

Experiments on a Quantum Hall Exciton Condensate

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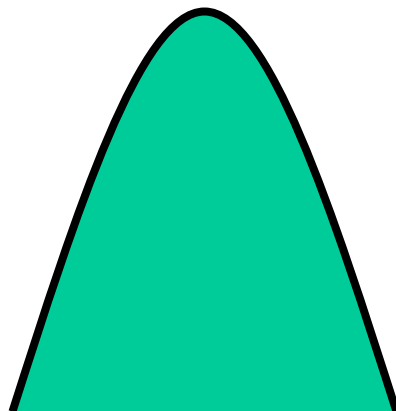
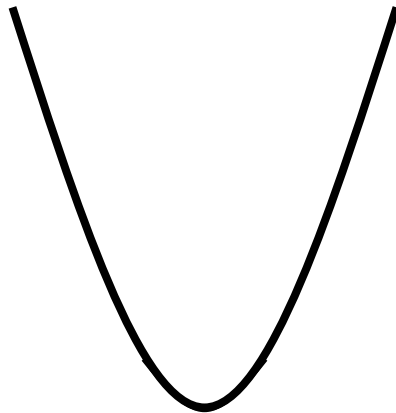


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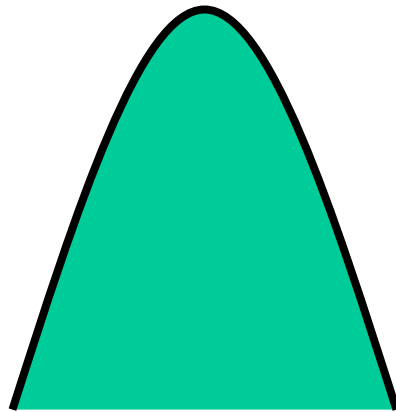
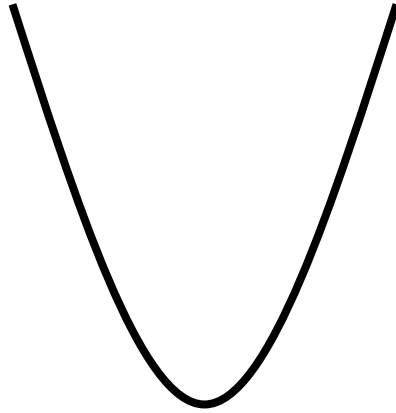


Division of Materials Sciences

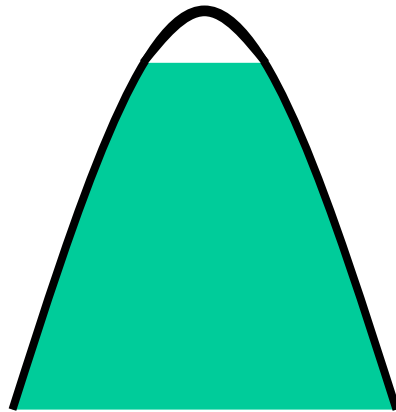
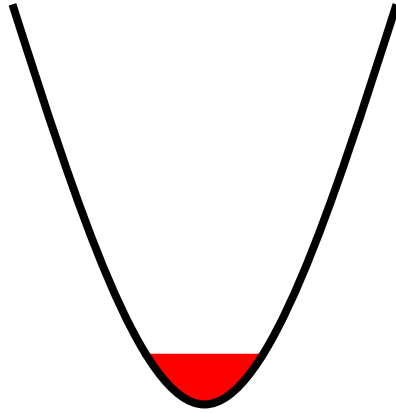
Intrinsic semiconductor at $T = 0$



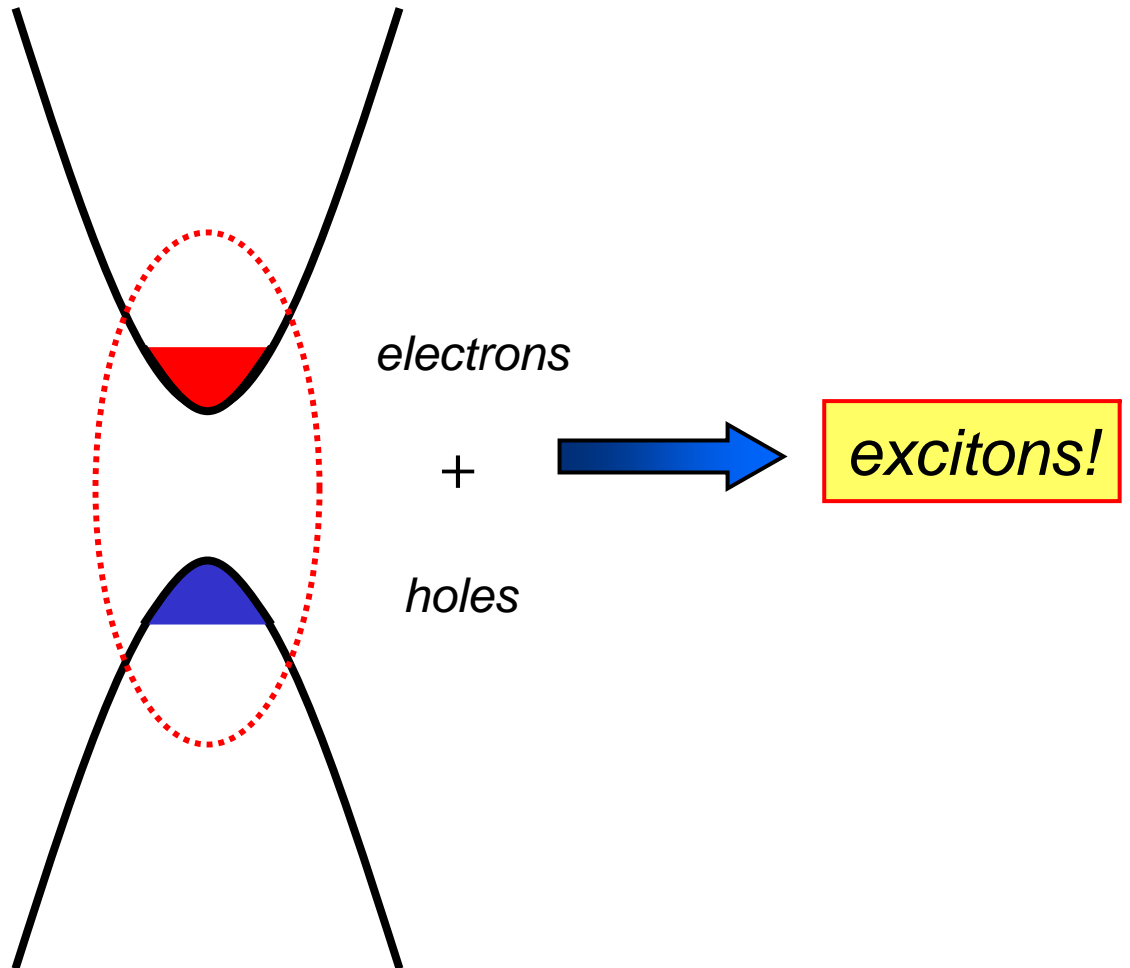
Add light



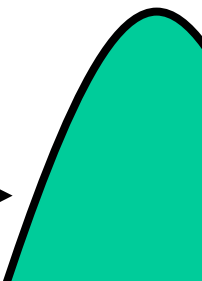
Only electrons



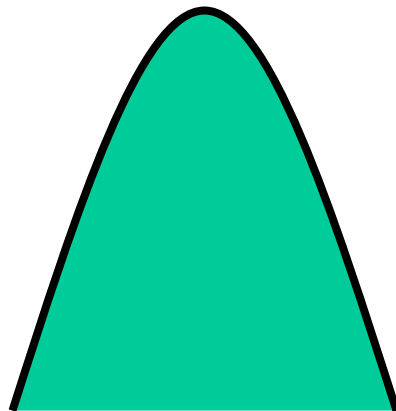
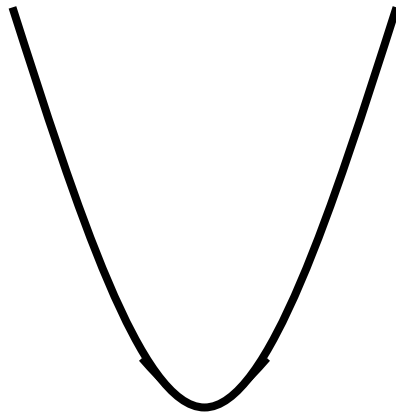
Particle-hole transformation on valence band



Embarrassing relative →



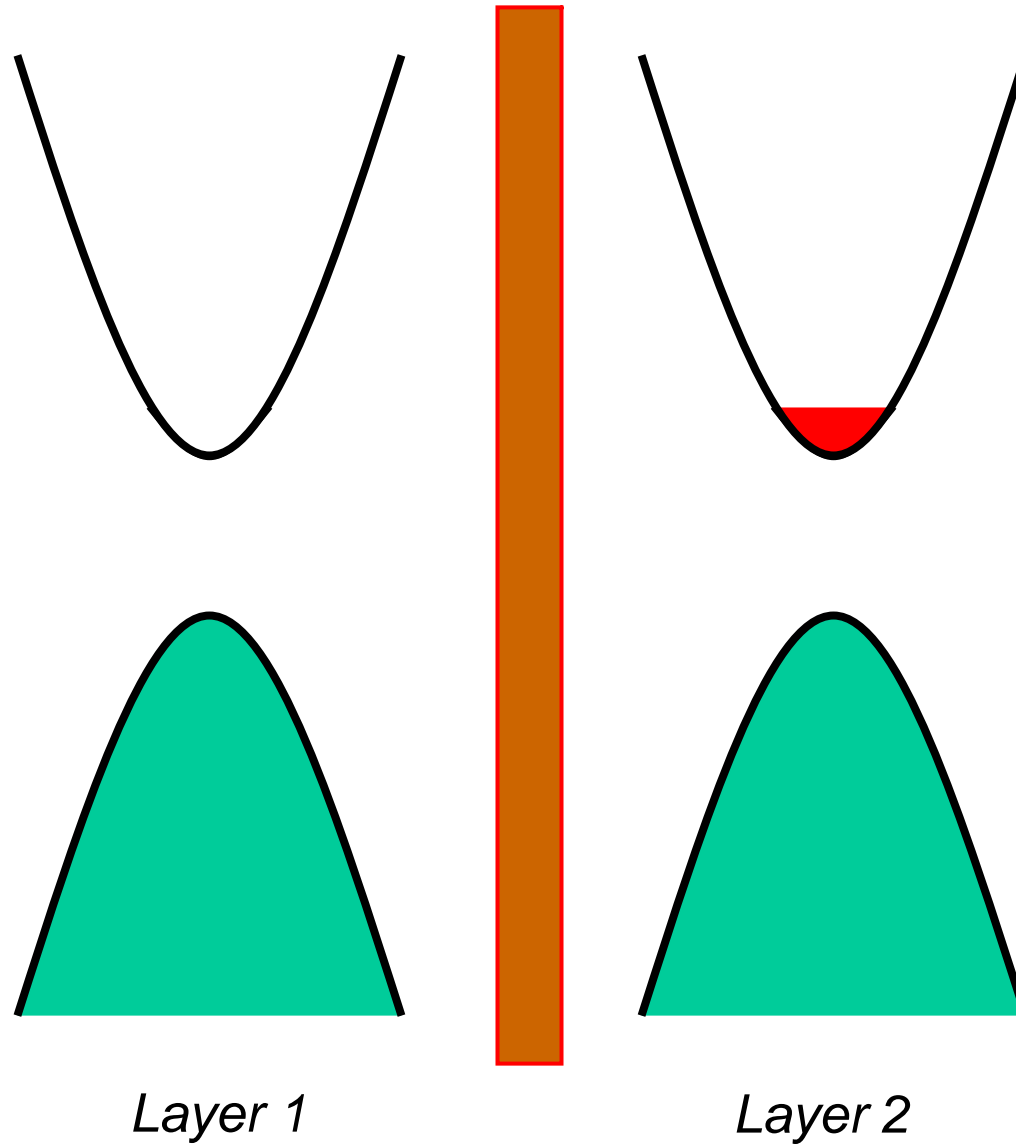
Recombination



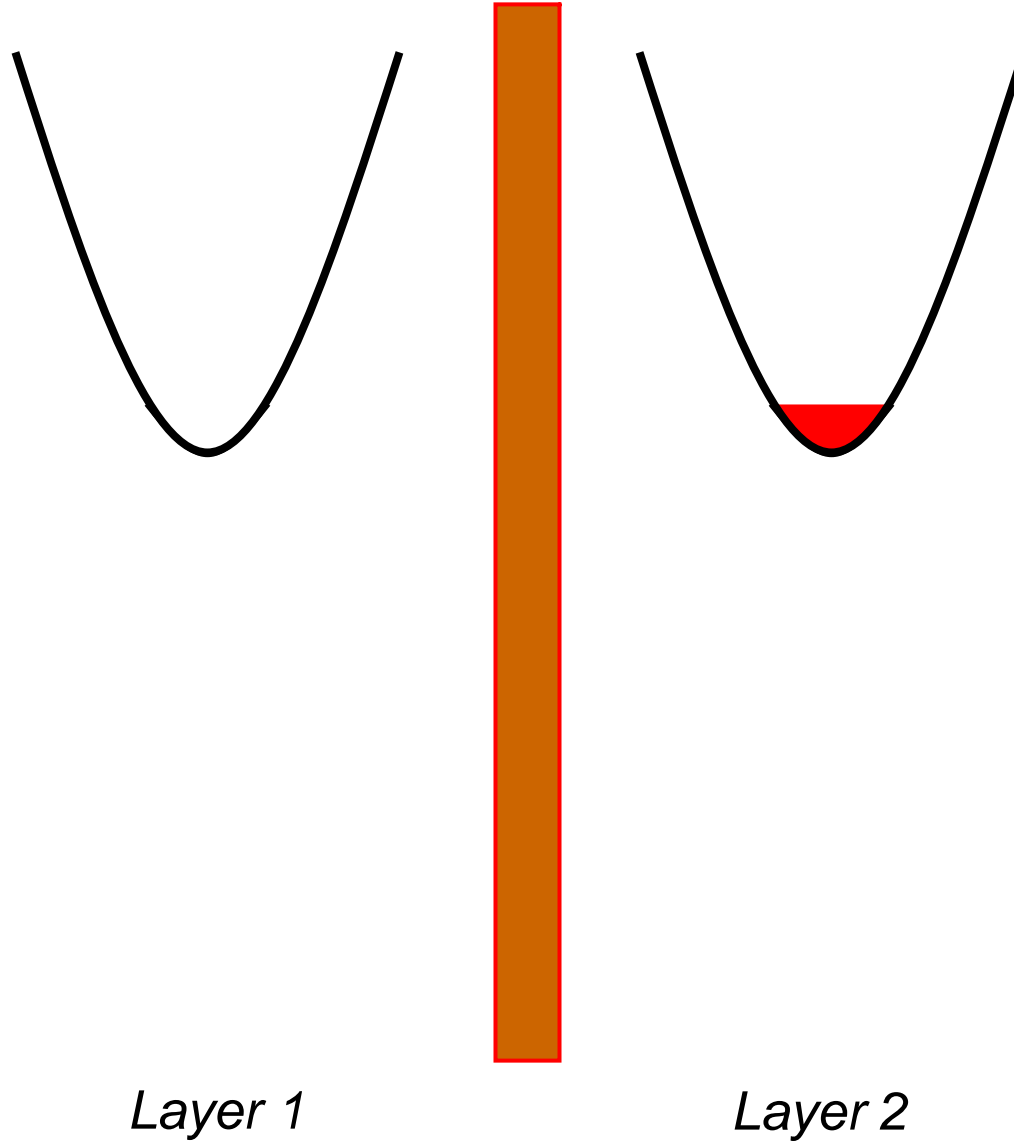


PHOTONS

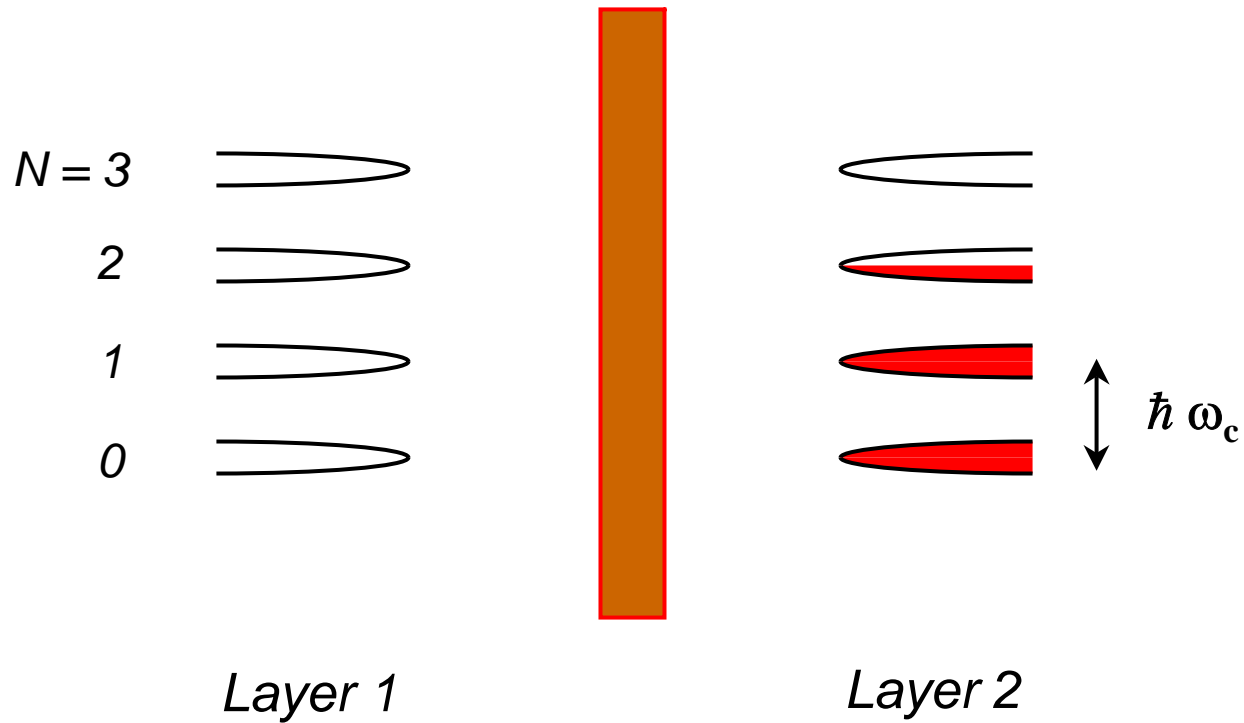
A different way: Doped bilayer semiconductor



Discard the valence band

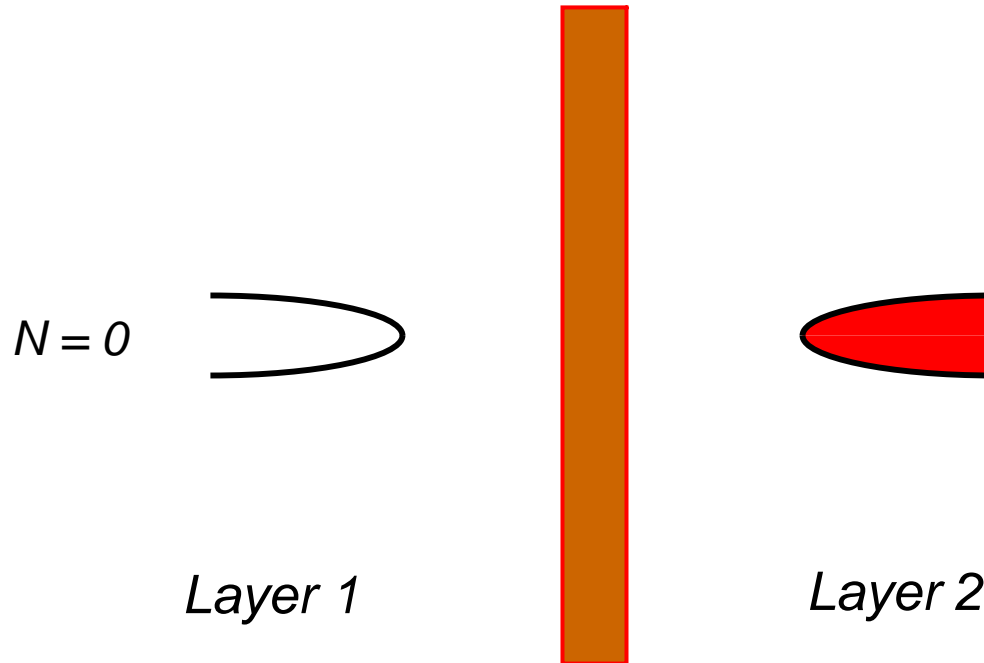


Add a magnetic field



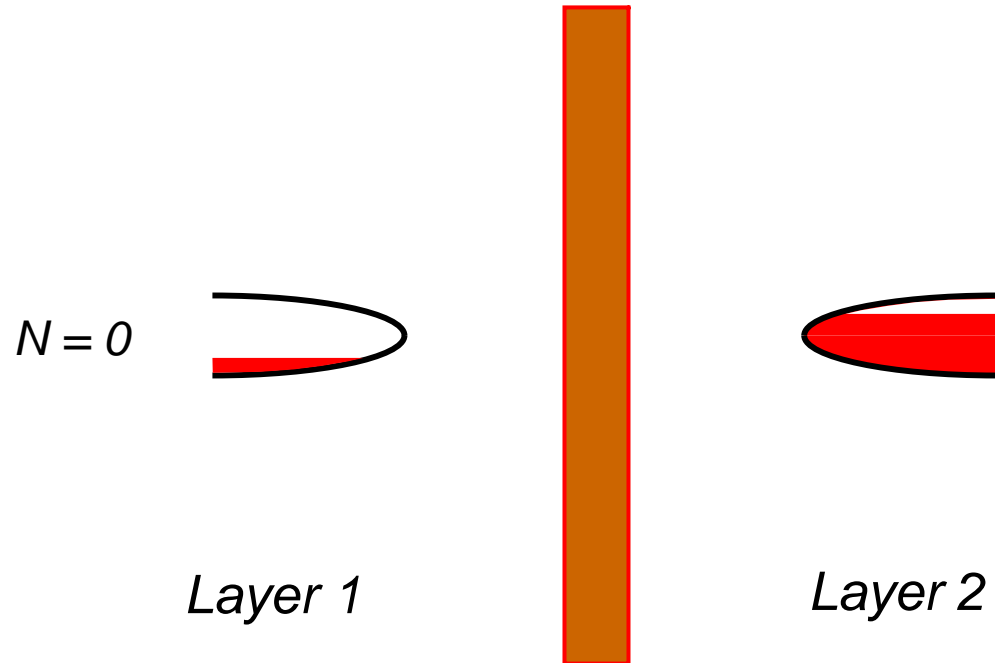
conduction band Landau levels

A bigger magnetic field: One filled Landau level



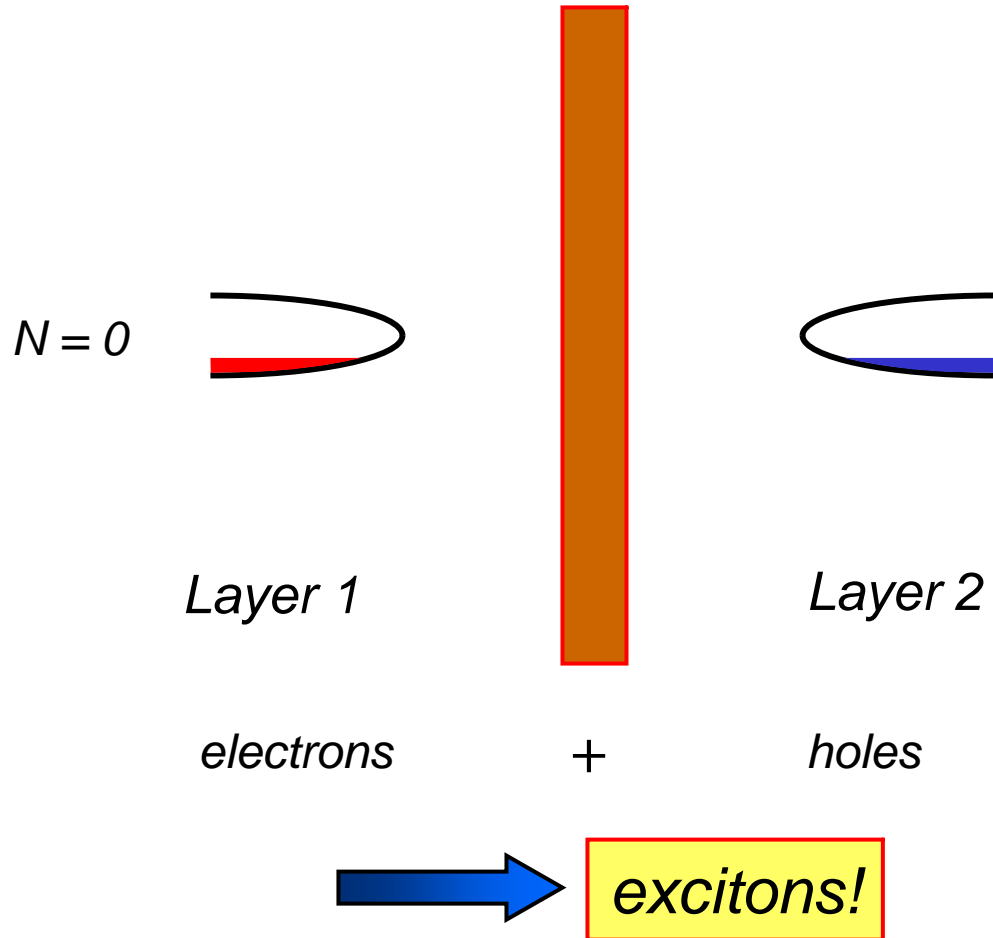
Analogous to intrinsic semiconductor at $T = 0$, but no gap!

Transfer charge by doping or gating



Total number of electrons unchanged

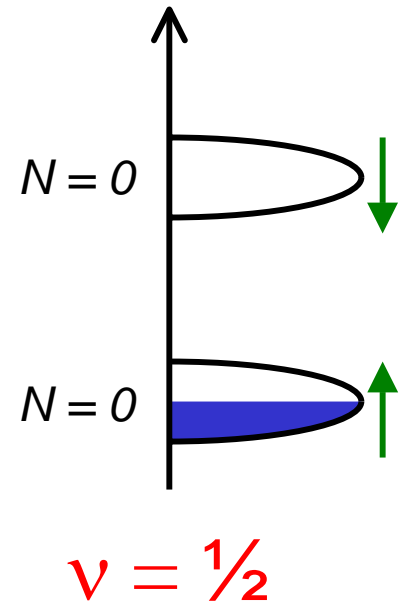
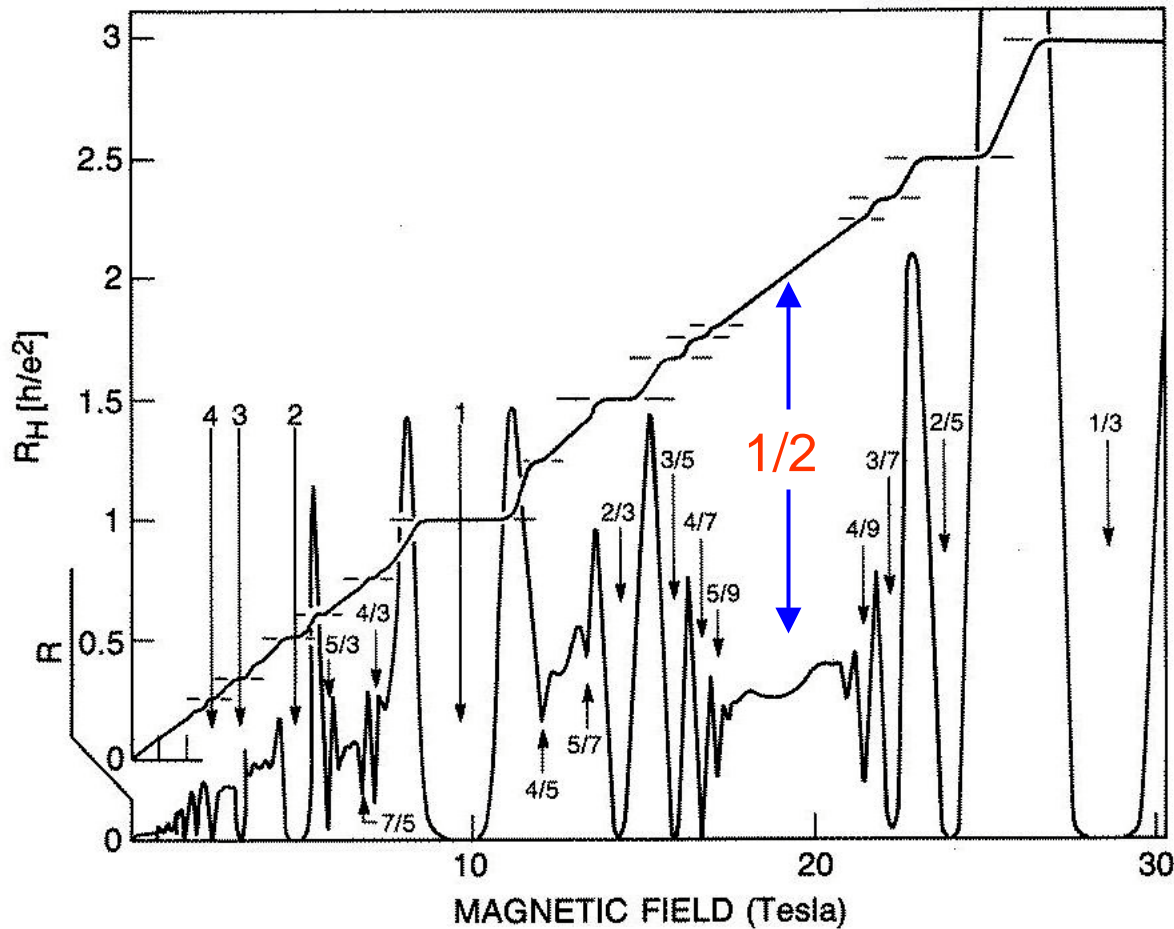
Particle-hole transformation on lowest Landau level in layer 2



Embarrassing relative →



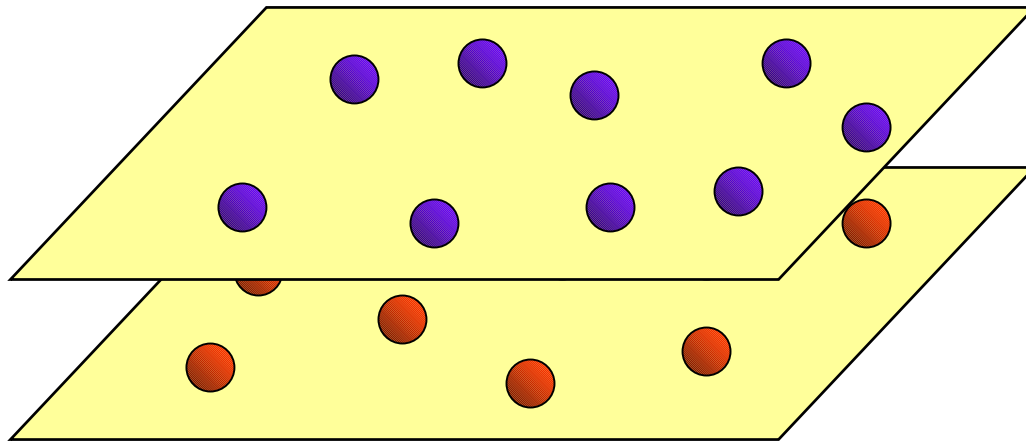
QHE Single Layer 2D System



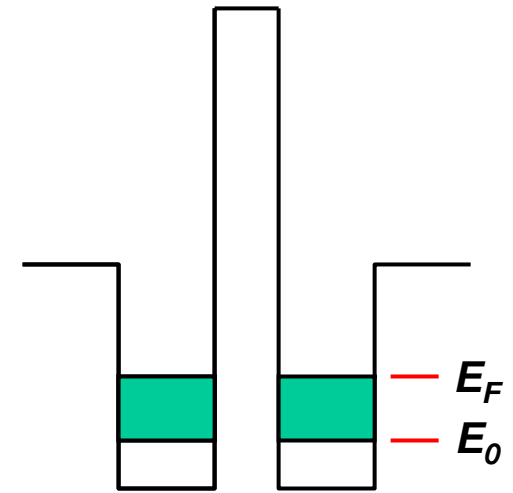
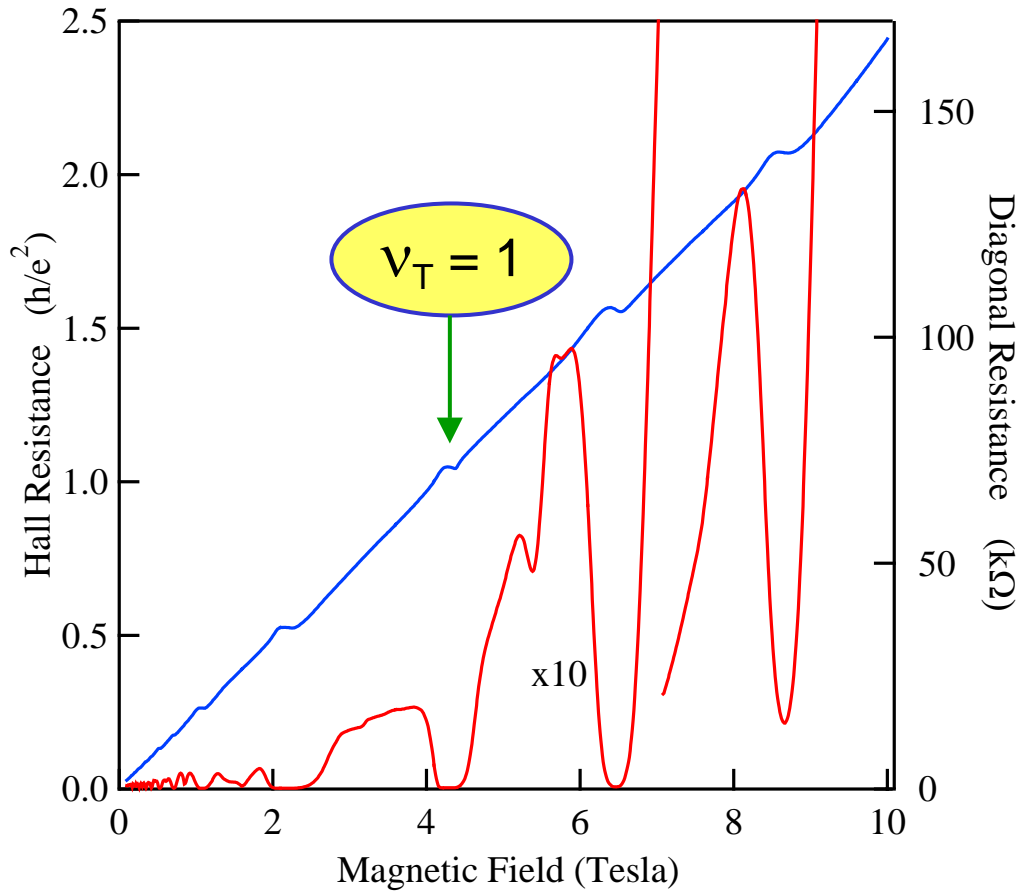
No QHE at half-filling of the lowest Landau level



Double Layer Two-Dimensional Electron Gas



QHE in Double Layer 2D System

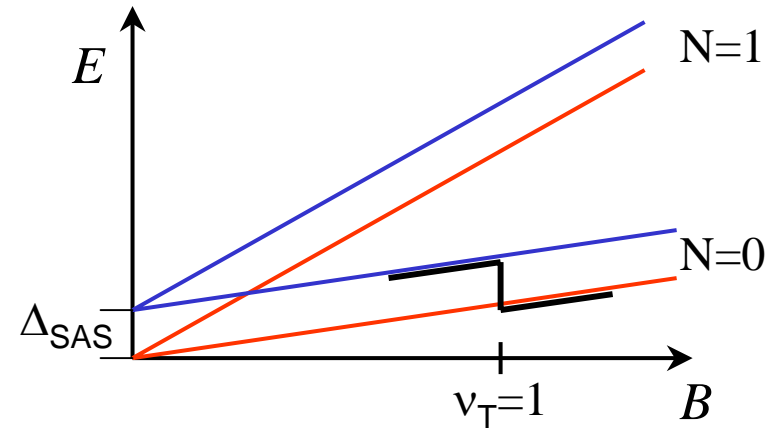
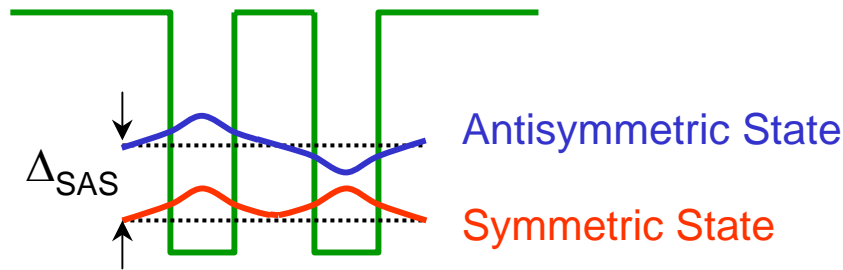


$$\nu_T = 1 = \frac{1}{2} + \frac{1}{2}$$

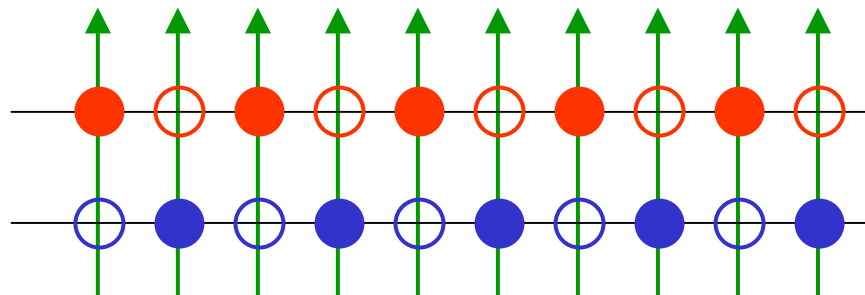


Single Particle and Many-body Origins for $\nu_T = 1$ QHE

1. Single Particle Tunneling



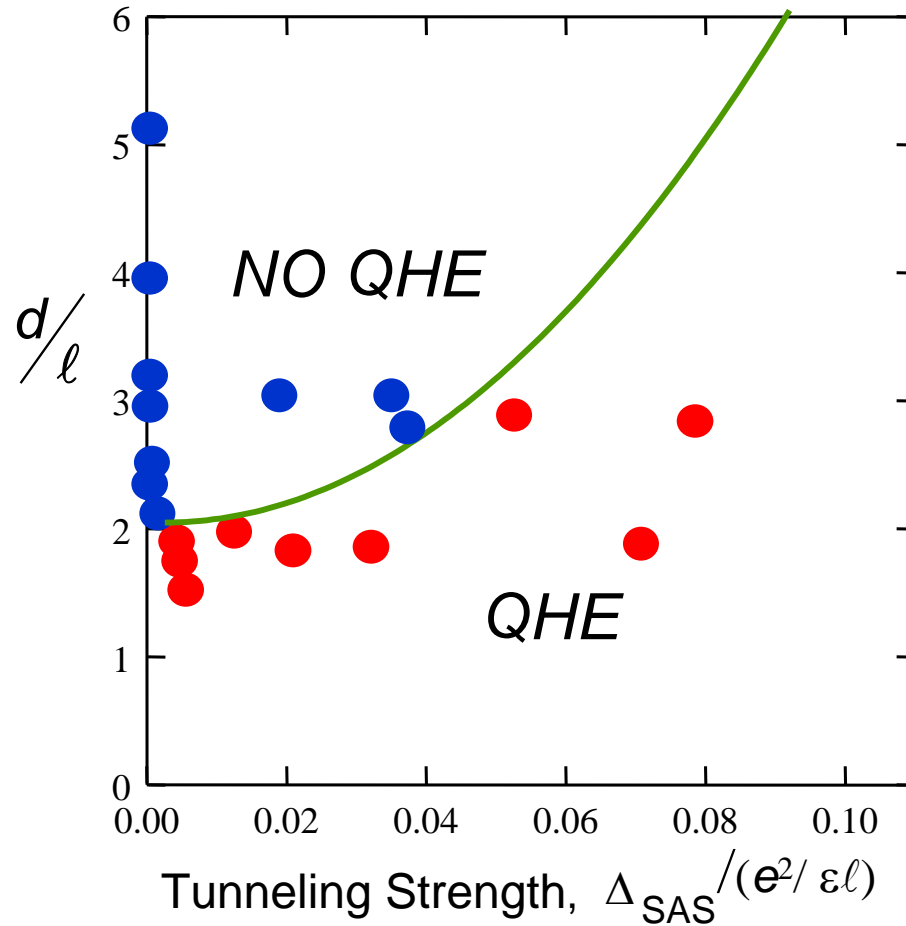
2. Pure Many-body Effect



$$\Psi \sim \prod_{i, \dots, n} (z_i - z_j) (w_k - w_l) (z_m - w_n)$$



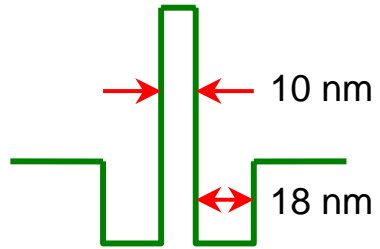
Phase Diagram



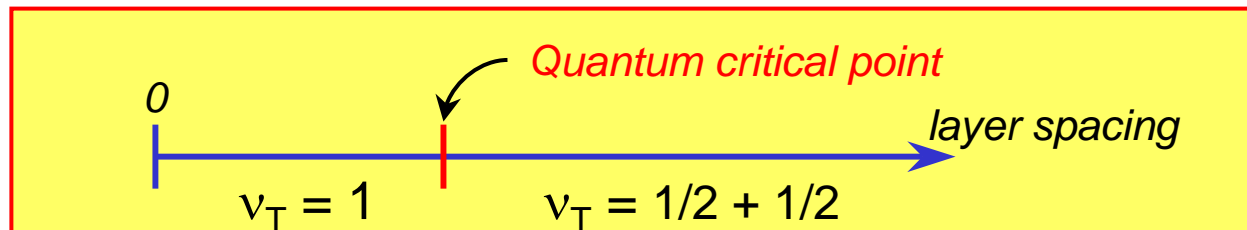
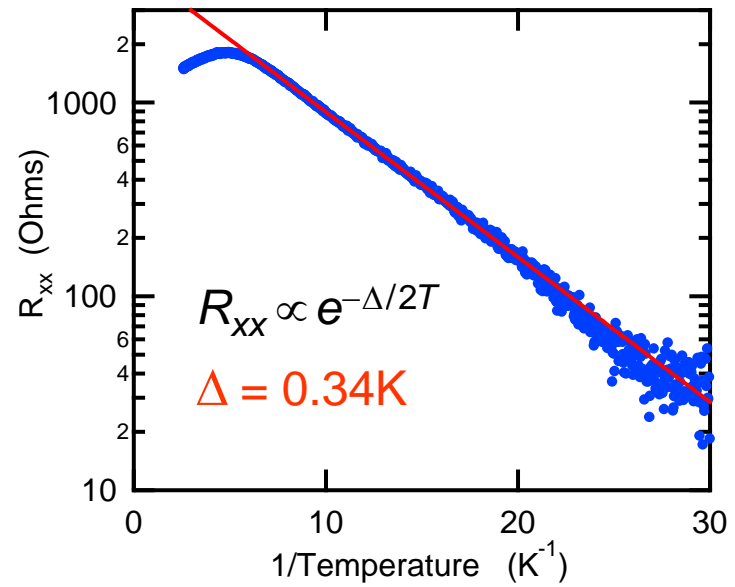
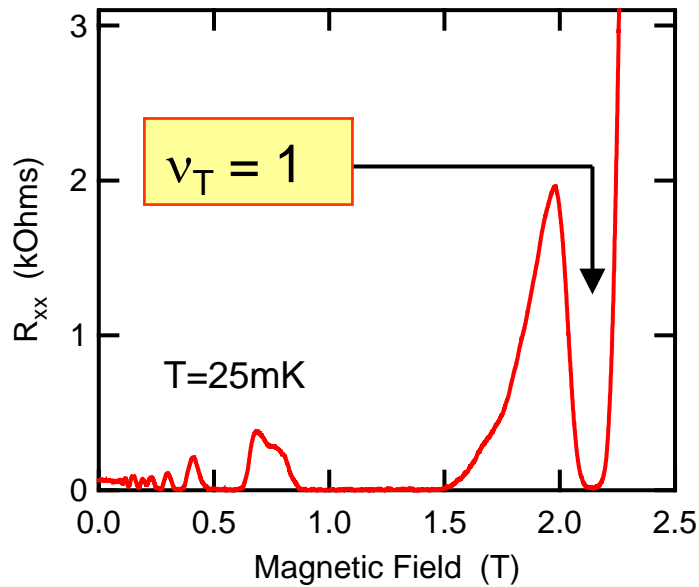
QHE persists in zero tunneling limit



Bilayer $\nu_T = 1$ QHE with Nearly Zero Tunneling

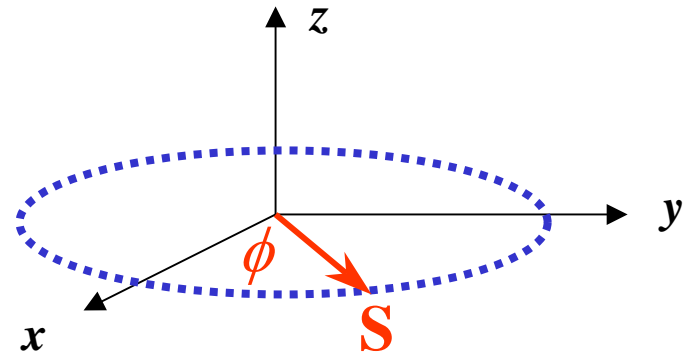
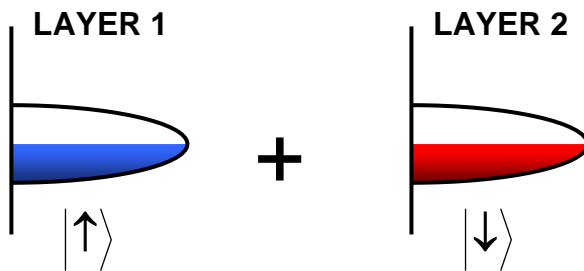


$$\Delta_{\text{SAS}} \approx 90 \mu\text{K} \approx 1.2 \times 10^{-6} \left(\frac{e^2}{\epsilon l} \right)$$



Easy-Plane Ferromagnet

layer index \rightarrow *pseudospin*



$$|\Psi\rangle = \prod_k |k\rangle \otimes \frac{1}{\sqrt{2}} \left[|\uparrow\rangle + e^{i\phi} |\downarrow\rangle \right]$$

Spontaneous interlayer phase coherence and “*which layer*” uncertainty.

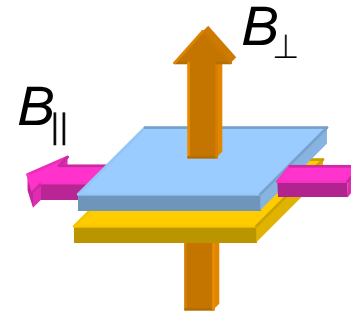
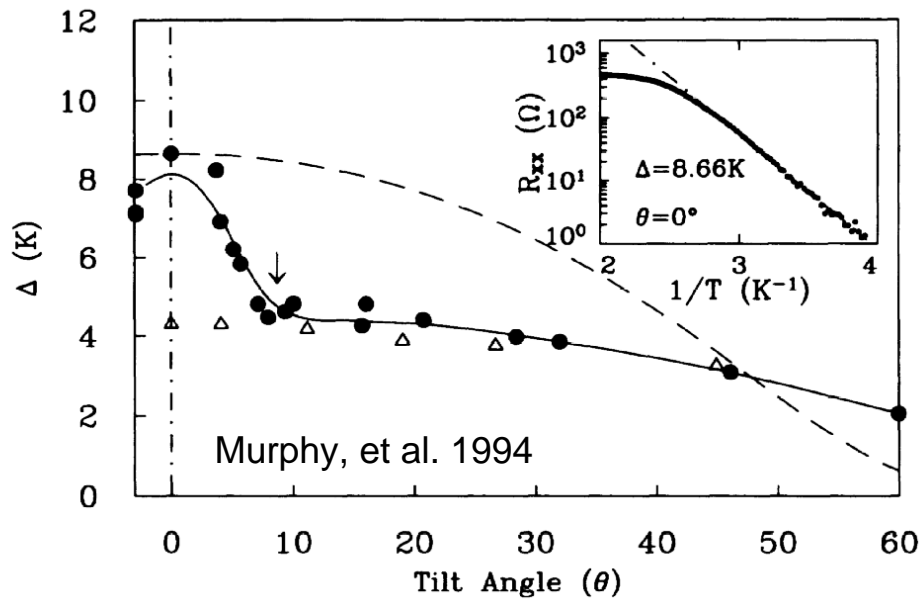
pseudo-spin waves, charged vortices, etc.



dissipationless pseudo-spin currents

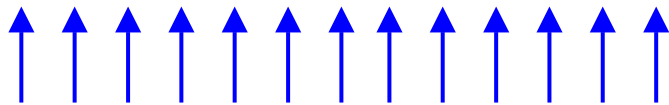


Pseudospin Textural Phase Transition



$$H = \frac{\rho_s}{2} |\nabla \varphi|^2 - \frac{\Delta_{SAS}}{4\pi \ell^2} \cos(\varphi - Qx)$$

$$Q = eB_{\parallel} d / \hbar$$



$$\lambda = h/eB_{\parallel}d$$



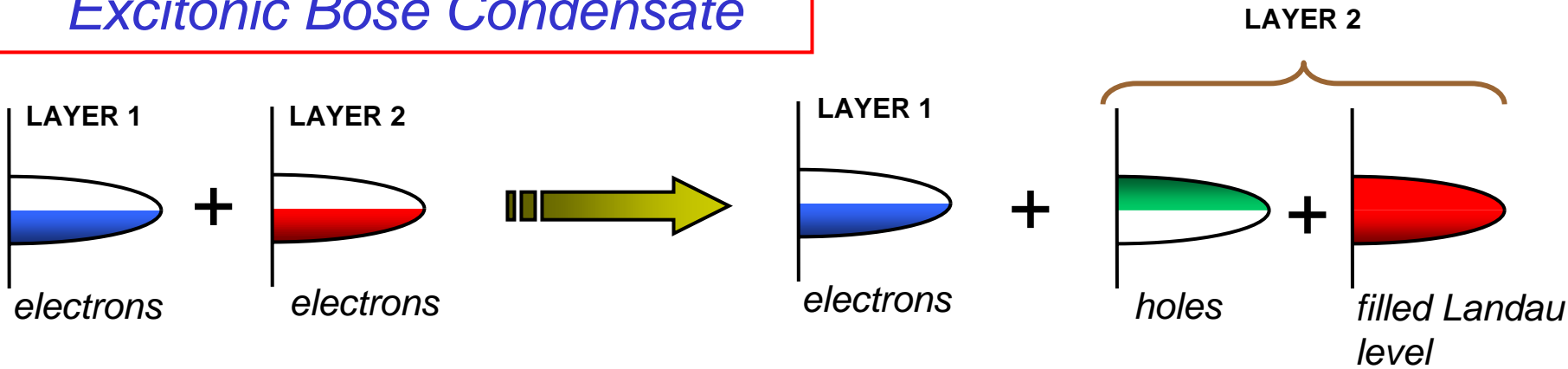
increasing B_{\parallel}



Yang, et al. 1994

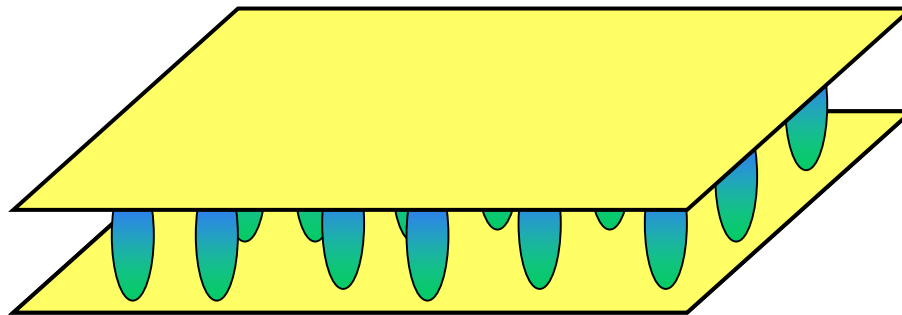


Excitonic Bose Condensate



$$|\Psi\rangle = \prod_k \frac{1}{\sqrt{2}} \left[1 + e^{i\phi} \underbrace{c_{k,1}^\dagger c_{k,2}}_{\text{exciton creation operator}} \right] |vac'\rangle$$

exciton creation operator



$\nabla\phi$

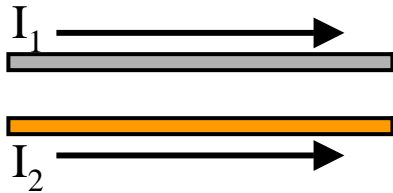


excitonic supercurrents

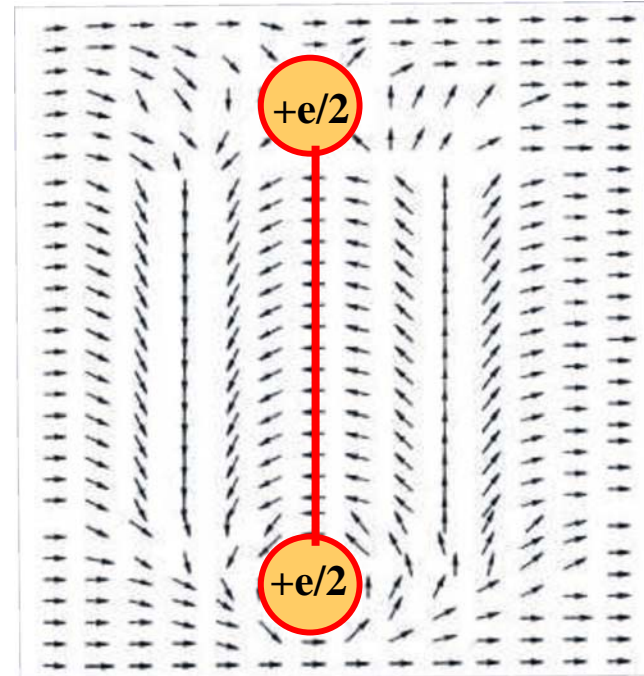


Two Transport Channels

1. Parallel Transport



meron / anti-meron pair

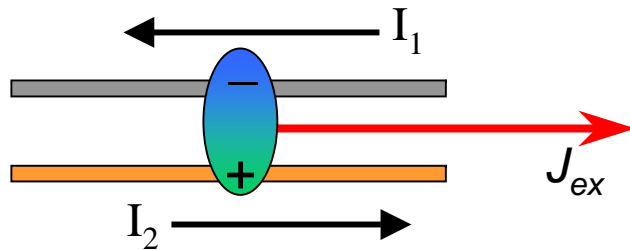


- vortex charged *excitations above gap*
- edge states
- quantized Hall effect
- dissipation $\sim \exp(-\Delta/T)$

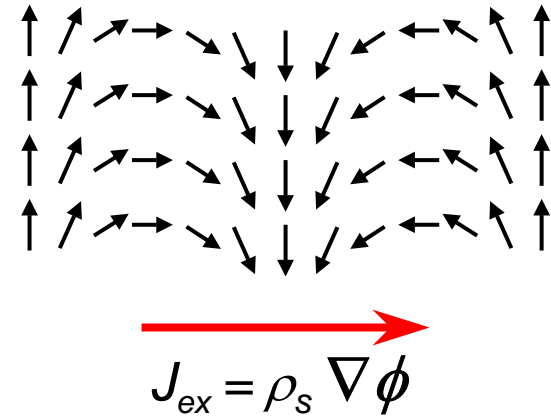


Two Transport Channels

2. Counterflow Transport



$$\nabla \phi = \text{constant}$$

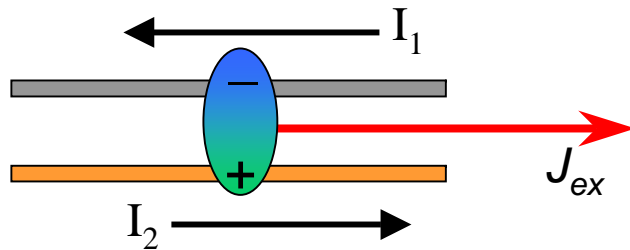


- *collective exciton transport in condensate*
- *bulk flow*
- *zero dissipation for $T < T_{KT}$*

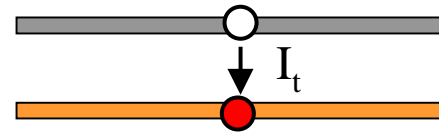


The really cool physics is in the antisymmetric channel!

counter-flow transport



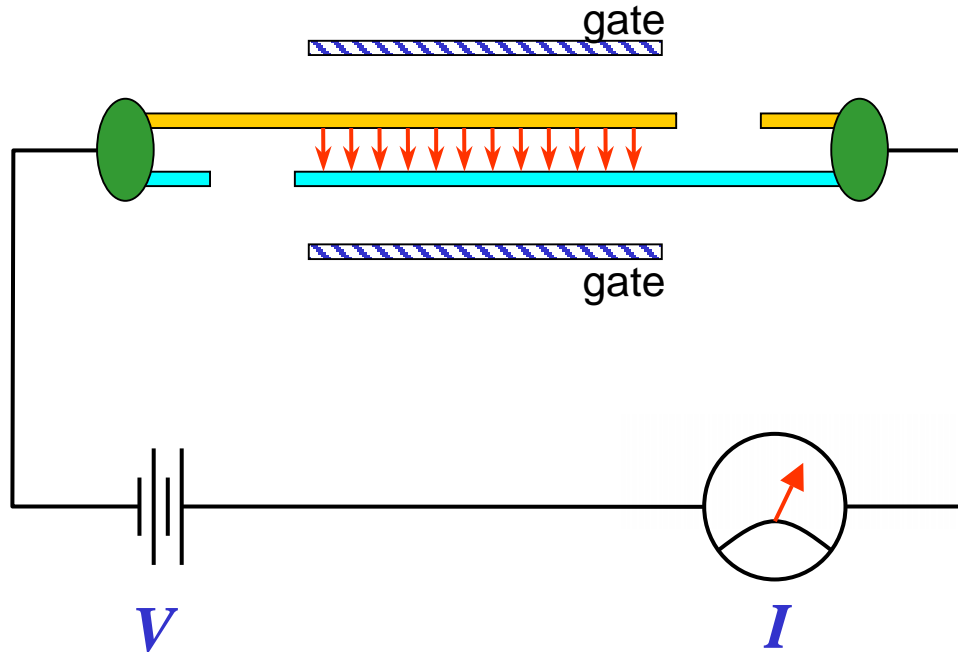
interlayer tunneling



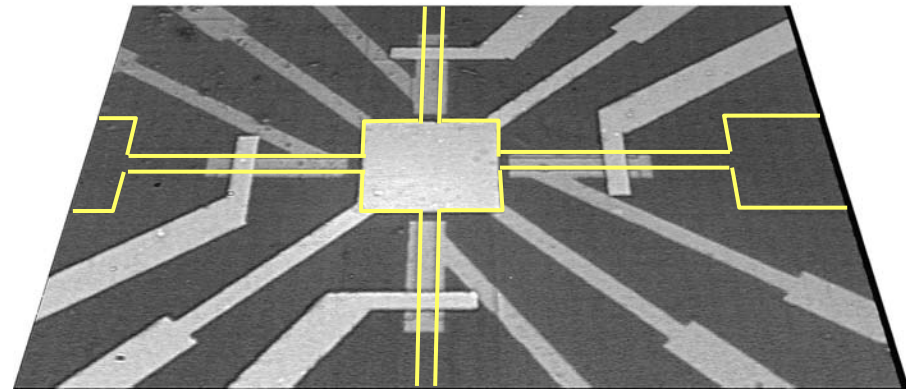
separate layer contacts are essential



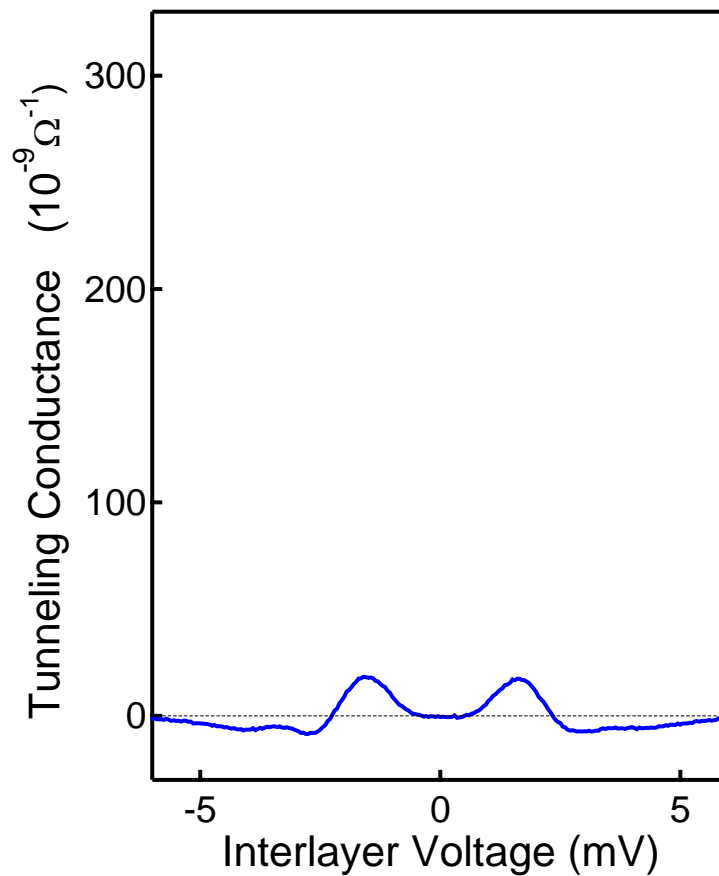
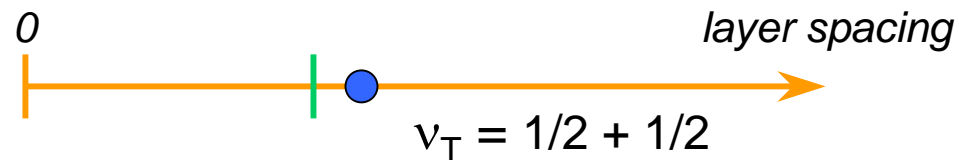
Tunneling Experiment



Symmetric gating allows continuous tuning of *effective* layer spacing in a single sample.



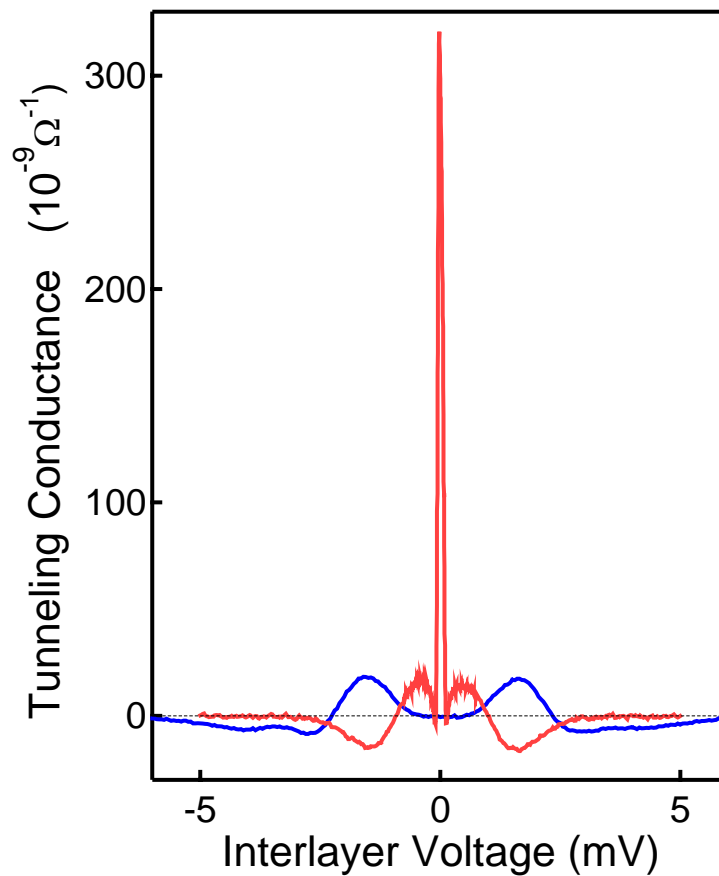
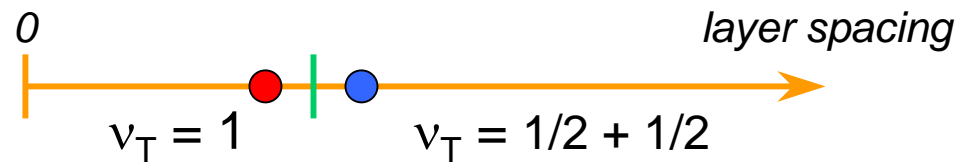
Crossing the phase boundary



Zero bias tunneling heavily suppressed by *intra*-layer correlations.



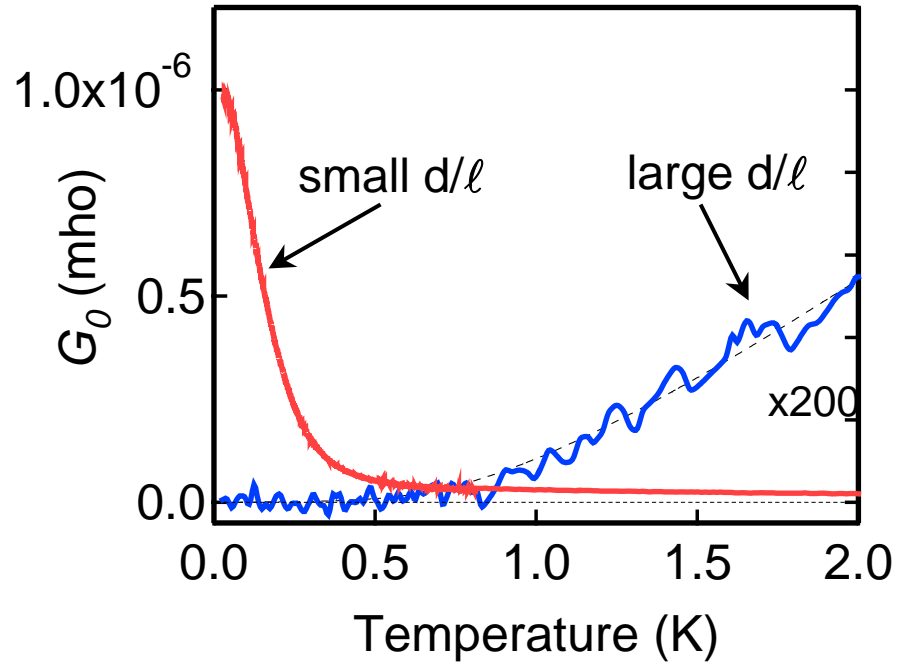
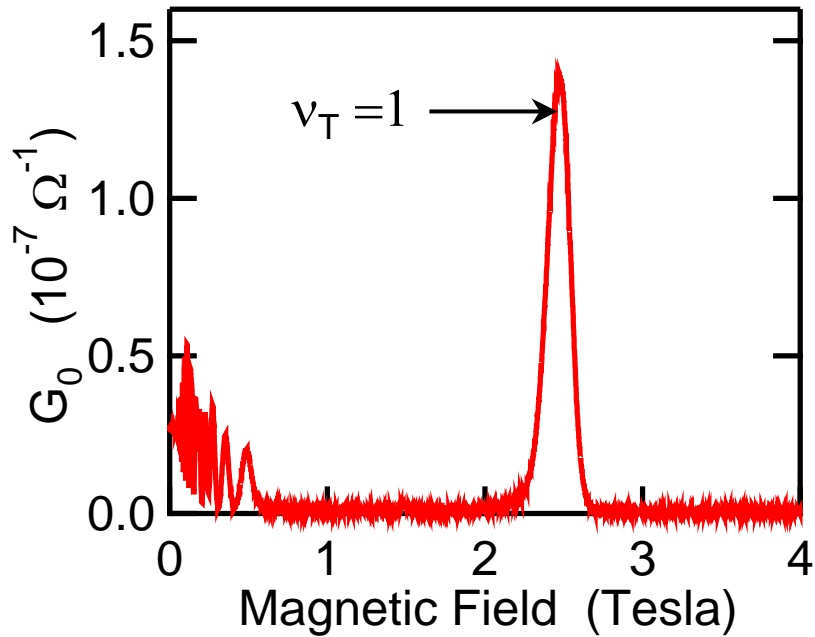
Crossing the phase boundary



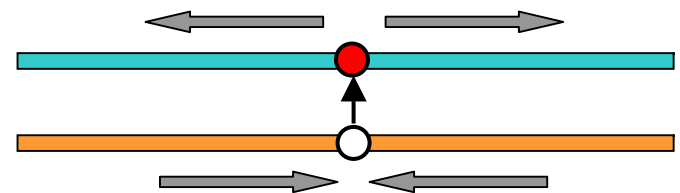
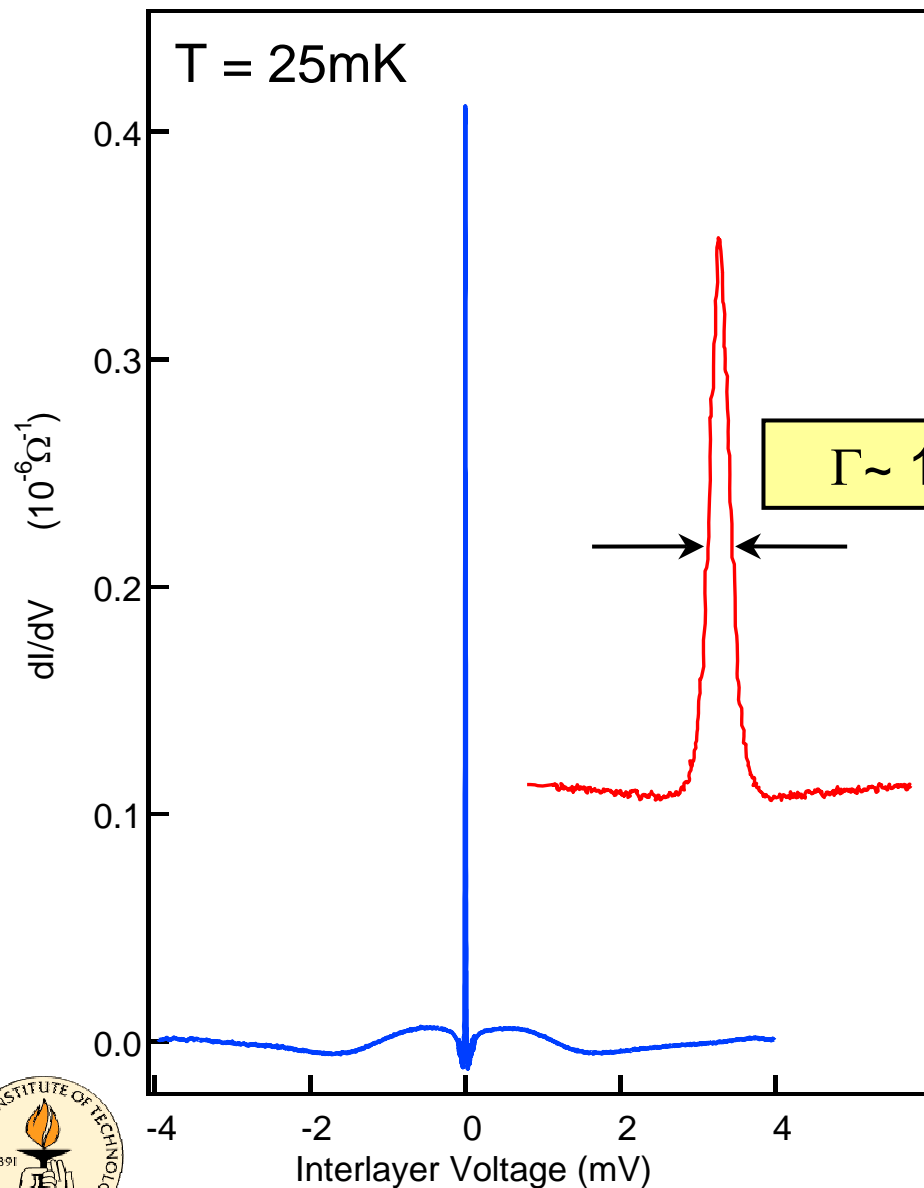
Coulomb gap replaced by resonant enhancement.



Magnetic Field and Temperature Dependences



Extremely Narrow Resonance



Counter-flow superfluidity rapidly relaxes charge defects created by tunneling.

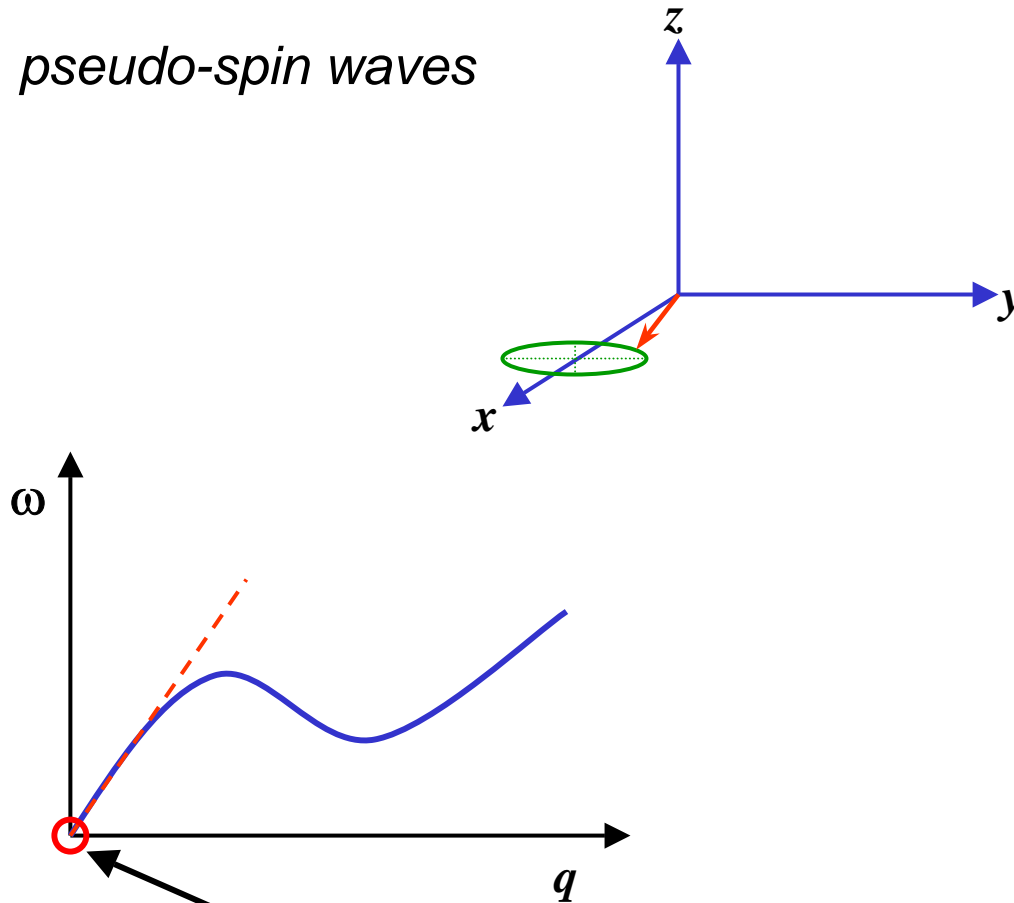


Collective Modes

Fertig '89

Wen & Zee '92

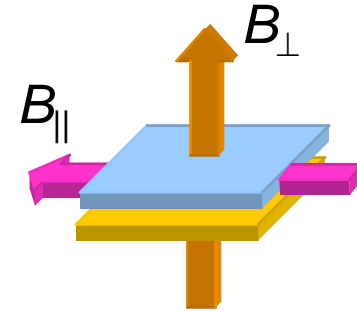
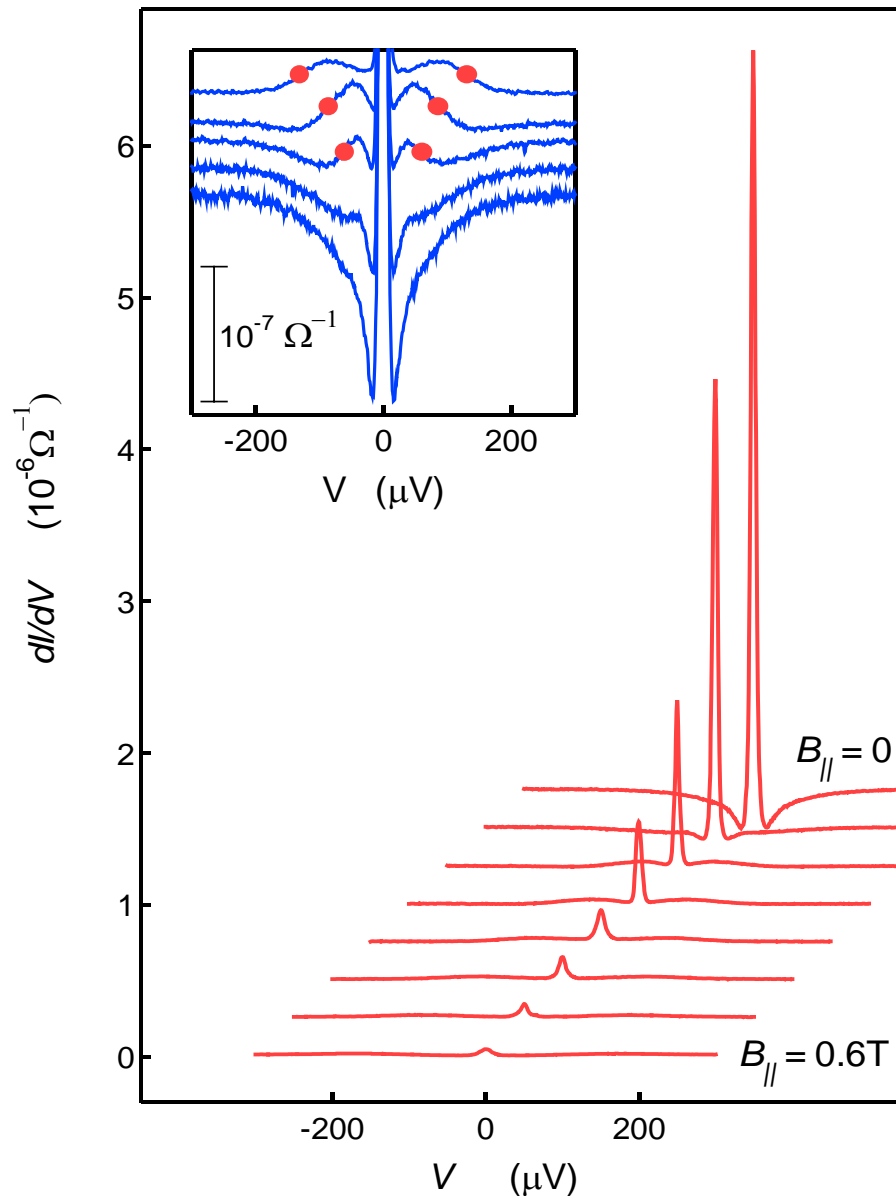
pseudo-spin waves



Tunneling detects the Goldstone mode of the broken symmetry ground state.



Parallel Field Spectroscopy

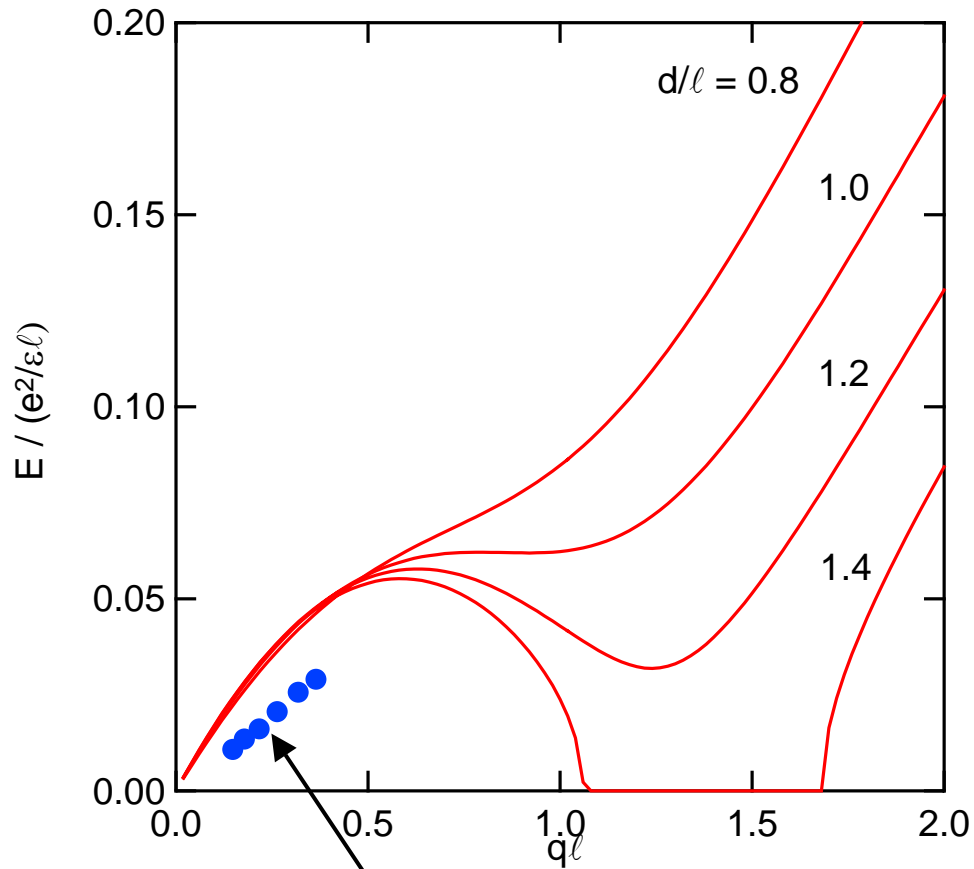


$$q = eB_{\parallel} d / \hbar$$

Peak is suppressed and satellites appear.



Measured Collective Mode Dispersion



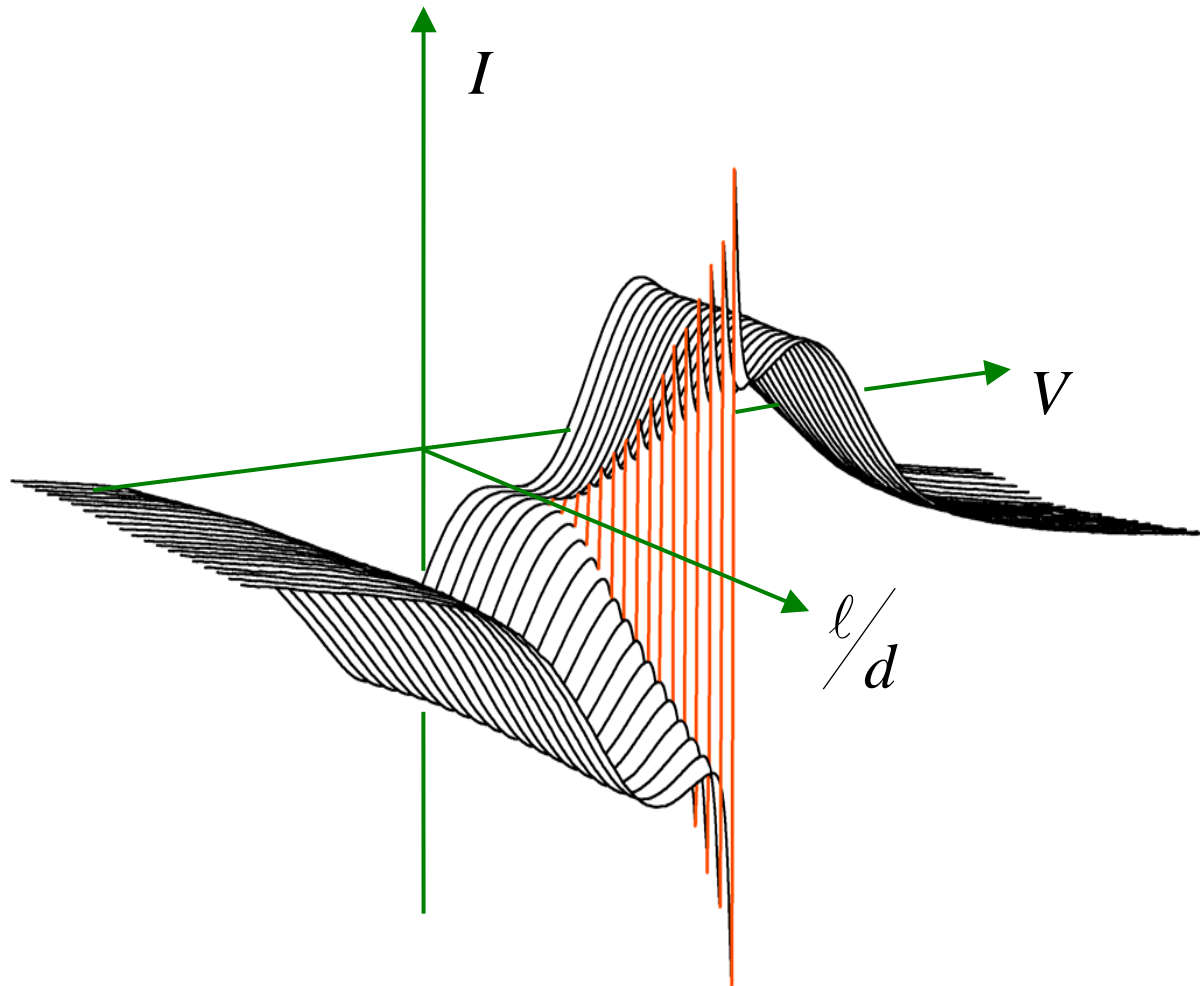
Direct observation of linearly dispersing Goldstone mode.

- red curves: A.H. MacDonald



$v \sim 15 \text{ km/sec}$

A quantum phase transition



Demonstration of spontaneous interlayer phase coherence.



Josephson Effect in a Non-Superconductor?

dc Josephson effect in a superconductor:

$$J_T = J_c \sin \Delta \phi$$

Interlayer tunneling in $\nu_T = 1$ bilayer QHE:

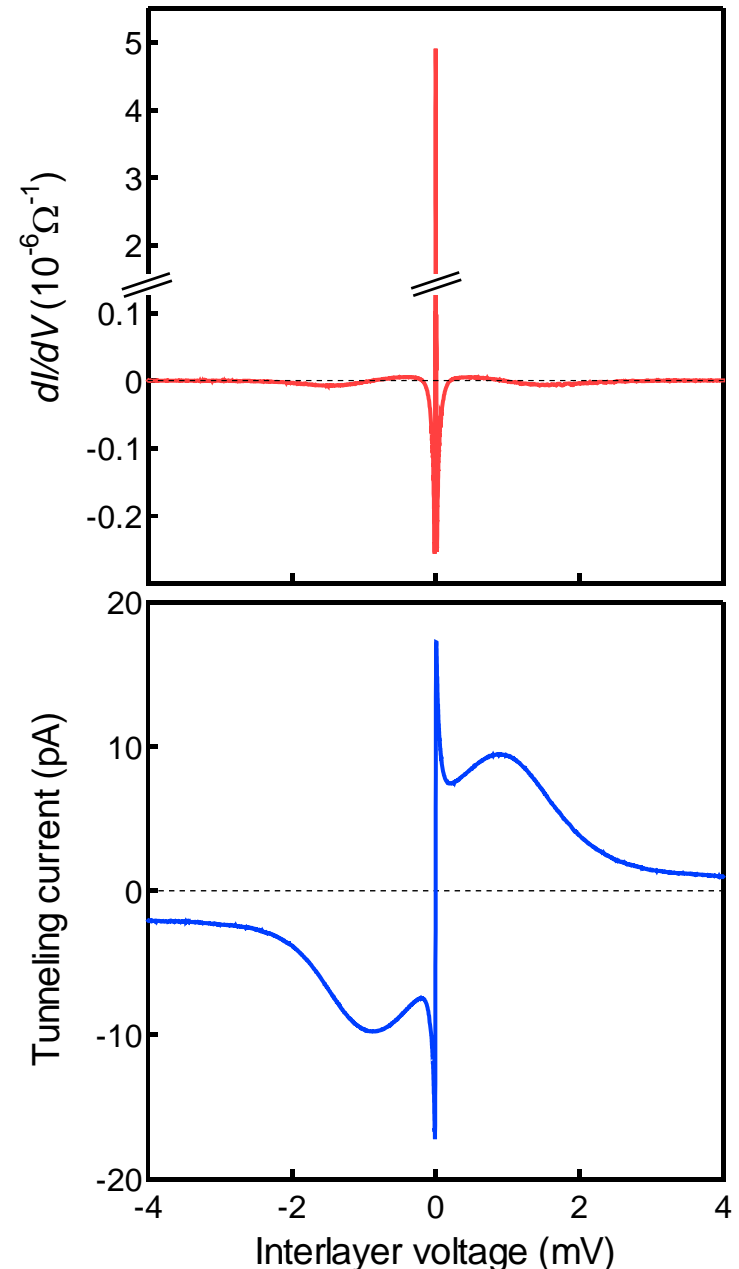
$$J_T = \frac{e}{4\pi\ell^2\hbar} \Delta_{\text{SAS}} \sin \phi$$

QHE critical current:

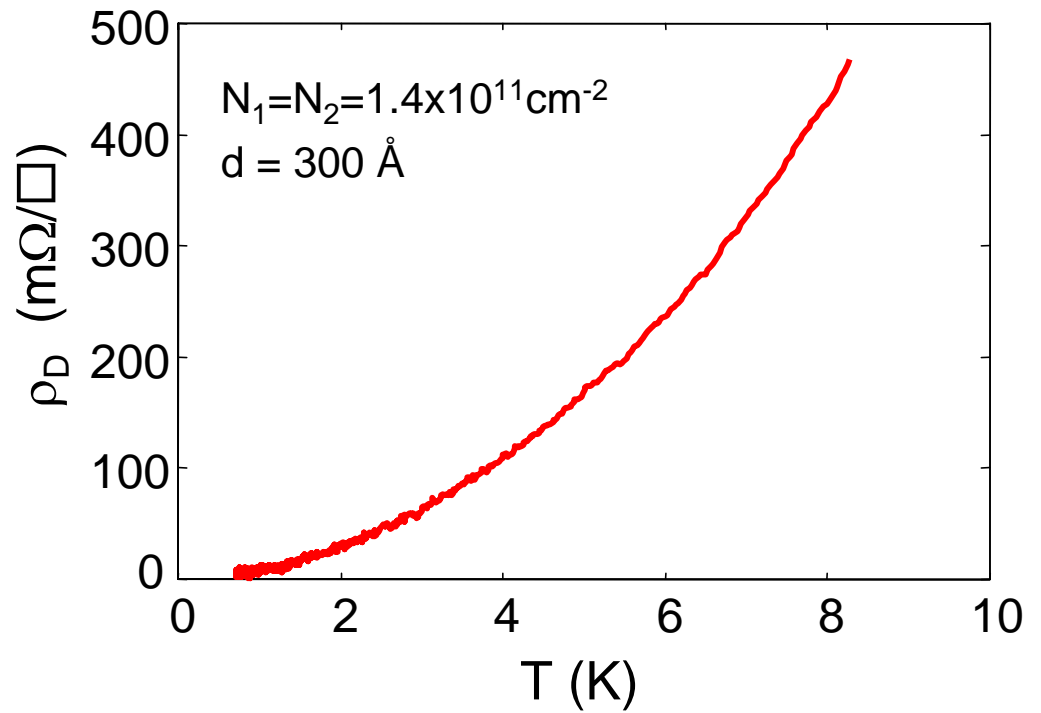
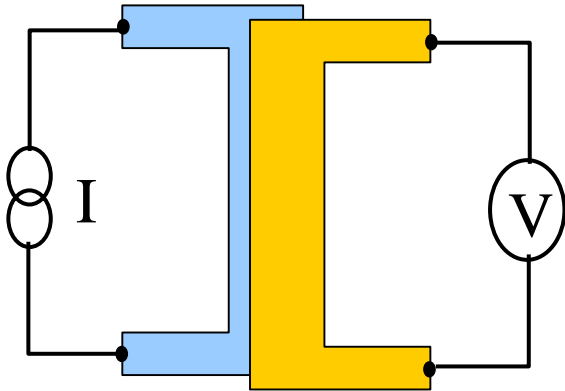
$$I_c = \frac{eA}{4\pi\ell^2\hbar} \Delta_{\text{SAS}} \approx 50 \mu\text{A}$$

Suggests that pseudospin field is heavily disordered on macroscopic scales

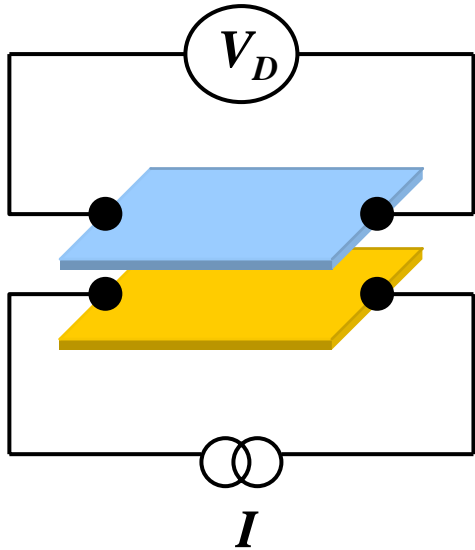
cf. Fertig and Straley, cond-mat/0301128



Coulomb Drag at $B=0$

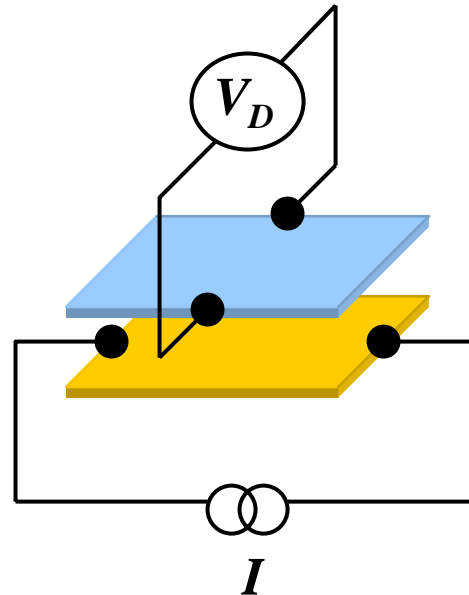


Two Kinds of Coulomb Drag in a Magnetic Field

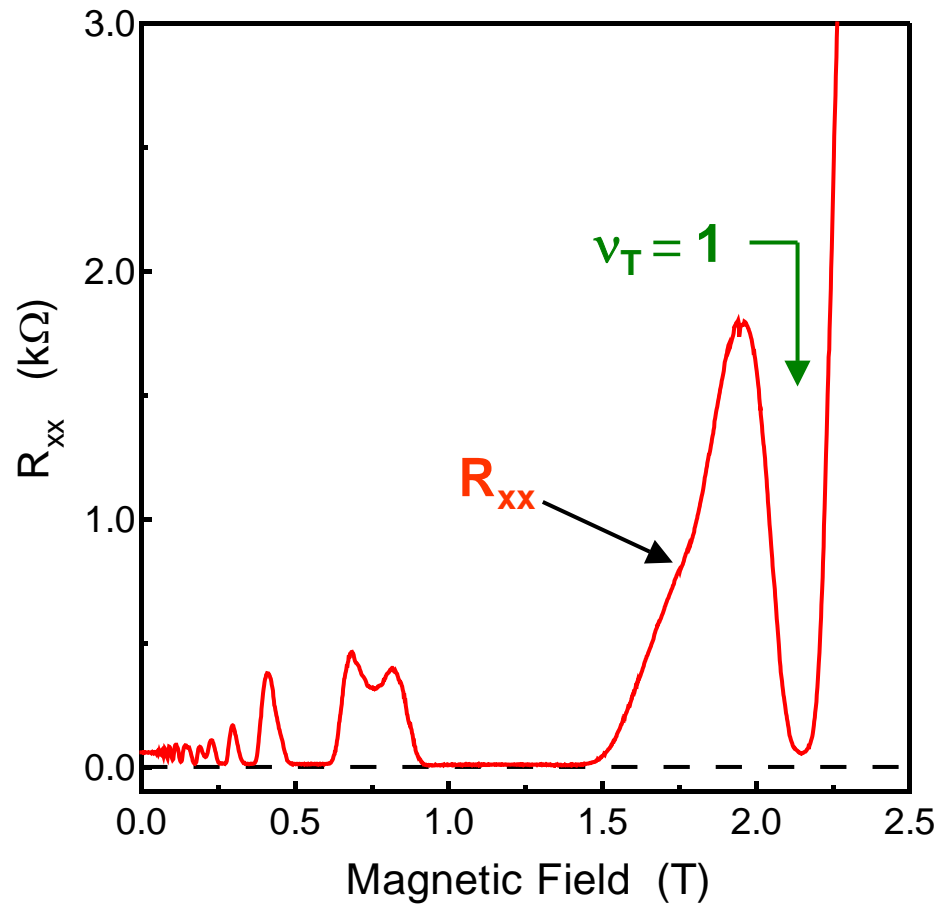


Longitudinal drag

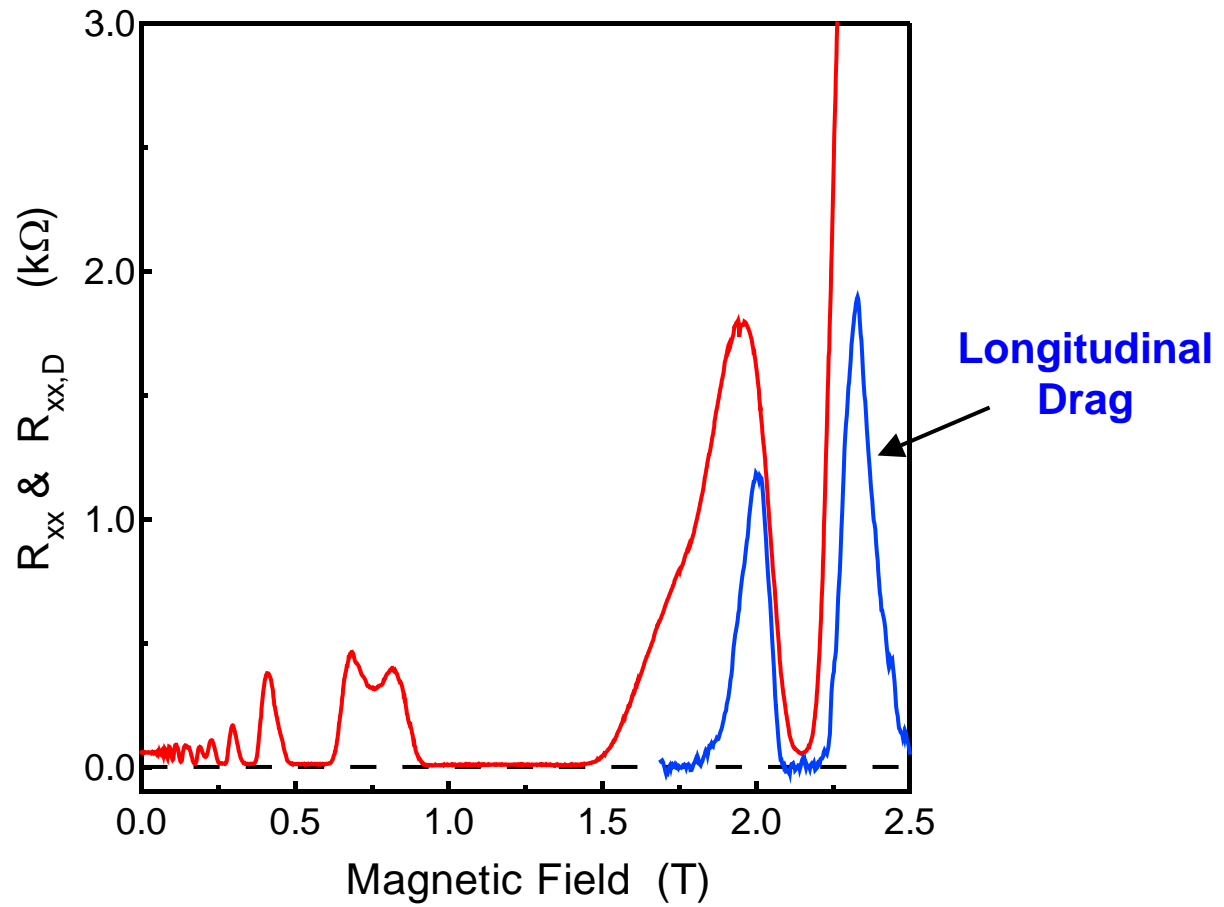
“Hall” drag



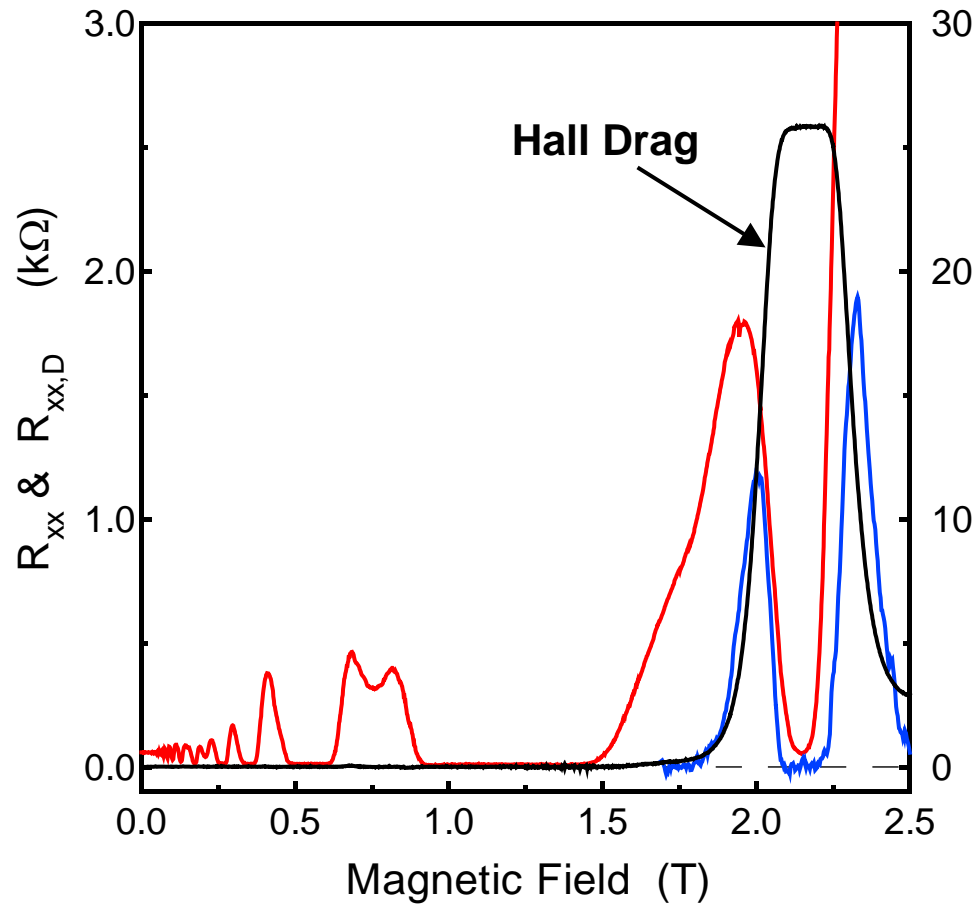
Transport coefficients



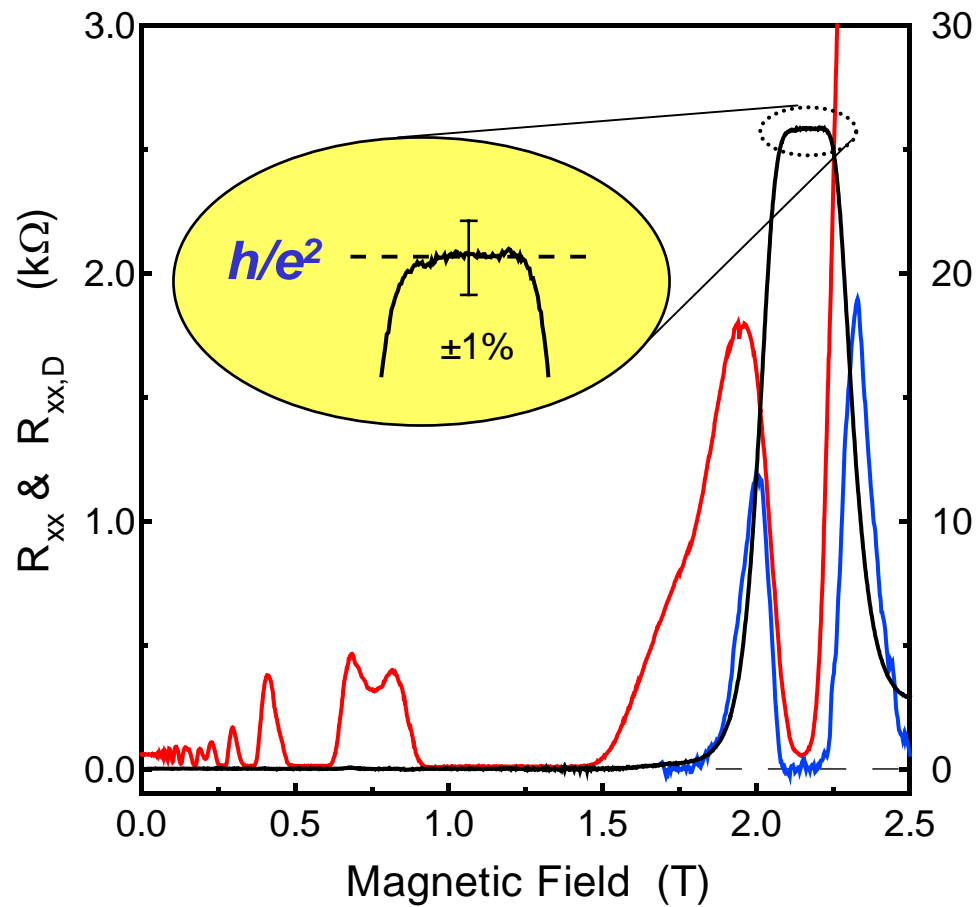
Transport coefficients



Transport coefficients



Transport coefficients

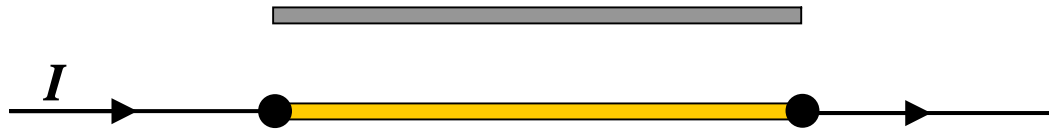


Exact quantization of Hall drag.



Two Transport Channels

In a *drag* measurement, current flows only in one layer:



The $\nu = 1$ excitonic state supports two kinds of transport:
charged quasiparticle currents and *neutral excitonic superflow*.

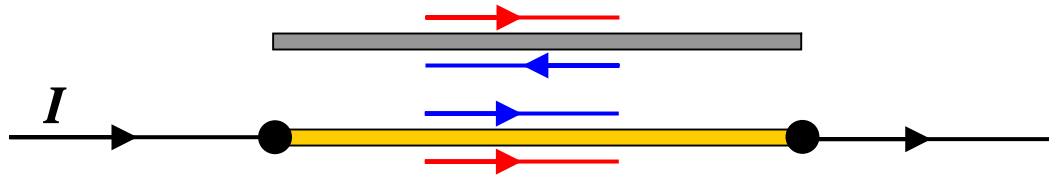


Two Transport Channels

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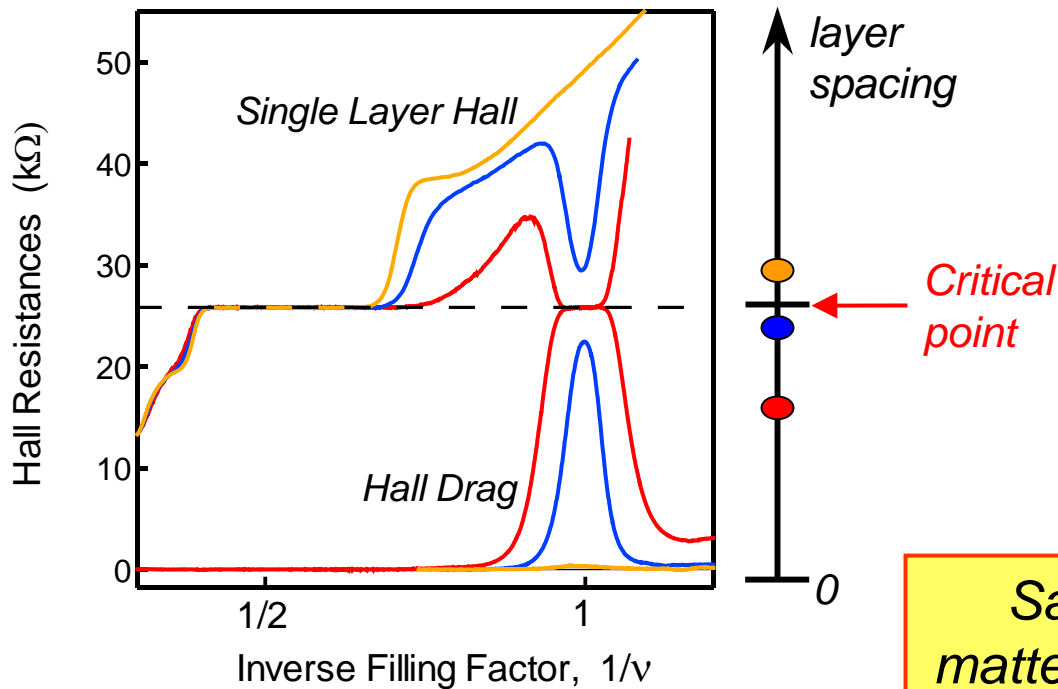
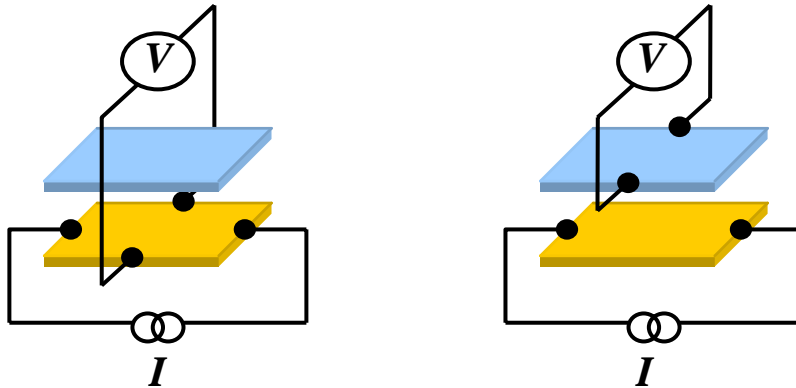
$$I_1 = I_{qp} + I_{ex} = I$$

$$I_2 = I_{qp} - I_{ex} = 0$$

“Explains” quantization of Hall drag, and...



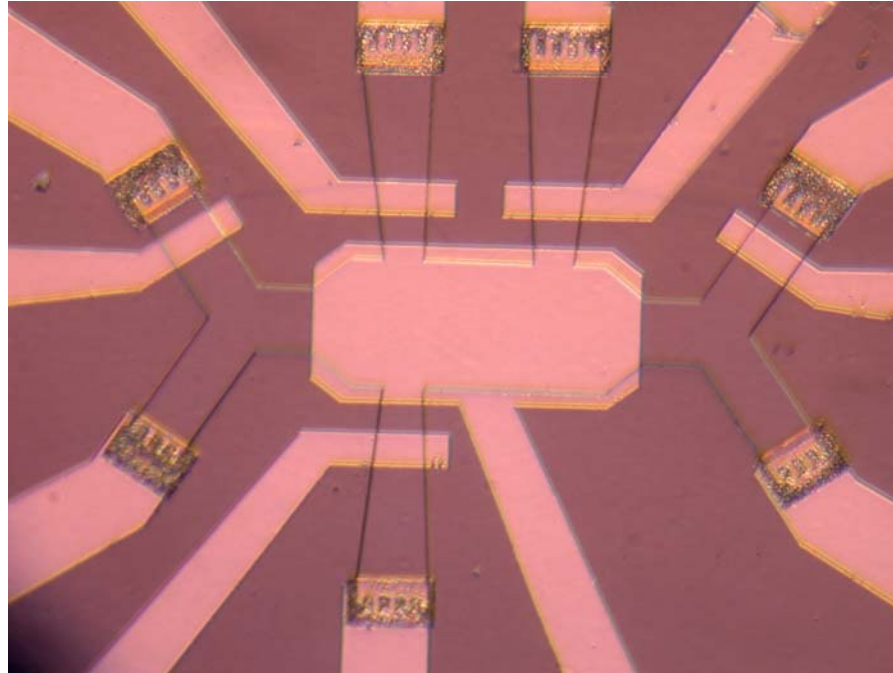
It's stranger still...



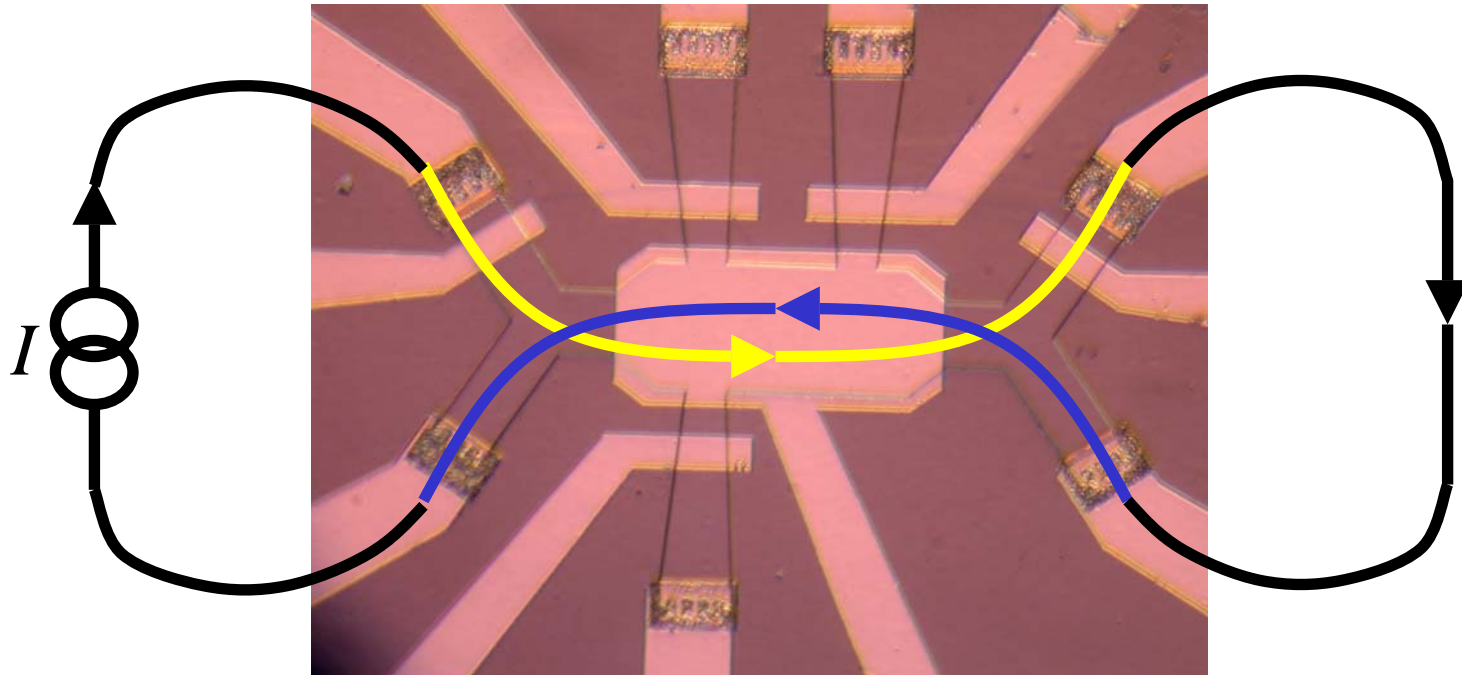
Same Hall voltages no matter how current is divided between the layers!



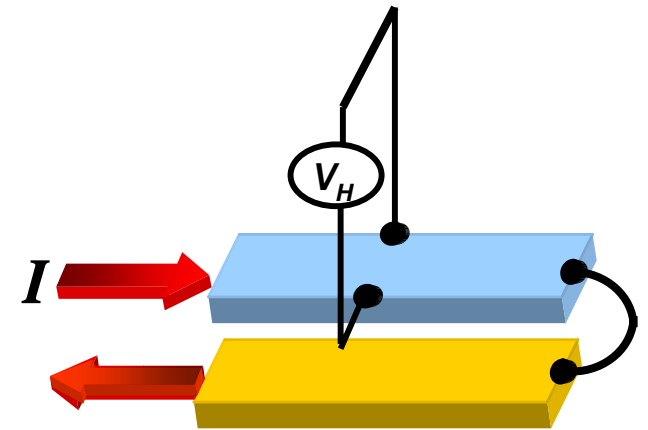
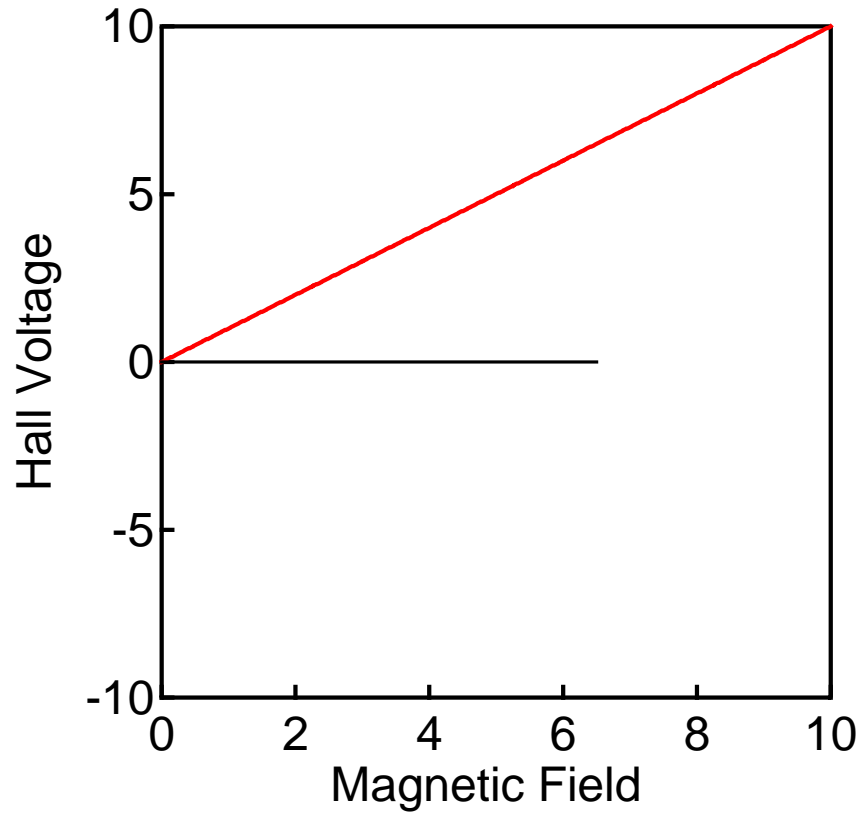
Counterflow Experiment



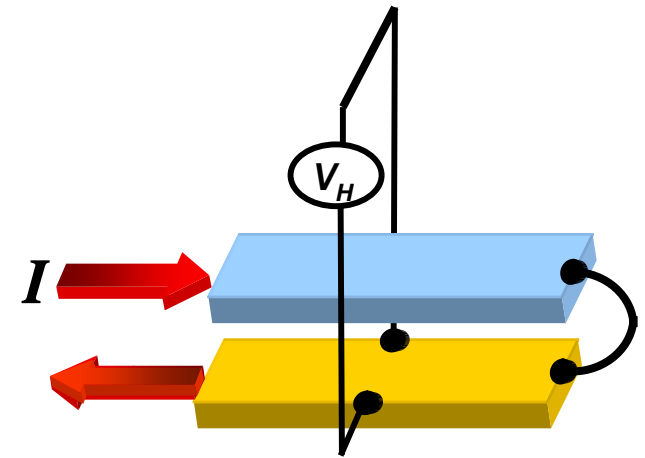
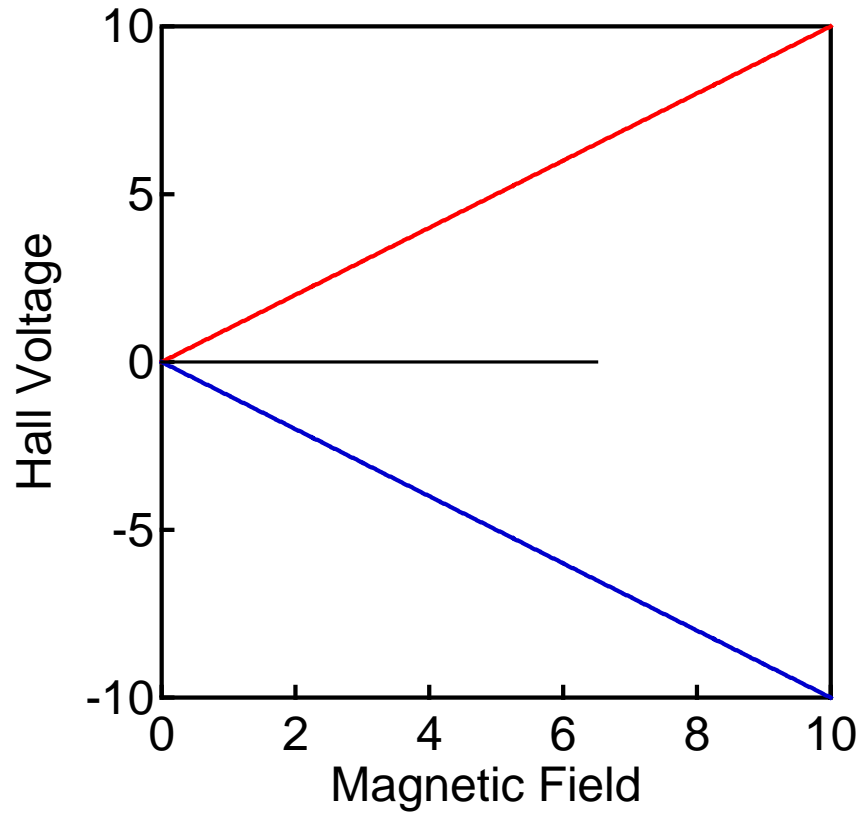
Counterflow Experiment



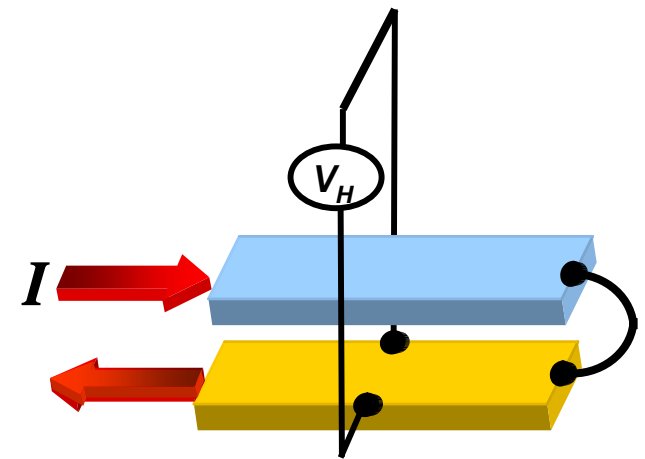
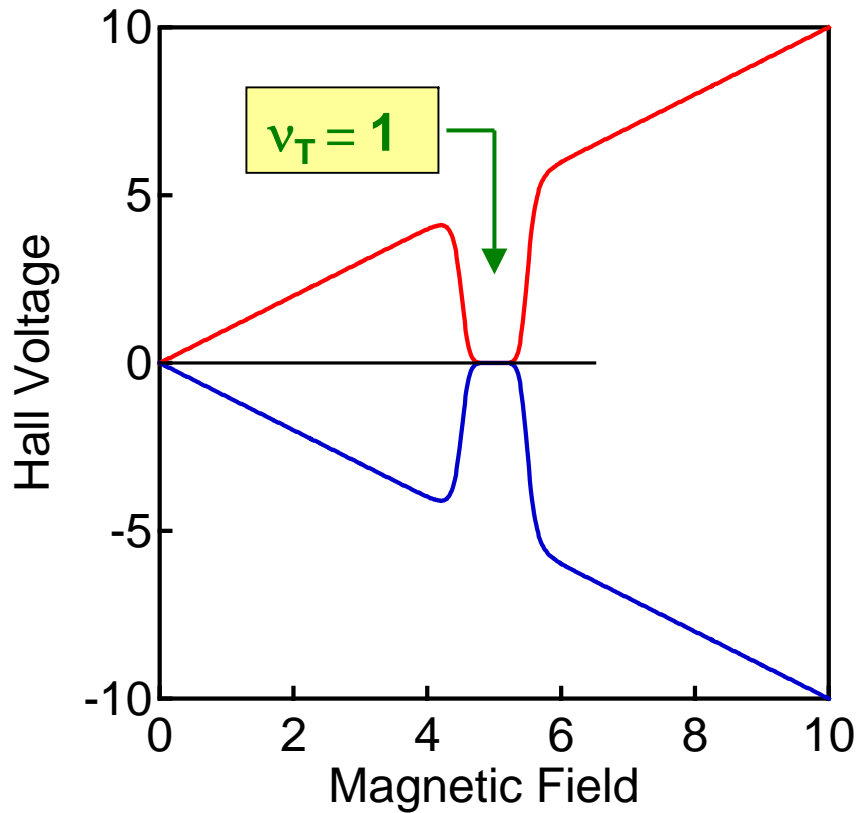
Counterflow Experiment - cartoon



Counterflow Experiment - cartoon



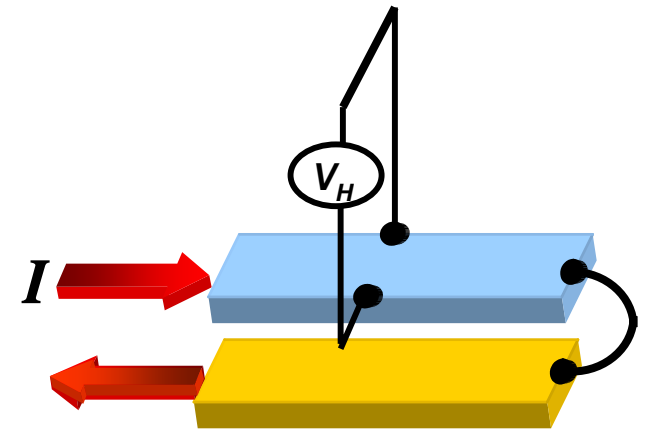
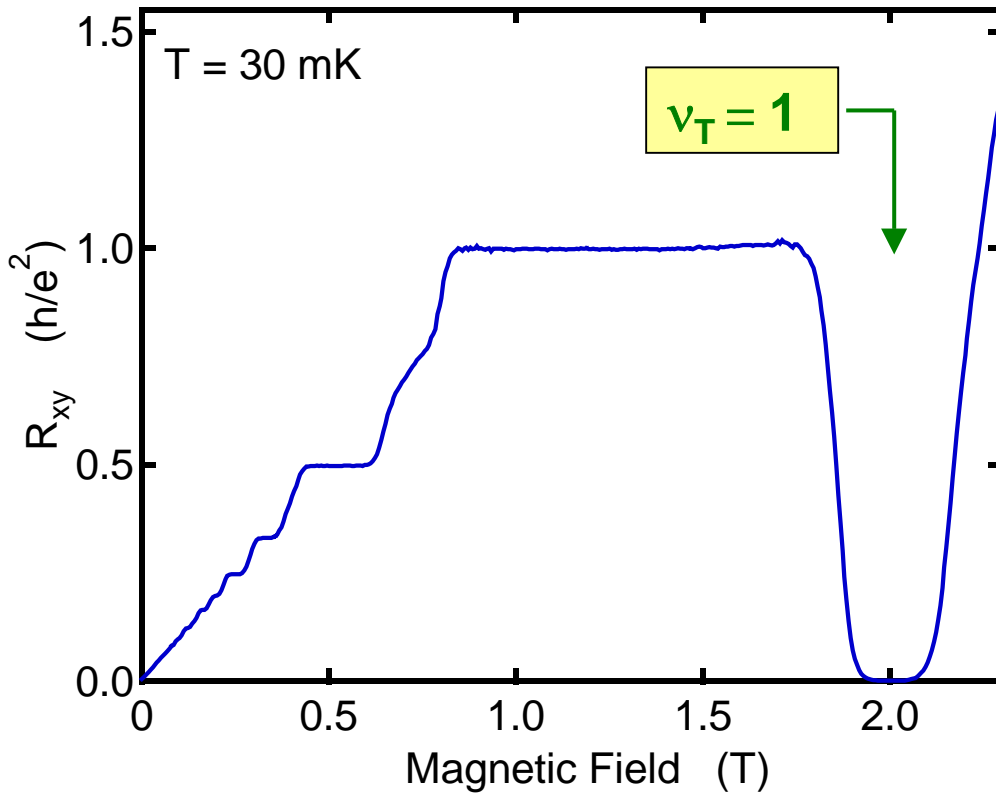
Counterflow Experiment - cartoon



Charge neutral excitons should feel no Lorentz force!



Counterflow Experiment - reality



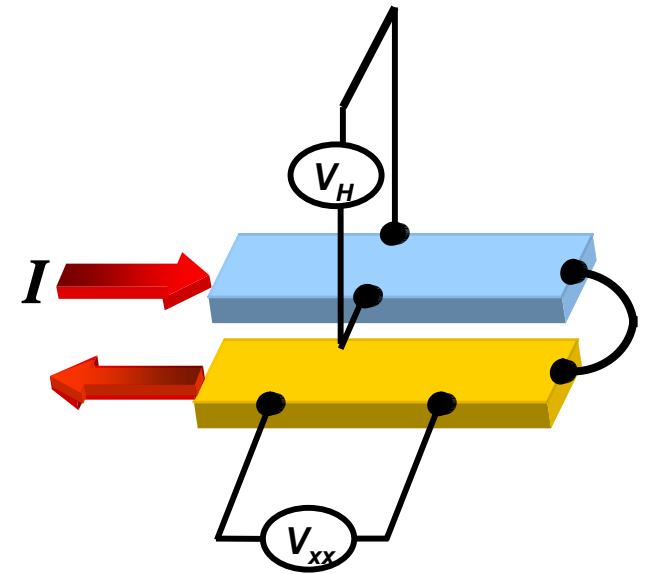
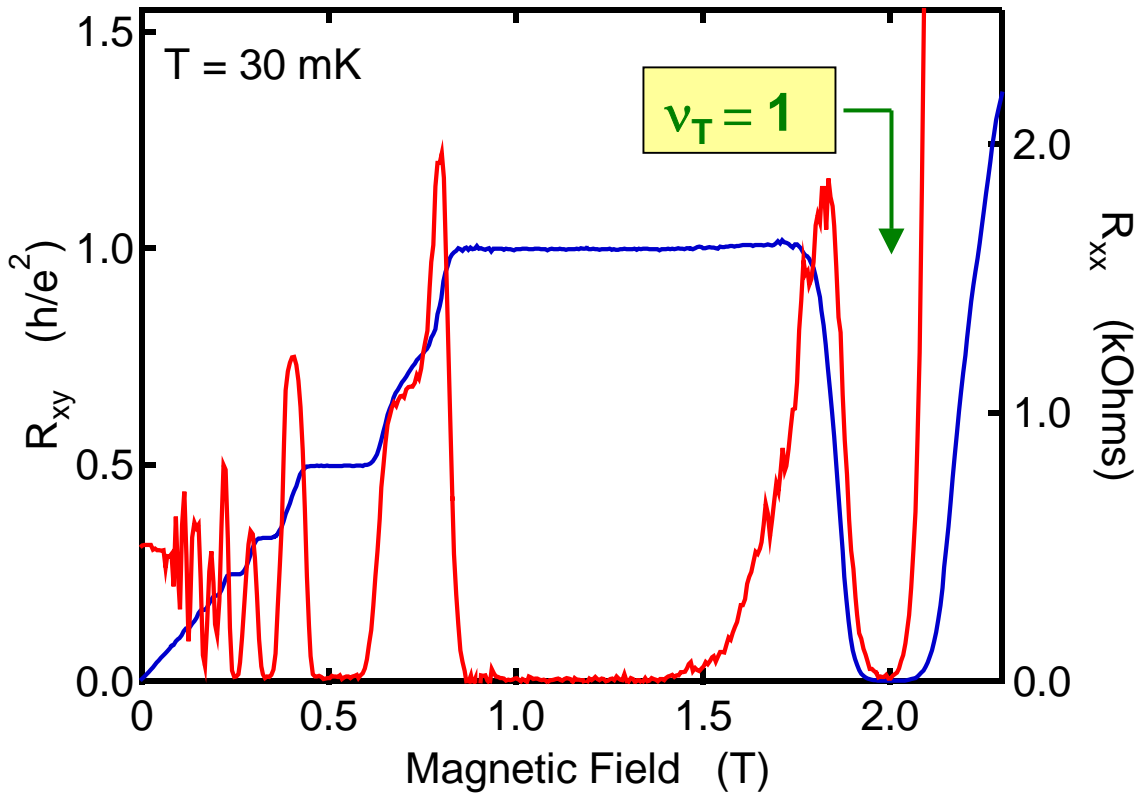
At $\nu_T = 1$ $R_{xy}^{CF} \rightarrow 0$ as $T \rightarrow 0$



exciton transport dominates counterflow



Counterflow Experiment - reality

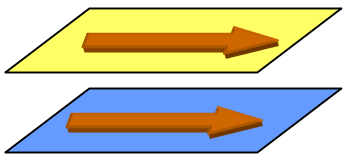


At $\nu_T = 1$: $R_{xy}^{CF} \rightarrow 0$ and $R_{xx}^{CF} \rightarrow 0$

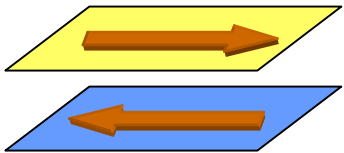


Parallel vs. Counterflow Transport

$$\sigma_{xx} = \frac{R_{xx}}{R_{xx}^2 + R_{xy}^2}$$



$$R_{xx}^{\parallel} \rightarrow 0 \quad \text{and} \quad R_{xy}^{\parallel} \rightarrow h/e^2 \quad \longrightarrow \quad \sigma_{xx}^{\parallel} \rightarrow 0$$

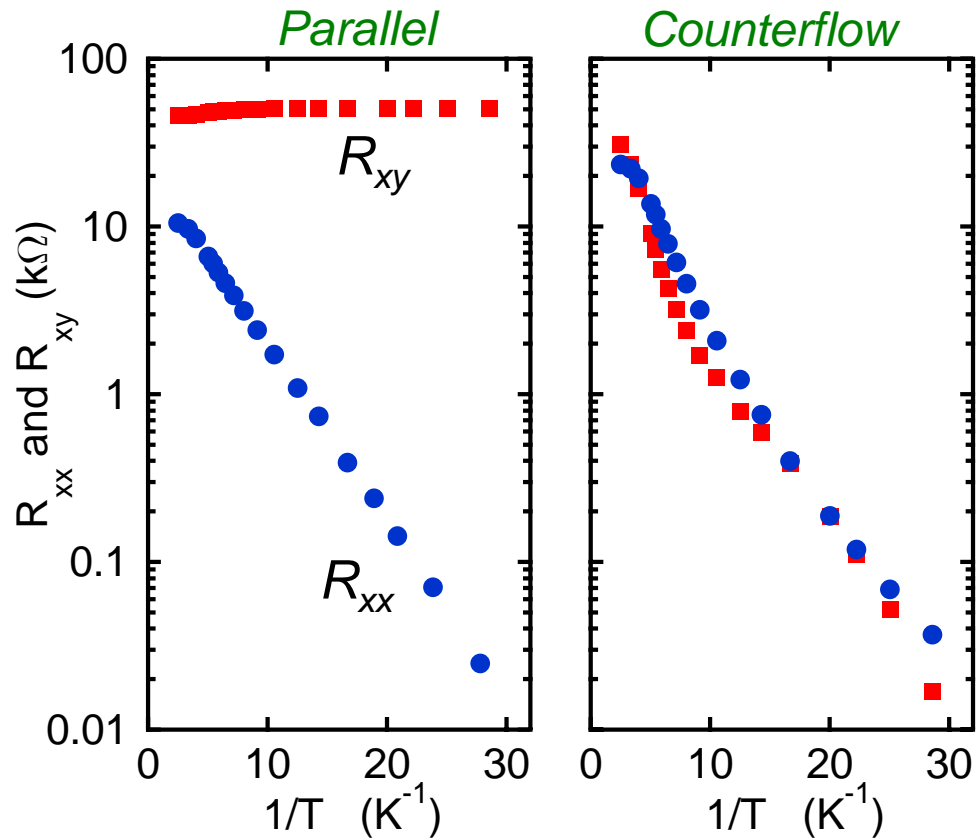


$$R_{xx}^{CF} \rightarrow 0 \quad \text{and} \quad R_{xy}^{CF} \rightarrow 0 \quad \longrightarrow \quad \sigma_{xx}^{CF} \rightarrow ?$$



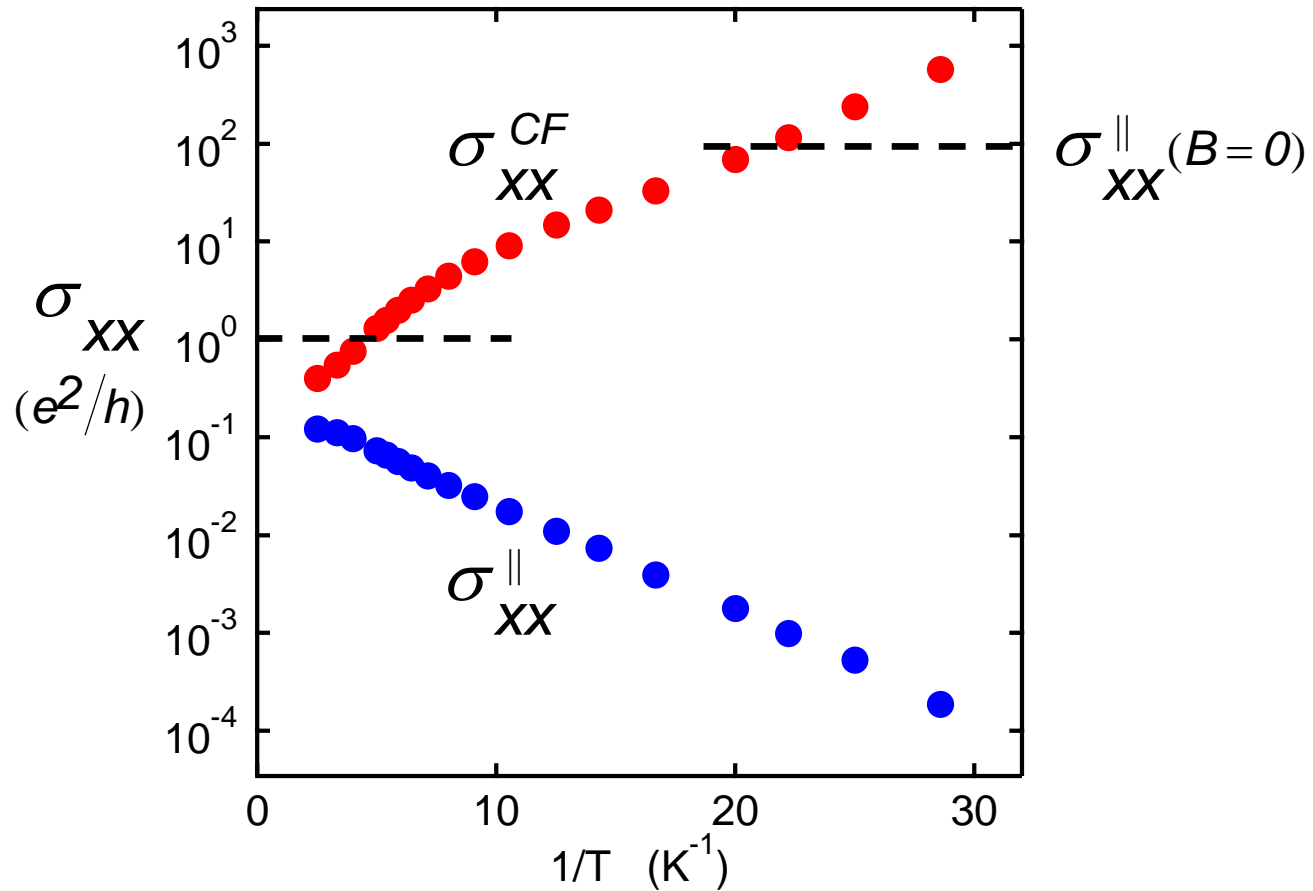
Temperature Dependences

$$d/\ell = 1.5$$



Conductivities

$$d/\ell = 1.5$$



Counterflow dissipation small but non-zero at all finite T .



Results

In closely-spaced bilayer 2D electron systems at $\nu_T = 1$:

Tunneling reveals a phase transition to an interlayer phase coherent state and unveils the Goldstone mode.

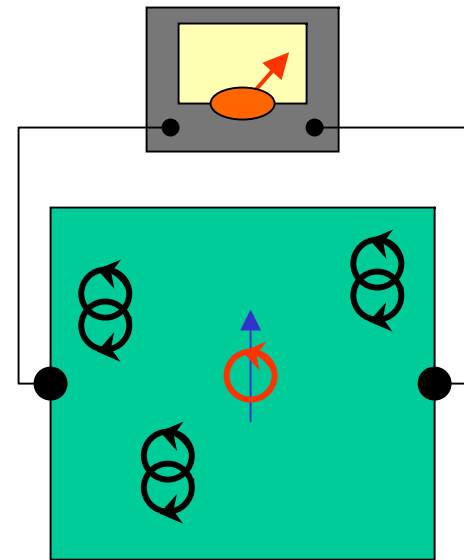
Counterflow and Coulomb drag reveal collective transport of excitons. Dissipationless as $T \rightarrow 0$.



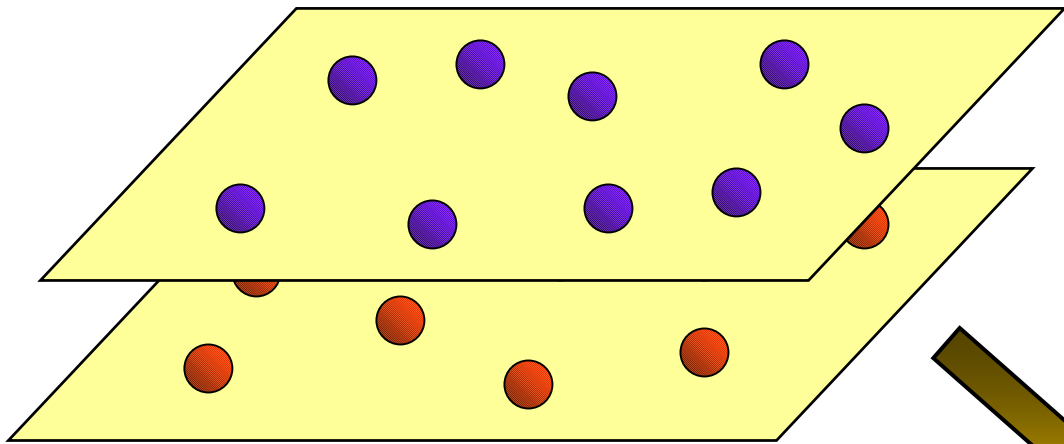
Puzzles

- *tunneling conductance peak too small*
- *$B_{||}$ -dependence of tunneling not understood*
- *no non-linearity observed in counterflow transport*
- *thermally-activated dissipation suggests $T > T_{KT}$*

Disorder induced free vortices present at all temperatures.



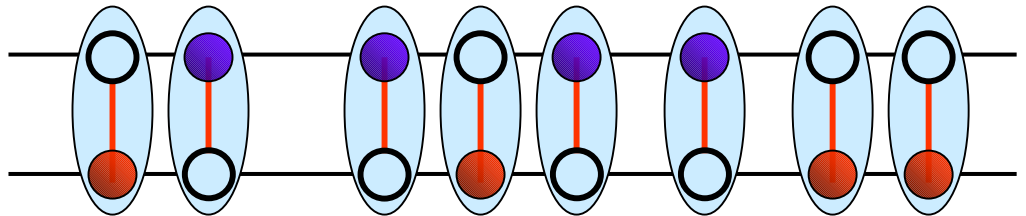
Quantum Hall Excitonic Bose Condensate



Start with a double layer 2D
electron gas

Add a magnetic field

*Presto! A BCS-like superfluid
comprised of interlayer excitons.*



Keldysh and Kopaev 1964; Lozovik and Yudson; Shevchenko 1976

