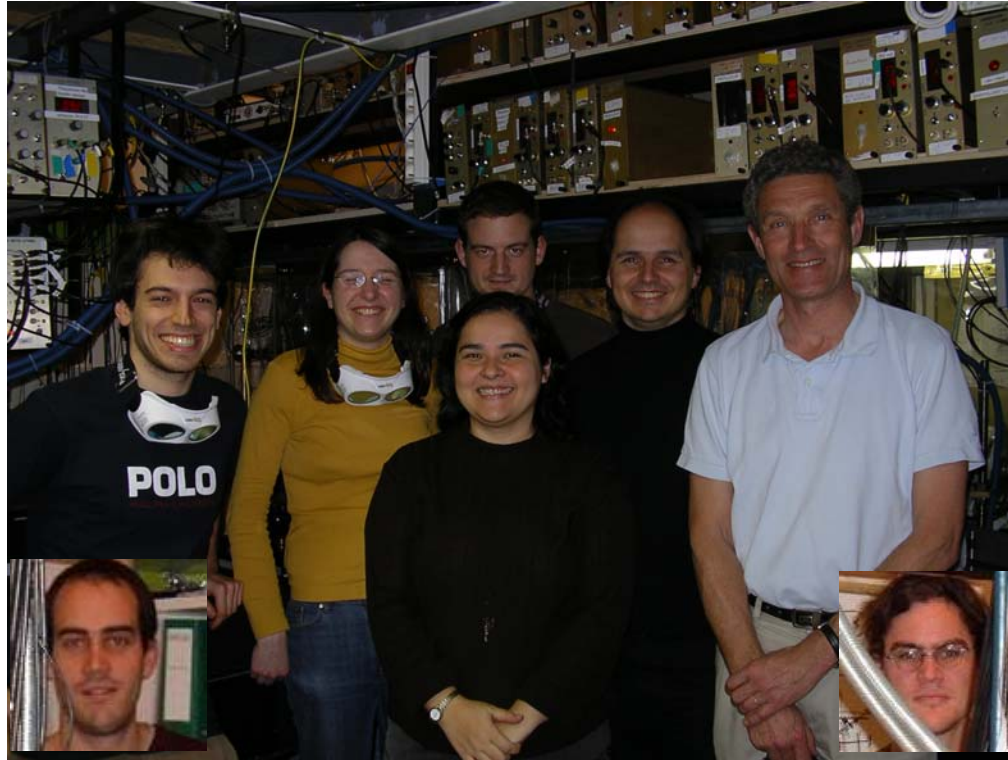


Unitary Fermi gas

Phase diagram and Raman spectroscopy



Collège de France



F. Chevy, M. Teichmann, L. Tarruell, J. McKeever,
K. Magalhães, S. Nascimbène, N. Navon and C. Salomon

Laboratoire Kastler Brossel, Paris

R. Combescot, C. Lobo, A. Recati, I. Carusotto,
S. Stringari, T. L Dao, A. Georges, J. Dalibard

Outline

Phase diagram of strongly interacting fermions at unitarity
With unbalanced spin populations

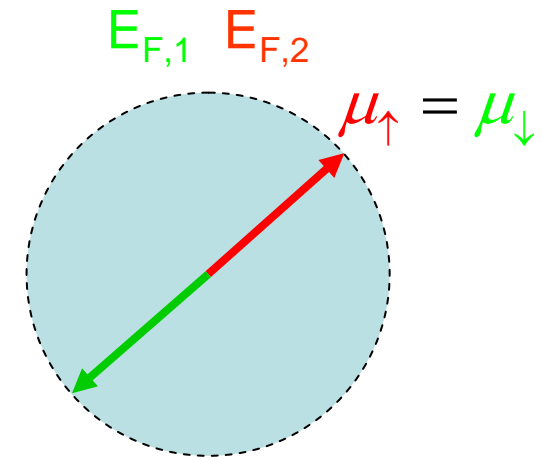
F. Chevy, *Phys. Rev. A* 74, 063628 (2006),
and cond-mat 0701350

Measurement of single particle excitations in strongly
interacting Fermi gas by stimulated Raman spectroscopy

T. L. Dao, A. Georges, J. Dalibard, C.S., I. Carusotto, cond-mat 0611206

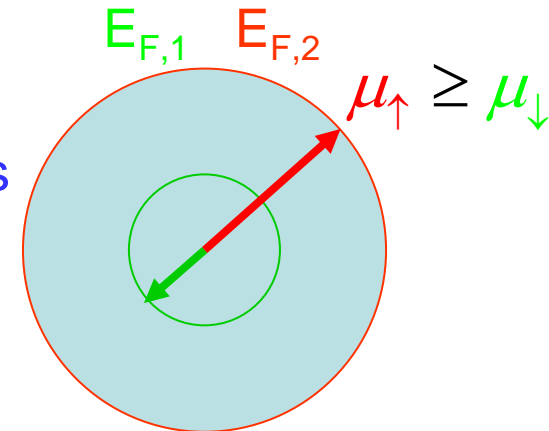
Motivation

Attractive Fermi gas with equal spin population
⇒ BCS theory, pairing at edge of Fermi surface



What is the nature and existence of superfluidity
when spin population is imbalanced ?

Mismatched density and/or pairing with different masses



Ex:

Superconductors in magnetic field or
quark matter

Cold gases: Mit and Rice expt

$$E_{F,i} = \frac{\hbar^2 k_{F,i}^2}{2m_i} = \frac{\hbar^2}{2m_i} \left(6\pi^2 n_i \right)^{2/3}$$

Overview of Theoretical scenarios

Chandrasekhar and Clogston: stability of the paired state : $\mu_{\uparrow} > \mu_{\downarrow}$

Conversion of a particle: $\downarrow \rightarrow \uparrow$

Decrease the grand potential $H - \mu_{\uparrow} N_{\uparrow} - \mu_{\downarrow} N_{\downarrow} : \mu_{\uparrow} - \mu_{\downarrow}$

Cost of pair breaking: Δ

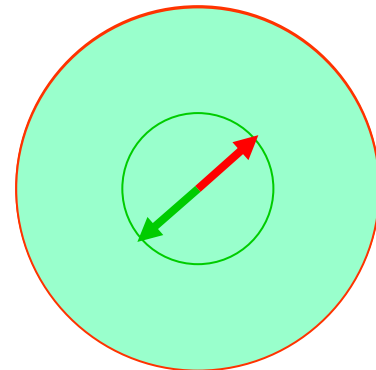
\Rightarrow Paired state stable for $\mu_{\uparrow} - \mu_{\downarrow} < \Delta$

And beyond?

Polarized phase : One spin species (Carlson, PRL **95**, 060401 (2005))

FFLO Phase (Fulde Ferrell Larkin Ovchinnikov) : pairing in $\mathbf{k}_{\uparrow} - \mathbf{k}_{\downarrow} \neq 0$
(C. Mora et R. Combescot, PRB **71**, 214504 (2005))

Sarma phase (internal gap) : pairing in $\mathbf{k}_{\uparrow} - \mathbf{k}_{\downarrow} = 0$
 \Downarrow opening of a gap in the Fermi sea of majority species. (Liu, PRL **90**, 047002 (2003))



Avalanche of recent publications !

P. Pieri and G.C. Strinati cond-mat/0512354 : diagrammatic method

Extrapolation from BEC regime

W. Yi and L.-M. Duan, cond-mat/0601006 : BCS at finite temperature

M. Haque and H.T.C. Stoof, cond-mat/0601321 : BCS at T=0

T.N. de Silva and E.J. Mueller, cond-mat/0601314 : BCS at T=0

D. Sheehy, L. Radzihovsky, PRL 06

A. Bulgac, M. McNeil Forbes '06

K. Levin et al., 06

M. Parish, Nature Physics 3 '07

.....

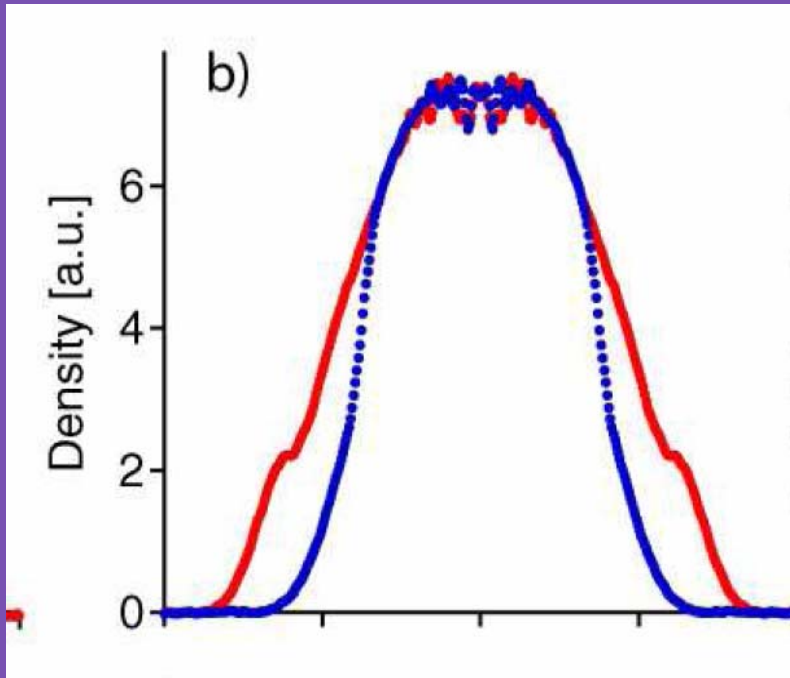
Assumptions:

1) Unitarity: universal parameter $\mu = (1 + \beta) E_F = \xi E_F$ known

2) Grand canonical description, Local density approx,

3) T=0 approach

Experimental results



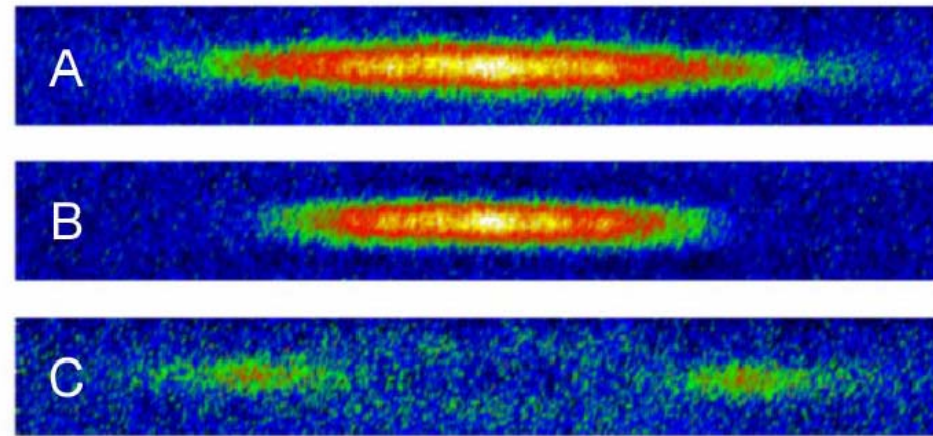
MIT: 3 phases

- Fully paired superfluid core
- Intermediate mixture
- Fully polarized rim

M.W. Zwierlein, *et al.*, *Science*, **311**
(2006) 492.

Rice: 2 phases

- Fully paired superfluid core
- Fully polarized rim



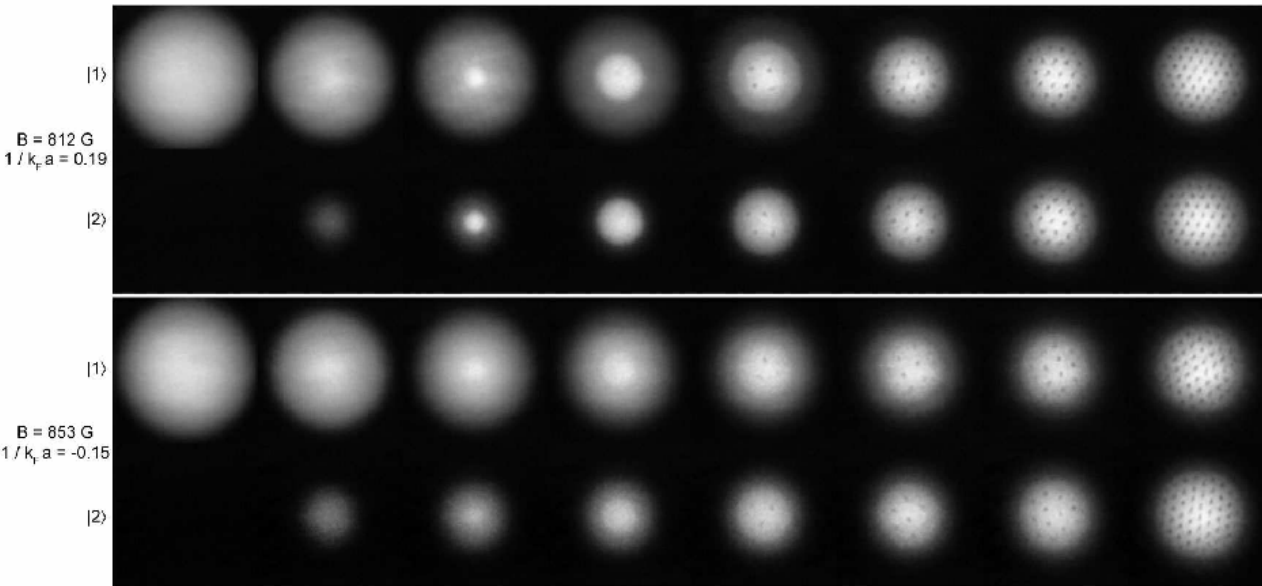
G. Partridge, W. Li, R.I. Kamar, Y.-A. Liao,
R.G. Hulet, *Science*, **311** (2006)

503.

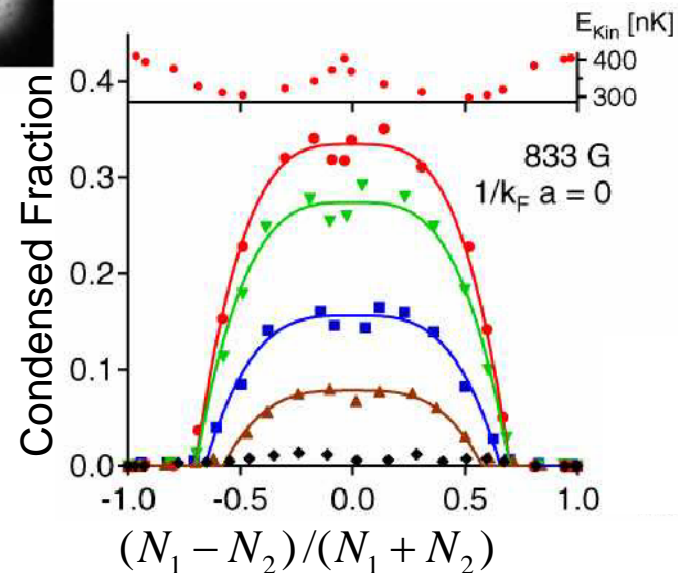
G. Partridge *et al.*, *Cond-mat* 0608455

MIT experiment

(Science Express, December 22, 2005)



Superfluidity observed in Time of flight
Loss of superfluidity for large
Spin population imbalance



Why Unitarity?

Dimensional analysis: for any *intensive* physical quantity Q (density, pressure...)

$$Q[V, \mu_{\uparrow}, \mu_{\downarrow}, m, \hbar, a] = Q_0[\mu_{\uparrow}, m, \hbar] f(\mu_{\downarrow} / \mu_{\uparrow}, 1/k_{F\uparrow} a)$$

Q_0 : value for the ideal Fermi gas;

$$\mu_{\uparrow} = \hbar^2 k_{F\uparrow}^2 / 2m$$

At Feshbach resonance, $a = \infty \Rightarrow Q/Q_0 = g(\mu_{\downarrow} / \mu_{\uparrow})$ only!

Application: universal equation of state of the balanced Fermi gas

For instance: Q =density, balanced Fermi gas ($\mu_{\uparrow} = \mu_{\downarrow}$)

$$n = \frac{1}{6\pi^2} \left(\frac{2m\mu_{\uparrow}}{\hbar^2} \right)^{3/2} \quad \text{x numerical factor}$$

$$\mu_{\uparrow} = \xi \frac{\hbar^2}{2m} (6\pi^2 n)^{2/3} = \xi E_F$$

Determination of ξ

Experiment	<i>ENS (⁶Li)</i>	0.41(15)	Theory	<i>BCS</i>	0.59
	<i>Rice (⁶Li)</i>	0.46(5)		<i>Astrakharchik</i>	0.42(1)
	<i>JILA(⁴⁰K)</i>	0.46(10)		<i>Perali</i>	0.455
	<i>Innsbruck (⁶Li)</i>	0.27(10)		<i>Carlson</i>	0.42(1)
	<i>Duke (⁶Li)</i>	0.51(4)		<i>Hausmann</i>	0.36

Universal phase diagram of the homogeneous unitary system

(F. Chevy, PRA Phys. Rev. A 74, 063628 (2006),
A. Bulgac, M. McNeil Forbes, cond-mat/0606043)

« Exact » eigenstates of the grand potential

- Single component ideal gas $\Omega = \Omega_0$
- Fully paired superfluid

$|\text{SF}\rangle_\mu$ eigenstate of the *balanced* grand-potential

$$\hat{\Omega}' = \hat{H} - \mu(\hat{N}_\uparrow + \hat{N}_\downarrow)$$

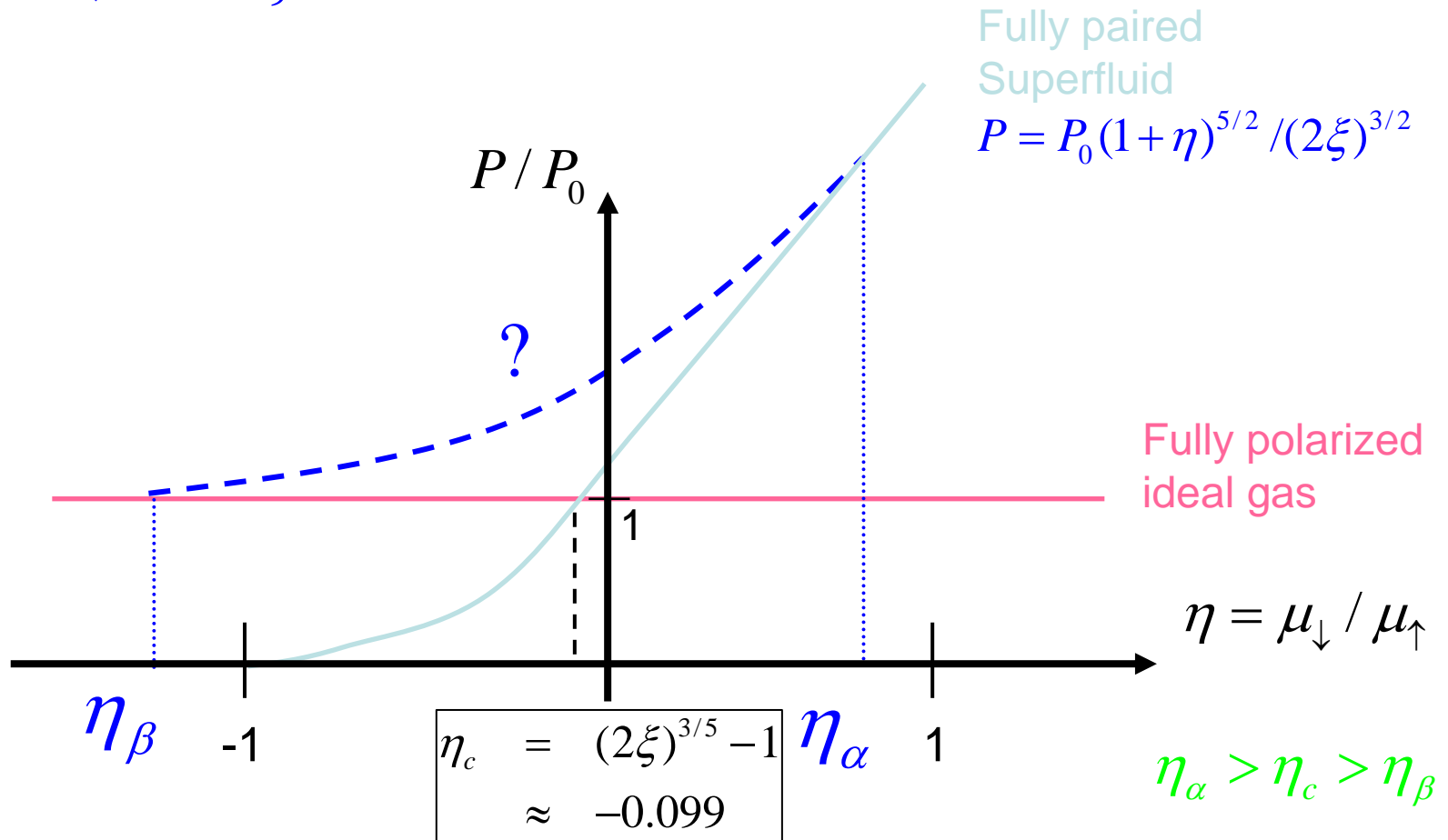
$$\hat{\Omega}' |\text{SF}\rangle_\mu = \Omega' |\text{SF}\rangle_\mu \quad \hat{N}_\uparrow |\text{SF}\rangle_\mu = \hat{N}_\downarrow |\text{SF}\rangle_\mu = N |\text{SF}\rangle_\mu$$

$$\hat{\Omega} = \hat{H} - \mu_\uparrow \hat{N}_\uparrow - \mu_\downarrow \hat{N}_\downarrow \quad \Rightarrow \quad \hat{\Omega} |\text{SF}\rangle_{(\mu_\uparrow + \mu_\downarrow)/2} = \Omega' |\text{SF}\rangle_{(\mu_\uparrow + \mu_\downarrow)/2}$$

$$\hat{\Omega} = \hat{H} - \frac{\mu_\uparrow + \mu_\downarrow}{2} (\hat{N}_\uparrow + \hat{N}_\downarrow) - \frac{\mu_\uparrow - \mu_\downarrow}{2} (\hat{N}_\uparrow - \hat{N}_\downarrow)$$

Universal phase diagram of the homogeneous unitary system (2)

$$\left. \begin{aligned} \Omega &= -PV \\ dP &= \sum_{\sigma=\uparrow\downarrow} n_{\sigma} d\mu_{\sigma} \end{aligned} \right\} \Rightarrow \text{Just need to know } n(\mu)$$

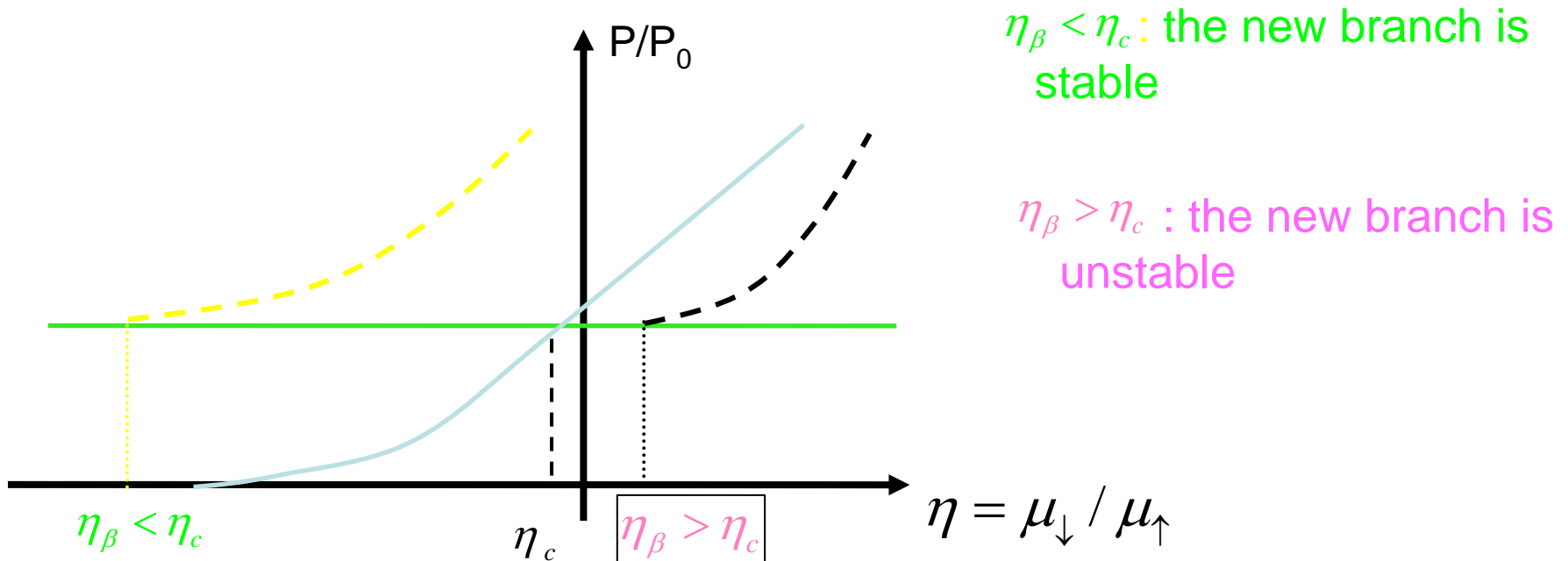


Theoretical evidence for an intermediate phase

General properties of a mixed branch?

Step 1: calculate the energy E of a single impurity atom immersed in a Fermi sea ($E = \mu_{\downarrow}$, with $n_{\downarrow} = 0^+$)

Step 2: $dP/d\mu_{\sigma} = n > 0$



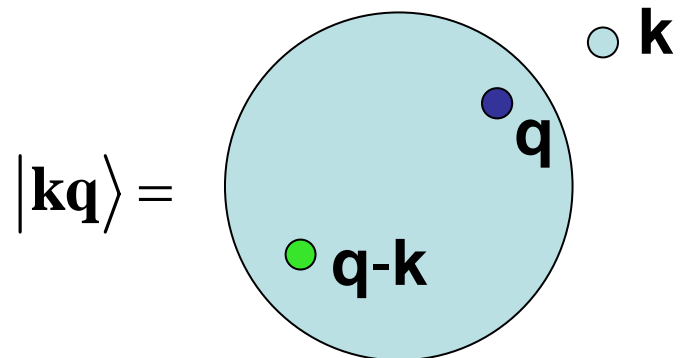
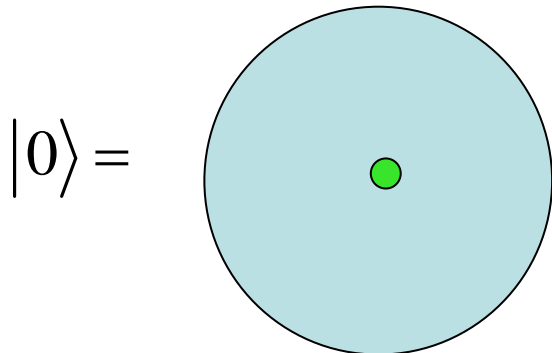
Variational upper bound for the N+1 body problem

$$\hat{H} = \sum_{\mathbf{k}\sigma} \varepsilon_{\mathbf{k}} \hat{a}_{\mathbf{k}\sigma}^\dagger \hat{a}_{\mathbf{k}\sigma} + \frac{g_b}{\Omega} \sum_{\mathbf{k}, \mathbf{k}', \mathbf{q}} \hat{a}_{\mathbf{k}'-\mathbf{q}, \uparrow}^\dagger \hat{a}_{\mathbf{k}+\mathbf{q}, \downarrow}^\dagger \hat{a}_{\mathbf{k}, \downarrow} \hat{a}_{\mathbf{k}', \uparrow}$$

$$\varepsilon_{\mathbf{k}} = \frac{\hbar^2 k^2}{2m} \quad \frac{1}{g_b} = \frac{m}{4\pi\hbar^2 a} - \frac{1}{\Omega} \sum_{\mathbf{k}} \frac{1}{2\varepsilon_{\mathbf{k}}}$$

One impurity: restrict the effect of interactions to the formation of a *single particle-hole pair*.

$$|\Psi\rangle = \varphi_0 |0\rangle + \sum_{\mathbf{k}, \mathbf{q}} \varphi_{\mathbf{k}, \mathbf{q}} |\mathbf{k}, \mathbf{q}\rangle$$



Comparison with Monte-Carlo simulations (C. Lobo et al. PRL. **97**, 200403 (2006))

Minimization of H with respect to φ_0 and φ_{kq}

$$E = \frac{1}{V} \sum_{q < k_F} \frac{m}{4\pi\hbar^2 a} - \frac{1}{V} \sum_{k < k_F} \frac{1}{2\varepsilon_k} + \sum_{k > k_F} \left(\frac{1}{E - (\varepsilon_k + \varepsilon_{q-k} - \varepsilon_q)} \right) \frac{1}{2\varepsilon_k} \quad \varepsilon_k = \frac{\hbar^2 k^2}{2m}$$

For $a = \infty$, $E = -0.606 E_F^*$ $\Downarrow \eta_\beta < -0.606 < \eta_c \sim -0.1$

Monte Carlo simulations : $\eta_\beta = -0.58(1)$

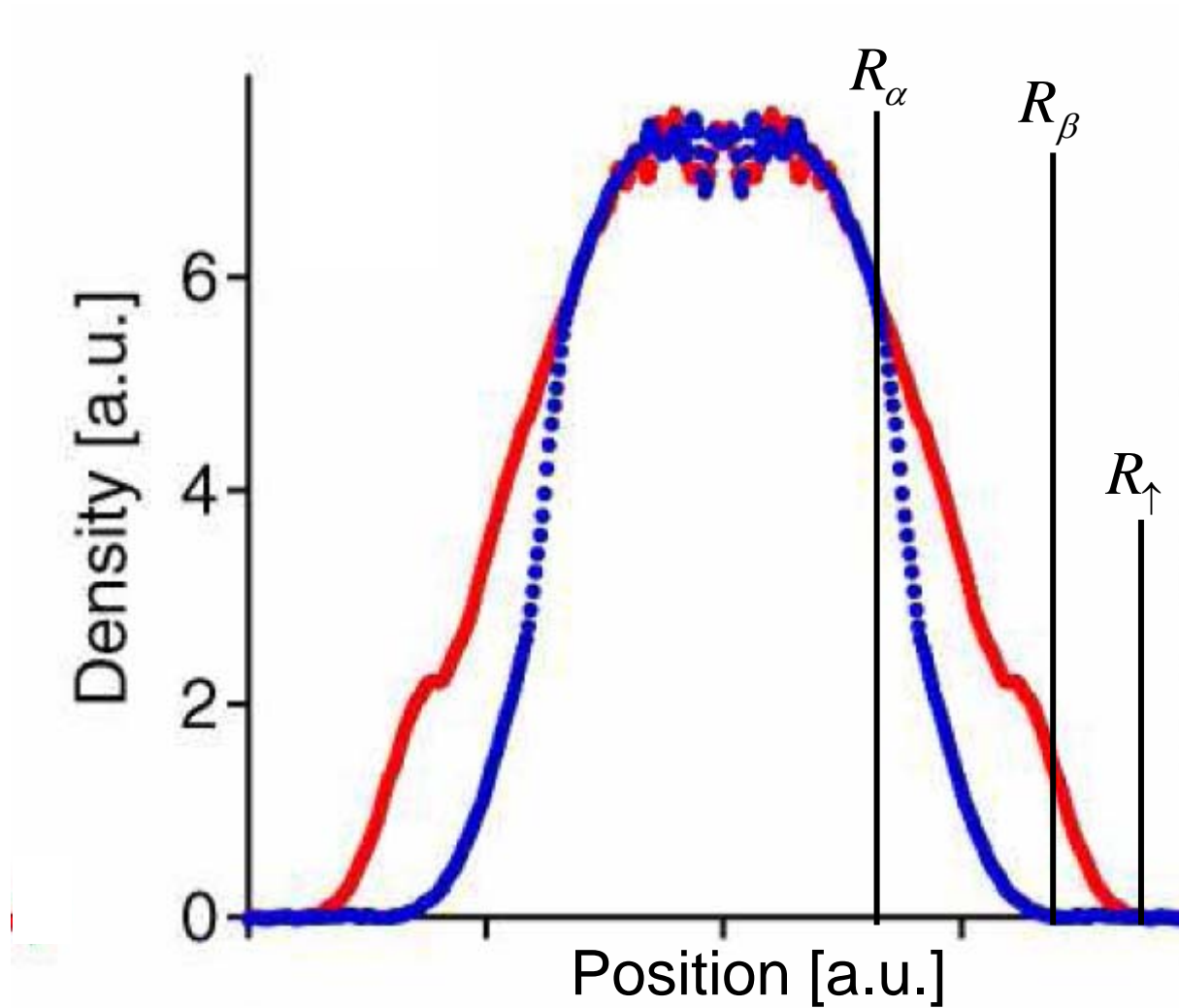
BCS theory: $\eta_\beta = 0$

Why such a good agreement?

Weak excitation of the pairs,
even at unitarity

$$\sum_{k,q} |\varphi_{k,q}|^2 = 0.2$$

Three phase mixture in a trap: MIT



Interpretation of MIT's experiment

Local density approximation in harmonic trap $V(r) \sim r^2$

$$\mu_{\uparrow}(\mathbf{r}) = \mu_{\uparrow}^0 - V(\mathbf{r})$$

$$\mu_{\downarrow}(\mathbf{r}) = \mu_{\downarrow}^0 - V(\mathbf{r})$$

R_{\uparrow} : outer radius of the majority component $\mu_{\uparrow}(R_{\uparrow}) = 0$

R_{β} : outer radius of the minority component $\mu_{\downarrow}(R_{\beta}) / \mu_{\uparrow}(R_{\beta}) = \eta_{\beta}$

R_{α} : radius of the superfluid core $\mu_{\downarrow}(R_{\alpha}) / \mu_{\uparrow}(R_{\alpha}) = \eta_{\alpha}$

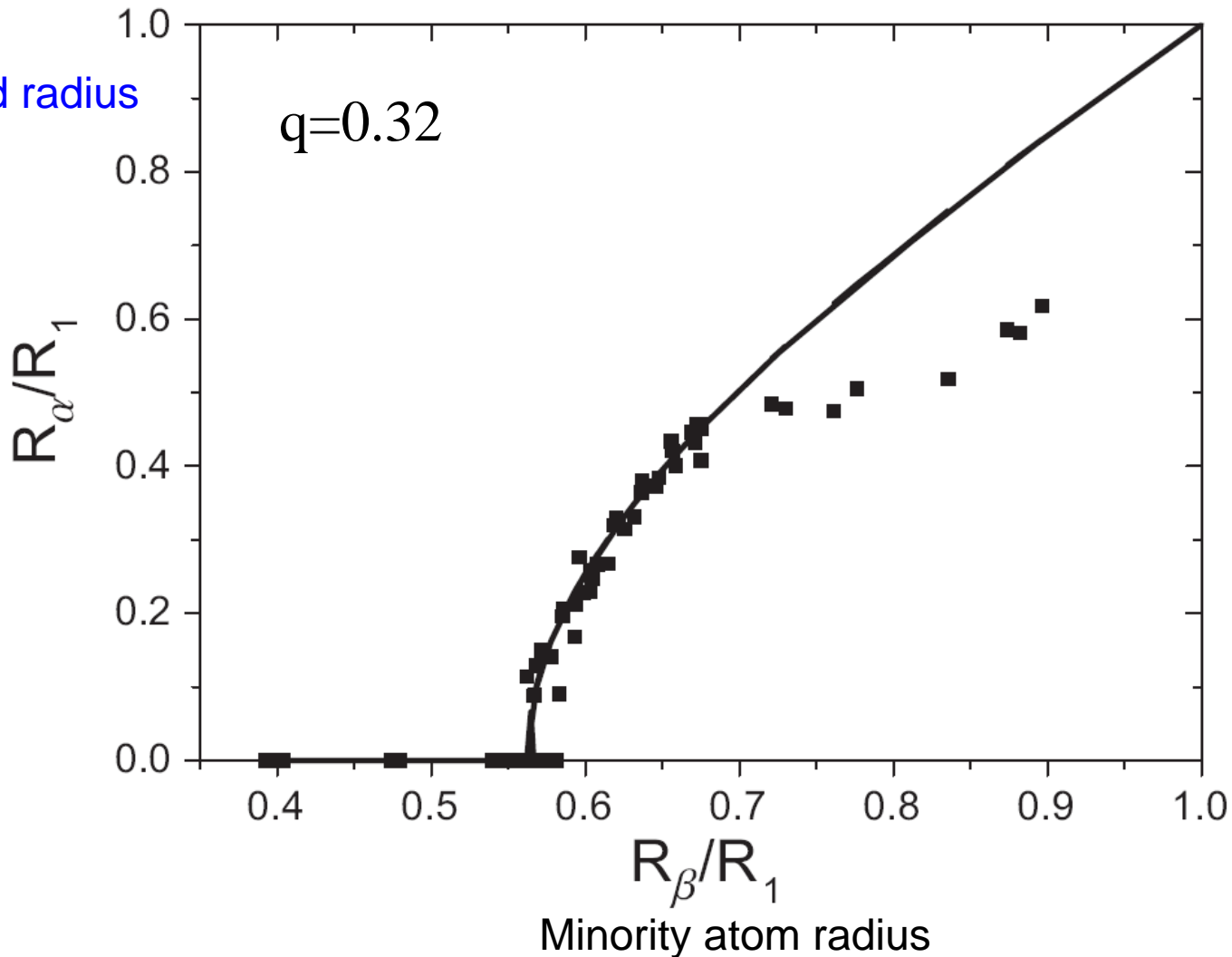
$$\frac{R_{\alpha}}{R_{\uparrow}} = \sqrt{\frac{(R_{\beta} / R_{\uparrow})^2 - q}{1 - q}}$$

$$q = \frac{\eta_{\alpha} - \eta_{\beta}}{1 - \eta_{\beta}}$$

Comparison with experimental data

(M. Zwierlein et al., Nature, 442,54 (2006))

Superfluid radius



Improved bounds for η_α and η_β

$$q=0.32$$

$$\eta_\alpha > -0.10$$

$$\eta_\beta < -0.60$$

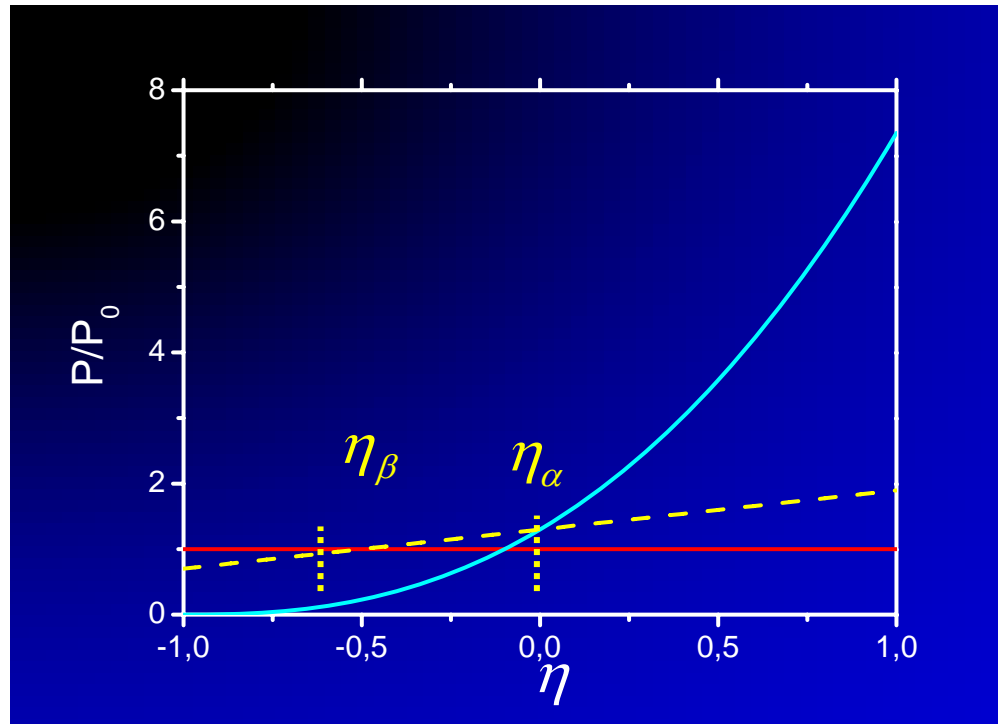
$$-0.62 < \eta_\beta < -0.60$$

$$-0.10 < \eta_\alpha < -0.088$$

$$\eta = \frac{\mu_\downarrow}{\mu_\uparrow}$$

BCS: $\eta_\alpha = 0.1$
 $\eta_\beta = 0$

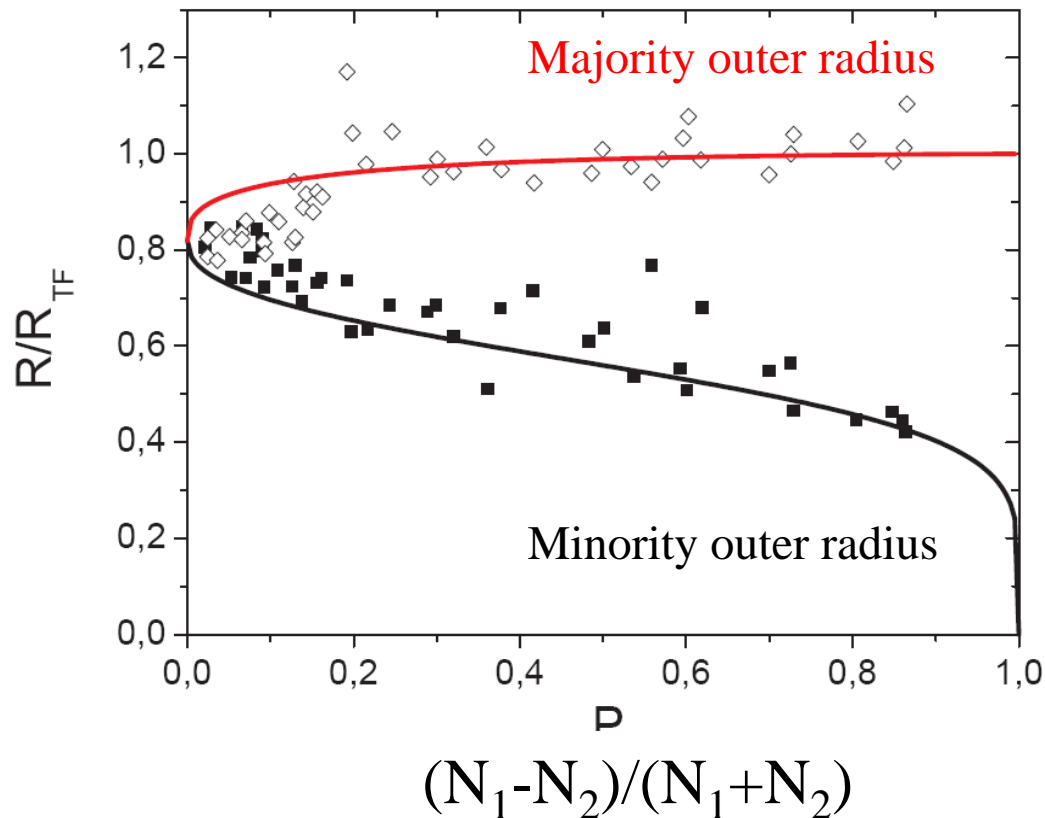
Convexity of P:
 real equation of state below



Unresolved mystery: Rice expt

Partridge et al. Science **311**, 503 (2006)

Experiment by Rice's group: fully compatible with 2 phase scenario (no intermediate phase+LDA). No adjustable parameters.



For $P=0.7$, $q \sim 0.16$ differs from MIT and contradicts theoretical bounds: $q > 0.31$ set by $\eta_\alpha > -0.1$ and $\eta_\beta < 0.60$

Summary: part 1

What have we demonstrated?

3 homogeneous phases in the phase diagram of imbalanced unitary Fermi gases

LDA valid ↻ agreement with MIT

But a host of unanswered questions!

Nature of the MIT/Rice's discrepancy (surface tension ?)

Microscopic nature of the intermediate phase?

Superfluid nature of the intermediate phase?

Dynamical properties, collective modes?

Extension to the BEC-BCS crossover (*in progress*)

Response to RF excitation (MIT experiment, *Schunck et al. cond-mat/0702066*)

Part 2. Measuring one-particle excitations In Fermi gases using Raman spectroscopy

T-L Dao, A. Georges, J. Dalibard, C. Salomon, I. Carusotto, Cond-mat/0611206

In Fermi liquid theory: low energy excitations are build out of quasiparticles
Dispersion relation on a given point of the Fermi surface:

$$\xi_k \sim \mathbf{v}_F(k_F) \cdot (\mathbf{k} - \mathbf{k}_F) + \dots \quad \text{Lifetime: } \Gamma_k^{-1}$$

Fermi surface: excitation energy vanishes: $\xi_{k_F} = 0$

Normal phase in Cuprate SC show strong deviations, anisotropic behavior.

Probe directly one particle correlator: $\langle \psi^\dagger(r, t) \psi(r', t') \rangle$

Bragg spectroscopy or noise corr. : $\langle \psi^\dagger(r, t) \psi(r, t) \psi^\dagger(r', t') \psi(r', t') \rangle$

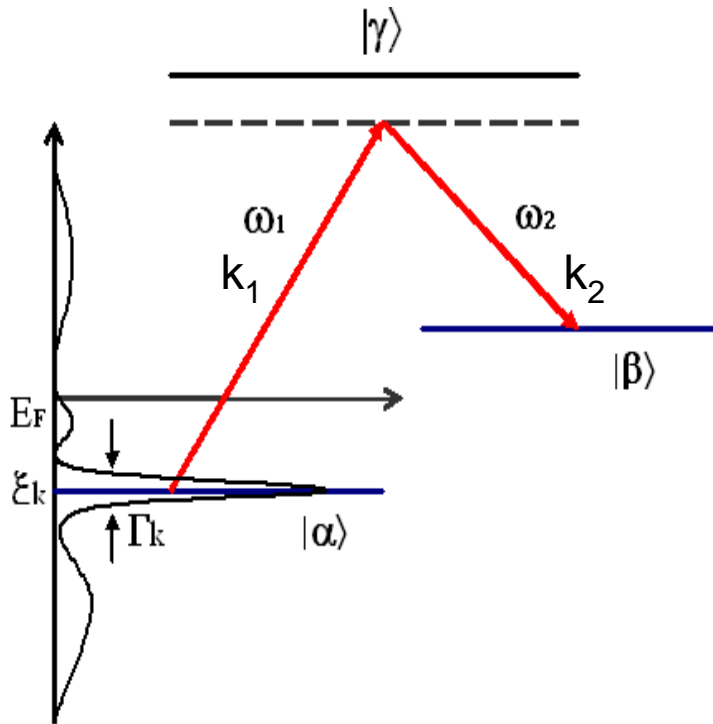
Stimulated Raman spectroscopy: widely used for Bose systems

Very interesting for fermions:

Probing Fermi surface of strongly interacting fermions (Time of Flight not adequate)

Momentum-resolved quasiparticles excitations

Two-photon Raman excitation



Interacting fermions: $|\alpha\rangle, |\alpha'\rangle$
 For instance ${}^6\text{Li}$ near F. Resonance

Third state empty: $|\beta\rangle$

No interaction with $|\alpha\rangle, |\alpha'\rangle$

Similar to ARPES in cond. Matter.

$$R(\mathbf{q}, \Omega) \sim \int_{-\infty}^{+\infty} dt \int d\mathbf{r} d\mathbf{r}' e^{i[\Omega t - \mathbf{q} \cdot (\mathbf{r} - \mathbf{r}')] } \times g_{\beta}(\mathbf{r}, \mathbf{r}'; t) \langle \psi_{\alpha}^{\dagger}(\mathbf{r}, t) \psi_{\alpha}(\mathbf{r}', 0) \rangle$$

Selectivity in $\mathbf{q} = \mathbf{k}_1 - \mathbf{k}_2$
 Selectivity in energy

Interacting Fermions in homogeneous 2D square Lattice

Raman resonance condition $\varepsilon_{\mathbf{k}+\mathbf{q},\beta} - \xi_{\mathbf{k}} = \Omega$

Threshold in Ω : $(\omega_1 - \omega_2)_T = \varepsilon_{\beta}^0 - \mu \sim \varepsilon_{\beta}^0 - \varepsilon_{\alpha}^0$

Ω close to threshold for $\mathbf{q} = -\mathbf{k}_F$

Shell surrounding Fermi Surface

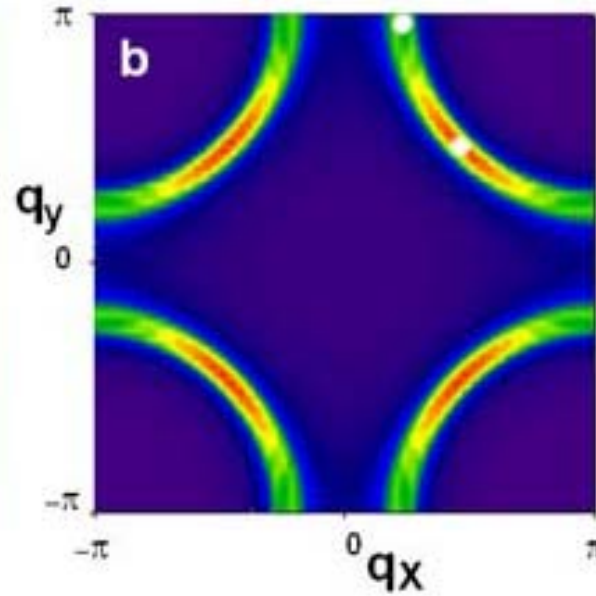
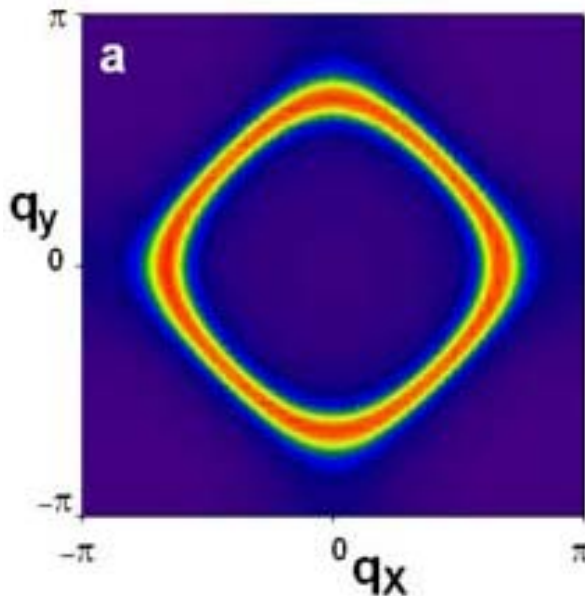
of width $\sqrt{2M\Delta\Omega} \sim \sqrt{2M\mathbf{v}_F(\mathbf{k}_F) \cdot (\mathbf{q} + \mathbf{k}_F)}$

Non interacting fermions

$$\xi_{\mathbf{k}} = -2t_{\alpha}(\cos k_x + \cos k_y) - \mu$$

Model pseudo-gap with d-wave symmetry

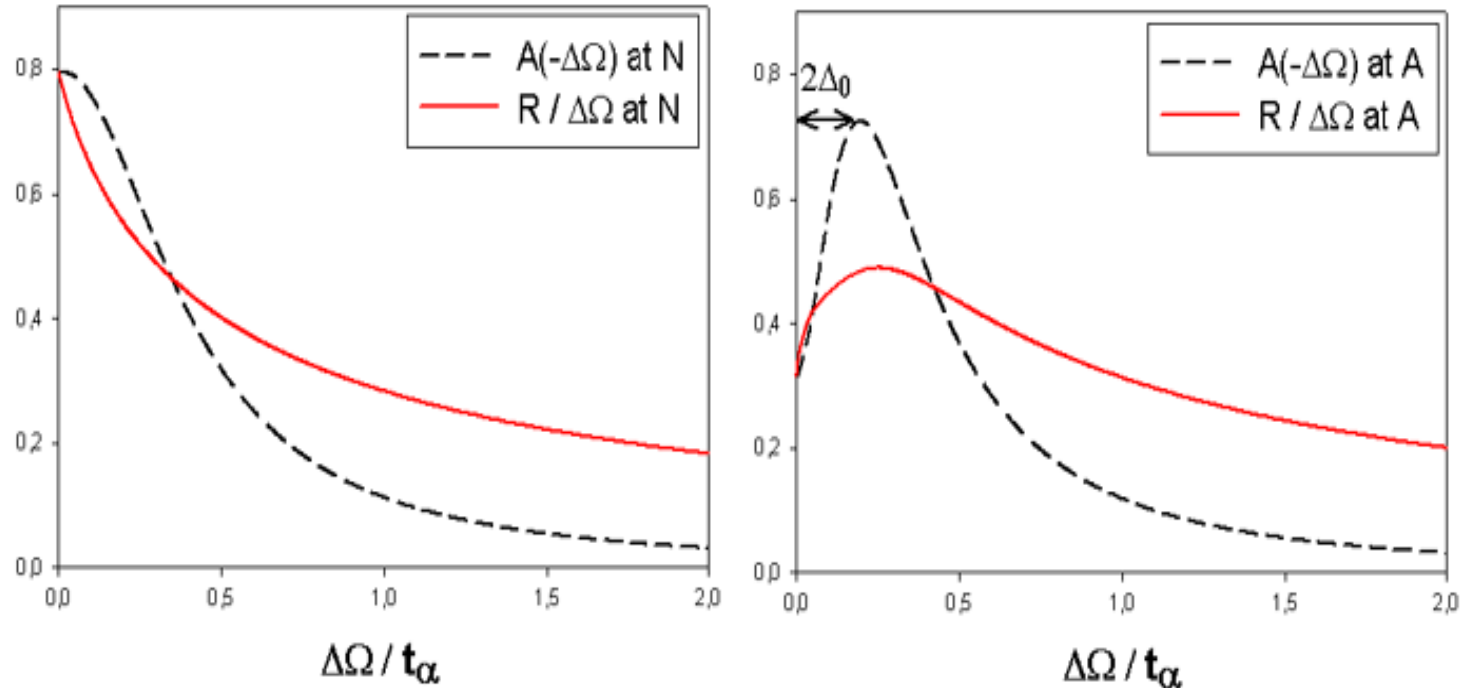
$$\Delta_{\mathbf{k}} = \Delta_0(\cos k_x - \cos k_y)$$



$$\begin{aligned} \Delta_0 &= 0.1 t_{\alpha} \\ N_{\alpha} &= 0.45 \\ \Gamma_0 &= 0.05 t_{\alpha} \\ \Gamma_1 &= 0.4 t_{\alpha} \end{aligned}$$

Lorentzian
spectral
Function
 $A(\mathbf{k}, \Omega)$
uniform in
k-space

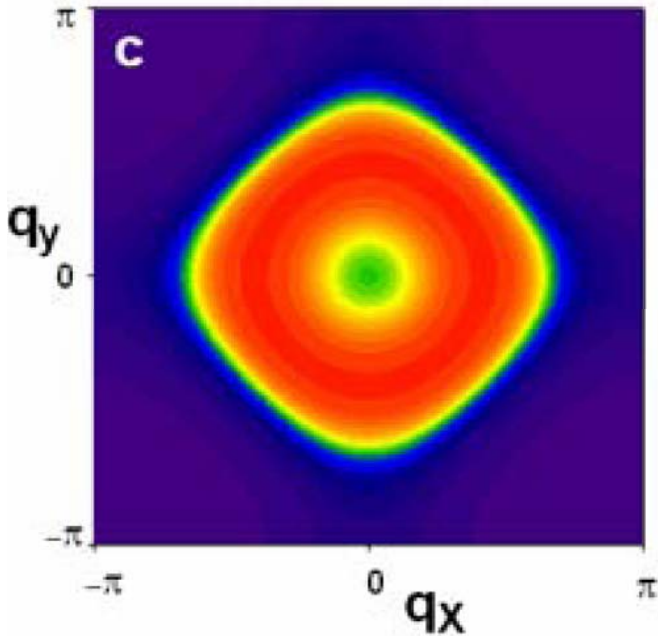
Rate $R(q, \Omega)$ and spectral function $A(k, \Omega)$



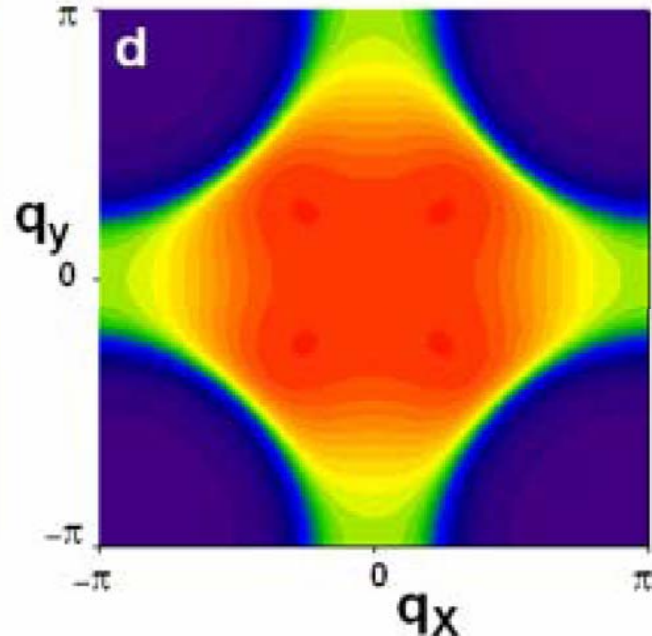
For fixed value of q , chosen near Nodal line or Anti-nodal
Scan Raman detuning

In Harmonic Trap

LDA: $\mu(R) = \mu_0 - M \omega_0^2 R^2 / 2$

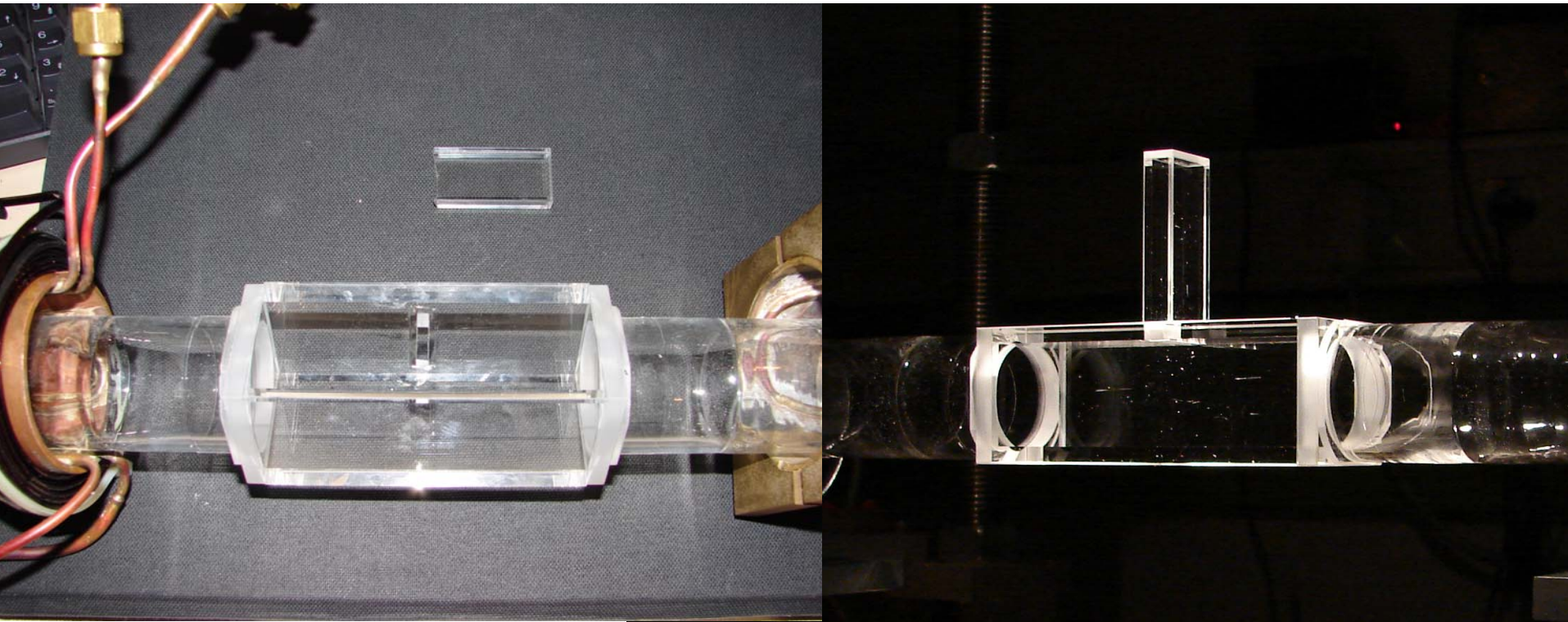


Non-interacting fermions



Model d-wave pseudogap state

New experimental setup



Enlarged glass cell

New laser sources: 120 mW diodes operating at 75 degree C

New Zeeman slower

More stable Ioffe-Pritchard trap

120 Watt far detuned optical trap (Fiber laser)

Access for 3D optical lattice

$3 \cdot 10^{10}$ ${}^7\text{Li}$ atoms in MOT \longrightarrow expected increase of $\times 10$ in ${}^6\text{Li}$ number

Ongoing: Transfer into magnetic trap

Thank you for your attention!

