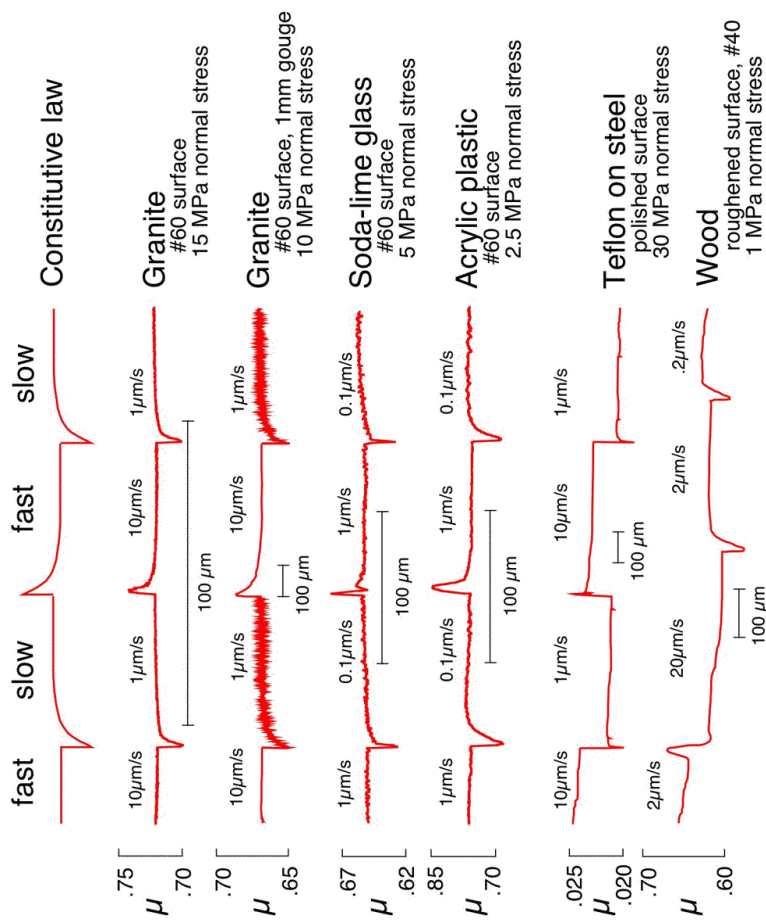
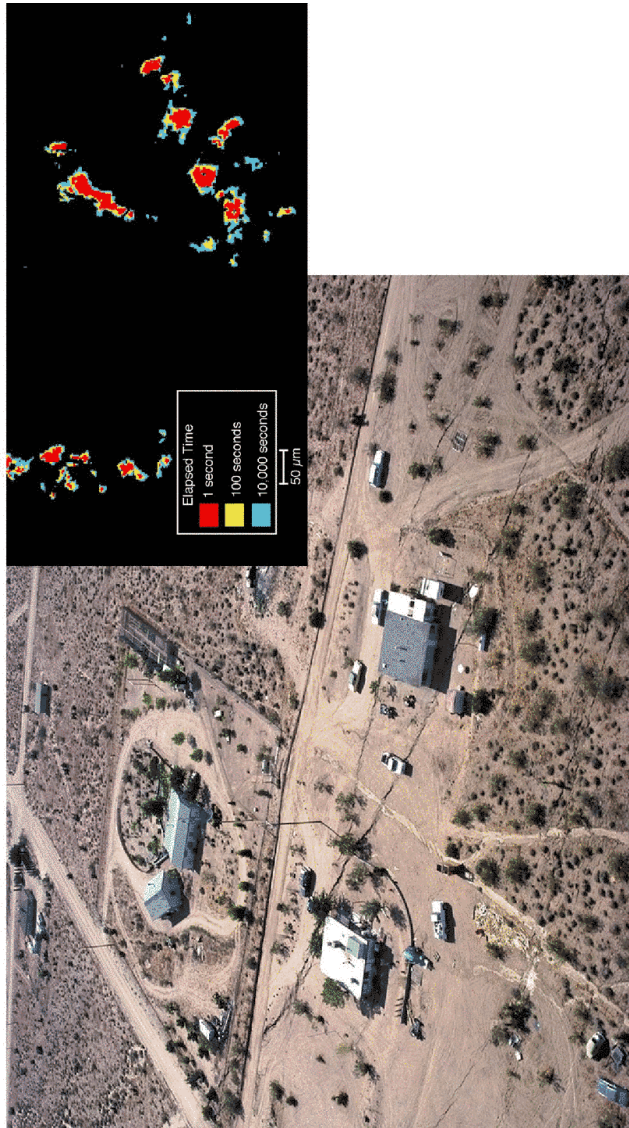


Macroscopic friction and earthquakes

Jim Dieterich, UC Riverside



Rate- and state-dependent friction

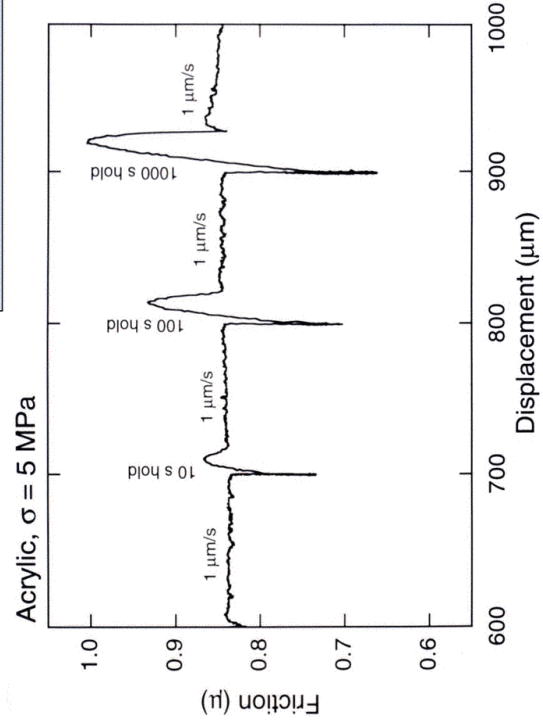
$$\frac{\tau}{\sigma} = \mu = \mu_0 + A \ln\left(\frac{V}{V^*}\right) + B \ln\left(\frac{\theta}{\theta^*}\right)$$

$$d\theta = dt - \frac{\theta}{D_c} d\delta - \frac{\alpha\theta}{B\sigma} d\sigma$$

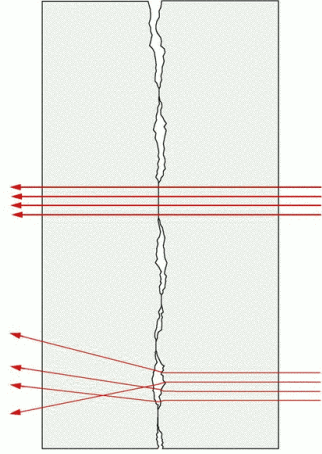
$$\frac{\tau}{\sigma} = \mu = \mu_0 + A \ln\left(\frac{V}{V^*}\right) + B \ln\left(\frac{\theta}{\theta^*}\right)$$

$$d\theta = dt - \frac{\theta}{D_c} d\delta - \frac{\alpha\theta}{B\sigma} d\sigma$$

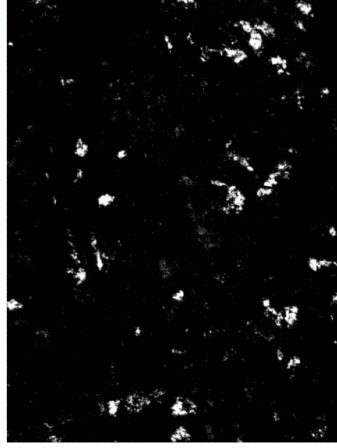
Slide – hold – slide experiments



Imaging contacts during slip



Schematic magnified view of contacting surfaces showing isolated high-stress contacts. Viewed in transmitted light, contacts appear as bright spots against a dark background.



Acrylic surfaces at 4MPa applied normal stress

Apparatus for viewing contacts during slip

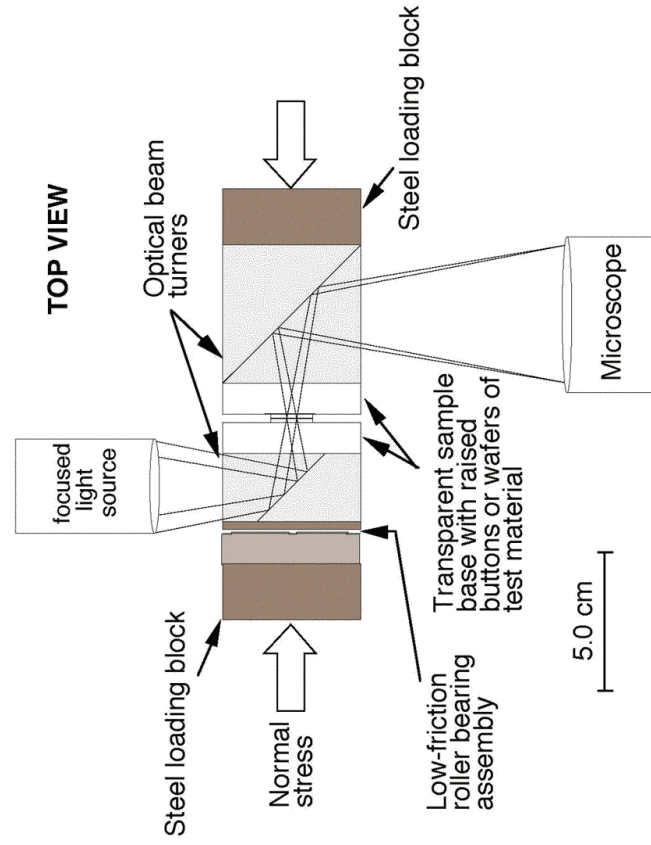
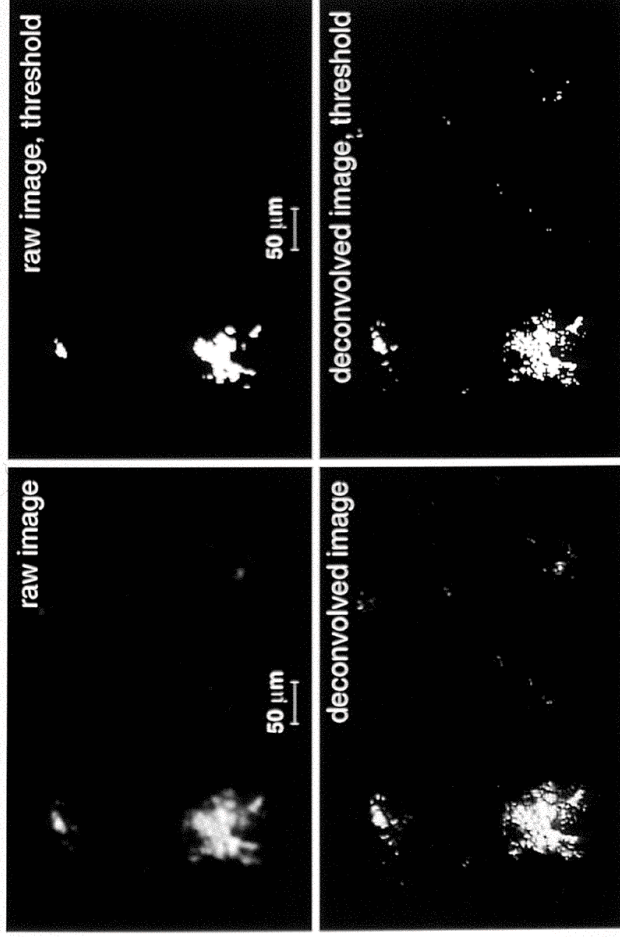
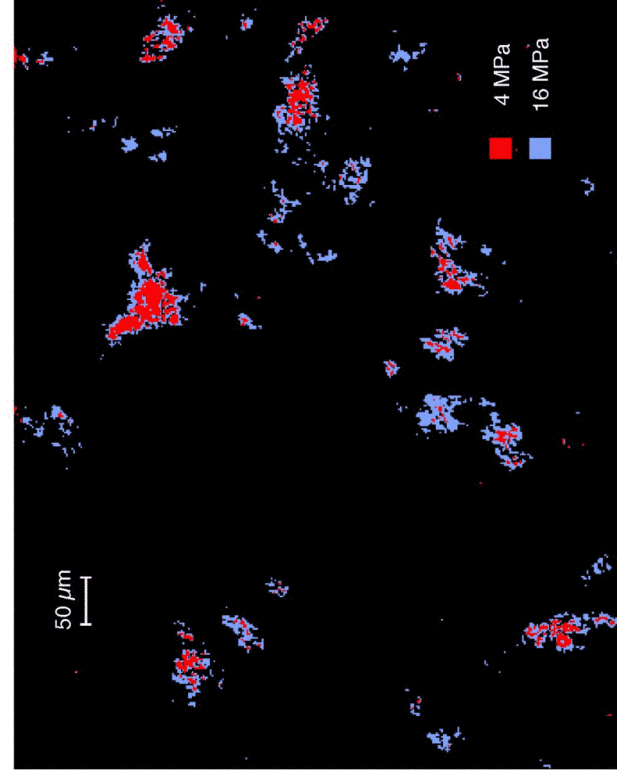


Image Processing

Images of contacts between #60 Quartz surfaces at 30 MPa



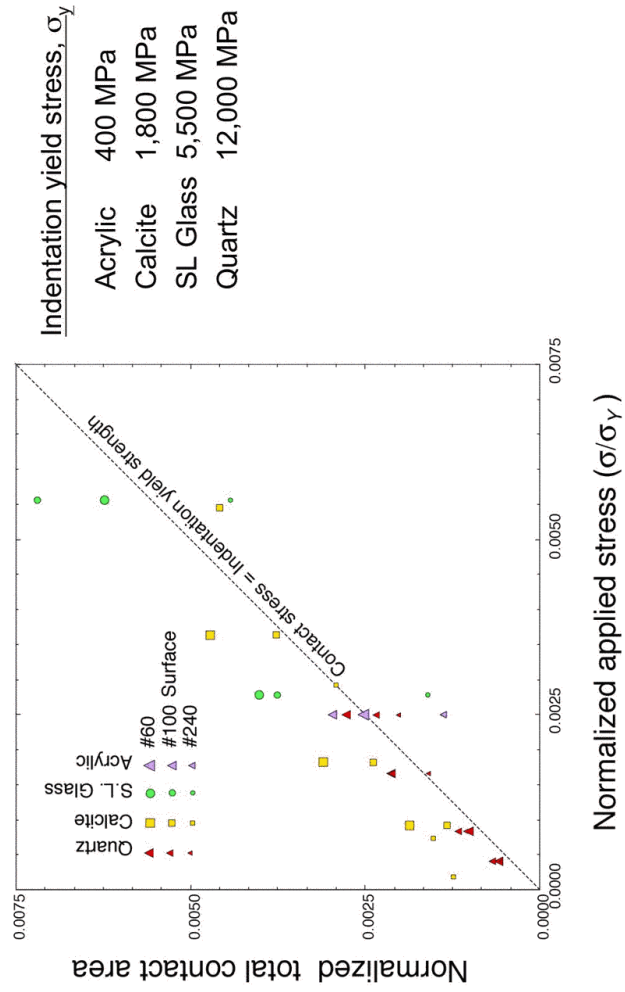
Change of contact area with normal stress



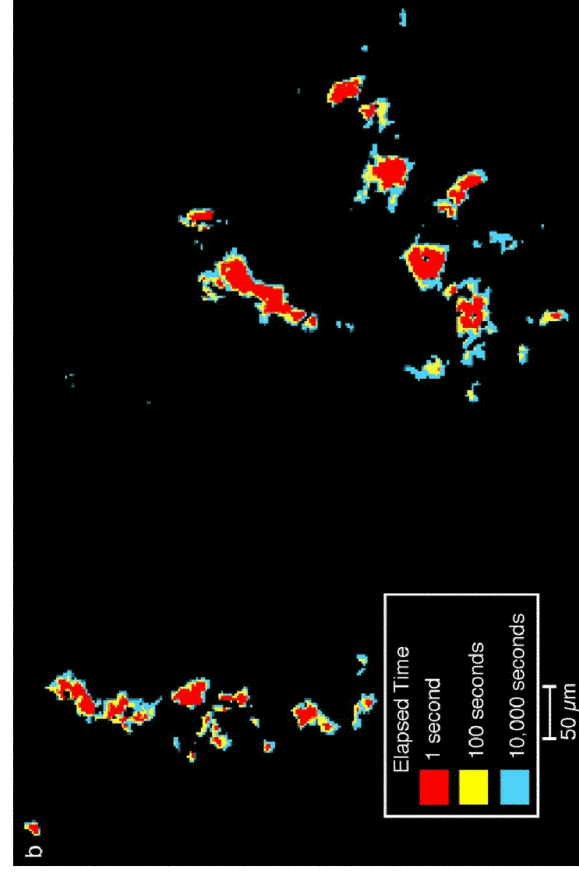
Acrylic plastic

Dieterich & Kilgore

Contact stresses

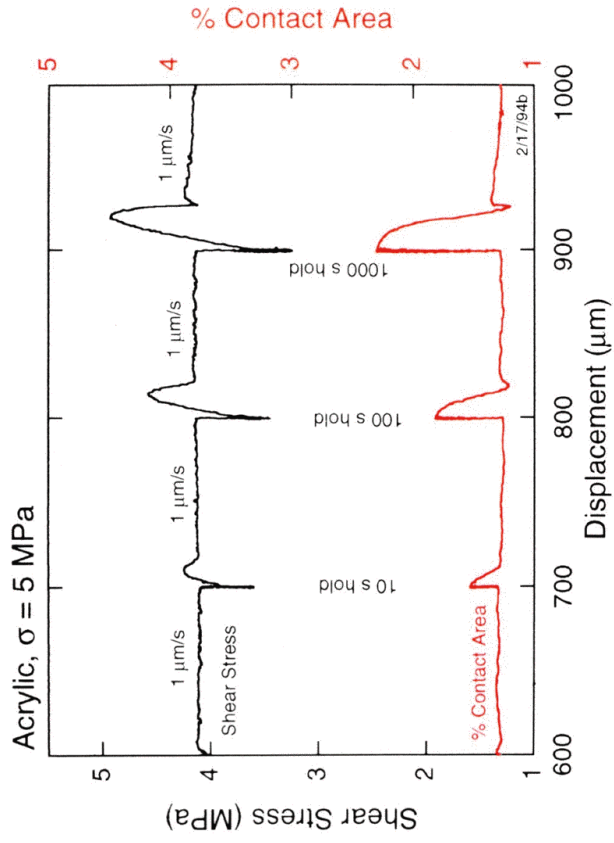


Increase of contact area with time



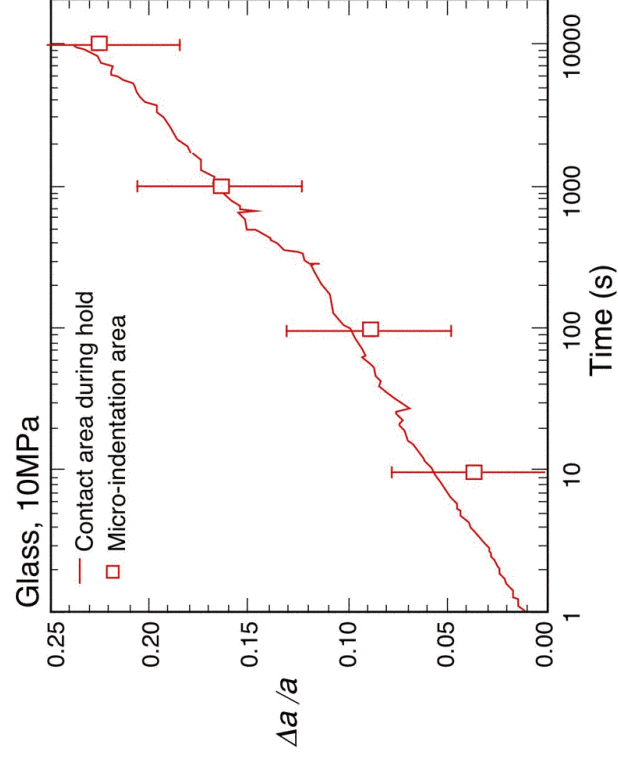
Acrylic plastic

Time dependent friction & contact area



Dieterich and Kilgore, PAGEOPH, 1994

Contact area during hold



Dieterich and Kilgore, PAGEOPH, 1994

Interpretation of friction terms

Bowden and Tabor adhesion theory of friction

Contact area: $\text{area} = c\sigma$

$c=1/$ indentation yield stress
 $g=$ shear strength of contacts

Shear resistance: $\tau = (\text{area}) (g), \tau/\sigma = \mu = cg$

Time and rate dependence of contact strength terms

Indentation creep: $c(\theta) = c_1 + c_2 \ln(\theta)$

Shear of contacts: $g(V) = g_1 + g_2 \ln(V)$

$$\mu = c_1 g_1 + c_1 g_2 \ln(V) + c_2 g_1 \ln(\theta) + \boxed{c_2 g_2 \ln(V+\theta)}$$

$$\mu = \mu_0 + A \ln(V) + B \ln(\theta) \quad \text{(Drop the high-order term)}$$

Slip instabilities – earthquake nucleation

Required:

Rate weakening at ss. $\mu_{ss} = \text{const.} + (A - B) \ln(V)$

Critical stiffness:

$$K_c = \frac{\sigma_c \xi}{D_c}$$

For slip on a fault patch, effective stiffness is given by elastic crack solution

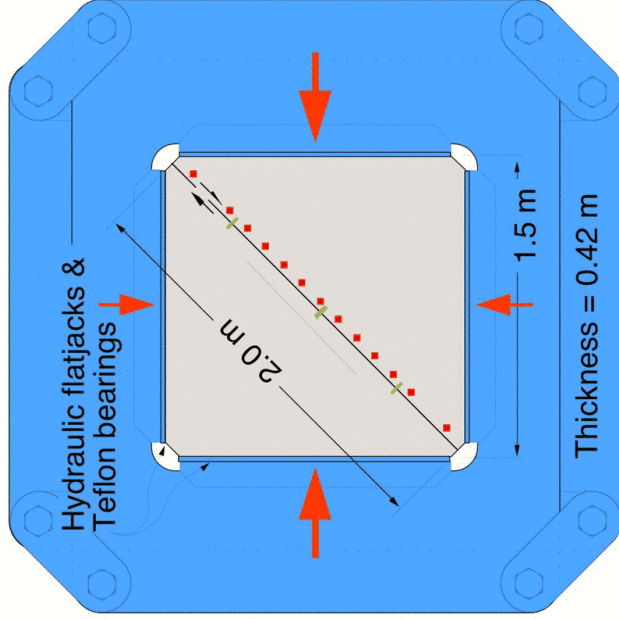
$$K = \frac{\Delta\tau}{d} = \frac{G\eta}{L}$$

Equating with critical stiffness:
 ($L > L_c$ for instability)

$$L_c = \frac{D_c G \eta}{\sigma_c \xi}$$

(η is crack geometry factor, $\eta \sim 1$)

Large-scale biaxial test of L_c



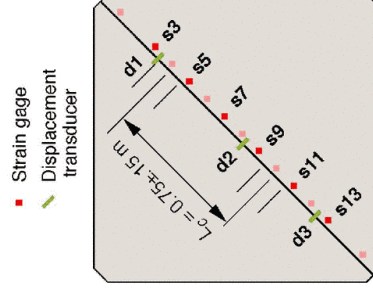
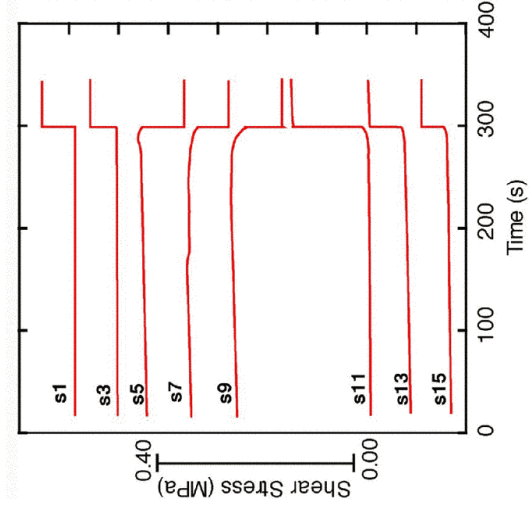
Minimum fault length for unstable slip

$$L_c = \frac{D_c G \eta}{\sigma \xi}$$

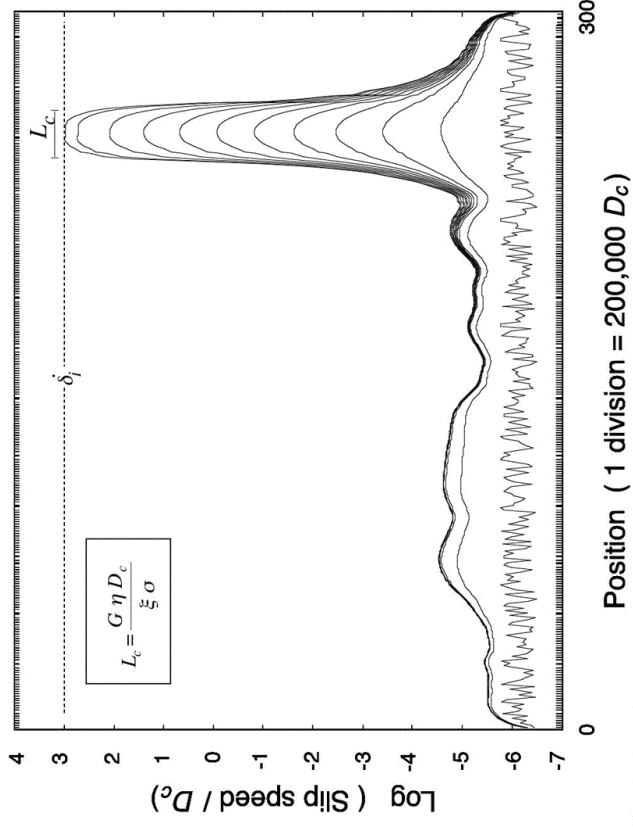
$G = 15000 \text{ MPa}$
 $\eta = 0.5$
 $D_c = 2 \mu\text{m}$
 $\sigma = 5 \text{ MPa}$
 $\xi = .4 B = 0.004$
 $L_c = .75 \text{ m}$

- Strain gage
- ◆ Displacement transducer

Confined Unstable Slip

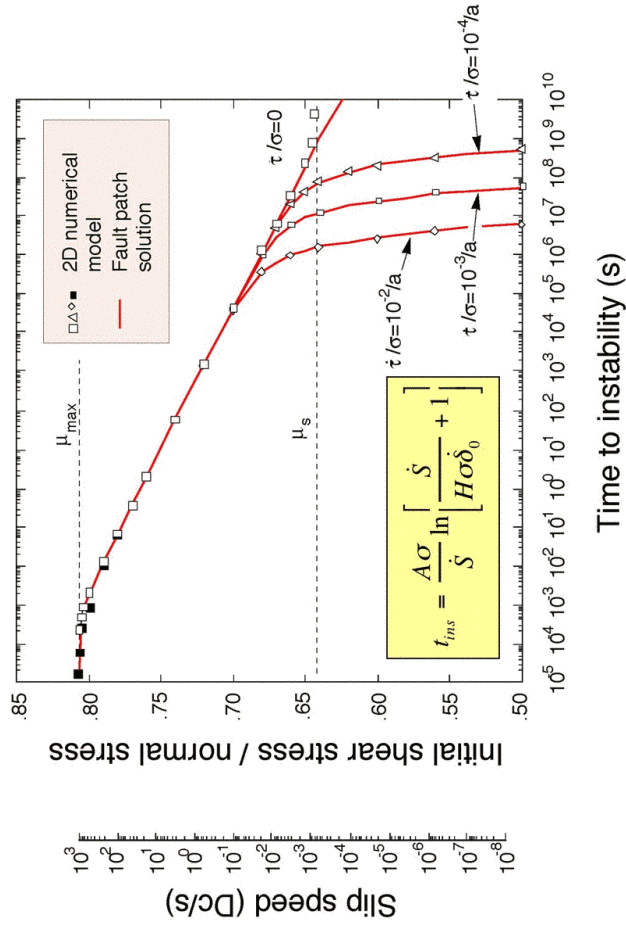


Earthquake nucleation – heterogeneous normal stress



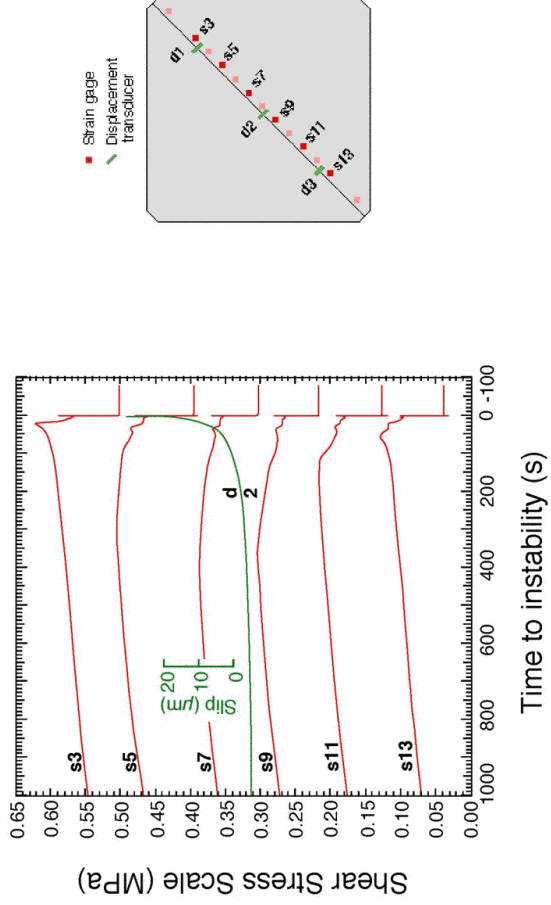
Dieterich, 1992, *Tectonophysics*

Time to instability

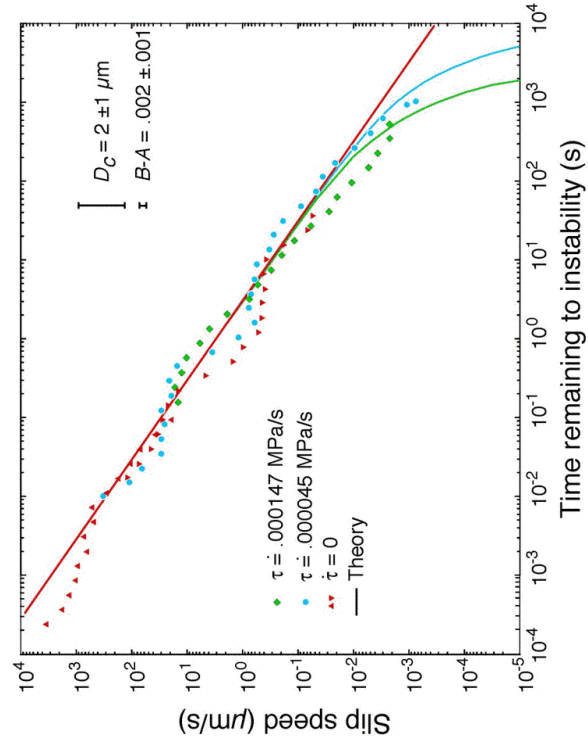


Dieterich, 1992, *Tectonophysics*

Accelerating slip prior to instability



Time to instability - Experiment and theory

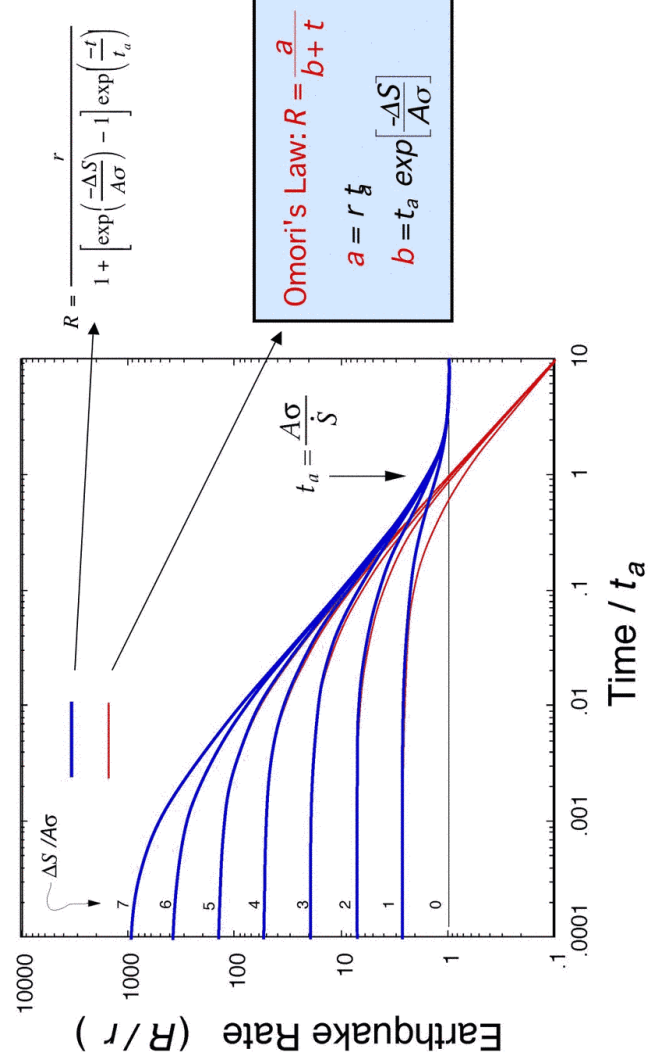


Formulation for earthquake rates

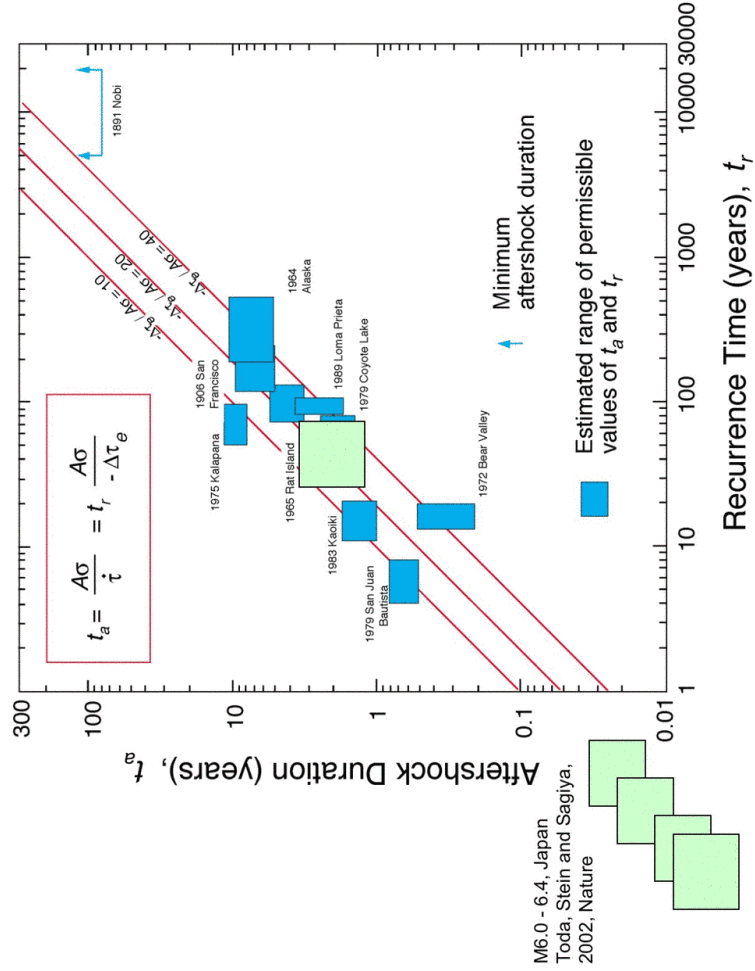
$$R = \frac{r}{\gamma \dot{\tau}_r}, \quad d\gamma = \frac{1}{A\sigma} \left[dt - \gamma d\tau + \gamma \left(\frac{\tau}{\sigma} - \alpha \right) d\sigma \right]$$

Dieterich, JGR (1994), Dieterich, Cayol, Okubo, Nature, (2000)

Earthquake rates following a stress step

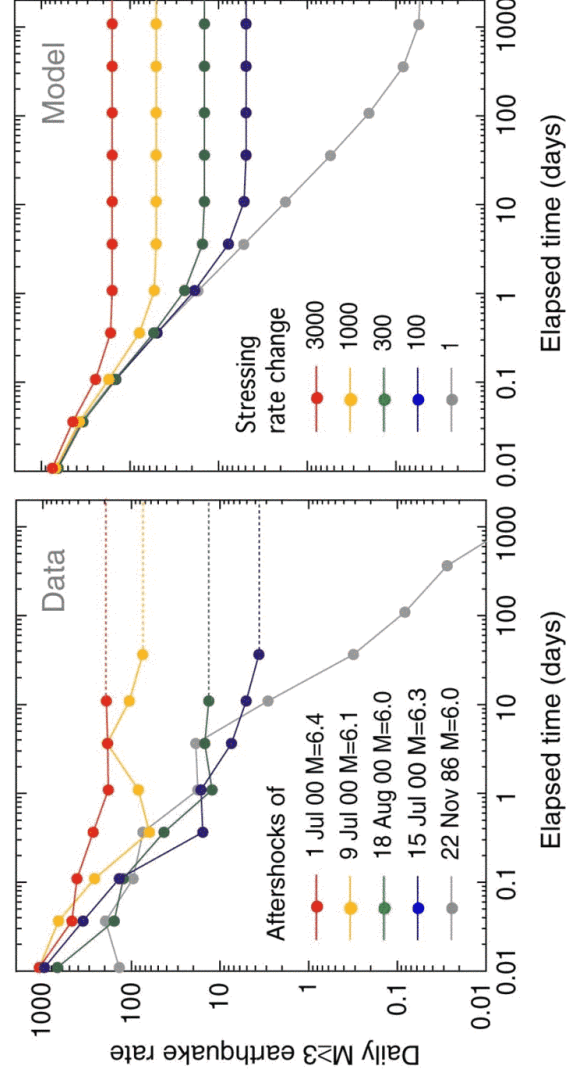


Aftershock duration



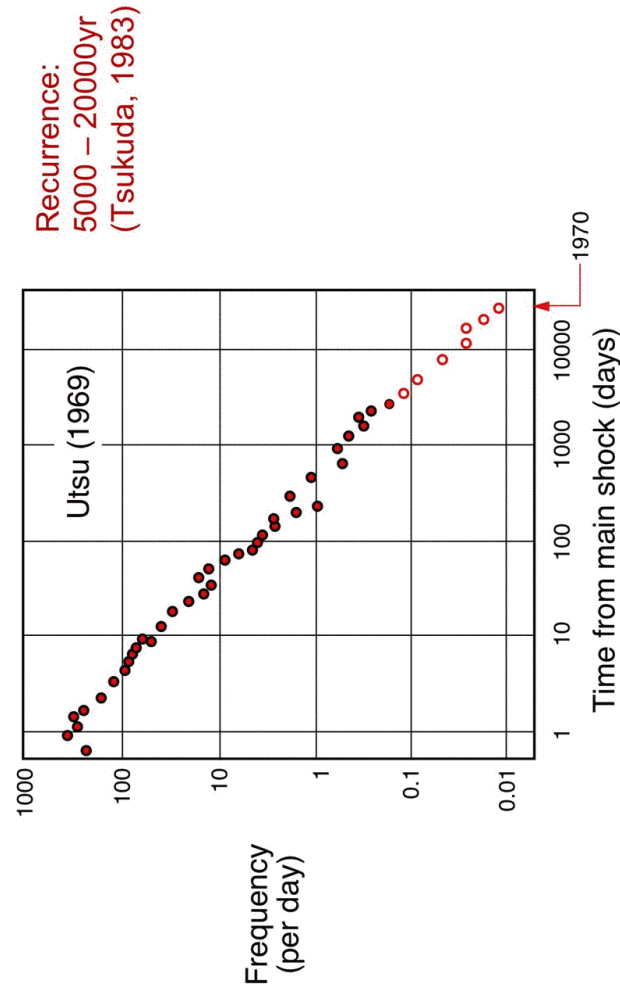
Short duration aftershock sequences

Izu Islands swarm in 2000:
7,000 $M \geq 3$ and 5 $M \geq 6$ earthquakes

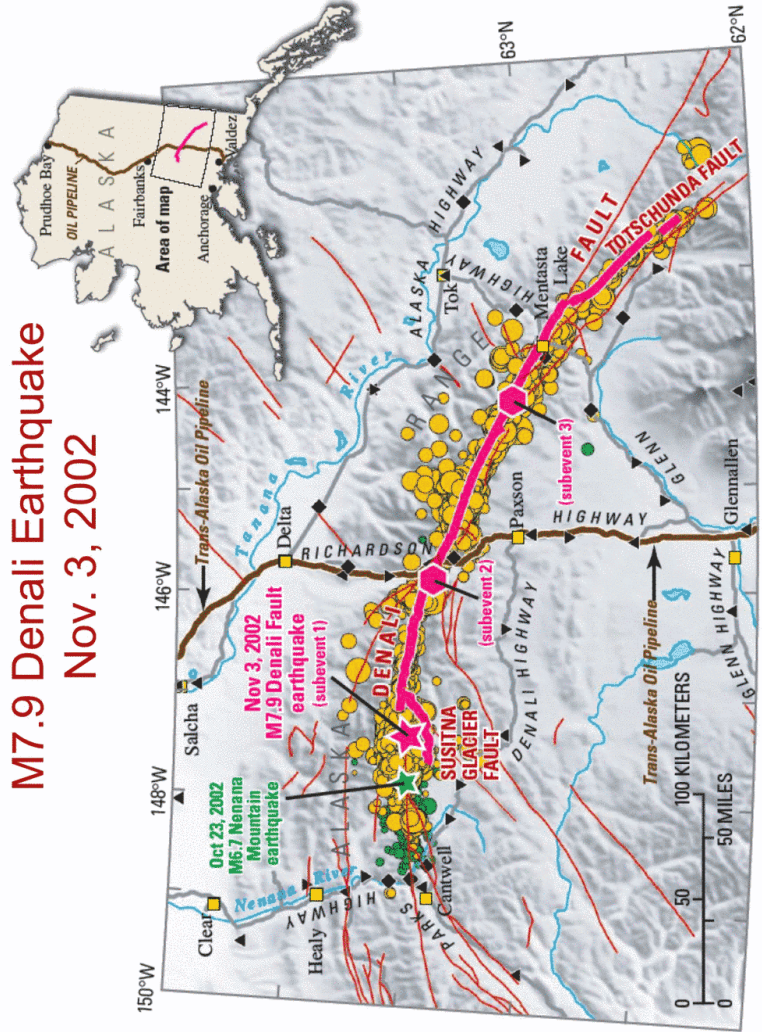


Toda, Stein and Sagiya, 2002, Nature

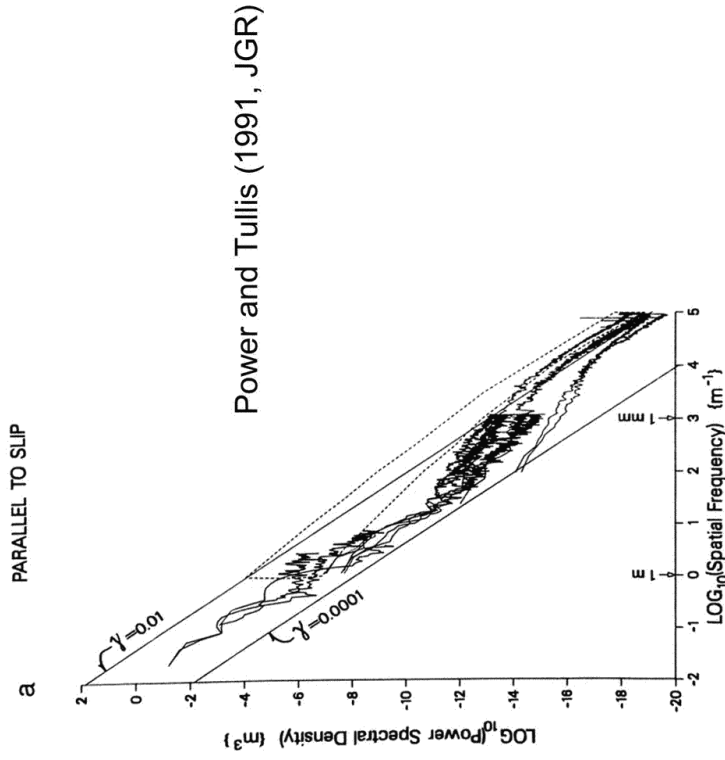
Nobi Earthquake, Oct 28, 1891



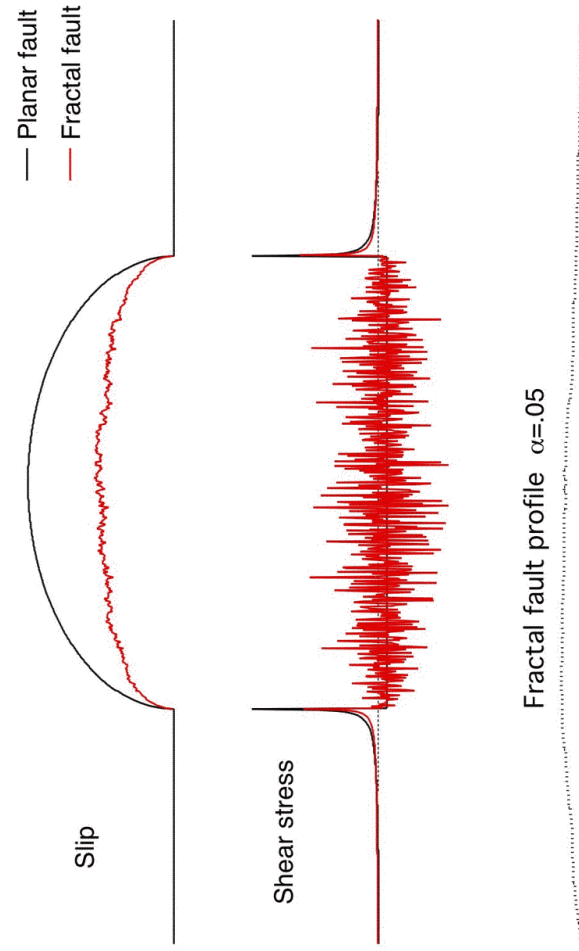
M7.9 Denali Earthquake Nov. 3, 2002



USGS, 2003

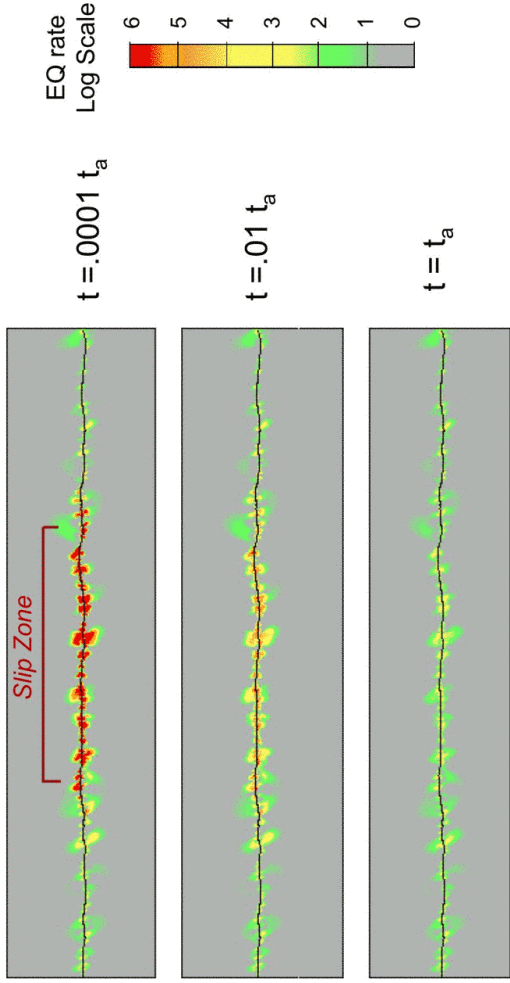


Slip of a fault patch

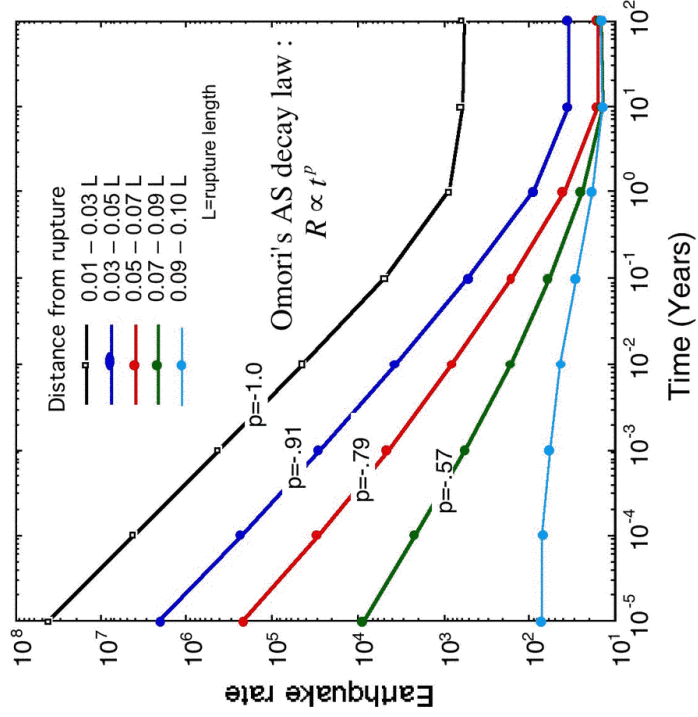


Stress relaxation: Seismicity following slip

$$R = \frac{r}{\gamma S_r} \quad d\gamma = \frac{1}{A\sigma} [dt - \gamma dS]$$

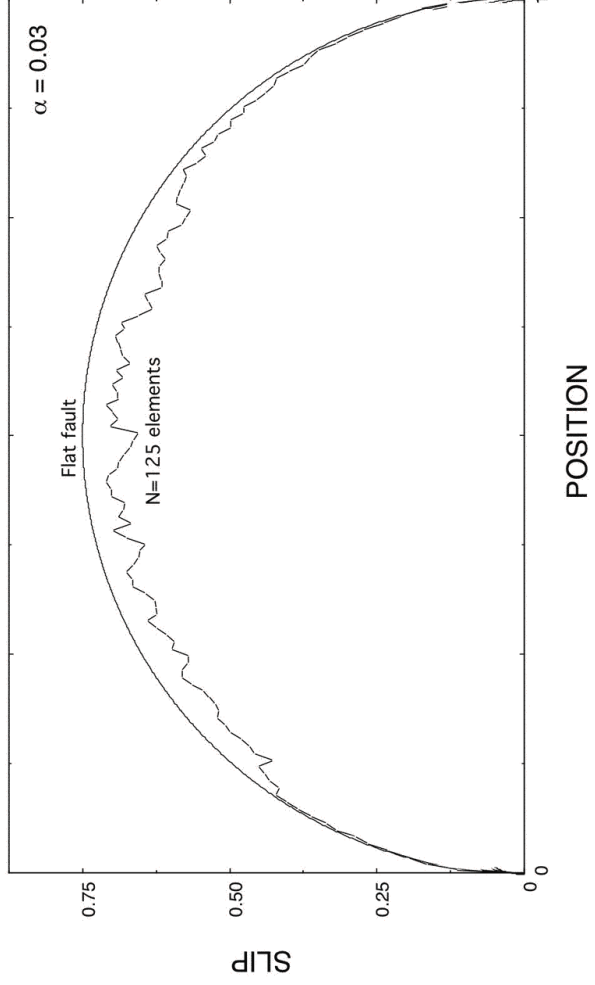


Aftershock rates as function of distance

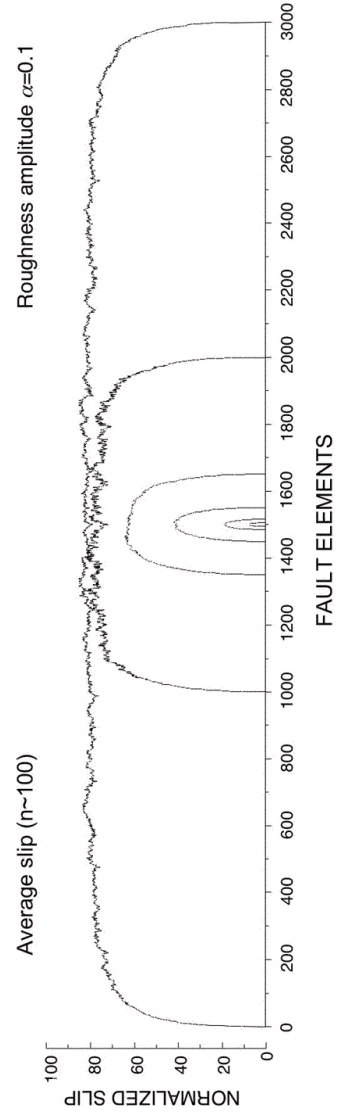


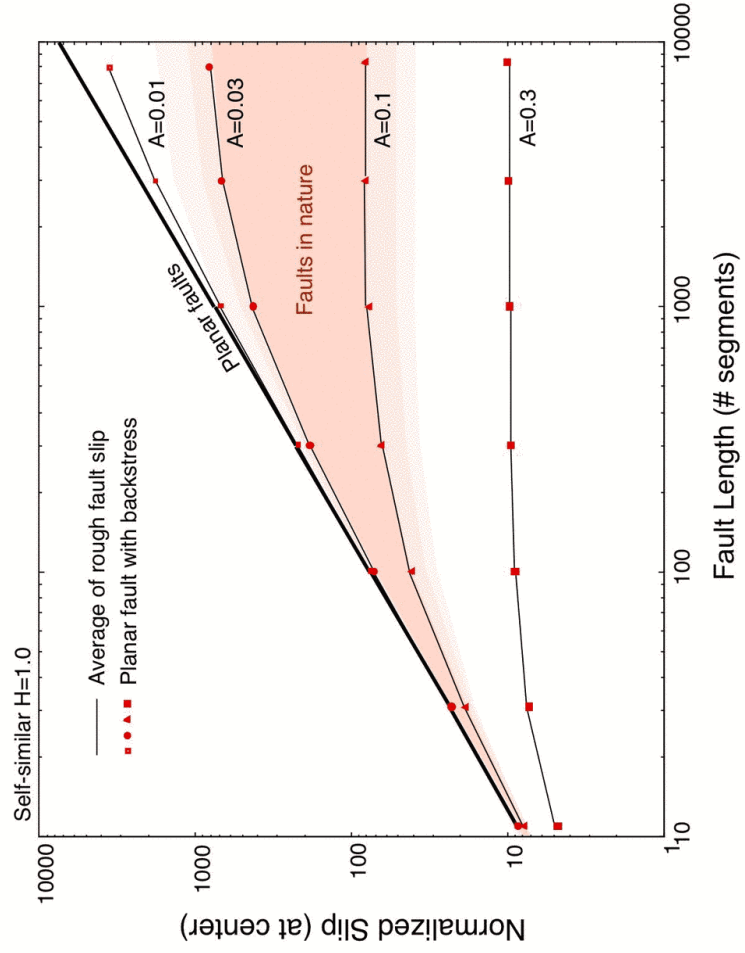
Effect of model resolution

Fault profile



Non-linear scaling of slip with fault length





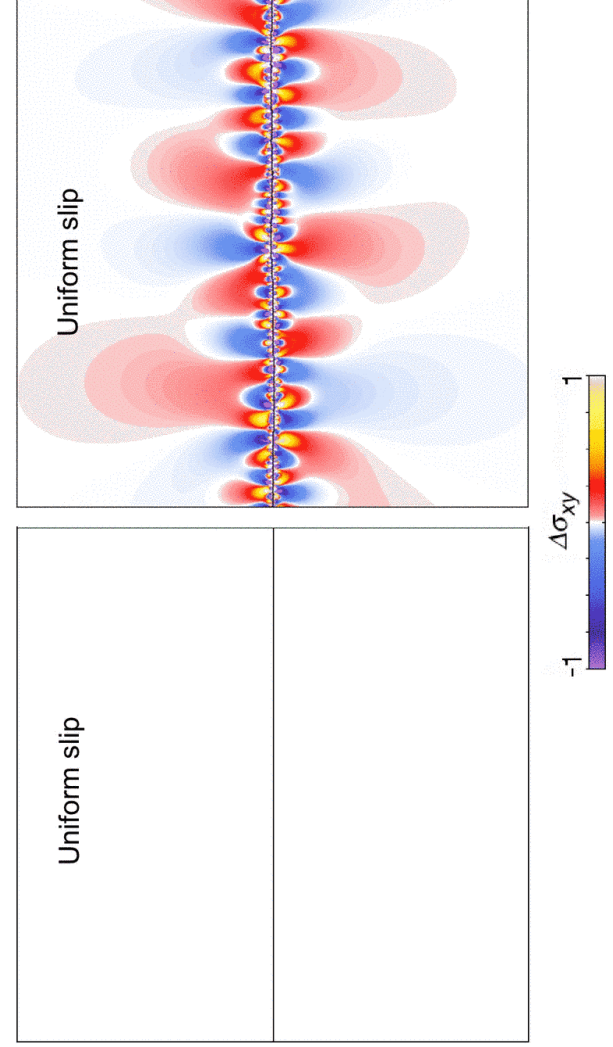
Fault slip and stress changes

Smooth fault

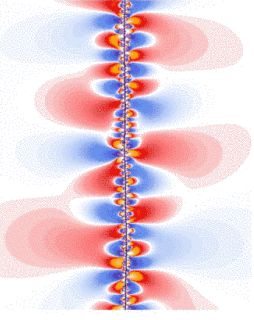
Uniform slip

Fault with self-similar roughness

Uniform slip



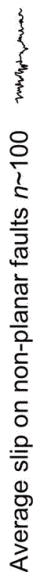
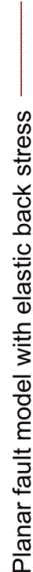
Origin of non-linear scaling and model scale-dependence

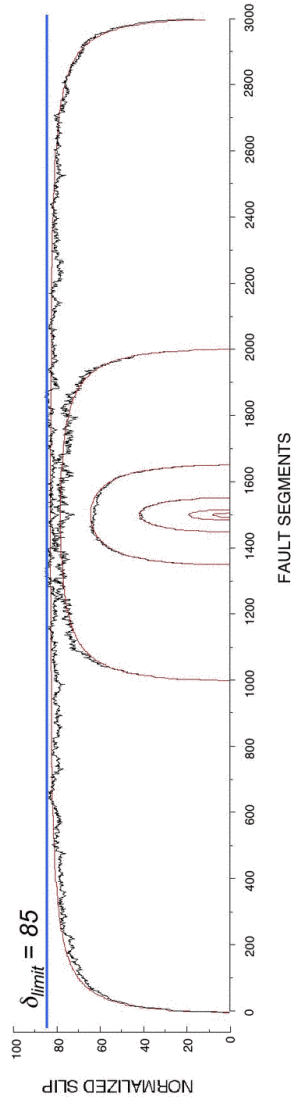


Geometric complexity inhibits to slip

Elastic strain energy increases with slip and requires greater work to slide.

Non-linear scaling of slip with fault length

Average slip on non-planar faults $n \sim 100$ 
 Planar fault model with elastic back stress 

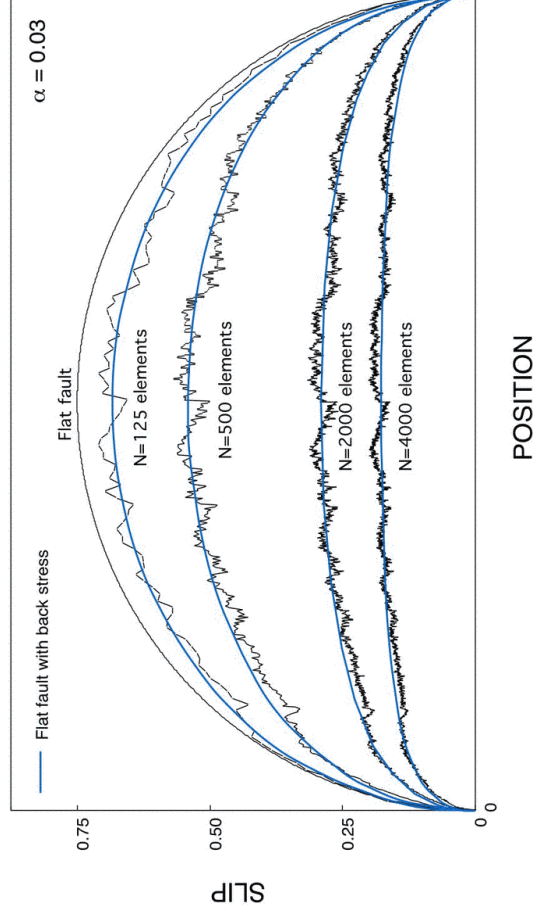


Roughness amplitude = .1

Back-stress also depends on number of fault elements N
for scaling with fixed L

$$S_{BACK} = \frac{NG}{\delta_{MAX}} d$$

Fault profile



Is the non-linear scaling applicable to faults?

Arises from increase of the elastic strain energy

In purely elastic models strain energy increases without limit

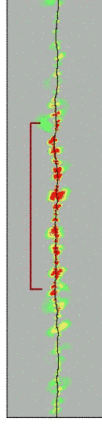
Real materials: limit to stresses and elastic strain energy

- Bulk yielding
- Slip on secondary faults or off-fault seismicity

Truncation of roughness at small wavelengths (by wear) may reduce but not eliminate the effect

Stress relaxation and off-fault seismicity Mechanisms and modeling

- 1) Bulk Aseismic bulk yielding: Viscoelasticity
 - Aseismic stress release at complexities between fault slip events (Nielsen and Knoppoff, 1998; Duan and Oglesby, 2005)
- 2) Seismicity rate simulations using the state-dependent rate eqn's
 - Simulates off-fault seismicity between fault slip events
 - Could be adopted to co-seismic represent co-seismic stress release
 - Appropriate time and stress dependence
- 3) Secondary faulting
 - Permits coseismic and interseismic stress relaxation
 - Complicated simulation to generate secondary faults
 - Finite model cannot fully represent off-fault seismicity



Stress relaxation: Secondary fault formation

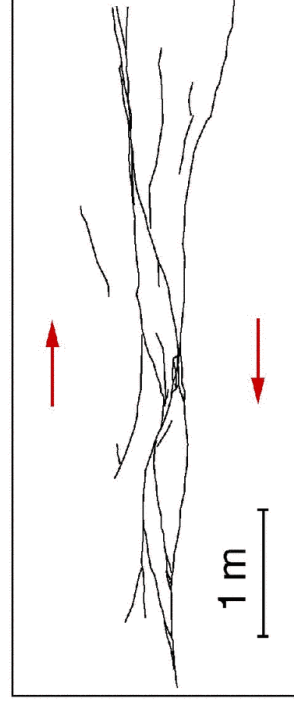
- Initial fault
- Secondary faults
- Simulation of secondary fault generation





Fault geometry

Individual faults exhibit approximately self-similar roughness (fractal dimension ~ 1).



Fault systems also appear to be scale-independent

