
Overview of a Spin-Liquid Candidate, $\text{Ba}_3\text{CuSb}_2\text{O}_9$

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Work done in collaboration with Patrick Lee



KITP, September 13, 2012

Zhou et. al., “**Spin liquid state in the $S = 1/2$ triangular lattice $Ba_3CuSb_2O_9$,**” Phys. Rev. Lett. **106**, 147204 (2011)

- low T linear specific heat, C_m
- magnetic susceptibility, χ
- no magnetic order down to ~ 0.2 K

Nakatsuji et. al., “**Spin-Orbital Short-Range Order on a Honeycomb-Based Lattice,**” Science **336**, 559 (2012)

- dumbbell structure
- χ , C_m , inelastic scattering
- Jahn-Teller distortions

Balents, “**The Impact of Ionic Frustration on Electronic Order,**” Science **336**, 547 (2012)

- dumbbells save frustration (from orb ordering)

Quilliam et. al., “**Singlet Ground State of the Quantum Antiferromagnet $Ba_3CuSb_2O_9$,**” Phys. Rev. Lett. **109**, 117203 (2012)

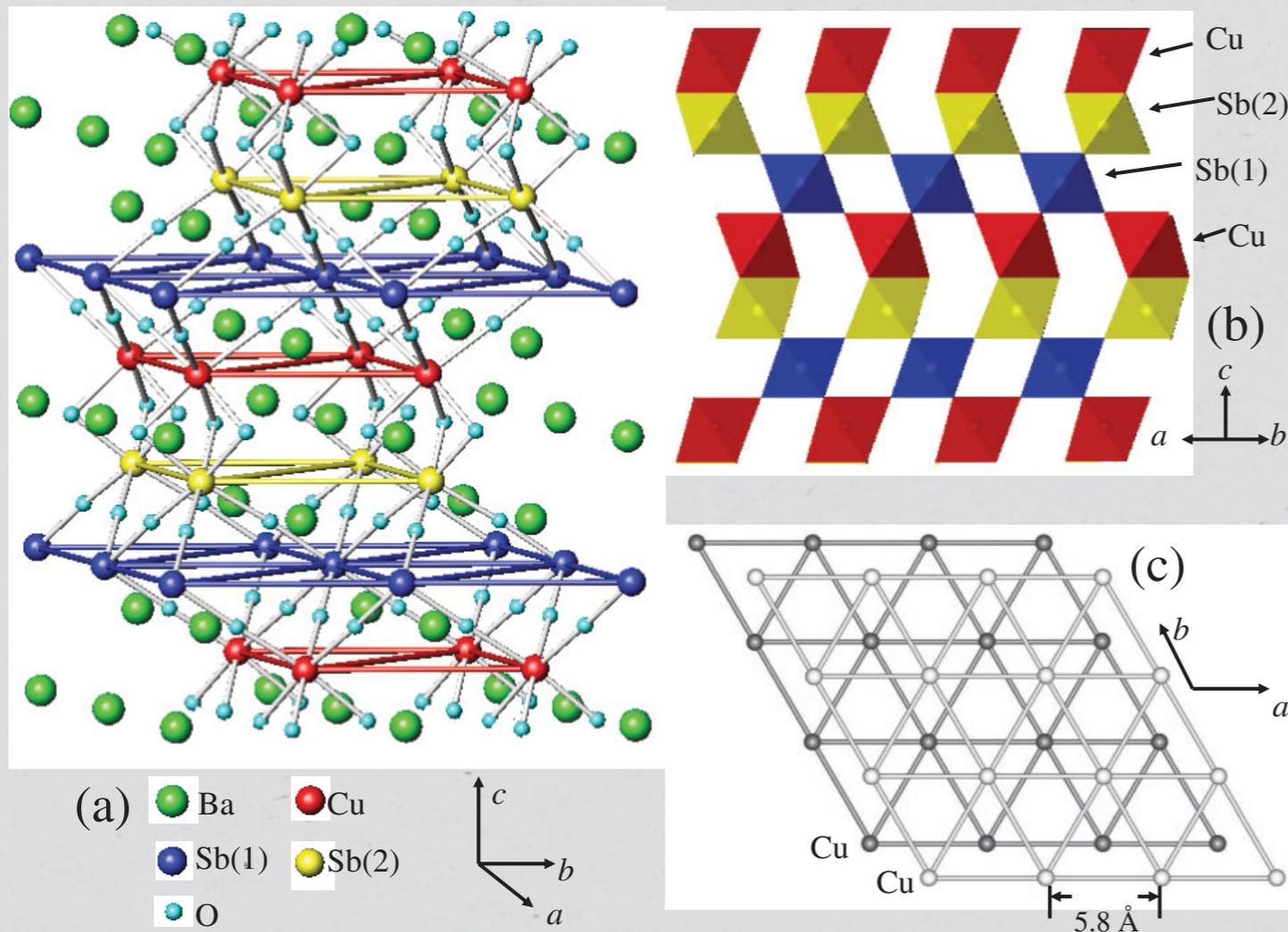
- NMR \rightarrow intrinsic χ
- relaxation $1/T_1 \sim$ excitation gaps
- singlets and (intrinsic) ‘orphan spins’

Nasu, Ishihara, “**Dynamical Jahn-Teller Effect in Spin-Orbital Coupled System,**” arXiv: 1209.0239 [cond-mat]

- competing orbital/AFM orders $\sim \rightarrow$ SO resonant states

Ba₃CuSb₂O₉ Structure

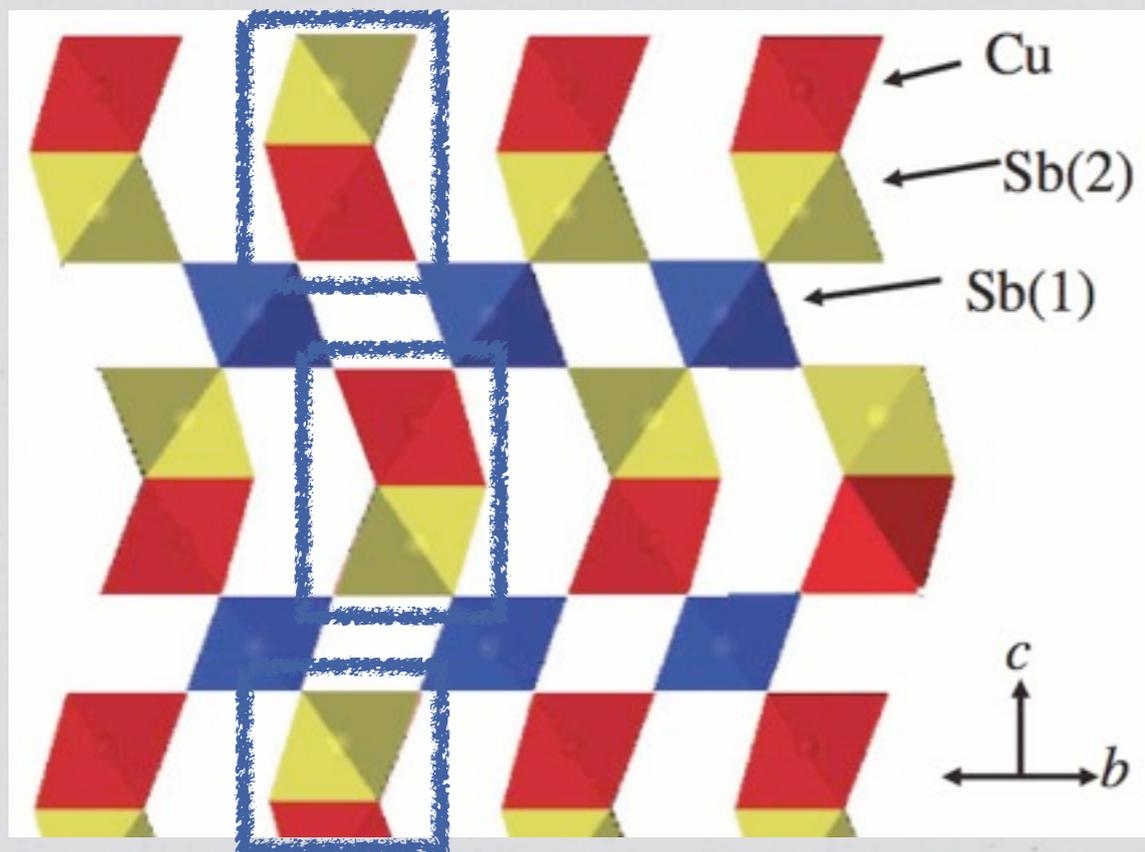
First, as proposed by Zhou et al:



- ▶ triangular layers of Cu²⁺
- ▶ isolated by 2 non-magnetic Sb⁵⁺ layers
- ▶ Cu sits in an octahedron of O²⁻: single hole is in degenerate e_g orbitals
- ▶ possible Jahn-Teller distortions may split the e_g orbitals

Ba₃CuSb₂O₉ Structure, cont'd

BUT! Cu²⁺/Sb⁵⁺ 'dumbbells' are electric dipoles
Nakatsuji et. al. measure no (discernible) pyroelectricity

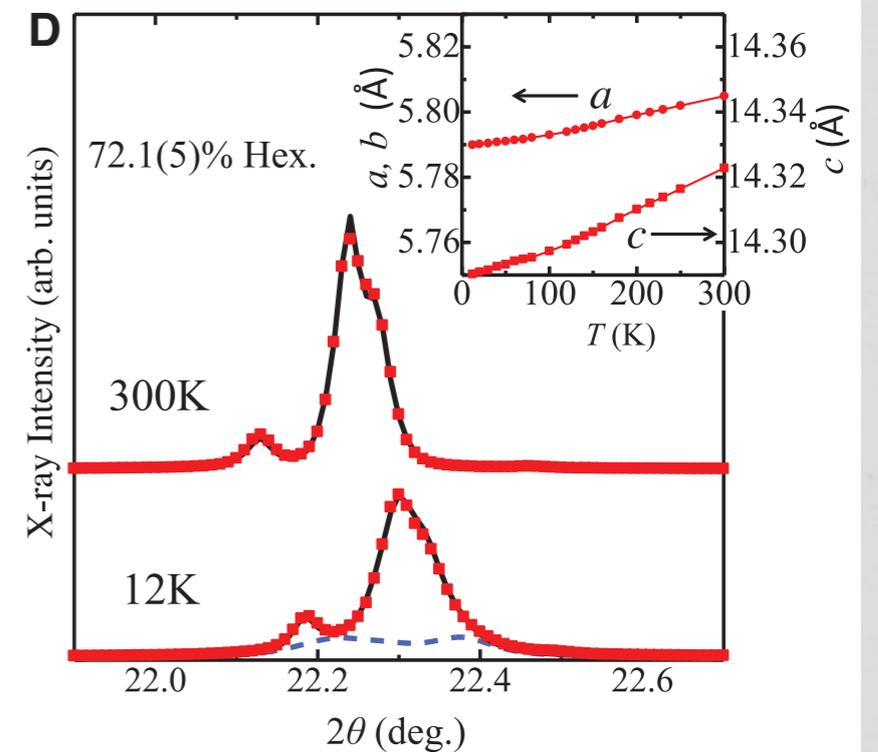
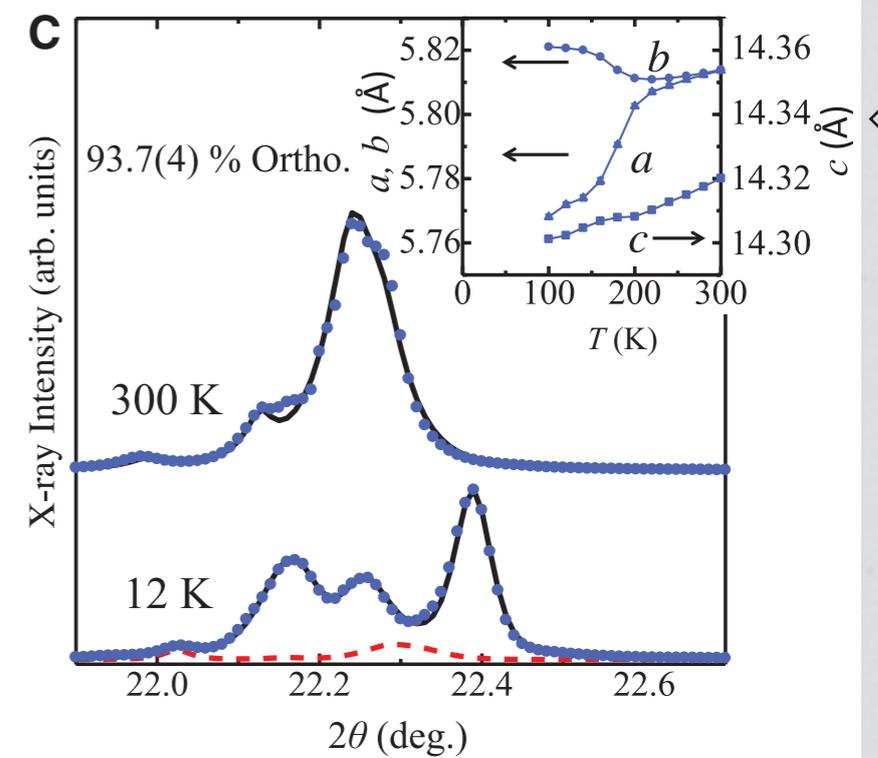
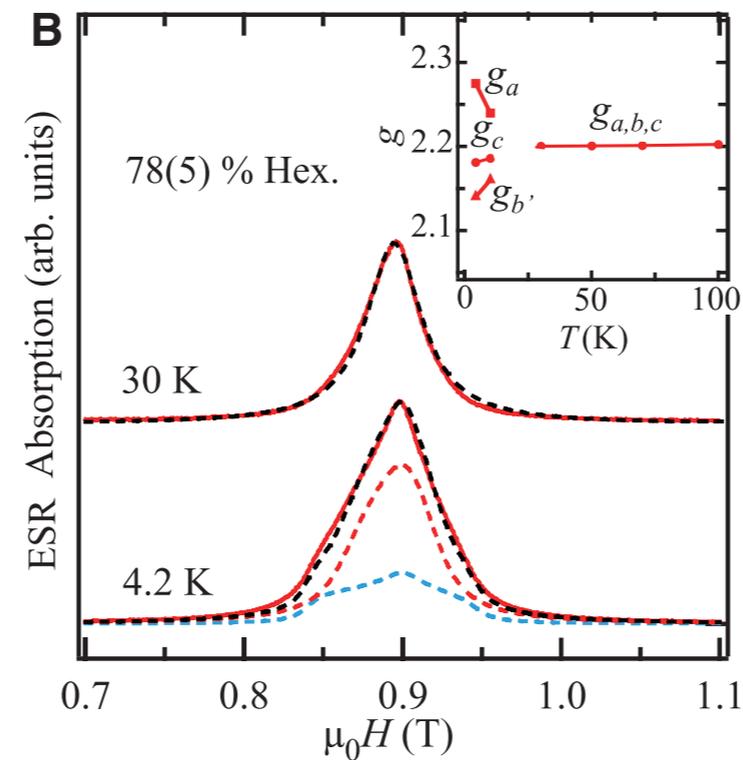
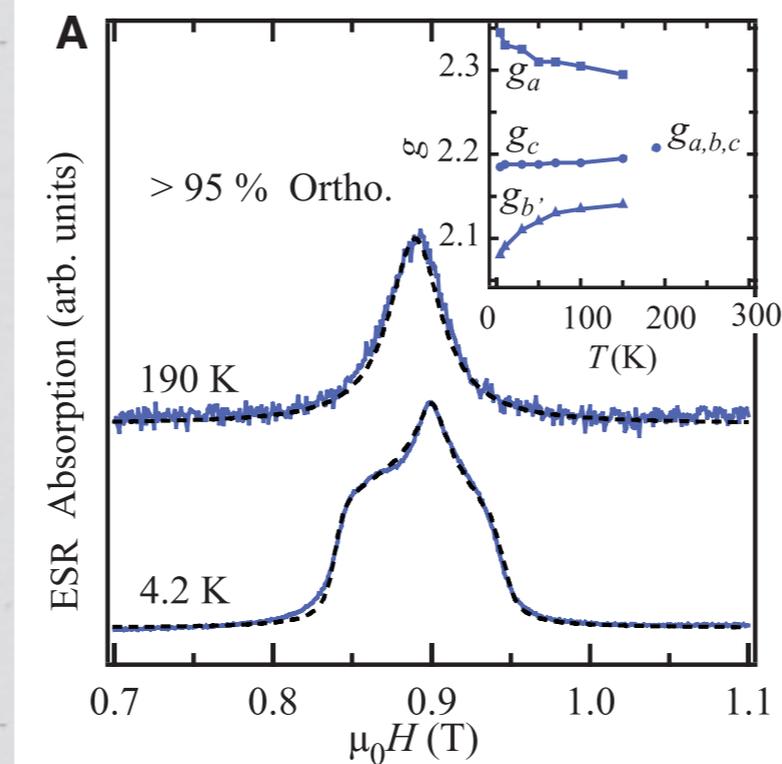


- Cu²⁺/Sb⁵⁺ dumbbells ~ Ising 'spins' on a triangular lattice: frustrated!
- are magnetic layers still isolated?
- what is the superexchange for same-plane vs out-of-plane Cu's?

Net electric dipoles should cancel.

Jahn-Teller distortions

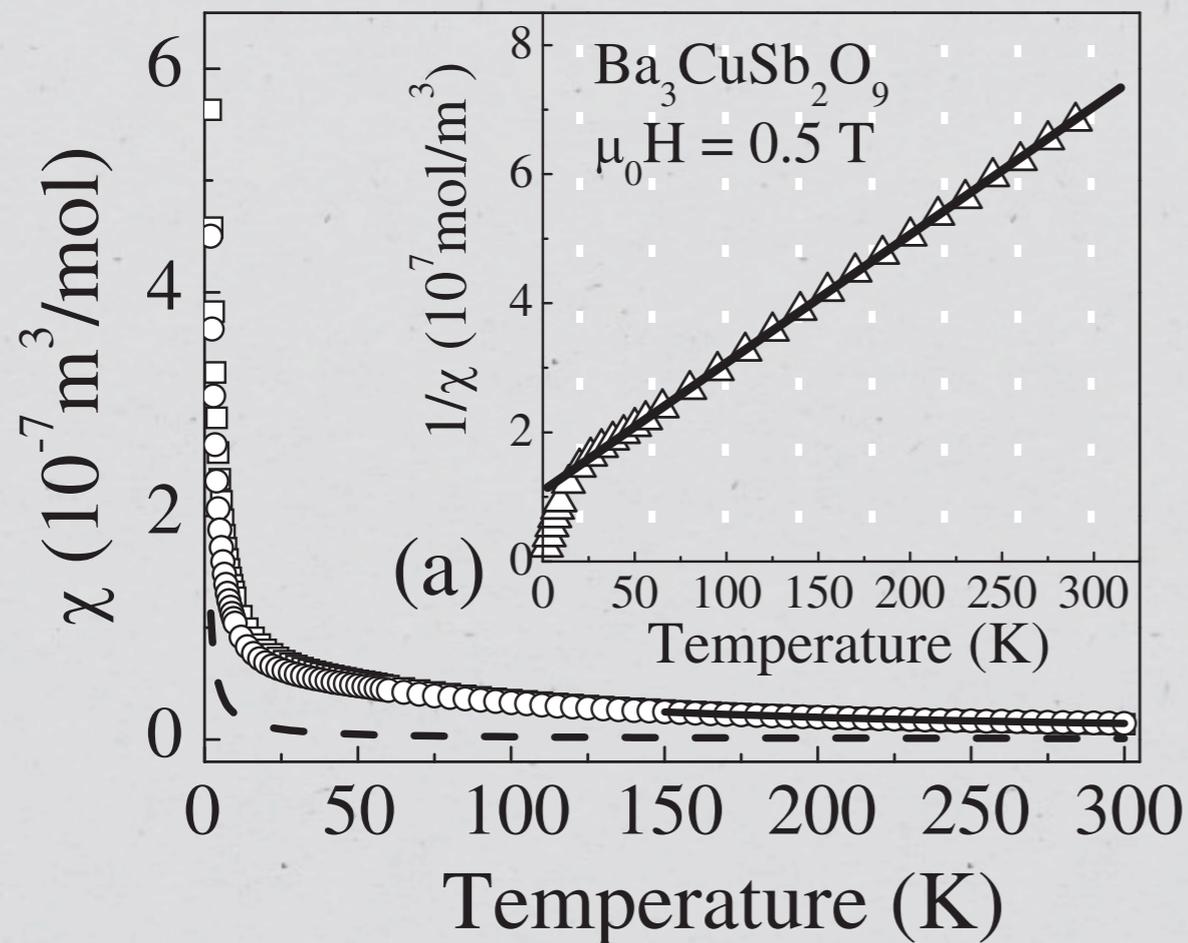
- ▶ for >8% off-stoichiometric Cu/Sb samples, find a 200K JT transition
- ▶ otherwise, JT 'remnants' <30K



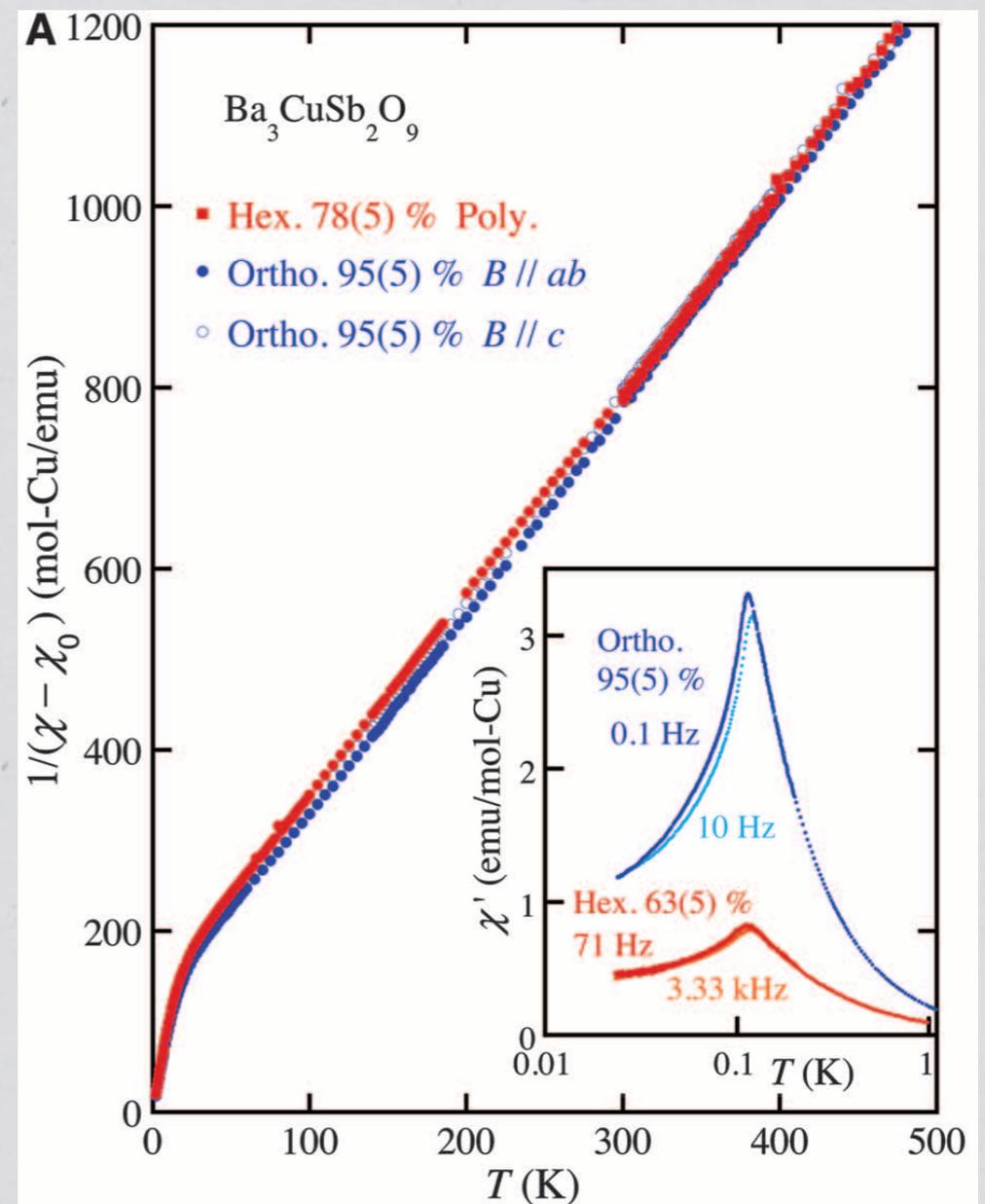
Nakatsuji et al (2012)

Magnetic Probes

- ▶ No magnetic ordering
- ▶ high T CW fit: $\theta_{CW} \sim -55\text{K}$ (afm)
- ▶ possible low T freezing of ortho phase orphan spins



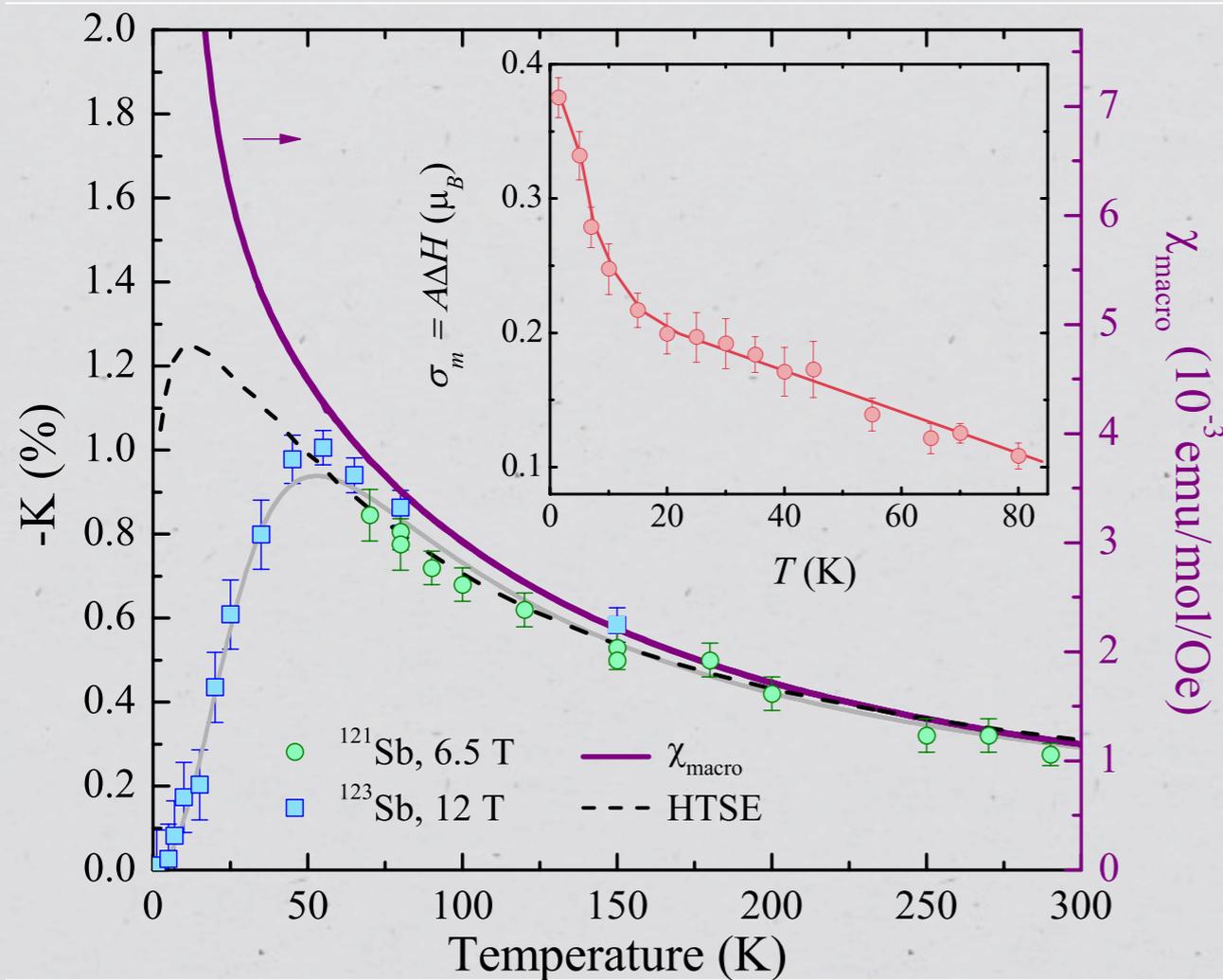
Zhou et al (2011)



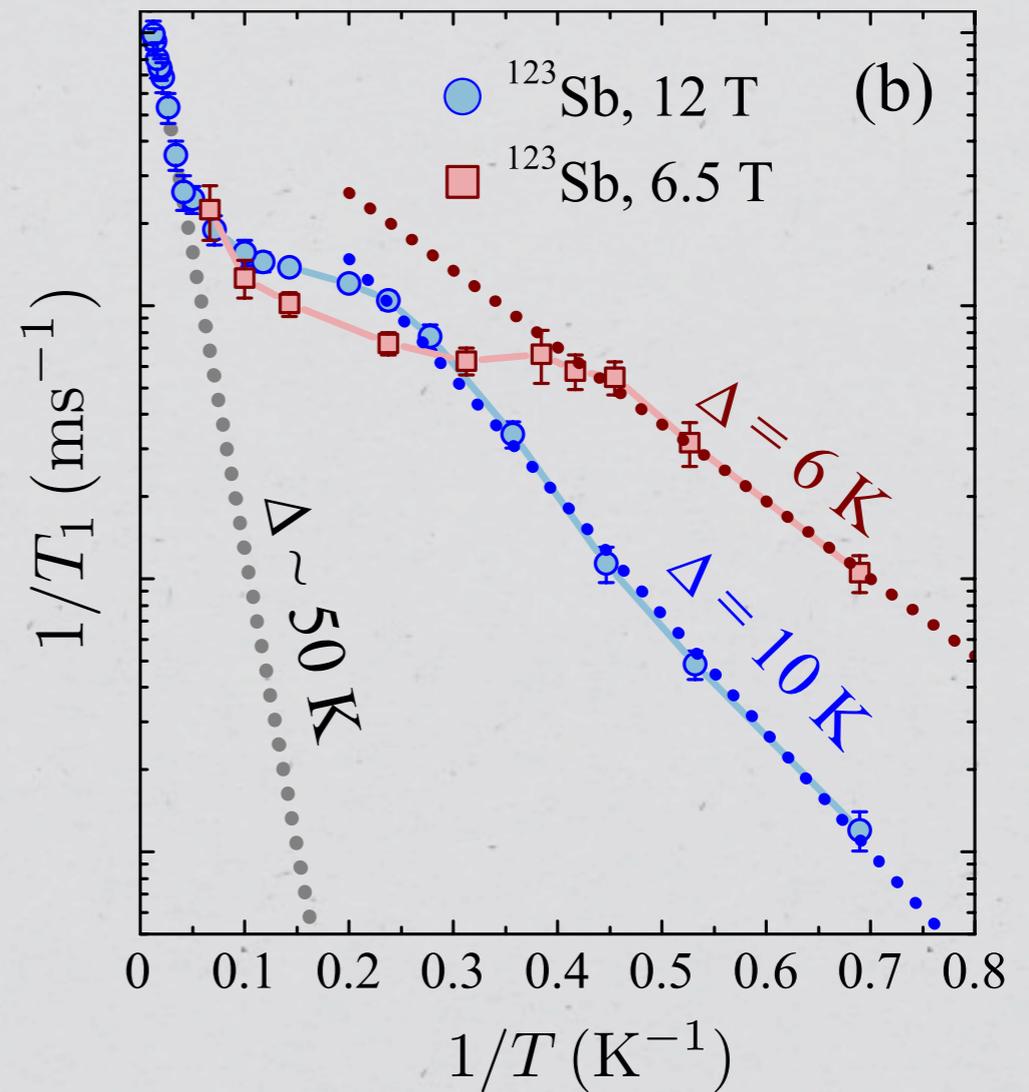
Nakatsuji et al (2012)

NMR (Quilliam et. al. 2012)

- ▶ ~singlet correlations at 50K (max) + diverging 'orphan spins'
- ▶ field dep (distribution) of gaps at lower energy
- ▶ extreme line broadening at low T

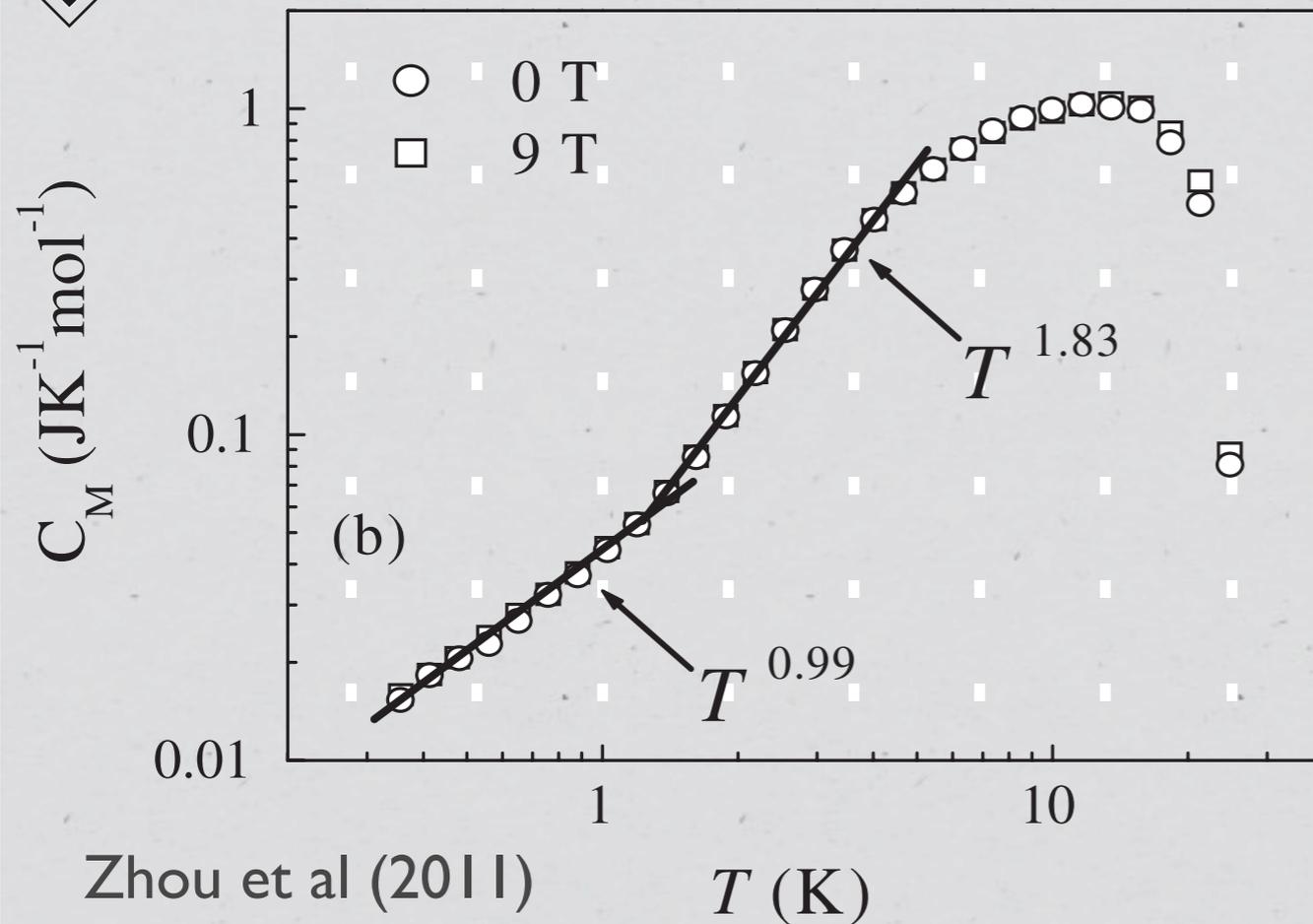


NMR lineshift, $K \sim \chi_{\text{int}}$



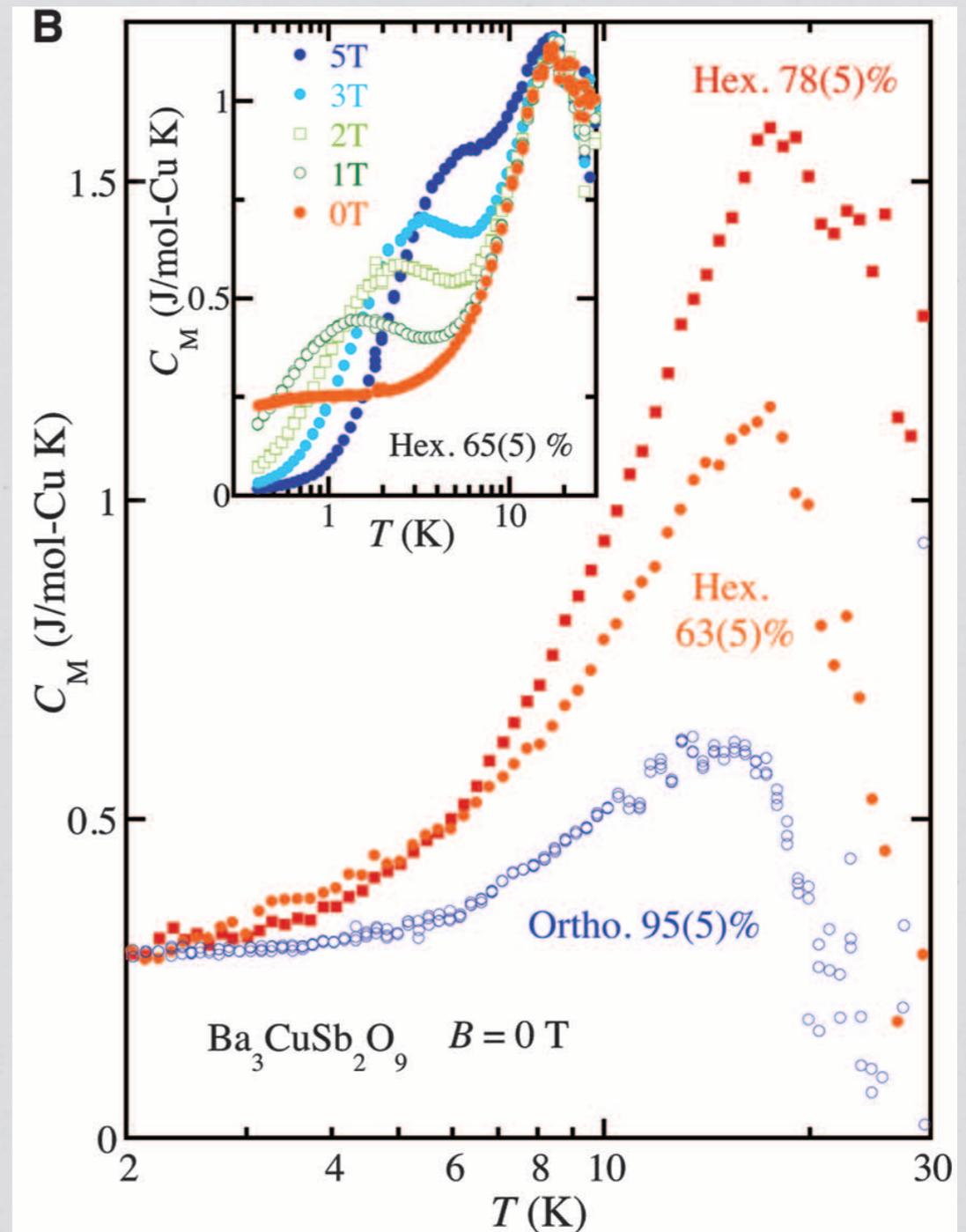
magnetic relaxation rate

Specific Heat



- broad max around 10-30 K $\approx J/2$
- more d.o.f. in hex phase
- local moments can be frozen in field

Nakatsuji et al (2012)

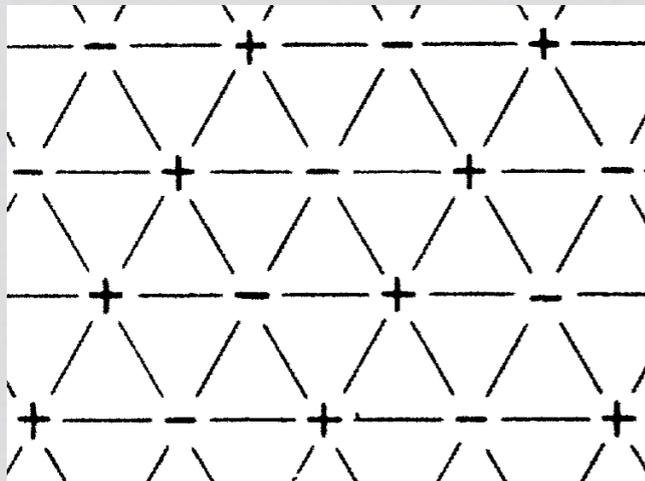


'Dumbbell' Order

- ▶ electric dipoles 'dumbbells' ~ Ising pseudospin;
Hamiltonian:

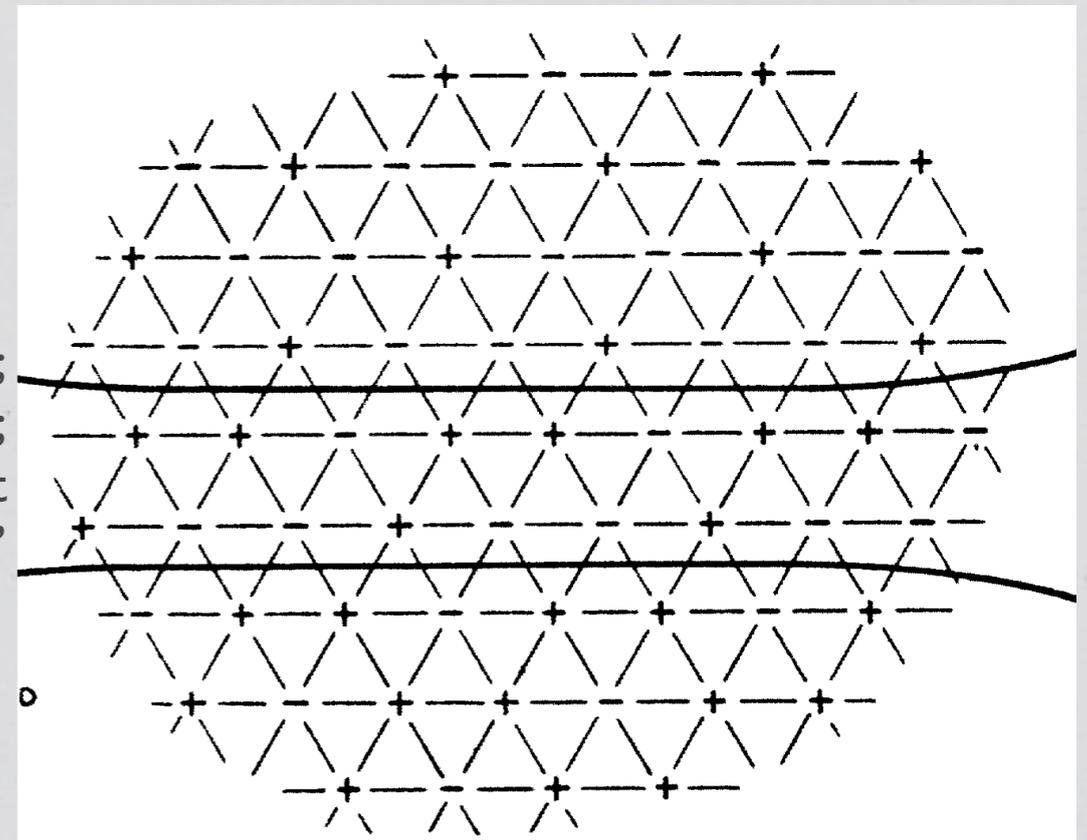
$$H = \sum_{\langle ij \rangle} J_{afm} S_i S_j - \sum_{\langle\langle ik \rangle\rangle} J_{fm} S_i S_k$$

stripes



honeycomb

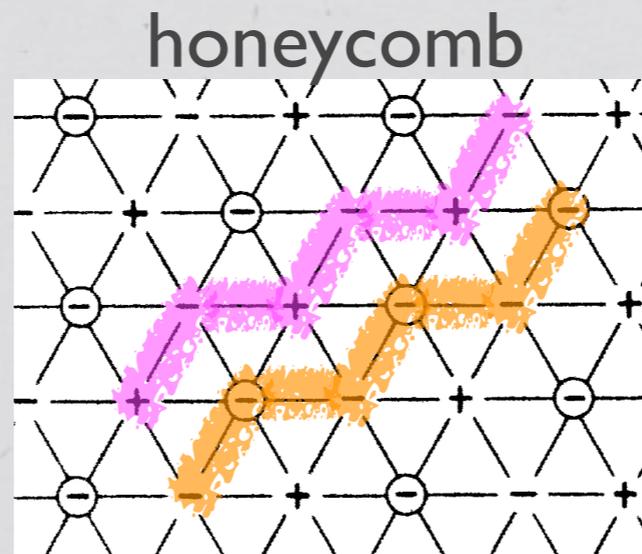
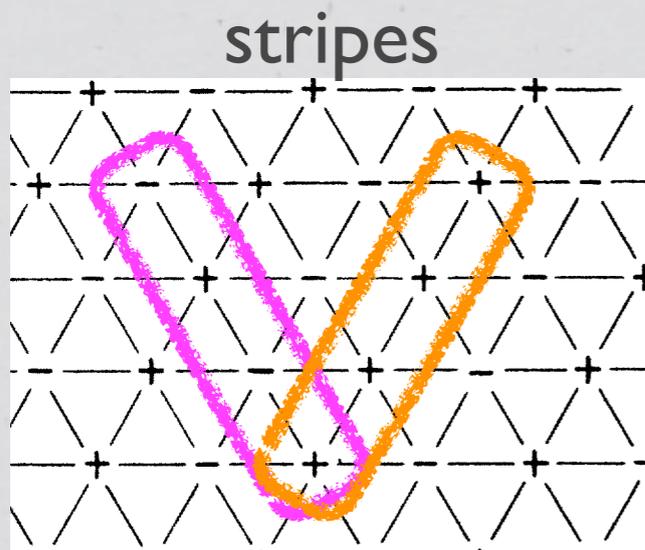
entropically
favoured:
zero E domain walls;
many domains;
honeycomb-like, but
rather 'disordered'



- ▶ agrees with "3 sub-lattice,
ferrielectric short-range order"
reported by Nakatsuji (2012)

'Dumbbell' Defects

- ▶ off-stoichiometry of Cu/Sb translates to a vacancy in the effective Ising lattice
- ▶ $\sim 1/9$ extra/fewer Cu means 1 vacancy per hexagon; should consider defect correlation effects



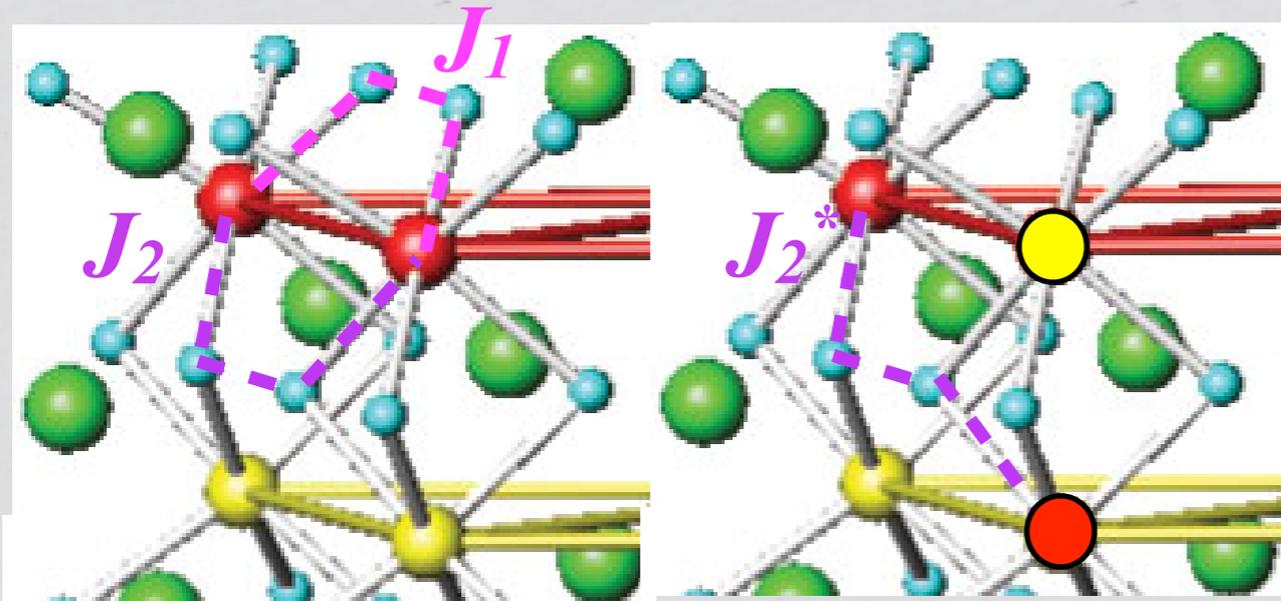
- ▶ single defects in honeycomb prefer maj. spin sites, \therefore net polarization
- ▶ in stripes, net zero polarization

- ▶ minimum energy 'defect line' in **stripes** when:

$$J_{afm} < nJ_{fm}$$

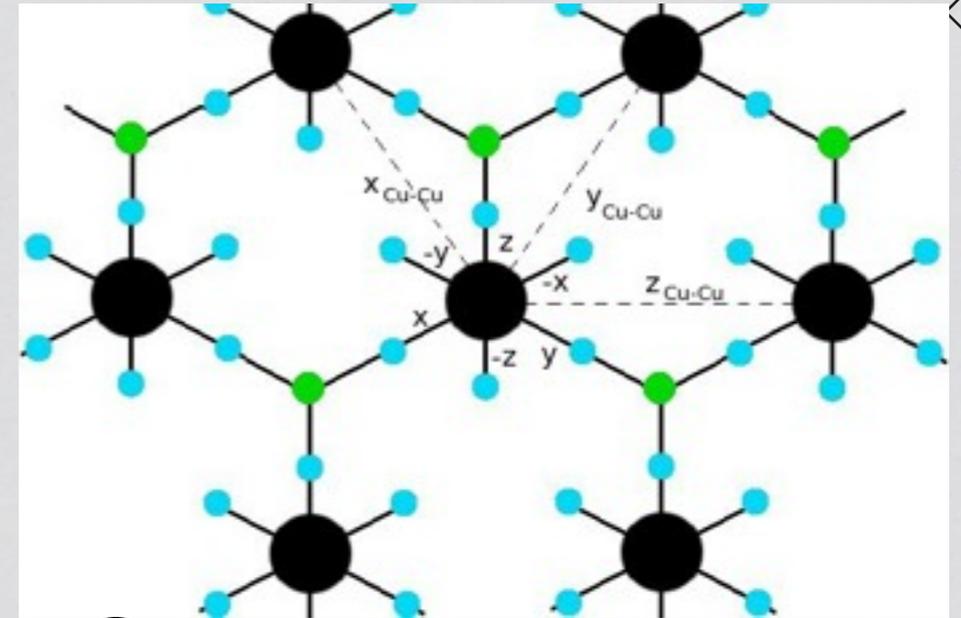
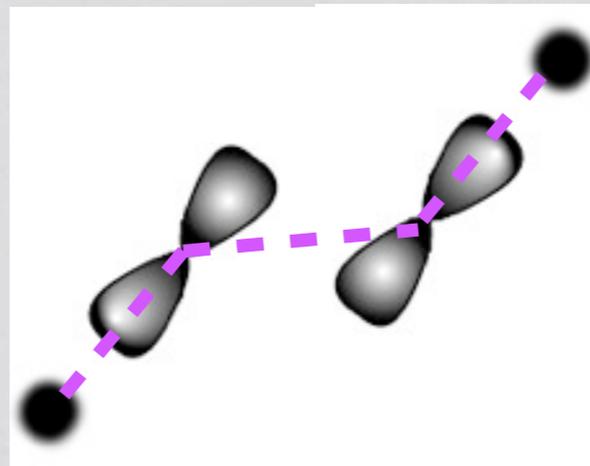
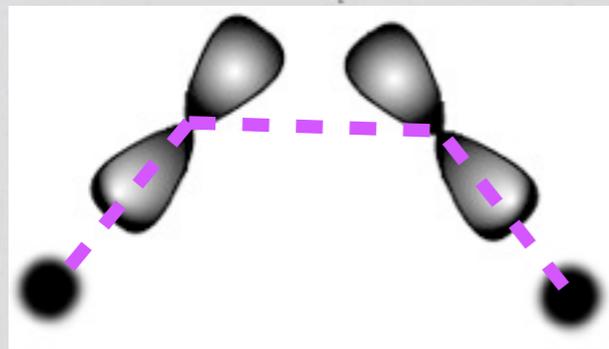
- ▶ with enough defects, expect a transition to (JT distorted) stripe phase

Cu-O-O-Cu Superexchange



parallel

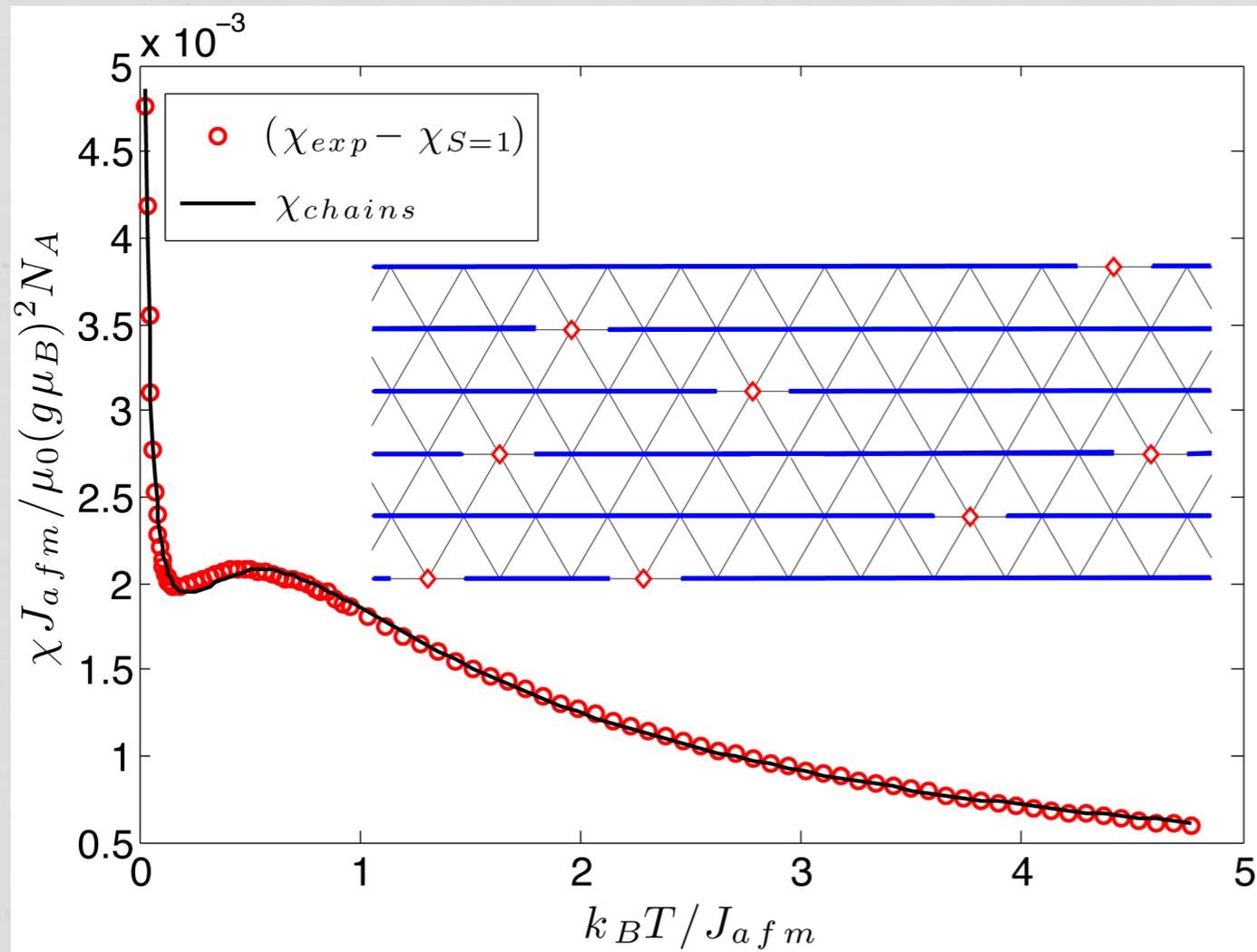
anti-parallel



- $\text{Cu}^{2+}/\text{Sb}^{5+}$ dumbbell
- O^{2-}
- Sb^{5+} : Sb(1) sites shown, Sb(2) sites sit directly under Cu sites

overlap is of *same order* and J_2^* (as opposed to arguments in Nakatsuji et. al.)

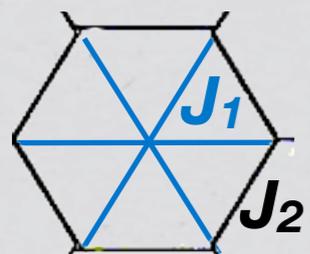
Spin chains in stripe order



- ▶ previously motivated by a triangular lattice of Cu & defects
- ▶ MF favours a uniform JT distortion leading to decoupled chains
- ▶ remarkable agreement in χ with $\sim 20\%$ $s=1/2$ 'orphan spins'
- ▶ BUT, JT distortion is not always present: what about hex phase?

arXiv:1202.5655 [cond-mat]

Disordered Honeycomb



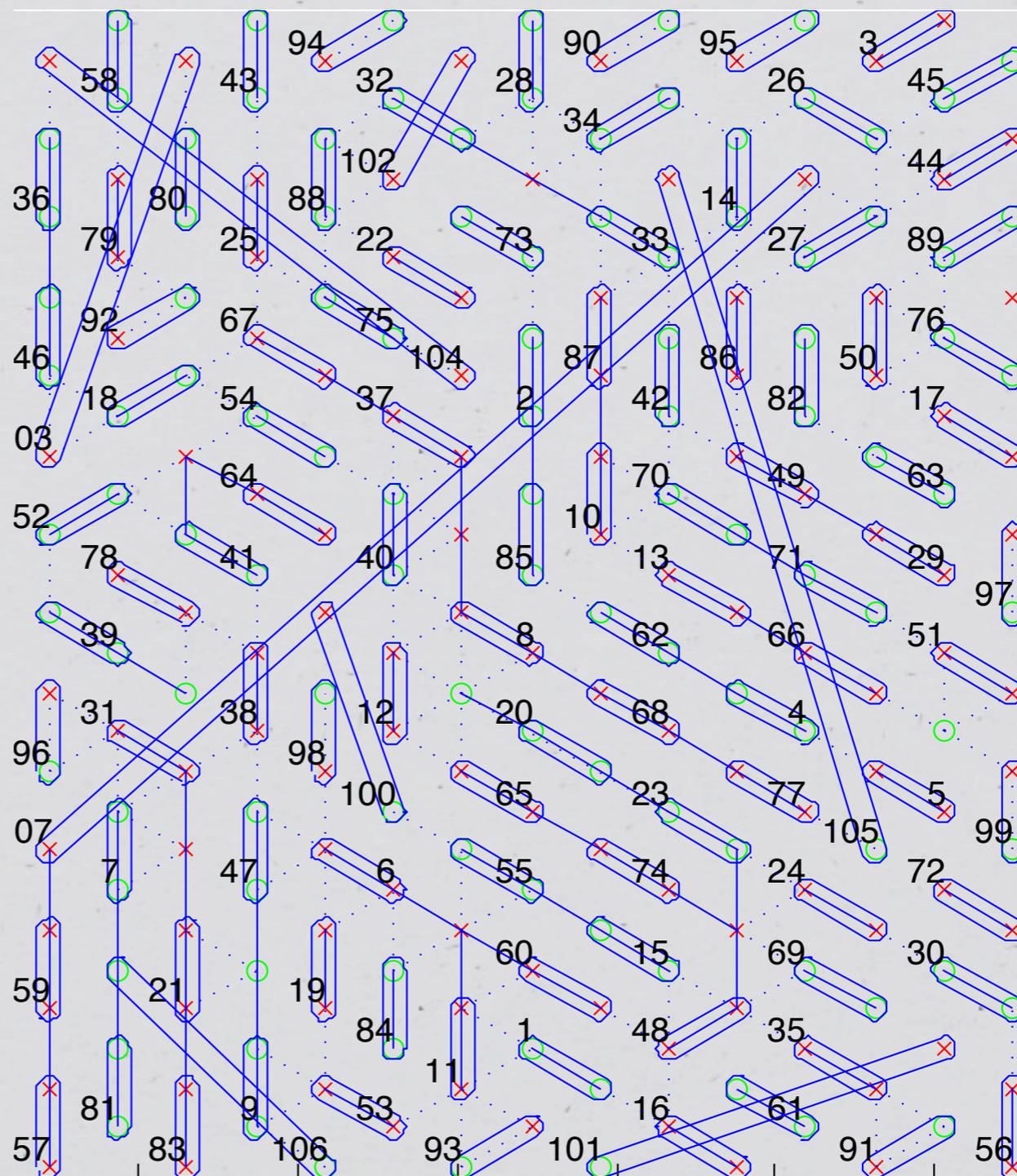
▶ with $J_1 \sim J_2/2$
when orbitals are
aligned

▶ relative alignment of e_g
orbitals far more important

▶ aligned orbitals naturally
decouple rows; relieve
frustration

▶ even in disordered phase,
expect (zigzag) spin chains

example output of
(unoptimized) singlet pairing ▶



Ba₃CuSb₂O₉ Summary

- No magnetic order.
- typical afm exchange of ~50-60 K
- broad peak in χ_{int} , C_m : singlets, spin chains?
- ~15-20% 'orphan spins' (intrinsic?)
- dumbbell order of in- and out-of-plane Cu on a triangular lattice ~ disordered honeycomb
- JT distorted in off-stoichiometry samples; remnant JT at low T in hex phase

