

A decorative grid of circles in two rows. The top row has 12 circles, with the 5th and 12th being red and the others grey. The bottom row has 11 circles, with the 2nd, 5th, 6th, and 10th being red and the others grey.

Paola Cappellaro

Localization and thermalization in nuclear spin systems

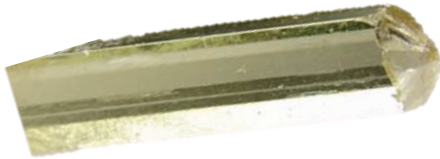
A decorative grid of circles in two rows. The top row has 10 circles, with the 3rd being red and the others grey. The bottom row has 10 circles, with the 2nd, 4th, and 6th being red and the others grey.

KITP,
August 28, 2018



Quantum Dynamics with Nuclear Spins

1. Buy a pendant on **ebay**™ ...

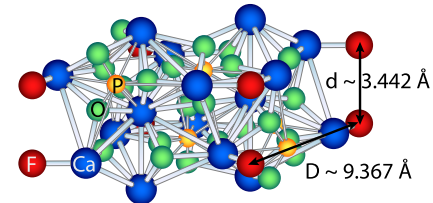


2. Use an old NMR spectrometer...

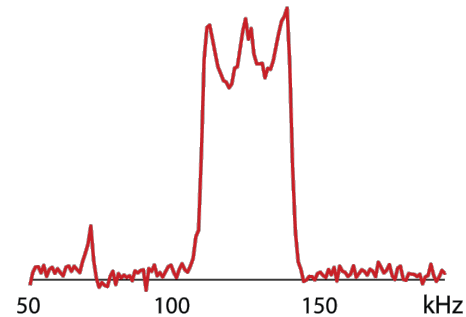


3. Observe localization & prethermalization

1. Fluorapatite Spin system & approximate model



2. Hamiltonian engineering & Multiple Quantum Coherences



3. Novel metric of localization & OTOC measurement

Poster this afternoon!

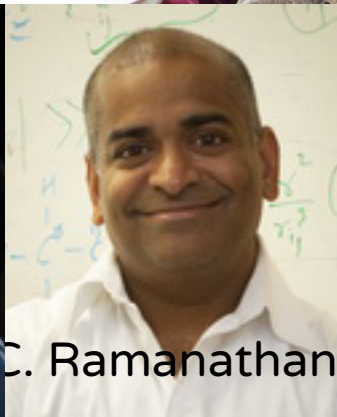


Ken Xuan Wei

Pai Peng



I. Marvian



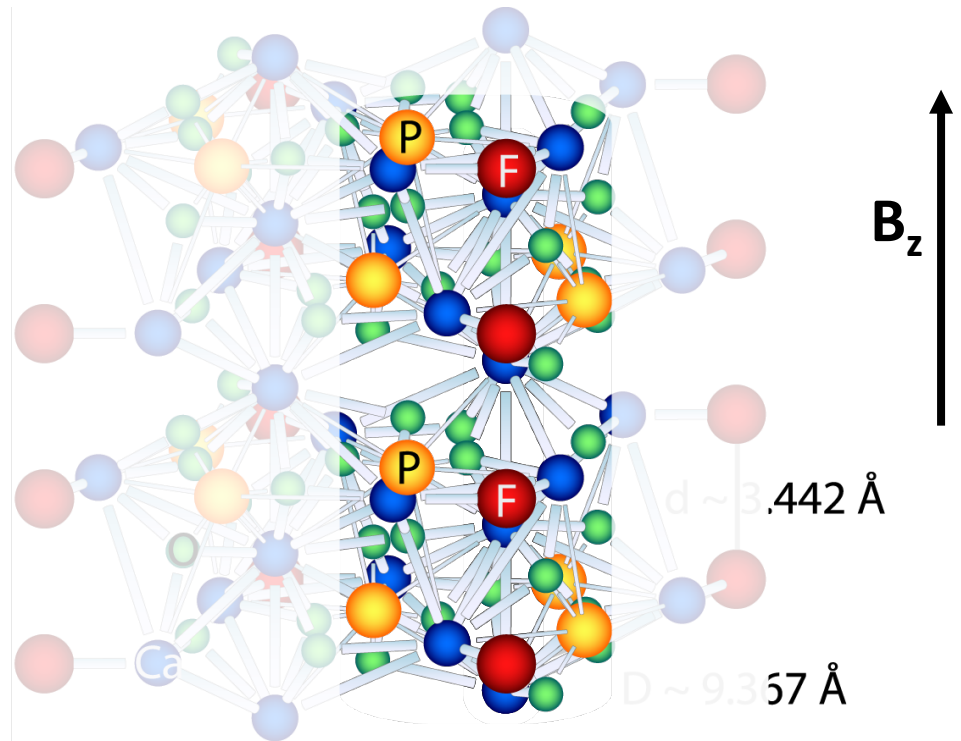
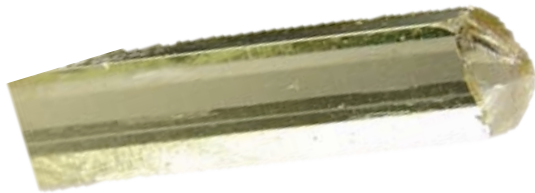
C. Ramanathan



S. Lloyd

1D system: FluorApatite

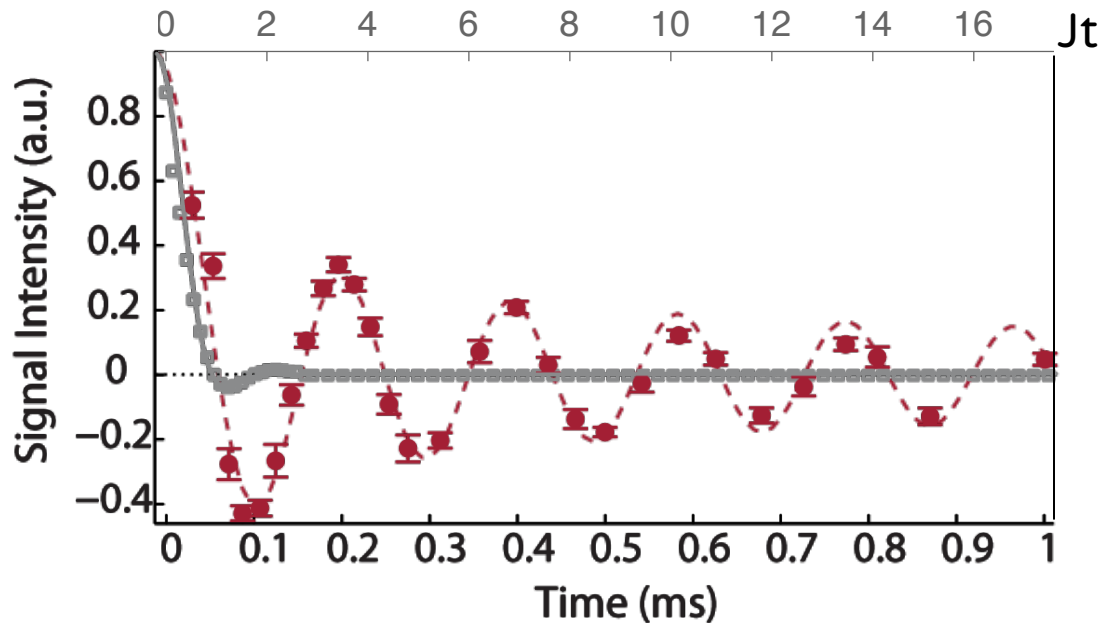
- Single-crystal, $\text{Ca}_5\text{F}(\text{PO}_4)_3$
- Quasi-1D system of ^{19}F spin $\frac{1}{2}$
 - Ratio of couplings:
 $C_{\text{in}}/C_{\text{x}} = D_{\text{x}}^3/d_{\text{in}}^3 \approx 40$
- Intrinsic disorder:
 - ^{31}P spin $\frac{1}{2}$
 - Other defects



Cho, Yesinowski, J. Phys. Chem. 1993
P.C., Ramanathan, Cory, PRL 2007

^{19}F spin dynamics

- Thermal state: $\rho = \frac{e^{-\beta\mathcal{H}}}{Z} \approx \frac{\mathbb{1}}{2^N} + \epsilon \sum_i \sigma_z^i$ ($\epsilon \approx 10^{-5}$)
- Many-body evolution: apparent loss of coherence in $<100\mu\text{s}$
- Control recovers the coherence for 1-2ms,
 - Signal shows distinct 1D, nearest-neighbor behavior



Van Lugt, Casper, Physica 1964, Zhang, PC, Viola et al. PRA 2009

Hamiltonian Engineering

- Dipolar Interaction with “long” range coupling $J \sim 1/r^3$

$$H_F = \sum_{j < k} J_{jk}^F (2\sigma_z^j \sigma_z^k - \sigma_x^j \sigma_x^k - \sigma_y^j \sigma_y^k) + \sum_{j, \kappa} h_{j\kappa} \sigma_z^j s_z^\kappa$$

- Use Average Hamiltonian* techniques to create Floquet Hamiltonian

$$H = u \sum J(\sigma_x^j \sigma_x^{j+1} + \sigma_y^j \sigma_y^{j+1})$$

integrable

$$+v \sum J(\sigma_z^j \sigma_z^{j+1})$$

interaction

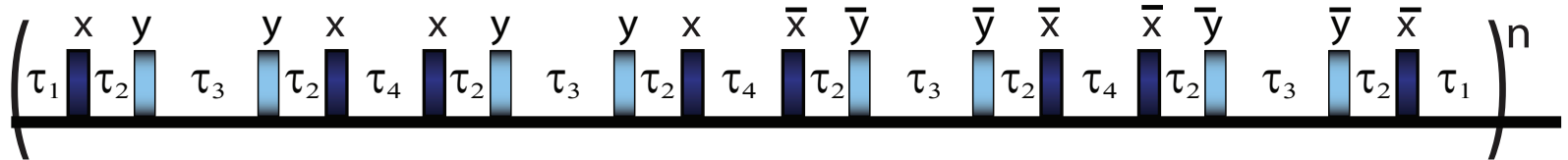
$$+g \sum h_j \sigma_z^j$$

disorder

*U. Haeblerlen, 1976

Floquet Hamiltonian

- Periodic pulse sequence stroboscopically creates desired H



$$\mathcal{T} \left\{ e^{-i \int_0^t [H_F + H_c(t')] dt'} \right\} \equiv e^{-i H t}$$

- Use Average Hamiltonian* techniques to create Floquet Hamiltonian

$$H = u \sum J(\sigma_x^j \sigma_x^{j+1} - \sigma_y^j \sigma_y^{j+1})$$

integrable

$$+v \sum J(\sigma_x^j \sigma_x^{j+1} + \sigma_y^j \sigma_y^{j+1} - 2\sigma_z^j \sigma_z^{j+1})$$

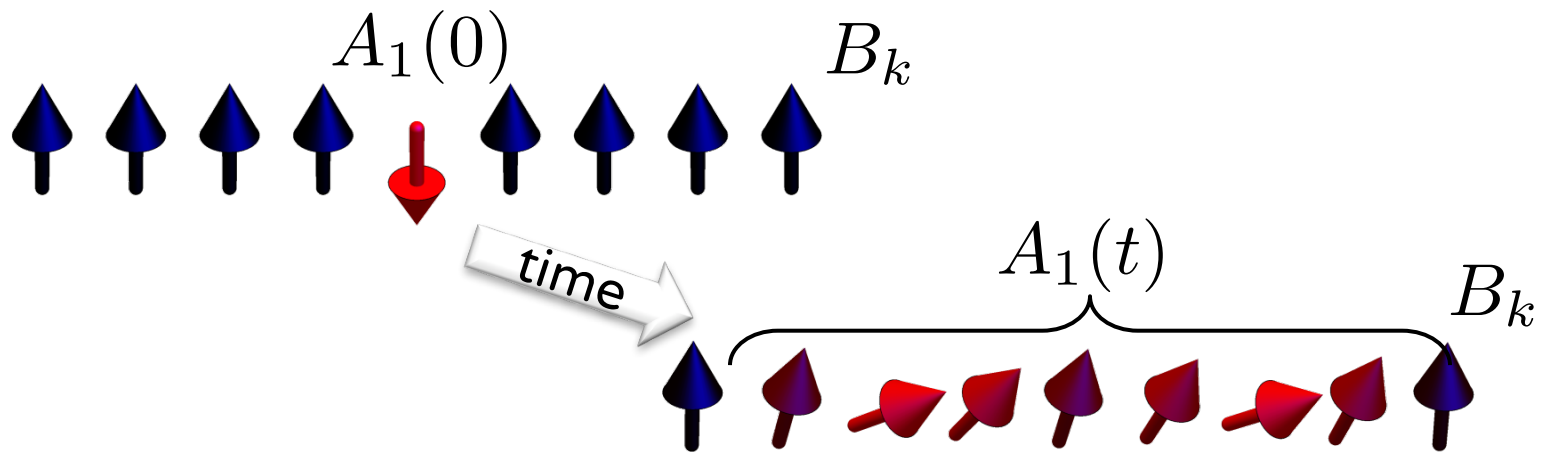
$$+g \sum h_j \sigma_z^j$$

$$+g' \sum \sigma_x^j$$

transverse field

Out-of-Time-Order Commutator

- Has a (traveling) perturbation affected a distant operator?
 - Do distant operators no longer commute, $[A_1(t), B_k(0)] \neq 0$?
 - Measure :
- $$\mathcal{C}(t) = \langle \|[A_1(t), B_k(0)]\|^2 \rangle_\beta$$
- Metric for localization, quantum information scrambling, quantum criticality, quantum chaos, ...



Out-of-Time-Order Commutator

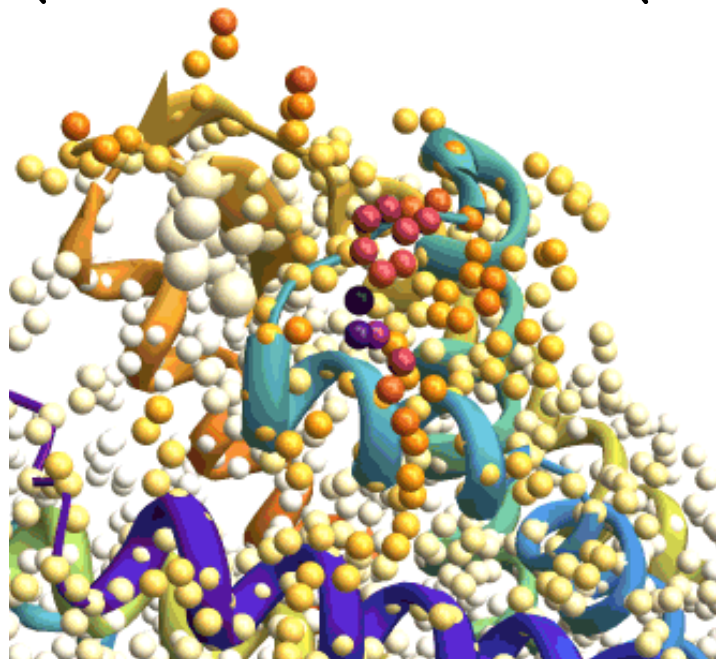
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- Metric for localization, quantum information scrambling, quantum criticality, quantum chaos, ...
- Extensive quantities, $B = \sum_k B_k$, reveal same effect

How to measure OTOc?

- “Spin counting” experiments in NMR:
 - How many correlated spins are there?
 - Gives information about sample geometry/structure
- Tool: Multiple Quantum Coherences (MQC)



Baum, Pines, JACS 1986, Munowitz, Pines, Mehring, JCP 1987

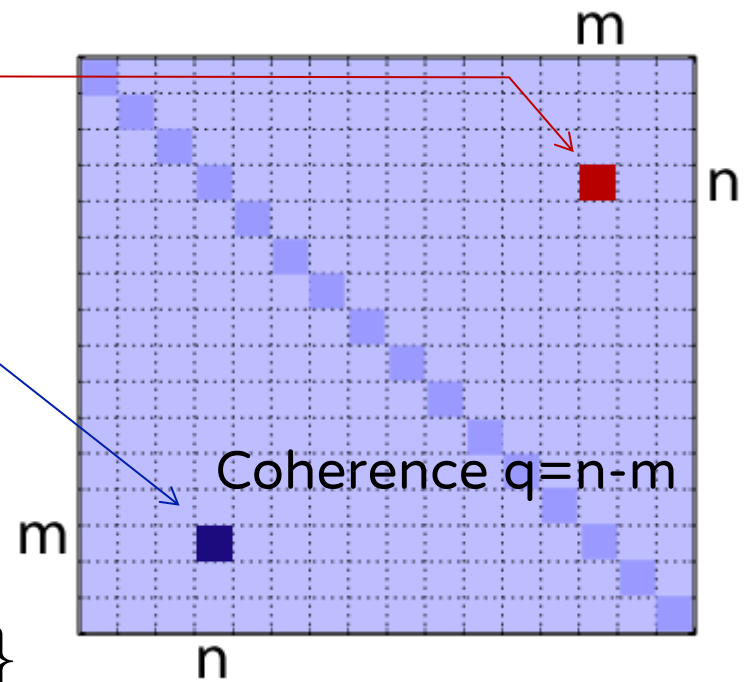
Multiple Quantum Coherences

- Tool: Multiple Quantum Coherences

– Coherent superposition: $|\psi\rangle = \frac{|m\rangle + |n\rangle}{\sqrt{2}} \xrightarrow{\sigma_z} \frac{|m\rangle + e^{i\varphi(n-m)} |n\rangle}{\sqrt{2}}$

$$\rho = |\psi\rangle\langle\psi| = \frac{1}{2} (|m\rangle\langle m| + |n\rangle\langle n| + |m\rangle\langle n| + |n\rangle\langle m|)$$

$$\rho = \sum_{q=-L}^L \rho_q \quad \rho_q \xrightarrow{\sigma_z} \rho_q e^{i\varphi q}$$



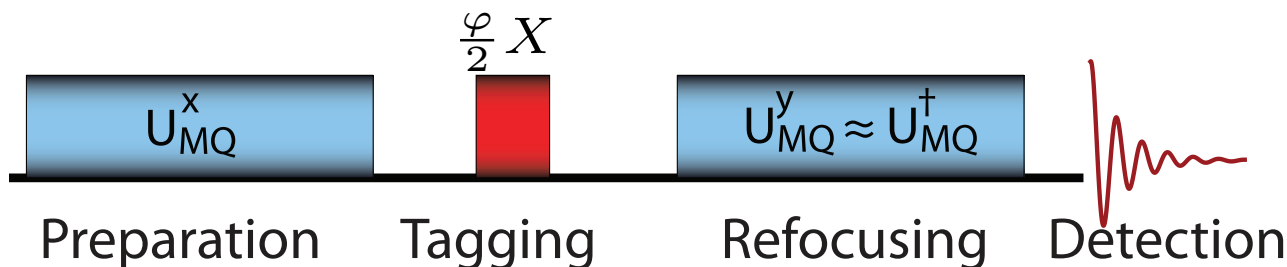
- MQC intensities: $I_q = \text{Tr} \{ \rho_q \rho_{-q} \}$

OTOC & MQC

- We can measure OTOC:

-> NMR has been doing it for the past 40 years!

1. Tweaking the pulse sequence we achieve **time reversal**
2. Exploiting initial **mixed state** $\rho_0 = \mathbb{1} + \epsilon Z$ we measure OTOC at $T = \infty$



- The signal is $S_\varphi(t) = \text{Tr}[\Phi^\dagger(t)\rho_0\Phi(t)Z]$
 $S_\varphi(t) = \text{Tr}[\Phi^\dagger(t)Z\Phi(t)Z]$
 with $\Phi(t) = U(t)e^{i\frac{\varphi}{2}X}U^\dagger(t)$

- OTO correlator for Φ at $\beta = 0$

Munowitz, Pines, Mehring, JCP 1987

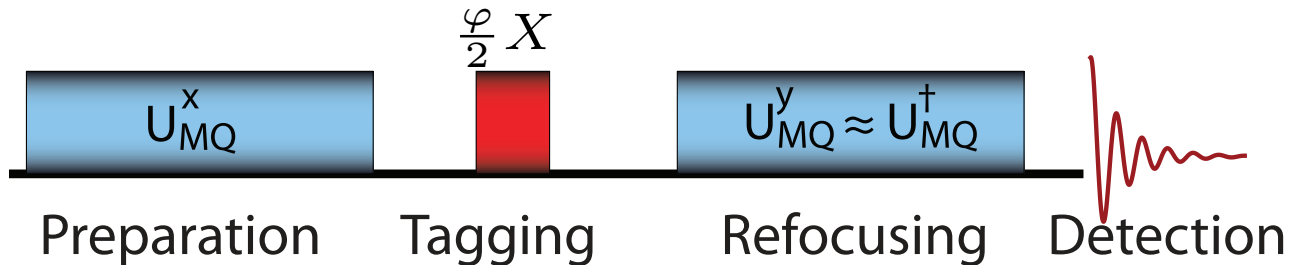
OTOC & MQC

- We can **better** measure OTOC:

-> by detecting NMR multiple quantum coherence distribution

1. Repeat the measurement varying φ

2. Fourier transform the signal to get MQC intensities, $I_q = \sum_{\varphi} S_{\varphi} e^{i\varphi q}$



- Expanding in power of φ , $S_{\varphi} \sim Z^2 + \|[Z(t), X]^2\| \varphi^2 + \dots$
which can be reliably extracted from MQC 2nd moment

$$\langle \|[Z(t), X]^2\| \rangle_{\beta=0} = 4 \sum_q q^2 I_q(X)$$

Correlation Length

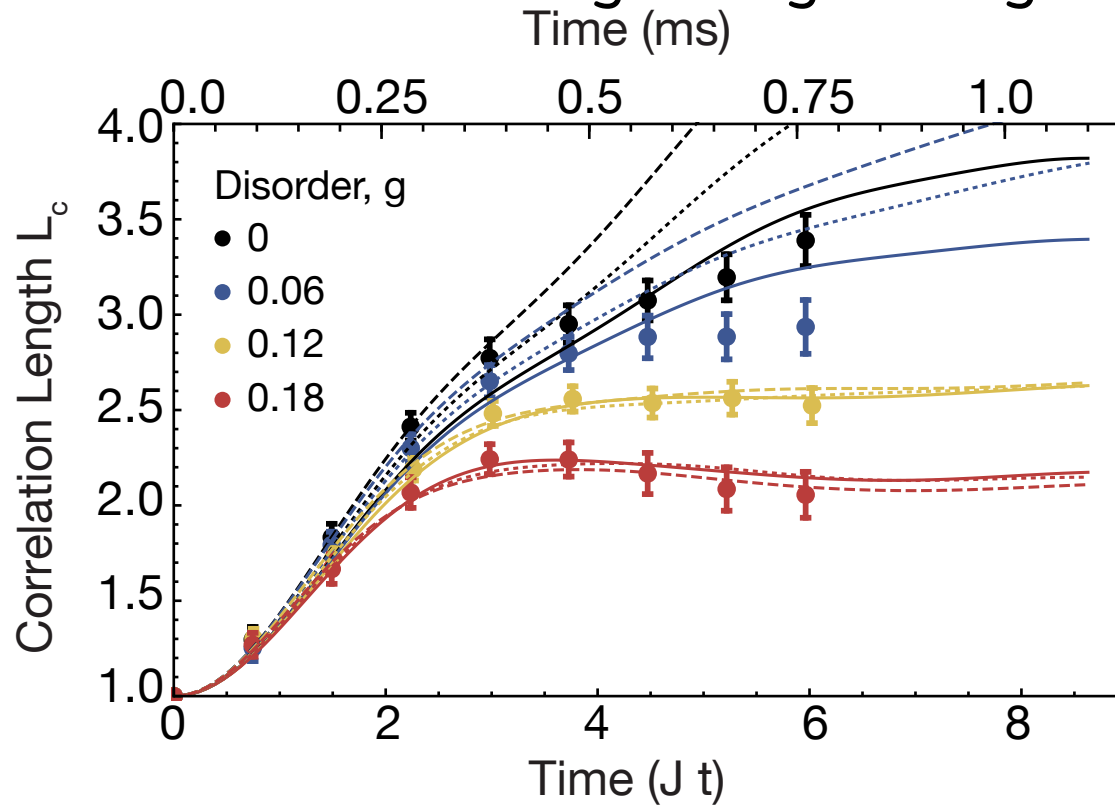
- For 1D systems, close to non-interacting, we can even relate MQC to the length over which correlations spread



- Defined average correlation length L_c for mixed states
- Explored differences in behavior between Anderson and Many-Body localization

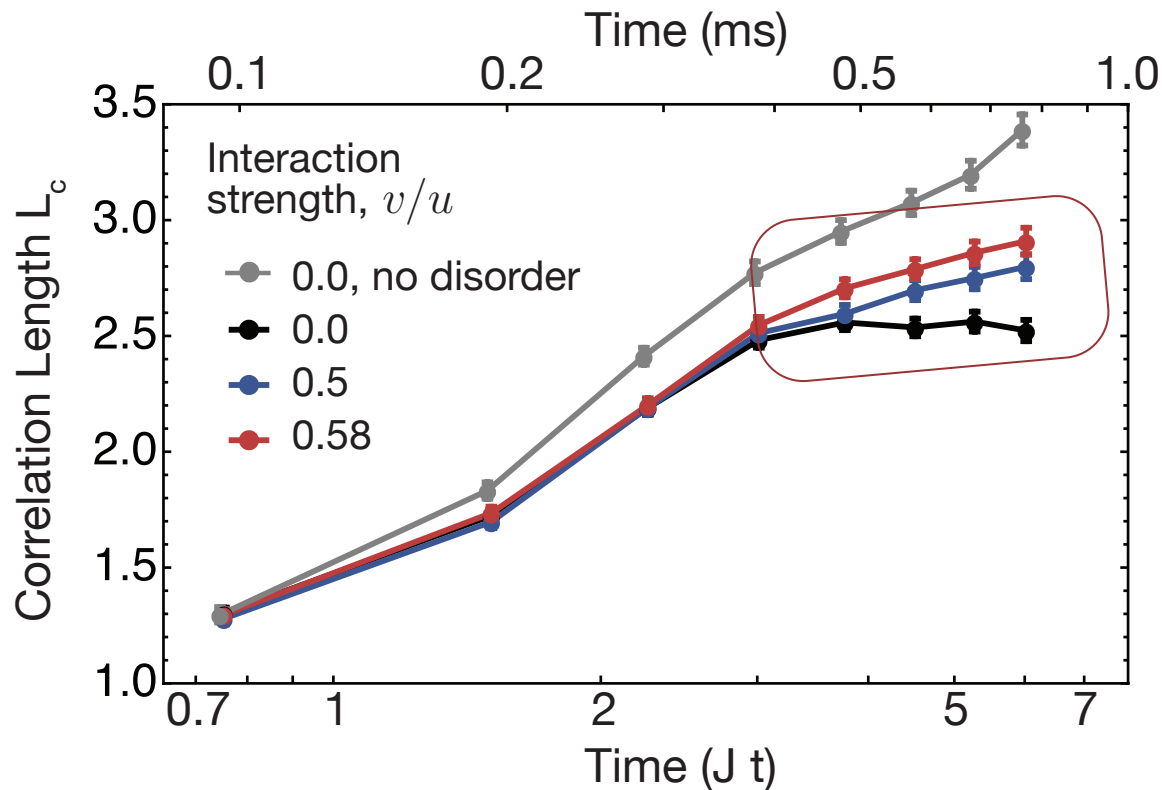
Experimental Results: AL

- L_c grows as a function of time.
- With disorder of increasing strength the growth is quenched



Experimental Results: MBL

- Now add the interaction term, $\sim v\sigma_z\sigma_z$:
- The correlation length keeps growing, but very slowly



Thermalization

- Dipolar Hamiltonian (XXZ, $J \sim r^{-3}$) leads to thermalization
 - Can a transverse field protect quantum information?

$$\mathcal{H}_{\text{TFD}} = \sum_{j,k} J_{jk} \left[S_z^j S_z^k - \frac{1}{2} (S_y^j S_y^k + S_x^j S_x^k) \right] + g \underbrace{\sum S_x^j}_X$$

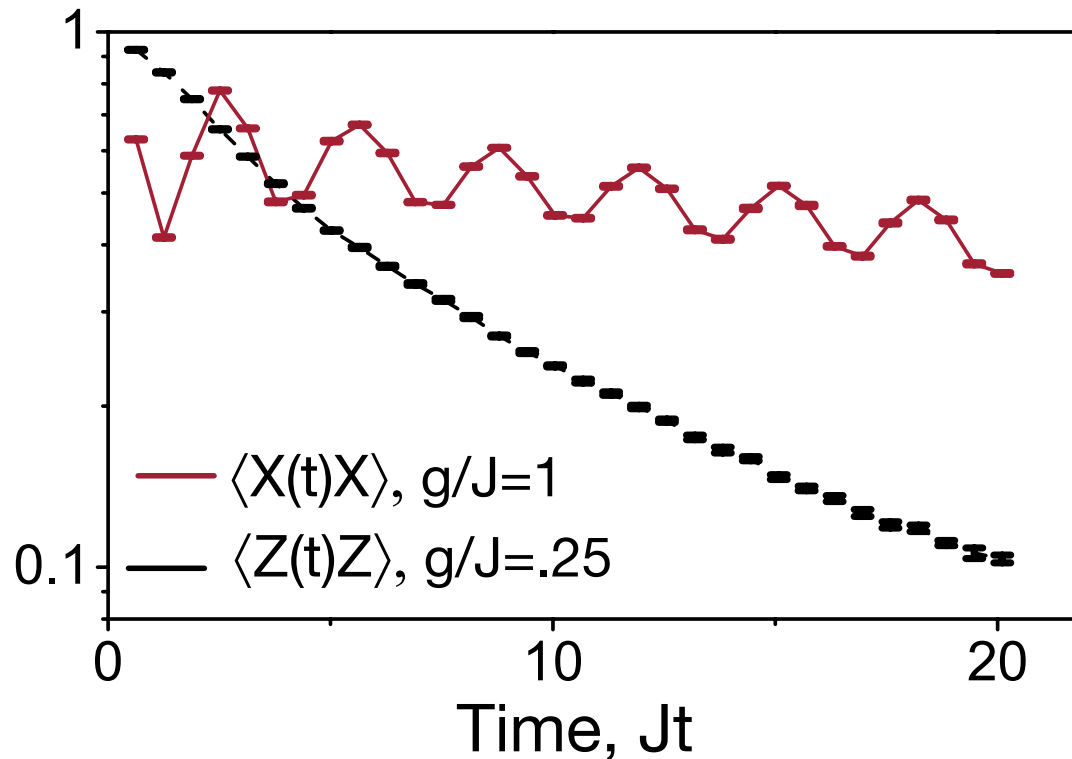
- In the limit of $g \gg J$, the transverse field becomes a quasi-conserved quantity X_{pre} , and the system enters a long quasi-equilibrium regime



under Floquet driving

Dynamics of Magnetization

- The transverse (X) and longitudinal (Z) magnetization show distinct behaviors, linked to the **thermal** and **prethermal** regimes



Thermalization

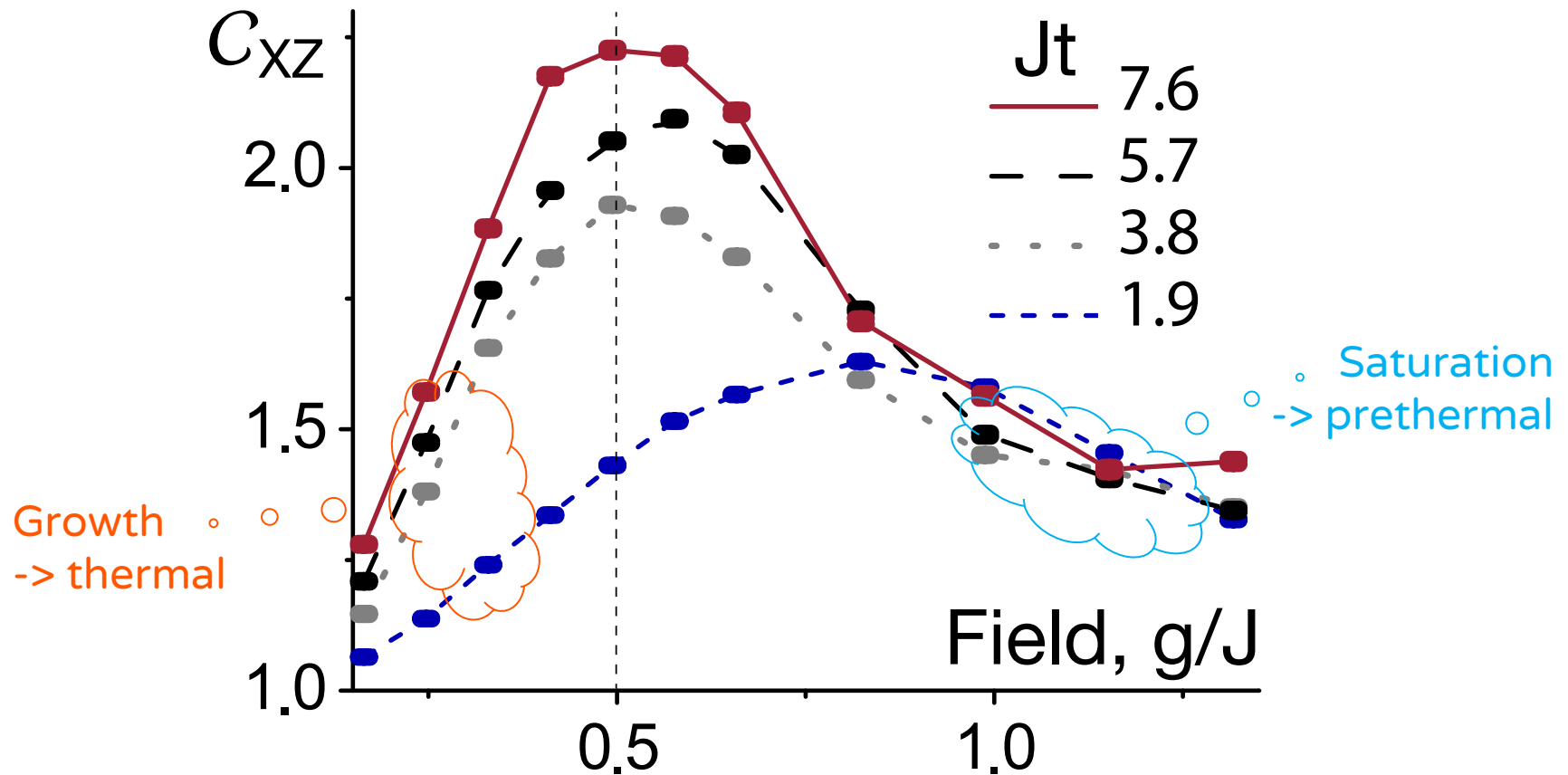
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- In the limit of $g \gg J$, the transverse field becomes a quasi-conserved quantity X_{pre} , and the system enters a long quasi-equilibrium regime
- Measuring OTOC's involving the conserved quantity, we can distinguish between **thermalization** and **prethermal** behavior
$$C_{\text{XZ}} = \langle |[X(t), Z]|^2 \rangle$$

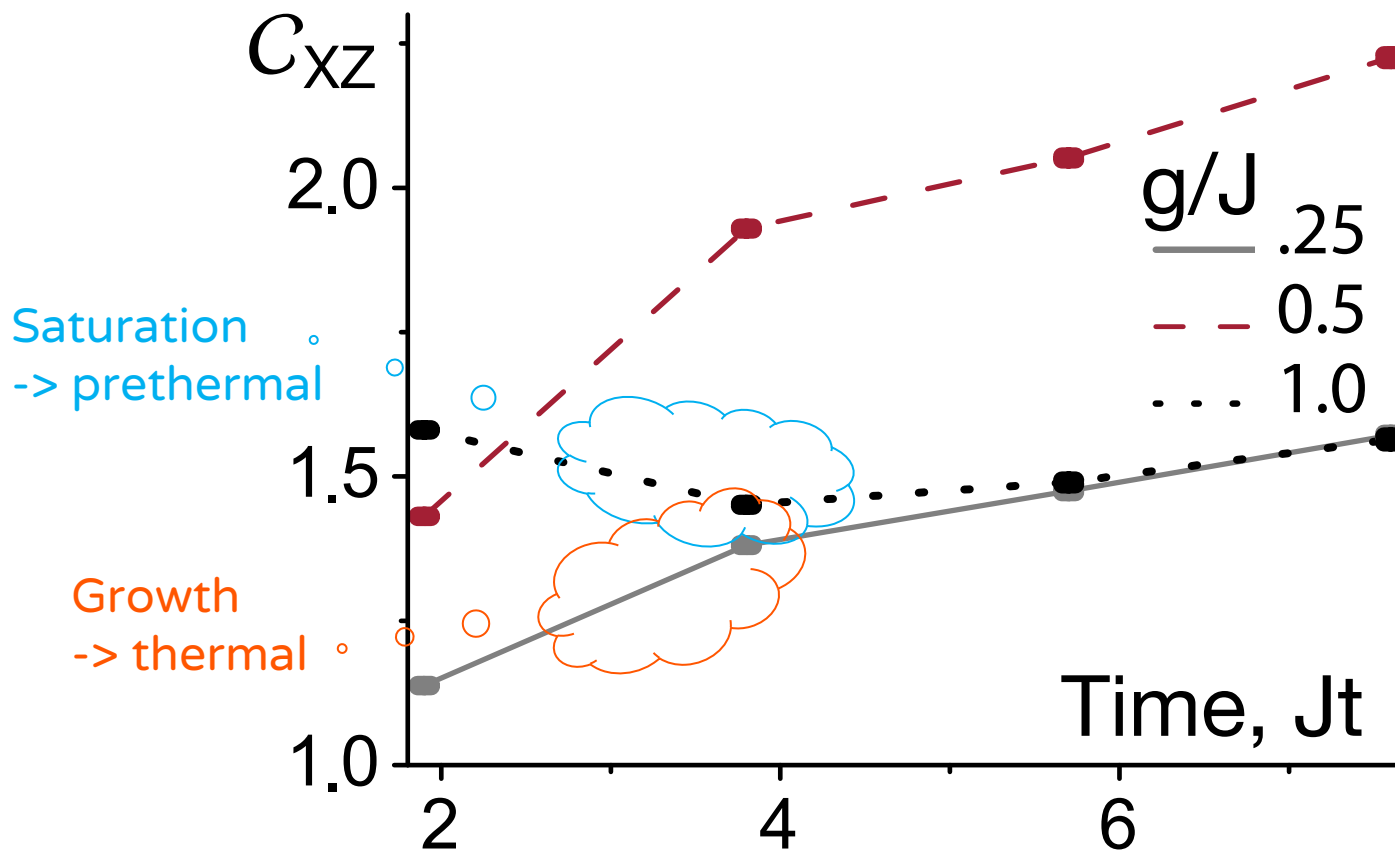
Prethermalization : Experiments

- Transverse field dipolar Hamiltonian (experiments)



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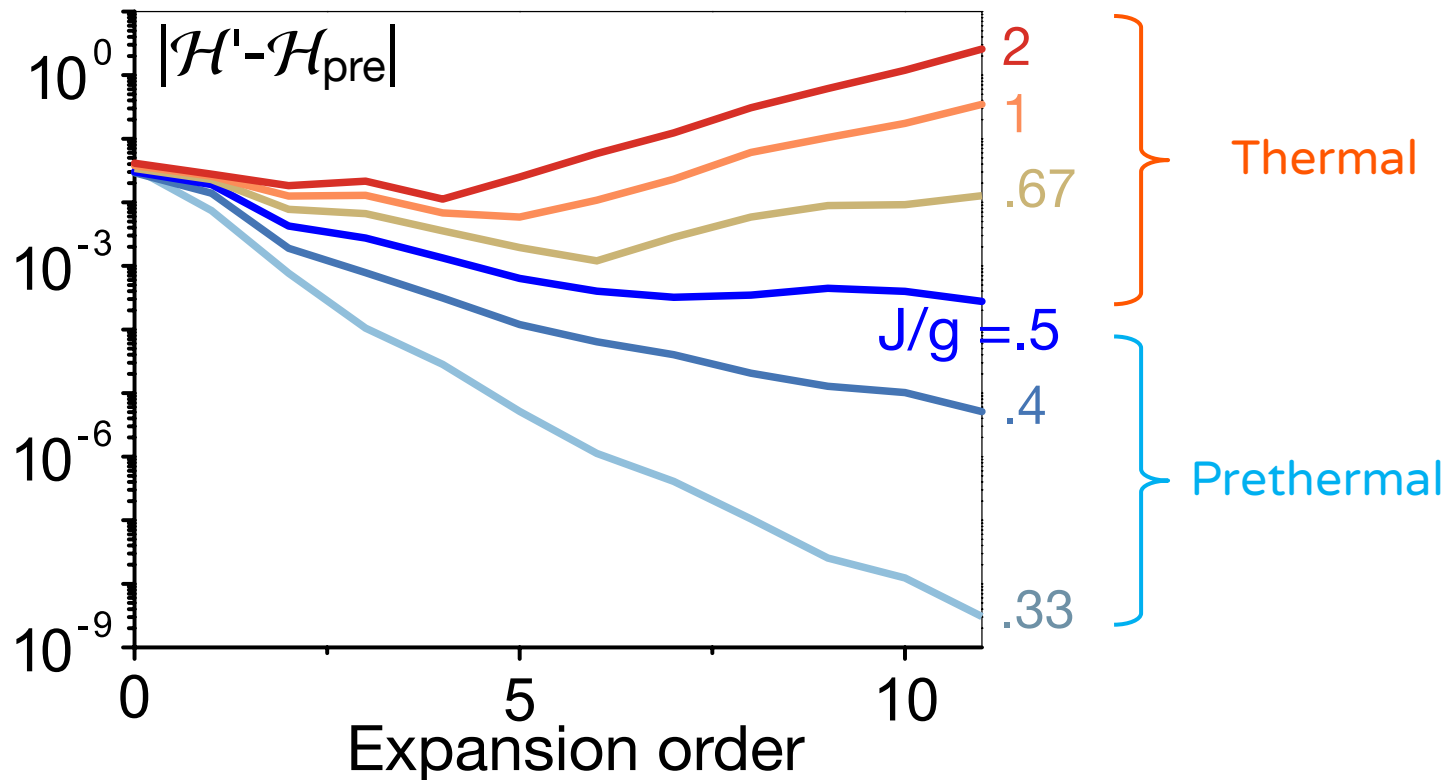


Prethermal Hamiltonian

- To check our interpretation we explicitly construct the prethermal Hamiltonian $\mathcal{H}_{\text{pre}} = \mathcal{H}' - \delta\mathcal{H}(e^{-J/g})$
 - We need to find the unitary transformation $\mathcal{H}' = U\mathcal{H}_{\text{TFD}}U^\dagger$
 - We use a recursive expansion in J/g : only if system is prethermal the expansion converges and $|\mathcal{H}_{\text{pre}} - \mathcal{H}'| \approx 0$

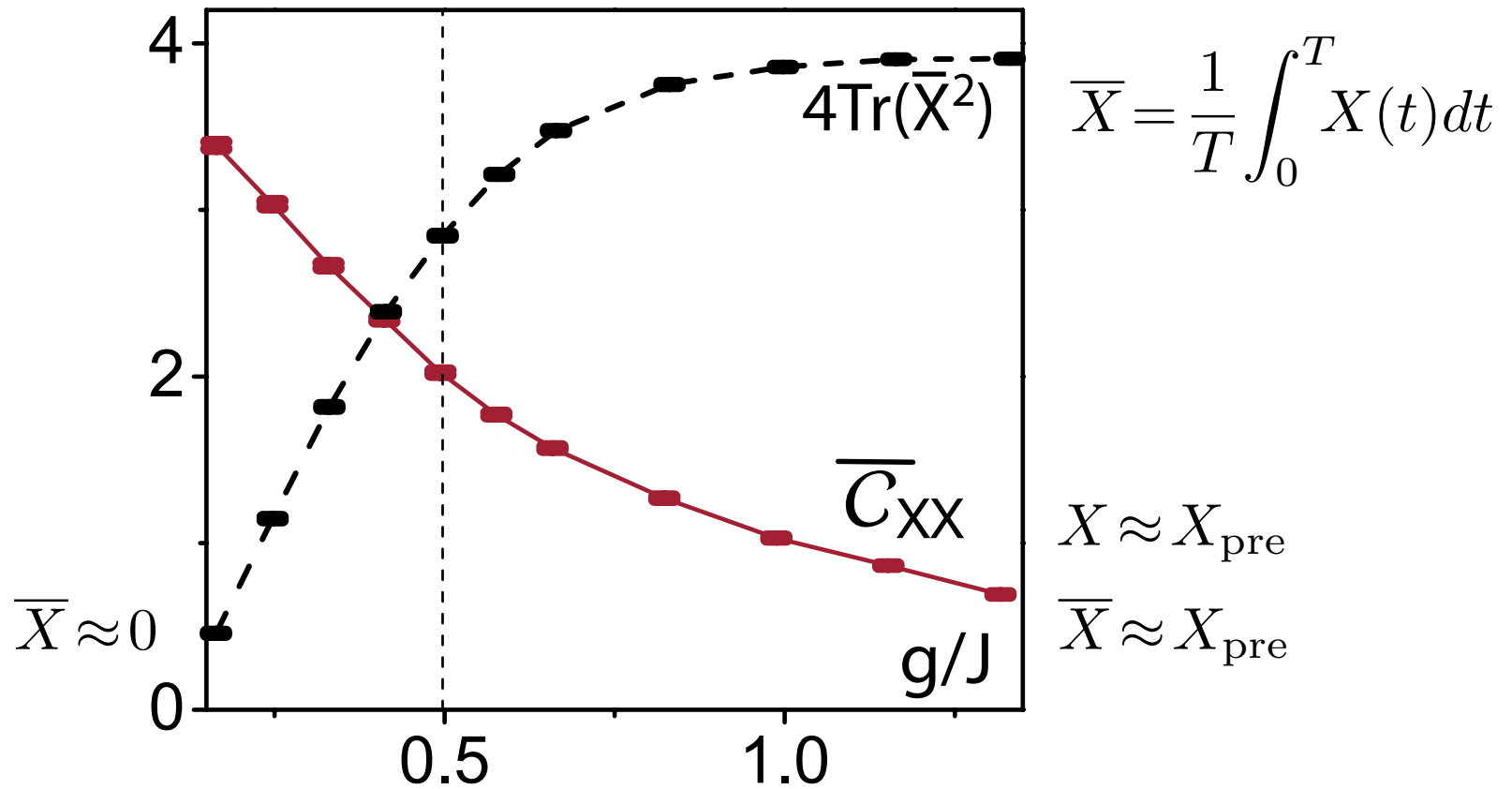
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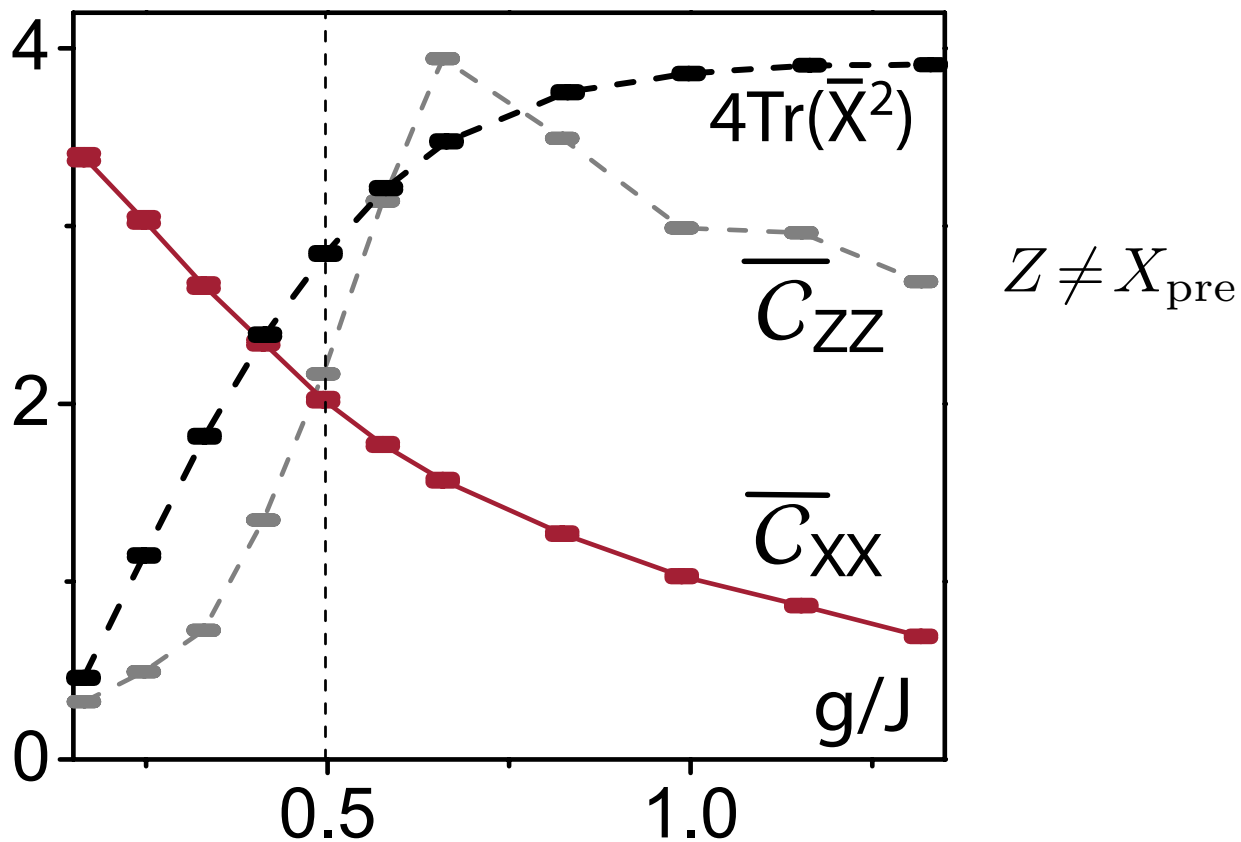
OTOC Dynamics

- Dynamics of other OTO commutators provides insight



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