

HK Moffatt: Discussion on 16 May '08.

$$\hat{u}(\underline{k}, \omega) = -\frac{A(\underline{k}, \omega)}{D(\underline{k}, \omega)} \hat{\theta}(\underline{k}, \omega)$$

where

$$A = \frac{-(\underline{k} \cdot \underline{B}_0)^2 \underline{k}_n (\underline{k}_n \underline{g})}{\eta k^2 - i\omega} - 2 \underline{k} \cdot \underline{\Omega} (\underline{k}_n \underline{g})$$

$$D = 4(\underline{k} \cdot \underline{\Omega})^2 + \frac{(\underline{k} \cdot \underline{B}_0)^4 k^2}{(\eta k^2 - i\omega)^2}$$

[ $D=0$  is dispersion relation for magnetostrophic waves damped by magnetic resistivity  $\eta$ ]

$$\hat{b}_i(\underline{k}, \omega) = \frac{i(\underline{B}_0 \cdot \underline{k}) \hat{u}_i}{\eta k^2 - i\omega}$$

$$i \underline{A}^* \cdot \underline{A} = \frac{-4(\underline{k} \cdot \underline{\Omega}) \omega (\underline{k} \cdot \underline{B}_0)^2 (\underline{k}_n \underline{g})^2 k}{\eta^2 k^4 + \omega^2} \quad (\text{real})$$

Hence helicity spectrum

$$H(\underline{k}, \omega) = \frac{4(\underline{k} \cdot \underline{\Omega}) \omega (\underline{k} \cdot \underline{B}_0)^2 k^2 (\underline{k}_n \underline{g})^2 \Gamma(\underline{k}, \omega)}{|\mathcal{D}|^2 (\eta^2 k^4 + \omega^2)}$$

$\Gamma(\underline{k}, \omega)$  = spectrum fu. of  $\theta$   
(assumed given)

and hence

$$\alpha = \frac{1}{3} \kappa_{ii} = -\frac{\eta}{3} \iint \frac{k^2 H(\underline{k}, \omega)}{\omega^2 + \eta^2 k^4} d\underline{k} d\omega \quad (R_m \leq 1) \quad \text{FOS}$$

$\kappa$ -quenching is contained in this expression

Buoyancy Flux

$$F_z = \langle \underline{u} \theta_z \rangle = \iint \frac{(\underline{k} \cdot \underline{B}_0)^2 \eta k^2 (\underline{k}_n \underline{g})^2 \Gamma(\underline{k}, \omega) d\underline{k} d\omega}{4(\underline{k} \cdot \underline{\Omega})^2 (\eta^2 k^4 + \omega^2) + (\underline{k} \cdot \underline{B}_0)^4 (\eta^2 k^4 + \omega^2) k^2}$$

Reynolds stress

$$\tau_{ij} = \langle u_i u_j \rangle = \iint \frac{\text{Re}[A_i^* A_j]}{|\mathcal{D}|^2} \Gamma(\underline{k}, \omega) d\underline{k} d\omega$$

$$\tau_{ij}^M = \langle b_i b_j \rangle = \dots$$

These drive differential rotation and meridional circulation

$$\frac{\partial \theta}{\partial t} + \underbrace{\underline{u}(\theta) \cdot \nabla \theta}_{\text{linear functional of } \theta'}$$

$$Pe = \frac{v l}{\kappa} \gg 1 \quad (\sim 10^8 \text{ in core})$$

So this nonlinearity dominates.

Example  $\underline{u} = \nabla_{\perp} \underline{A}$  (so  $\nabla \cdot \underline{u} = 0$ )

Let  $\underline{A} = a(x) \theta$

$$\underline{u} = \nabla_{\perp} \underline{A} = -a_{\perp} \nabla \theta + \underline{t} \theta \quad (\underline{t} = \nabla_{\perp} a)$$

$$\underline{u} \cdot \nabla \theta = \theta \underline{t} \cdot \nabla \theta$$

Suppose  $\underline{t} = (0, 0, t)$  with  $t$  constant

Then we get

$$\frac{\partial \theta}{\partial t} + t \theta \frac{\partial \theta}{\partial x} = \kappa \nabla^2 \theta$$

Burger's equation: sharp fronts form in a finite time!  
So not Kolmogorov type turbulence??