



Tunneling density of states of graphene

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SPP 1243





E. Mariani

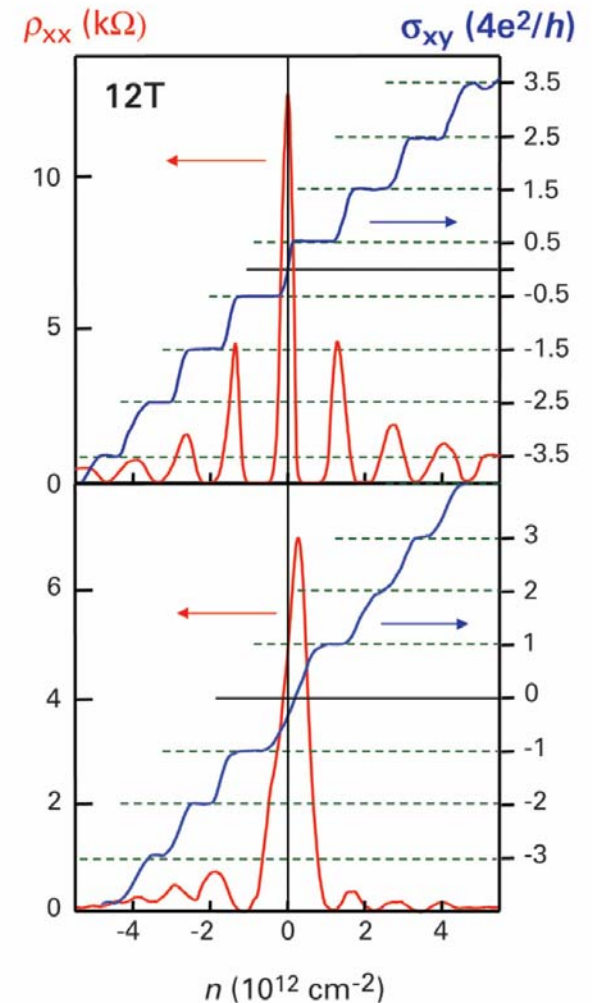
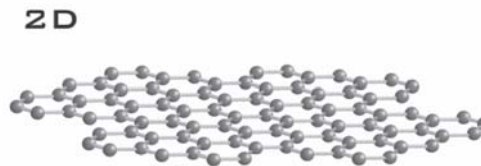
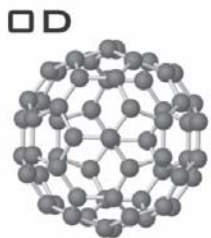
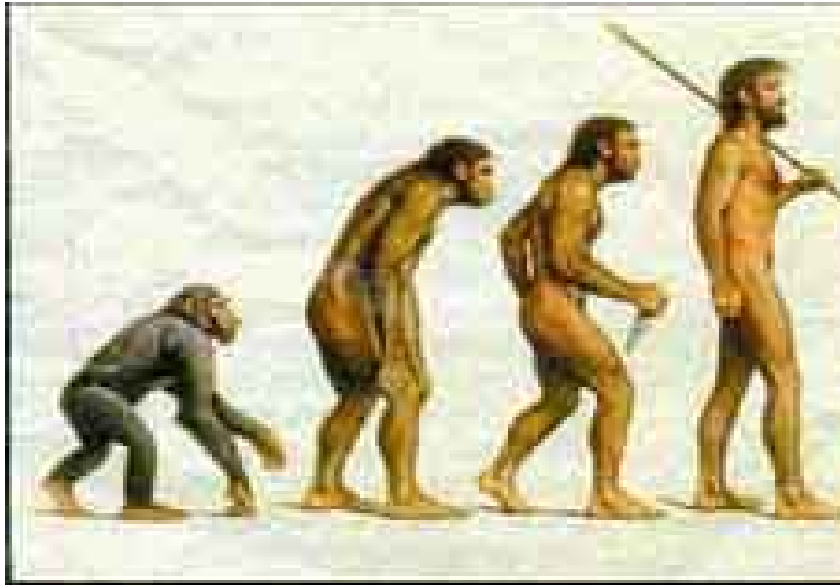


L. Glazman



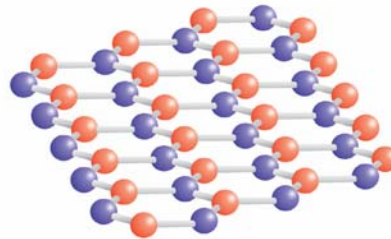
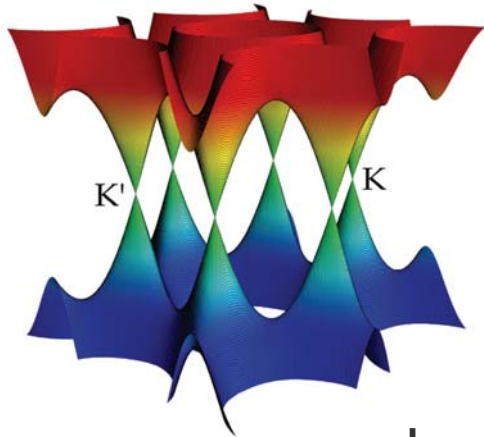
A. Kamenev

Evolution



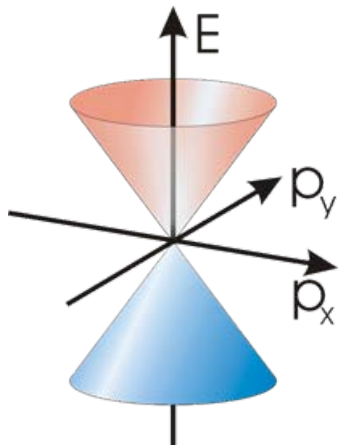
Novosolov *et al.*, Nature 2005
Zhang *et al.*, Nature 2005

Tight-binding description



$$H = -t \sum_{\langle ij \rangle} c_i^\dagger c_j$$

Low-energy Dirac Hamiltonian



Dirac points

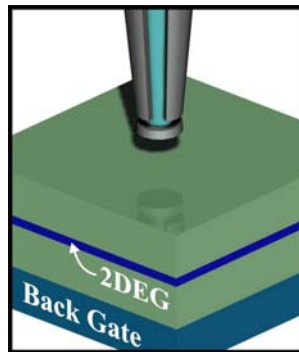
$$H = v \boldsymbol{\Sigma} \cdot \mathbf{p}$$

$$H = v \begin{bmatrix} 0 & p_x - ip_y & 0 & 0 \\ p_x + ip_y & 0 & 0 & 0 \\ 0 & 0 & 0 & -(p_x - ip_y) \\ 0 & 0 & -(p_x + ip_y) & 0 \end{bmatrix} \begin{bmatrix} A_1 \\ B_1 \\ B_2 \\ A_2 \end{bmatrix}$$



Graphene is exposed at surface:

- directly accessible by local probes
- enhanced spatial resolution



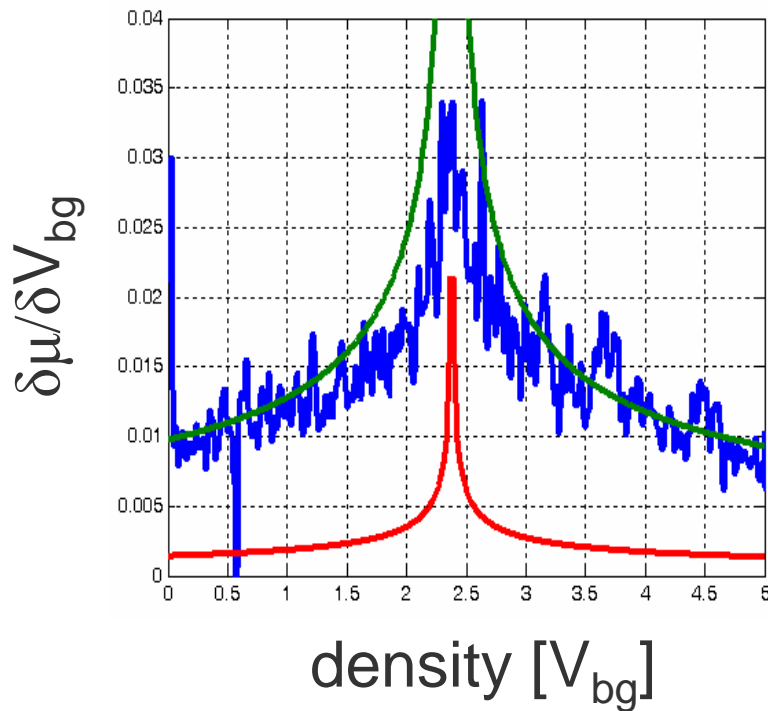
Scanning SET
(A. Yacoby et al.)



STM
(e.g. talk by E. Andrei
next week)

Scanning SET:
(A. Yacoby et al.)

(local) compressibility



$$\kappa^{-1} \sim \left[\frac{d\mu}{dn} \right]$$

Hartree-Fock:

$$\mu = \hbar v (2\pi n_0)^{1/2} - \sqrt{\frac{2}{\pi}} \frac{e^2}{2\pi\epsilon_{\text{eff}}} n^{1/2} + E_{\text{corr}}$$

measured from
Dirac pt

all electrons

see MacDonald et al. (2007)

STM: (local) tunneling density of states



STM

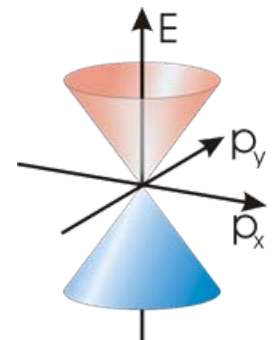
Dirac equation describes
coarse grained local TDOS

Dirac dispersion

$$E_p = v p$$



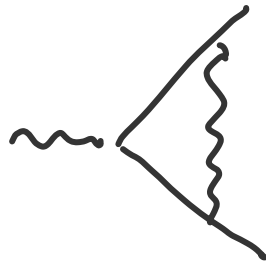
$$\nu(E) = \frac{|E|}{2\pi\hbar^2 v^2}$$



“Fine structure constant” $g = \frac{e^2}{\hbar v} \sim 1$



$$E(p) = [v + (e^2 / 4) \ln(D / v p)] p$$



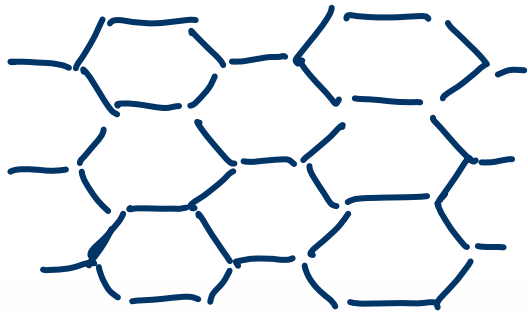
$$\delta e = 0$$

TDOS

$$\nu(E) = \frac{|E|}{2\pi\hbar^2 [v + (e^2 / 4) \ln(D / v p)]^2}$$

J. González et al., Nucl. Phys. B 424, 595 (1994)

Tight binding



$$H = -t \sum_{\langle ij \rangle} c_i^\dagger c_j$$

- site energies
- random hopping
- lattice defects

Peres et al. PRB 2006

Dirac theory

general T-invariant local operator:

$$\delta H = u_0(\mathbf{r})\mathbf{1} + \sum_{ij} u_{ij}(\mathbf{r})\Lambda_i\Sigma_j$$

where

$$\Lambda_{x,y} = \pi_{x,y} \otimes \sigma_z \quad \Lambda_z = \pi_z \otimes \sigma_0$$

and

Λ_i, Σ_j odd under T-reversal

McCann et al. PRL 2006

	scalar potential	(fictitious) gauge field
intra- valley	u_0, u_{zz}	u_{zx}, u_{zy} (abelian)
inter- valley	u_{xz}, u_{yz}	$u_{xx}, u_{xy}, u_{yx}, u_{yy}$ (nonabelian)

McCann et al. PRL 2006

Neglecting intervalley scattering

$$H = v \boldsymbol{\Sigma} \cdot (\mathbf{p} + \mathbf{a}(\mathbf{r})) + V(\mathbf{r})$$

fictitious gauge field
due to random hopping

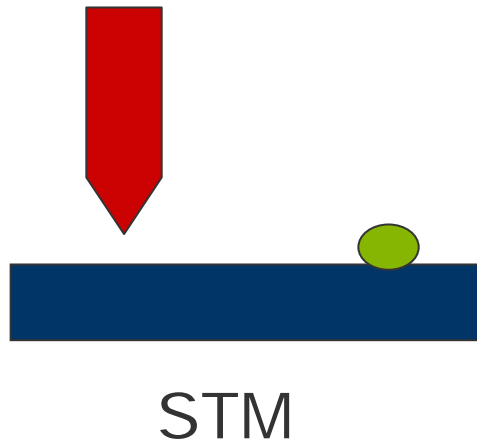
random site energies

Time-reversal symmetry:

$$\mathbf{a}(\mathbf{r}) = \Pi_z \otimes \mathbf{1} \begin{pmatrix} u_{zx}(r) \\ u_{zy}(r) \end{pmatrix}$$



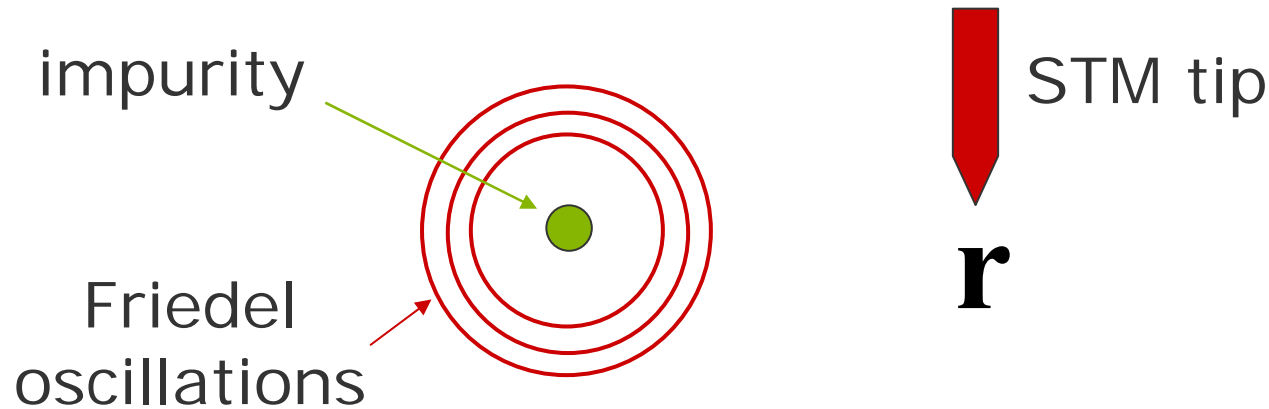
Weak disorder (ballistic regime):



chiral symmetry
↓
no backscattering

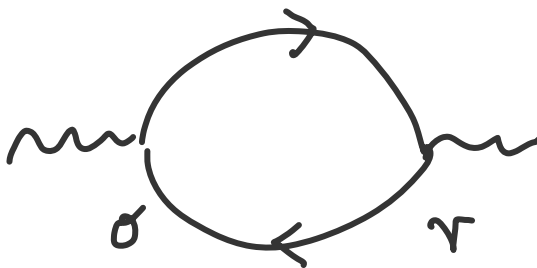
$$G(\mathbf{r}, \mathbf{r}; E) = G_0(\mathbf{r}, \mathbf{r}; E) + G_0(\mathbf{r}, \mathbf{0}; E) \hat{u} G_0(\mathbf{0}, \mathbf{r}; E) + \dots$$

$$\delta \nu(E) \sim \text{tr} \delta G(\mathbf{r}, \mathbf{r}; E) \sim \text{tr} [s_r \hat{u} s_{-r}] \sim 0$$



- Fate of zero-bias anomaly as E_F approaches Dirac pt ?
 - wavelength of Friedel oscillations diverge
 - strength of electron-electron interaction increases

- Probe of fictitious gauge field ?

$$\delta\rho(\mathbf{r}) \sim \text{bubble diagram} \sim \int d\varepsilon \text{tr} G(\mathbf{r}, \mathbf{0}; \varepsilon) \hat{u} G(\mathbf{0}, \mathbf{r}; \varepsilon)$$


Green function:

$$\mathbf{G}_\varepsilon^R(\mathbf{p}) = \frac{\hat{s}_{-\mathbf{p}}}{\varepsilon + vp + i\eta} + \frac{\hat{s}_{\mathbf{p}}}{\varepsilon - vp + i\eta}$$

$$\mathbf{G}_\varepsilon^R(\mathbf{r}, 0) \simeq -\frac{e^{i\pi/4} p_\varepsilon}{\sqrt{2\pi v}} \frac{e^{ip_\varepsilon r}}{\sqrt{p_\varepsilon r}} \left[\hat{s}_{\mathbf{r}} + \frac{i}{4p_\varepsilon r} \boldsymbol{\Sigma} \cdot \frac{\mathbf{r}}{r} \right]$$

extra 1/r

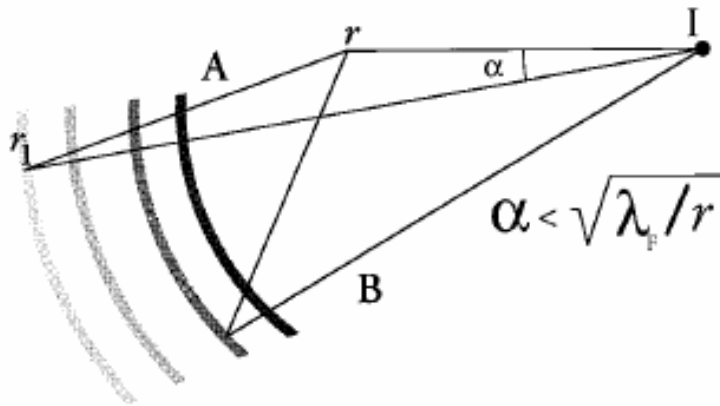


$\delta\rho(\mathbf{r}) \sim \frac{1}{r^3}$ as opposed to $1/r^2$ in usual 2d electron gas

Cheianov *et al.* PRL 2006

Ballistic regime:
 perturbation theory in disorder and interactions

Rudin et al. PRB 1997



$$\hat{V}_F(\mathbf{r}_1, \mathbf{r}_2) = \frac{1}{2} V(\mathbf{r}_1 - \mathbf{r}_2) \delta\rho(\mathbf{r}_1, \mathbf{r}_2).$$

$$\hat{H}_{HF}(\mathbf{r}_1, \mathbf{r}_2) = \hat{V}_H(\mathbf{r}_1) \delta(\mathbf{r}_1 - \mathbf{r}_2) - \hat{V}_F(\mathbf{r}_1, \mathbf{r}_2)$$

$$\delta \mathbf{G}_\epsilon^R(\mathbf{r}, \mathbf{r}) = \int d\mathbf{r}_1 d\mathbf{r}_2 \mathbf{G}_\epsilon^R(\mathbf{r}, 0) \hat{u} \mathbf{G}_\epsilon^R(0, \mathbf{r}_1) \hat{H}_{HF}(\mathbf{r}_1, \mathbf{r}_2) \mathbf{G}_\epsilon^R(\mathbf{r}_2, \mathbf{r})$$

Global TDOS (finite density n_i of impurities):

$$\frac{\delta\nu_\omega}{\nu_F} \simeq \frac{4\nu_F^2 n_i}{\sqrt{\pi} k_F^2} \text{Tr} \left[2\hat{u}^2 - \sum_{\alpha=x,y} \hat{\Sigma}_\alpha \hat{u} \hat{\Sigma}_\alpha \hat{u} \right] \ln \left(\frac{\hbar\omega}{E_F} \right)$$

(valid for $1/\tau < \omega < v k_F$)

- logarithmic singularity tied to Fermi energy (zero-bias anomaly)
- relative strength of anomaly independent of Fermi energy (since $\nu_F \sim k_F$)

Potential disorder
(e.g. u_{zz})

$$\frac{\delta\nu_\omega(\mathbf{r})}{\nu_F} \simeq -\frac{4\nu_F^2}{\pi^{3/2}} \frac{u_{zz}^2}{(k_F r)^2}$$

(valid for $\lambda_F < r < v/\omega$)

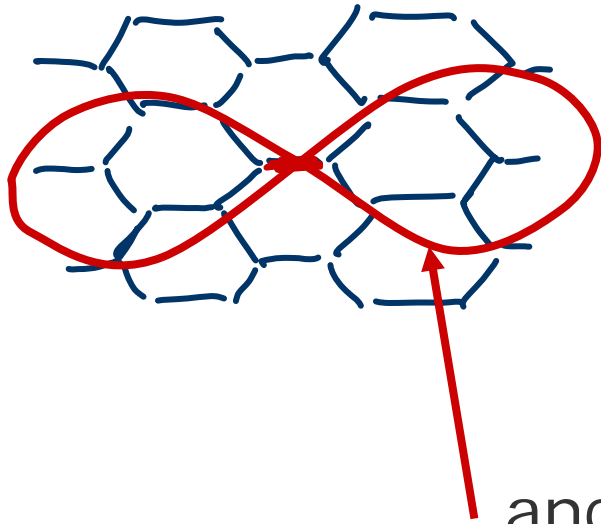
Fictitious gauge field
(e.g. u_{zx})

$$\frac{\delta\nu_\omega(\mathbf{r})}{\nu_F} \simeq -\frac{4\nu_F^2}{\pi^{3/2}} \frac{u_{zx}^2 \sin^2 \phi}{(k_F r)^2}$$

angular dependence!

- u_{zy} gives $\cos^2 \phi$ dependence
- analogous for nonabelian gauge field

e.g. random hopping



contributes to u_{zy} (abelian),
 u_{yz} and u_{xz} (nonabelian)

angular dependence for intravalley
contribution

- Logarithmic zero-bias anomaly in tunneling density of states due to interactions and disorder
- Originates from Fock contribution alone
- Relative strength of zero-bias anomaly is independent of doping
- Angular dependence of tunneling density of states around impurity is signature of fictitious gauge field

E. Mariani, L. Glazman, A. Kamenev, FvO, unpublished