

GRAPHENE: SYMMETRIES, TRANSITIONS

HALL EFFECT

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- IH, PRL 97, 146401 '06
- IH, COND-MAT/0610349

ISSUES:

- $U \sim 10$ eV : WHY NOT STRONGLY CORRELATED?
- WHAT IF IT WERE? (GROSS-NEVEU CRIT. POINT.)
- WEAK INTERACTION & $B \neq 0$: $\nu = 0$, EVEN, 1

LINEARIZE :

$$H_0 = \int_{\vec{k}=\vec{k}+\vec{p}} d^2\vec{k} (U^\dagger(\vec{p}), V^\dagger(\vec{p})) \begin{pmatrix} v_F p_x \tau_x - v_F p_y \tau_y \\ \end{pmatrix} \begin{pmatrix} U(\vec{p}) \\ V(\vec{p}) \end{pmatrix}$$

$$+ \int_{\vec{k}=-\vec{k}+\vec{p}} d^2\vec{k} (U^\dagger(\vec{p}), V^\dagger(\vec{p})) \begin{pmatrix} -v_F p_x \tau_x - v_F p_y \tau_y \\ \end{pmatrix} \begin{pmatrix} U(\vec{p}) \\ V(\vec{p}) \end{pmatrix}$$

OR, IN ACTION: $S = \int_0^{1/\hbar\omega_T} dt \int d^2\vec{x} L_0$

$$L_0 = \sum_{z=\pm 1} \Psi_z^\dagger(\vec{x}, t) \left(\frac{\partial}{\partial t} + M_1 \frac{\partial}{\partial x} + M_2 \frac{\partial}{\partial y} \right) \Psi_z(\vec{x}, t)$$

WHERE

$$\Psi_z(\vec{x}, t) = T \sum_{\omega_n} \int_0^{2\pi} \frac{d\vec{k}}{(2\pi a)^2} e^{i\omega_n t + i\vec{k}\cdot\vec{x}} \begin{pmatrix} U_z(\vec{k}+\vec{k}, \omega_n) \\ V_z(\vec{k}+\vec{k}, \omega_n) \\ U_z(-\vec{k}+\vec{k}, \omega_n) \\ V_z(-\vec{k}+\vec{k}, \omega_n) \end{pmatrix}$$

$$M_1 \approx \frac{1}{a}, \quad v_F = ta \frac{\sqrt{3}}{2} (\rightarrow 1)$$

MATRICES :

$$M_1 = -i \frac{\tau_x}{-2x}, \quad M_2 = i \frac{\tau_y}{2y} \quad ; \quad \{M_1, M_2\} = 0$$

FIND SOME \mathcal{F}_0 : $\mathcal{F}_0^2 = 1, \{M_i, \mathcal{F}_0\} = 0$
SO THAT :

$$L_0 = \sum_{z=\pm 1} \Psi_z^\dagger \mathcal{F}_0 (\mathcal{F}_0 \partial_t + \mathcal{F}_1 \partial_x + \mathcal{F}_2 \partial_y) \Psi_z$$

$$\mathcal{F}_1 = \mathcal{F}_0 M_1$$

$$\mathcal{F}_2 = \mathcal{F}_0 M_2$$

FREEDOM OF CHOICE (IN \mathcal{F}_0) :

$$\mathcal{F}_0 = a \mathbf{I}_2 \otimes \tau_z + b \tau_x \otimes \tau_x + c \tau_y \otimes \tau_x,$$

OR

$$\mathcal{F}_0 = \tau_z \otimes \tau_z,$$

$$a^2 + b^2 + c^2 = 1.$$

"CHIRAL" SYMMETRY : (ROTATIONS OF (q, s, c))

• FOR EACH SPIN : $SU(2) = \left\{ \mathcal{J}_3, \mathcal{J}_5, i\mathcal{J}_3\mathcal{J}_5 \right\}$

EXAMPLE: $\mathcal{J}_0 = \frac{z_x}{z_z}$, $\mathcal{J}_3 = \frac{z_y}{z_x}$, $\mathcal{J}_5 = \frac{-iz_y}{z_x}$

• SPIN - 1/2 : $U(4)$, 16 GENER.

$\left\{ I_2, \vec{b} \right\}_{2 \times 2} \otimes \left\{ I_4, \mathcal{J}_3, \mathcal{J}_5, \mathcal{J}_{35} \right\}_{4 \times 4}$

EXACT : $\vec{b} \otimes I_4$; SPIN ROTAT. $SU(2)$
 $I_2 \otimes I_4$; $U(1)$, NUMBER CONS.

THE REST "EMERGING" !

BROKEN BY :

$L \sim m \Psi_2^+ \mathcal{J}_0 \Psi_2$ (- $\mathcal{J}_3, \mathcal{J}_5$ BROKEN)
 - \mathcal{J}_{35} PRESERVED)

- $q=1 \Rightarrow$ CDW
- $q=0 \Rightarrow$ "NEUTLE" (HOU, CHANOU, MUDRY, COND-MAT/0609740)
- $q=2, AF$

INTERACTION :

$H_1 = \frac{U}{2} \sum_{\vec{r}, \sigma, \sigma'} n_{\sigma}(\vec{r}) n_{\sigma'}(\vec{r})$; $\vec{X} = \vec{A}, \vec{B} = \vec{A} + \vec{b}$

DECOMPOSITION : n - density
 u - magnetization
 cov

$H_1 = \frac{U}{8} \sum_{\vec{r}} \left\{ \begin{aligned} & (n(\vec{A}) + u(\vec{A} + \vec{b}))^2 + (n(\vec{A}) - u(\vec{A} + \vec{b}))^2 - (u(\vec{A}) + m(\vec{A} + \vec{b}))^2 \\ & - (u(\vec{A}) - m(\vec{A} + \vec{b}))^2 \end{aligned} \right\}$

NEAR $\pm \vec{V}$:

$\Psi_2^+(\vec{r}, \sigma) \Psi_2(\vec{r}, \sigma) \approx \frac{1}{2} \bar{\Psi}_2(\vec{r}, \sigma) (I_2 \otimes e^{i2\vec{r} \cdot \vec{r}_2} \frac{z_x}{z_z} \otimes (z_x \otimes I_2)) \Psi_2$
 $v=U$
 $v=V$

WITH INTERACTION: $L = L_0 + \sum_{x=d,e,t,a} L_x$

$$L_x = g_x \left(\sum_z W_{x,z} \bar{\Psi}_z M_x \Psi_z \right)^2 + \tilde{g}_x \sum_{\mu=3,5} \left(\sum_{z=\pm 1} W_{x,z} \bar{\Psi}_z M_x \mathcal{F}_\mu \Psi_z \right)^2$$

$$\begin{aligned} W_{d,2} &= W_{e,2} = 1 \\ W_{t,2} &= W_{a,2} = 2 \\ M_d &= M_t = \mathcal{F}_0 \\ M_c &= M_a = \mathbf{I} \end{aligned}$$

("RELATIVISTIC MASS")

- d - density
- t - ferromagnetic
- c - charge den. wave
- a - anti fermion

→ $g_c = (U - V) \frac{q^2}{8}$

$g_d = (U + V) \frac{q^2}{8}$

$g_t = -U \frac{q^2}{8}$

→ $g_a = -U \frac{q^2}{8}$

V - NEXT NN REPULSION

$$\left(\tilde{g}_x = -\frac{g_x}{2} \right)$$

RG FOR LARGE N (HARTREE): $1 \rightarrow \frac{1}{6}$

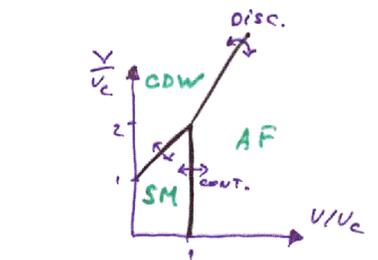
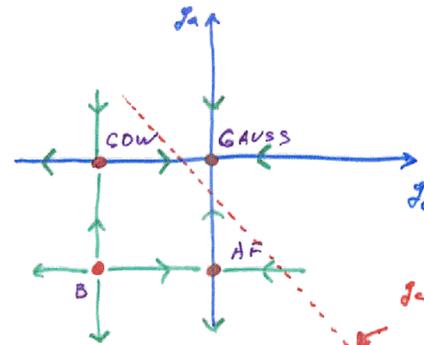
$$\frac{dg_x}{d \ln b} = -g_x - C_x g_x^2$$

$$\left(\frac{d\tilde{g}_x}{d \ln b} = -\tilde{g}_x + 2\tilde{g}_x^2 \right)$$

$$\begin{aligned} C_{d,t} &= 0 \\ C_{c,a} &= 4 \end{aligned}$$

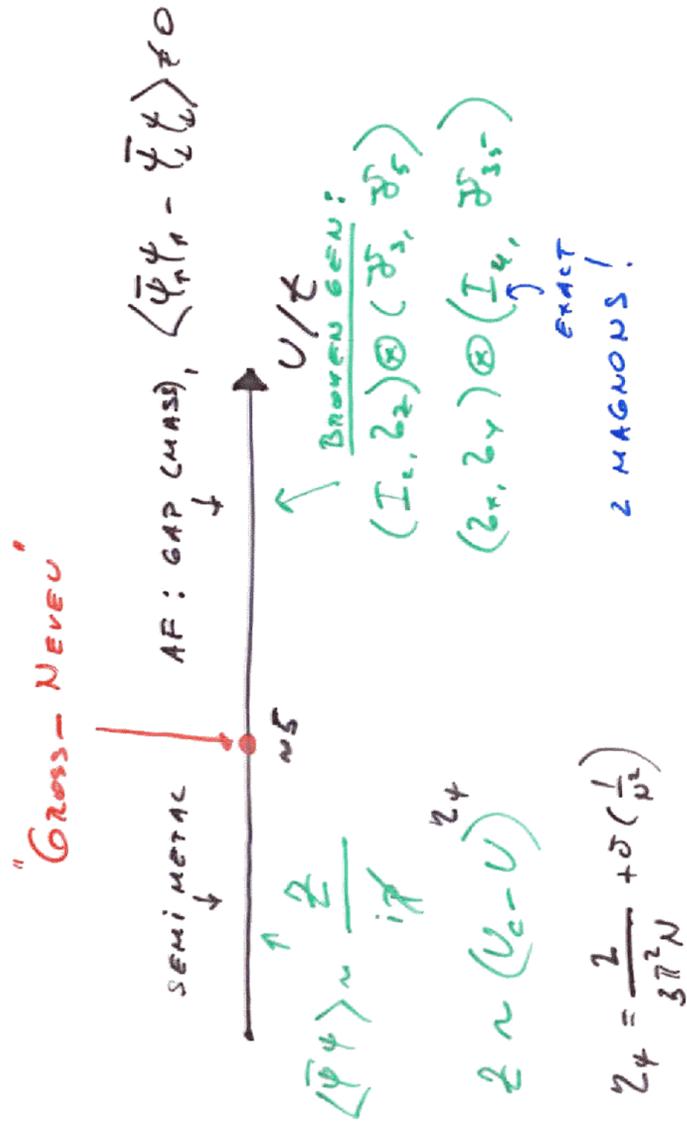
$$\left(\frac{2g_x \frac{1}{\pi} \rightarrow g_x \right)$$

Flow:



$$g_c + g_a = -\frac{V q^2}{8}$$

(IH, PRL '06)



Higgs-like PHENOMENON!
(AT STRONG COUPLING)

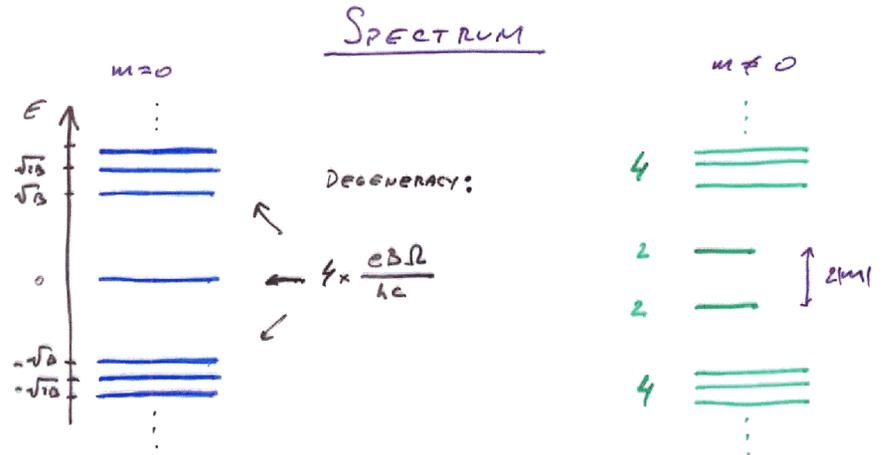
MAGNETIC FIELD : QHE

$B \approx 10 \text{ T} \Rightarrow \frac{l_B}{a} = \sqrt{\frac{B_0}{B}} \approx 100$
 $(B_0 = \frac{1}{a^2} \approx 10^5 \text{ T})$

SO LAB. FIELDS ARE WEAK!

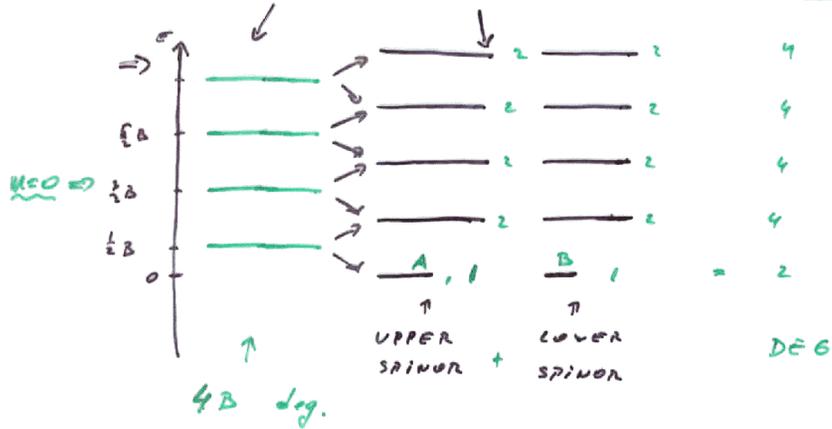
CONTINUUM (LOW-ENERGY) APPROX:
 (NON-INTERACTING, NO ZEEMAN FIRST)

$H_0 \rightarrow M_1(-i\partial_x - A_x) + M_2(-i\partial_y - A_y) + m\mathcal{F}_0$



WHY? TAKE ONE SPIN STATE:

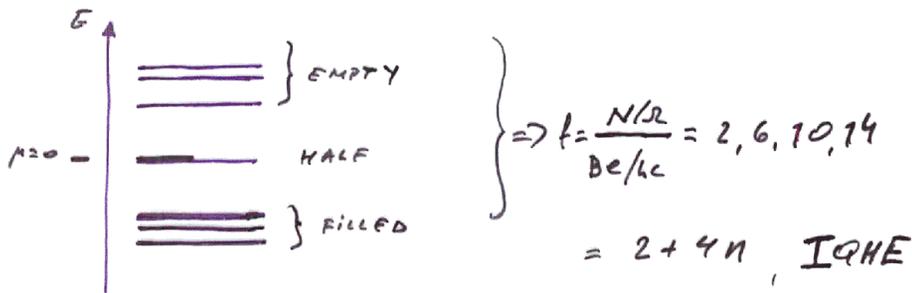
$$H_0^2 = -(\nabla - i\vec{A})^2 + B(z_x \otimes z_x) + m^2 \rightarrow \underline{mB + m^2}$$



$$H_0 \rightarrow \pm \sqrt{mB + m^2}$$

a) $m \neq 0$ $2 = \text{deg}$
 b) $m = 0$: $+m \Rightarrow \text{deg} = 1$
 $-m \Rightarrow \text{deg} = 1$

SO FOR $m = 0$ (GAPLESS), WITH SPIN:



IQHE AT $\nu = \pm 1 \Rightarrow$ INTERACTIONS

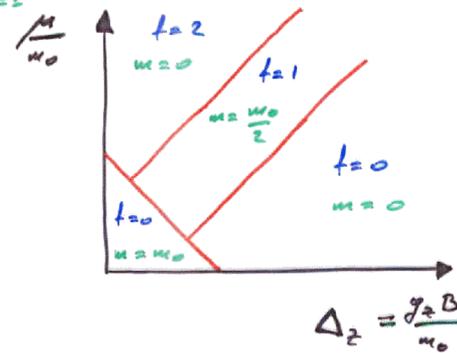
FOR $l_B \gg a$, FIRST RESCALE TO $\lambda = \frac{1}{l_B}$,

\Rightarrow SINGLE DOMINANT (LEAST IRRELEVANT) INTERACTION!

IF $-g_a > -g_c$ ($V=0$):

$$L = i \sum_c \bar{\Psi}_2 \delta_{\mu\nu} \partial_\mu \Psi_2 - \sum_c (\mu + 2g_2 B) \Psi_2^+ \Psi_2 + g_a \sum_c (2 \bar{\Psi}_2 \Psi_2)^2$$

$$M = g_a \sum_{\nu=1}^2 \nu \langle \bar{\Psi}_2 \Psi_2 \rangle$$



DEGENERACY +

$$M_0 = g_a \left(\frac{2B_\perp}{\pi} \right)$$

Y. ZHANG ET AL, PRL '06 :

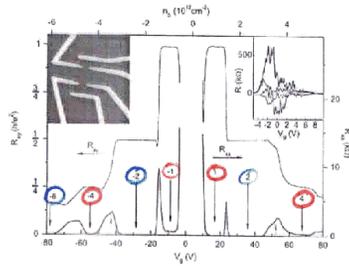


FIG. 1 (color online). R_{xx} and R_{xy} measured in the device shown in the left inset, as a function of V_g at $B = 45$ T and $T = 1.4$ K. $-R_{xy}$ is plotted for $V_g > 0$. The numbers with the vertical arrows indicate the corresponding filling factor ν . Gray arrows indicate developing QH states at $\nu = \pm 3$. n_s is the sheet carrier density derived from the geometrical consideration. Right inset: R_{xx} (dark solid lines) and R_{xy} (light solid lines) for another device measured at $B = 30$ T and $T = 1.4$ K in the region close to the Dirac point. Two sets of R_{xx} and R_{xy} are taken at different time under the same condition. Left inset: an optical microscope image of a graphene device used in this experiment.

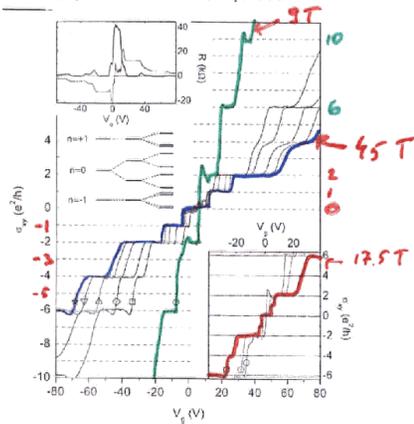


FIG. 2 (color online). σ_{xy} as a function of V_g at different magnetic fields: 9 T (circle), 25 T (square), 30 T (diamond), 37 T (up triangle), 42 T (down triangle), and 45 T (star). All the data sets are taken at $T = 1.4$ K, except for the $B = 9$ T curve, which is taken at $T = 30$ mK. Left upper inset: R_{xx} and R_{xy} for the same device measured at $B = 25$ T. Left lower inset: a sche-

ENERGY SCALES :

- LL SEPARATION $\sim t \sqrt{\frac{B_{\perp}}{B_0}} \sim 100 \text{ meV}$
- Zeeman : $\sim g_2 \left(\frac{B}{10^4 \text{ T}} \right) \text{ eV} \sim 1 \text{ meV}$
- INTERACTION $U_0 \sim U \frac{B_{\perp}}{B_0} \sim 1 \text{ meV}$

AT A LARGE ZEEAMAN TERM ($\sim \sqrt{B_{\perp}^2 + B_{\parallel}^2}$):

- $\Delta_{f=0} \sim B_{\parallel}$
- $\Delta_{f=1} \sim B_{\perp}$

SUMMARY:

- 1) $B \neq 0$: METAL - INS, HIGGS - LIKE TRANSITION AT LARGE INTERACTION
- 2) $B \neq 0$: "CATALYSIS" OF THE "RELATIVISTIC" GAP \Rightarrow IQHE AT $t = \pm 1$, t -EVEN
- 3) ZEEMAN + INTERACTION: **PINE STRUCTURE!**

FUTURE ISSUES ($U \neq 0$)

- $U/t \gtrsim 5$: - OTHER ORDERS, SPIN LIQUIDS?
- FATE OF FERMIONS?
- TOPOLOGICAL DEFECTS AND ZERO MODES

• COULOMB:

$$\frac{d\epsilon^2}{d\ln b} = -\frac{e^4}{8\pi}$$

LOG-CORRECTIONS?

- DOPING: **SUPERCONDUCTIVITY**, ...?
- MASS VS. QH FERROMAGNETISM, $B \neq 0$?
- TUNING U/t IN GRAPHENE?