

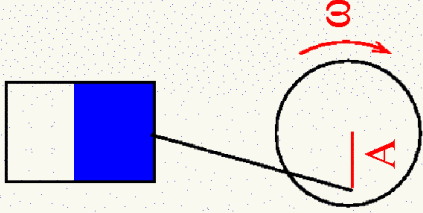
Vibrofluidization

Thorsten Pöschel, Thomas Schwager, Simon Renard

Charité, Berlin

Clara Salueña

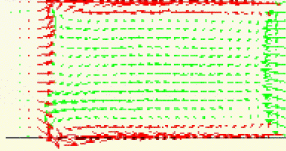
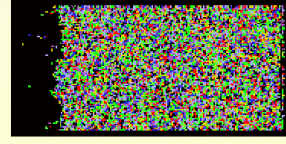
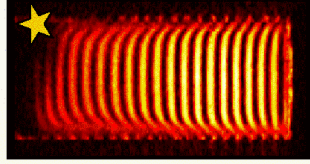
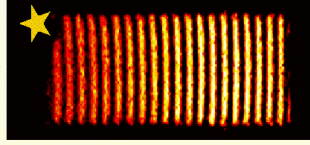
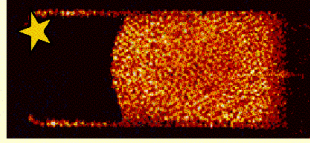
Universitat de Tarragona



1. Vertically vibrated granular material, some effects
2. Conditions for convection: Theory
3. Experiment

1. Vibrated granular material: effects

- convection
- segregation
- pattern formation
- oscillons
- bubbles
- compaction
- ...

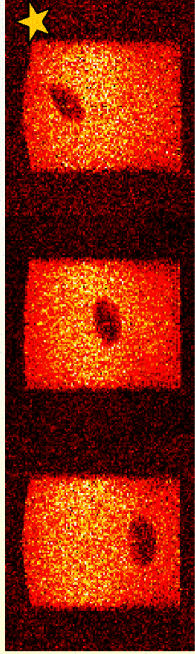


Ehrichs, Jaeger, Karczmar, Knight, Kuperman, Nagel, Science'95

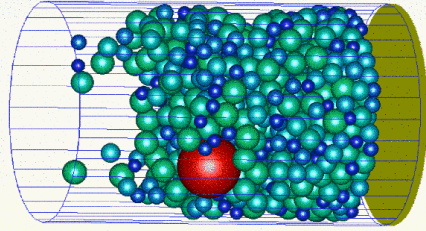
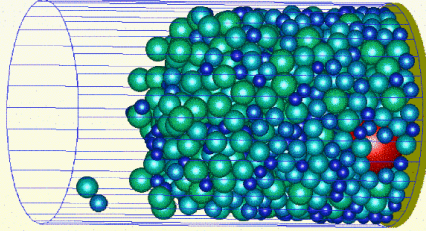
Salueña, Pöschel, Esipov'97

1. Vibrated granular material: effects

- convection
- segregation
- pattern formation
- oscillons
- bubbles
- compaction
- ...



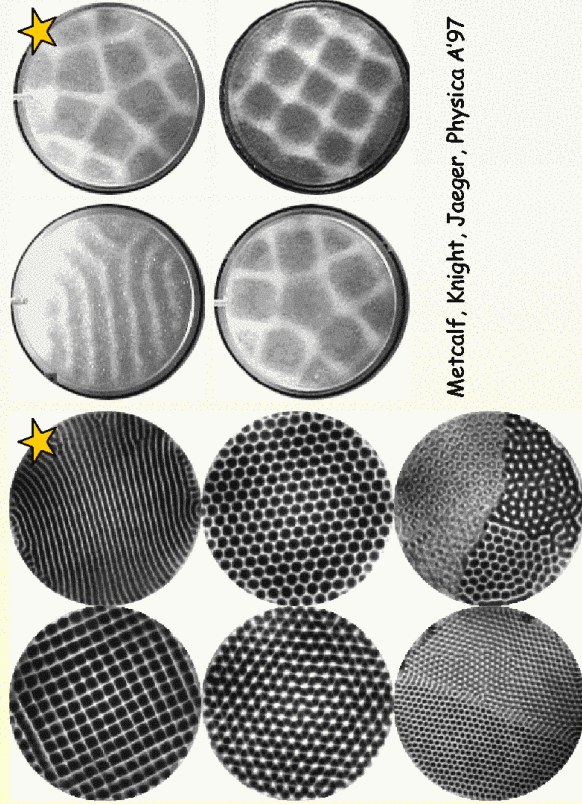
Knight, Jaeger & Nagel, PRL'93: Convection Connection



Gallas, Herrmann, Pöschel, Sokołowski, J. Stat. Phys'96

1. Vibrated granular material: effects

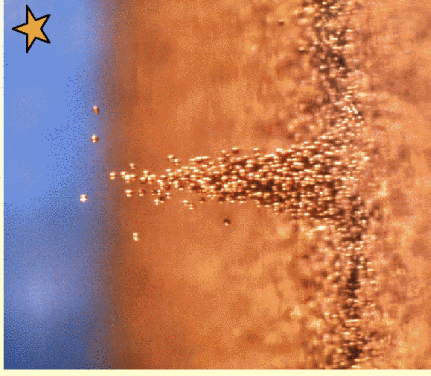
- convection
- segregation
- pattern formation
- oscillons
- bubbles
- compaction
- ...



Metcalfe, Knight, Jaeger, Physica A'97

Melo, Umbanhowar, Swinney, PRL'95

1. Vibrated granular material: effects



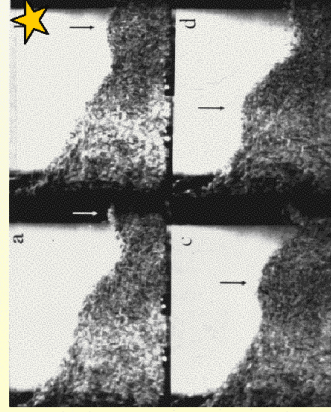
- convection
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- bubbles
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- ...

Umbanhowar, Melo, Swinney, Nature'96

1. Vibrated granular material: effects

- convection
- segregation
- pattern formation
- oscillons
- bubbles
- compaction
- ...

vibrated annular container



travelling waves

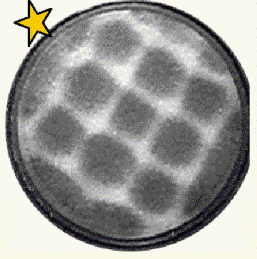
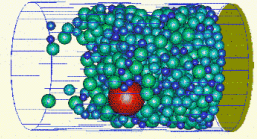
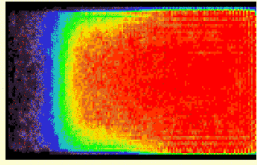
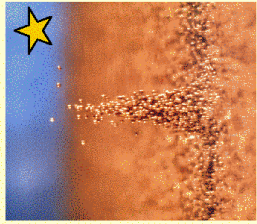
bubbles

Pak & Behringer, PRL'93

Pak & Behringer, Nature'94

1. Vibrated granular material: effects

What have all these effects in common?



For all these effects the particles have to lose contact for part of the period. → Vibrofluidization

motion of the vibrating table: $z_0(t) = A_0 \cos \omega_0 t$

necessary condition:

$$\left(\frac{d^2 z_0}{dt^2} \right)_{\max} = A_0 \omega_0^2 > g$$

$$\Gamma \equiv A_0 \omega_0^2 / g > 1$$

Froude number

???

2. Threshold for the onset of fluidization

Kitchen table experiment: shake sand with tracers

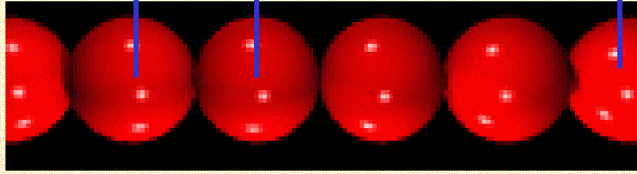
$$f = 8\text{Hz}, \quad \Gamma = 0.96g$$



1 frame / min
(details later)

2. Threshold for the onset of fluidization

We are looking for a necessary condition.
Therefore, we chose the system of smallest energy dissipation.



$$\left. \begin{array}{l} z_{k+1}(t) \\ z_k(t) \end{array} \right\} \xi_{k,k+1} = 2R - |z_k(t) - z_{k+1}(t)|$$

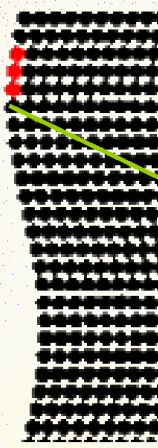
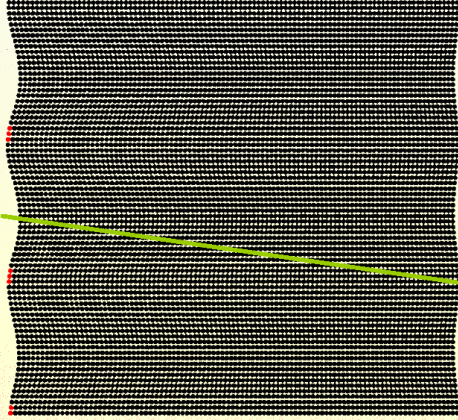
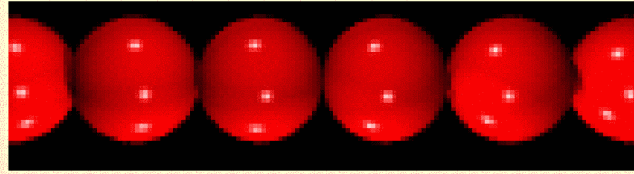
define response function

$$\hat{R}(\omega_0) \equiv \frac{A_N}{A_0} = \frac{\text{amplitude of top particle at frequency } \omega_0}{\text{amplitude of driving}}$$

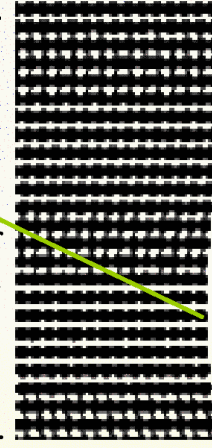
For $\hat{R}(\omega_0) > 1$ $\left(\begin{array}{l} A_N \omega_0^2 > g \\ A_0 \omega_0^2 < g \end{array} \right)$ is possible!

$$z_0(t) = A_0 \cos \omega_0 t$$

2. Threshold for the onset of fluidization



MD simulation (unrealistic parameters, only for illustration)



time →

time →

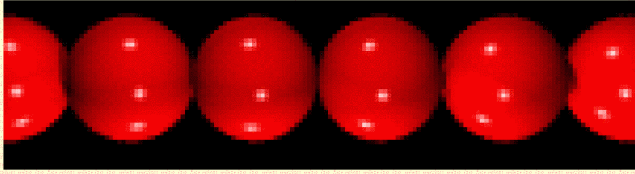
$$\hat{R}(\omega_0) \equiv \frac{A_N}{A_0}$$

For $\hat{R}(\omega_0) > 1$ $\left(\begin{array}{l} A_N \omega_0^2 > g \\ A_0 \omega_0^2 < g \end{array} \right)$ is possible!

2. Threshold for the onset of fluidization

$$\frac{d^2 z_k}{dt^2} = \frac{1}{m} (F_{k,k+1} - F_{k-1,k}) - g$$

Newton's equation of motion



particle model: viscoelastic spheres



$$F(t) = F^{el}(t) + F^{dis}(t)$$

$$F^{el} = -\mu \sqrt{R} \xi^{3/2} \quad \mu = \frac{1}{3} \frac{Y}{1-\nu^2} \quad \text{H. Hertz 1881}$$

$$F^{dis} = -\alpha \sqrt{R} \sqrt{\xi} \frac{d\xi}{dt} \quad \alpha = \frac{\sqrt{2}}{3} \frac{YA}{1-\nu^2} \quad \text{Brilliantov, Spahn, Hertzsch, Pöschel, PRE'96}$$

define $u_k = z_k - 2Rk \quad (k=0, \dots, N)$

$$\xi_{k,k+1} = u_k - u_{k+1}$$

top particle loses contact if $u_N > u_{N-1}$

2. Threshold for the onset of fluidization

$$\frac{\partial^2 u}{\partial t^2} = -g - \frac{\partial}{\partial z} \left[\kappa \left(-\frac{\partial u}{\partial z} \right)^{3/2} \right] - \beta \frac{\partial^2 u}{\partial t \partial z} \sqrt{-\frac{\partial u}{\partial z}} \quad \left. \frac{\partial u}{\partial z} \right|_{z=L} = 0$$

$$\kappa = 2Y/\pi\rho(1-\nu^2) \quad \beta = A\kappa \quad \rho = \text{density}$$

Newton's equation of motion for the displacement of the particles

Limit of elastic particles

new variables $x = 1 - \frac{z}{L} \quad \tau = \left(\frac{g\kappa^2}{L^5} \right)^{1/6} t \quad \gamma = \left(\frac{g^2 L^5}{\kappa^2} \right)^{1/6}$

$$\frac{\partial^2 u}{\partial \tau^2} = -\gamma^2 + \frac{1}{\gamma} \frac{\partial}{\partial x} \left[\left(\frac{\partial u}{\partial x} \right)^{3/2} \right] \quad \left. \frac{\partial u}{\partial x} \right|_{x=0} = 0$$

amplification factor

Pöschel, Schwager, Salgueira, PRE2000

$$\hat{R}^{-1} = \left(\frac{6}{25} \right)^{1/5} \Gamma \left(\frac{3}{5} \right) \Omega_0^{2/5} \left| J_{-2/5} \left(\frac{2}{5} \right) \sqrt{6\Omega_0} \right| \quad \Omega = \left(\frac{L^5}{g\kappa^2} \right)^{1/6} \omega$$

2. Threshold for the onset of fluidization

Results for the limit of elastic particles

amplification factor

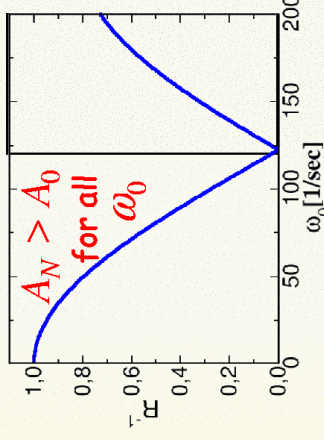
$$\hat{R}^{-1} = \left(\frac{6}{25}\right)^{1/5} \Gamma\left(\frac{3}{5}\right) \Omega_0^{2/5} \left| J_{-2/5} \left(\frac{2}{5}\right) \sqrt{6} \Omega_0 \right|$$

Taylor expansion:

$$\hat{R}^{-1} = 1 - \frac{2}{5} \left(\frac{L^5}{g \kappa^2}\right)^{1/3} \omega_0^2 + \dots$$

one particle: $A_0 = \frac{g}{\omega_0^2}$

a column: $A_0 = \frac{g}{\omega_0^2} - \frac{2}{5} \left(\frac{L^5}{g \kappa^2}\right)^{1/3}$



(parameters for rubber, $Y = 4 \times 10^7$ Pa, $L = 0.6$ m)

$$\kappa = 2Y / \pi \rho (1 - \nu^2)$$

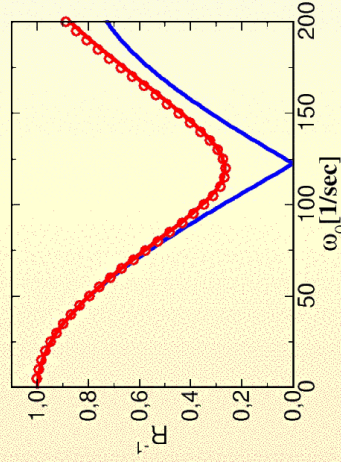
2. Threshold for the onset of fluidization

$$\frac{\partial^2 u}{\partial t^2} = -g - \frac{\partial}{\partial z} \left[\kappa \left(-\frac{\partial u}{\partial z} \right)^{3/2} \right] + \beta \frac{\partial^2 u}{\partial t \partial z} \sqrt{-\frac{\partial u}{\partial z}} \quad \frac{\partial u}{\partial z} \Big|_{z=L} = 0$$

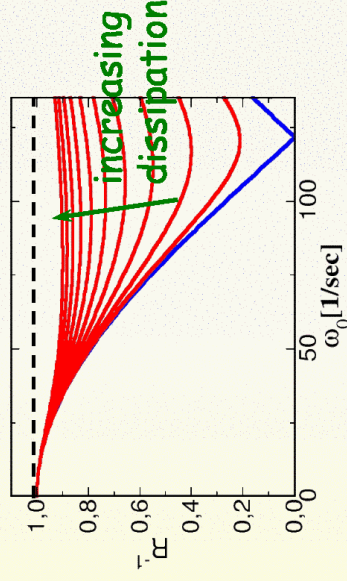
Newton's equation of motion for the displacement of the particles

$$\hat{R}^{-1} = 1 - \omega_0^2 \frac{2}{5} \left(\frac{L^5}{g \kappa^2}\right)^{1/3} + \omega_0^4 \left(\frac{L^5}{g \kappa^2}\right)^{2/3} \left[\frac{3}{100} + \frac{8}{45} \beta^2 \left(\frac{g}{\kappa^4 L^5}\right)^{1/3} \right]$$

Pöschel, Schwager, Salgueira, PRE 2000

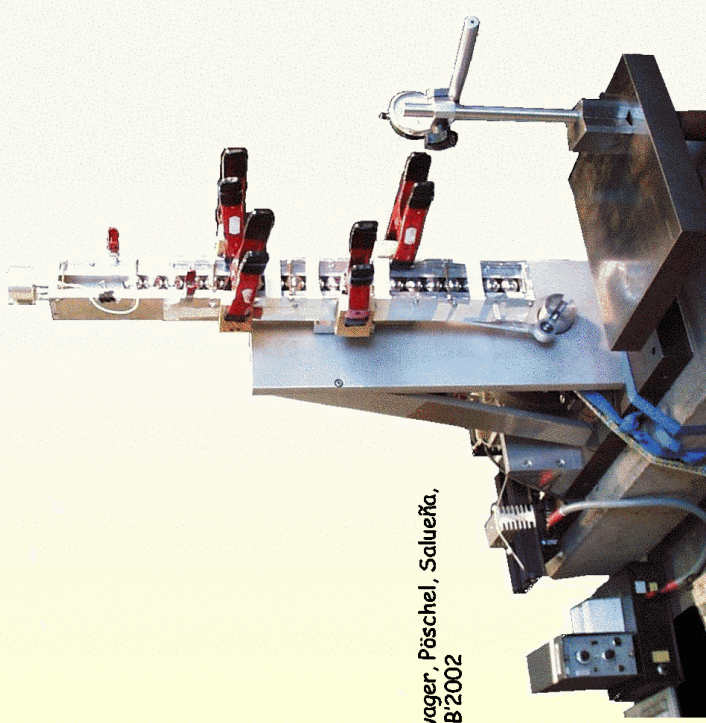
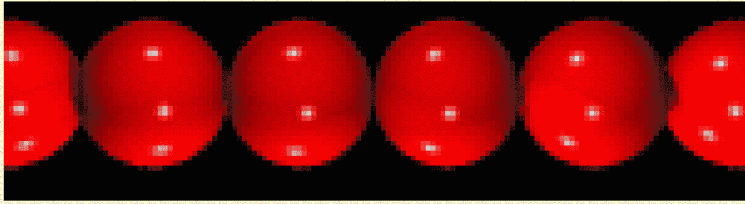


— elastic particles
— dissipative particles
○ MD simulation



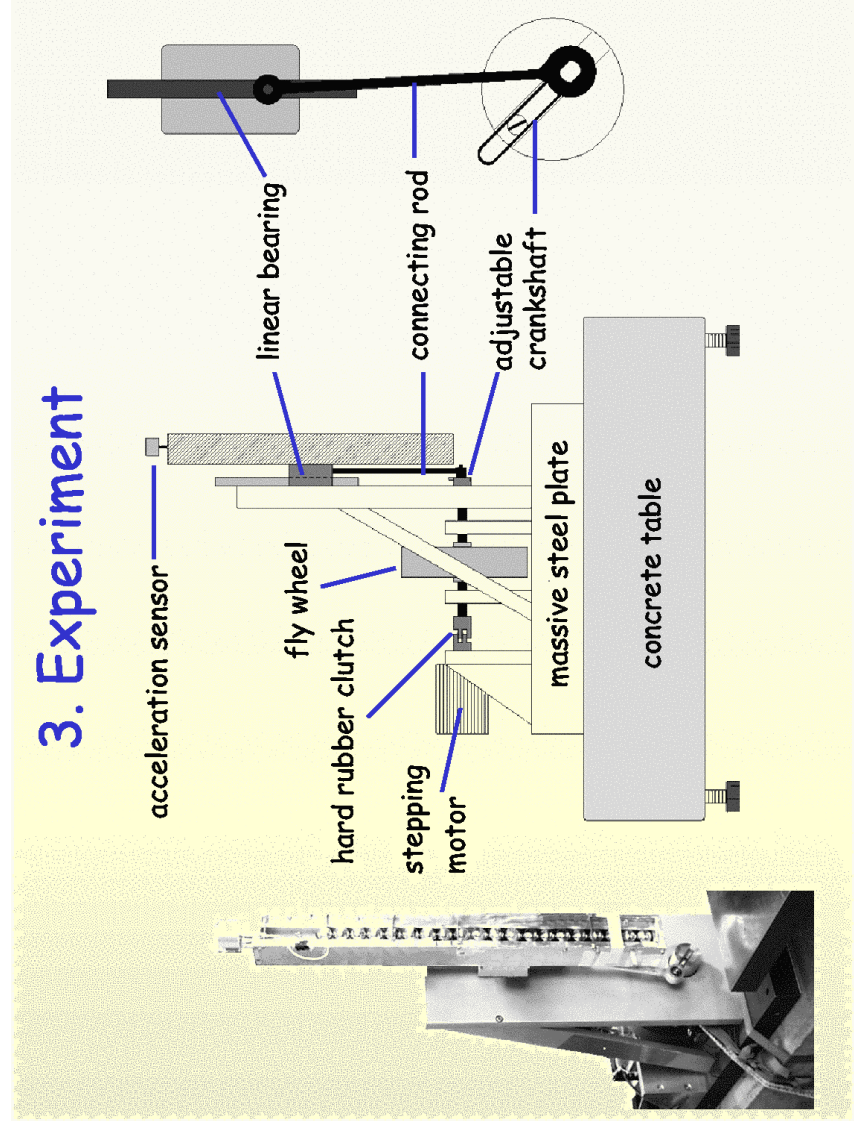
Even for large dissipation there is always a region of amplification.

3. Experiment



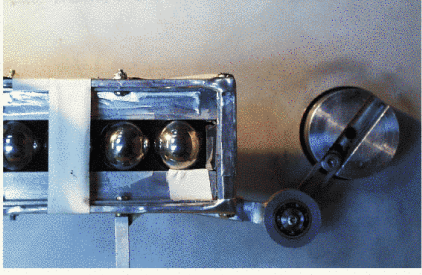
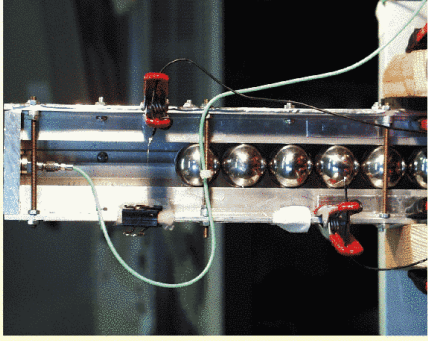
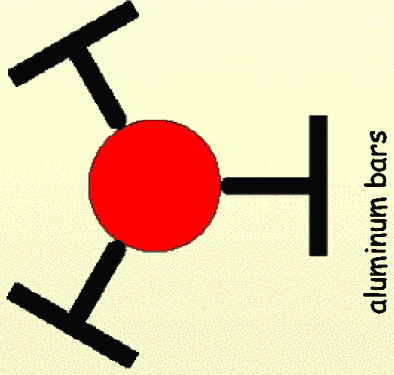
Renard, Schwager, Pöschel, Salueña,
Eur. Phys. J. B 2002

3. Experiment

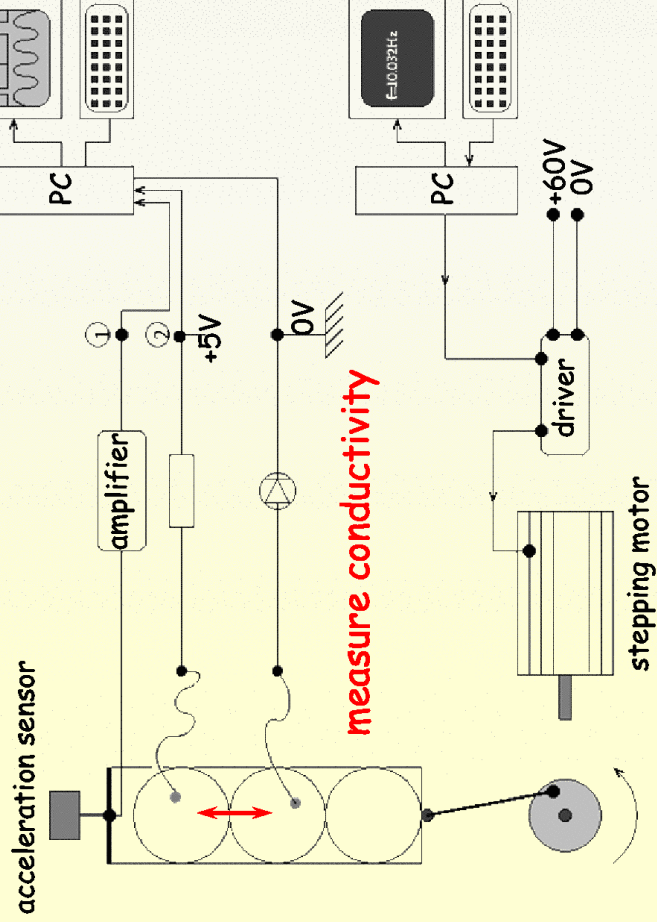


3. Experiment

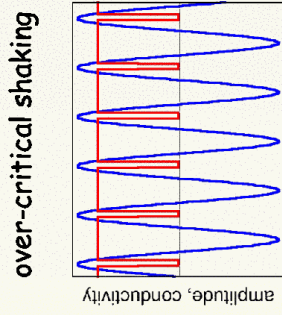
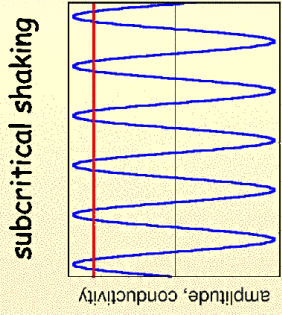
view from top



3. Experiment

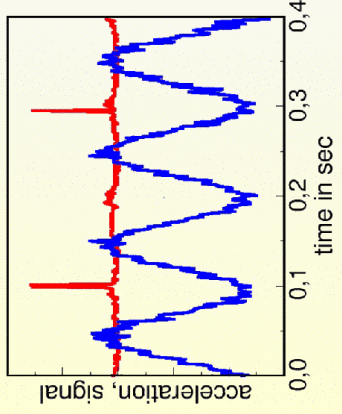
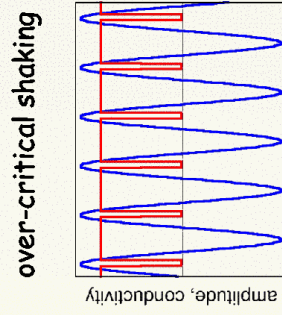
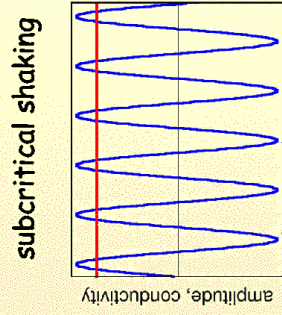


3. Experiment



conductivity
amplitude

3. Experiment



typical signal:
high-frequency
noise!!

3. Experiment

What do we expect?

- 1 sphere: fluidization for $\Gamma \equiv A\omega^2/g > 1$
- 20 spheres: fluidization for $\frac{A\omega^2}{g} > \Gamma \left(\frac{3}{5} \right) \left| M^{2/5} J_{-2/5}(2M) \right|$ $M = \text{material and system constants}$

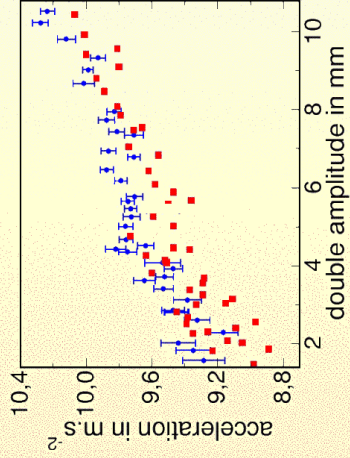
Taylor expansion: $\frac{A\omega^2}{g} > 1 - B_2\omega^2 + B_4\omega^4$

$$B_2 = \frac{18^{2/3}}{45} \left(\frac{\pi^2 \rho^2 L^5}{g\mu^2} \right)^{1/3}$$

$$B_4 = \frac{1}{18^{2/3}} \left(\frac{\pi^2 \rho^2 L^5}{g\mu^2} \right)^{2/3} \left(\frac{3}{100} + \frac{4 \cdot 324^{2/3}}{405} \alpha^2 \left(\frac{g\pi^4}{\mu^4 \rho^2 L^5} \right)^{1/3} \right)$$

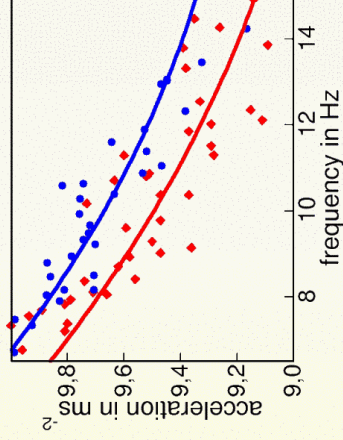
3. Experiment

determine for different amplitudes the critical frequency (acceleration)



1 sphere → systematic errors

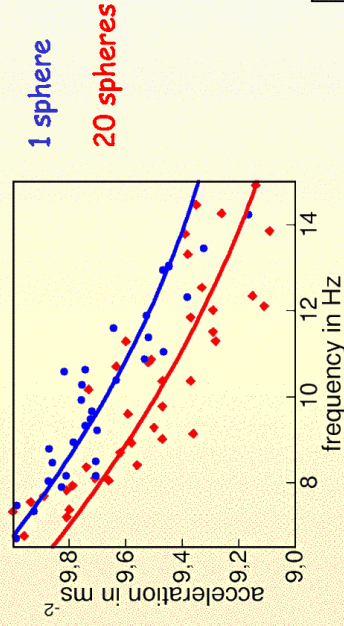
20 spheres



same data in different representation together with best fit

$$f(\omega) = (a_2\omega^2 + a_1\omega + a_0)^{-1}$$

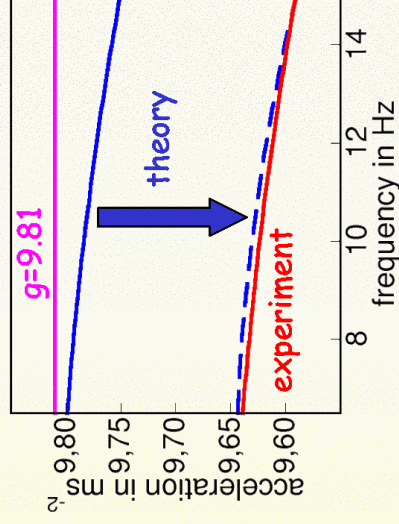
3. Experiment



Eliminate systematic errors by dividing both curves.

amplification factor:

$$\frac{(a_2\omega^2 + a_1\omega + a_0)^{-1}}{(a_1\omega^2 + a_0\omega + a_0)^{-1}}$$



Summary

- fluidization implies separation of particles
- column of viscoelastic spheres:
analytical: There is always a region of amplification for the driving frequency
- MD: agreement
Experiment: agrees too
- The effect has not been confirmed for real (3D) granular systems
- There are MD simulations which describe fluidization for $A\omega^2 < g$

Taguchi, PRL 69, 1367 (93); EPL 24, 203 (93)

Luding et al., PRE 50, 3100 (94)

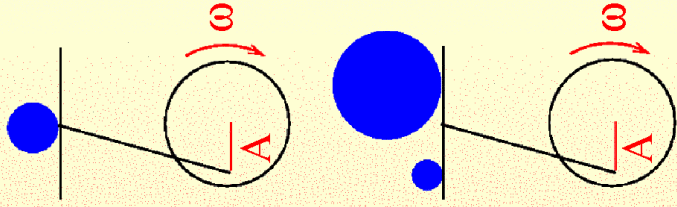
Banker, Mehta, PRA 45, 3435 (92)

Gallas, Herrmann, Sokolowski, J. Physique II 2, 1389 (92)

Claim: $\Gamma \equiv A\omega^2 / g > 1$ is insufficient to describe the onset of fluidization (and possibly other effects too).
Instead, we propose

$$\Gamma \equiv A\omega^2 / g > R^{-1} \quad \text{with} \quad R^{-1} = R^{-1}(\omega)$$

Comment 1: Does Γ describe a single particle?



particle starts to jump at

$$A\omega^2 = g$$

Large particle jumps for smaller frequency.

Reason: Finite size body behaves like a non-linear oscillator, even if the material is linear.

Comment 2: What about gravity?

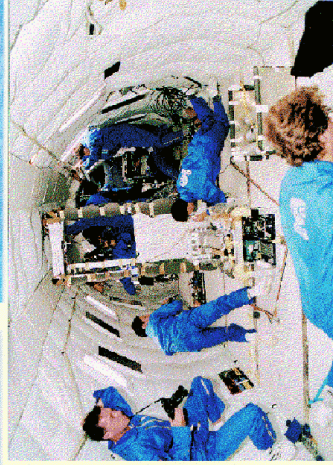
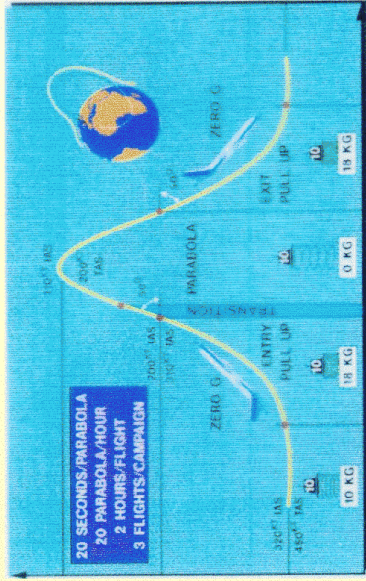
Both condition for the onset of fluidization depend on the value of gravity:

$$\frac{A\omega^2}{g} > 1$$

$$\frac{A\omega^2}{g} > 1 - B_2\omega^2 + B_4\omega^4$$

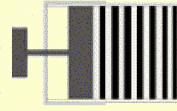
What do we expect in the absence of gravity?

Comment 2: What about gravity?



Comment 2: What about gravity?

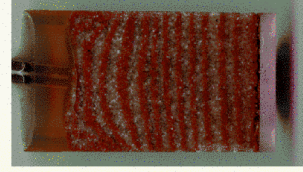
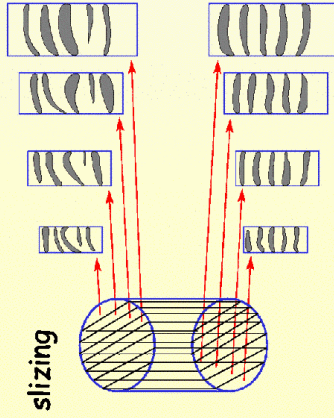
Preparation: layers of black and white sand.
Fixed by a piston



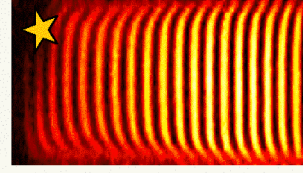
During the shaking: Fixating by liquid resin



Ratkai. Pow. Techn.'76

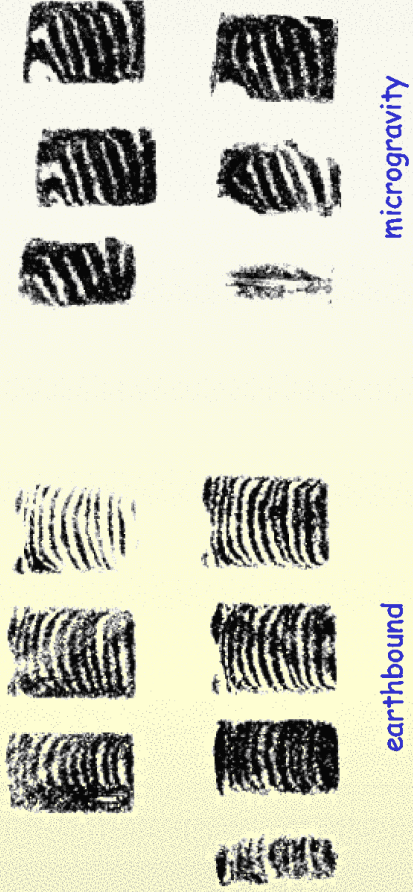


Pöschel, Freund'96



Ehrichs et al'95

Comment 2: What about gravity?



microgravity

$$\frac{A\omega^2}{g} \ll 1$$

earthbound

Pöschel, Freund'96

Comment 2: What about gravity?

