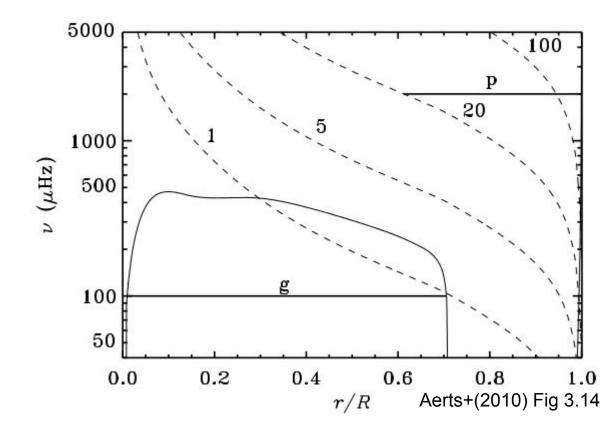
Fluid Instabilities in Neutron Star Mergers

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Merging Visions
Kavli Institute for Theoretical Physics
June 26, 2019

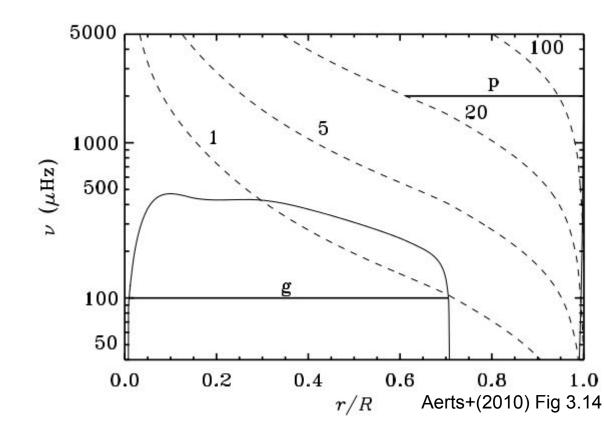
basic overview of modes

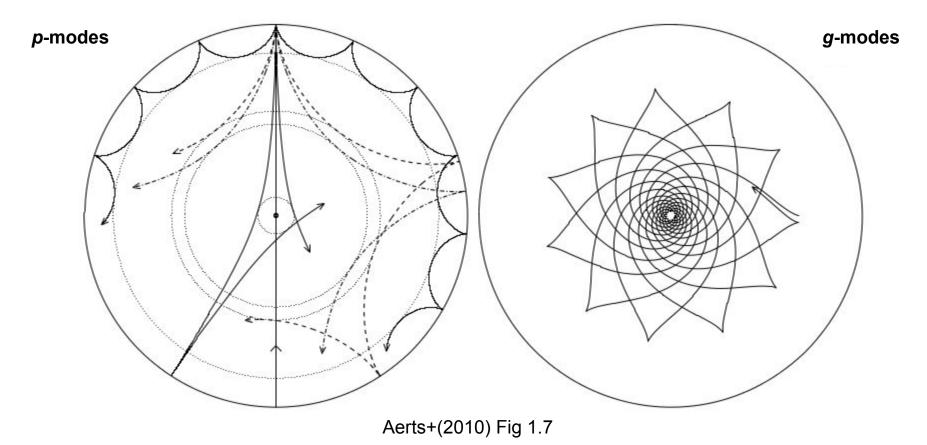
- p-modes
 - high frequencies
 - ω∝n (radial order)



basic overview of modes

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- g-modes
 - low frequencies
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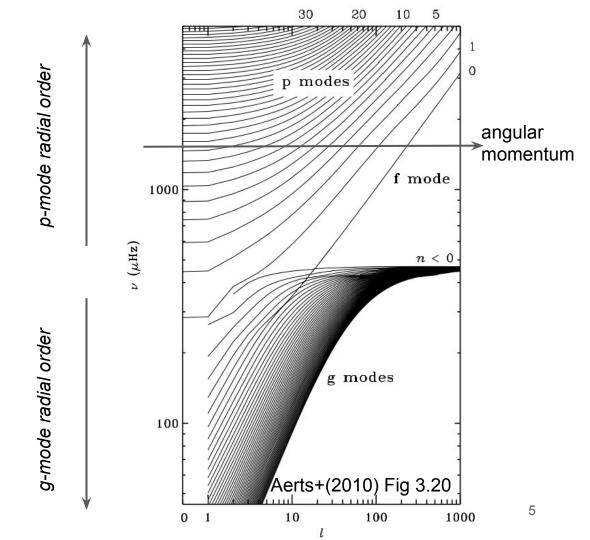




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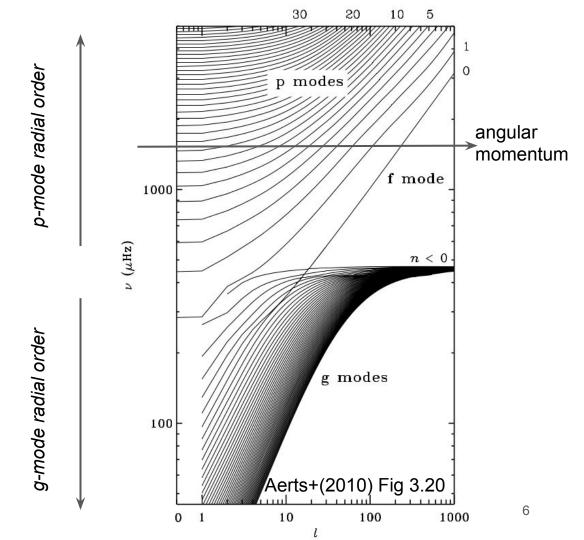
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 - rotational/inertial modes
 - o often "mixed" with *g*-modes

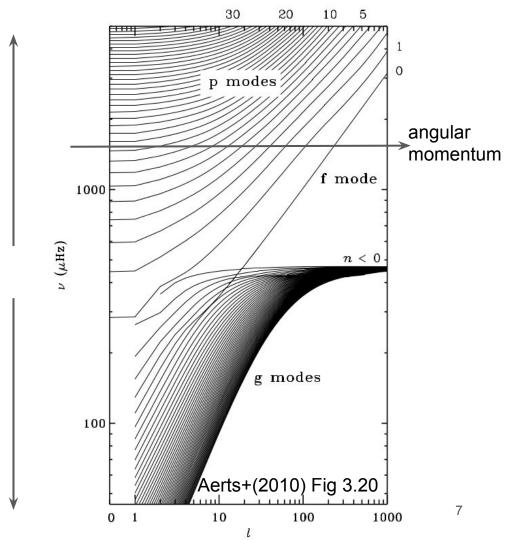


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 - o often "mixed" with *g*-modes
- w-modes
 - metric oscillations
 - o only present in relativistic stars

p-mode radial order

g-mode radial order



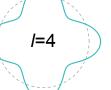
multi-mode interactions

equations of motion for mode amplitudes can be described via a Galerkin decomposition

based on the linear eigenmodes.









multi-mode interactions

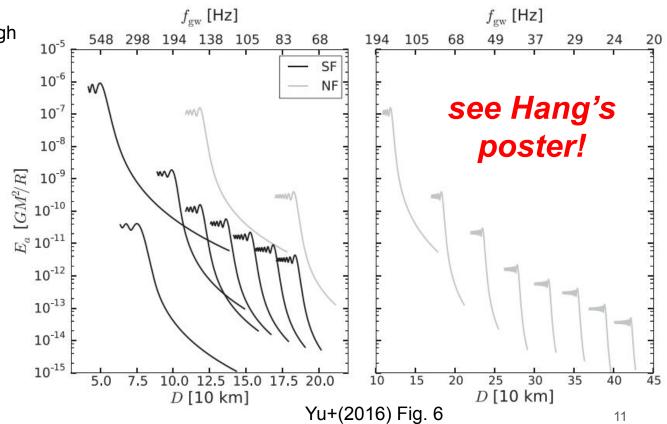
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linear dynamical tides and (resonant) nonlinear dynamical tides

GW-driven inspiral sweeps through resonances too fast to efficiently transfer energy through resonant interactions.



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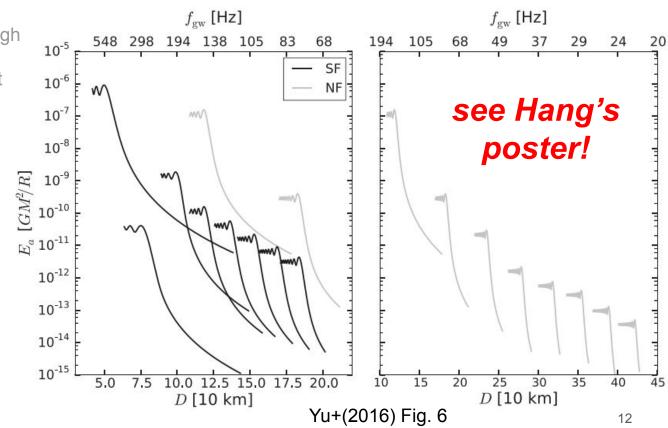
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high-frequencies

- stronger coupling
- less time near resonance

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- weaker coupling
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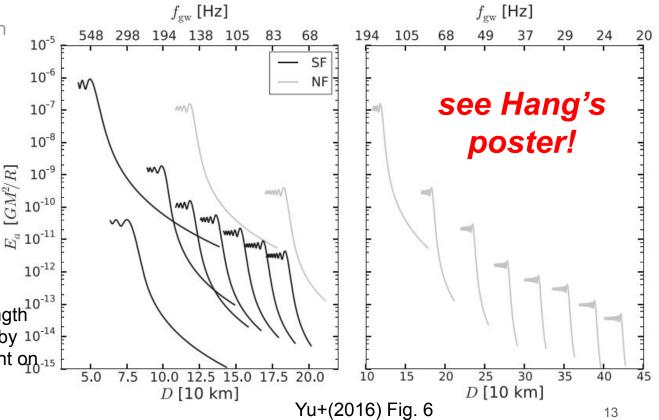
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large spins may shift long-wavelength modes to lower frequencies, thereby 10⁻¹⁴ increasing the amount of time spent on resonance (Ho+1999).



(nonresonant) *p-g* secular instabilities

• instability of the *linear tidal bulge* coupled to a *high-frequency p-mode* and a *low-frequency g-mode* (Weinberg+(2012)).

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 - o coupling coefficients in the Galerkin decomposition can be large.
 - o could be important for compact systems containing either Neutron Stars or White Dwarfs.

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 - o coupling coefficients in the Galerkin decomposition can be large.
 - o could be important for compact systems containing *either* Neutron Stars or White Dwarfs.
- *nonresonant* and active whenever the linear tidal perturbation is above some threshold.
 - equivalently, when orbital separation is below some threshold or when orbital frequency is above some threshold.
 - 4-mode couplings are also important (Venumadhav+(2016)) and can dynamically cancel part of the instability, but an instability still exists for dynamical tidal fields.
 - finite-frequency and other non-adiabatic effects on mode shapes spoil the cancellation between 3- and 4-mode interactions, resulting in smaller but still possibly relevant growth timescales (Weinberg (2016)).
- difficult to simulate
 - spatial grid required to resolve high-order g-modes is prohibitively expensive.
 - larger number of relevant coupled modes makes Galerkin amplitude equations difficult to simulate.

(nonresonant) p-g secular instabilities

- phenomenological model
 - o dissipation by *p-g* instability modifies orbital evolution and Gravitational-Wave phase (Essick+(2016)).

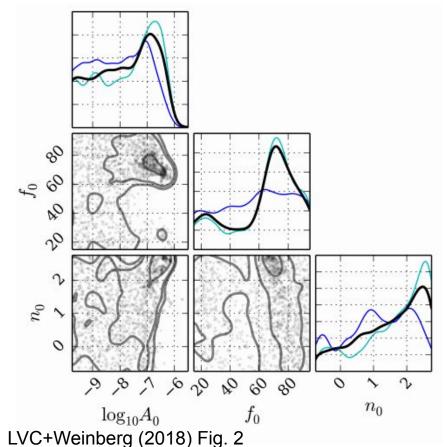
 A_0 : overall amplitude of induced phase shift

 f_0 : saturation frequency ~ instability threshold assuming modes grow quickly

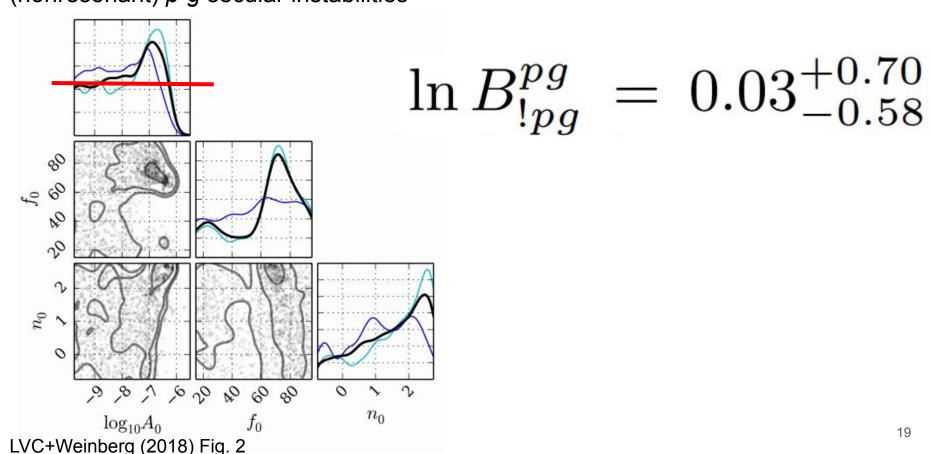
 n_0 : scaling of energy dissipated as a function of frequency

worked example: *p-g* instabilities with GW170817

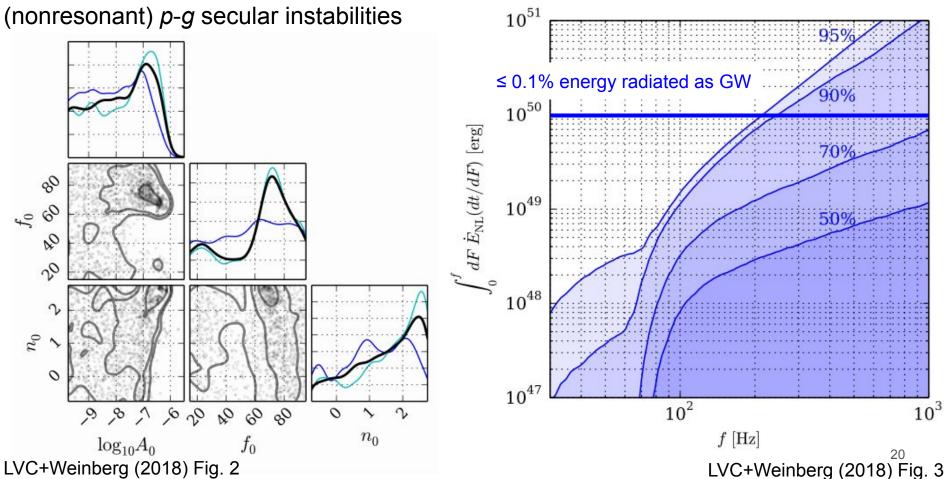
(nonresonant) p-g secular instabilities



worked example: *p-g* instabilities with GW170817 (nonresonant) p-g secular instabilities



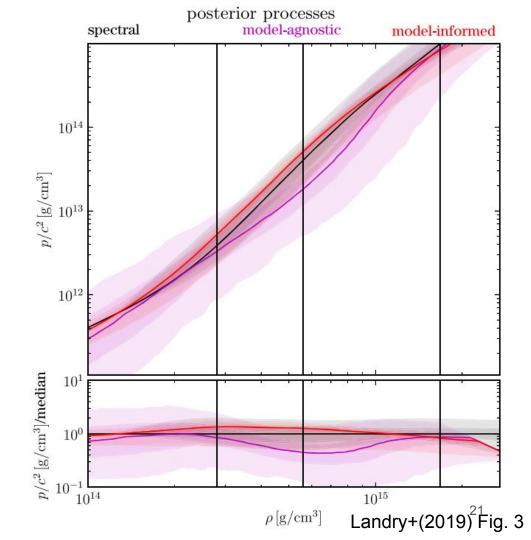
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Can Equation of State Constraints help?

linear resonant tides

- yes!
- knowledge of the EOS specifies the mode spectra and shapes.
- but these effects are likely to be negligible anyway...



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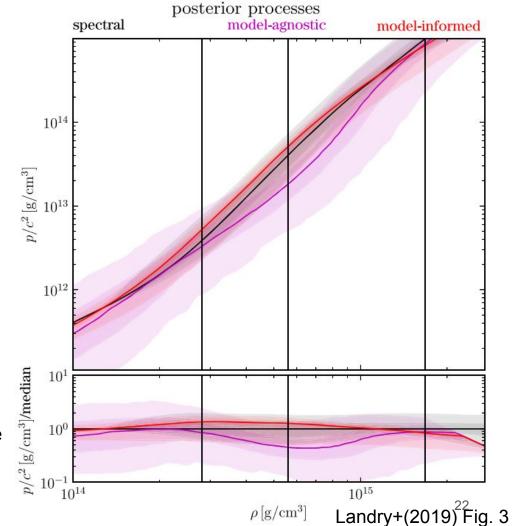
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- yes!
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nonlinear (non)resonant tides, CFS instabilities

- not clear... (e.g., Zhou+(2017))
- uncertain physics within NS core
 - o damping mechanisms
 - saturation mechanism
- difficult calculations

Observations and phenomenological models may be the fastest way to constrain these effects...



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