



strongly interacting Fermi gases: recent results from Innsbruck

Rudolf Grimm

**“Center for Quantum Optics”
Innsbruck**



University



Austrian Academy of Sciences



I: collective oscillations in the BEC-BCS crossover

II. Fermi-Fermi mixture of atoms (^6Li - ^{40}K)

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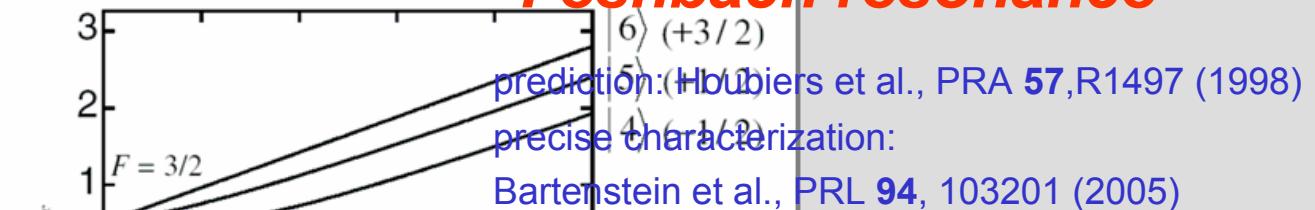
^6Li spin mixture



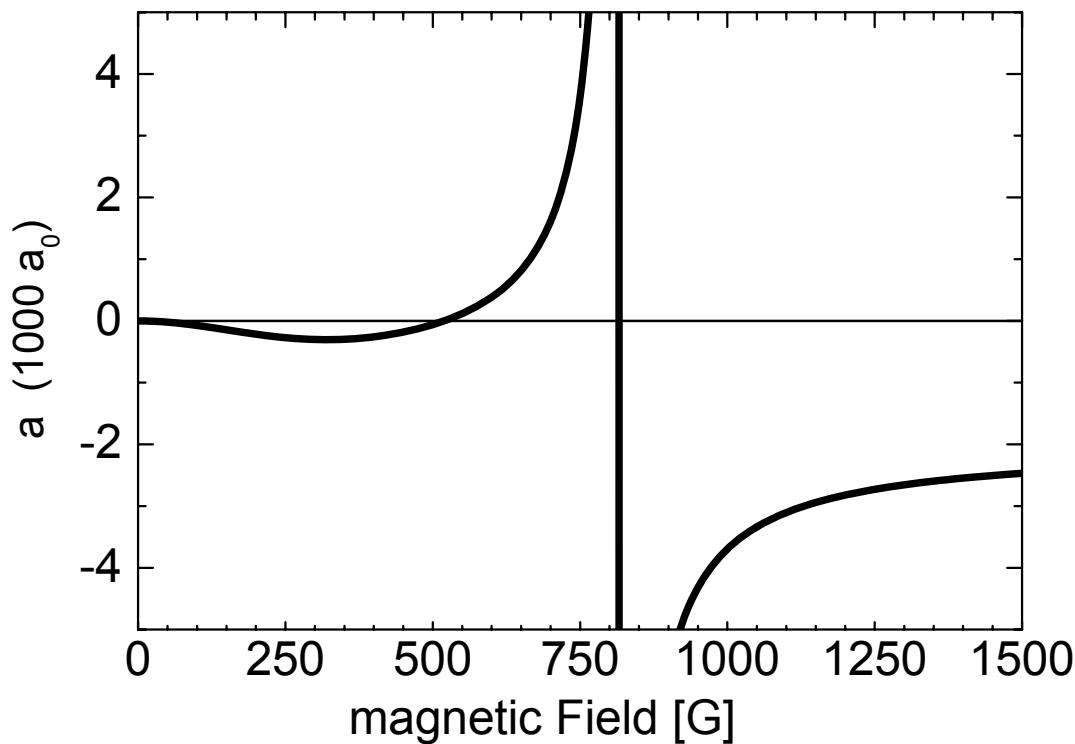
Ultracold atoms
quantum gases

^6Li ground state in a magnetic field

Feshbach resonance

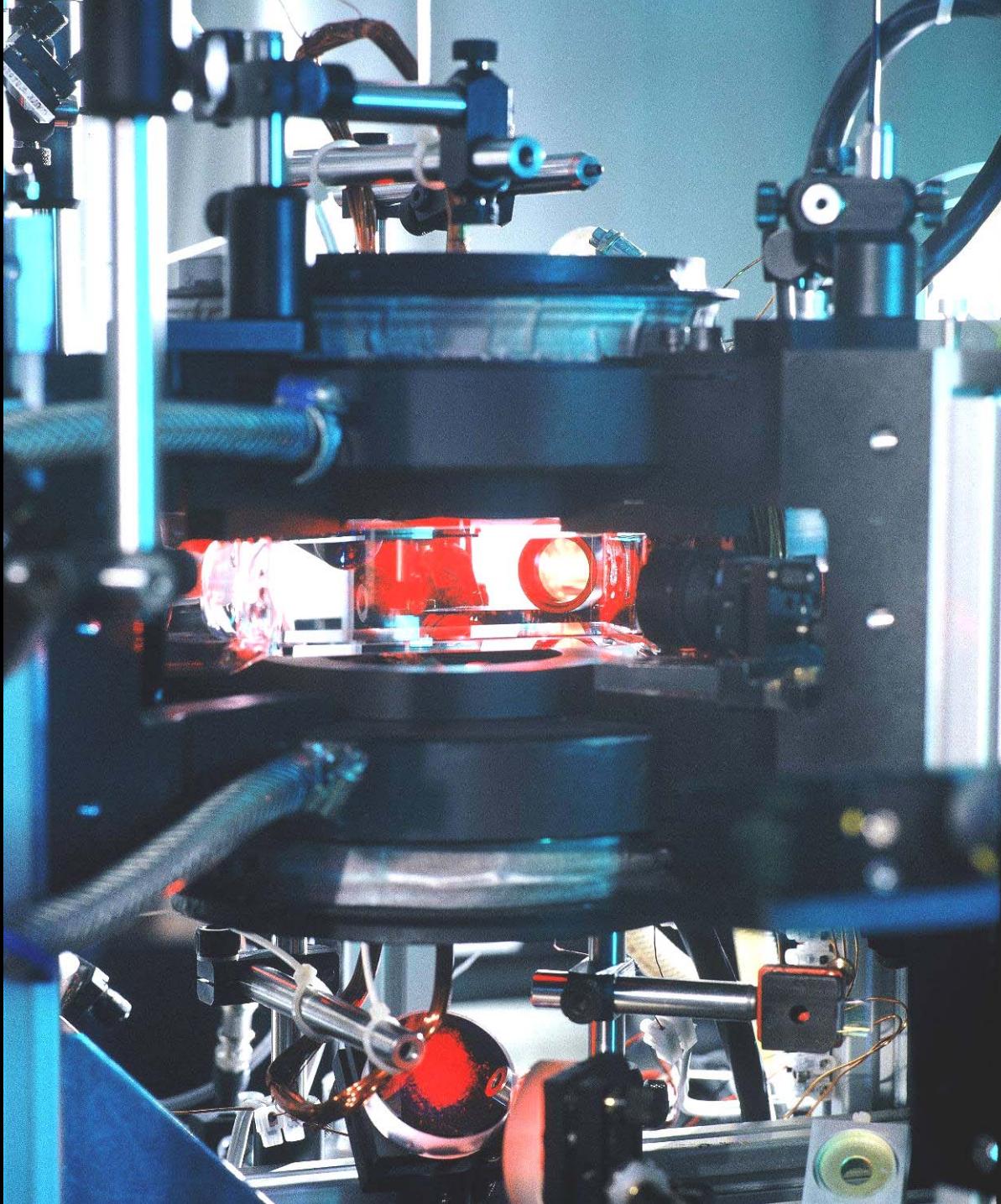


spin mix
two low
stable a
two-bo

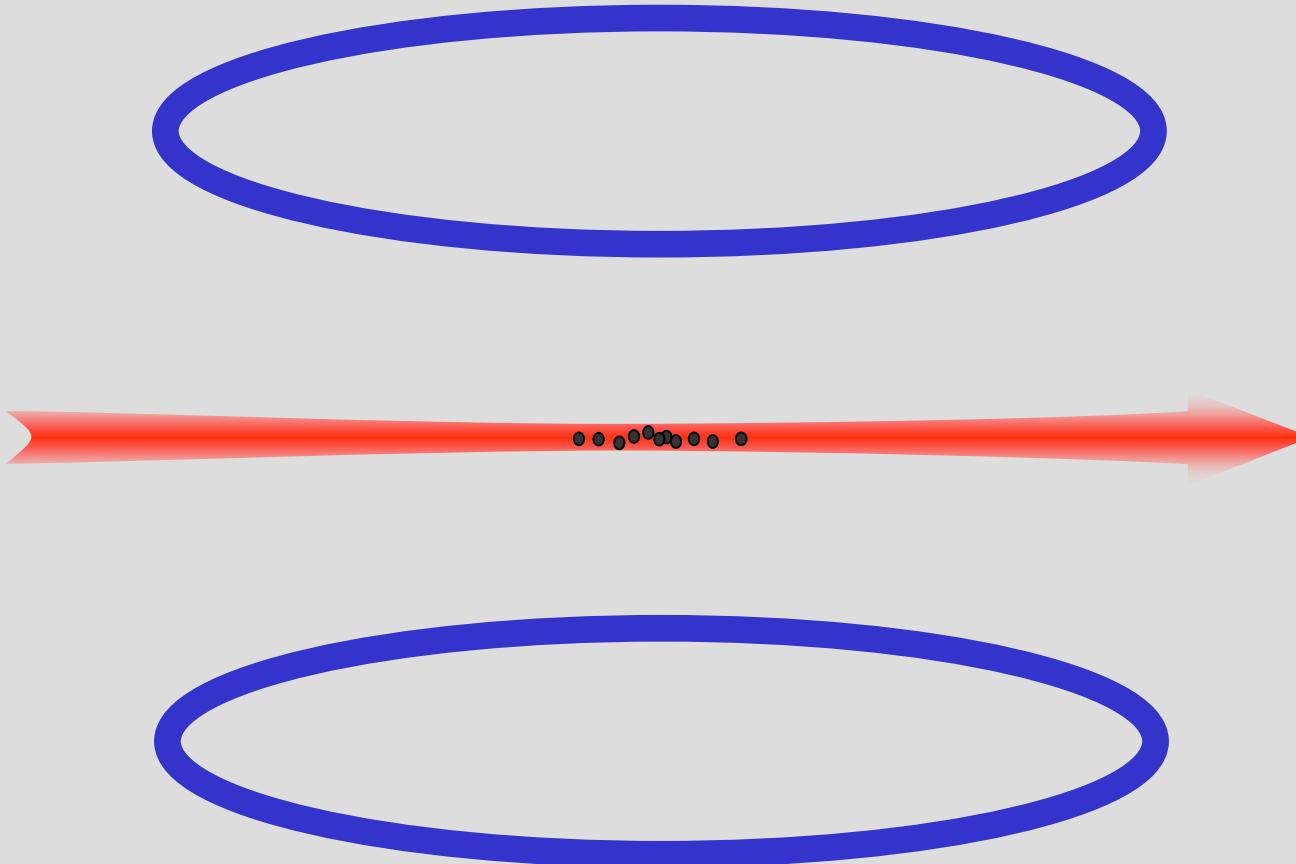


the ${}^6\text{Li}$ team





optical trap for evaporative cooling



collective modes

- Stringari, Europhys. Lett. **65**, 749 (2004)
- Hu, Minguzzi, Liu, Tosi, PRL **93**, 190403 (2004)
- Heiselberg, PRL **93**, 040402 (2004)
- Combescot, Leyronas, Europhys. Lett. **68**, 762 (2005)
- Manini, Salasnich, PRA **71**, 033625 (2005)
- Bulgac, Bertsch, PRL **94**, 070401 (2005)
- Kim, Zubarev, PRA **72**, 011603(R) (2005)
- Astrakharchik, Combescot, Leyronas, Stringari, PRL **95**, 030404 (2005)
- ... *and many more recent papers...*

theory

collective modes: elementary situation

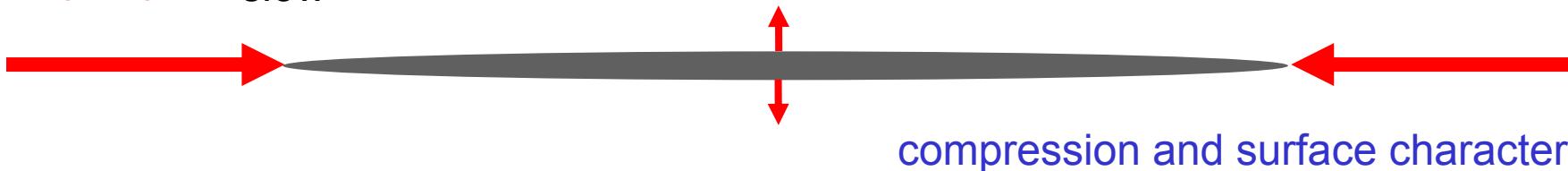


elongated trap limit ($\omega_r \gg \omega_z$ well fulfilled for all our expts.)

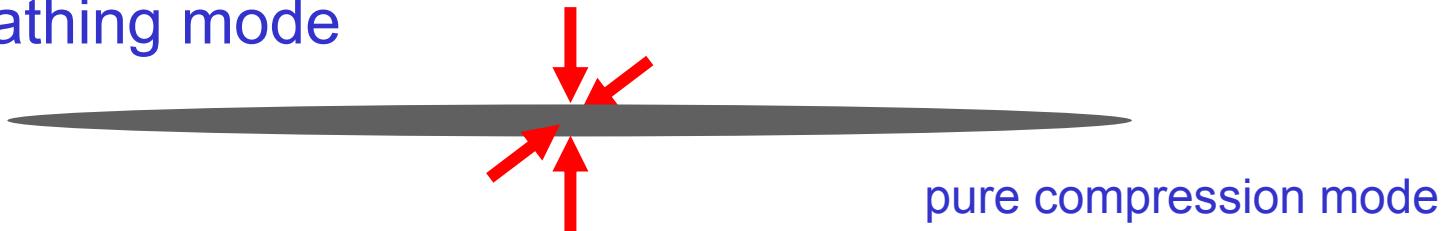
cylindrical symmetry (radial trap freq. ω_r),

3D harmonic potential

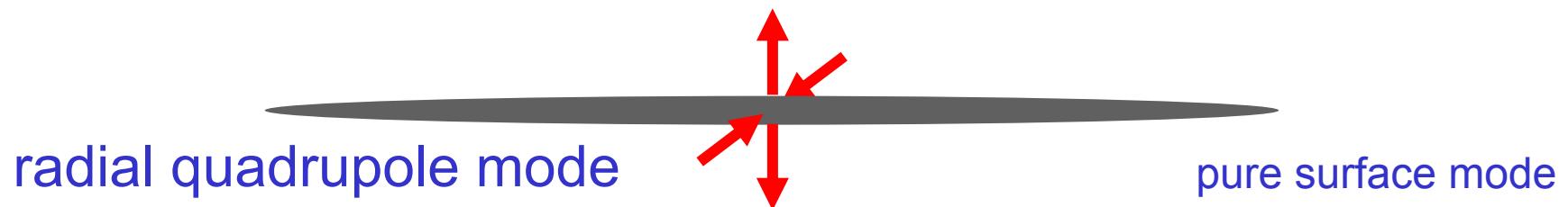
axial slow



radial breathing mode



radial fast



collective mode freq. in BEC-BCS crossover



polytropic index !!

$$\mu \propto n^\gamma$$

equation of state

general	BEC	unitarity (and BCS limit)
frequencies of low-lying collective modes	γ	$\gamma = 1$

axial ω_z	$\sqrt{3 - 1/(\gamma + 1)}$	$\sqrt{5/2}$ $= 1.581$	$\sqrt{12/5}$ $= 1.549$
radial ω_\perp breathing	$\sqrt{2(\gamma + 1)}$	2	$\sqrt{10/3}$ $= 1.826$
radial ω_\perp quadrupole	$\sqrt{2}$	$\sqrt{2}$	$\sqrt{2}$

hydrodynamic behavior

2
2
2
collisionless
oscillation

axial mode

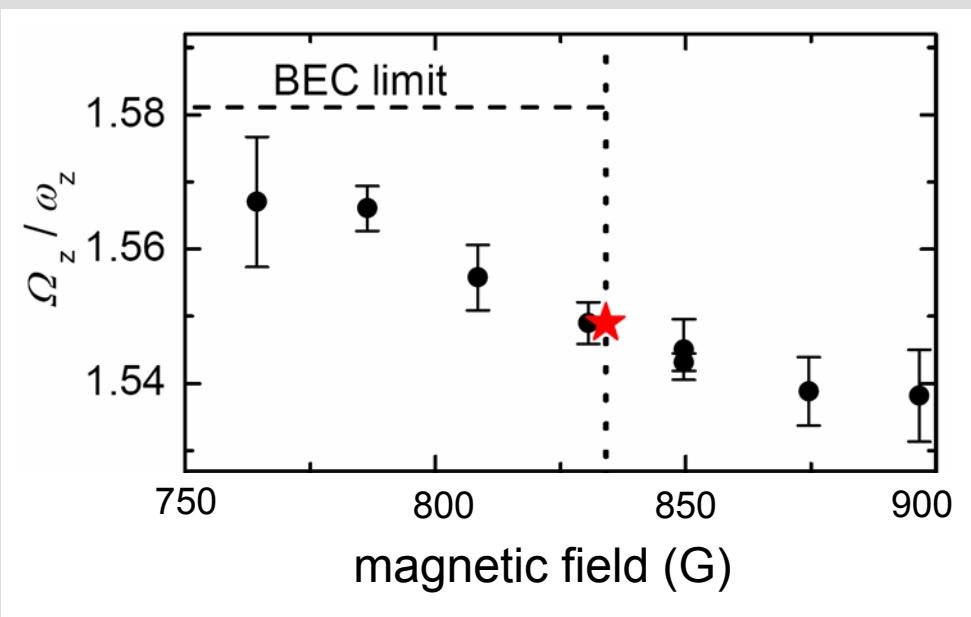


axial slow



compression and surface character

Bartenstein et al., PRL **92**, 203201 (2004)



unitarity

$$\Omega_z / \omega_z = \sqrt{12/5}$$

confirms equation of state

$$\mu \propto n^{2/3}$$

**extremely weak damping
in resonance region !!**

hint on superfluidity

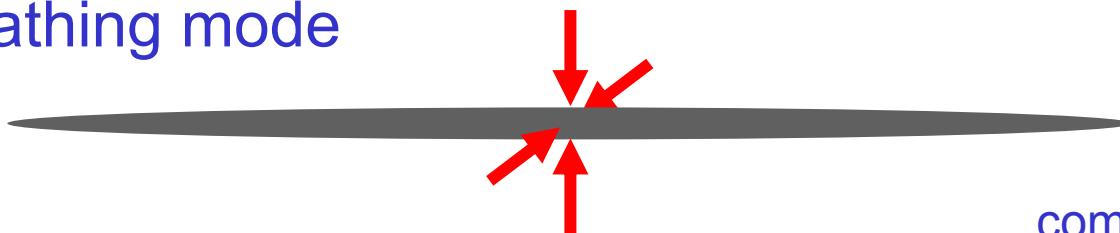
radial compression mode



radial breathing mode

fast

compression mode



Innsbruck

Bartenstein et al., PRL **92**, 203201 (2004)
see also Altmeyer et al., cond-mat/0611285

Altmeyer et al., PRL **98**, 0404401 (2007)

Duke

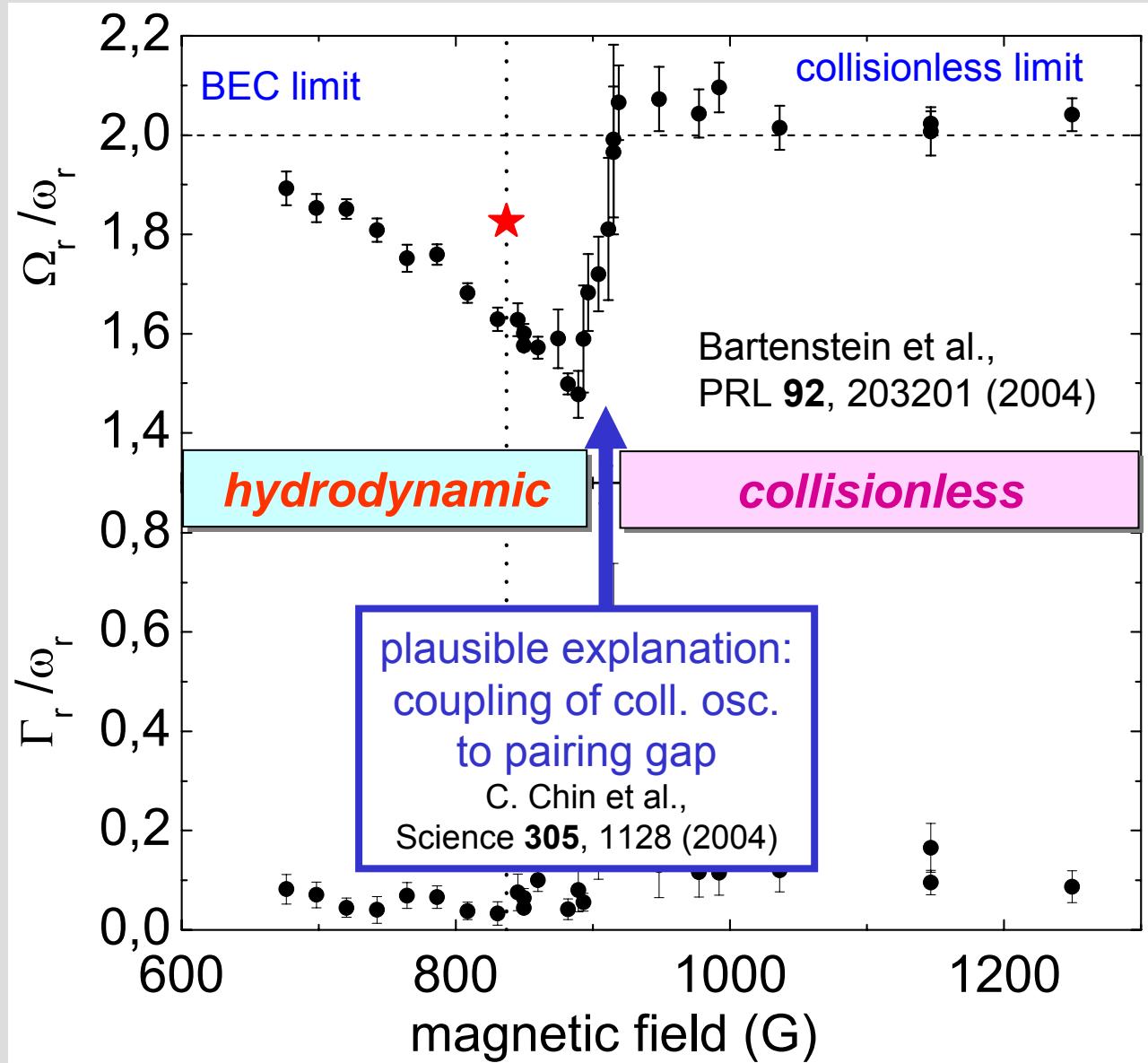
Kinast et al., PRL **92**, 150402 (2004)
Kinast et al. PRA **70**, 051401(R) (2004)
Kinast et al., PRL **94**, 170404 (2005)

radial compression mode



frequency
(normalized
to sloshing
mode)

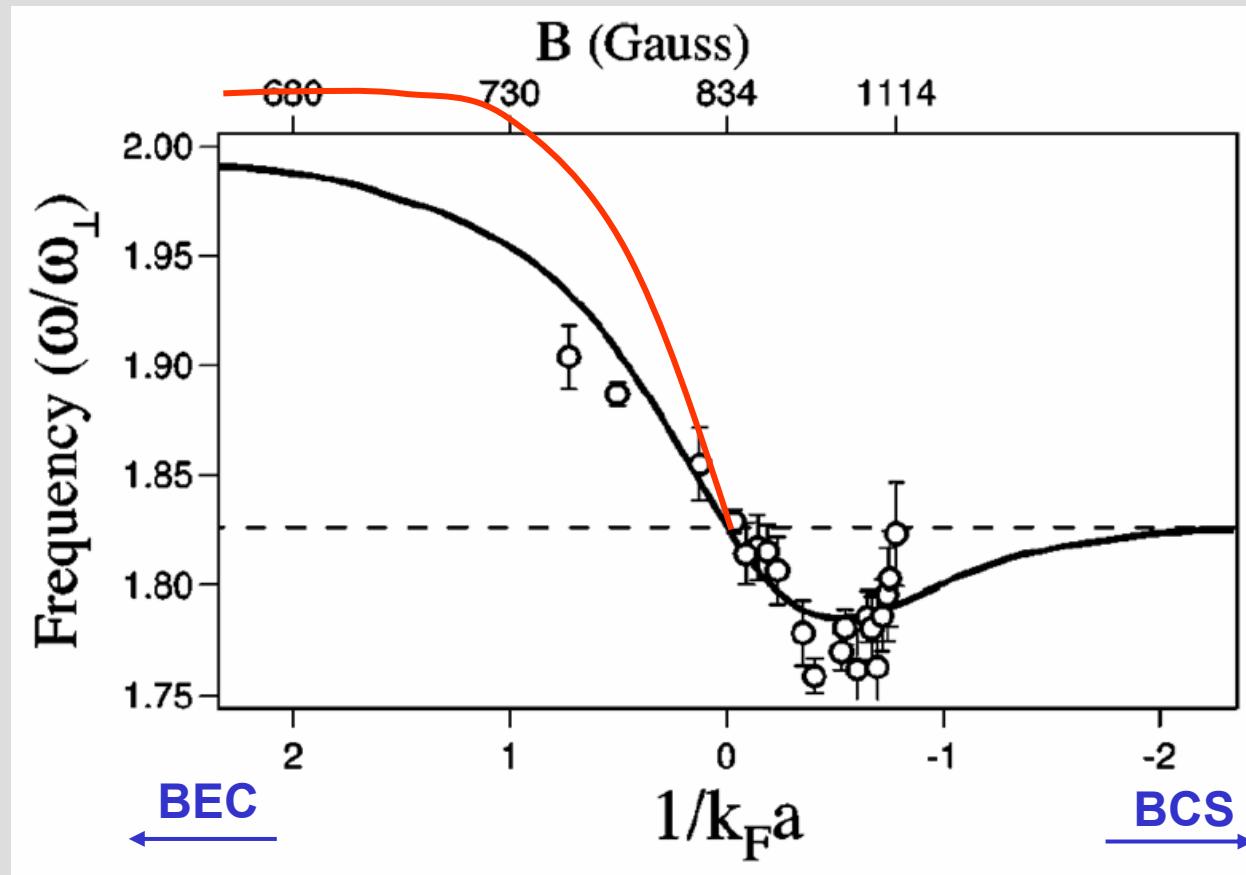
damping



radial mode



exp. data from John Thomas group at Duke Kinast et al, PRA 70, 051401 (2004)



theory: mean field BCS à la Leggett, Nozières & Schmitt-Rink

Hu et al., PRL 93, 190403 (2004)

quantum Monte Carlo, Astrakharchik et al., PRL 95, 030404 (2005)

see also Manini and Salasnich, PRA 71, 033625 (2005)

quantum Monte Carlo



Ultracold atoms
quantum gases

Equation of State of a Fermi Gas in the BEC-BCS Crossover: A Quantum Monte Carlo Study

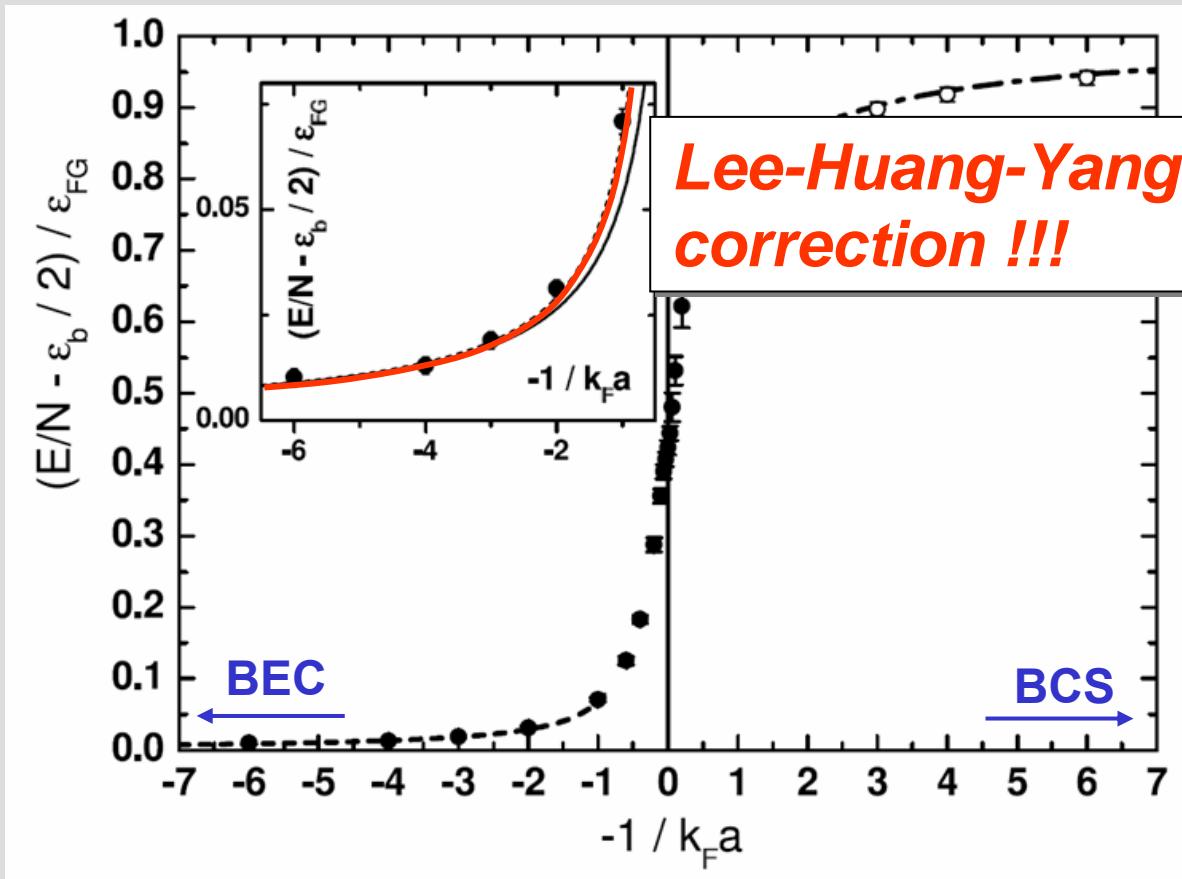
G. E. Astrakharchik,^{1,2} J. Boronat,³ J. Casulleras,³ and S. Giorgini¹

¹Dipartimento di Fisica, Università di Trento and BEC-INFM, I-38050 Povo, Italy

²Institute of Spectroscopy, 142190 Troitsk, Moscow region, Russia

³Departament de Física i Enginyeria Nuclear, Campus Nord B4-B5, Universitat Politècnica de Catalunya, E-08034 Barcelona, Spain

(Received 4 June 2004; published 10 November 2004)





Lee-Huang-Yang correction

Phys. Rev. 105, 1119 (1957)

Many-Body Problem in Quantum Mechanics and Quantum Statistical Mechanics

T. D. LEE, Columbia University, New York, New York

AND

C. N. YANG, Institute for Advanced Study, Princeton, New Jersey

(Received December 10, 1956)

PHYSICAL REVIEW

VOLUME 106, NUMBER 6

JUNE 15, 1957

Eigenvalues and Eigenfunctions of a Bose System of Hard Spheres and Its Low-Temperature Properties

T. D. LEE, Columbia University, New York, New York

AND

KERSON HUANG AND C. N. YANG, Institute for Advanced Study, Princeton, New Jersey

(Received March 19, 1957)

It is shown that the pseudopotential method can be used for an explicit calculation of the first few terms in an expansion in power of $(\rho a^3)^{\frac{1}{2}}$ of the eigenvalues and the corresponding eigenfunctions of a system of Bose particles with hard-sphere interaction. The low-temperature properties of the system are discussed.

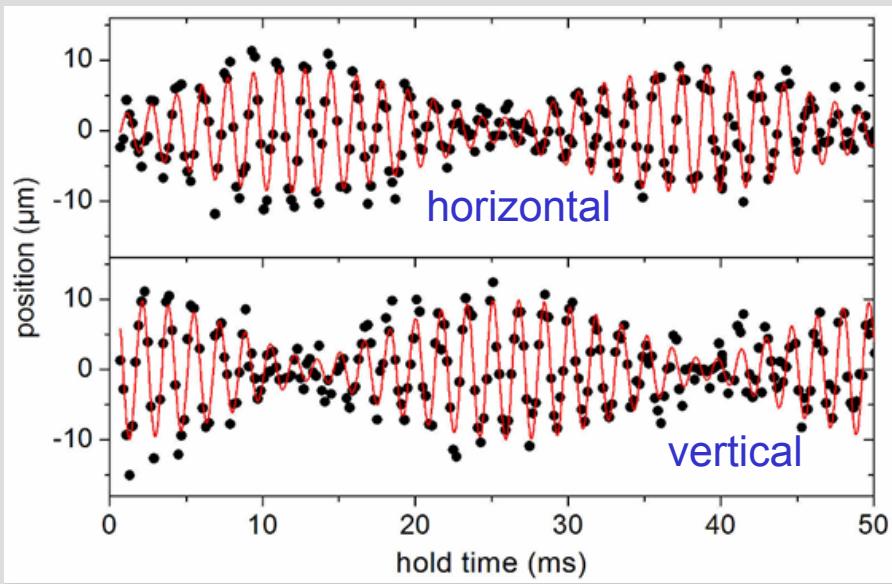
(E) The ground-state energy per particle for a Boltzmann gas and for a Bose gas at a finite density and infinite volume is

$$E/N = 4\pi a\rho \left[1 + 128(\rho a^3)^{\frac{1}{2}} / 15\pi^{\frac{1}{2}} + O(\rho a^3) \right]. \quad (7)$$

leading correction is positive

→ upshift of collective-mode frequency in mBEC regime !

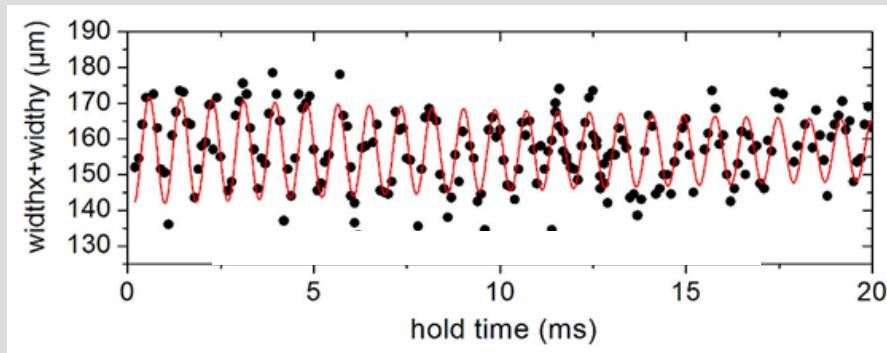
precision measurements



sloshing modes

beat reveals
trap ellipticity of $\sim 6\%$

can be measured with
 $\sim 10^{-3}$ uncertainty



compression mode

accurate determination
of frequency needs
very low damping
→ optimized cooling !

anharmonicity effects
in Gaussian trap potential ~2%

suppressed to few 10^{-3}
by normalization to sloshing mode

precision measurements



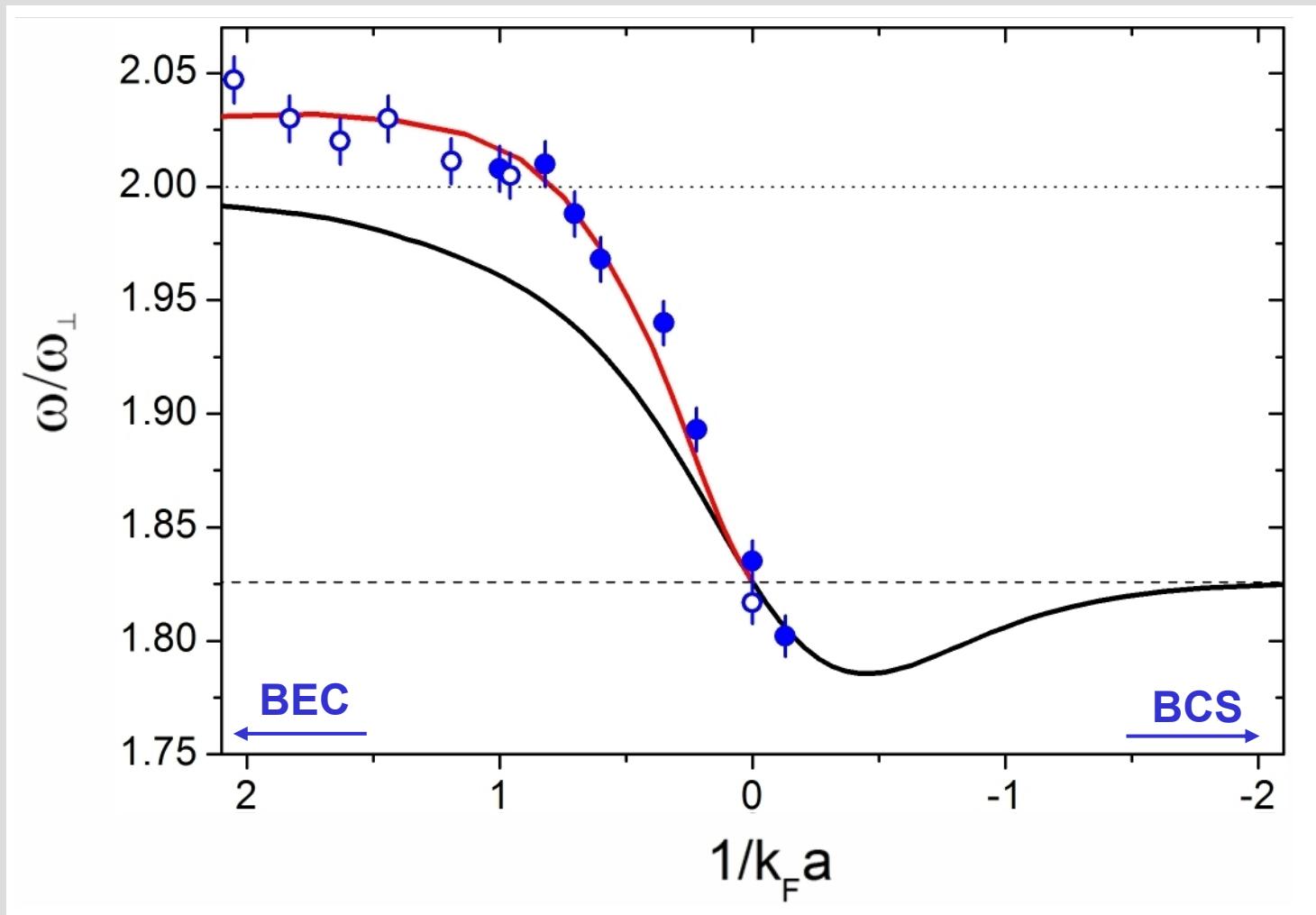
B (G)	Sloshing			Compression		Correction	
	$1/k_F a$	$\omega_\perp/2\pi$ (Hz)	ϵ	$\omega_c/2\pi$ (Hz)	γ/ω_\perp	$\kappa\epsilon^2$ (10^{-4})	$b\alpha$
727.8	2.21	292.7(5)	0.083(3)	596.3(6)	0.007(2)	48	20
735.1	1.96	298.6(5)	0.091(3)	602.8(8)	0.008(3)	60	26
742.5	1.75	294.5(5)	0.067(3)	593.2(7)	0.005(2)	63	28
749.8	1.55	296.3(4)	0.073(3)	599.0(7)	0.006(2)	38	28
760.9	1.27	296.0(4)	0.088(2)	592.3(7)	0.009(2)	58	24
771.9	1.03	293.6(7)	0.074(5)	586.2(6)	0.007(3)	41	27
834.1	0	287.5(7)	0.073(5)	559.4(9)	0.014(3)	55	94
757.2	1.07	605.0(9)	0.065(3)	1210.9(12)	0.010(2)	32	13
768.2	0.87	592.5(7)	0.060(2)	1186.6(12)	0.012(2)	36	16
775.6	0.75	590.2(4)	0.060(1)	1170.2(21)	0.007(4)	28	14
782.2	0.64	604.8(9)	0.061(3)	1187.1(16)	0.006(3)	29	16
801.3	0.28	586.8(7)	0.063(2)	1135.2(12)	0.010(2)	33	24
812.3	0.24	586.5(7)	0.058(2)	1106.9(16)	0.014(3)	30	33
834.1	0	596.3(9)	0.070(3)	1089.0(12)	0.010(2)	48	40
849.1	-0.14	583.2(7)	0.052(2)	1046.7(37)	0.007(2)	29	47

this took us a year!

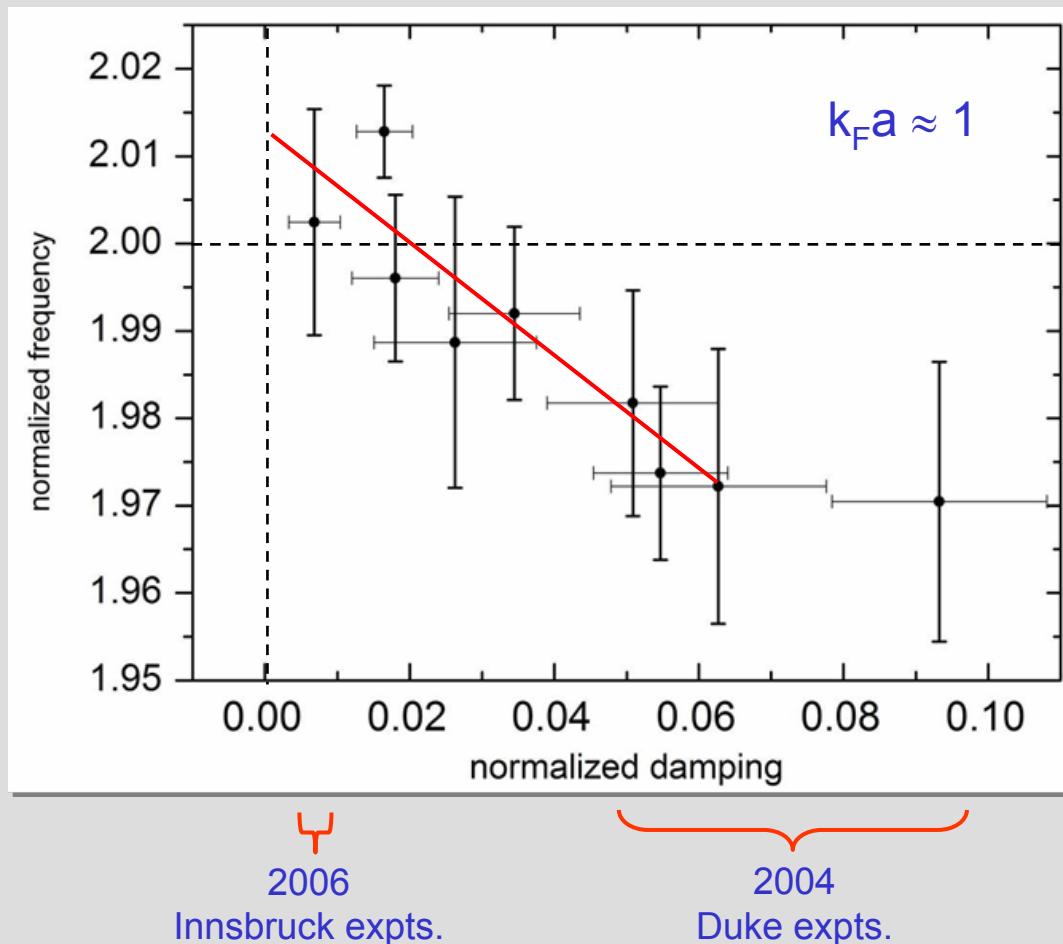
precisions measurements



Altmeyer et al., PRL 98, 0404401 (2007)



temperature shift



temperature shifts explain
previous "agreement" of expt. data with (too simple) BCS mean field theory

measurements on radial breathing mode

confirm

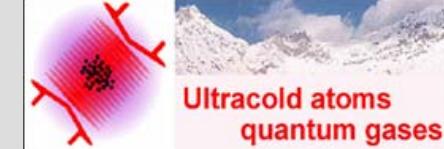
**equation of state from quantum Monte-Carlo calculations
(including LHY correction)**

rule out

simple mean-field BCS theory for crossover

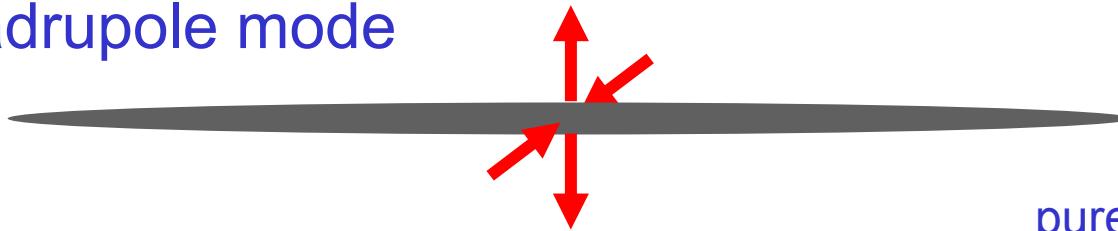
precision test of many-body theories !

radial quadrupole mode



Altmeyer et al., manuscript in preparation
(should appear on the arXiv this week)

radial quadrupole mode



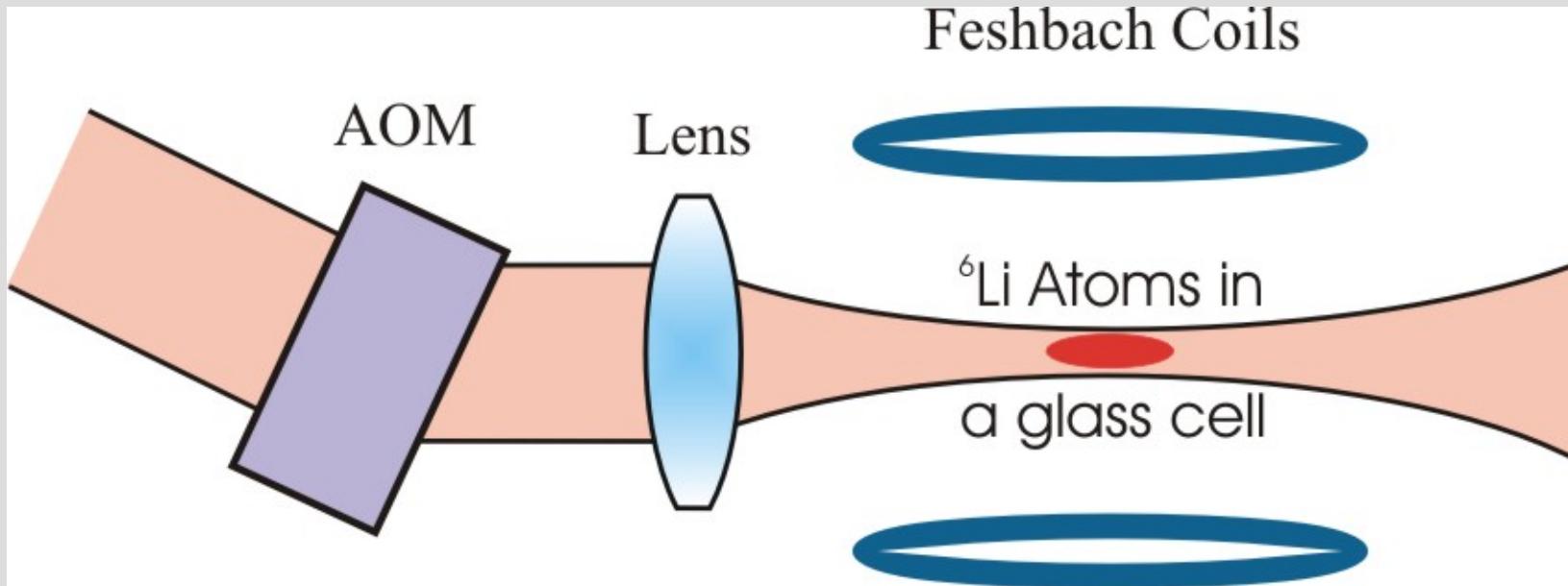
pure surface mode

pure test of hydrodynamics !

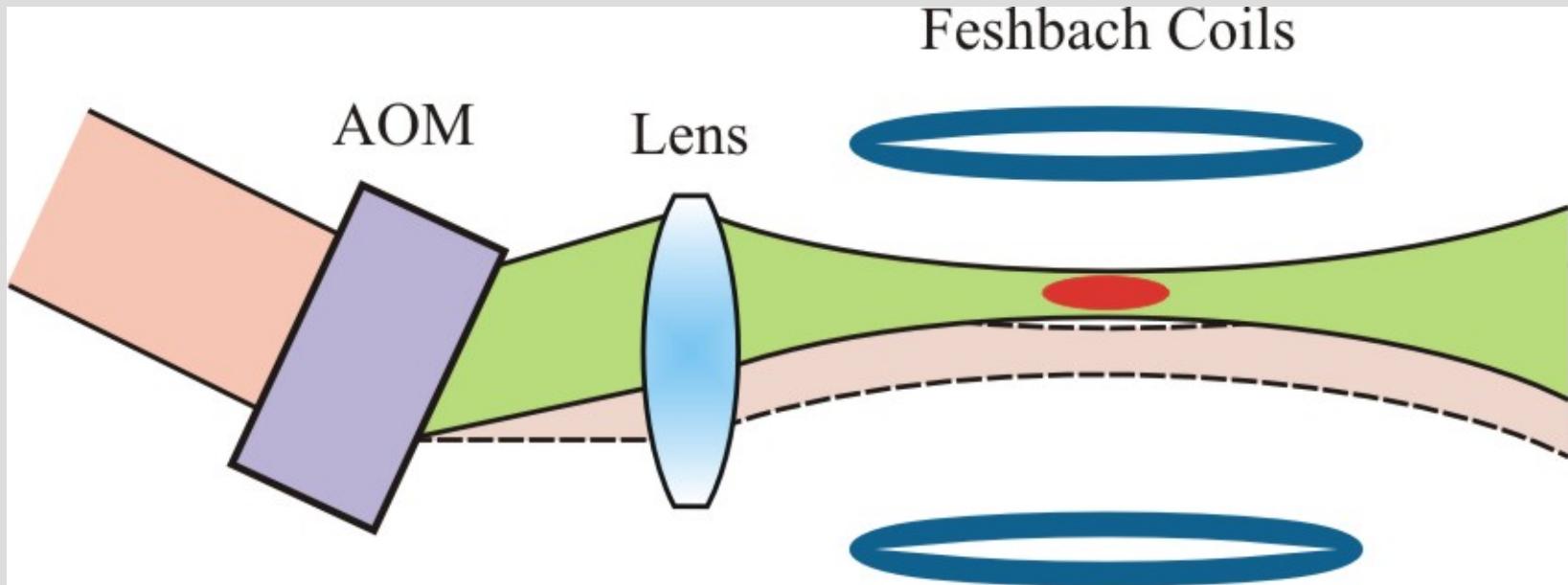
(mode freq. independent of eq. of state)

hydrodynamic freq. $\sqrt{2}\omega_r$ collisionless freq. $2\omega_r$

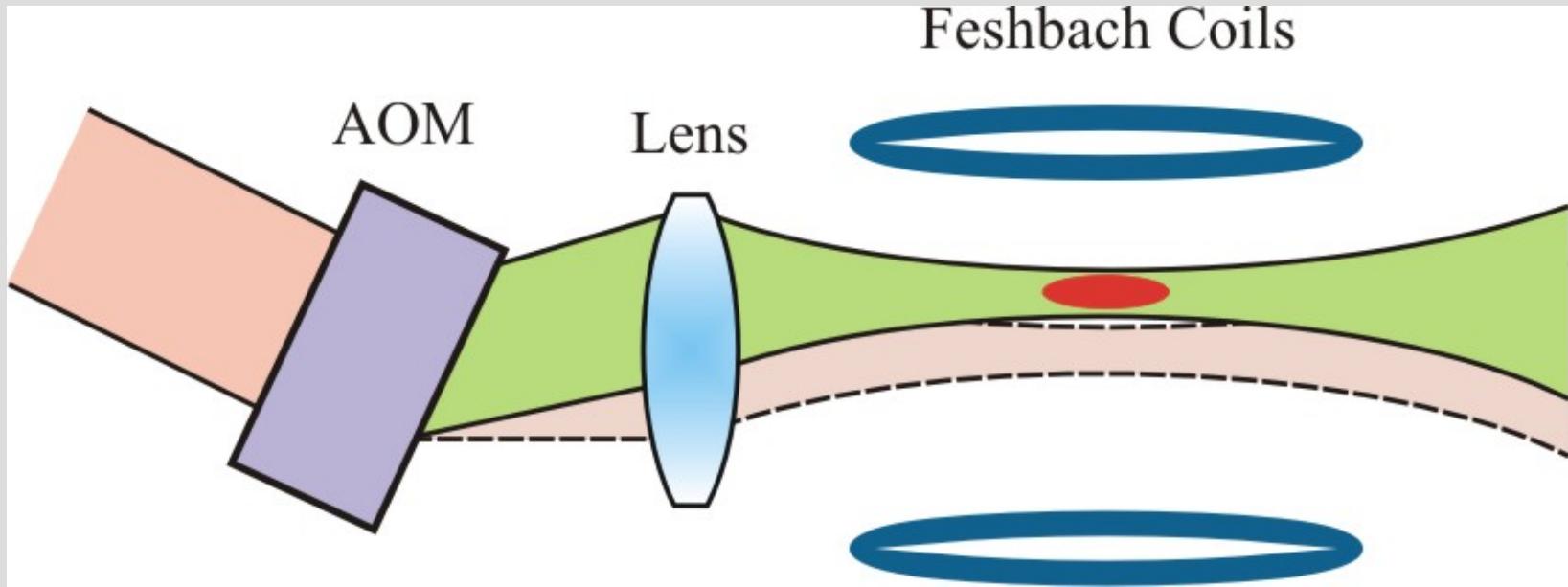
scanning the trap beam



scanning the trap beam

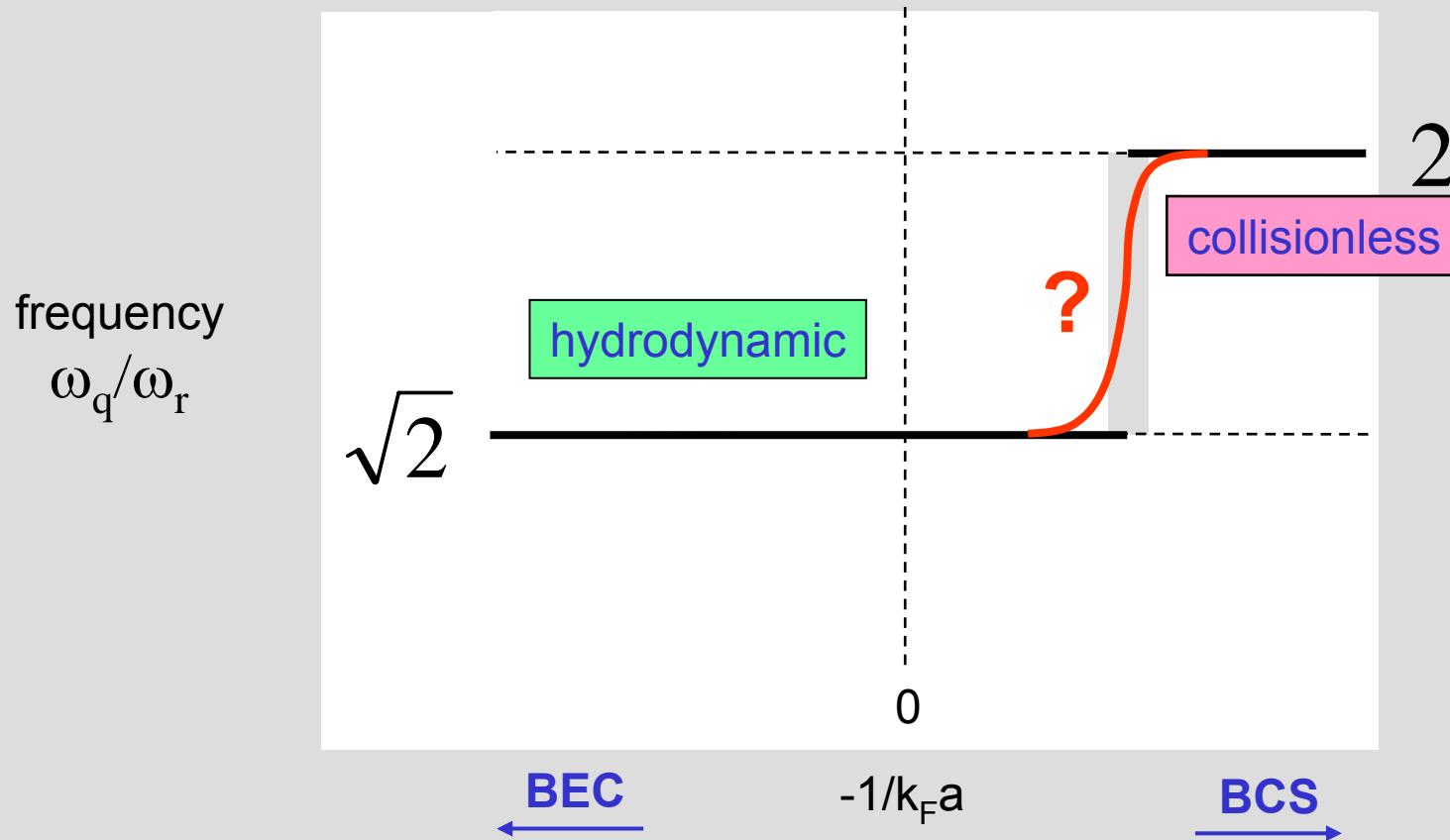


scanning the trap beam



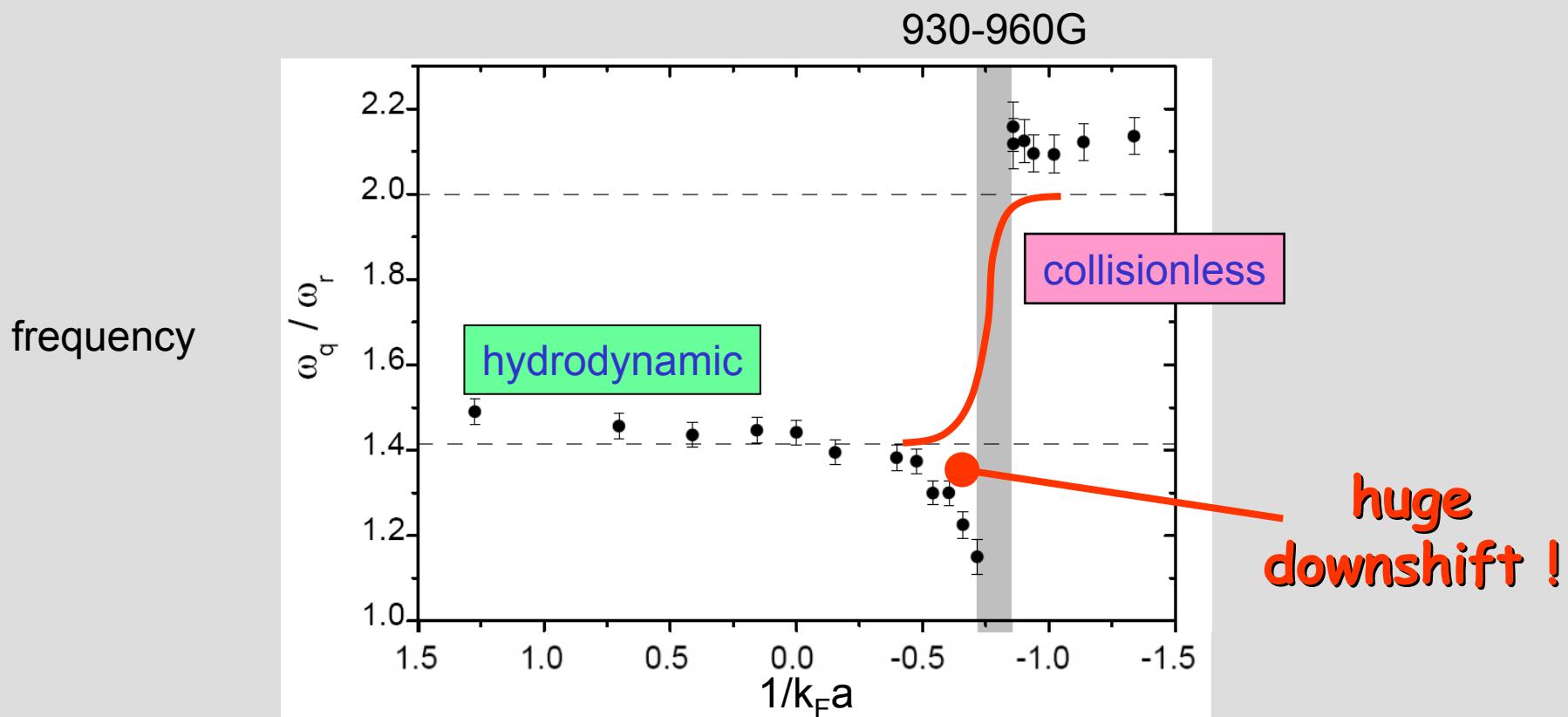
**time-averaged potentials:
a powerful tool for controlled trap deformations**

radial quadrupole mode: expectations

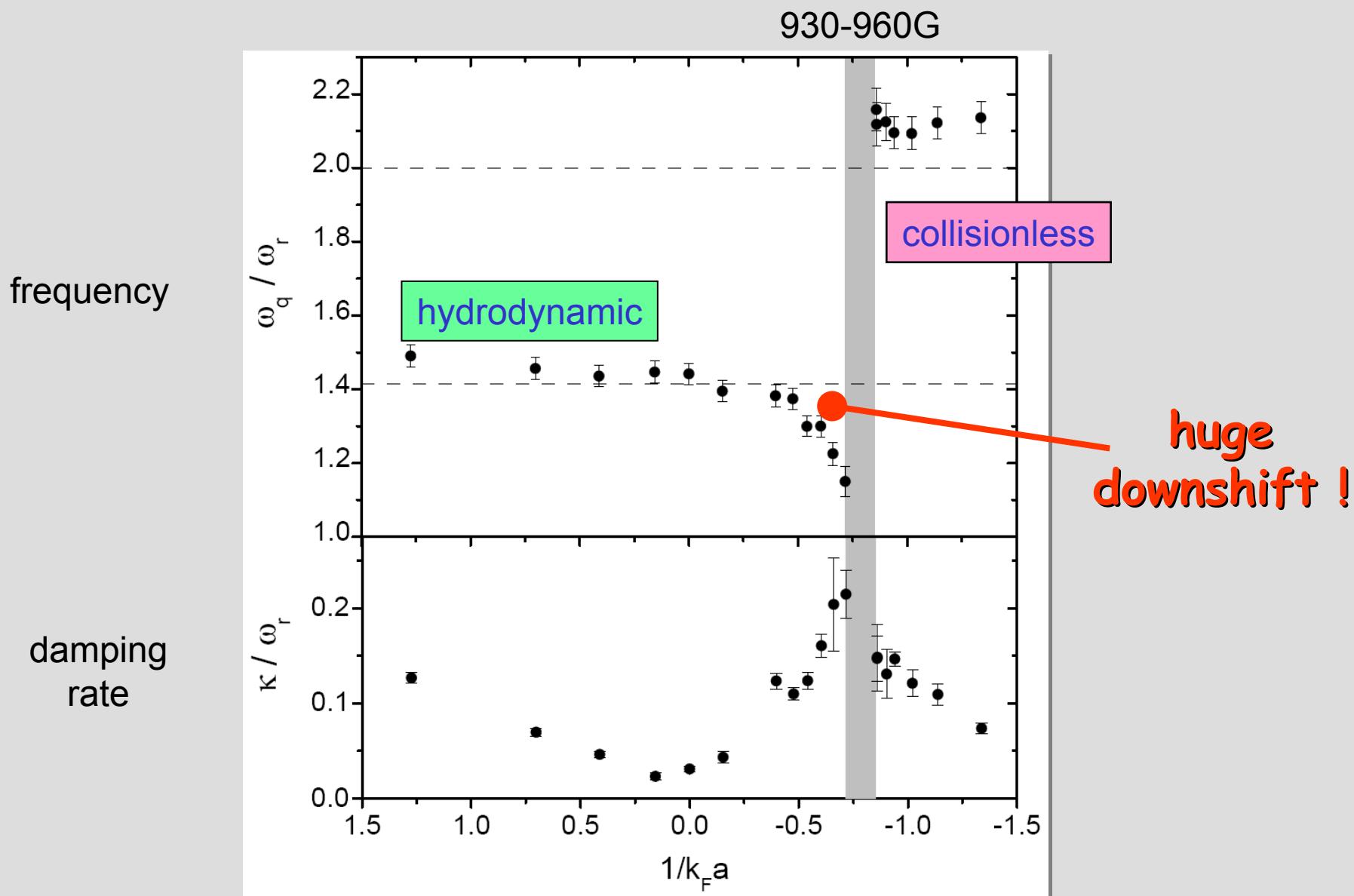


- how sharp is the transition ?
- does it show any structure ?
- smooth change between the two frequencies ?

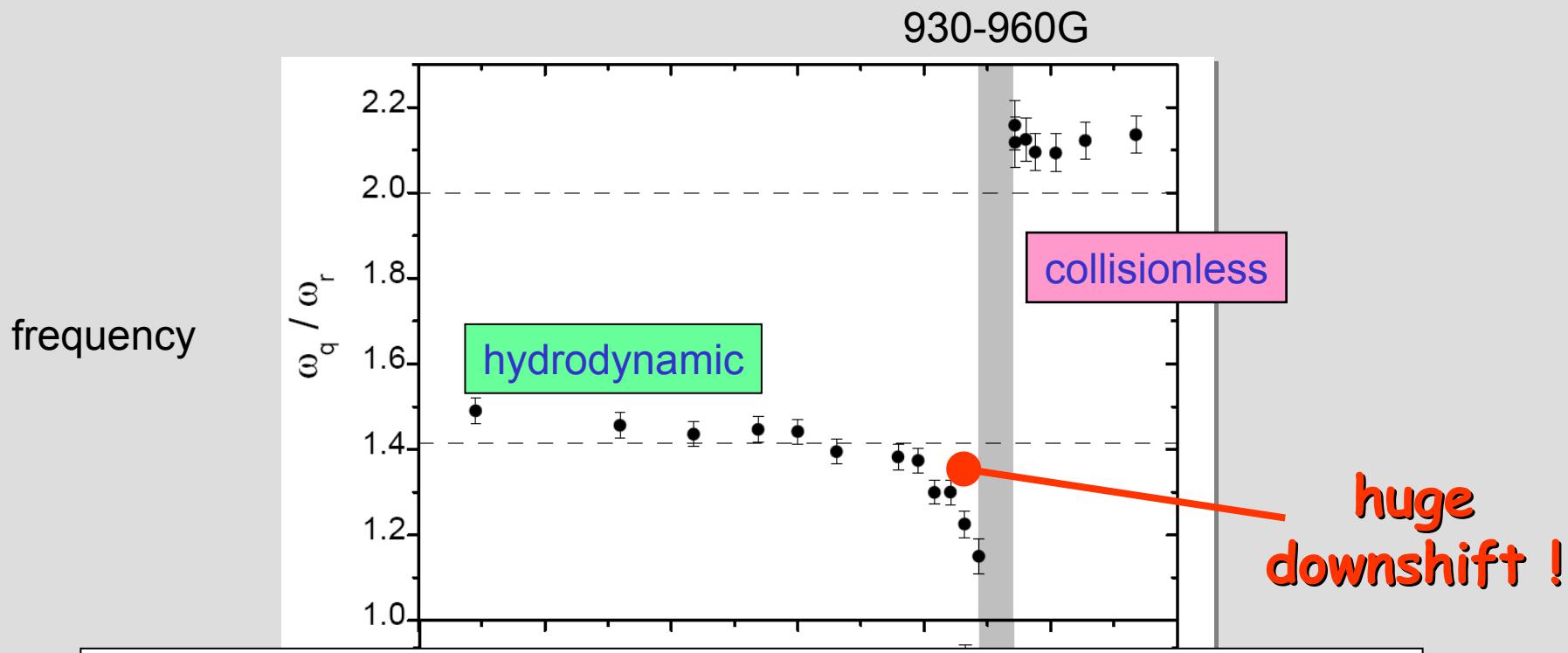
results on radial quadrupole mode



results on radial quadrupole mode



results on radial quadrupole mode



standard hydrodynamic theory breaks down !

coupling of oscillations to pairing gap ?

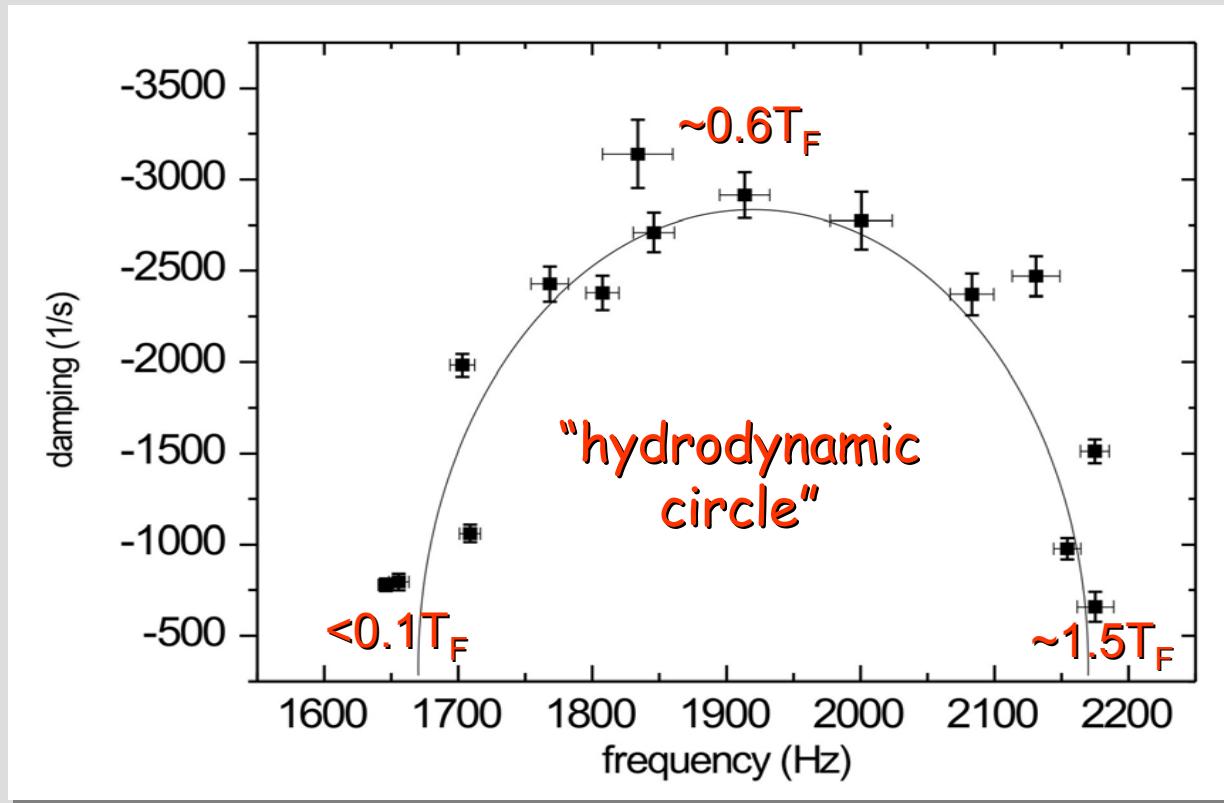
R. Combescot and X. Leyronas, PRL 93, 138901 (2005)

theory needed !!!

temperature-driven transition



at unitarity ($1/k_F a = 0$)



temperature-driven transition from hydodyn. to coll'less
shows “normal” behavior
(for atomic BEC see Buggle et al., PRA 72, 043610 (2005))

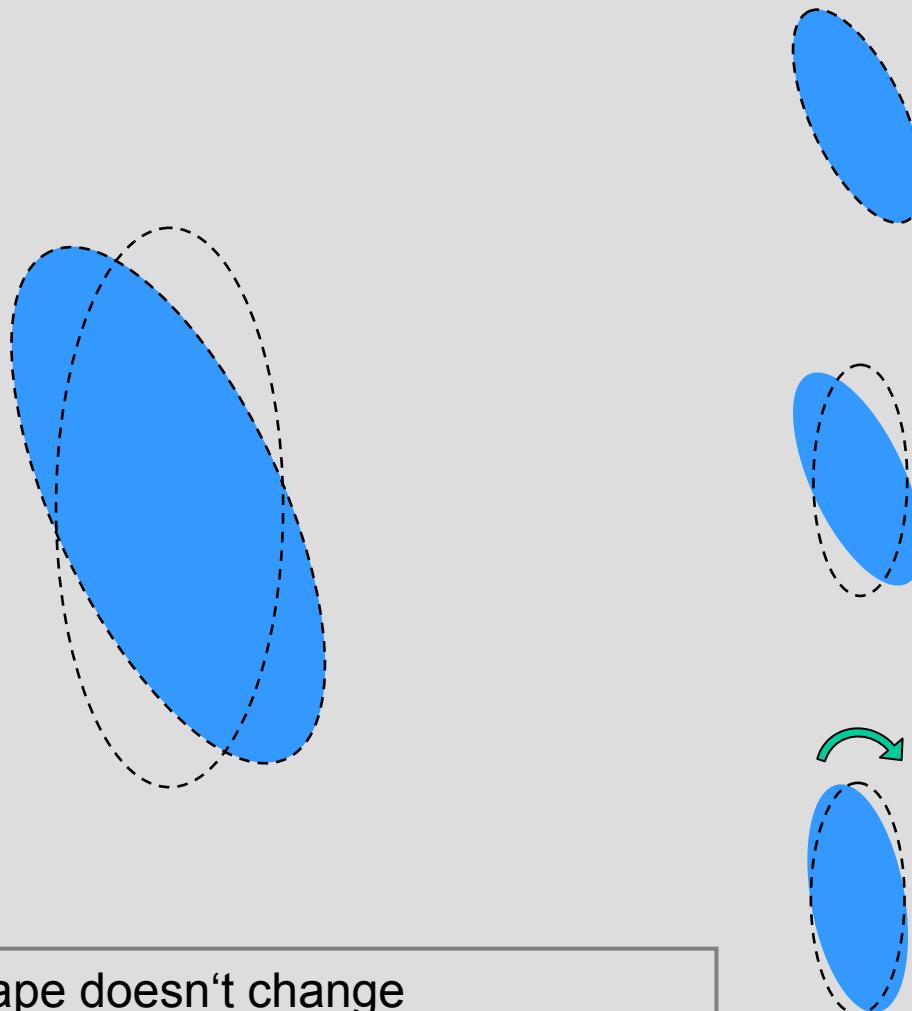
exploring BEC-BCS crossover physics in ultracold Fermi gases

zero-temperature behavior quite well investigated
but, temperature-dependent phenomena
widely unexplored !

what is our best experimental tool to do this ?

scissors modes

break cylindrical
symmetry of trap !



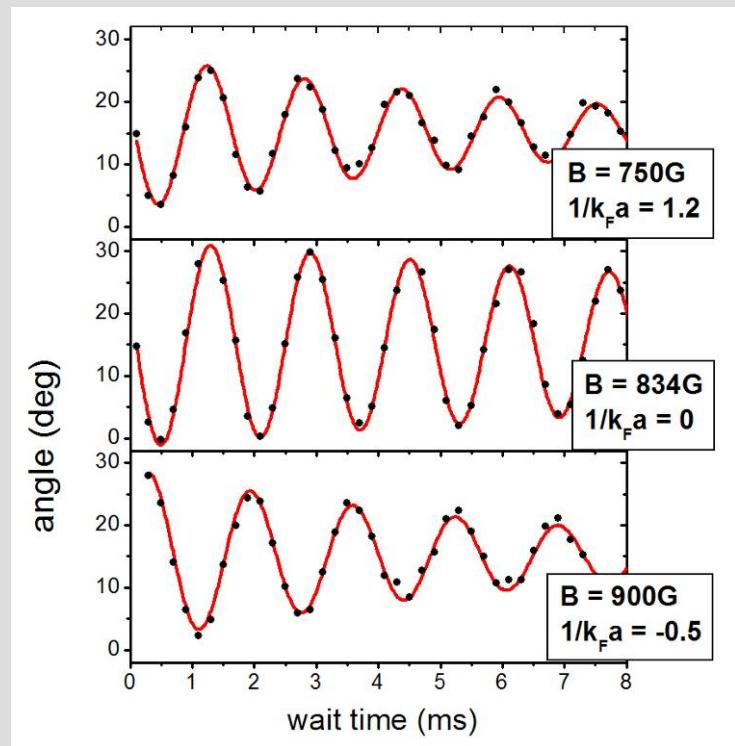
- cloud shape doesn't change
- qualitative difference in the oscillation
for hydrodynamic and collisionless regime

atoms in elliptic potential
aspect ratio ~ 2

sudden change of angle
of the elliptic potential

angle of elliptic cloud
oscillates

scissors oscillation at low T

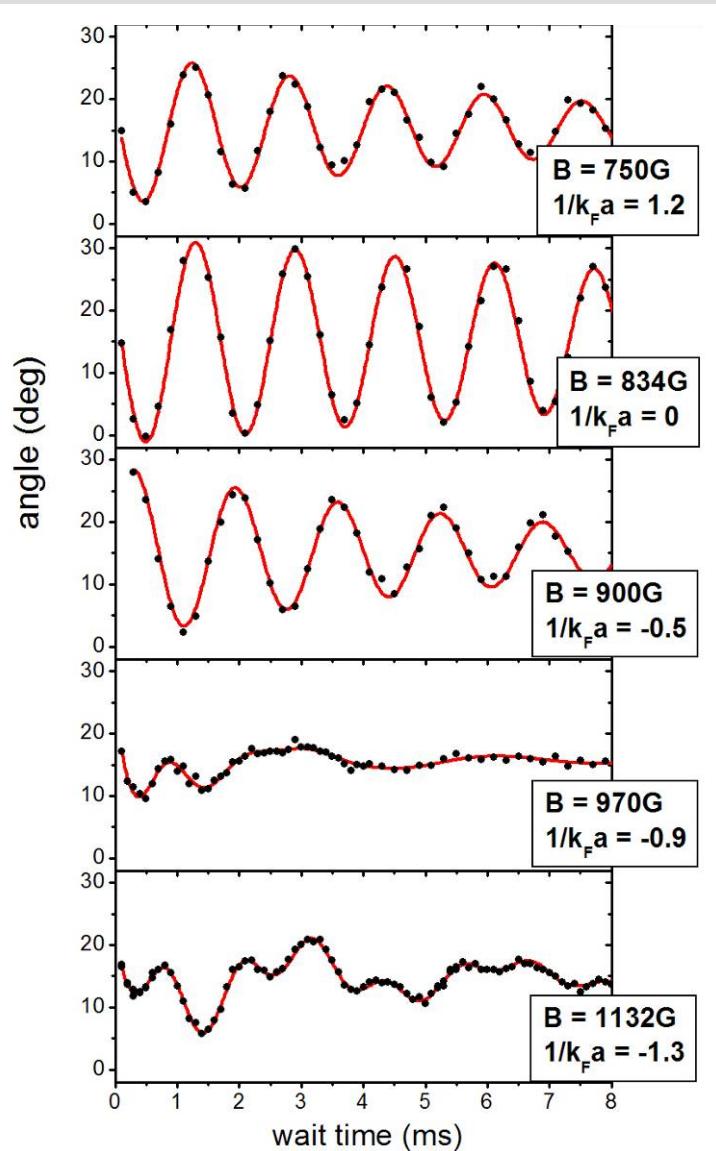


hydrodynamic gas

single frequency oscillation

$$\sqrt{\omega_x^2 + \omega_y^2}$$

scissors oscillation at low T



hydrodynamic gas

single frequency oscillation

$$\sqrt{\omega_x^2 + \omega_y^2}$$

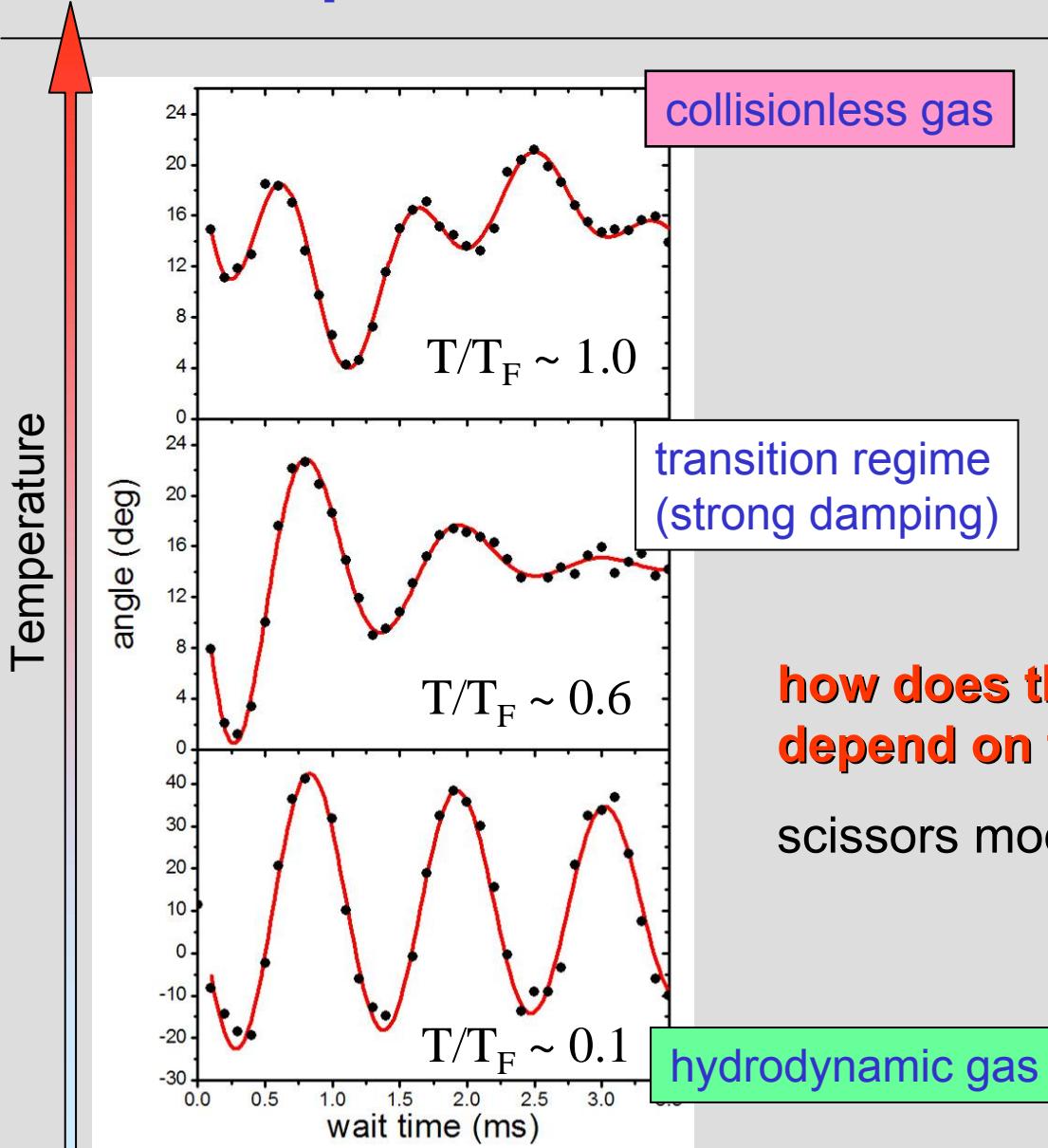
collisionless gas

oscillation with two frequencies

$$\omega_x + \omega_y$$

$$\omega_x - \omega_y$$

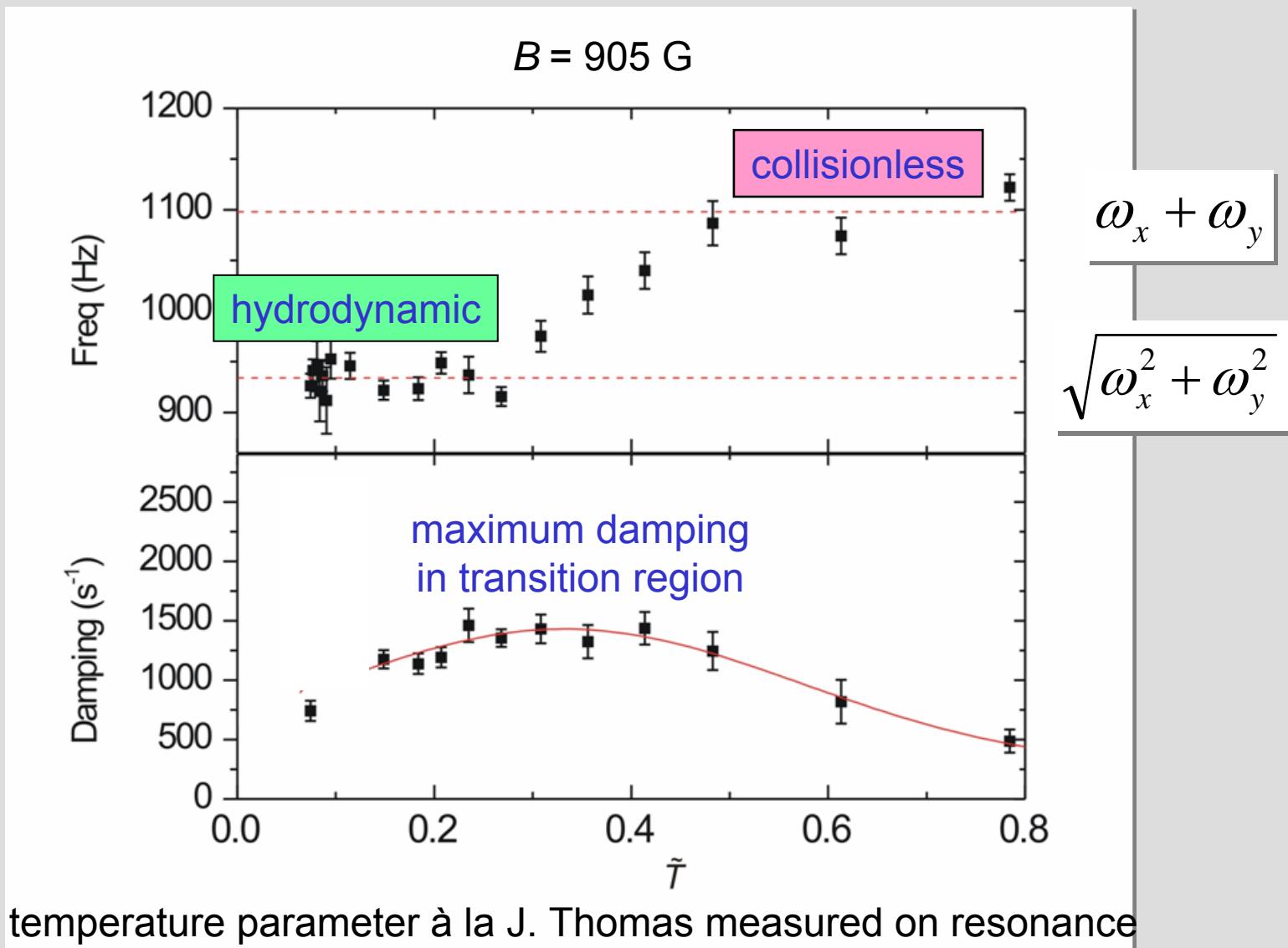
finite temperature



how does the transition temperature depend on the interaction strength ?

scissors mode excellent tool to probe this!

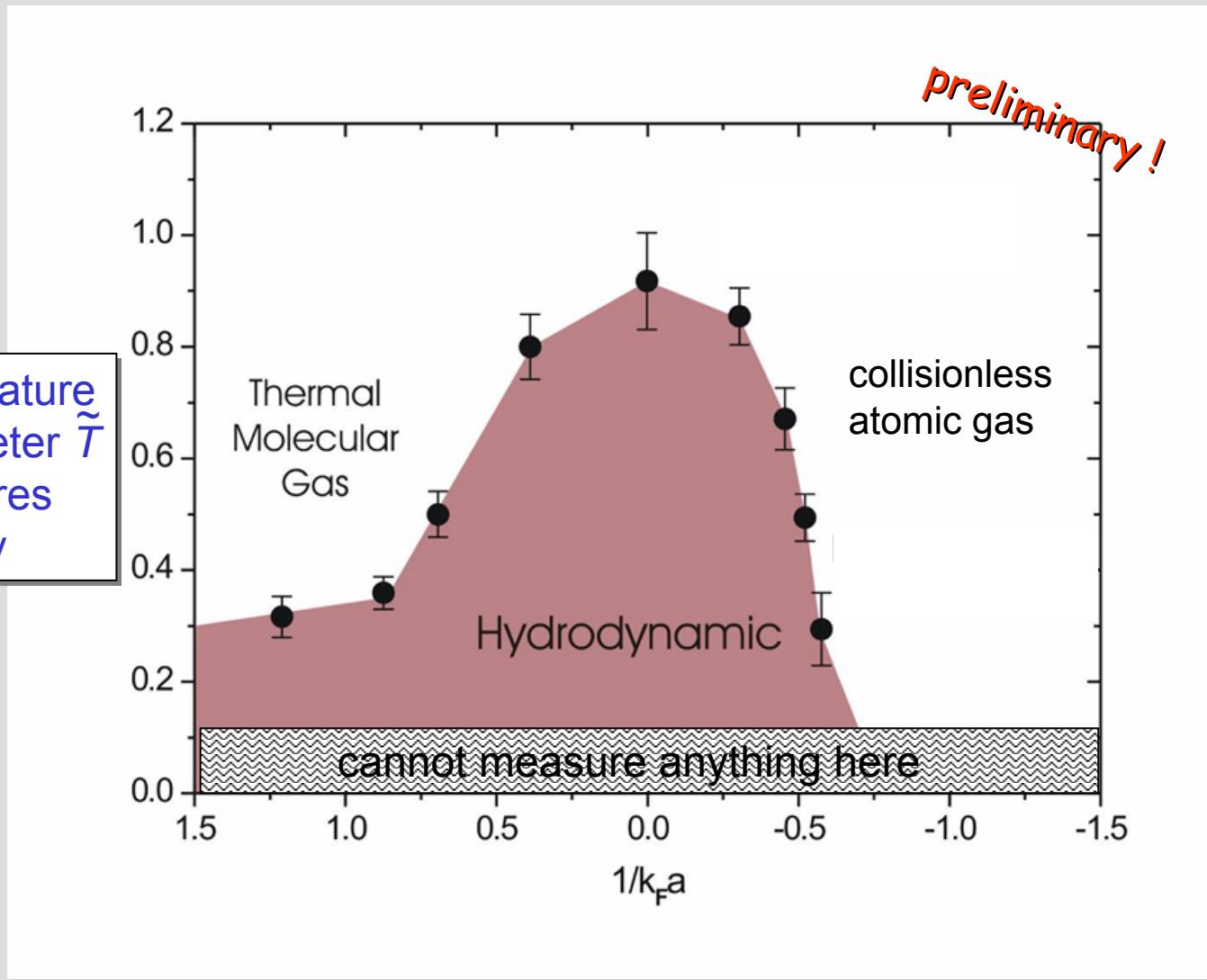
damping vs temperature



transition temperature



temperature parameter \tilde{T}
measures entropy

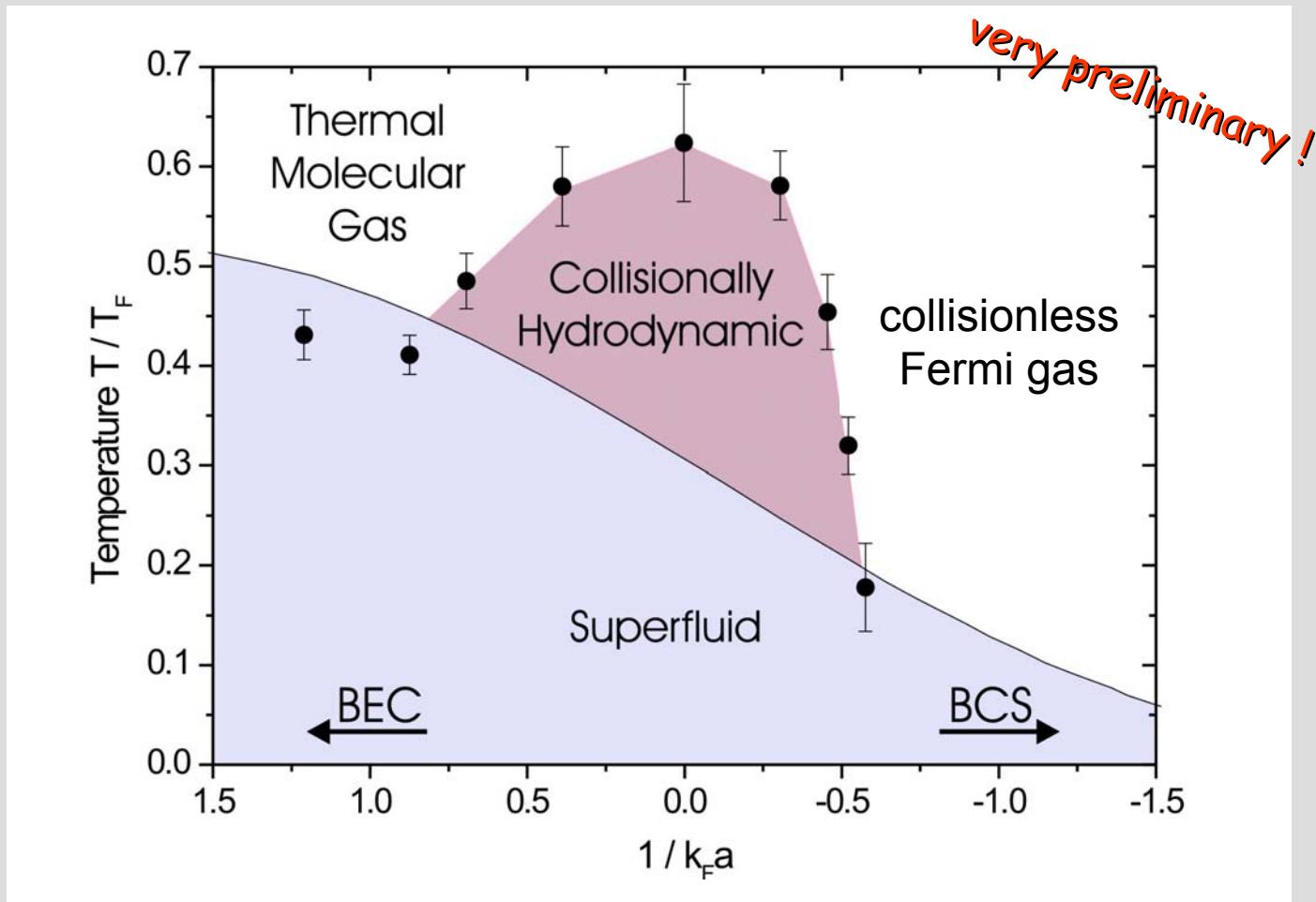


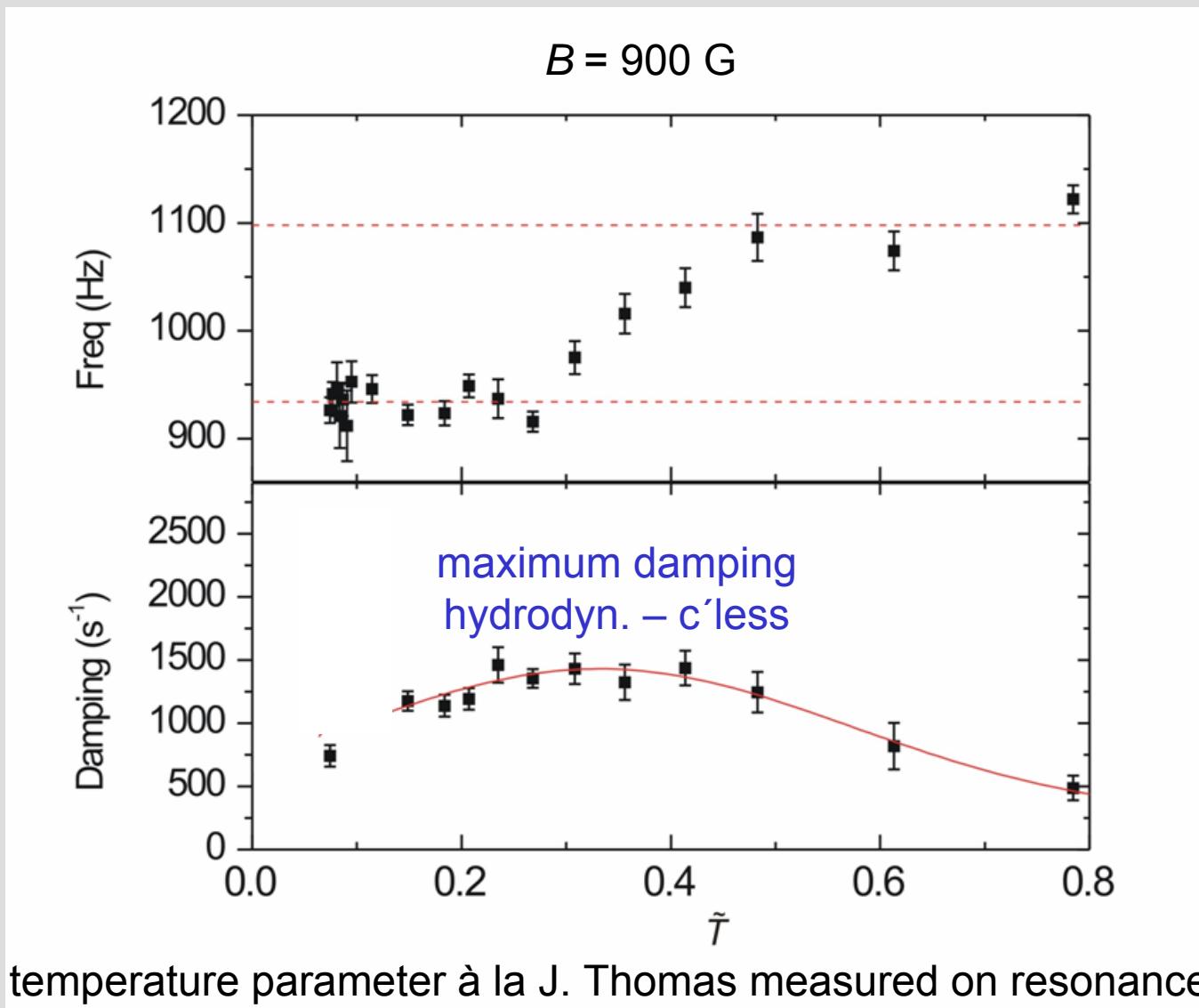
phase diagram



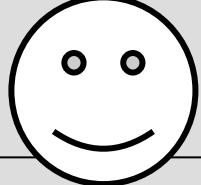
Ultracold atoms
quantum gases

convert the data into a phase diagram
using entropy calculations from K. Levin group (PRL 95, 260405 (1995))

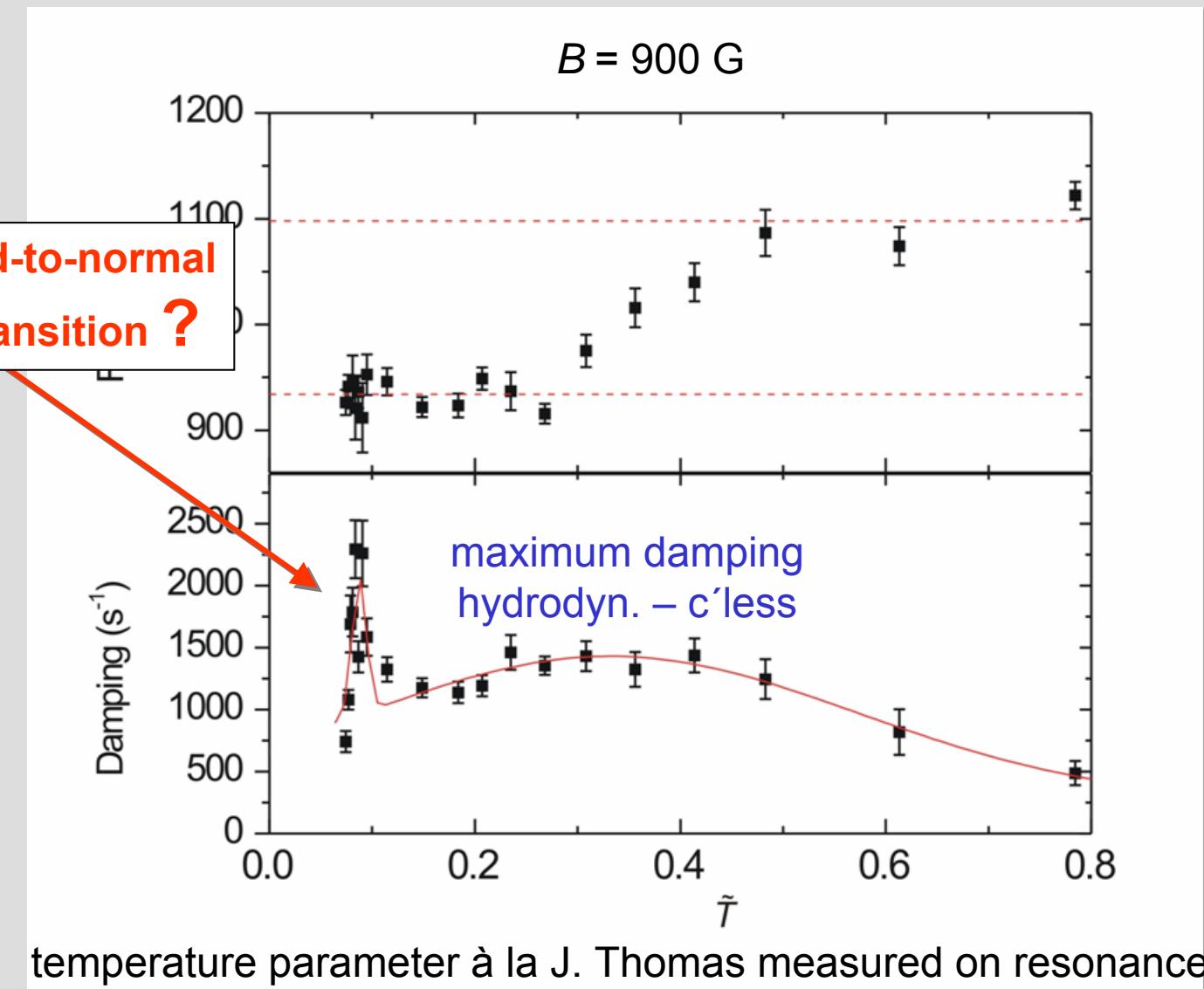




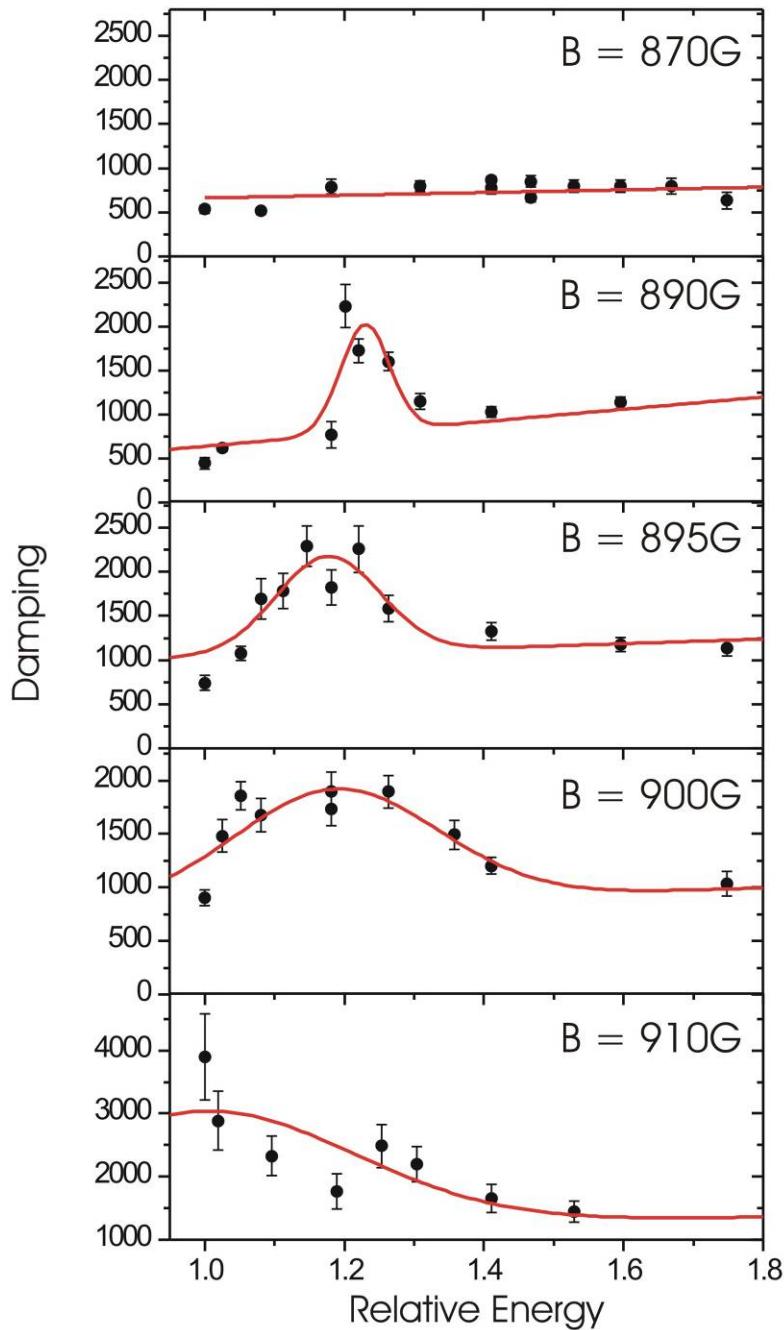
surprise



superfluid-to-normal
phase-transition ?

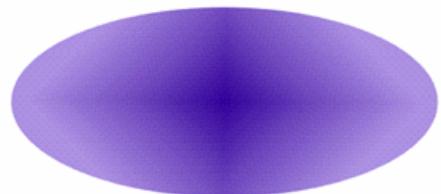


various B



damping feature
only observable
in a narrow B -field
range
on BCS side of
resonance

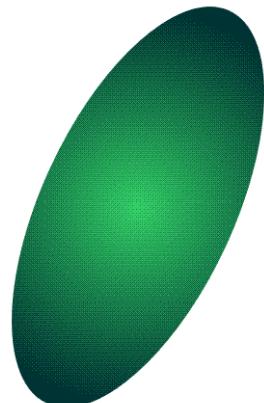
conclusion on surface modes



hydrodynamic-to-collisionless transition
with large change of frequency

strong damping

transition is not smooth!
puzzle of down-shift occurring as a precursor



hydrodynamic-collisionless transition
single frequency – two frequencies

weak damping

temperature-induced transition:
phase diagram for hydrodynamic behavior in crossover

striking damping peak at low temperature:
superfluid phase transition?

The background of the slide features a wide-angle photograph of a snowy mountain range under a clear blue sky. In the lower-left foreground, there's a dark, winding track or path through the snow. The mountains in the distance have patches of snow and rocky peaks.

very preliminary !

^6Li - ^{40}K : first results with a Fermi-Fermi mixture of atoms



Devang
Naik

Gabriel
Kerner

Rudi
Grimm

Florian
Schreck

Frederik
Spiegelhalder

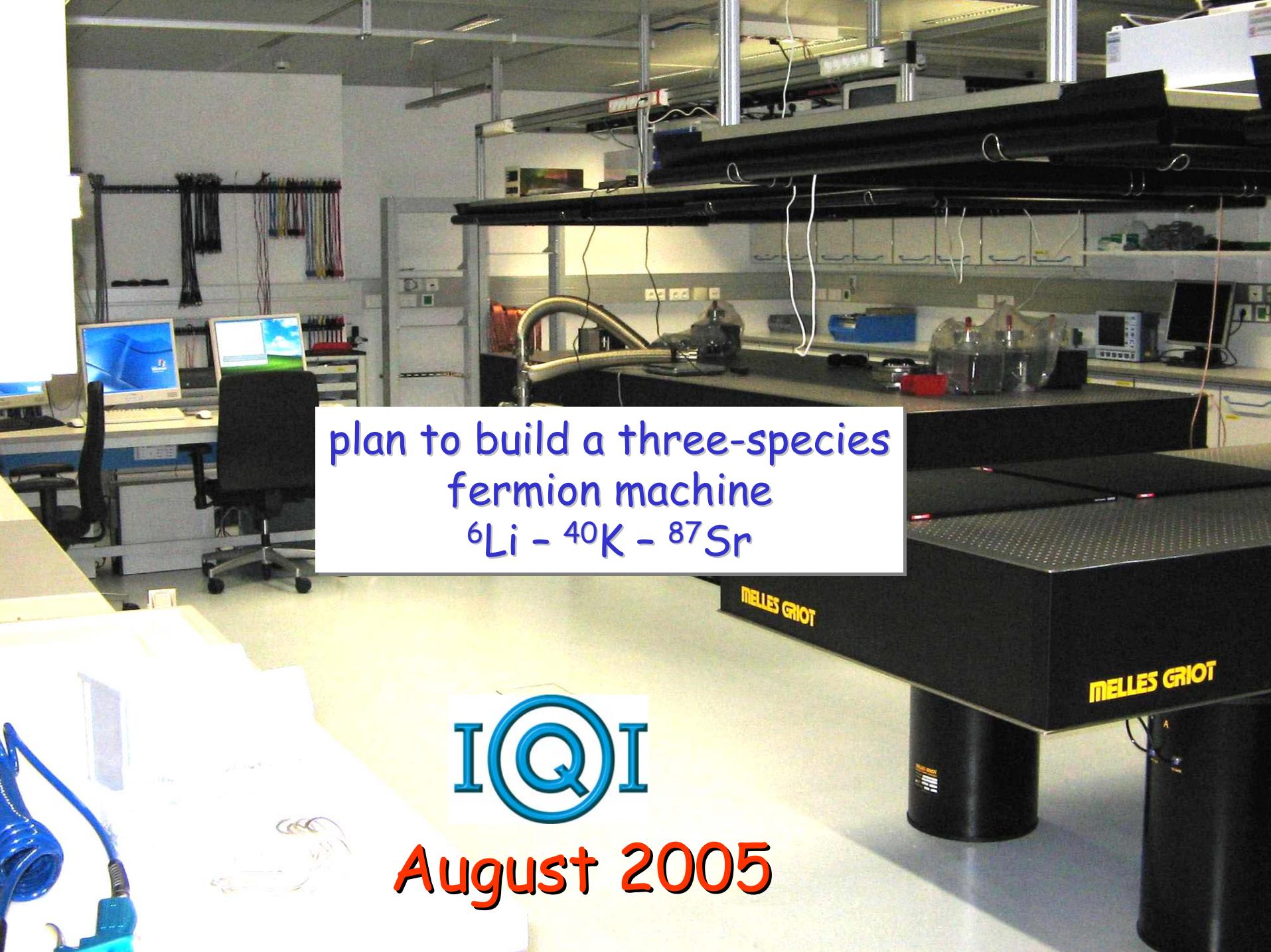
Eric
Wille

Clarice
Aiello

Eileen
Spiegelhalder

Gerhard
Hendl

Andreas
Trenkwalder



plan to build a three-species
fermion machine
 ^6Li - ^{40}K - ^{87}Sr



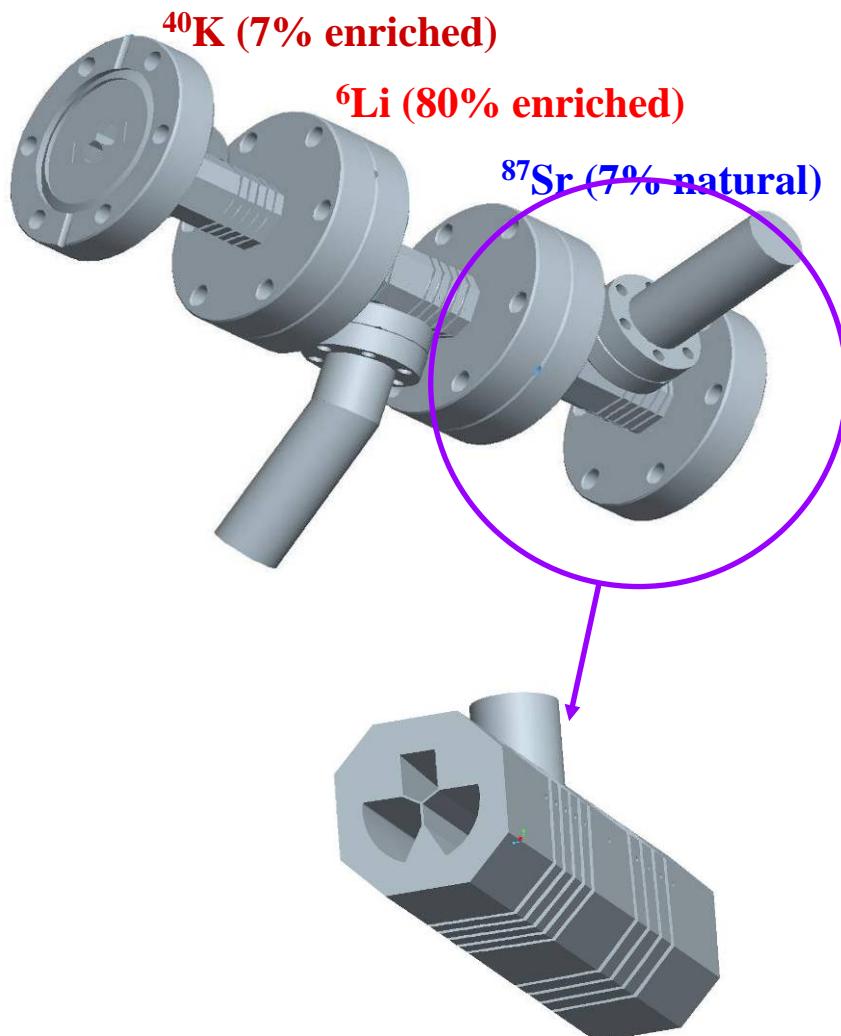
August 2005

A photograph of a laboratory control room and experimental setup. On the left, a person sits at a desk with multiple computer monitors displaying data. Behind them is a rack of electronic equipment and sensors. To the right, a large, complex optical bench or experimental apparatus is visible, with various lenses, mirrors, and cables. The bench is labeled "MELLES GRIOT".

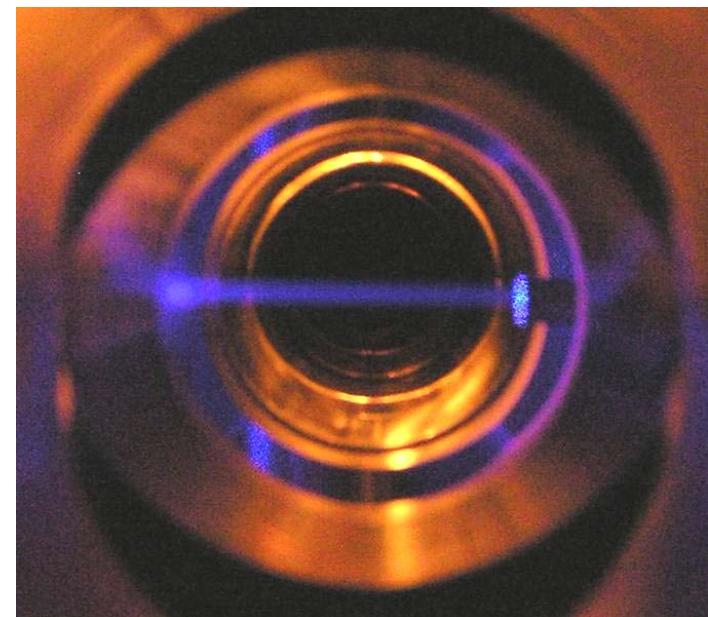
IQI

Februar 2007

three-species oven



Microtubes inside

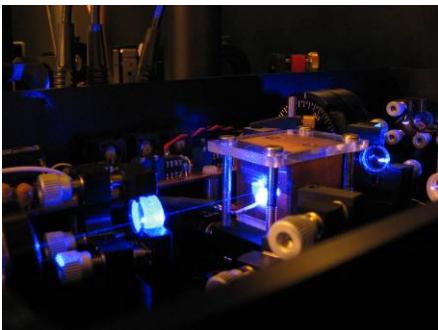


Strontium atomic beam

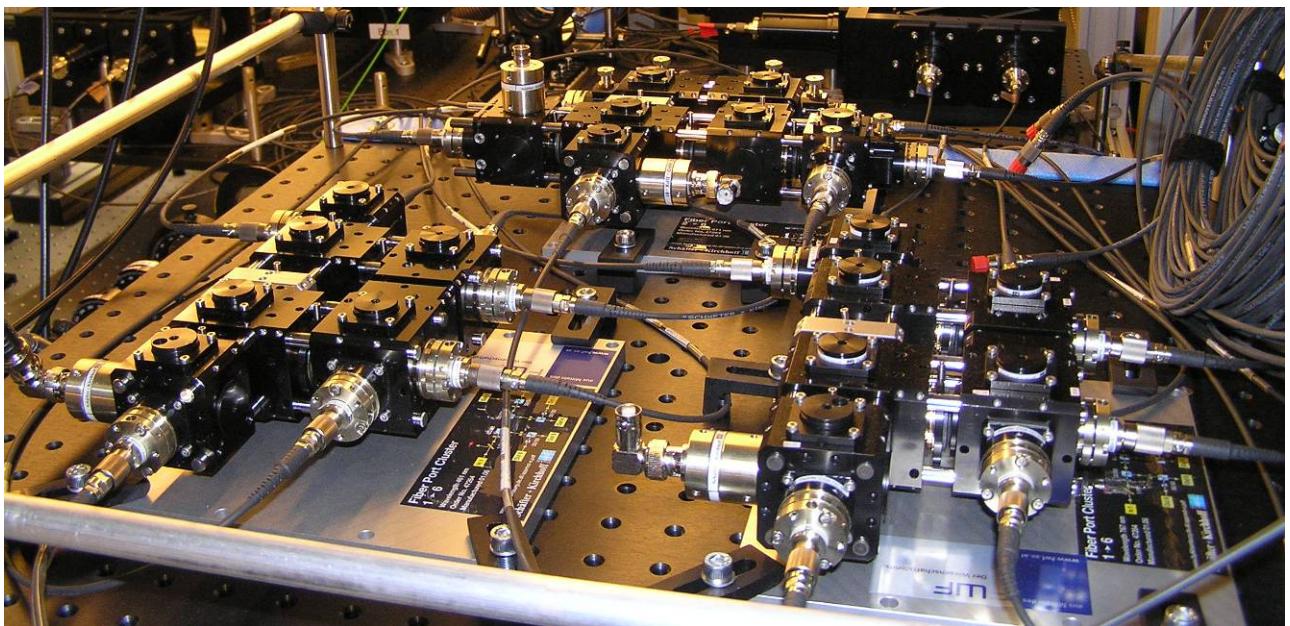
cooling laser system



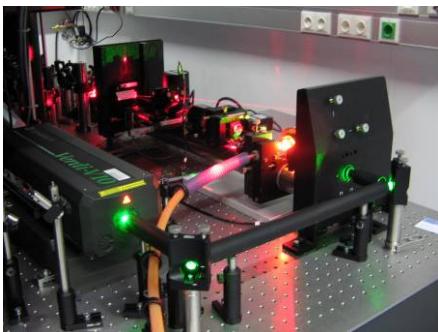
Sr, 461nm: doubled diode laser



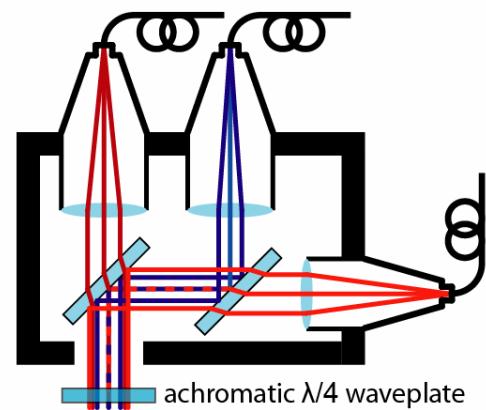
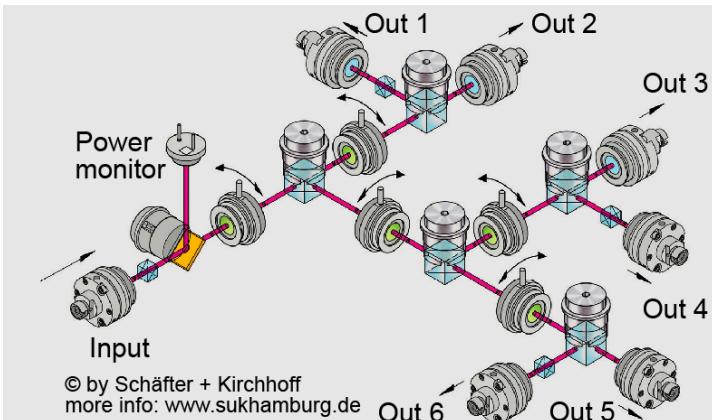
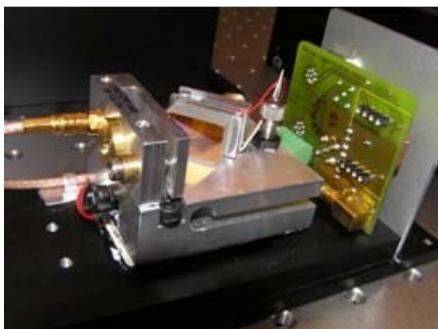
MOT beam delivery:



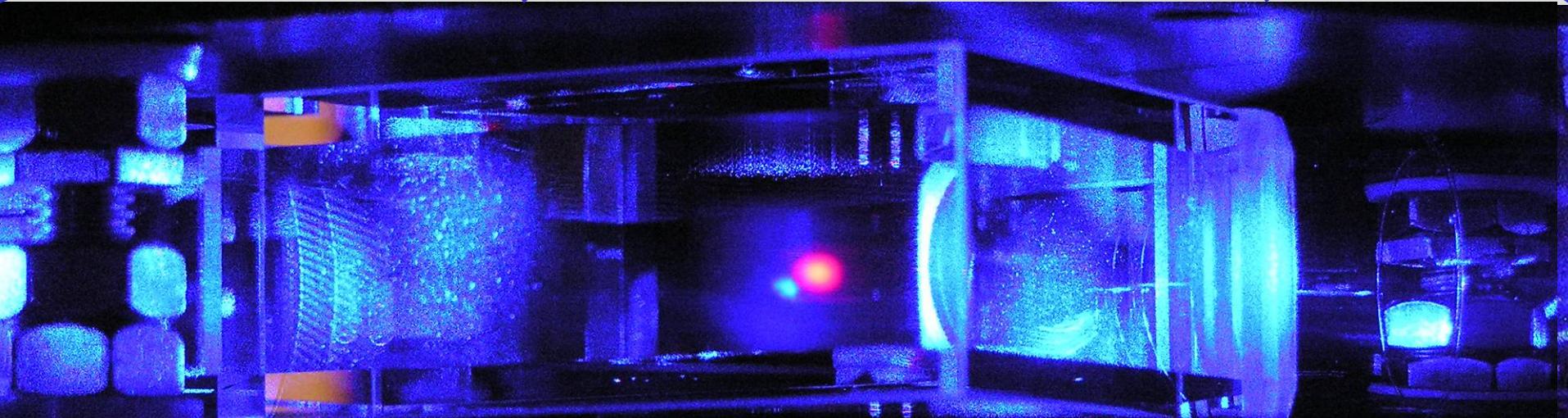
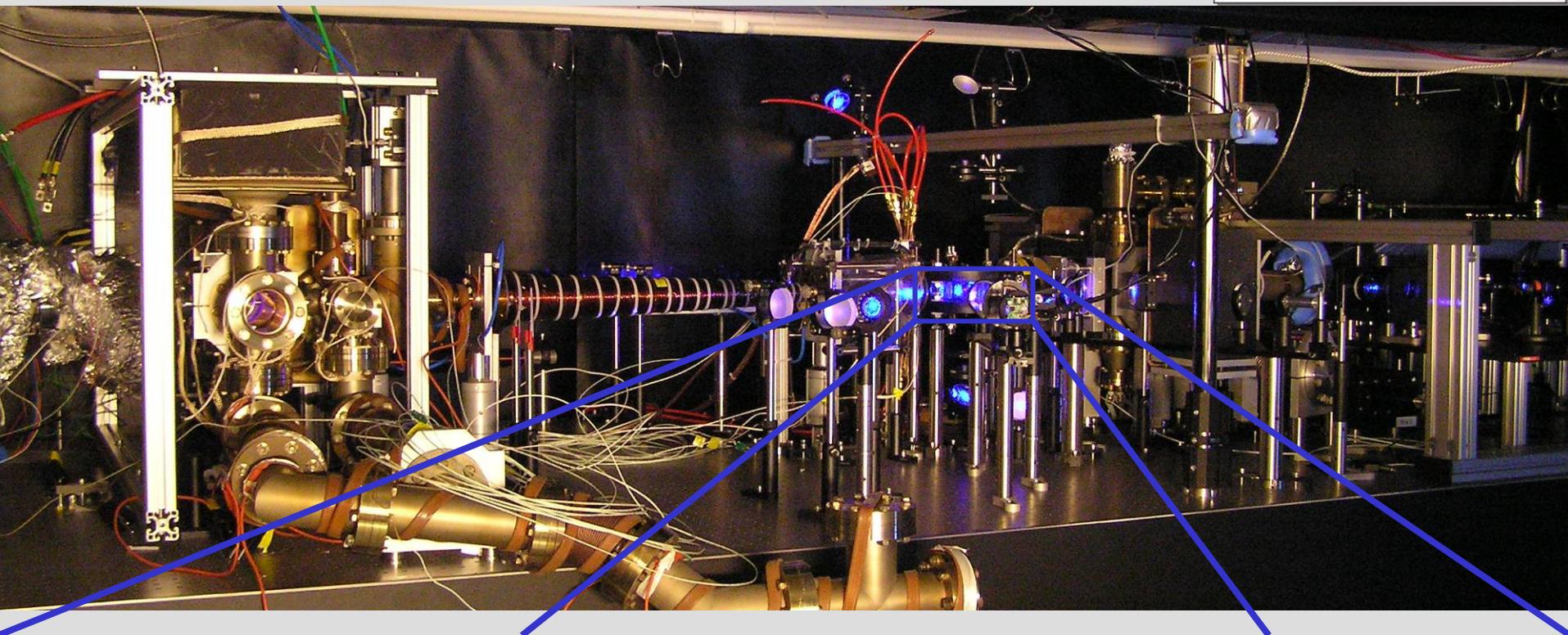
Li, 671nm: dye laser



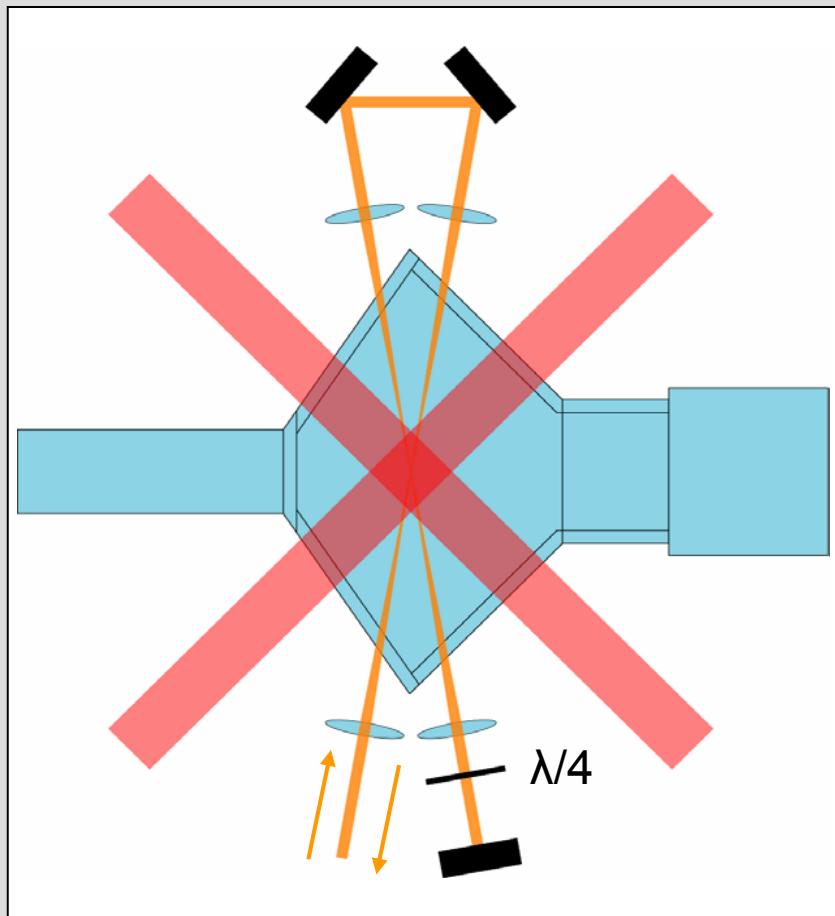
K, 767nm: diode lasers



Lithium-Potassium-Strontium machine



from MOT to dipole trap

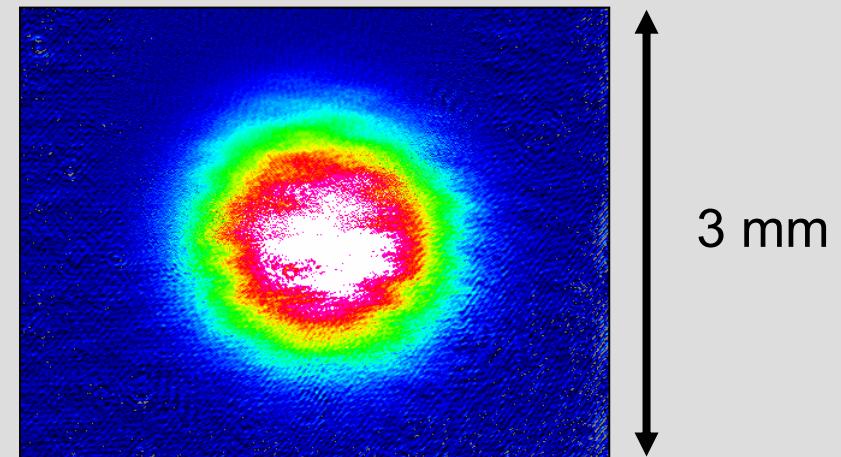


Dipole trap (100W 1075nm laser):

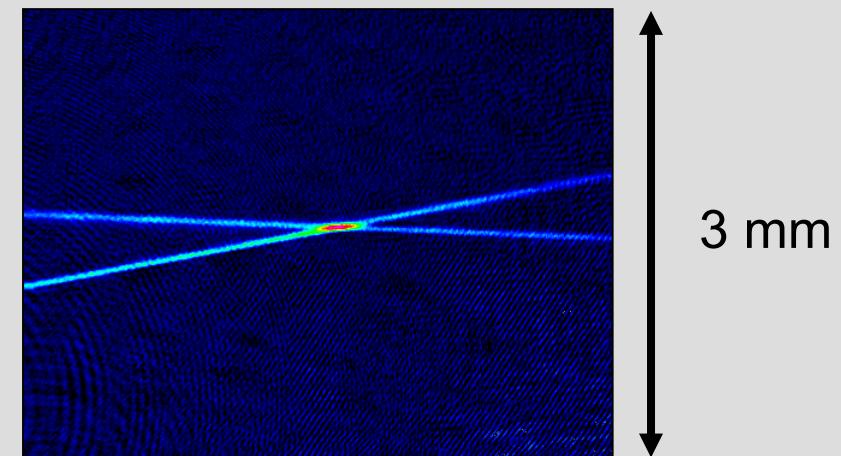
$U \sim k_B 1 \text{ mK}$

$w \sim 60 \mu\text{m}$

${}^6\text{Li}$ MOT: $N \sim 10^9$ $T \sim 300 \mu\text{K}$



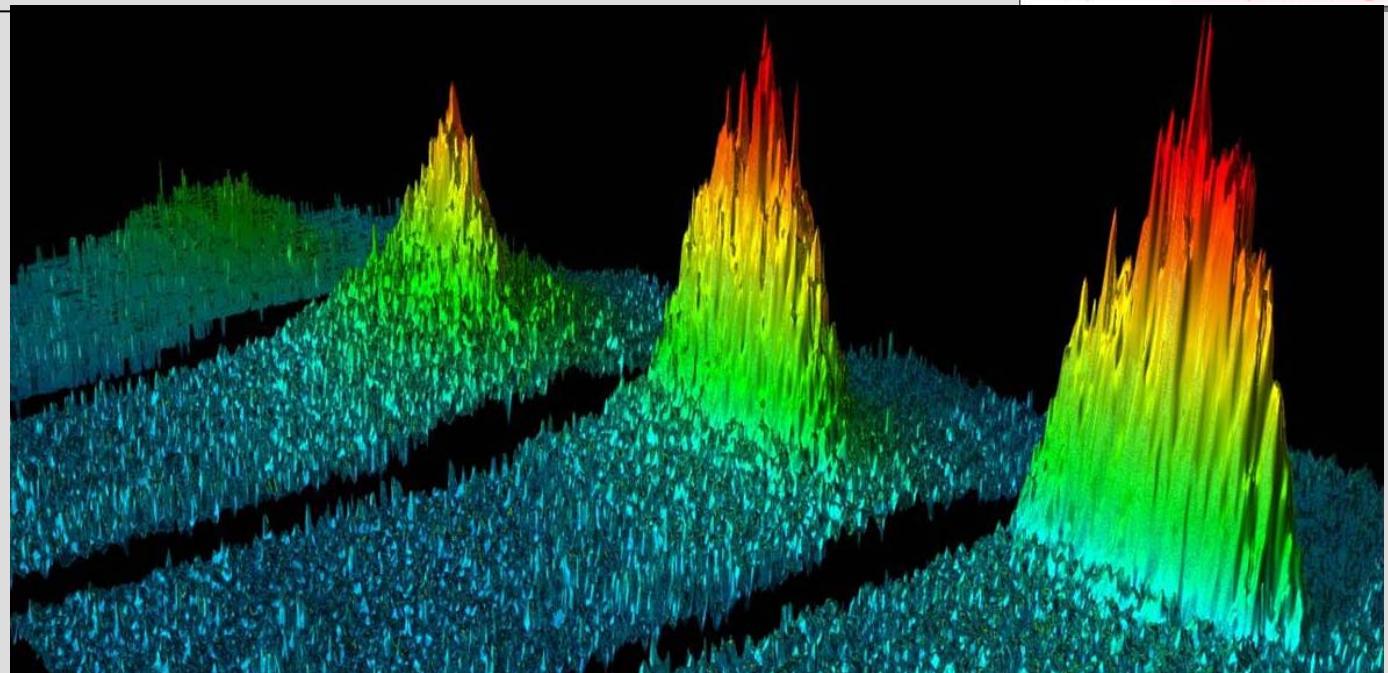
Dipole trap: $N > 10^6$



${}^6\text{Li}_2$ molecular Bose condensation



Ultracold atoms
quantum gases



in dipole trap

After 10ms time of flight:

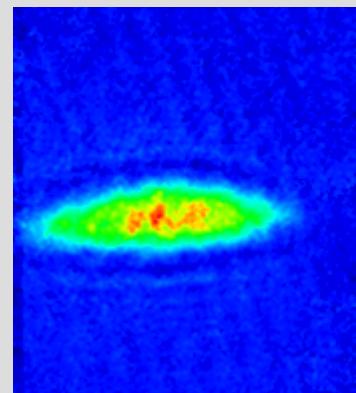
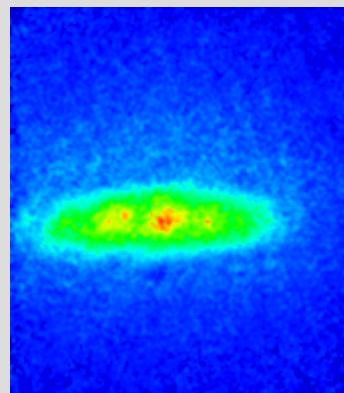
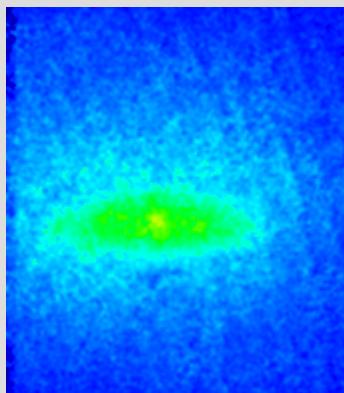
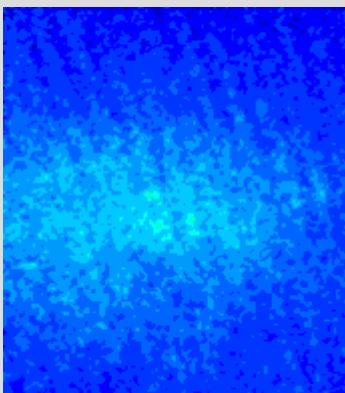
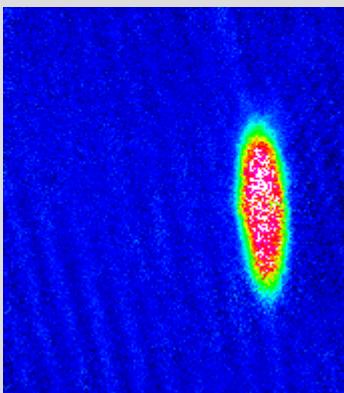
PURE BEC!

4.3 sec evap

4.7 sec evap

4.8 sec evap

5.1 sec evap

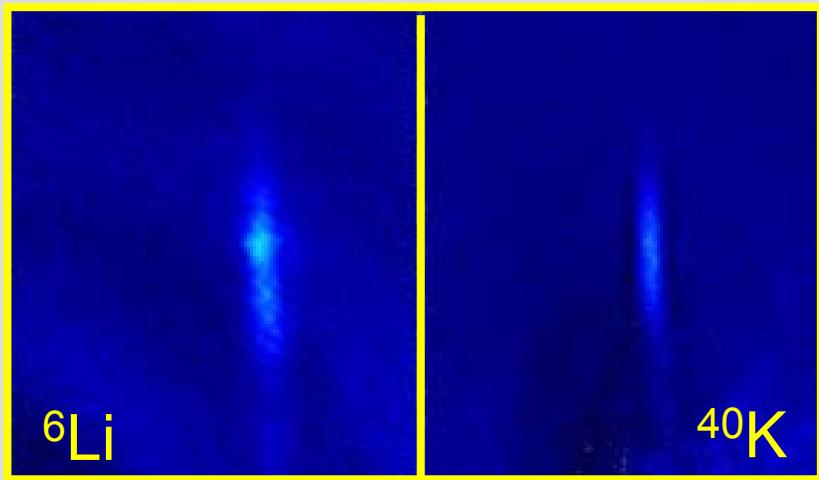


heteronuclear FF mixture !



absorption images of ^6Li and ^{40}K atoms

after 3.9 s of forced evaporative cooling at 750G



$26 \mu\text{K}$ trap depths $55 \mu\text{K}$

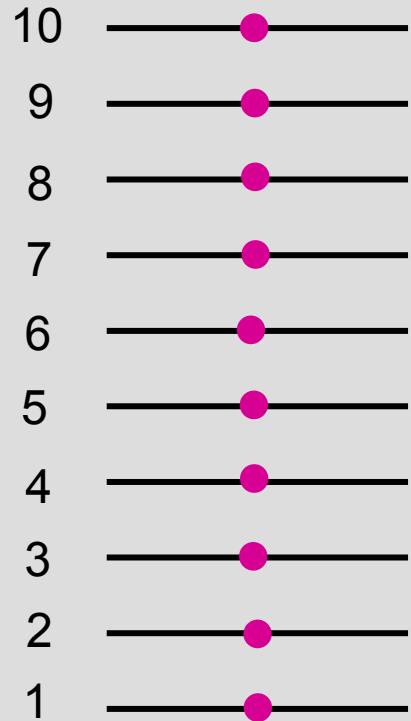
estimated temperature $\sim 2 \mu\text{K}$

heteronuclear Fermi-Fermi mixture
(stable up to the point where $^6\text{Li}_2$ dimers are formed)

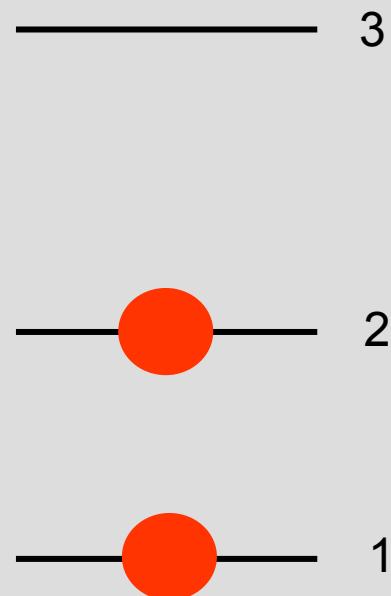
spin relaxation



^{40}K



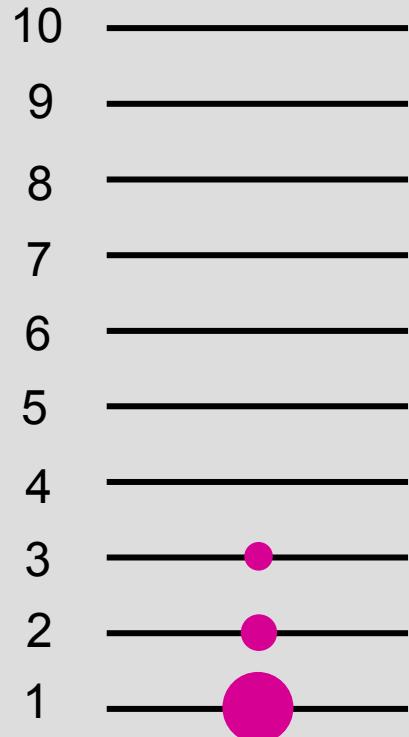
^6Li



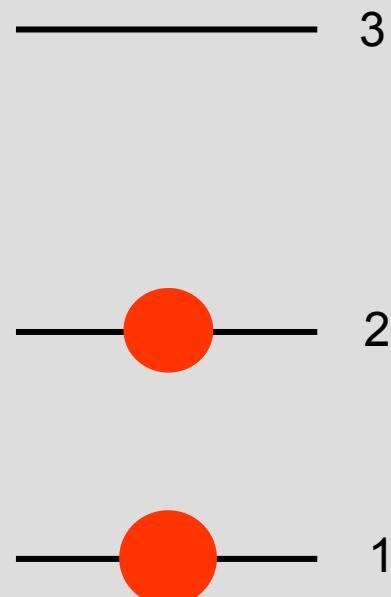
spin relaxation



^{40}K

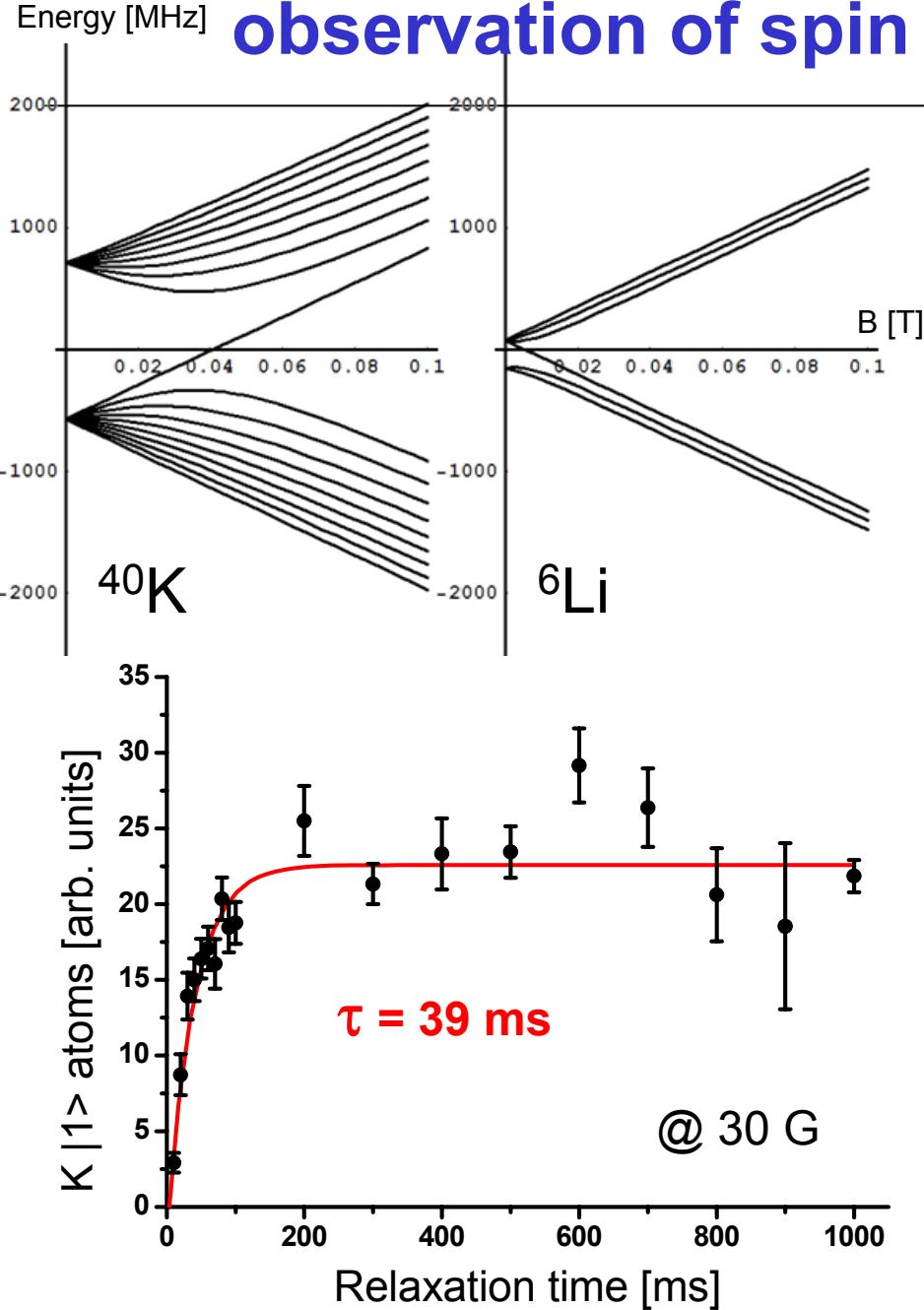


^6Li

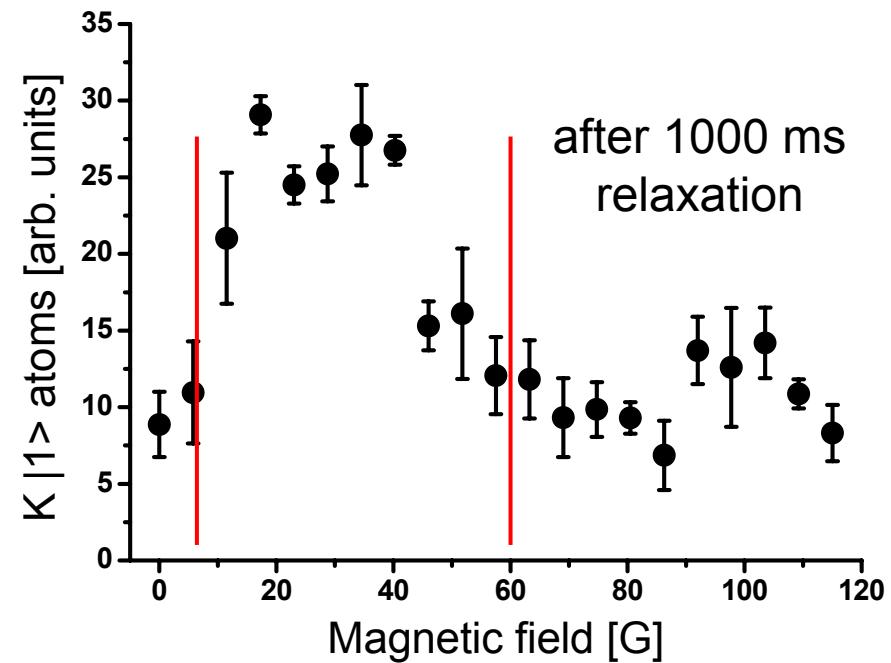




observation of spin relaxation



- Energy splitting @ 6 G: $\sim k_B 100 \mu\text{K}$
- Energy transfer @ 60 G: $\sim k_B 400 \mu\text{K}$ to the ${}^{40}\text{K}$ atom for a m_f changing collision
- Temperature: $\sim 300 \mu\text{K}$
- Trap depth: $\sim 3000 \mu\text{K}$ for ${}^{40}\text{K}$



spin relaxation



^{40}K

^6Li

stable mixtures can be created
if one of the species is fully polarized
into the lowest state !

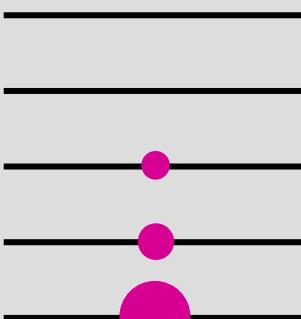
$\Delta m = -1$

$-5/2$

$-7/2$

$-9/2$

m



$-3/2$

$-1/2$

$\Delta m = +1$

$1/2$

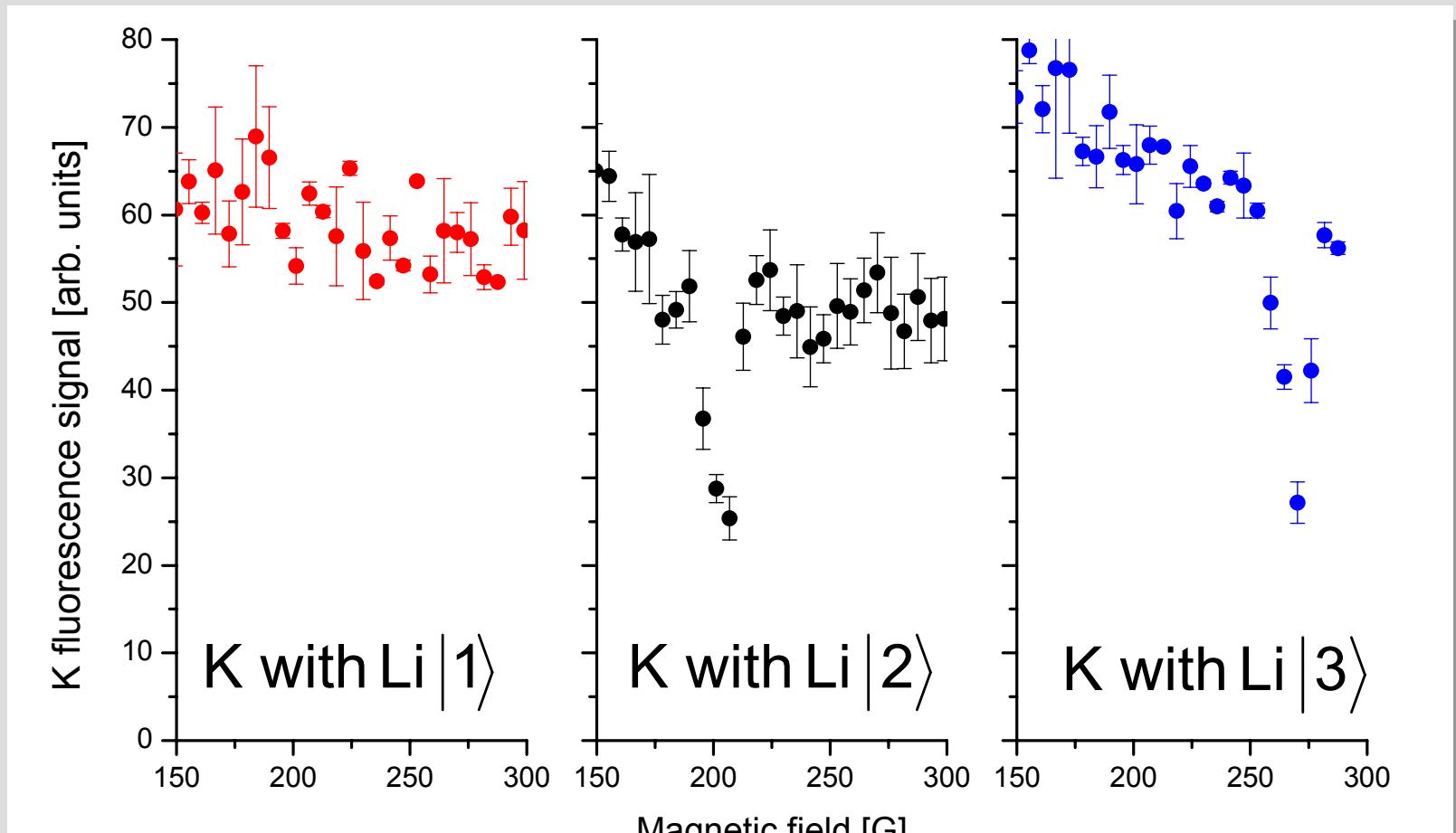
m

loss resonances



pure spin states of ${}^6\text{Li}$ prepared (optical cleaning, rf transfer)

${}^{40}\text{K}$ in a mixture of different states



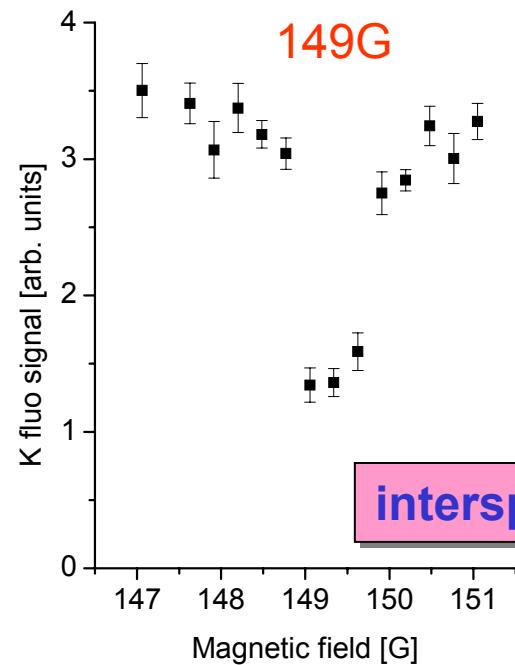
further resonances



Ultracold atoms
quantum gases

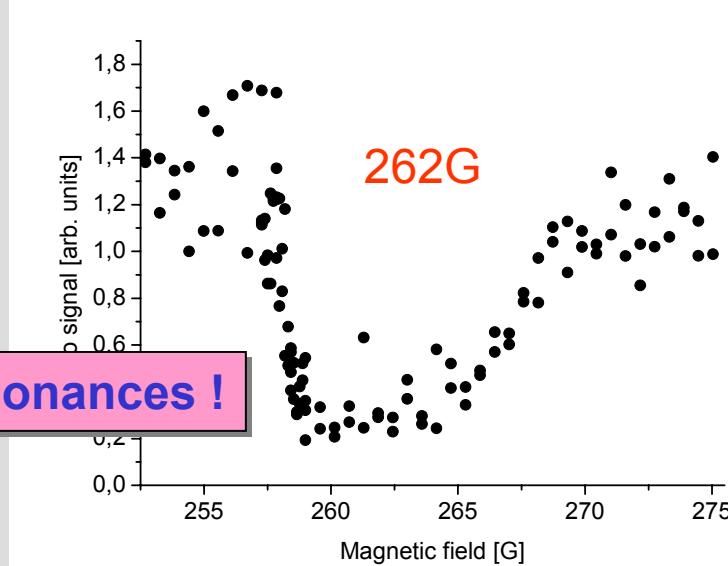
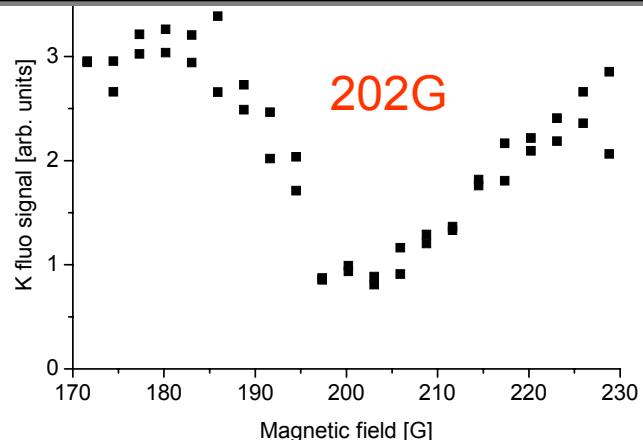
three-component mixture
of ${}^6\text{Li}$ (lowest state)
with ${}^{40}\text{K}$ (lowest two states)

monitor ${}^{40}\text{K}$ atom number



interspecies resonances !

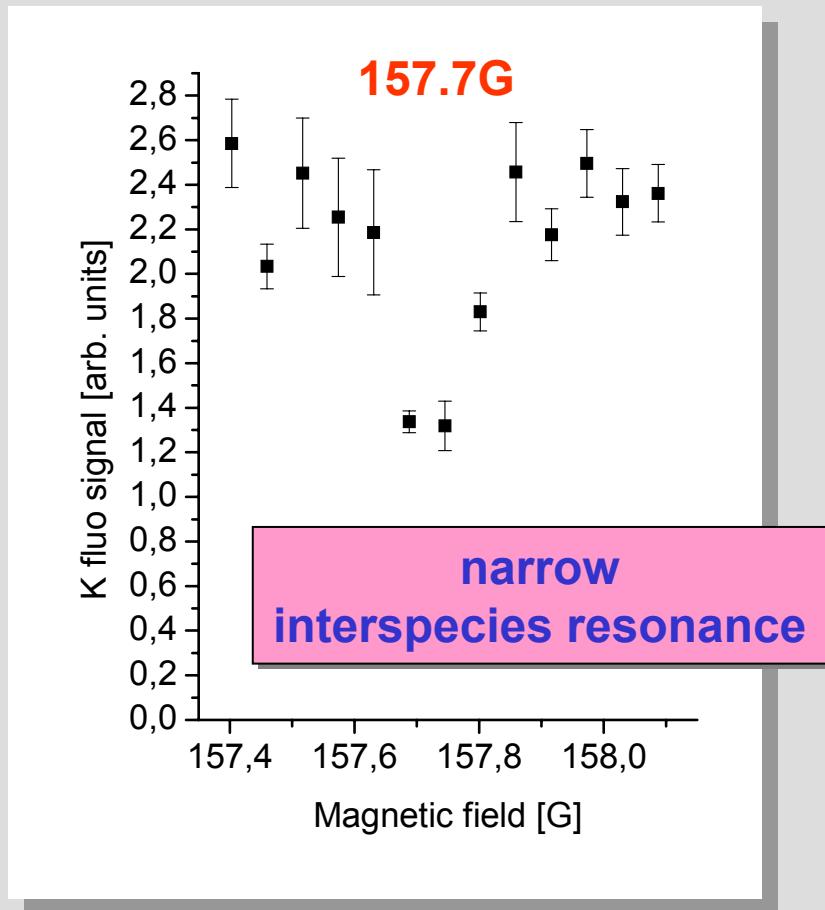
the “famous” ${}^{40}\text{K}$ resonance (homonuclear)



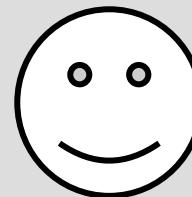
further resonances



two-component mixture
of ${}^6\text{Li}$ (lowest state) with ${}^{40}\text{K}$ (lowest state)



there is more
to come (soon)...



preliminary conclusions on ${}^6\text{Li}-{}^{40}\text{K}$



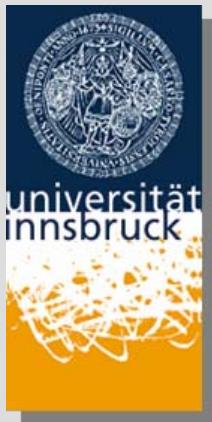
Fermi-Fermi mixtures available in the lab!

stable mixture if one species is in lowest state state
(otherwise spin relaxation)

sympathetic cooling of ${}^{40}\text{K}$ by ${}^6\text{Li}$ works well
at the large- a side of the Li-resonance
(at least until molecules are formed)

several interspecies resonances are observed
(${}^{40}\text{K}$ spin states tbd)

experiments in progress, right now in the lab in Innsbruck...



the end



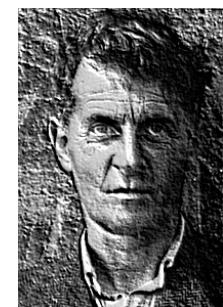
funding



Der Wissenschaftsfonds.



bm:bwk



Ludwig Wittgenstein
(1889-1951)
Austrian Philosopher