Inayoshi+ 16)

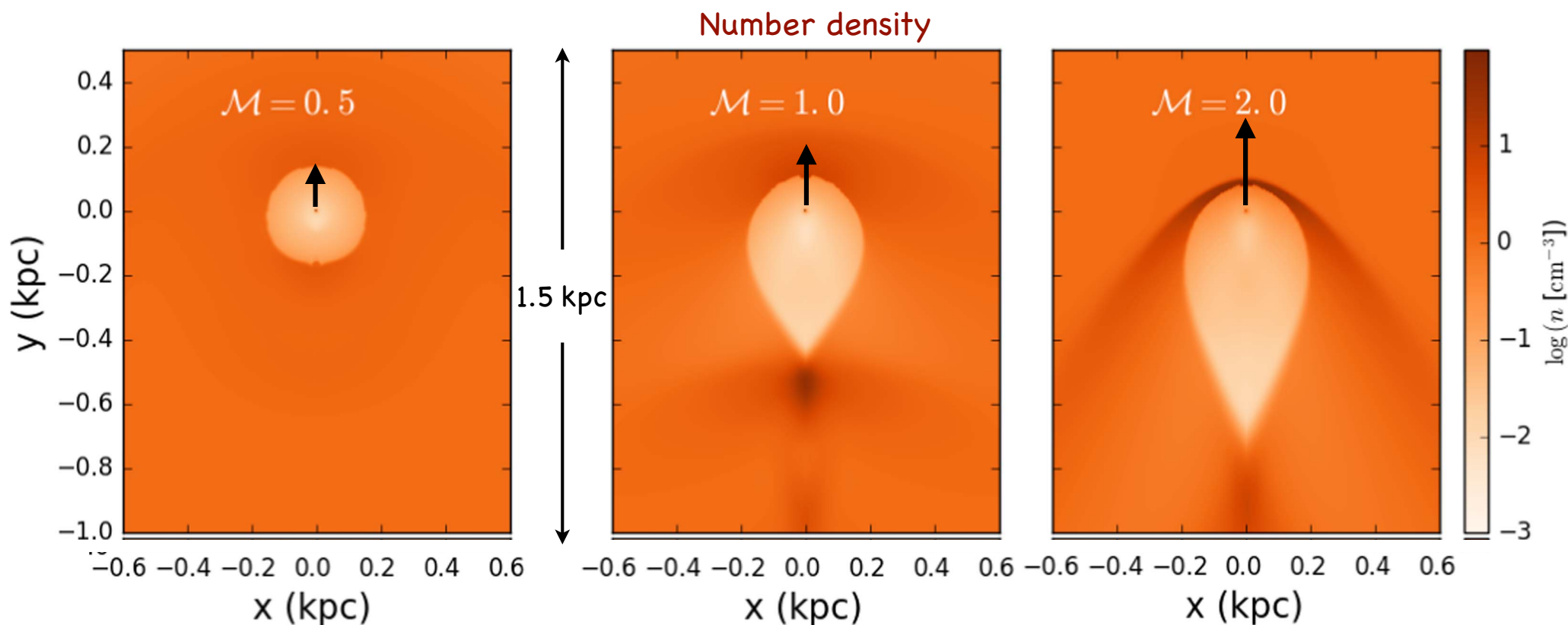


This happens when:

$$\dot{M}_{\text{B}} \gtrsim 5000 \dot{M}_{\text{Edd}} \quad \longleftrightarrow \quad \left( \frac{n_{\infty}}{10^5 \text{ cm}^{-3}} \right) \gtrsim \left( \frac{M_{\text{BH}}}{10^4 M_{\odot}} \right)^{-1} \left( \frac{T_{\infty}}{10^4 \text{ K}} \right)^{3/2}$$

# Efficiency of Gas Dynamical Friction in Presence of Radiative Feedback

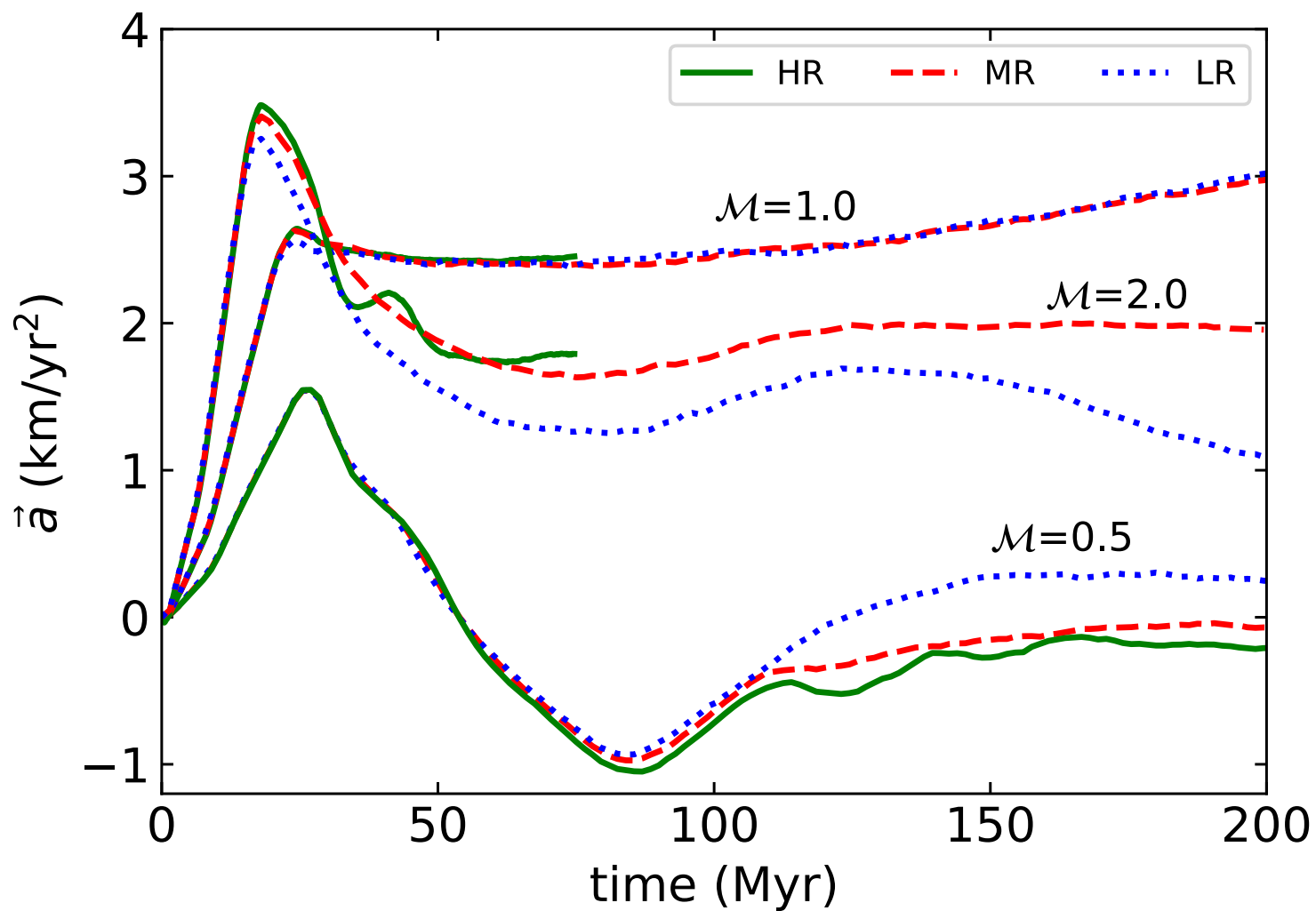
2D RHD simulations of  $10^6 M_\odot$  MBH moving through a uniform, neutral background medium  
- KwangHo Park & TB 17, 19 (the latter in prep.) -



wake evacuated when

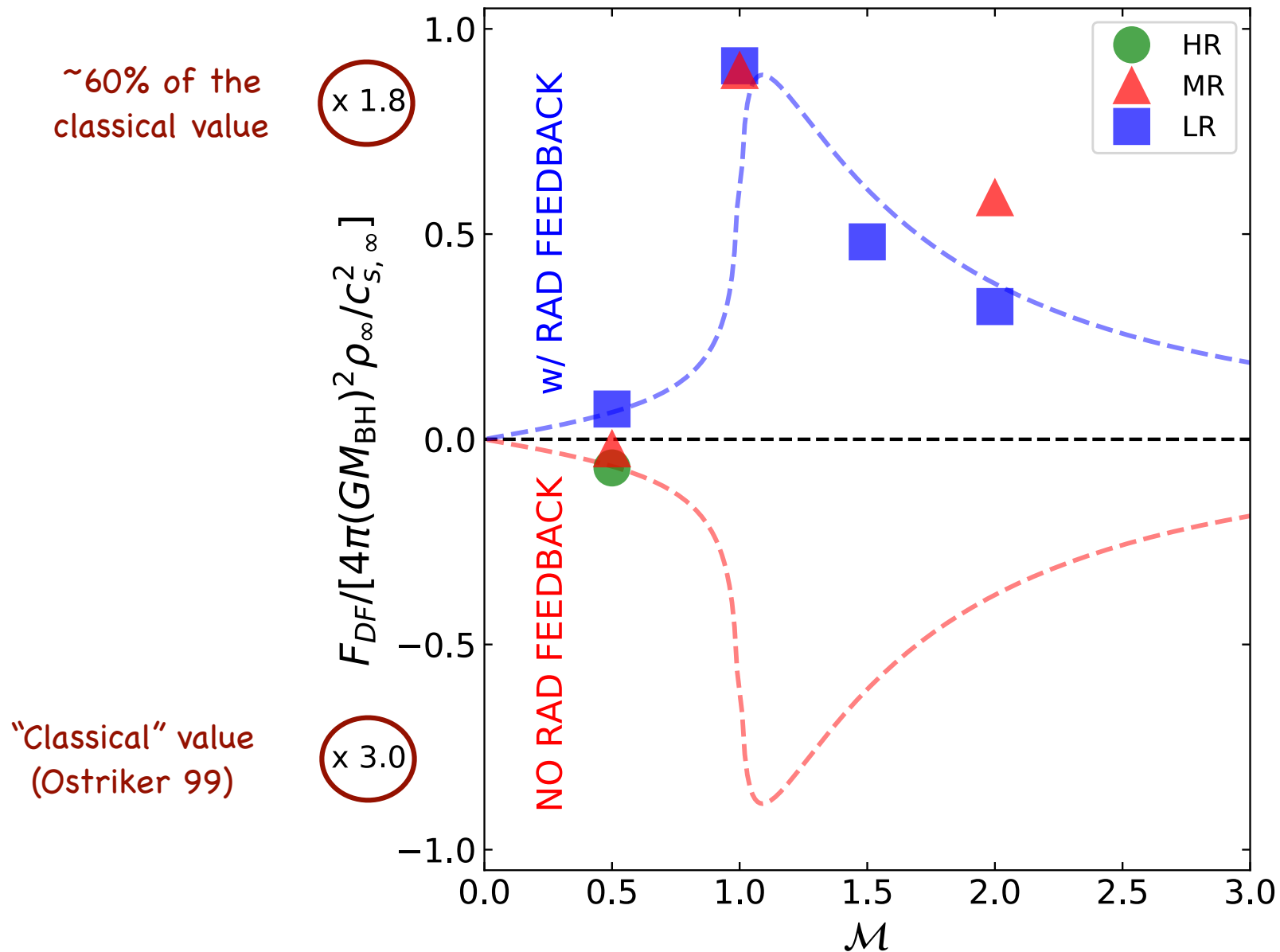
$$\mathcal{M} < 4 \text{ and } (1 + \mathcal{M}^2) M_{\text{BH}} n_\infty < 10^9 M_\odot \text{ cm}^{-3}$$

# MBH Acceleration as a Function of Time



(Park & TB 17, 19)

# MBH Acceleration as a Function of the Mach Number



# Conclusions



- Ionizing radiation from MBHs gives rise to negative gas DF for a range of physical scenarios. Stellar DF may still operate unaffected but would have to work harder.
- The effect is more severe at the low mass end of BH spectrum => BHs with masses  $<10^7 M_{\odot}$  have fewer means to reach the centers of merged galaxies.

$$\mathcal{M} < 4 \text{ and } (1 + \mathcal{M}^2) M_{\text{BH}} n_{\infty} < 10^9 M_{\odot} \text{ cm}^{-3}$$

- Prescription for sub-resolution model of gas DF for large-scale simulations.
- **Not taken into account:** gas inhomogeneities, anisotropic outflows, curved BH trajectories, other sources of shocks / radiation, magnetic fields, etc.