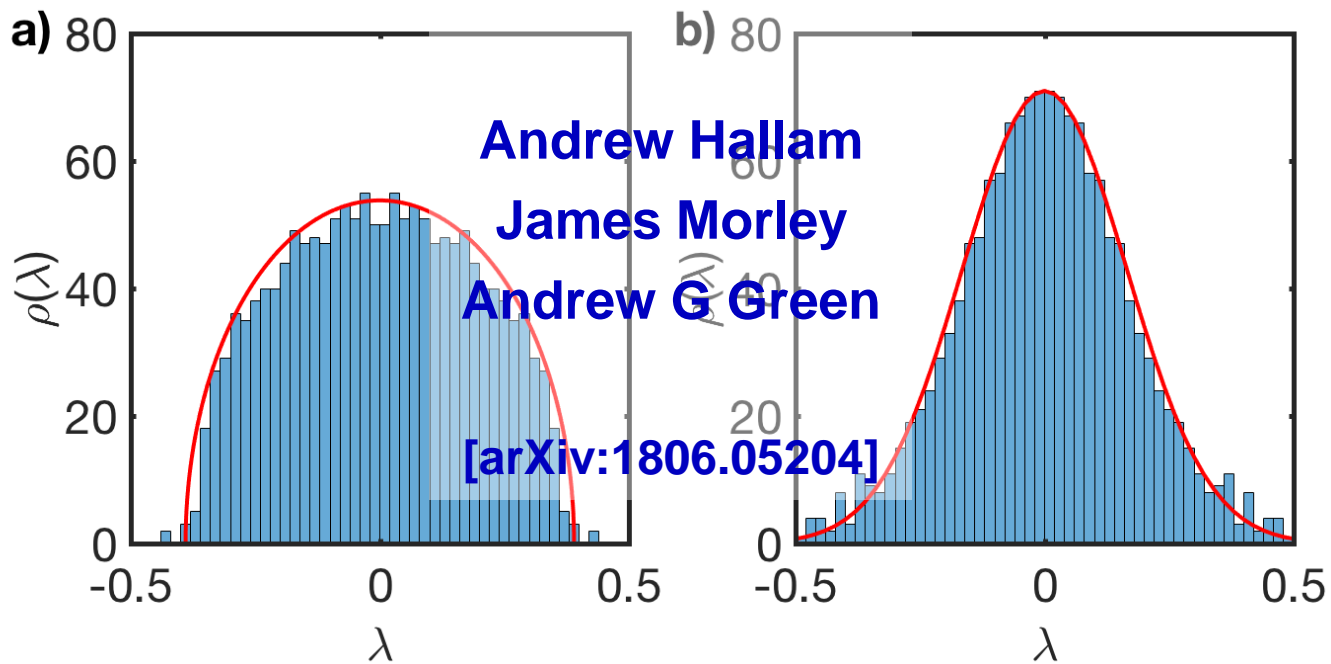


The Lyapunov Spectrum of Quantum Thermalisation



What connects quantum and classical many-body thermalization?

- **Mapping Quantum to Classical Dynamics**
 - TDVP of wavefunction MPS
 - TDVP of thermofield MPS
- **Extracting the Lyapunov Spectrum**
- **Results: Ising with tilted field**
 - Entanglement growth vs Kolmogorov-Sinai
 - A semi-circle law for Lyapunov spectrum
- **Discussion**



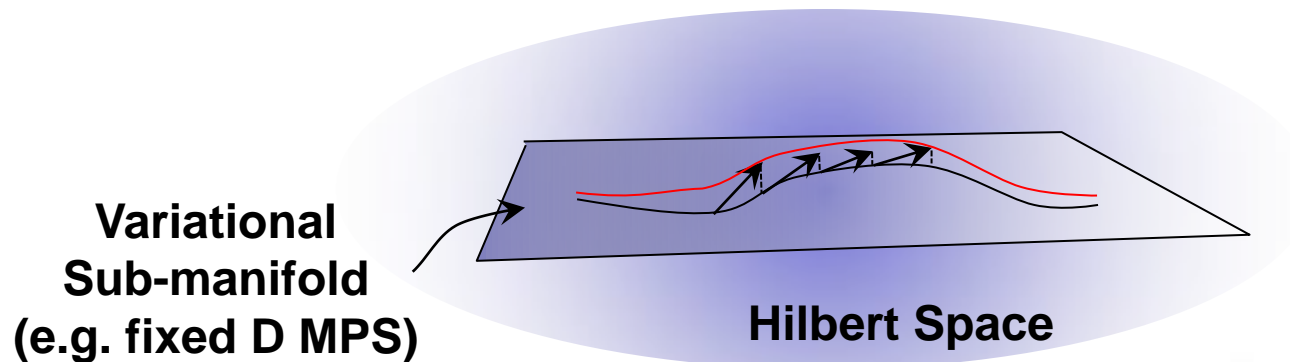
TDVP for Wavefunction

- **Variational parametrization** $|\psi(X)\rangle$

$$\langle \partial_{X_i} \psi | \partial_{X_j} \psi \rangle \dot{X}_j = i \langle \partial_{X_i} \psi | \mathcal{H} | \psi \rangle$$

- **Project dynamics / Optimize fidelity**

$$\text{Max}_X |\langle \psi(\mathbf{X} + d\mathbf{X}) | e^{i\hat{\mathcal{H}}dt} | \psi(\mathbf{X}) \rangle|^2$$



[Dirac, Proc. Cam. Phil. Soc. 1930]

[Haegeman et al PRL 2011]

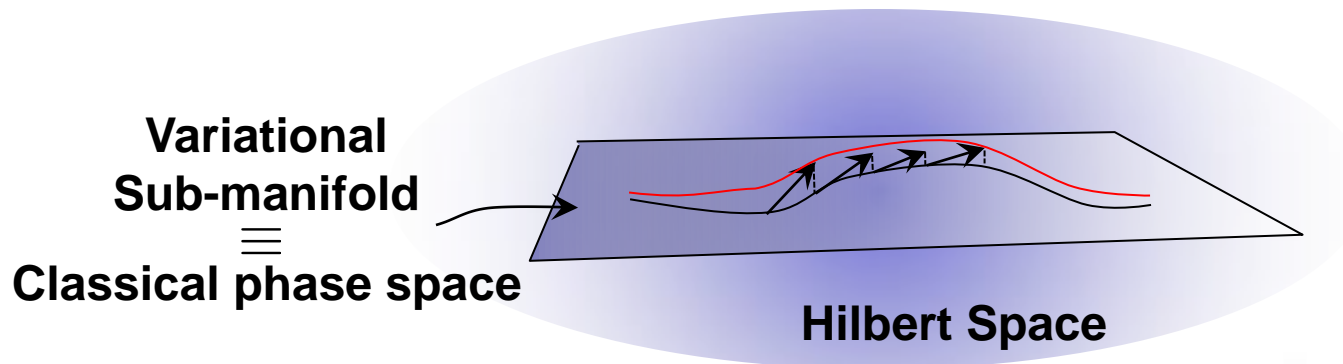
TDVP for Wavefunction

- Variational parametrisation $|\psi(X)\rangle$

$$\langle \partial_{X_i} \psi | \partial_{X_j} \psi \rangle \dot{X}_j = i \langle \partial_{X_i} \psi | \mathcal{H} | \psi \rangle$$

- g_{ij} →
- Classical Hamiltonian dynamics - Conserved quantities

$$q_i = \sqrt{2} \text{Im} \left[g_{ij}^{-1/2} X_j \right] \quad p_i = \sqrt{2} \text{Re} \left[g_{ij}^{-1/2} X_j \right]$$



Thermalisation => Dynamical Chaos

TDVP for Thermofield Double

$$\hat{\rho} = \sum_{\alpha} \gamma_{\alpha} |\alpha\rangle\langle\alpha| \quad \Leftrightarrow \quad |\psi\rangle = \sum_{\alpha} \gamma_{\alpha}^{1/2} |\alpha, \alpha\rangle$$

$$\mathcal{H} = \mathcal{H} \otimes \mathbf{1} + \mathbf{1} \otimes \mathcal{H}$$

Thermofield MPS:

- MPS approximation to $|\psi\rangle$ with $A_{ij}^{\sigma\delta}$ ↖ Doubled physical index
- TDVP optimizes fidelity $\text{Max}_{\mathbf{X}} \text{Tr} \left[\sqrt{\hat{\rho}(\mathbf{X} + d\mathbf{X})} \sqrt{e^{i\mathcal{H}dt} \hat{\rho}(\mathbf{X}) e^{-i\mathcal{H}dt}} \right]$
- Optimizes a set of observations
- Pure states evolve to mixed states

TDVP for Thermofield Double

$$\hat{\rho} = \sum_{\alpha} \gamma_{\alpha} |\alpha\rangle \langle \alpha| \quad \Leftrightarrow \quad |\psi\rangle = \sum_{\alpha} \gamma_{\alpha}^{1/2} |\alpha, \alpha\rangle$$
$$\mathcal{H} = \mathcal{H} \otimes \mathbf{1} + \mathbf{1} \otimes \mathcal{H}$$

Aside

Thermofield MPS:

- Expectations should be same on two sub-spaces
- Numerical drift fixed by symmetrising

$$\mathbb{A}_{i \otimes i', j \otimes j'}^{\sigma \delta} = \mathbb{A}_{i' \otimes i, j' \otimes j}^{\delta \sigma}$$

- Achieve by tangent space gauge fixing

[Haegeman et al PRL107, 070601 (2011)] [Hallam, Morley, Green arXiv:1806.05204]

Outline:

- Mapping Quantum to Classical Dynamics
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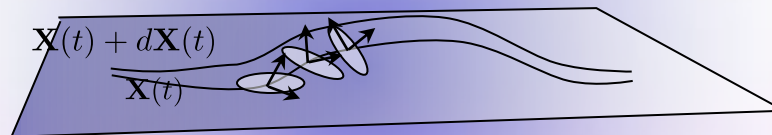
Calculation

- Trajectory – TDVP

$$\langle \partial_{X_i} \psi | \partial_{X_j} \psi \rangle \dot{X}_j = i \langle \partial_{X_i} \psi | \mathcal{H} | \psi \rangle$$

- Distance measure $dS^2 = 1 - |\langle \psi(\mathbf{X} + d\mathbf{X}) | \psi(\mathbf{X}) \rangle|^2$
- Hamiltonian dynamics on classical phase space ...
... use tools from classical dynamical systems

[Geist et al, Prog Theor Phys 83, 875 (1990)]



Hilbert Space

Calculation

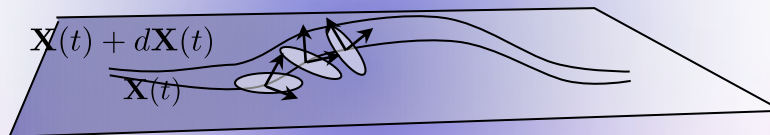
- Trajectory – TDVP

$$\langle \partial_{X_i} \psi | \partial_{X_j} \psi \rangle \dot{X}_j = i \langle \partial_{X_i} \psi | \mathcal{H} | \psi \rangle$$

- Divergence of $\mathbf{X}(t)$ and $\mathbf{X}(t) + d\mathbf{X}(t)$ -linearised TDVP

$$-i \langle \partial_{X_i} \psi | \partial_{X_j} \psi \rangle dX_j = \underbrace{\langle \partial_{X_i} \partial_{X_j} \psi | \mathcal{H} | \psi \rangle}_{\text{Generates non-zero Lyapunovs}} d\bar{X}_j + \langle \partial_{X_i} \psi | \mathcal{H} | \partial_{X_j} \psi \rangle dX_j$$

- Growth in some direct's
- Reduce in others
- Direct's change
- Average along trajectories



Calculation

- Trajectory – TDVP

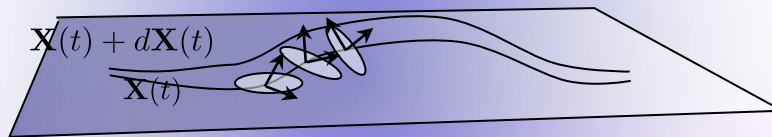
$$\langle \partial_{X_i} \psi | \partial_{X_j} \psi \rangle \dot{X}_j = i \langle \partial_{X_i} \psi | \mathcal{H} | \psi \rangle$$

- Instantaneous Lyapunov Exponents

$$d\dot{X}_i = M_{ij} dX_j \Rightarrow \{dX_i(t=0)\} = Q(t=0)$$

$$Q(t+dt)R(t+dt) = \exp[M(t)dt]Q(t) \Rightarrow \lambda_i = \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T \lambda_i(t)$$

- Lyapunov Exponents $\lambda_i(t) = \lim_{dt \rightarrow 0} \frac{1}{dt} \log R_{ii}(t)$



Hilbert Space

Aside

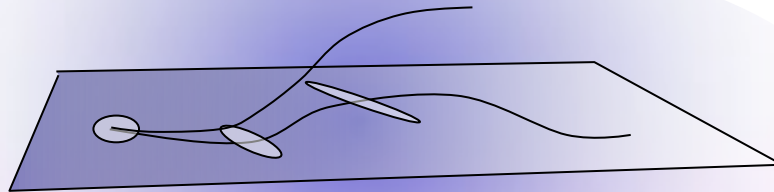
Relationship to Zero-point Fluctuations

- Initial wavepacket – evolve with truncated Wigner

- Path integral over MPS [arXiv:1607.01778] $e^{i \int \mathcal{H} dt} = \int DA e^{iS[A]}$

$$S[A] = S[A_0] + Tr \left[\underbrace{i\bar{Y}\dot{Y} + \bar{Y}Y \langle \partial_{\bar{Y}} \psi | \mathcal{H} | \partial_Y \psi \rangle + \bar{Y}\bar{Y} \langle \partial_{\bar{Y}} \partial_{\bar{Y}} \psi | \mathcal{H} | \psi \rangle + YY \langle \psi | \mathcal{H} | \partial_{\bar{Y}} \partial_{\bar{Y}} \psi \rangle}_{\text{Anomalous terms}}$$

Anomalous terms: Generate zero-point fluctuations about saddle point



Hilbert Space

Calculation

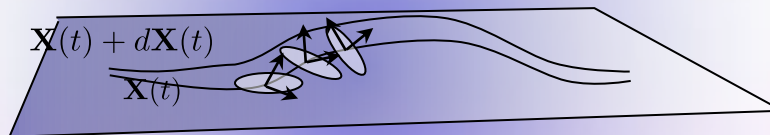
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Outline:

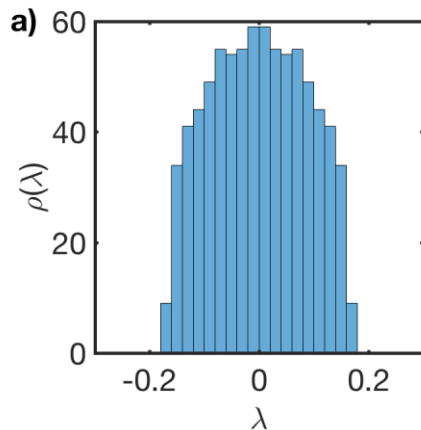
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Wavefunction MPS

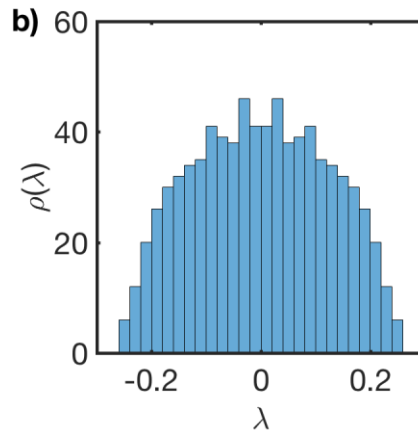
$$\mathcal{H} = \sum_i \left[J \sigma_i^z \sigma_{i+1}^z + h^x \sigma_i^x + h^z \sigma_i^z \right]$$

Spectra: Evolve from product state $|\psi(t=0)\rangle_i = 0.540|\uparrow\rangle + 0.841|\downarrow\rangle$



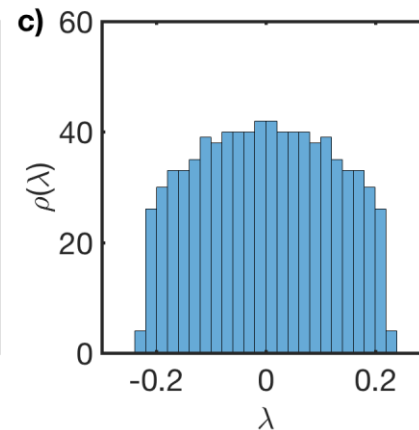
Non-integrable

$$J = 1, h^z = 1, h^x = 0.5$$



Integrable

$$J = 1, h^z = 1, h^x = 0$$



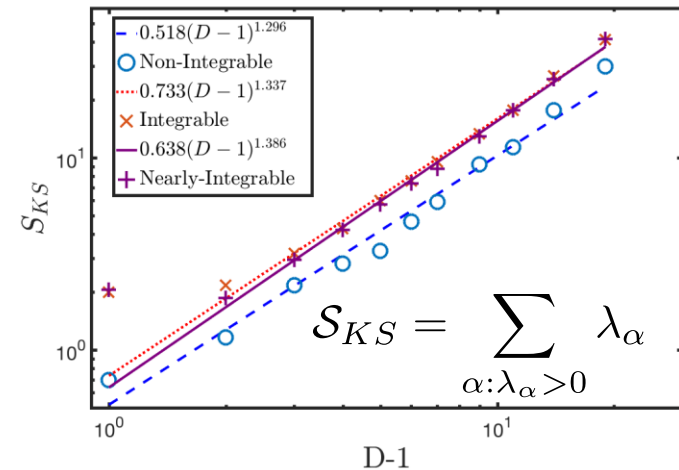
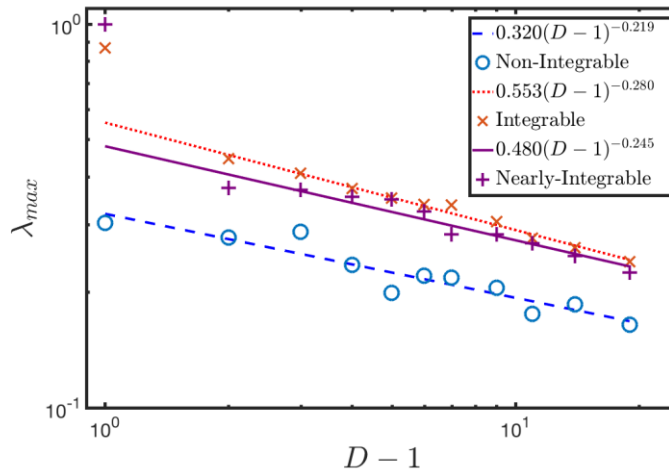
Nearly - Integrable

$$J = 1, h^z = 1, h^x = 0.1$$

Wavefunction MPS

$$\mathcal{H} = \sum_i \left[J \sigma_i^z \sigma_{i+1}^z + h^x \sigma_i^x + h^z \sigma_i^z \right]$$

Bond Order Dependence:

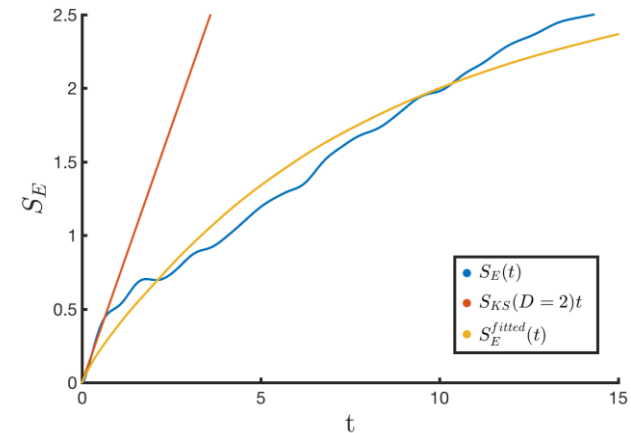
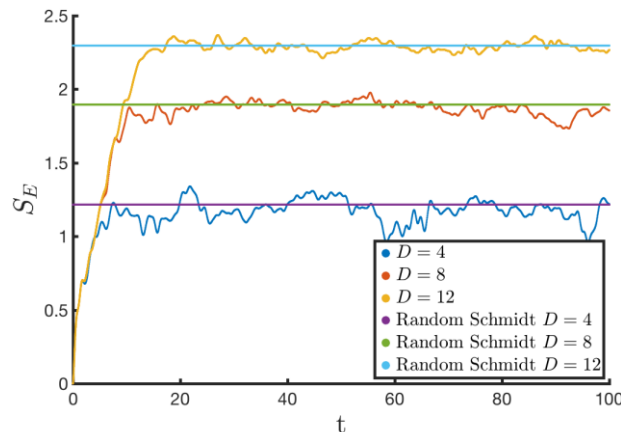


- Spectra do not converge with D
- How then do we relate to physical properties?

Wavefunction MPS

$$\mathcal{H} = \sum_i \left[J \sigma_i^z \sigma_{i+1}^z + h^x \sigma_i^x + h^z \sigma_i^z \right]$$

Entanglement growth:



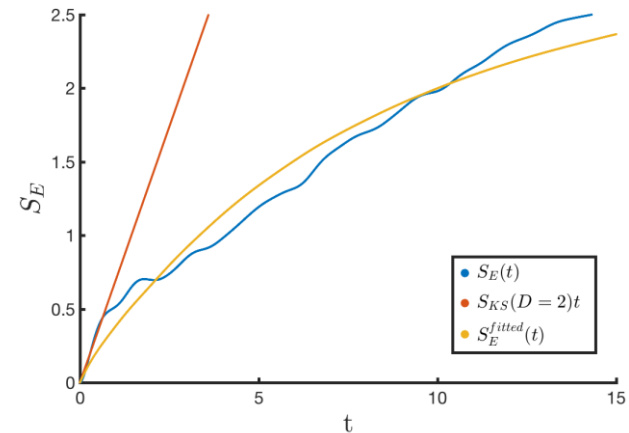
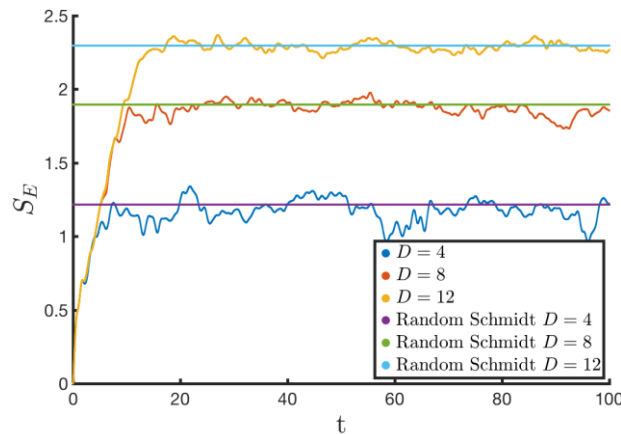
- **Initial entanglement growth:** $\dot{S}_E(t=0) = \mathcal{S}_{KS}(D=2)$
- **Bond order grows with time:** $S_E(t) = \mathcal{S}_E^*(D) \Rightarrow D(t)$

[Zurek and Paz, Physica D (1995)] [Miller and Sarkar PRE (1999)]

Wavefunction MPS

$$\mathcal{H} = \sum_i \left[J \sigma_i^z \sigma_{i+1}^z + h^x \sigma_i^x + h^z \sigma_i^z \right]$$

Entanglement growth:



Generalise:

$$\dot{S}_E(t) = \frac{\mathcal{S}_{KS}(D(t))}{(D(t) - 1)^2}$$

Aside

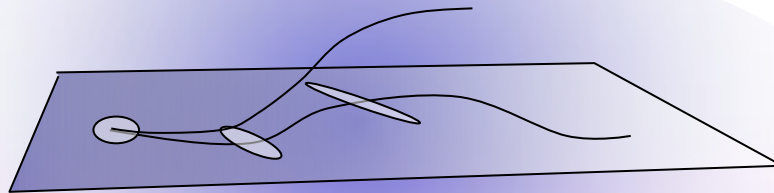
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$$\mathcal{S}[A] = \mathcal{S}[A_0] + Tr \left[\begin{array}{c} i\bar{Y}\dot{Y} + \bar{Y}Y \langle \partial_{\bar{Y}} \psi | \mathcal{H} | \partial_Y \psi \rangle \\ + \bar{Y}\bar{Y} \langle \partial_{\bar{Y}} \partial_{\bar{Y}} \psi | \mathcal{H} | \psi \rangle + YY \langle \psi | \mathcal{H} | \partial_{\bar{Y}} \partial_{\bar{Y}} \psi \rangle \end{array} \right]$$

- \mathcal{S}_{KS} related to fluctuation determinant averaged on trajectory?



Hilbert Space

Wavefunction MPS

$$\mathcal{H} = \sum_i \left[J \sigma_i^z \sigma_{i+1}^z + h^x \sigma_i^x + h^z \sigma_i^z \right]$$

Comments:

- **Crossover between complementary pictures**
- **Low bond order – chaotic thermalisation**
- **High bond order – dephasing thermalization**

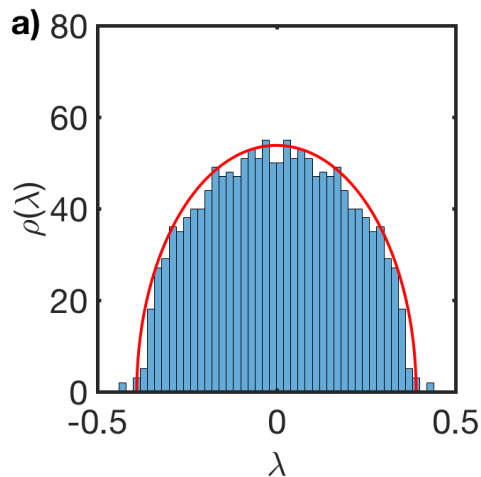
- **Lyapunov Spectra don't converge**
- **Effective bond order grows $D(t)$**

$$\dot{S}_E(t) = \frac{\mathcal{S}_{\mathcal{K}\mathcal{S}}(D(t))}{(D(t) - 1)^2}$$

Thermofield MPS

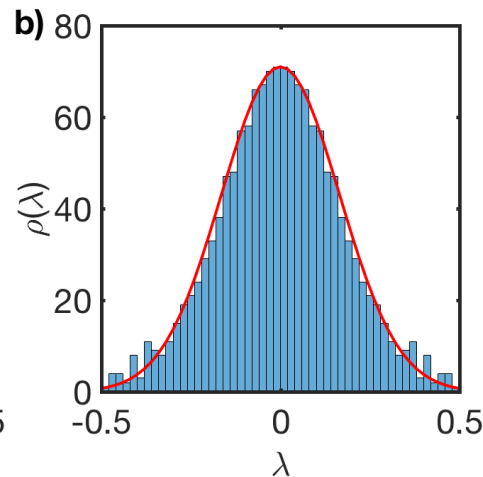
$$\mathcal{H} = \sum_i \left[J \sigma_i^z \sigma_{i+1}^z + h^x \sigma_i^x + h^z \sigma_i^z \right]$$

Lyapunov Spectra: Start from product state near centre of spectrum



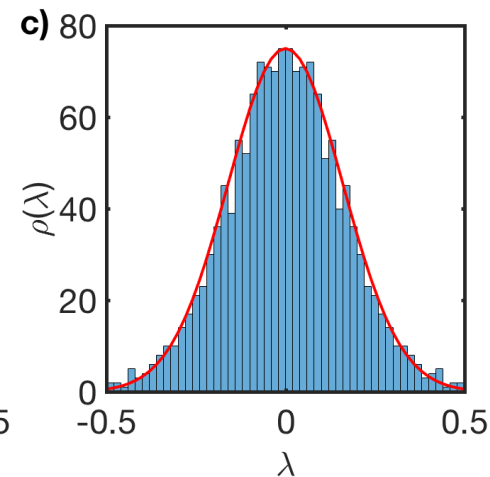
Non-integrable

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Integrable

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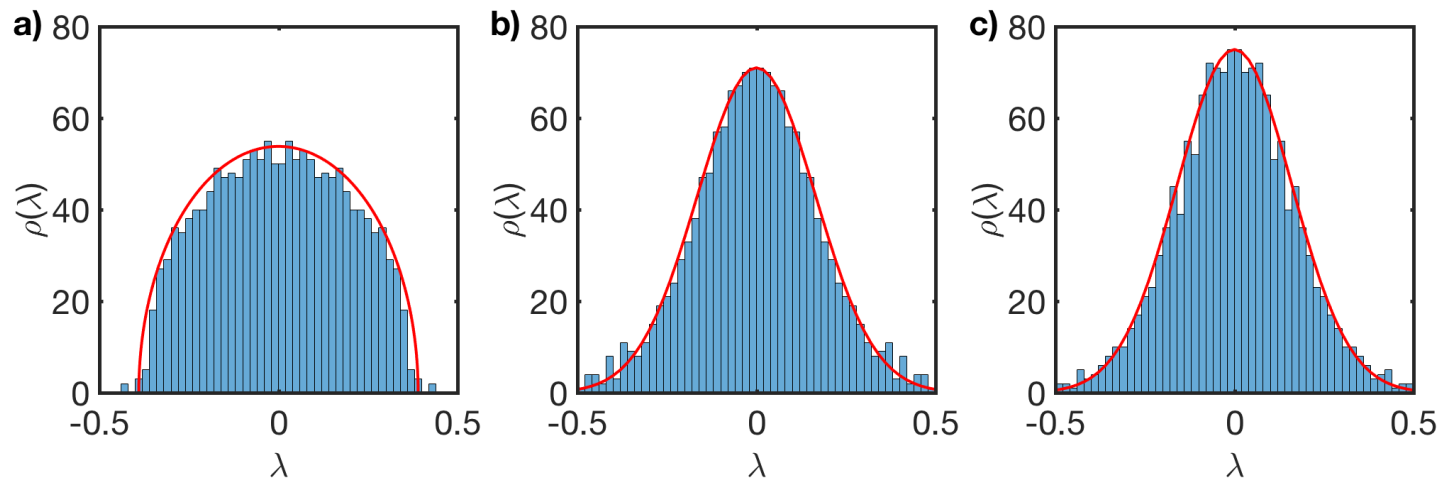
Nearly-integrable

$$J = 1, h^z = 1, h^x = 0.1$$

Thermofield MPS

$$\mathcal{H} = \sum_i \left[J \sigma_i^z \sigma_{i+1}^z + h^x \sigma_i^x + h^z \sigma_i^z \right]$$

Lyapunov Spectra: Start from product state near centre of spectrum

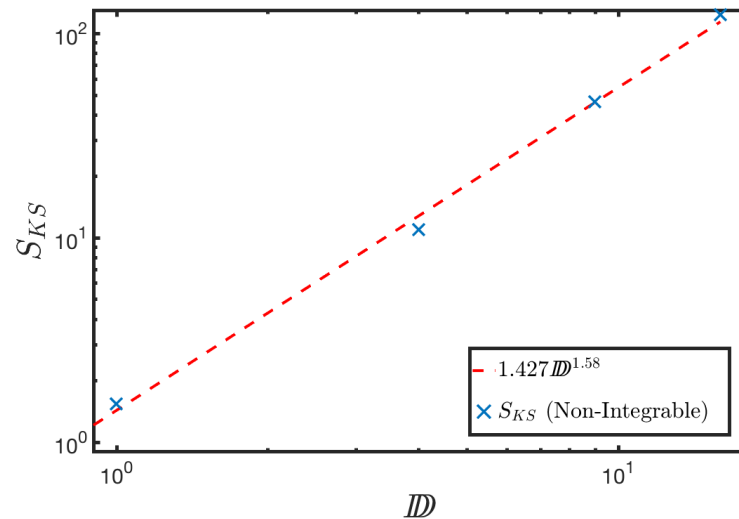


- **Semi-circle distribution also in semiclassical limit of matrix model**
[Gur-Ari, Hanada, and Shenker, JHEP('16), Hanada, Shimada, and Tezuka, PRE ('18)]

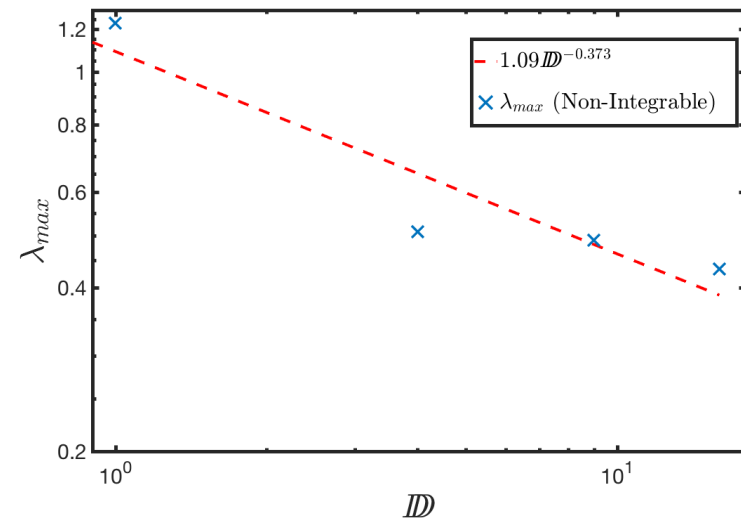
Thermofield MPS

$$\mathcal{H} = \sum_i \left[J \sigma_i^z \sigma_{i+1}^z + h^x \sigma_i^x + h^z \sigma_i^z \right]$$

Bond Order Dependence:



Kolmogorov-Sinai Entropy

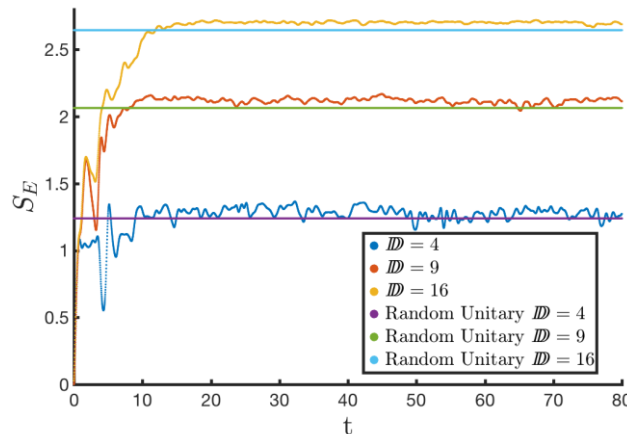


Max Lyapunov

Thermofield MPS

$$\mathcal{H} = \sum_i \left[J \sigma_i^z \sigma_{i+1}^z + h^x \sigma_i^x + h^z \sigma_i^z \right]$$

Bond Order Dependence:



$$\mathbb{D} = 1 : \quad A^{\sigma\delta} = \delta^{\sigma\delta} / \sqrt{2}$$

$$\mathbb{D} > 1 : \quad A_{IJ}^{\sigma\delta} = \frac{1}{\sqrt{d}} \sum_{\gamma=1}^d U_{(\sigma i),(\gamma j)} U_{(\delta i'),(\gamma j')}$$

$$U \in SU(d\sqrt{\mathbb{D}})$$

Apply arbitrary unitary to each half of Hilbert space

- **Operator entanglement: saturates near random unitary values**
- **Equivalent representations of infinite temperature state**
- **Seed of how to compress thermofield MPS**

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Conclusions:

- Eigenstate Thermalisation vs Dynamical Chaos
- Map quantum to classical Hamiltonian dynamics
 - TDVP of Wavefunction MPS
 - TDVP of Thermofield MPS
- Extract the full Lyapunov spectrum

Wavefunction MPS

$$\dot{S}_E(t) = \frac{S_{\mathcal{K}\mathcal{S}}(D(t))}{(D(t) - 1)^2}$$

Thermofield MPS

