



Granular Flow in a Modified Split Bottom Couette Cell

Jeremy B. Lechman and Gary S. Grest
Sandia National Laboratories

Xiang Chen, Heinrich M. Jaeger, Greg S. Karczmar, Matthias E. Möbius, and Sidney R. Nagel
University of Chicago

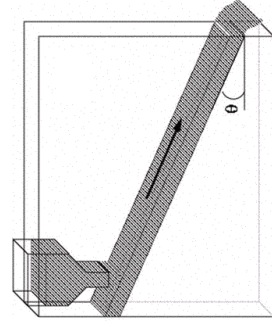
Antonio F. Barbero
University of Almeria, Spain



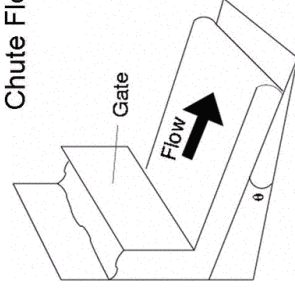
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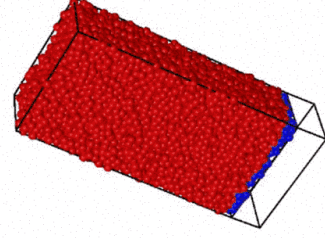
Dense Gravity Driven Flow



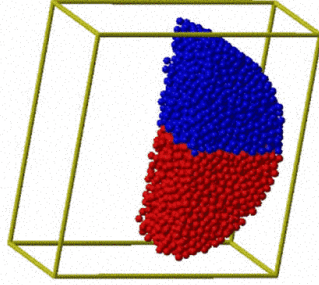
Chute Flow



Heap Flow



Dense Granular Flows



Rotating Drum

- Boundary Driven Flows
 - Couette Flow
 - Shear Flow

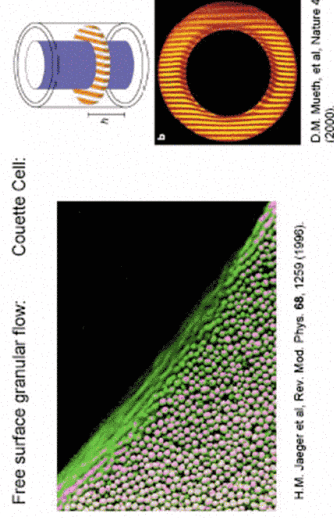
Shear Bands in Dense Granular Flow

- Shear bands: narrow and distinct bands of high rates of shear deformation (localization of energy dissipation)

– Phenomenon plays an important role in many applications

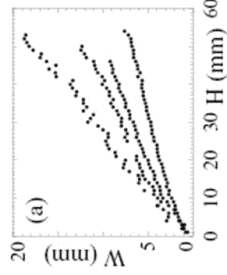
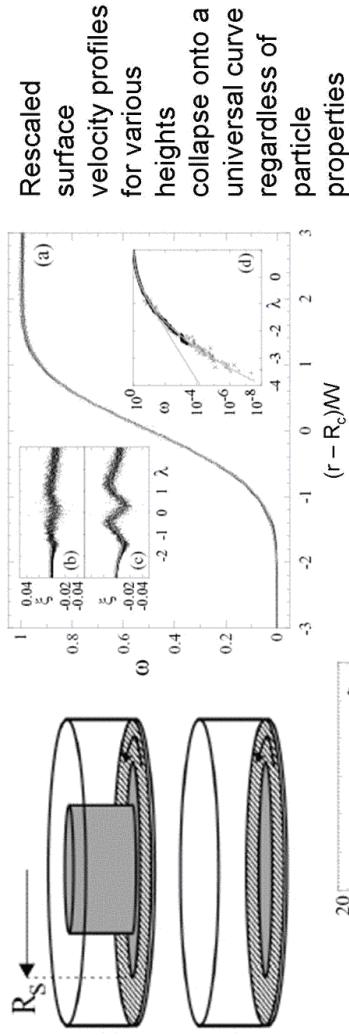
- ballistic impact
- explosive fragmentation
- high speed machining
- metal forming
- interfacial friction
- powder compaction
- soil failure
- seismic events
- **granular flow**

- What is the role of micro-structure?
- Difficult to access range of a flowing states to test flow theories
- Non-universality



exponential velocity profiles

Split Bottom Couette Cell



Shear zones are detached from the walls and become wide with increasing total height H .

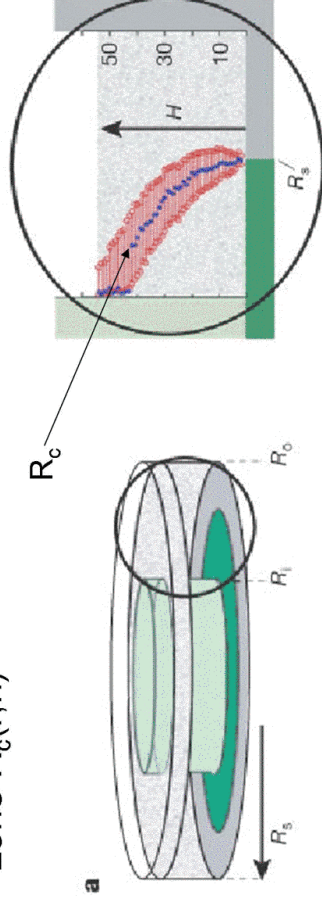
Fenistein *et al.* (Nature **425**, 256 (2003) ; PRL **92**, 094301 (2004))

- Address questions of velocity fluctuations, particle diffusion, and granular temperature?
- Relationship between micro and macro scales?

focus attention on 'shallow' packs

Shape of the Shear Zones

- Parameters involved in rescaling center $R_c = f(R_s, H)$ and width $W = f(H, \text{particle})$ of shear zone appear to have separate length scales
- Theoretical description/predictions for shape of shear zone $R_c(r, h)$



Fenistein *et al.* (Nature **425**, 256 (2003); PRL **92**, 094301 (2004))

Discrete Element Simulations

- Allows observation of bulk behavior away from influence of side walls
- Allows detailed measurements of microscopic quantities (e.g., inter-particle forces)
- Observe bulk behavior, go beyond “shallow” regime, test theory

Integrate Newton’s equations-translational and rotational d.o.f.

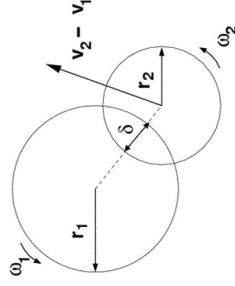
$$\mathbf{F}_n = f(\delta/d)(k_n \delta \mathbf{n}_{ij} - \frac{m}{2} \gamma_n \mathbf{v}_n)$$

$$\mathbf{F}_t = f(\delta/d)(-k_t \Delta s_t - \frac{m}{2} \gamma_t \mathbf{v}_t)$$

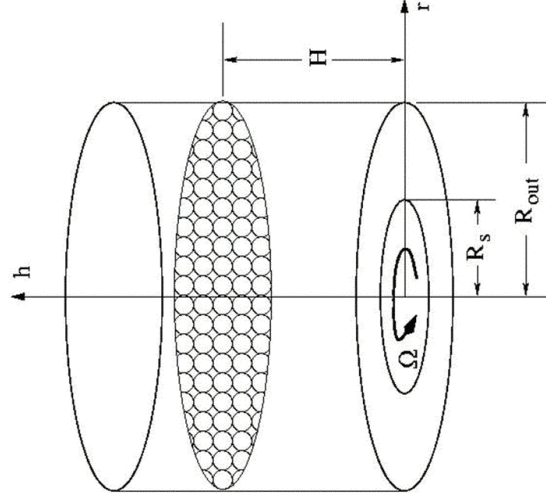
$$f(x) = \sqrt{x} \quad \text{Hertzian springs}$$

Δs_t Elastic tangential displacement

$F_t \leq \mu F_n$ Coulomb Failure Criterion



DEM System Parameters



$$R_s = 30.0d$$

$$R_{out} = 37.8d$$

$$\Omega = 0.014 \text{ rad}/\tau \quad \text{where } \tau = \sqrt{d/g}$$

$$5.4d \leq H \leq 34.2d$$

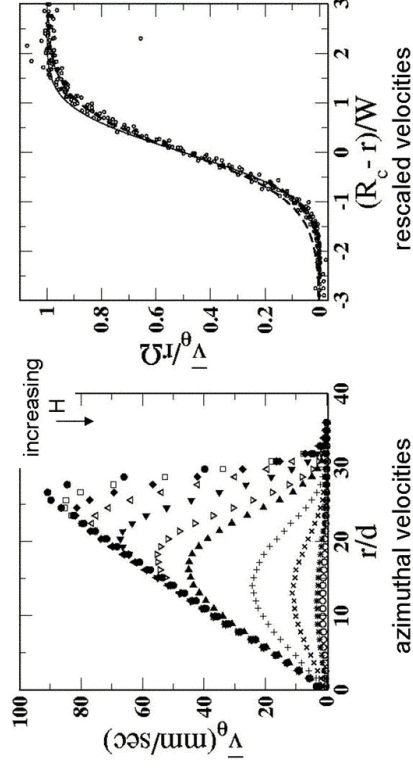
20,000 – 180,000 particles

rough bottom composed of layer of glued particles

- All experiments, simulations are in hydrostatic regime ($H < 6R_{out}$)

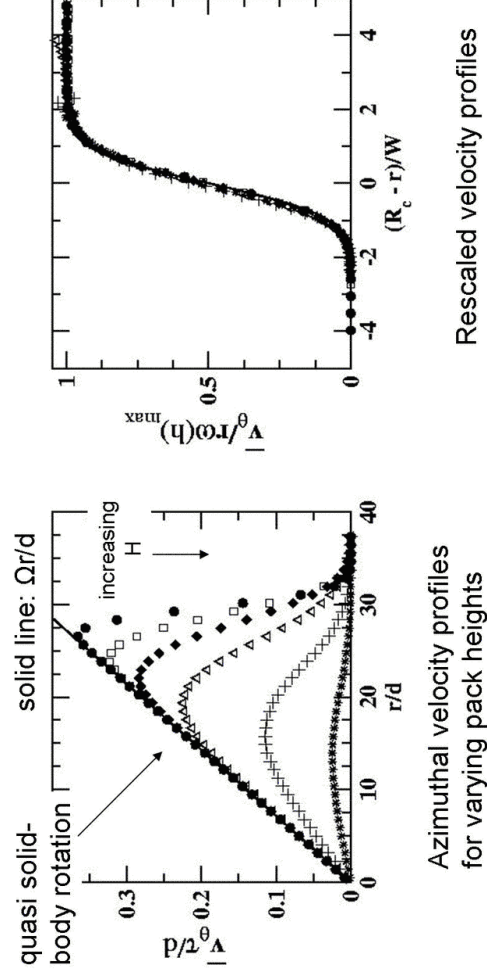
- Values picked to exactly match experimental system for poppy seeds ($d=0.8\text{mm}$) - particle tracking at surface
- MRI – poppy and rajagara seeds – $R_{out}=51.6d$, $R_s=40.6d$

Experimental Surface Velocity Profiles – Poppy Seeds

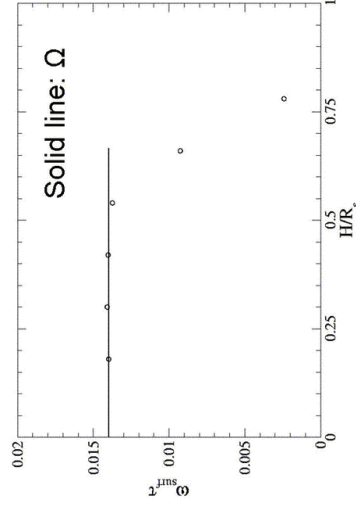


- Linear azimuthal velocity profile near center for shallow packs (regime of previous work)
- Slight asymmetry in the rescaled velocities

Simulation Surface Velocity Profiles



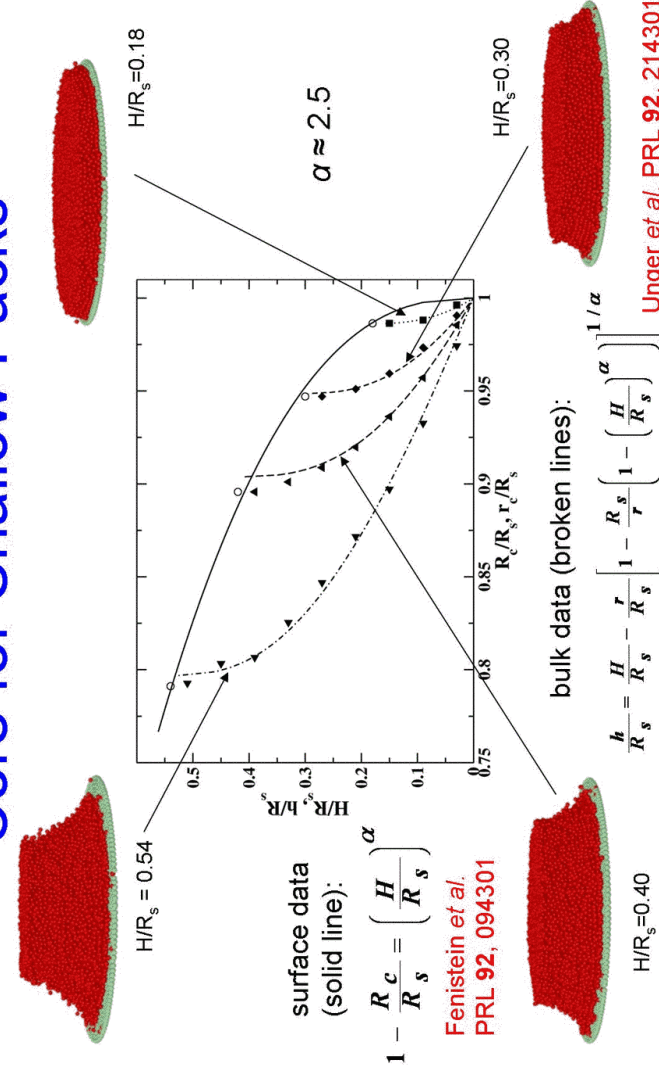
Shallow vs. Deep Packs



- As observed from the surface: qualitative change at $H/R_s \approx 0.5$ in agreement with previous work

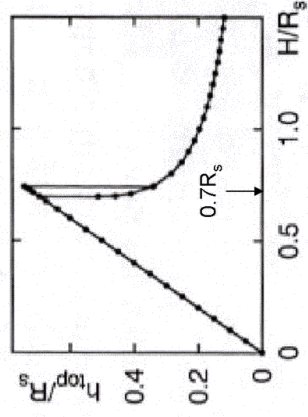
What happens in the bulk?

Shape of Shear Zone and Inner Core for Shallow Packs

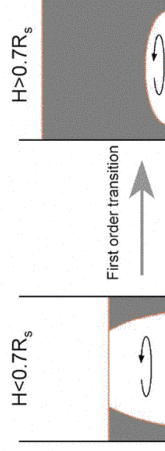


Proposed Theory – Unger et al.

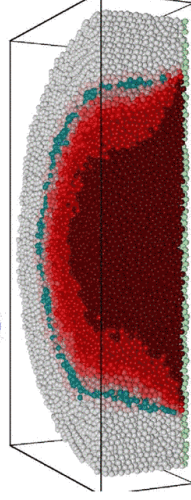
- Least dissipation (minimum torque)
 - Infinitely thin shear surface between two bulk solid regions
 - Hydrostatic pressure
 - Coulomb friction between solid regions
- Describes shape of shear zone for shallow packs based on bulk stress state
- Predicts for tall packs:
 - transition in shape of shear zone (open → closed)
 - first order accompanied by hysteresis
 - height of the shear zone, h_{top} is proportional to R_s/H



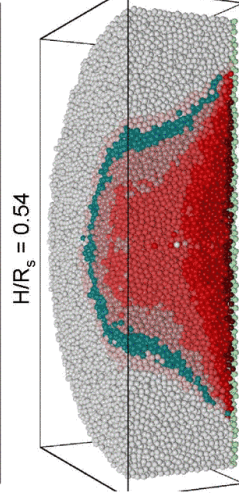
Unger et al. PRL 92, 214301 (2004)



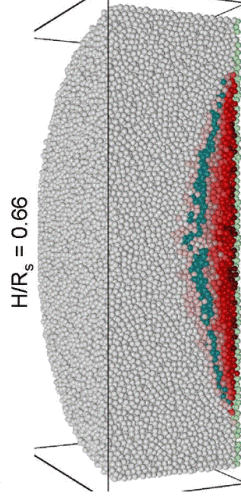
Normalized Angular Velocity



“open” shape

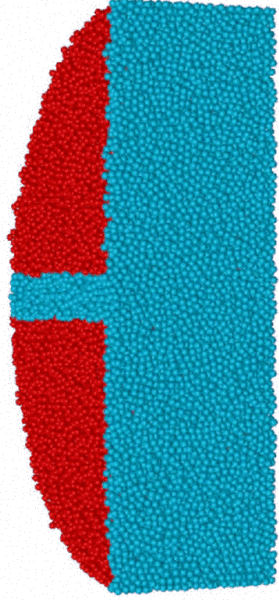


- $\frac{\omega}{\Omega} \geq 0.95$ dark red
- $0.45 \leq \frac{\omega}{\Omega} < 0.55$ teal
- $\frac{\omega}{\Omega} < 0.35$ white



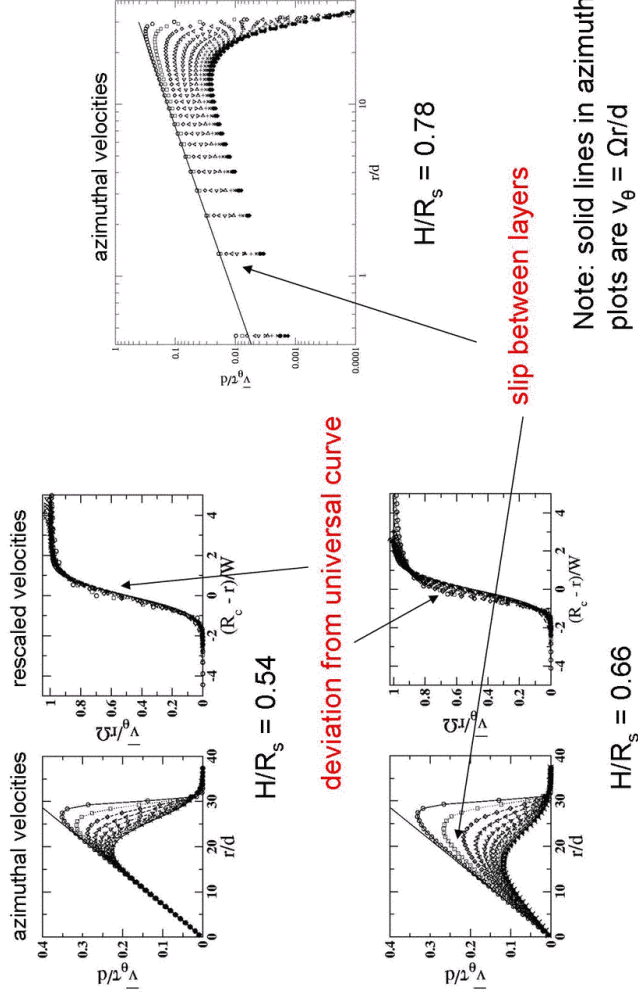
“closed” shape

Shear in Deep Pack



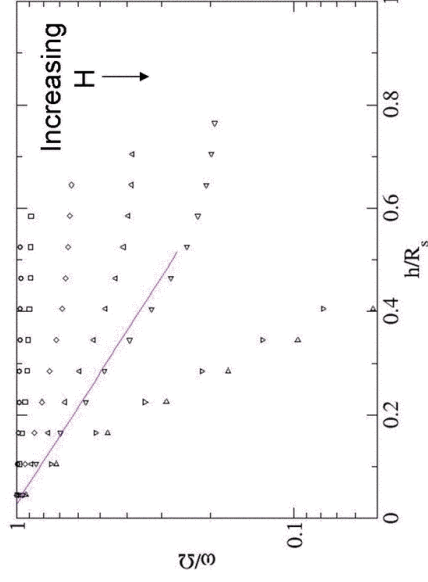
$H/R_s = 0.78$

Bulk Velocity Profiles



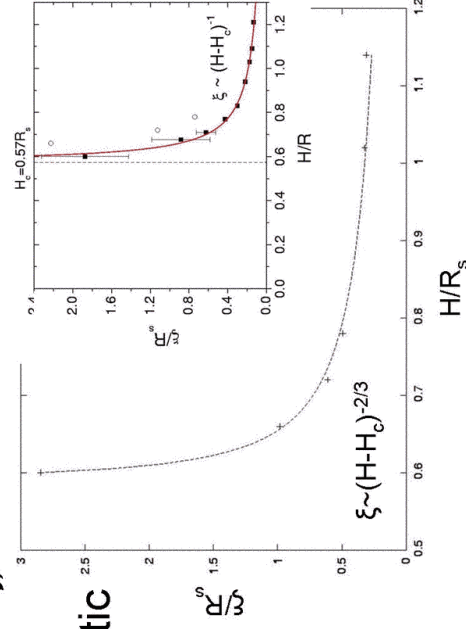
Slip Between Layers

- For $H/R_s > 0.5$ slip between layers increases with H
- MRI give good data for deep packs and longer times



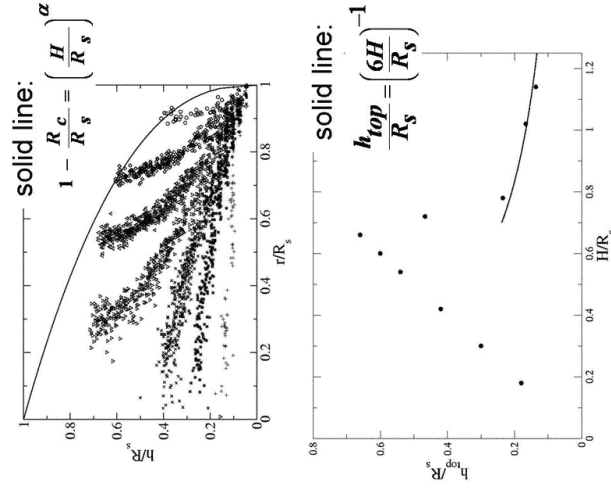
Slip Transition?

- From $w(h) \sim \exp(-H/\xi)$
- ξ gives characteristic length scale
- ξ diverges with decreasing H

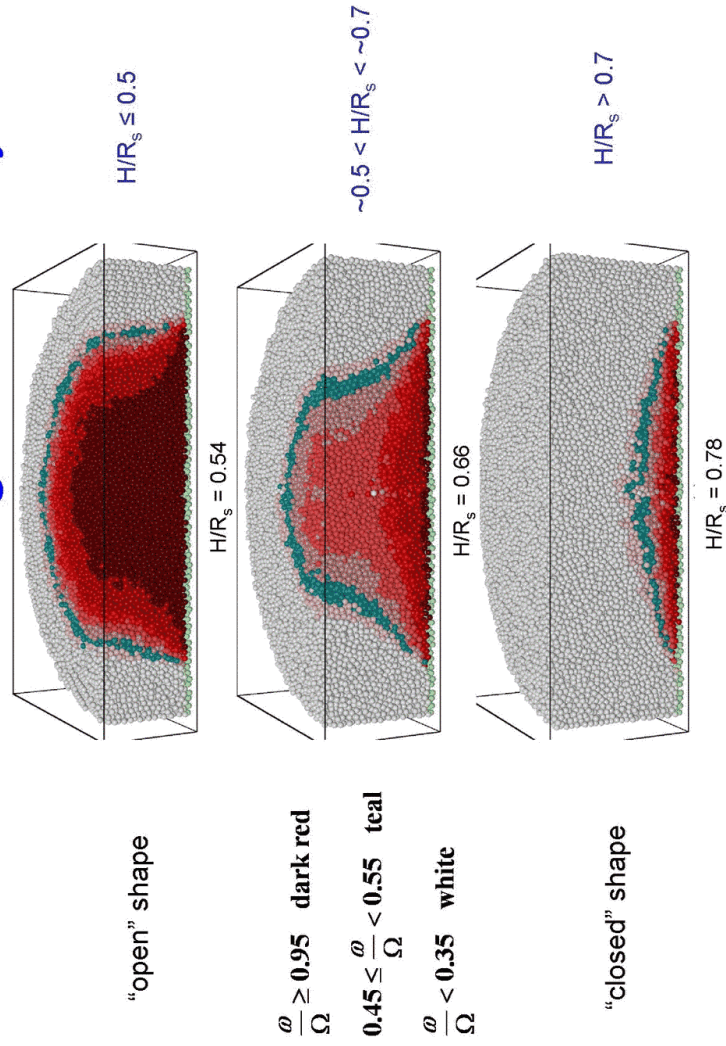


Nature of the Transition

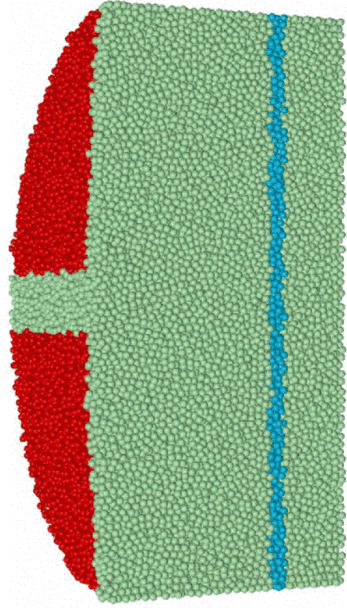
- Basic assumptions violated
 - Shear zone has finite width
 - Smooth transition between moving and stationary regions
 - Slip between layers
- Direct test: How to define top of shear zone?
 - Choose $\frac{\omega}{\Omega} = 0.5$?



Normalized Angular Velocity

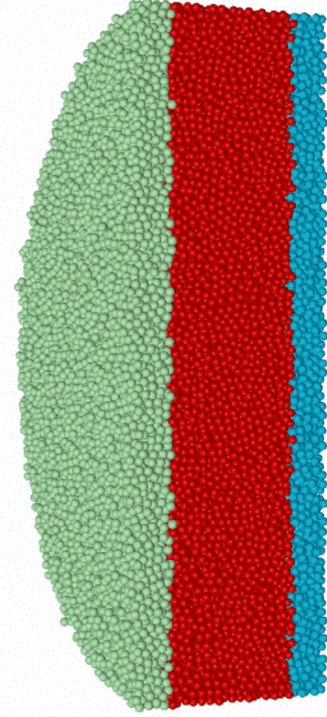


Shear in Deepest Pack



$$H/R_s = 1.14$$

Convection in Deep Pack



$$H/R_s = 0.78$$

Conclusions

- Slip between layers is an increasingly significant mode of deformation for packs of $H > 0.5 R_s$
 - apparently related to torsion failure of the inner core
- Results indicate a second order slip transition
- Theory captures bulk stress state, but misses details of the flow.
 - Finite width of the shear zone, torsion failure
 - A “shape transition” may still be present
 - For deepest packs, trend is $h_{top} \propto R_s/H$
- Attraction interactions reduce the value of transition to lower H
- Convection observed for intermediate piles
- Can variational theory be extended to account for these or is there a more fundamental difference in the nature of the flow?