Quantum computational complexity, phase transitions and glassiness

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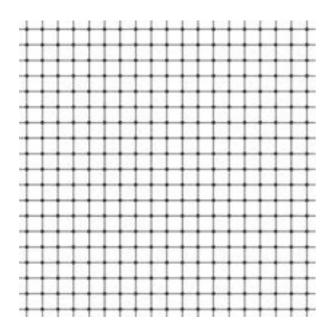
KITP, August 23, 2010



- Complexity theory for physics
- Physics for complexity theory

Complexity theory: the lightning intro

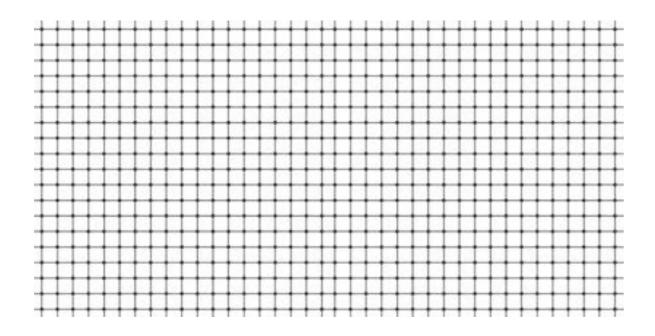
Complexity theory classifies problems according to the scaling of the resources a computer requires to solve large instances





Complexity theory: the lightning intro

Complexity theory classifies problems according to the scaling of the resources a computer requires to solve large instances

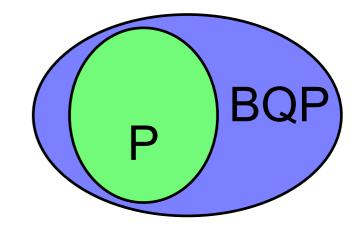




Complexity theory: the lightning intro

Complexity theory classifies problems according to the scaling of the resources a computer requires to solve large instances

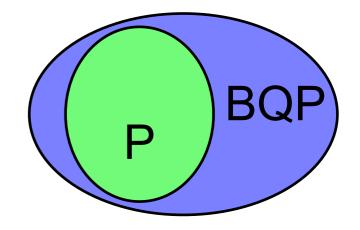
- P -- solvable in polynomial time by a classical computer
- BQP -- solvable in polynomial time by a quantum computer



P and BQP

Can identify efficiently solvable (easy!) problems directly: find polynomial time algorithms

- Classical (P):
 - What is energy of given configuration in 3D Ising model?
 Arithmetic ~500AD
- Quantum (BQP):
 - Is there a factor 1



Shor 1994

NP and QMA

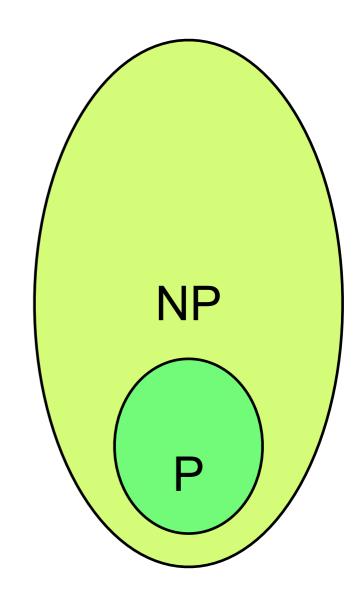
Identify reasonable problems by finding polynomial times algorithms to check them

- NP -- checkable in poly time by a classical computer
 - Is the ground state energy of Ising model below E?
- QMA -- checkable in poly time by a quantum computer
 - Is the ground state energy of a local Hamiltonian below E?
 Kitaev 200²

m		
	NP	QMA
	1 11	



There are hard problems.*

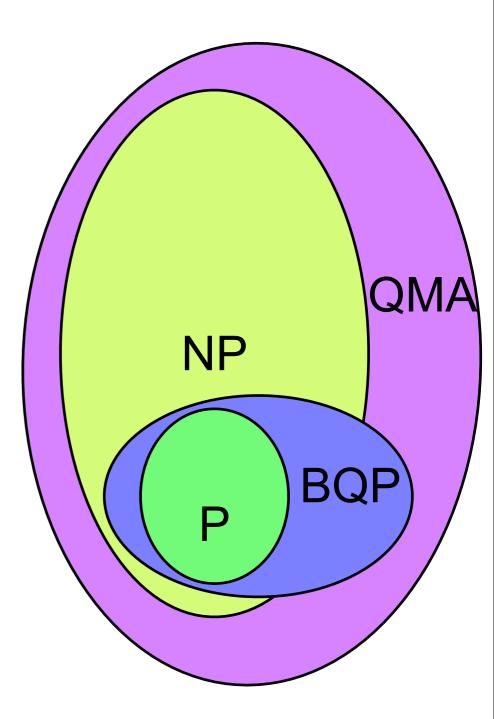


* Conjecture



There are hard problems.*

(even with quantum computers)



* Conjecture

Strong Church-Turing hypothesis

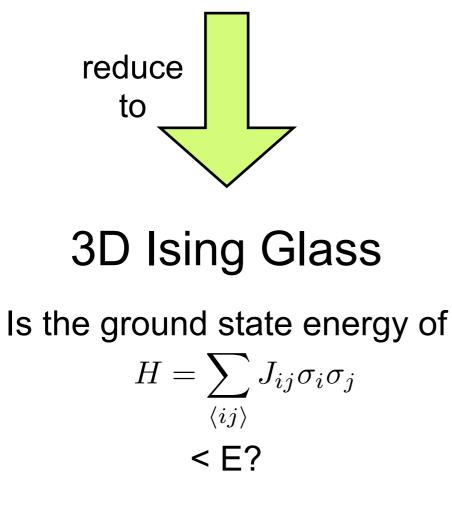
A computer can *efficiently* simulate any physical model of computation.

- All physical models of computation are equivalent
- Any physical object undergoing natural dynamics can be viewed as a computer.
- If $P \neq NP$ there must be glassy physical systems

Guilt by association: -completeness

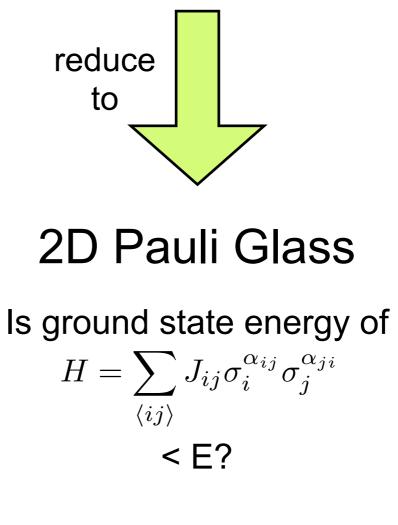
Identify hard problems by circumstantial evidence





Barahona 1982

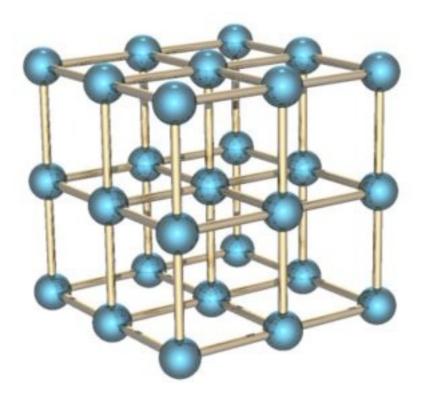
All classical verifiable problems (NP) All quantum verifiable problems (QMA)



Oliveira, Terhal 2005

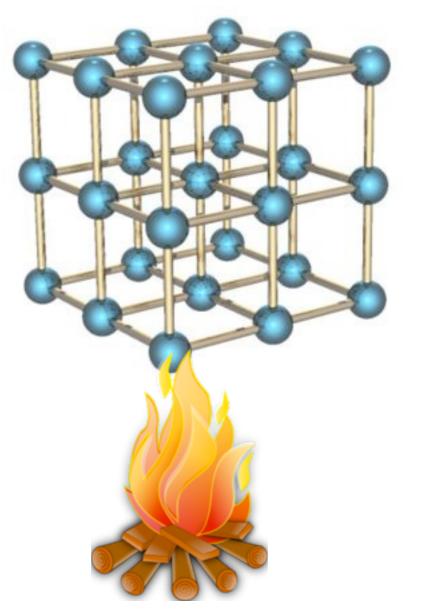
Extended hardness

• No physical process can find ground state of H efficiently.



Extended hardness

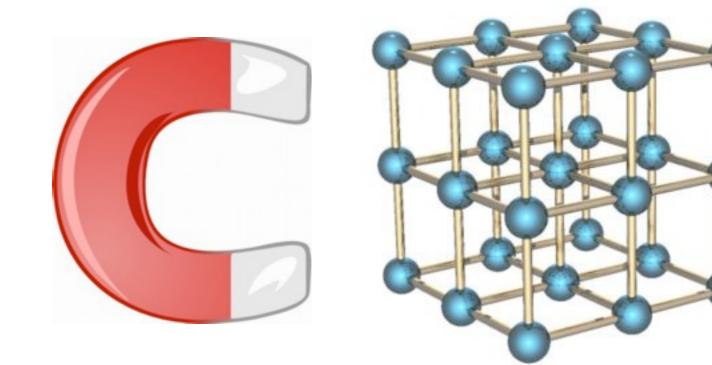
• No physical process can find ground state of H efficiently.



Thermal annealing

Extended hardness

• No physical process can find ground state of H efficiently.



Thermal annealing

Adiabatic annealing

Some -complete problems for physicists

NP-complete ground state energy problems:

k-SAT

q-state Potts

2D translationinvariant tiling Planar Ising glass in a field

2D Pauli Glass

3D Ising glass

QMA-complete ground state energy problems:

k-Local Hamiltonian

k-QSAT*

ID translation-

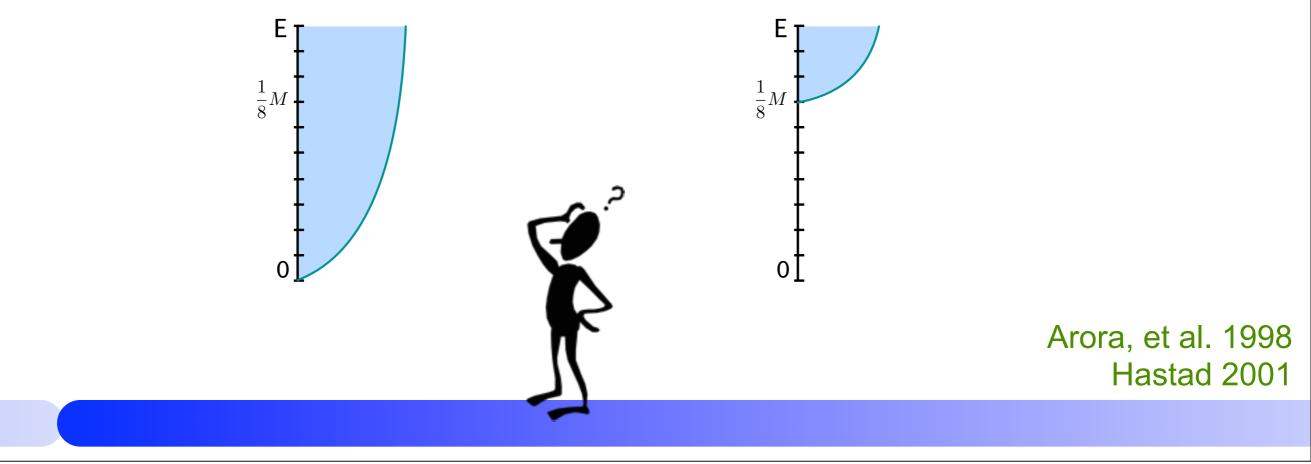
invariant Hamiltonian

Cook 1971; Levin 1973 Garey, Johnson 1979 Barahona 1982 Kitaev 2001 Bravyi 2006 Oliveira, Terhal 2005 Gottesman 2009

*QMA₁-complete

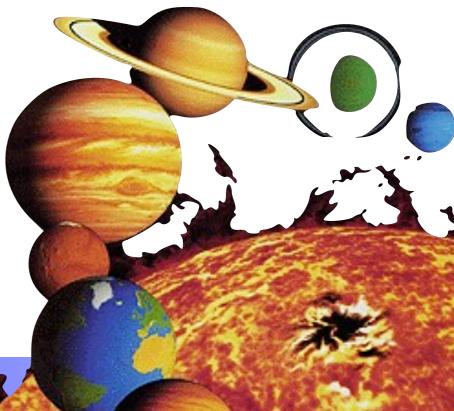
Finite temperature slowness and PCP

- The PCP theorem asserts that distinguishing certain optimization problems with zero ground state energy from those with extensive ground state energy is NP-complete.
- Annealing to a finite temperature must be slow for these.
- Best current result: 3-SAT cannot be cooled to any finite temp.



$P \neq NP$ as physical principle

- No physical process can solve NP-complete problems in polynomial time.
 - Quantum mechanics must be linear
 - Closed time-like curves are forbidden
 - No hidden variables (almost).
 - etc



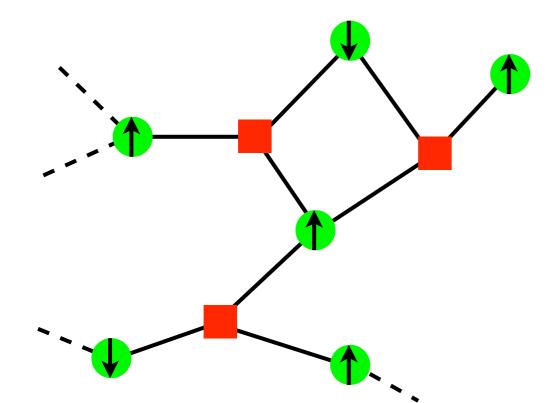
Aaronson 2005

Statistical mechanics of hard problems

- Complexity theory shows hardness of -complete problems
- Suggests optimization problems which should exhibit glassiness
- Does not reveal mechanisms or features underlying slowness

 Studying ensembles of -complete problems not only shows glassiness but also glass transitions

Classical 3-SAT: An 'Ising' model



N bits $\vec{\sigma} \in \{\pm 1\}^N$ M constraints $E^m = \delta_{\sigma_{m_1},\phi_1^m} \delta_{\sigma_{m_2},\phi_2^m} \delta_{\sigma_{m_3},\phi_3^m}$ $H = \sum E^m(\sigma_{m_1},\sigma_{m_2},\sigma_{m_3})$

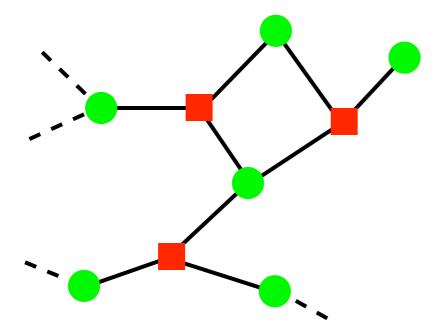
 E^m is 0 for satisfying states

 $m \in G$

Is the ground state energy zero? $\exists \vec{\sigma} \text{ s.t. } E^m(\sigma_{m_1}, \sigma_{m_2}, \sigma_{m_3}) = 0 \ \forall m \in G?$

Ensemble of 3-SAT: Average complexity

Random 3-SAT



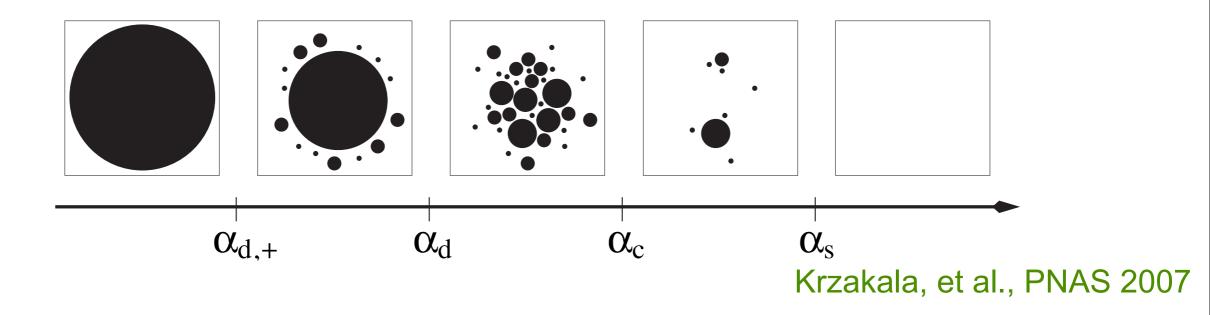
Clause density $\alpha = M/N$

Random graph Disordered couplings

- Which instances are hard?
- Ensembles of 'typical' instances
- Control parameters
- Spin glass physics

Fu, Anderson 1985; Levin 1986

Classical glass theorist's phase diagram

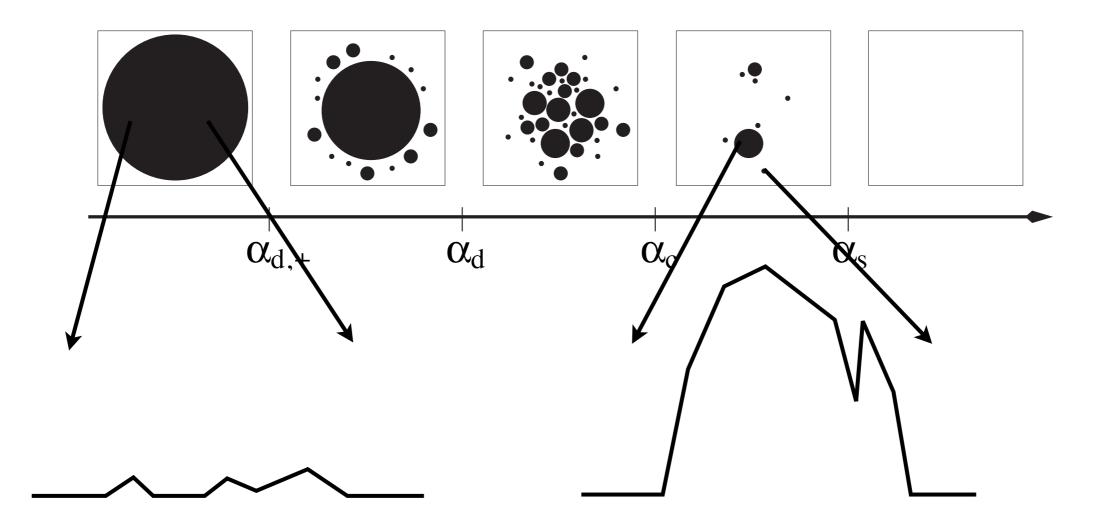


- Qualitative phase diagram of random constraint satisfaction problems
- Phase transitions: clustering of satisfying assignments
- Based on RSB cavity methods
- Quantum cavity methods?

cf. Mezard, et al., Science, 2002

CRL, et al., PRB 2008 Hastings, PRB 2007 Leifer, Poulin, Ann Phys 2008

Clustering



- Barrier to local search dynamics and relaxation classically
- Many-body localization in adiabatic algorithm
 Altshuler, et al., 2009

Statistical physics of constraint satisfaction

	Classical		
Worst-case complexity	NP-completeness		
Statistical physics	Satisfiability transitions Dynamical transitions Clustering transitions		
Heuristic algorithms	Simulated annealing Belief propagation Survey propagation		

Statistical physics of constraint satisfaction

	Classical	Quantum
Worst-case complexity	NP-completeness	QMA-completeness
Statistical physics	Satisfiability transitions Dynamical transitions Clustering transitions	?
Heuristic algorithms	Simulated annealing Belief propagation Survey propagation	??

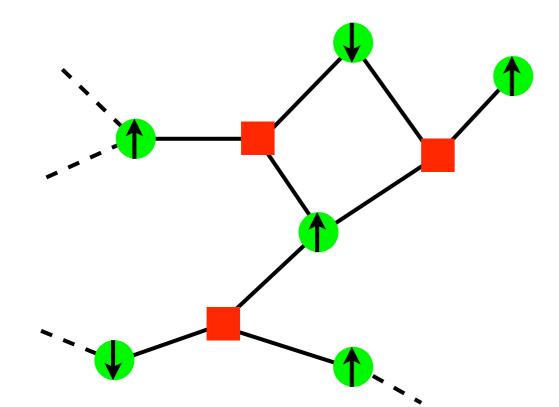
Statistical physics of constraint satisfaction

	Classical	Quantum
Worst-case complexity	NP-completeness	QMA-completeness
Statistical physics	Satisfiability transitions Dynamical transitions Clustering transitions	Satisfiability transitions Entanglement transitions
Heuristic algorithms	Simulated annealing Belief propagation Survey propagation	Adiabatic algorithm? Quantum Metropolis?

Quantum Satisfiability

- Natural quantum generalization of classical satisfiability (k-SAT)
- Quantum hard worst-case complexity: QMA₁-complete
- Are typical instances hard?
- Motivated by classical story, but has its own features...

Quantum k-QSAT: A k-local qubit model



N qubits $\mathcal{H} = (\mathbb{C}^2)^{\otimes N}$ M constraints $\Pi^m = |\phi^m\rangle\langle\phi^m|$ $H = \sum_{m \in G} \Pi^m$

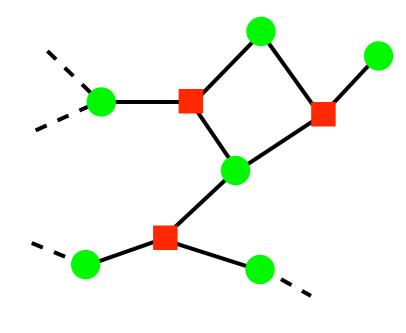
 Π^m penalizes 1 out 2^k states

Is the ground state energy zero?

 $\exists |\psi\rangle \in \mathcal{H} \text{ s.t. } \Pi^m |\psi\rangle = 0 \ \forall m \in G?$

Ensemble of k-QSAT: Average complexity

Random graph



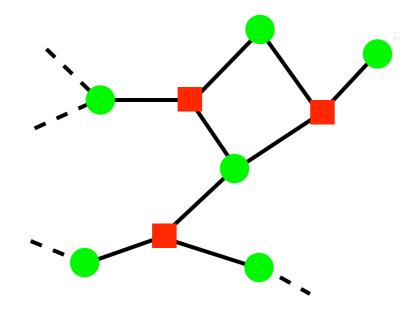
Discrete

Clause density $\alpha = M/N$ Place edges w.p. $p = \alpha/\binom{N}{k}$

CRL, et al., 2009

Ensemble of k-QSAT

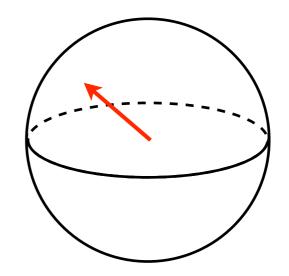
Random graph



Discrete

Clause density $\alpha = M/N$ Place edges w.p. $p = \alpha/\binom{N}{k}$

Generic projectors

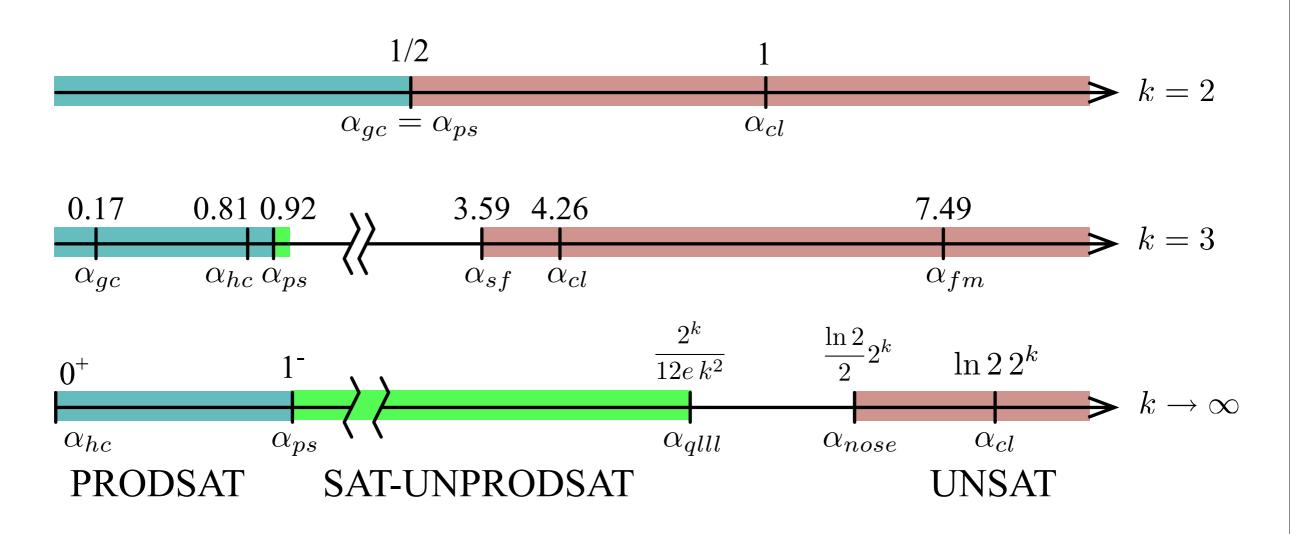


Continuous

$$\Pi^m \leftarrow \mathbb{CP}^{2^k - 1}$$

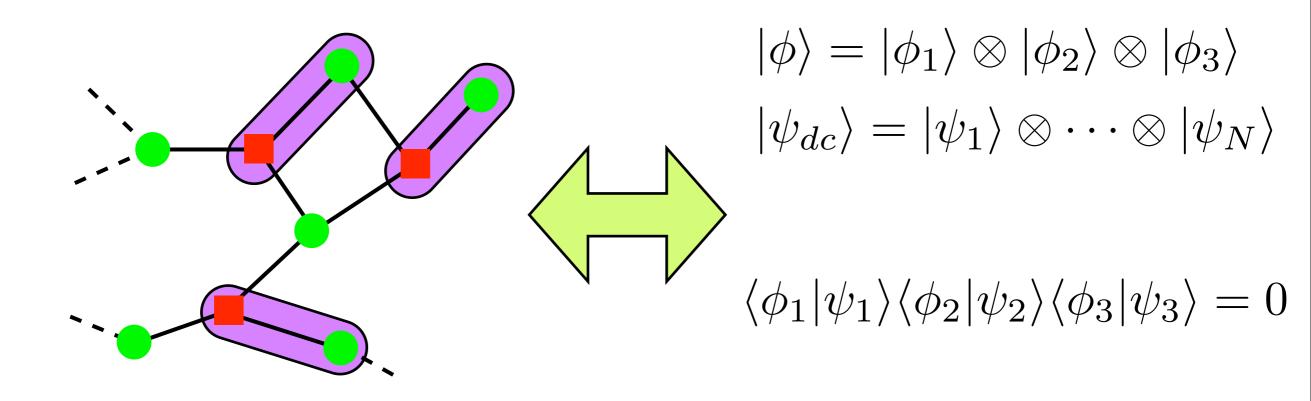
CRL, et al., 2009

Phase diagram of k-QSAT



- k=2 has direct PRODSAT-UNSAT transition
- Large k has entangled SAT phase, barrier to description of GS
- Numerics for k=3 (small sizes): $\alpha_c \approx 1 \pm 0.06$

PRODSAT: dimer covering states



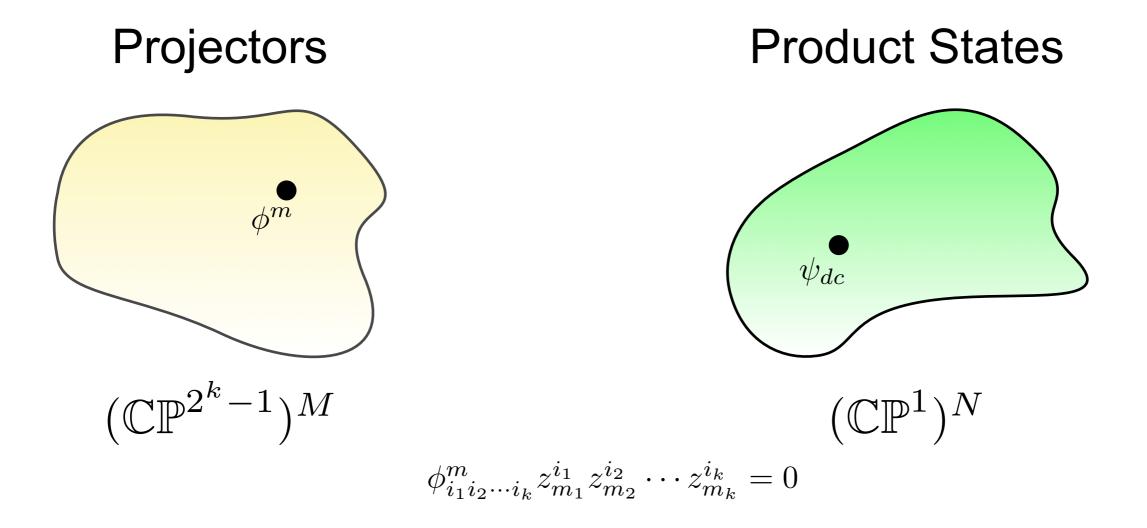
Dimer covering

Satisfying product state

- Simple argument for product projectors
- Product state perturbation theory for generic projectors

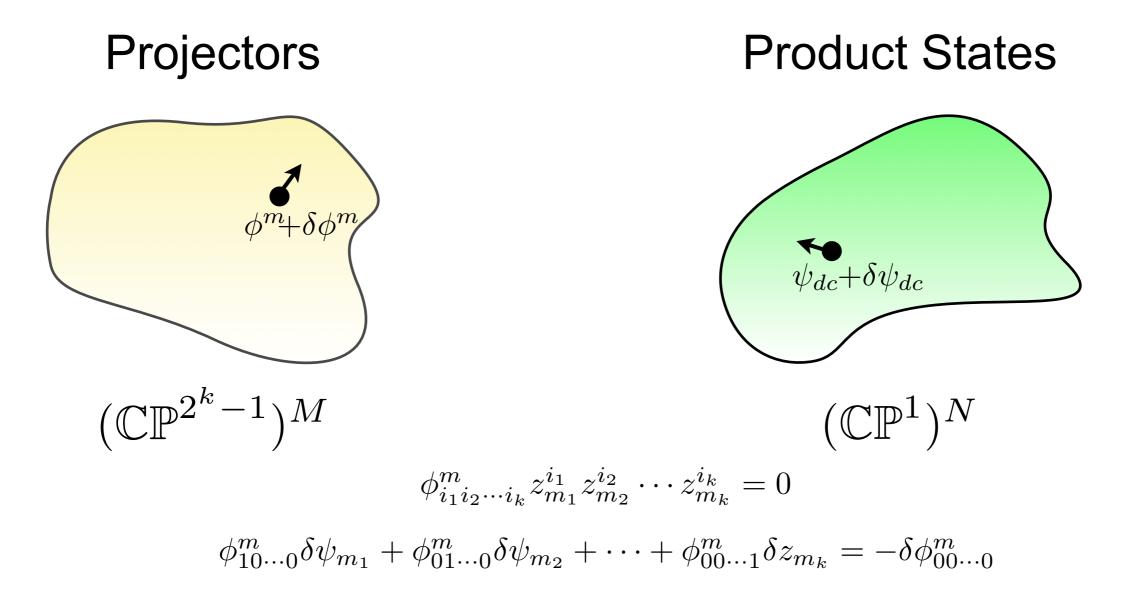
CRL, et al., 0910.2058, 2009

Product state perturbation theory



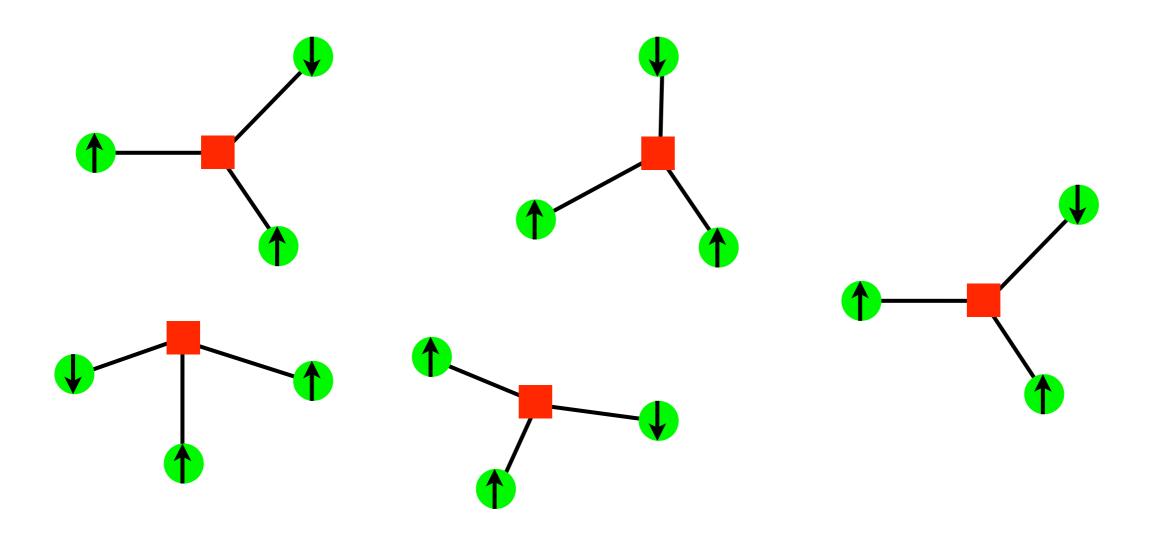
- Linearization of product state satisfiability condition
- Generically solvable if and only if G has dimer covers

Product state perturbation theory



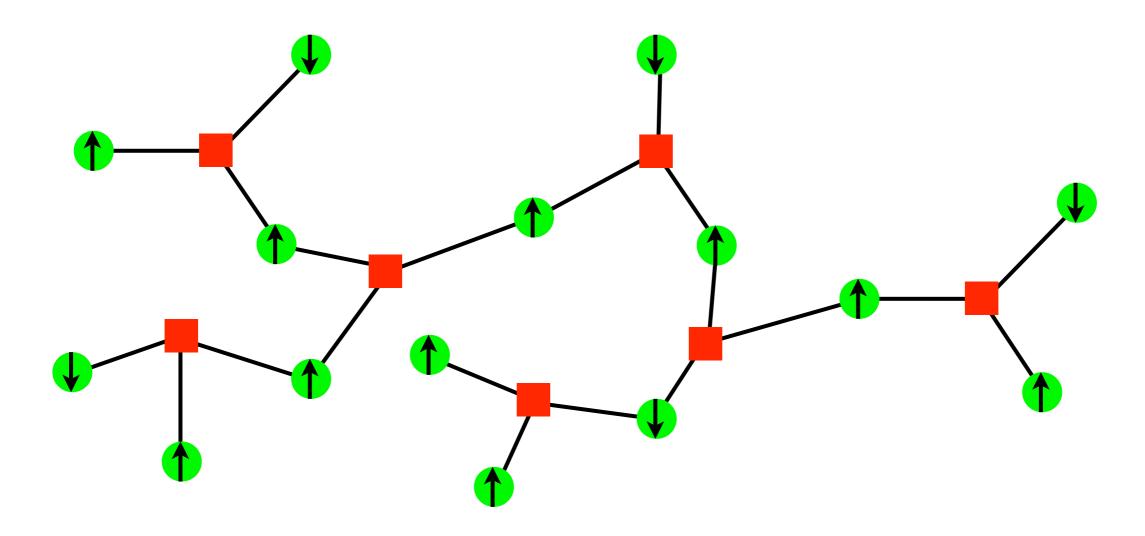
- Linearization of product state satisfiability condition
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Quantum Lovász Local Lemma



• Satisfying subspace exists for nonintersecting projectors

Quantum Lovász Local Lemma



- Satisfying subspace exists for nonintersecting projectors
- Lovász: if $D \le 2^k/e 1$ then the subspace still exists

Ambainis, Kempe, Sattath, 2009

Random QSAT: Open questions

- Where's the SAT-UNSAT phase transition?
- What is the complexity of computing R_g(G)? Does generic QSAT have a classical test?
- Is the entanglement transition a thermodynamic transition in the entropy?
- What about higher rank projectors? Chains at higher rank.
- Glass physics: Quantum analogs of clustering/dynamical phase transitions from classical glass problems? Many-body localization?
- Stat mech: Cavity methods for QSAT problem? Finite energy/ temperature behavior? Does it anneal?
- Quantum algorithms: Quantum generalizations of belief propagation, survey propagation? How does the adiabatic algorithm fare? Quantum metropolis?

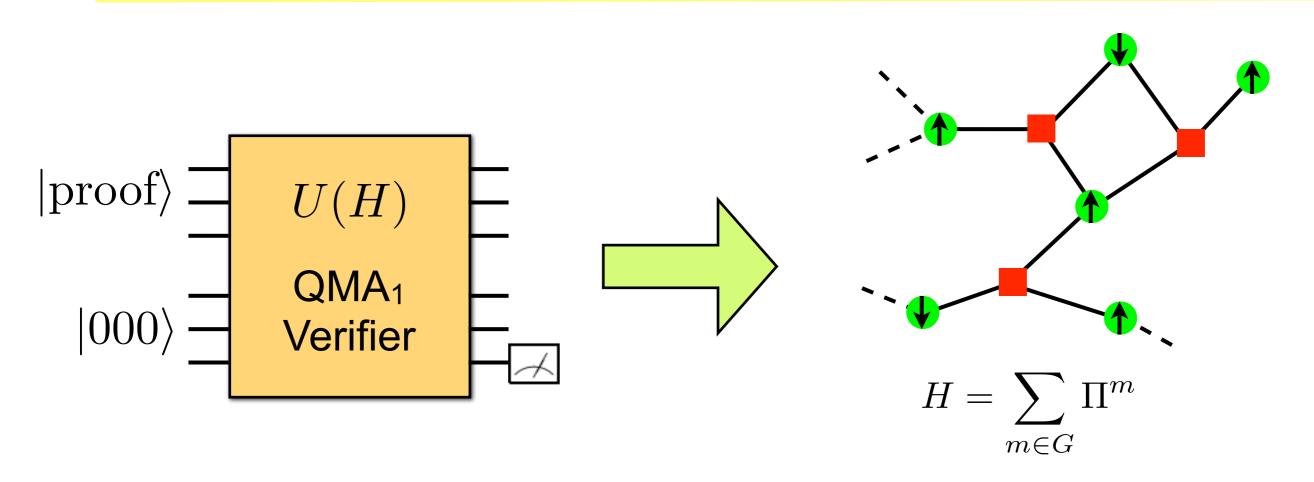
CRL, et al., QIC 10, 2010 CRL, et al., PRA 81, 2010

Conclusion

- Complexity theory attempts to identify problems which are hard.
- Hard problems in this sense are unsolvable by any physical process.
- Such hardness results provide indirect existence proofs for glassy dynamics.
- Random ensembles of -complete problems provide direct insight into those features which may lead to hardness and slow down.
- Phase transitions, clustering and replica symmetry breaking are characteristics of the classical statistical treatments of SAT.
- Quantum SAT is a quantum hard problem whose random ensemble exhibits a novel kind of 'clustered' phase: the entangled SAT phase.

Introductory Les Houches Lecture Notes to appear on arXiv soon with R. Moessner, A. Scardicchio and S. L. Sondhi.

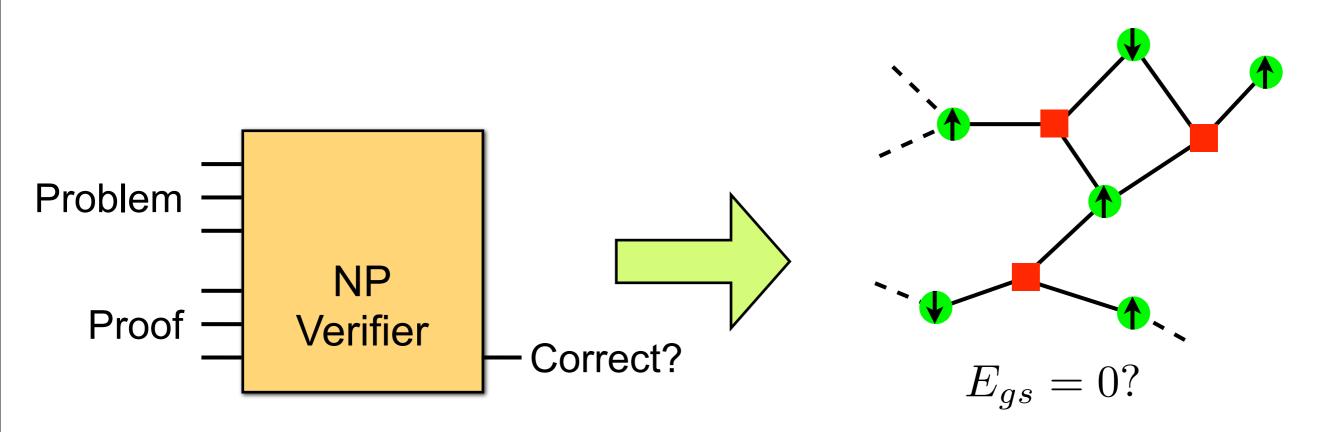
k-QSAT: Worst-case complexity



- 2-QSAT in P -- classical algorithm to solve -- Easy!
- k-QSAT (k > 3) is QMA₁-complete -- Hard!

Bravyi, 2006

3-SAT: Worst-case complexity



- 3-SAT can encode the operation of the verification circuit.
- 3-SAT is NP-complete: solve 3-SAT efficiently and you could solve all of NP efficiently (P=NP)

Cook 1971; Levin 1973