

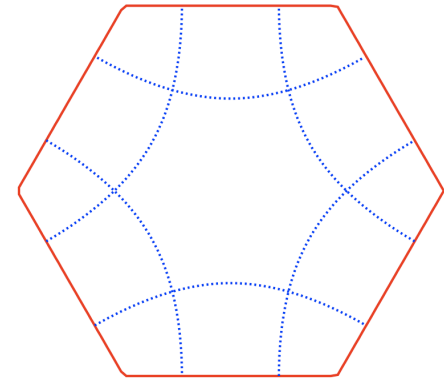
Mott Metal to Spin-Liquid Transition

MPA Fisher (with **O. Motrunich**, **Ryan Mishmash**, Matt Block, Ivan Gonzalez, Donna Sheng, **Roger Melko**)

KITP Conference on Exotic Phases of Frustrated Magnets 10/7/12

Mott metal-insulator transition in 2d

- Can the Mott transition be a 2nd order T=0 quantum phase transition?
- If yes, what is nature of the insulating phase?
- What is the nature of the quantum critical point?



- Here, seek to address using controlled numerical calculations in quasi-1d systems (DMRG, VMC, bosonization of gauge theory)
- Next talk by Senthil – Analytic field theories and scaling approach to 2d Mott transition

(Some) References

Theory of a continuous Mott transition in two dimensions
T. Senthil, PRB (2008).

Mott transition between a spin-liquid insulator and a metal in three dimensions
D. Podolsky, A. Paramekanti, Yong Baek Kim, **T. Senthil**, PRL (2009).

A controlled expansion for certain non-Fermi liquid metals
D. F. Mross, J. McGreevy, H. Liu, **T. Senthil**, PRB (2010)

Decohering the Fermi liquid: A dual approach to the Mott Transition
D. F. Mross, **T. Senthil**, PRB (2011)

Universal transport near a quantum critical Mott transition in two dimensions
W. Witczak-Krempa, P. Ghaemi, **T. Senthil**, Y.B. Kim, arXiv:1206.3309

Spin Bose-Metal phase in a spin-1/2 model with ring exchange on a two-leg triangular strip,
D.N. Sheng, O.I. Motrunich, MPAF, PRB (2009)

Spin Bose-Metal and Valence Bond Solid phases in a spin-1/2 model with ring exchanges
on a four-leg triangular ladder, M.S.Block, D.N. Sheng. O.I. Motrunich, MPAF, PRL (2011)

Mott metal to Spin Liquid transition in quasi-1d electron models,
R.V. Mishmash, I. Gonzalez, R. Melko, O. Motrunich, MPAF (in the works!)

2d Metal and Mott insulator phases

2d Hubbard model, one electron/site:

$$\mathcal{H} = -t \sum_{\langle ij \rangle} [c_{i\alpha}^\dagger c_{j\alpha} + h.c.] + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

Metal for small U/t (band theory)

Mott insulator for large U/t , electrons localize

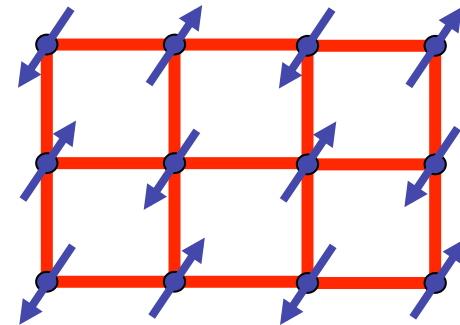


Sir Neville Mott



Mott Insulator \rightarrow effective spin Hamiltonian

$$H_{spin} = J \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j$$



(Almost!) Invariably the Spins order (AFM) in the Mott insulator

Mott metal-insulator *transition*??



Nature of $T=0$ transition?

- First order, always possible
- 2nd order (continuous) is problematic

**For continuous Metal to AFM insulator transition
need insulating behavior and magnetism to turn on **together**. *Not generic!***

(Simplest scenario - intervening AFM ordered metal, between metal and AFM insulator)

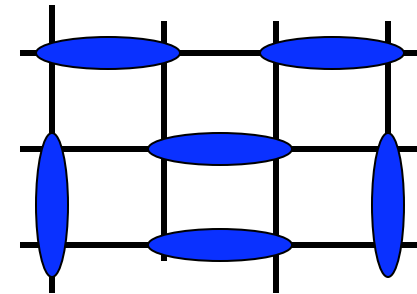
But what if Mott insulator is in a *Spin Liquid* (no magnetic order)?

**Is a 2nd order metal to spin-liquid Mott transition possible?
And if so, **which** spin-liquid?**

Not all 2d Spin liquids alike!

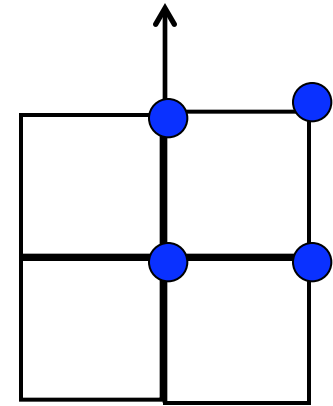
Topological Spin Liquids

- Fully gapped
- Fractionalized excitations, eg $s=1/2$ spinon



Algebraic Spin Liquids

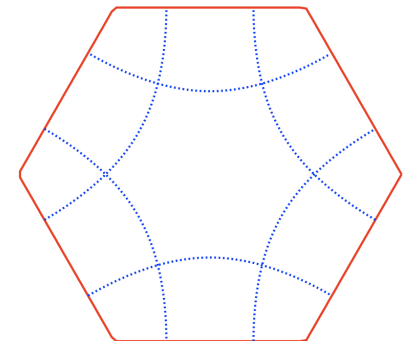
- Stable gapless phase
- Power-law correlations at finite set of discrete momenta



“Spin Bose-Metals”

- Stable Gapless phase
- Spin correlations singular on **surfaces** in momentum space

“Bose Surfaces”

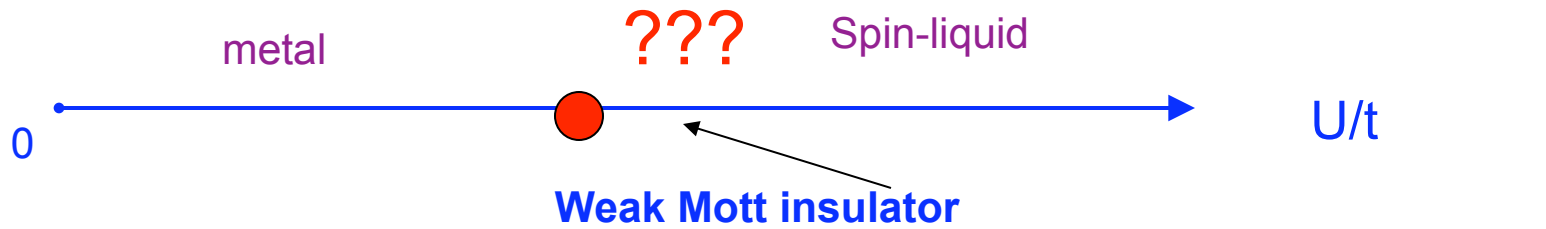


Continuous Mott transition: Into which Spin-liquid?

If metal-insulator transition is 2nd order, then have a

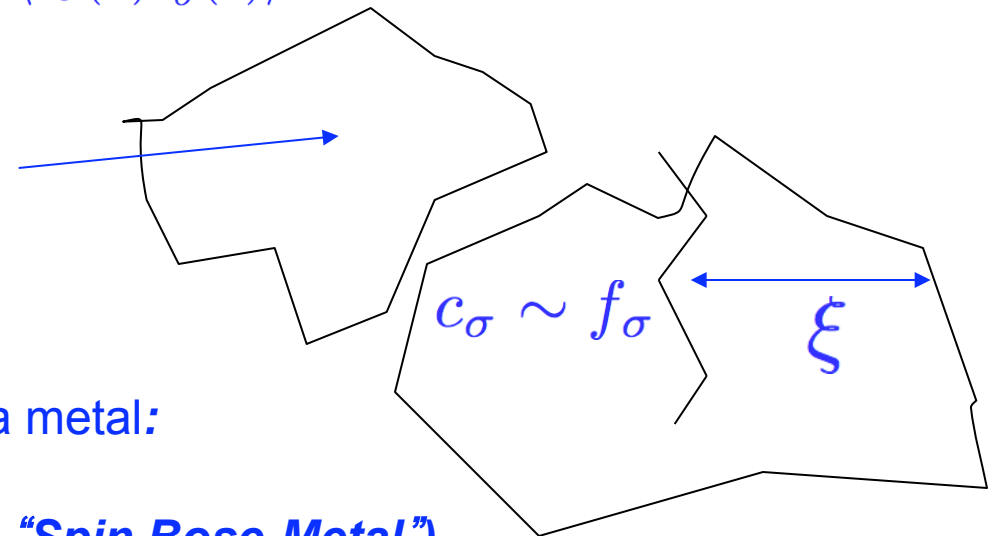
Weak Mott insulator

Long charge correlation length



$$\langle c_\sigma(x)c_\sigma^\dagger(0) \rangle \sim e^{-x/\xi} \quad \xi \gg a$$

Inside correlation region electrons do not "know" they are insulating

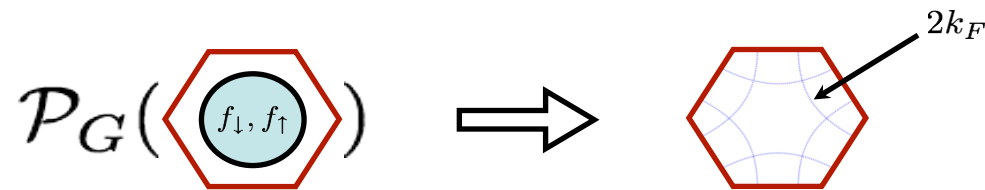


Candidate spin liquid that "resembles" a metal:

Gutzwiller projected Fermi sea (a "Spin Bose-Metal")

Gutzwiller projected Fermi sea (spin Bose-metal)

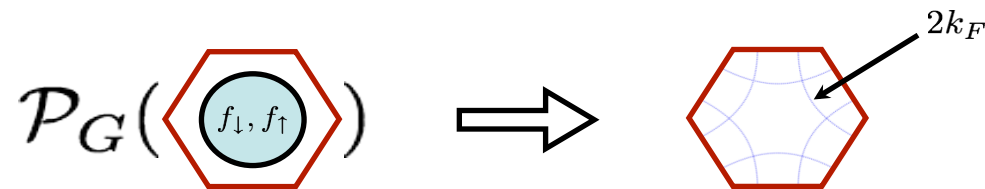
Construct spin-wavefunction from free Fermi gas with one Fermion/site



$$\Psi_0 = c_1 \begin{array}{|c|c|c|c|c|} \hline \uparrow & \downarrow & \uparrow & \uparrow & \\ \hline \downarrow & \uparrow & \uparrow & \downarrow & \downarrow \\ \hline \downarrow & & \uparrow & \uparrow\downarrow & \downarrow \\ \hline \downarrow & \downarrow & \uparrow\downarrow & \downarrow & \downarrow \\ \hline \uparrow & \downarrow & \downarrow & \downarrow & \uparrow \\ \hline \end{array} + c_2 \begin{array}{|c|c|c|c|c|} \hline \uparrow & \downarrow & \uparrow & \uparrow & \uparrow \\ \hline \downarrow & \uparrow & \uparrow & \downarrow & \downarrow \\ \hline \downarrow & \uparrow & \uparrow & \downarrow & \downarrow \\ \hline \downarrow & \downarrow & \downarrow & \downarrow & \downarrow \\ \hline \uparrow & \downarrow & \downarrow & \downarrow & \uparrow \\ \hline \end{array} + c_3 \begin{array}{|c|c|c|c|c|} \hline \uparrow & \downarrow & \uparrow & \uparrow & \downarrow \\ \hline \downarrow & \uparrow & \uparrow\downarrow & \downarrow & \uparrow \\ \hline \downarrow & \downarrow & \uparrow & & \downarrow \\ \hline \downarrow & & \downarrow & \downarrow & \downarrow \\ \hline \uparrow & \uparrow\downarrow & \downarrow & \downarrow & \uparrow \\ \hline \end{array} + \dots$$

Gutzwiller projected Fermi sea (spin Bose-metal)

Construct spin-wavefunction from free Fermi gas with one Fermion/site



$$\Psi = c_1 \begin{array}{|c|c|c|c|c|} \hline \uparrow & \downarrow & \uparrow & \uparrow & \\ \hline \downarrow & \times & \uparrow & \downarrow & \downarrow \\ \hline \downarrow & & \uparrow & \downarrow & \downarrow \\ \hline \downarrow & \uparrow & \downarrow & \downarrow & \downarrow \\ \hline \uparrow & \downarrow & \downarrow & \downarrow & \uparrow \\ \hline \end{array} + c_2 \begin{array}{|c|c|c|c|c|} \hline \uparrow & \downarrow & \uparrow & \uparrow & \uparrow \\ \hline \downarrow & \uparrow & \uparrow & \downarrow & \downarrow \\ \hline \downarrow & \uparrow & \uparrow & \downarrow & \downarrow \\ \hline \downarrow & \downarrow & \downarrow & \downarrow & \downarrow \\ \hline \uparrow & \downarrow & \downarrow & \downarrow & \uparrow \\ \hline \end{array} + c_3 \begin{array}{|c|c|c|c|c|} \hline \uparrow & \downarrow & \uparrow & \uparrow & \downarrow \\ \hline \downarrow & \uparrow & \downarrow & \downarrow & \uparrow \\ \hline \downarrow & \downarrow & \uparrow & \downarrow & \downarrow \\ \hline \downarrow & \downarrow & \downarrow & \downarrow & \downarrow \\ \hline \uparrow & \uparrow & \downarrow & \downarrow & \uparrow \\ \hline \end{array} + \dots$$

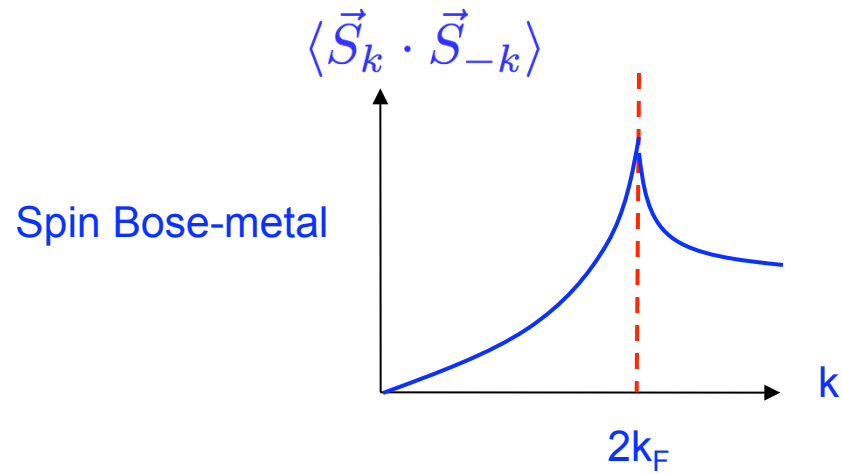
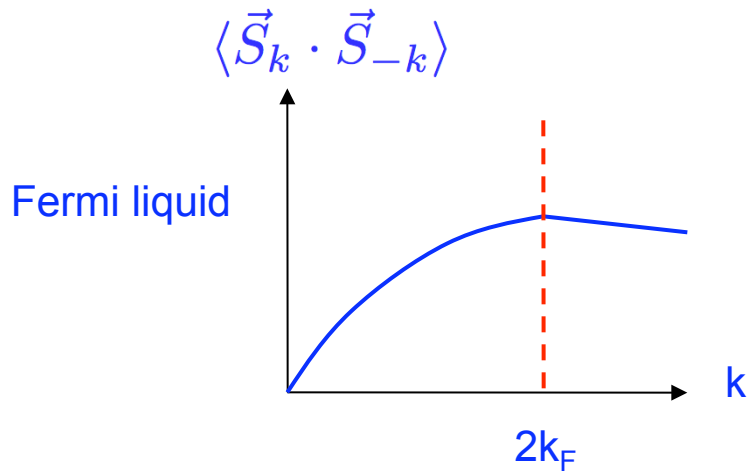
Quantum Spin Liquid Wavefunction

$$\Psi = P_G \Psi_0 \quad \text{Throw away all parts of wf with empty/doubly-occupied sites}$$

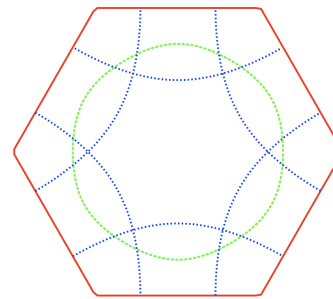
Effective field theory: Fermionic spinons coupled to U(1) gauge field

Phenomenology of Spin Bose-Metal (from wf and Gauge theory)

Singular spin structure factor at “ $2k_F$ ” in Spin Bose-Metal
(more singular than in Fermi liquid metal)



$2k_F$ “Bose surface” in
triangular lattice Spin Bose-Metal



Experimental candidate Spin Bose-Metals

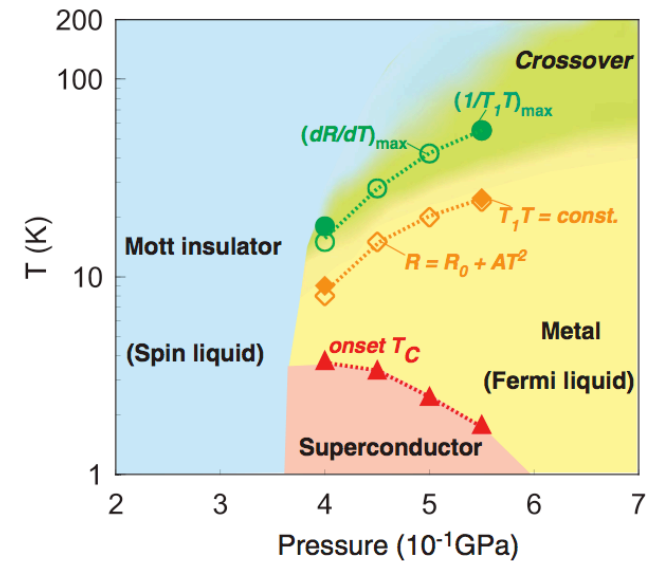
Organic triangular-lattice Mott materials

$k\text{-(ET)}_2\text{Cu}_2(\text{CN})_3$ Kanoda et. al. PRL 91, 177001 (2005)

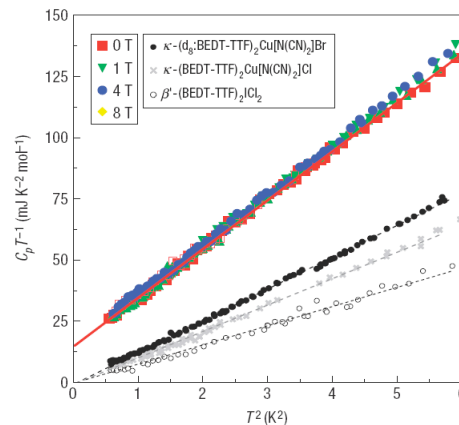
$\text{EtMe}_3\text{Sb}[\text{Pd}(\text{dmit})_2]$ Ito, Kato et. al. PRB 77, 104413 (2008)

- Hubbard at half-filling
- Weak Mott insulator – metal under pressure
- No magnetic order for $T \ll J$
- Pauli spin susceptibility

Motrunich (2005) , S. Lee and P.A. Lee (2005)
 “spinon Fermi sea” proposed

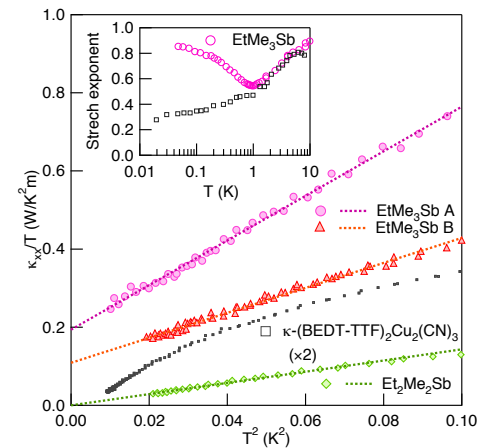


Metallic specific heat, $C \sim T$, Wilson ratio of order one



S. Yamashita et al (2008)

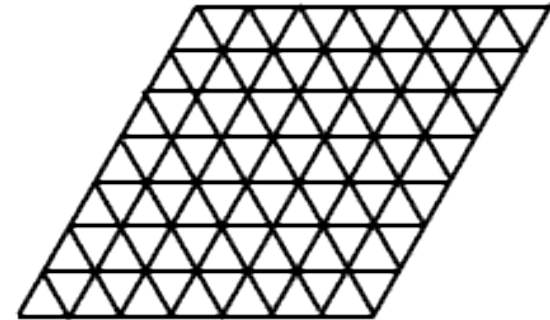
Thermal conductivity $K/T \sim \text{const}$ in $(\text{dmit})_2$



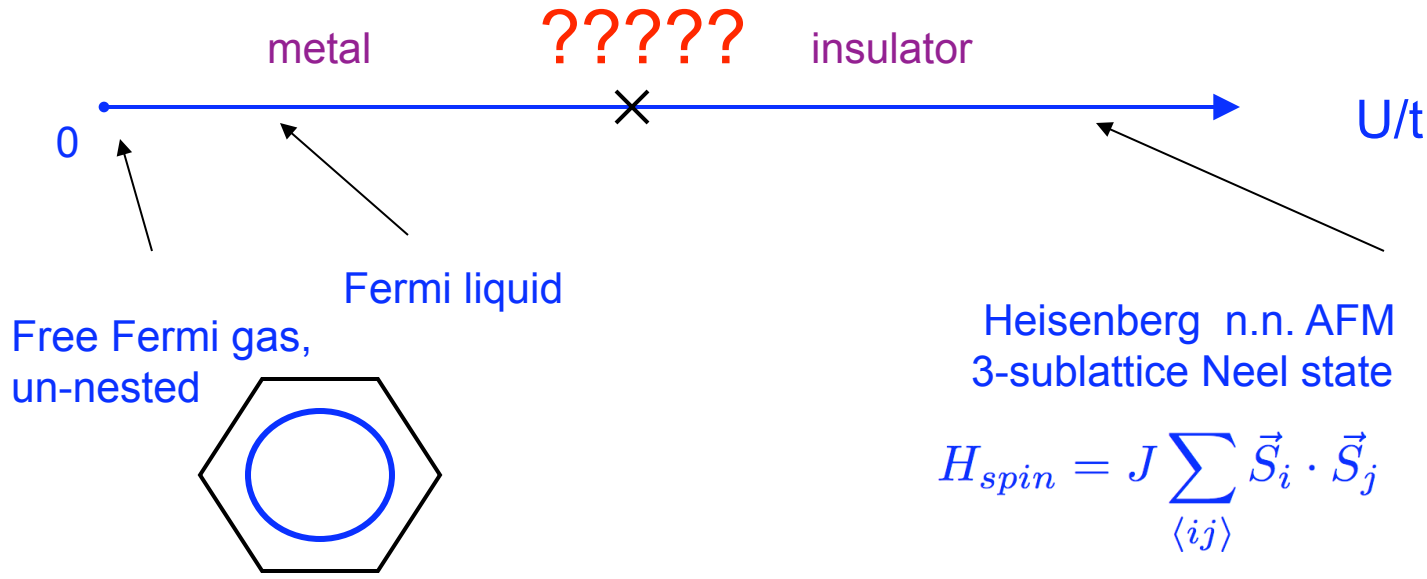
M. Yamashita et al (2010)

Hubbard model phase diagram on triangular lattice?

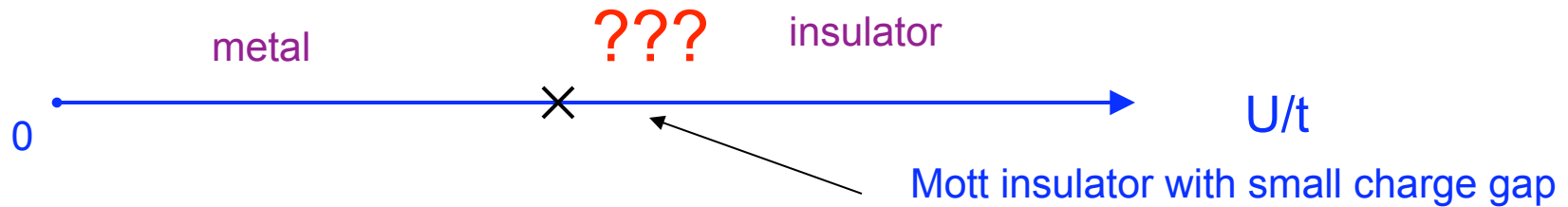
$$\mathcal{H} = -t \sum_{\langle ij \rangle} [c_{i\alpha}^\dagger c_{j\alpha} + h.c.] + U \sum_i n_{i\uparrow} n_{i\downarrow}$$



Phase diagram at Half filling?

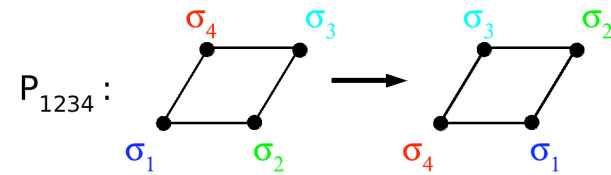


“Weak” Mott insulator on triangular lattice - Ring exchange



Expansion in t/U from Hubbard: spin model with ring exchange (mimics charge fluctuations)

$$\hat{H}_{\text{eff}} = \frac{2t^2}{U} \sum_{ij} \mathbf{S}_i \cdot \mathbf{S}_j + \frac{20t^4}{U^3} \sum_{\text{rhombi}} (P_{ijkl} + \text{h.c.}) + \dots$$



Triangular Lattice Heisenberg Model with ring exchange

$$\hat{H}_{\text{ring}} = J_2 \sum_{ij} \mathbf{S}_i \cdot \mathbf{S}_j + J_4 \sum_{\text{rhombi}} (P_{ijkl} + \text{h.c.})$$

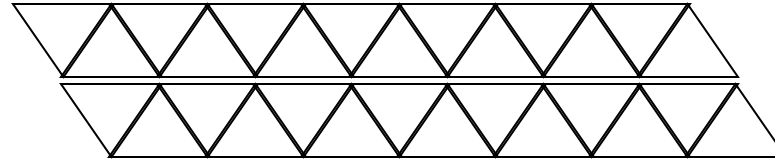
Is ground state a spin-liquid? A projected Fermi sea state?

VMC suggests yes for $J_4/J_2 > 0.3$, O. Motrunich - 2005

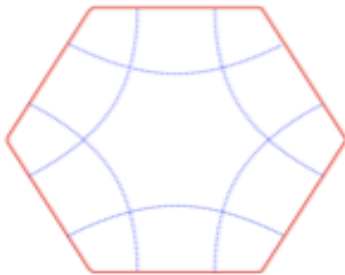
But, VMC is biased,
 QMC has a sign problem,
 DMRG is problematic in 2d so we have a quandry

Quasi-1d route to “Spin Bose-Metals”

Triangular strips:

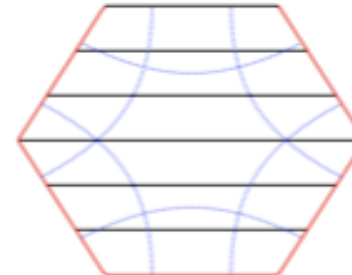


2d



Singular $2k_F$ surface
in spin correlators
("Bose surface")

Quasi-1d ladders

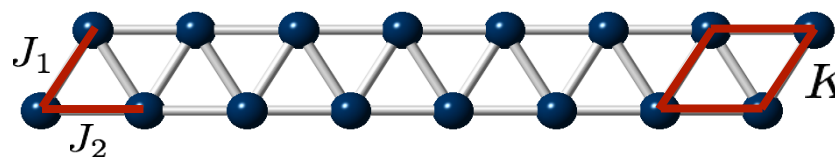
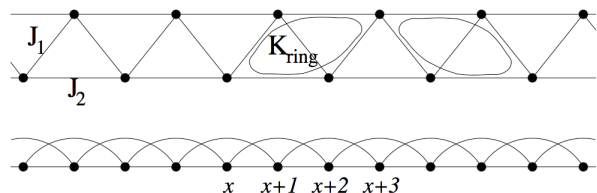


Fingerprint of 2d singular surface
on ladders - many gapless 1d modes,
of order N

2d Spin Bose-Metal should have quasi-1d descendent states

2-leg zigzag strip: SBM Descendent states?

View as 1d model



$$P_{1234}|\sigma_1, \sigma_2, \sigma_3, \sigma_4\rangle \rightarrow |\sigma_4, \sigma_1, \sigma_2, \sigma_3\rangle$$

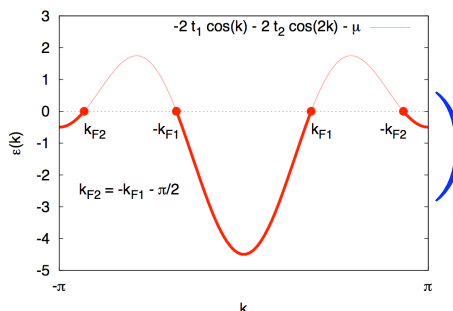
$$H_{\text{Heis}} + H_{\text{ring}} = \sum_x [2J_1 \mathbf{S}(x) \cdot \mathbf{S}(x+1) + 2J_2 \mathbf{S}(x) \cdot \mathbf{S}(x+2) + K(P_{x,x+2,x+3,x+1} + \text{h.c.})]$$

Analysis of J_1 - J_2 - K model on zigzag strip

D.N. Sheng, O.I. Motrunich, MPAF, PRB (2009)

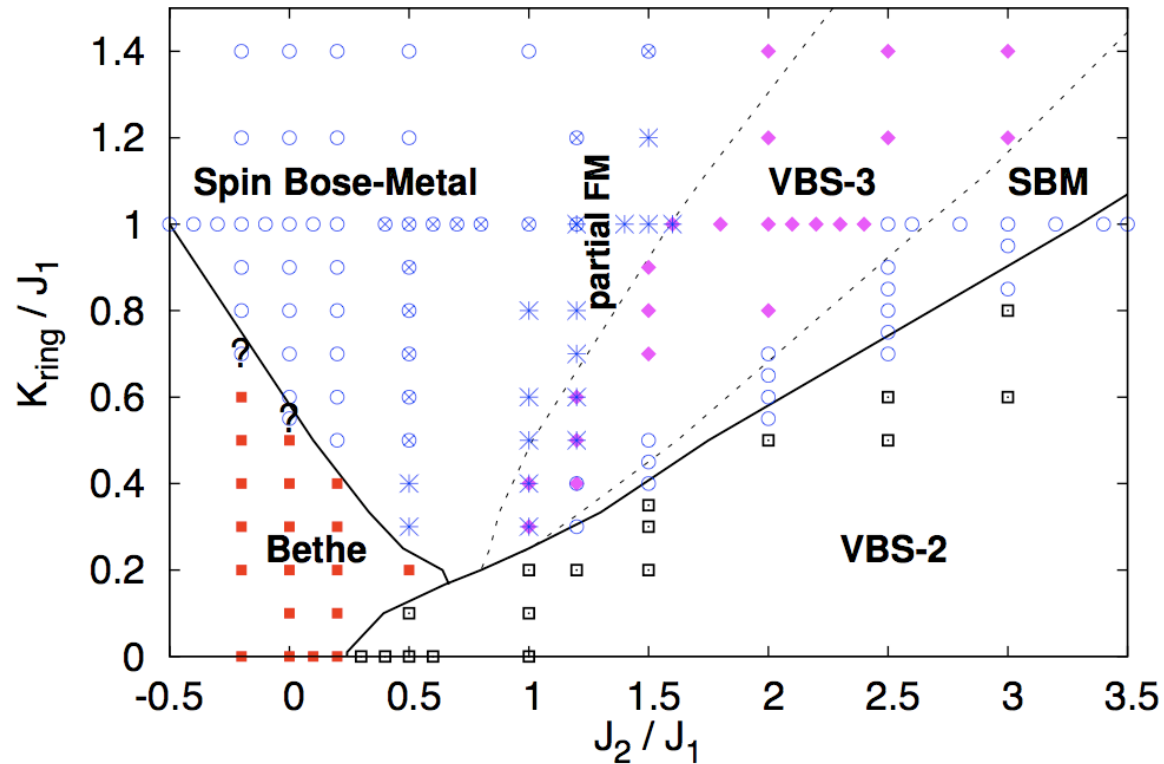
- DMRG
- Bosonization of gauge theory
- Gutzwiller wavefunctions: study with VMC

$$\Psi_{VMC} = \mathcal{P}_G(\dots)$$



Single Variational parameter,
size of 2nd Fermi sea

DMRG Phase Diagram of 2-leg zigzag ring model



Spin-Bose-Metal phase (Gutzwiller projected Fermi sea) over large region

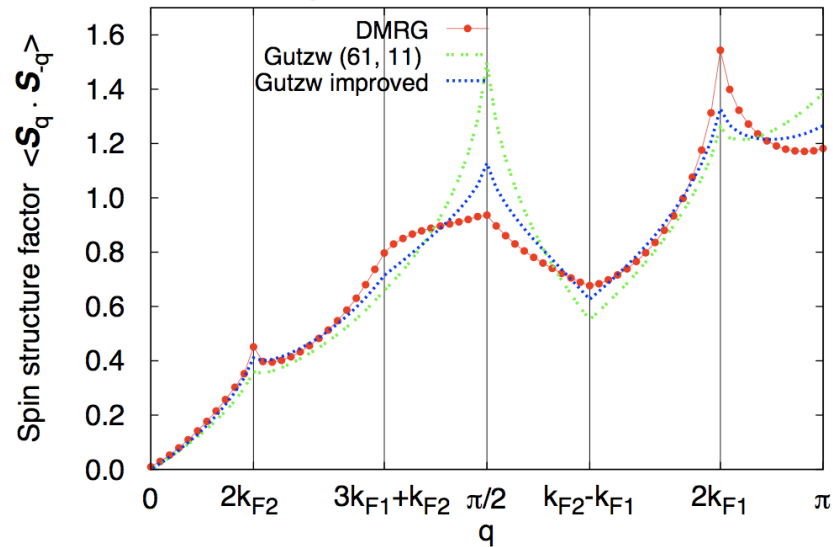
D.N. Sheng, O.I. Motrunich and MPAF, PRB (2009)

Spin Structure Factor in Spin Bose-Metal

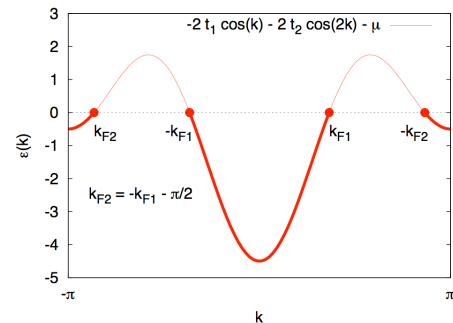
Singularities in momentum space locate the “Bose” surface (points in 1d)

$$\langle \vec{S}_k \cdot \vec{S}_{-k} \rangle$$

$K_{\text{ring}} = J_1 = 1, J_2 = 0, J_3 = 0; L=144$



Singular momenta can be identified with $2k_{F1}, 2k_{F2}$ which enter into Gutzwiller wavefunction

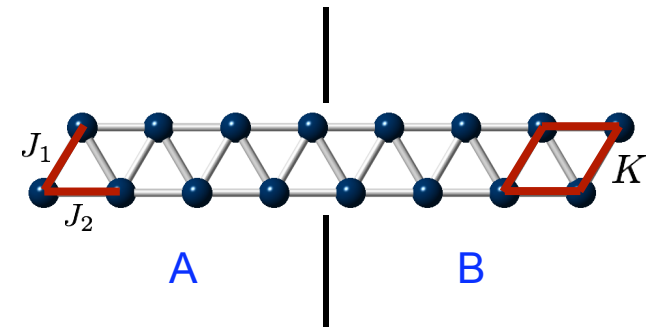


Entanglement in SBM

Reduced density matrix for sub-system A $\rho_A = \text{Tr}_B[\rho]$

Entanglement entropy $S_A = -\text{Tr}_A[\rho_A \ln \rho_A]$

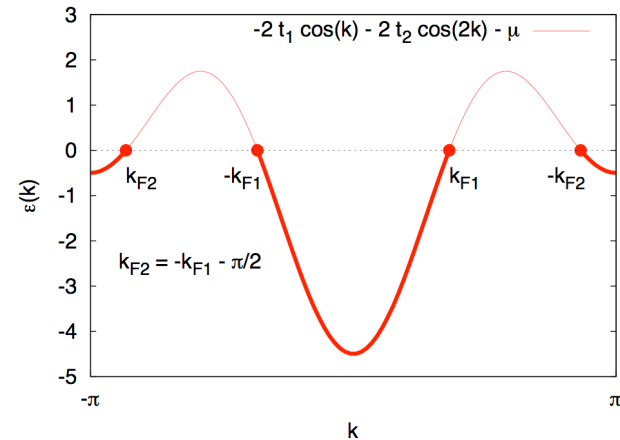
1d Gapless system (Conformal field theory, CFT) $S(X) = \frac{c}{3} \ln(X)$ $c = \text{central charge}$



Quasi-1d Gauge Theory for SBM

Two Fermi points - 2 charge and 2 spin modes

Gauge projection; Kills overall charge mode, 3 modes remain



Zigzag Spin Bose-Metal

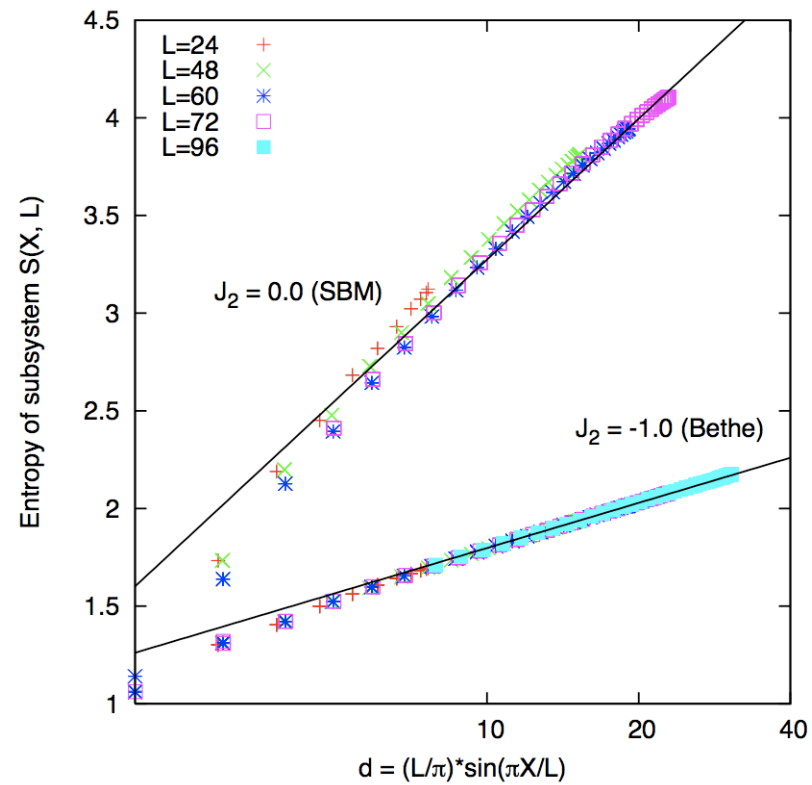
**Two gapless spin modes and one gapless spin-chirality mode:
Central charge $c=3$ predicted**

Entanglement Entropy from DMRG?

$$S(X, L) = \frac{c}{3} \log \left(\frac{L}{\pi} \sin \frac{\pi X}{L} \right) + \dots$$

nn 1d Heisenberg AFM
gives $c=1$, as expected

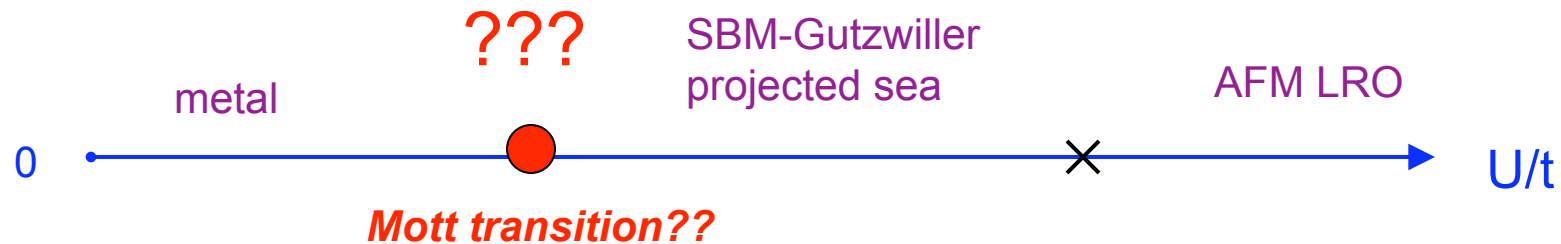
Spin Bose-Metal is
much more entangled



Entanglement entropy in SBM consistent with $c=3$

Metal to SBM Mott transition?

DMRG on 2-leg and 4-leg zigzag ladders - evidence for SBM state in 2d triangular lattice



Can 2d metal to spin-Bose-metal Mott transition be 2nd order?

2d Effective field theories say yes! (see next talk, Senthil)

Electron split into $s=1/2$ fermionic (spinon) and charge 1 bosonic ("chargon"), coupled to a $U(1)$ gauge field

$$c_{\sigma} = b f_{\sigma}$$

Mott transition driven by superfluid-insulator transition of b

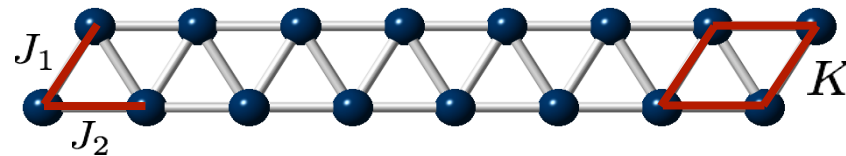
Quantum Criticality of Mott Transition?

2d: Challenging spinon-chargeon gauge theory

- Spinon fermi sea plus gauge field (SBM phase) Mross, McGreevy, Liu and Senthil (2010) and references therein
- Chargons/Spinons plus gauge field (Mott metal-SBM transition)

T. Senthil (2008); Mross and Senthil (2011),
Witczak-Kremo, Ghaemi, Senthil, Kim, cond-mat:1206.3309

Toy Problem: Metal to Spin-Bose-metal transition on 2-leg zigzag



(I) Bosonize Hubbard model on triangular strip:

(II) DMRG solution of Hubbard-type model on 2-leg zigzag

Bosonization of Metal-SBM transition on 2-leg zigzag

- Metal: 2 band Luttinger liquid (4 gapless modes)
- Eight fermion Umklapp drives transition into insulator (at strong coupling)

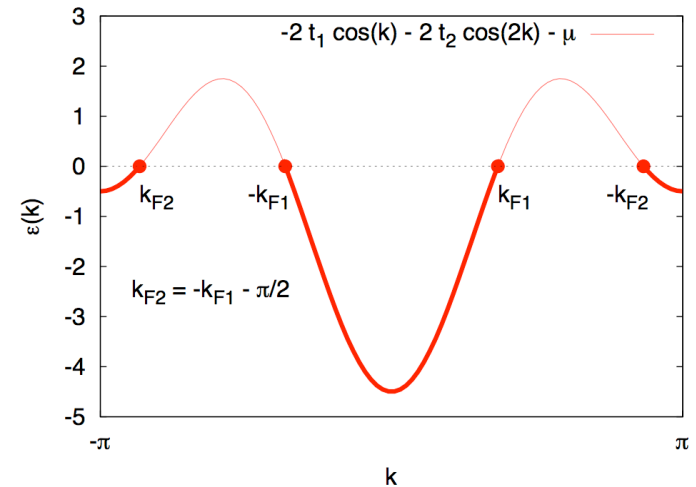
$$H_8 = v_8 (c_{R1\uparrow}^\dagger c_{R1\downarrow}^\dagger c_{R2\uparrow}^\dagger c_{R2\downarrow}^\dagger c_{L1\uparrow} c_{L1\downarrow} c_{L2\uparrow} c_{L2\downarrow} + \text{H.c.}),$$

- Kosterlitz-Thouless transition gaps overall charge mode, leaving 3 mode SBM
- Universal signature of KT Mott transition, Density correlator: $\langle \delta n_q \delta n_{-q} \rangle = g_\rho (2q/\pi)$

$$g_\rho \geq 1/2; \quad \text{metal} \quad (g_\rho = 1; \quad \text{free fermions})$$

$$g_\rho = 0; \quad \text{insulator}$$

$$g_\rho^* = 1/2; \quad \text{KT transition}$$

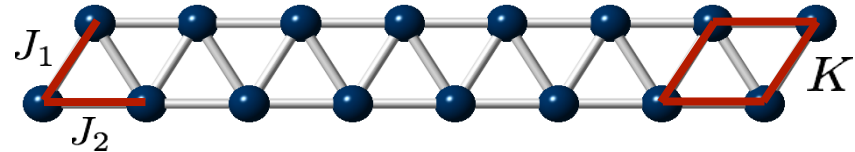


DMRG for Metal-SBM transition on 2-leg zigzag

Kinetic energy: Electrons hopping on 2-leg zigzag

R.V. Mishmash, I. Gonzalez, R. Melko, O. Motrunich, MPAF (iunpublished)

$$\mathcal{H}_0 = - \sum_i [t c_{i\sigma}^\dagger c_{i+1\sigma} + t' c_{i\sigma}^\dagger c_{i+2\sigma} + h.c.]$$

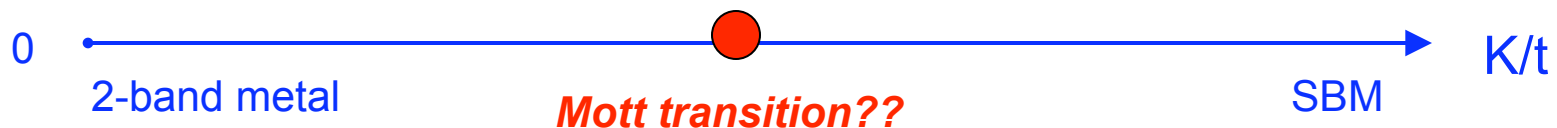


Interactions: 2 models

- Hubbard $\mathcal{H}_U = U \sum_i n_{i\uparrow} n_{i\downarrow}$
- 2-Spin and 4-spin ring exchanges

$$\mathcal{H}_{JK} = \sum_i [J_1 \vec{S}_i \cdot \vec{S}_{i+1} + J_2 \vec{S}_i \cdot \vec{S}_{i+2}] + K (P_{i,i+2,i+3,i+1} + h.c.)$$

Here, we will fix $K/J_1 = 0$, $J_2/J_1 = 0$, $t'/t = 0.75$ and vary the single dimensionless parameter K/t :



DMRG for electron zigzag t-JK model

R.V. Mishmash, I. Gonzalez,
R. Melko, O. Motrunich,
MPAF (iunpublished)

Spin-spin and
Density-density
correlators

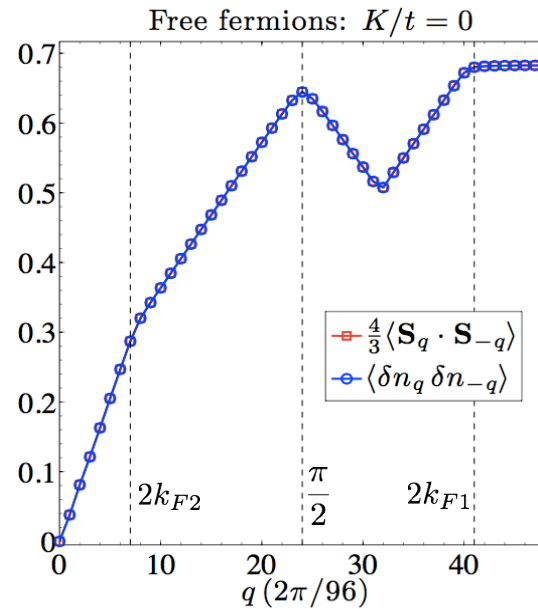
$$\langle \delta n_q \delta n_{-q} \rangle$$

$$\langle \vec{S}_q \cdot \vec{S}_{-q} \rangle$$

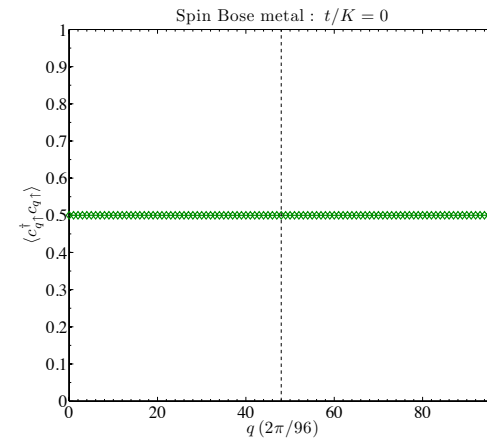
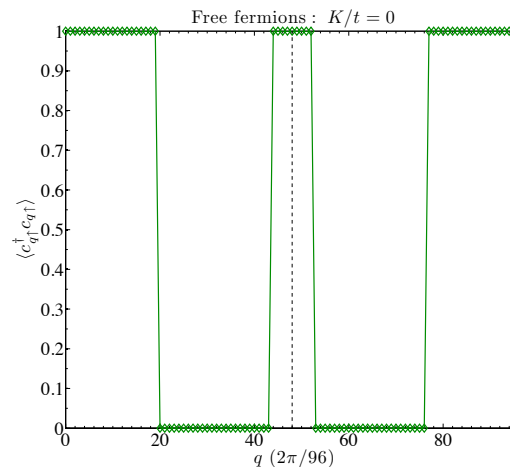
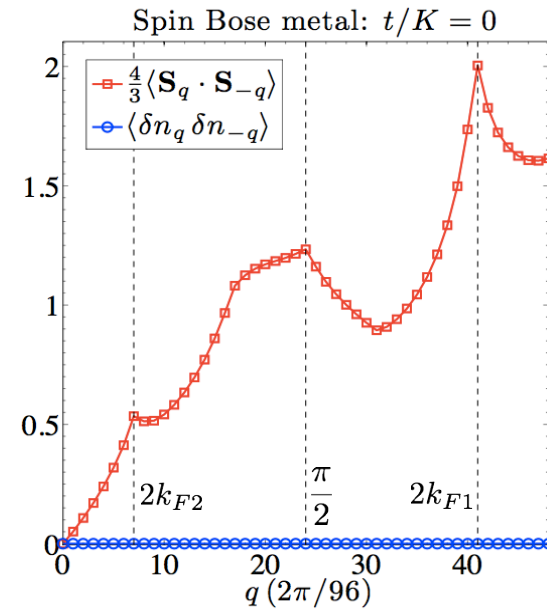
Electron momentum
distribution functions

$$\langle c_{q\uparrow}^\dagger c_{q\uparrow} \rangle$$

Free electron Metal: $K/t=0$



SBM Insulator: $K/t = \text{infinity}$



Quantum Criticality at Mott transition?

Evidence for Kosterlitz-Thouless
from density-density structure factor

$$\langle \delta n_q \delta n_{-q} \rangle = g_\rho (2q/\pi)$$

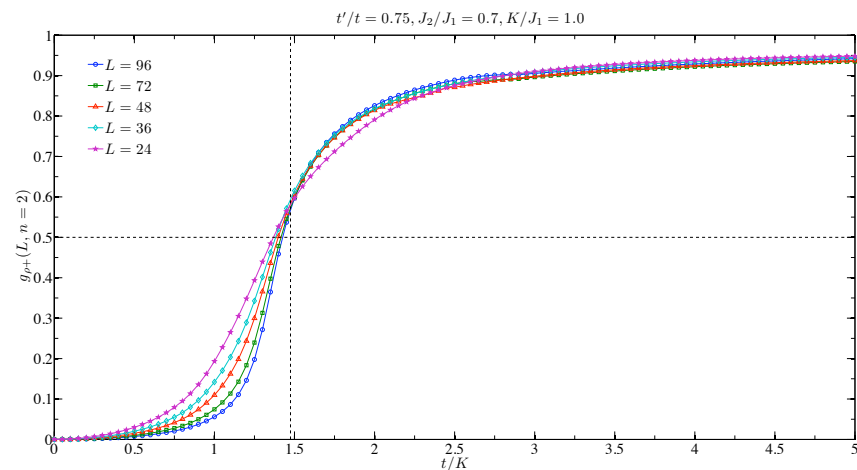
KT Theory prediction:

$g_\rho \geq 1/2$; *metal* ($g_\rho = 1$; *free fermions*)

$g_\rho = 0$; *insulator*

$g_\rho^* = 1/2$; *KT transition*

DMRG data for $L \rightarrow \infty$
consistent with $g_\rho = 0$
in insulator, and $g_\rho^* = 1/2$
at KT transition

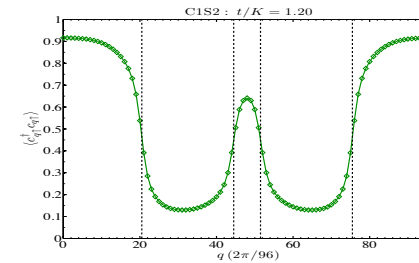


Singular *charge-density* fluctuations in SBM Insulator!

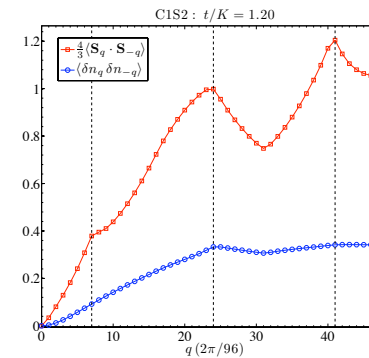
(Charge Friedel Oscillations in a Mott insulator:
D.F. Mross and T. Senthil PRB (2011))

Electron gapped in insulator,
smooth electron momentum
distribution function

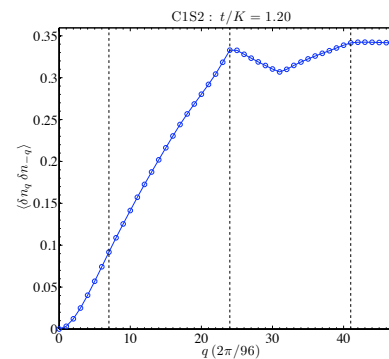
$$\langle c_{q\uparrow}^\dagger c_{q\uparrow} \rangle$$



Spin structure factor singular at $2k_F$
(signature of “spinon Fermi-surface”)



Density structure factor also
singular at $2k_F$, despite being
insulating $\langle \delta n_q \delta n_{-q} \rangle \sim q^2$



DMRG for 2-leg zigzag Hubbard

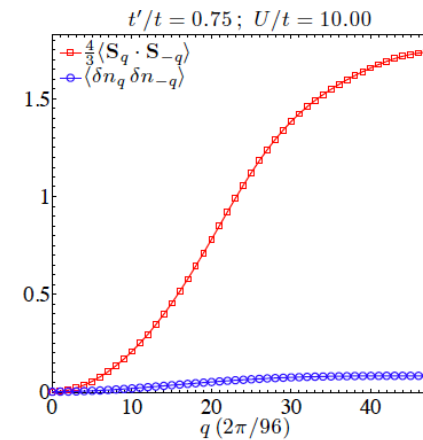
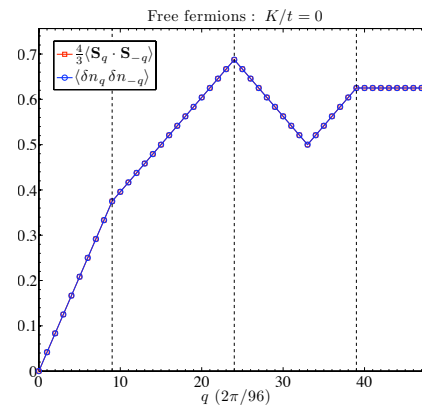
Free fermions $U/t=0$

Insulator $U/t=10$

Spin-spin and
Density-density
correlators

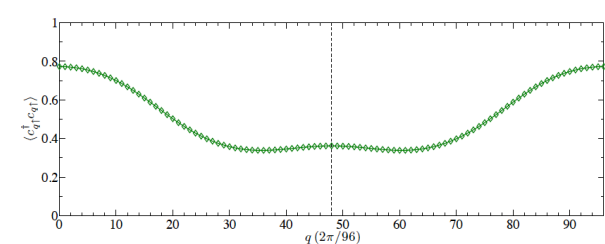
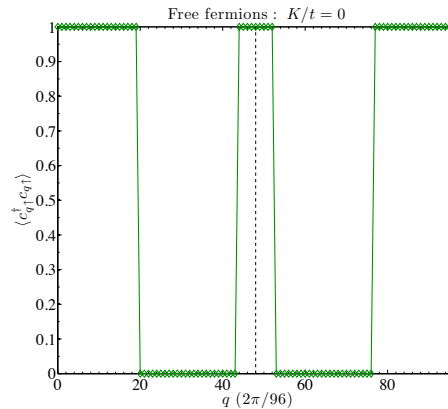
$$\langle \delta n_q \delta n_{-q} \rangle$$

$$\langle \vec{S}_q \cdot \vec{S}_{-q} \rangle$$



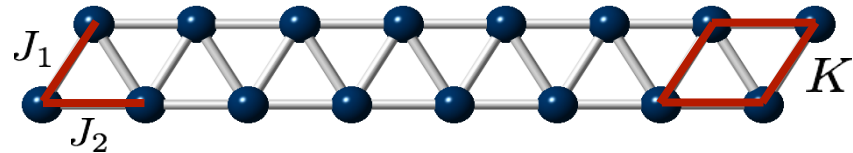
Electron momentum
distribution functions

$$\langle c_{q\uparrow}^\dagger c_{q\uparrow} \rangle$$



Summary and Outlook

- Continuous 2d Mott transition from metal to Spin-Bose-metal insulator (Gutzwiller projected Fermi sea state) is possible
- DMRG on 2-leg zigzag finds quasi-1d analog: Kosterlitz-Thouless transition from 2-band metal to quasi-1d SBM (Gutzwiller projected sea)



Numerics (DMRG)

- Continuous Mott transition on 2-leg zigzag Hubbard model?
- Mott transition accessible on multi-leg triangular lattice Hubbard model?
- Access 2d criticality?

Analytics: Field theory of 2nd order 2d Mott metal-insulator transition (Senthil et. al.)

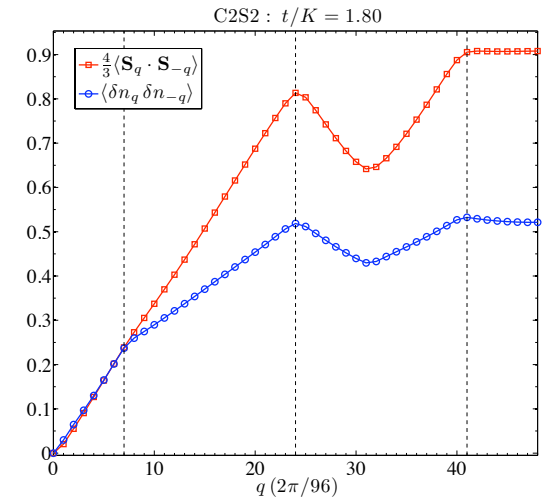
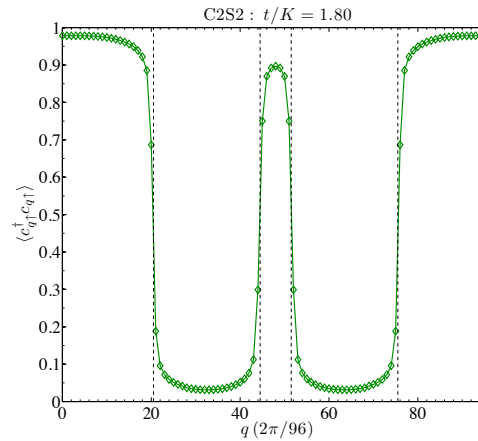
Experimental manifestation of continuous Mott transition?
(Organics, perhaps weak first order metal-insulator transition?)

DMRG features

2-Band Luttinger liquid

Electron momentum distribution quite sharp despite O(1) interactions

Expected singular features in spin and density momentum structure factors



Evidence for Kosterlitz-Thouless Criticality at Mott transition

$$\langle \delta n_q \delta n_{-q} \rangle = g_\rho(2q/\pi)$$

$$L \rightarrow \infty, \quad \Delta g_\rho \approx 1/2$$

