

Bose gas in Flatland

Berezinskii-Kosterlitz-Thouless Physics in an Atomic Gas

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***now at: University of Cambridge**



1530
Collège de France

Outline of the talk

Bose gases in 2D

Berezinskii-Kosterlitz-Thouless transition

Homogeneous vs. trapped & ideal vs. interacting gas

Critical point of an interacting 2D gas

BEC vs. BKT

Vortices and quasi-long-range coherence

Long-range order in reduced dimensionality

more vulnerable to fluctuations, disorder...

c.f. classical transport:



1D - impossible



3D - easy



2D - marginal

BEC, coherence, and superfluidity in 2D

Homogeneous 2D Bose fluid in the thermodynamic limit

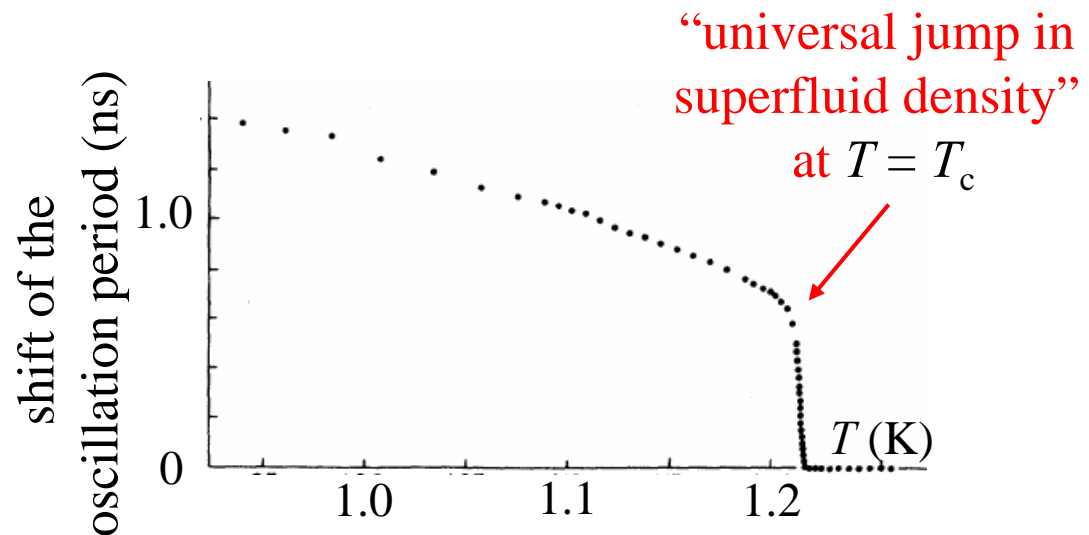
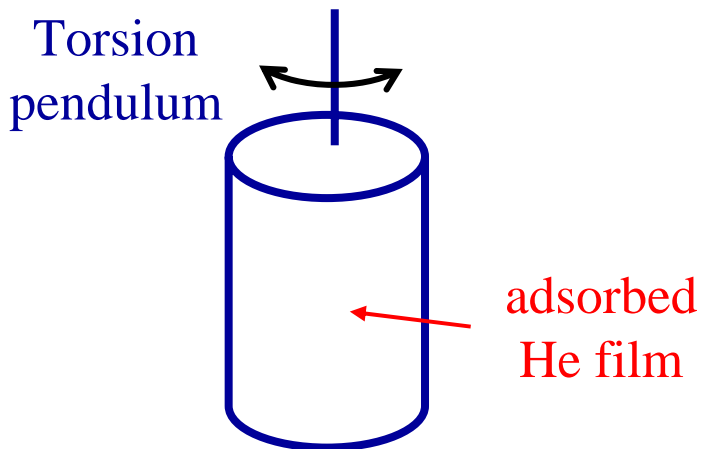
No BEC in an ideal gas

No true long-range order in an interacting gas at finite T

(Mermin-Wagner-Hohenberg theorem)

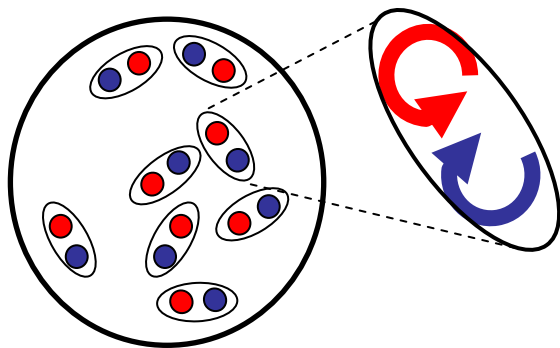
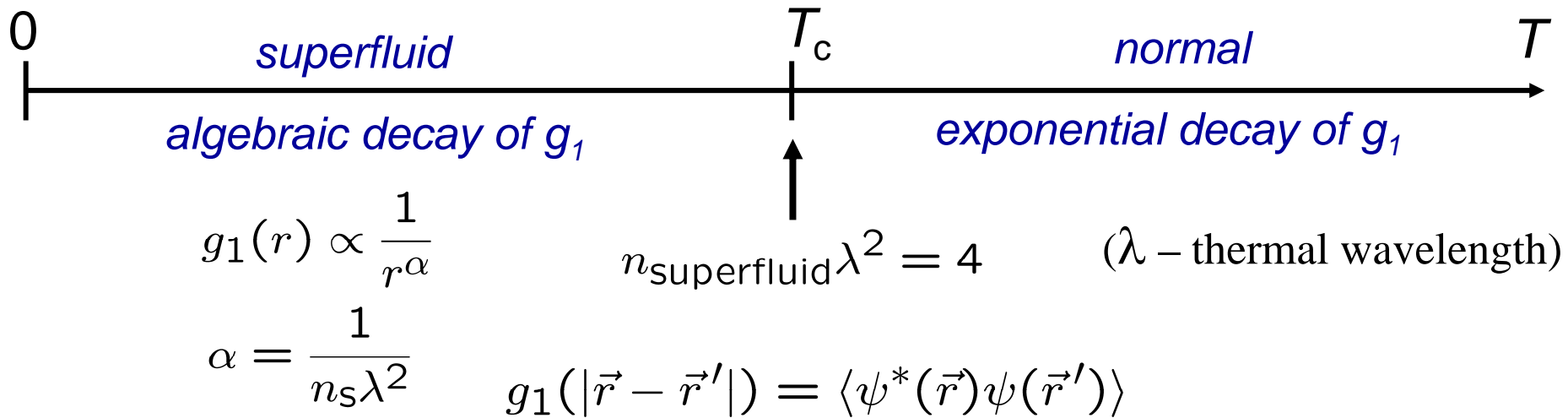
But still a superfluid transition at finite T

Bishop and Reppy (1978), superfluidity in liquid He films :



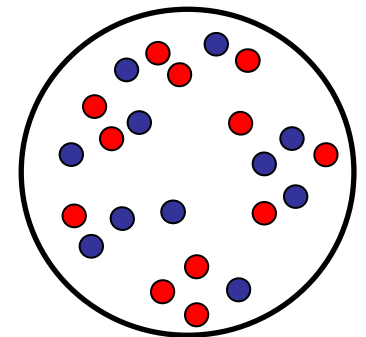
Berezinskii & Kosterlitz – Thouless (1971-73)

Phase transition without spontaneous symmetry breaking



Bound vortex-antivortex pairs

Unbinding of
vortex pairs



Proliferation of free vortices

(Ideal gas) In a harmonic trap...

Homogeneous system:

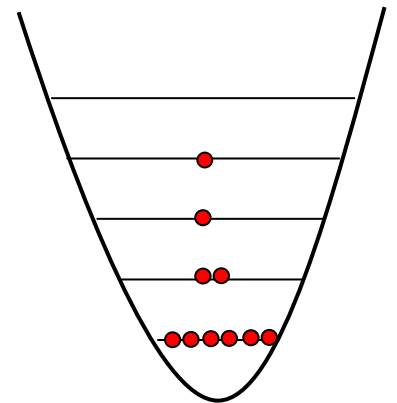
3D: BEC occurs when the phase space density reaches $n\lambda^3 = 2.6$

2D: no BEC for any phase space density $n\lambda^2$

In a harmonic trap:

3D: BEC occurs when $N = 1.2 \left(\frac{k_B T}{\hbar\omega} \right)^3$

2D: BEC occurs when $N = 1.6 \left(\frac{k_B T}{\hbar\omega} \right)^2$



Does harmonic trapping make 2D boring?

What about interactions?

The effect of (weak) interactions on BEC

3D harmonic trap:

Repulsive interactions slightly decrease the central density, for given N and T

For an ideal gas, the central density at condensation point is:

$$N = 1.2 \left(\frac{k_B T}{\hbar \omega} \right)^3 \longleftrightarrow n(0) \lambda^3 = 2.6 \quad (\text{semi-classical})$$

Just put in a bit more atoms to obtain the needed $n(0)$

2D harmonic trap:

The same procedure completely fails:

$$N = 1.6 \left(\frac{k_B T}{\hbar \omega} \right)^2 \longleftrightarrow n(0) \lambda^2 = \infty$$

where

$$n(r) = \int \rho(r, p) \frac{d^2 p}{h^2} \quad \rho(r, p) = \left[e^{\beta \left(\frac{p^2}{2m} + \frac{m\omega^2 r^2}{2} \right)} - 1 \right]^{-1}$$

The effect of (weak) interactions on BEC

Treat the interactions at the mean field level: $V_{\text{eff}}(r) = \frac{m\omega^2 r^2}{2} + 2gn_{\text{mf}}(r)$

where the mean field density is obtained from the self-consistent equation

$$n_{\text{mf}}(r) = \int \rho_{\text{mf}}(r, p) \frac{d^2 p}{h^2} \quad \rho_{\text{mf}}(r, p) = \left[e^{\beta(\frac{p^2}{2m} + V_{\text{eff}}(r))} - 1 \right]^{-1}$$

Two remarkable results

- One can accommodate an arbitrarily large atom number. Badhuri et al
- The effective frequency deduced from $V_{\text{eff}}(r) \simeq m\omega_{\text{eff}}^2 r^2 / 2$ tends to zero when $\mu \rightarrow 2gn_{\text{mf}}(0)$ Holzmann et al.

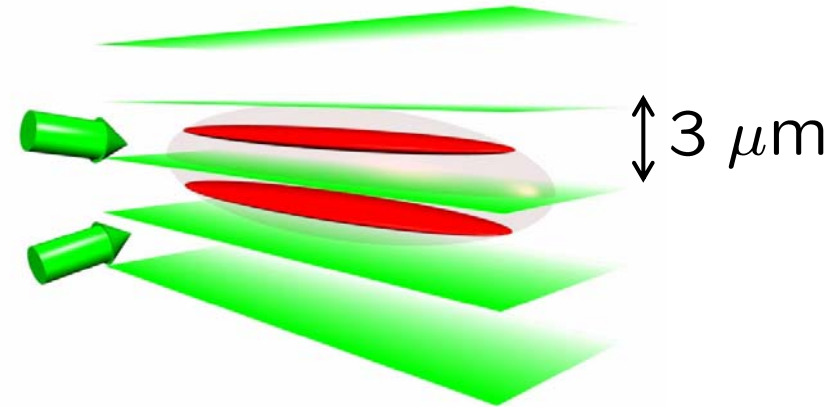
Similar to a 2D gas in a flat potential...
...BEC suppressed, expect BKT (?)

How to make an ultracold 2D Bose gas

3D BEC + 1D optical lattice

2 independent 2D clouds
(no tunnelling)

10^5 atoms/plane
plane thickness: $0.2 \mu\text{m}$, separation: $3 \mu\text{m}$



(other 2D experiments at MIT, Innsbruck, Oxford, Florence, NIST, Heidelberg etc.)

Why 2 planes?

Crucial info in the *phase* of Ψ , and $g_1(x, y) = \langle \psi^*(x, y) \psi(0) \rangle$

→ accessible in an interference experiment

2.

Critical point of an interacting 2D Bose gas

P. Krüger, Z. H. and J. Dalibard, cond-mat/0703200

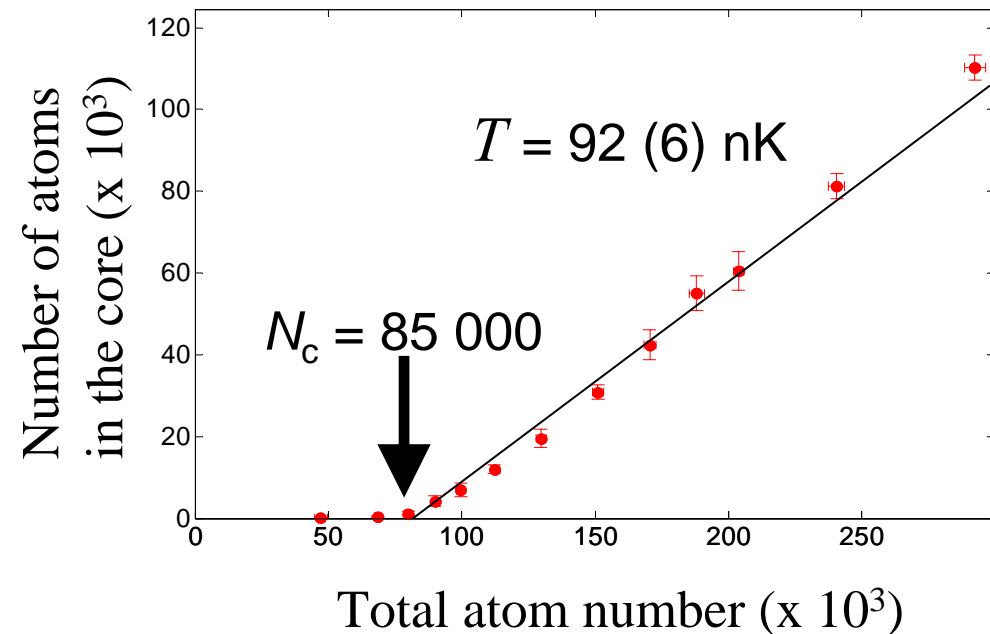
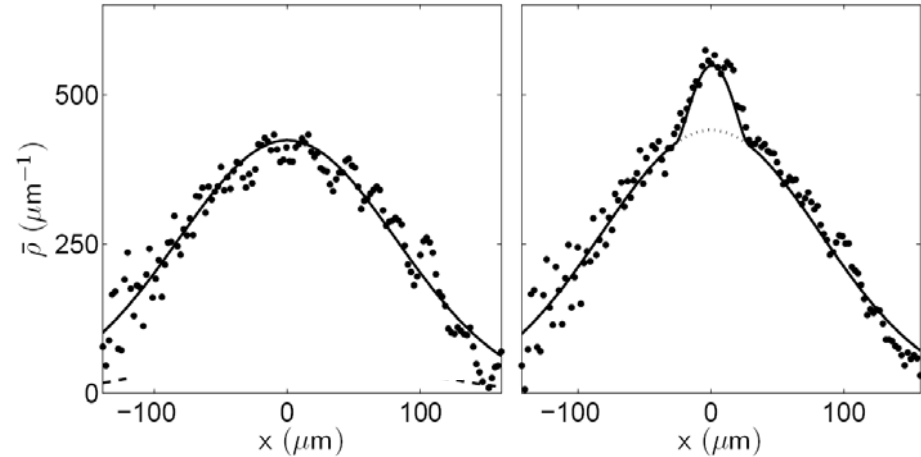
Phase transition in a 2D atomic gas

Fix the temperature T

Vary the atom number N

Bimodal distribution

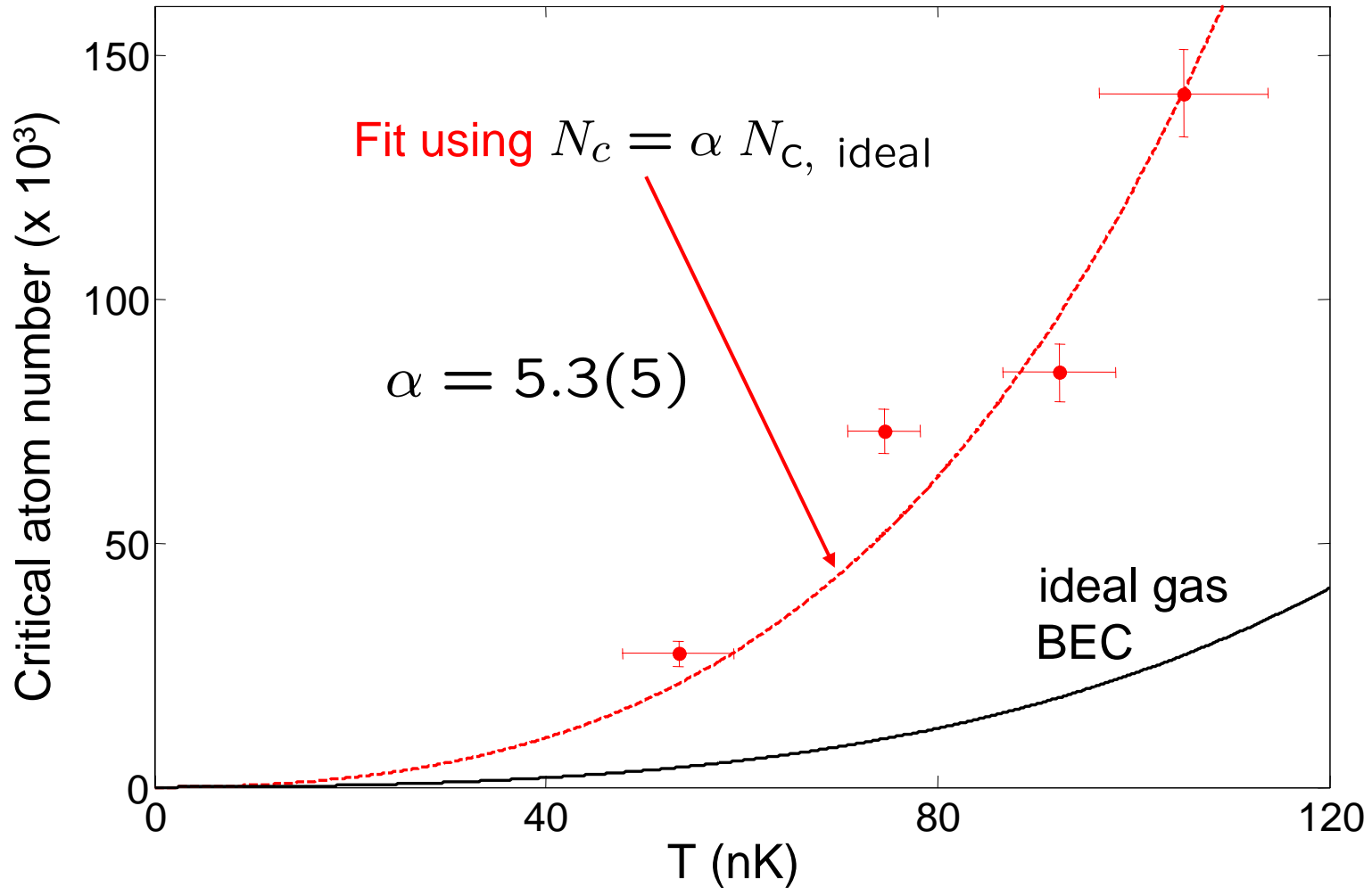
for $N > N_c$



Similar signature to 3D BEC

Dense core follows the Thomas-Fermi law in time-of-flight expansion, characteristic of superfluid hydrodynamics

Critical atom number vs. T



5.3 times larger than the ideal gas BEC prediction!

Can it be the Kosterlitz-Thouless critical point?

$n_{\text{superfluid}}\lambda^2 = 4$ is universal and elegant, but not the whole story

Total critical density depends on microscopics (long standing problem!)

Fisher & Hohenberg + Prokof'ev et al.:

$$n_{\text{total}}\lambda^2 = \ln\left(\frac{C}{\bar{g}}\right)$$

$$C = 380 \pm 3$$

$$\bar{g} = \frac{mg}{\hbar^2}$$

dimensionless
interaction strength

For our setup:

$$\bar{g} = 0.13$$



$$n_{\text{total}}\lambda^2 \simeq 8.0$$

Extract from the experiment:

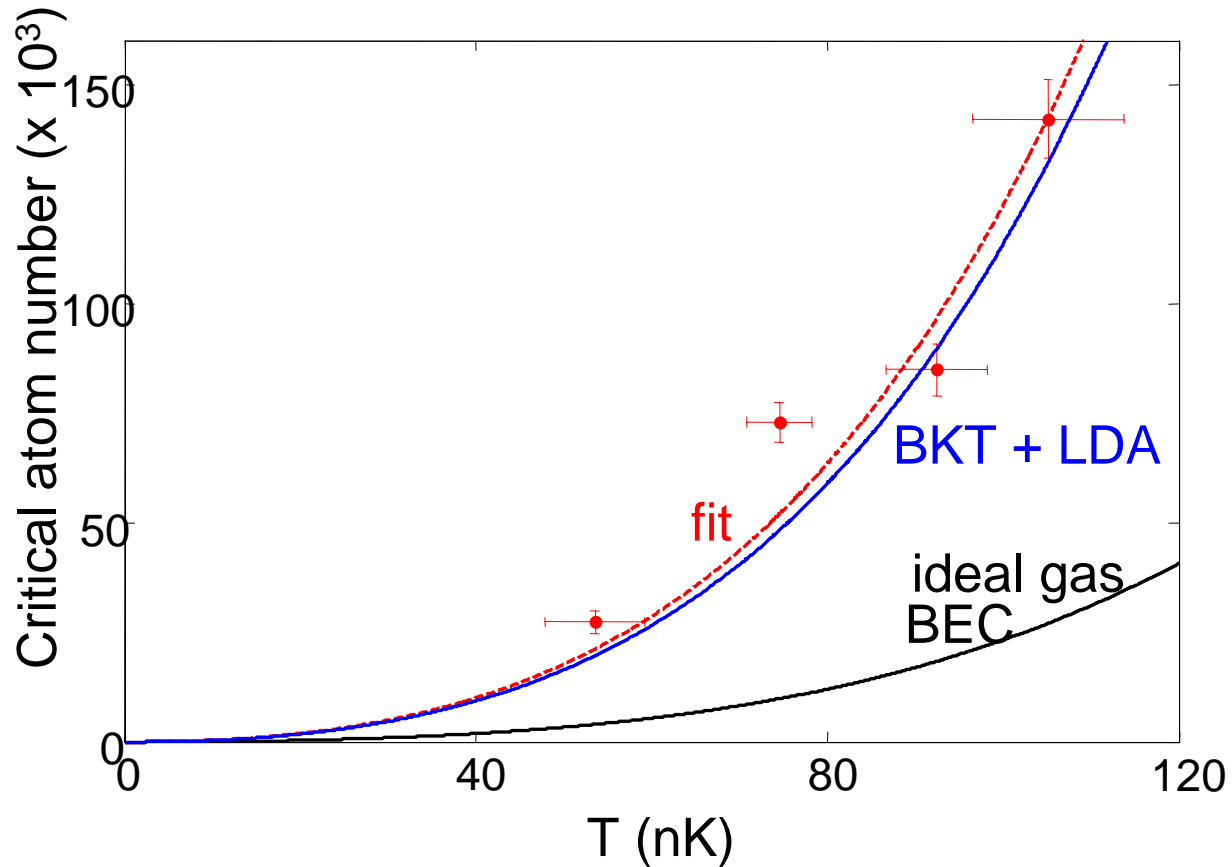
$$n_c\lambda^2 = 8.6 \pm 0.8$$

(in the center of the cloud)

Critical atom number vs. T

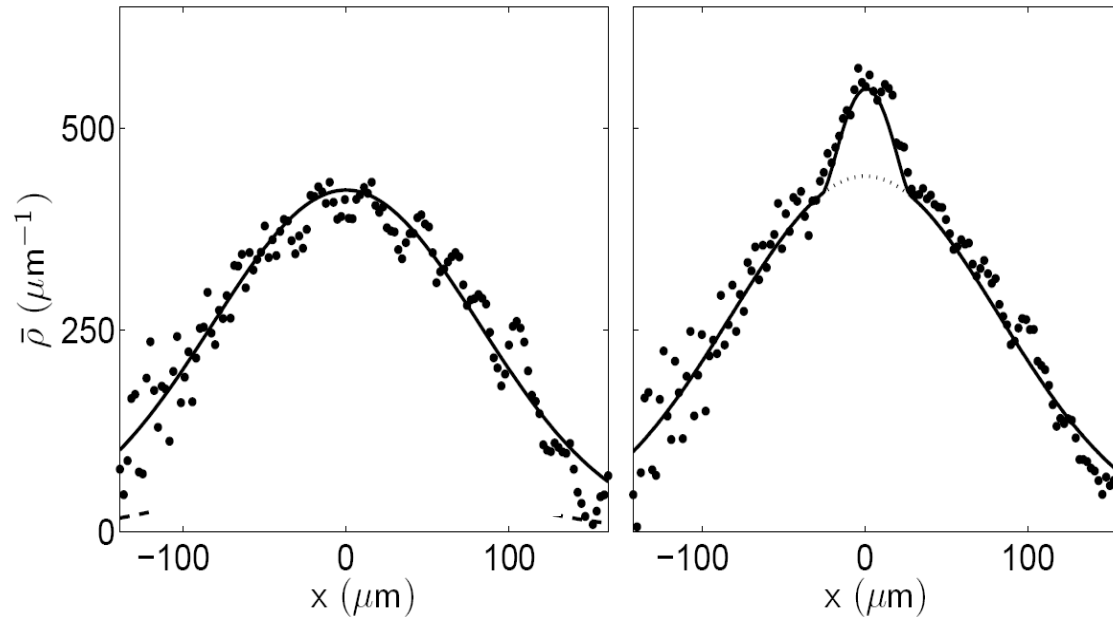
BKT + LDA + experimentally observed Gaussian profiles:

$$N_{C, \text{KT}} = 4.9 N_{C, \text{ideal}}$$



Not bad...

Equation of state?



Bimodal distribution fitted well by **Gaussian** + **Thomas-Fermi**

...but why?

3.

Coherence of an interacting 2D Bose gas

Z. H., P. Krüger, M. Cheneau, B. Battelier, S. Stock, and J. Dalibard

Phys. Rev. Lett. **95**, 190403 (2005)

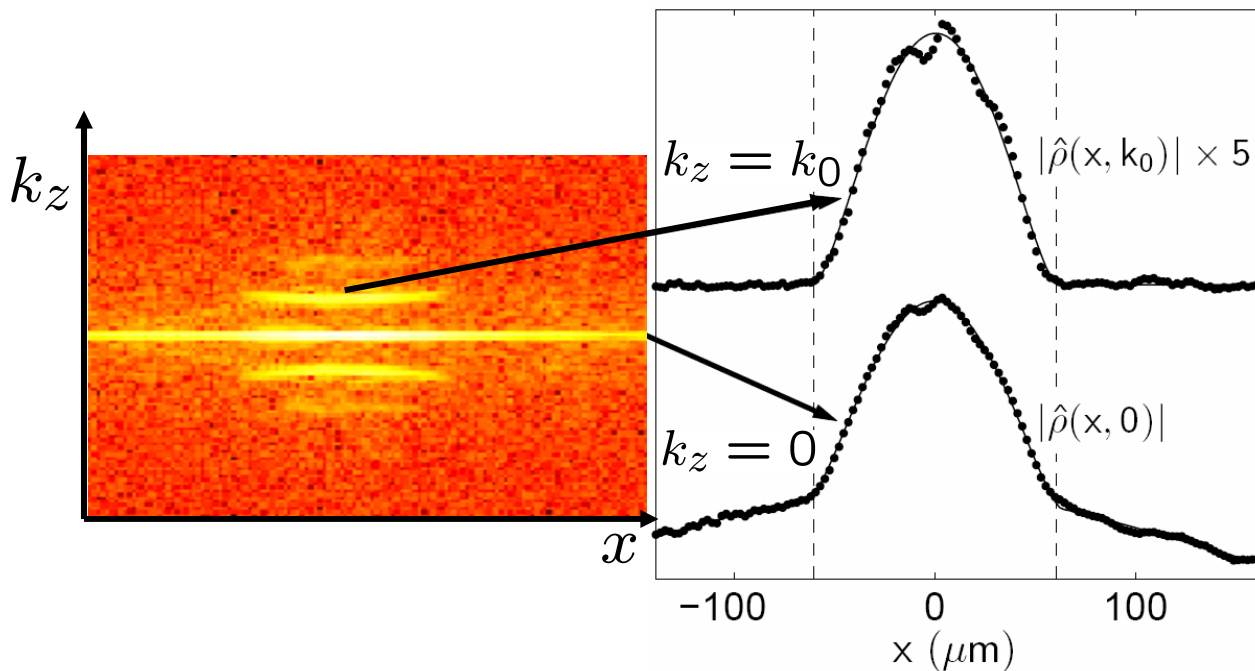
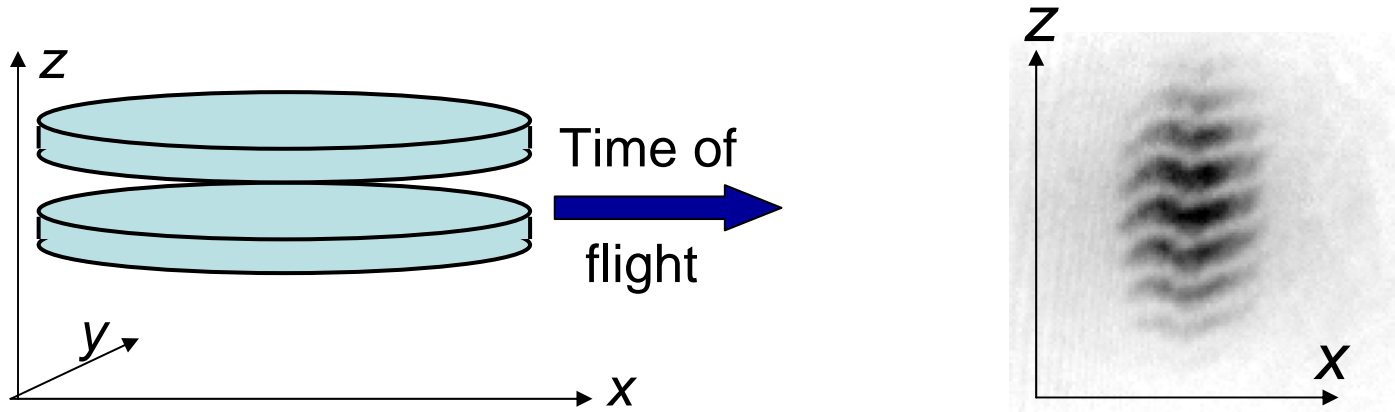
Nature **441**, 1118 (2006)

cond-mat/0703200

+ Schweikhard, Tung and Cornell, cond-mat/0704.0289

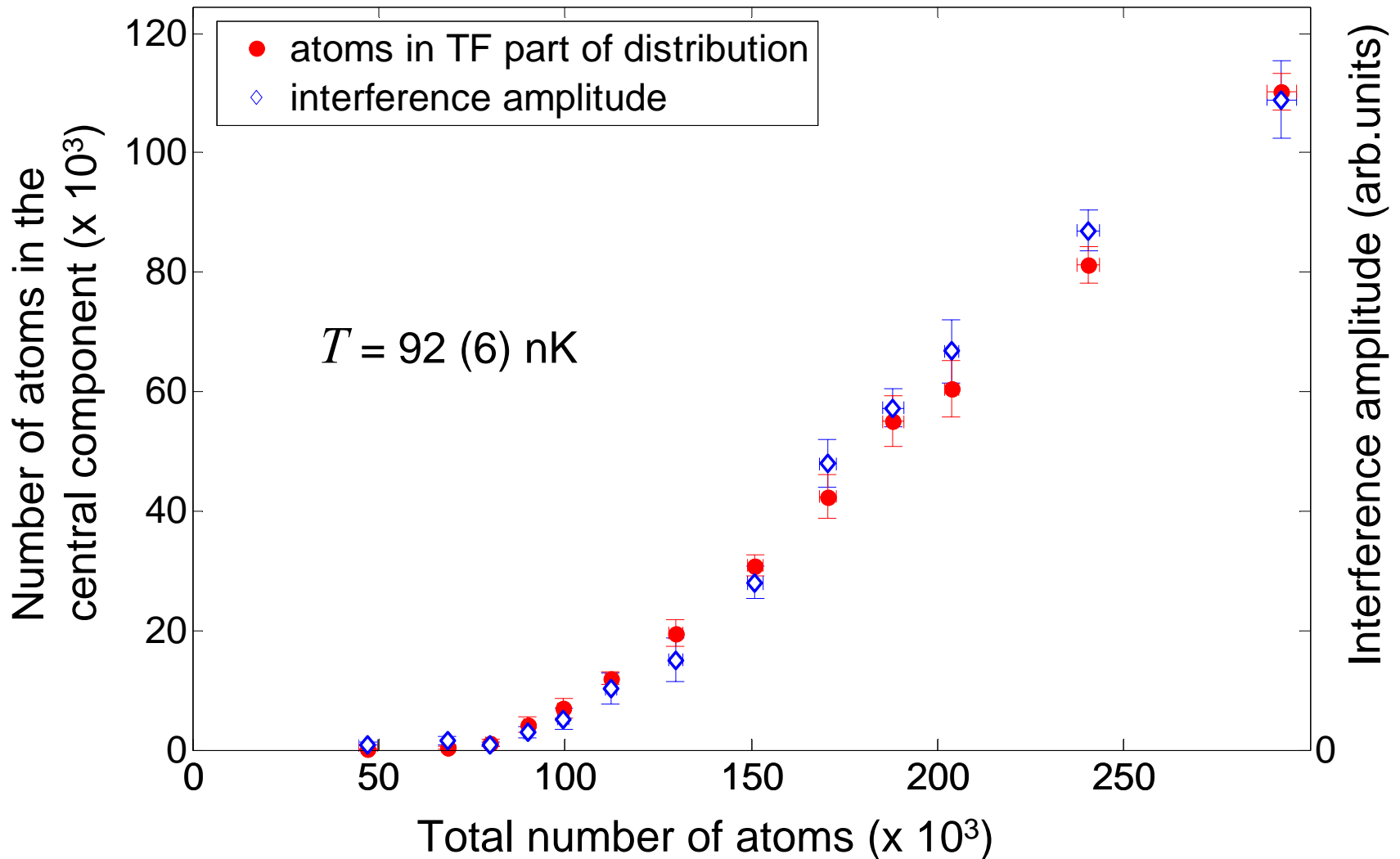
Theory: Shlyapnikov-Gangardt-Petrov, Holtzman *et al.*, Kagan *et al.*,
Stoof *et al.*, Mullin *et al.*, Simula-Blackie, Hutchinson *et al.*
Polkovnikov-Altman-Demler

Interference of two 2D gases



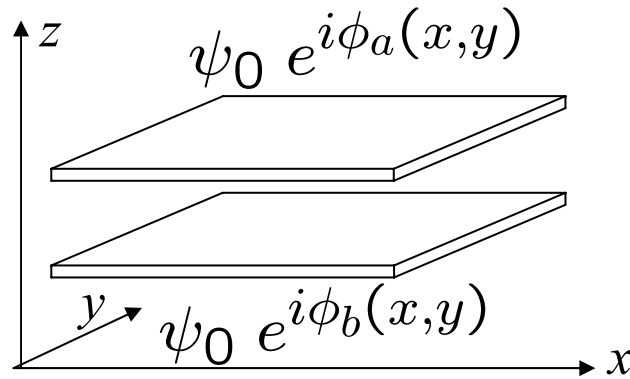
The interfering part coincides with the central part of the bimodal distribution

Bimodality and interferences

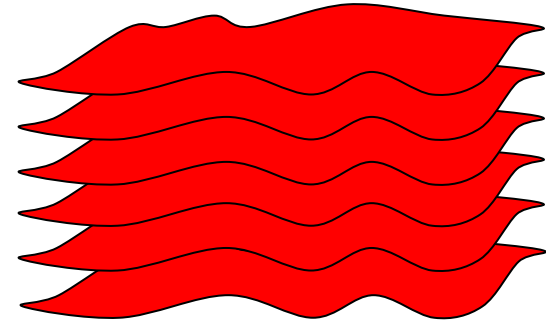



Within our accuracy, onsets of bimodality and interference coincide

Local vs. long-range coherence



Time of
flight

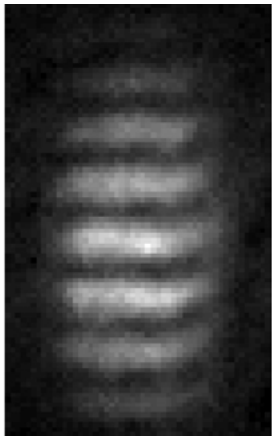


Phonons (“spin waves”)

→ smooth phase variations

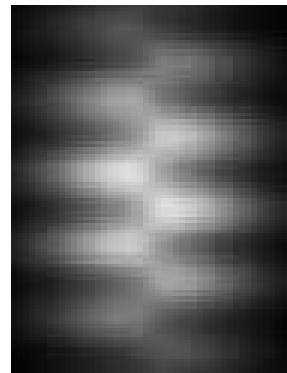
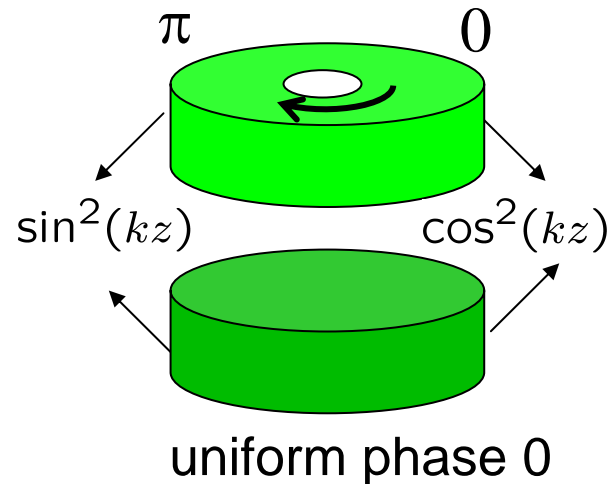
cold

hot



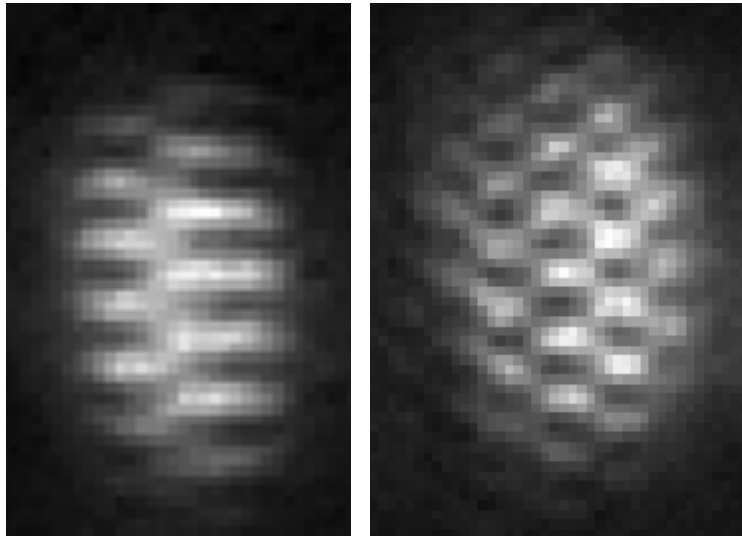
Vortices

→ sharp dislocations

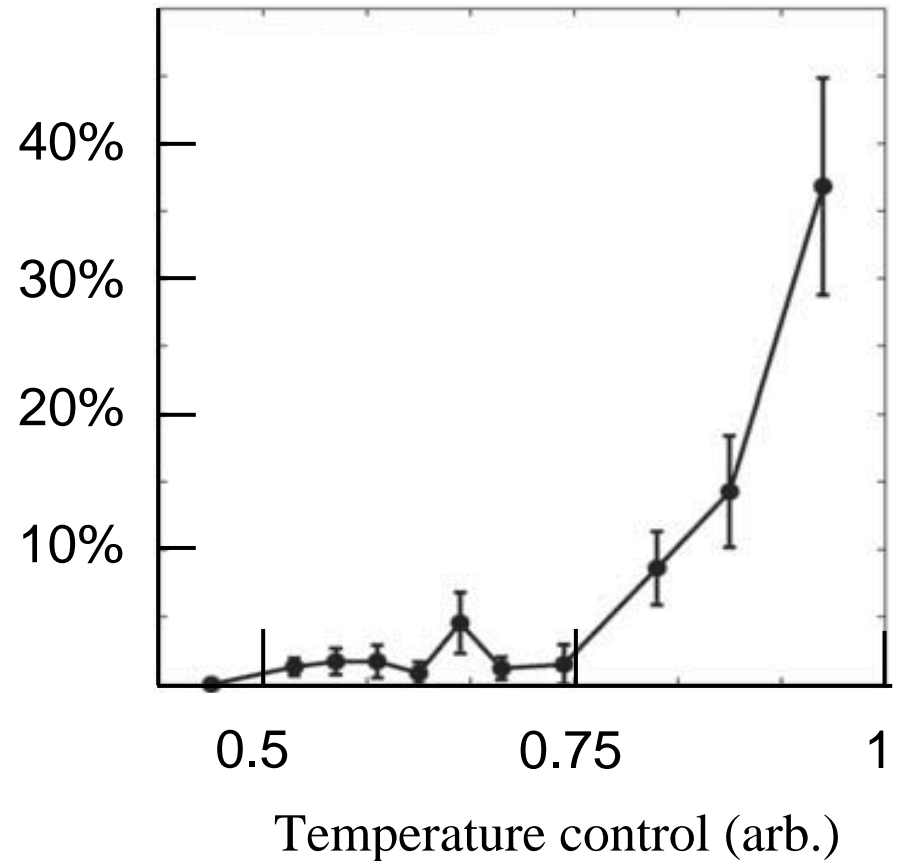


Free vortices in 2D clouds

Fraction of images showing at least one dislocation in the central region:



(Similar results at NIST)

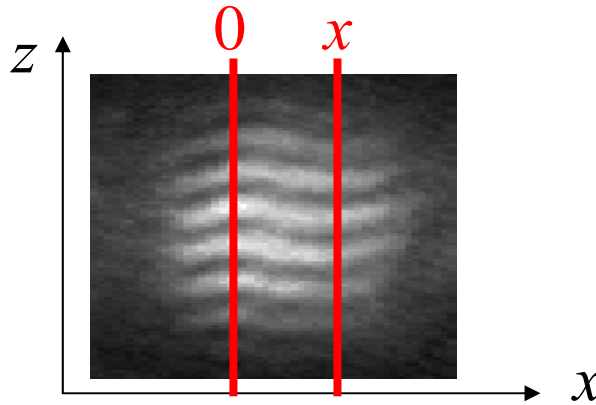


Long-range coherence

$$\text{Embedded in: } g_1(\vec{r}) = \langle \psi^*(\vec{r}) \psi(0) \rangle$$

The interference signal between $\psi_a(x, y)$ and $\psi_b(x, y)$ gives:

$$|\psi_a|^2 + |\psi_b|^2 + \underbrace{\psi_a^* \psi_b}_{\kappa(x, y)} e^{ikz} + \psi_a \psi_b^* e^{-ikz}$$



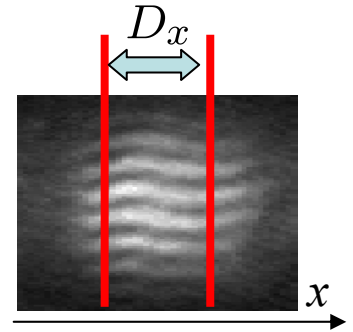
$$\begin{aligned} \langle \kappa(x, y) \kappa(0) \rangle &= \langle \psi_a^*(x, y) \psi_b(x, y) \psi_a(0) \psi_b^*(0) \rangle \\ &= \langle \psi_a^*(x, y) \psi_a(0) \rangle \langle \psi_b(x, y) \psi_b^*(0) \rangle = |g_1(x, y)|^2 \end{aligned}$$

Long-range coherence

Polkovnikov, Altman, Demler:

Integrated contrast: $C(D_x) = \frac{1}{D_x} \int_0^{D_x} \kappa(x) dx$

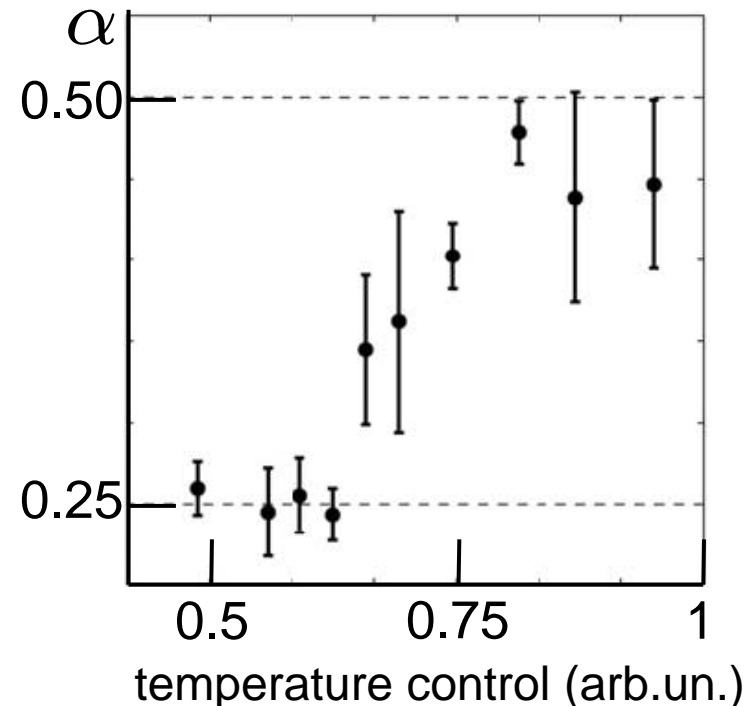
$\langle C^2(D_x) \rangle \sim \frac{1}{D_x} \int_0^{D_x} (g_1(x, 0))^2 dx$ scales as: $1/(D_x)^{2\alpha}$



"universal jump in superfluid density"

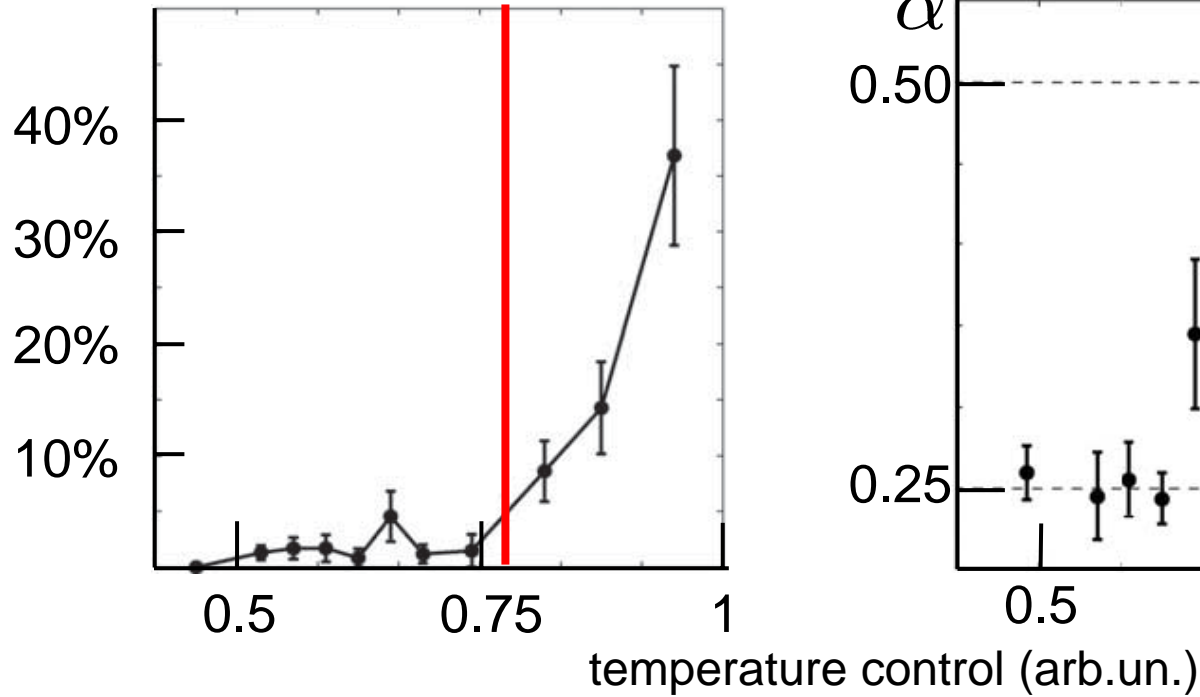
➔ drop in α from 0.5 to 0.25

(in an infinite uniform system)

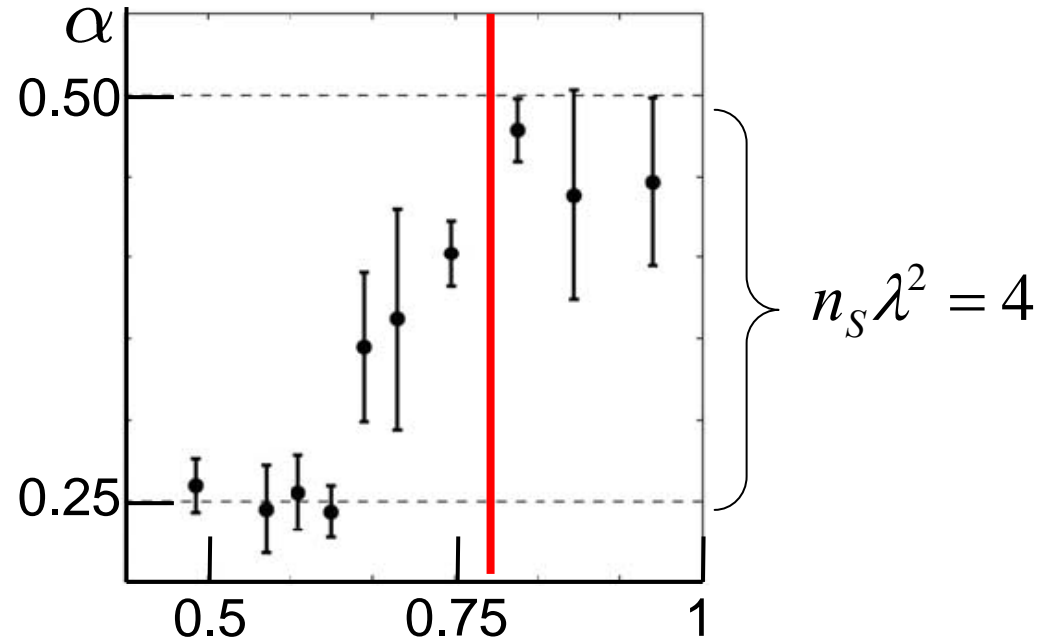


Vortices vs. Correlations vs. Temperature

vortices:



first order coherence:



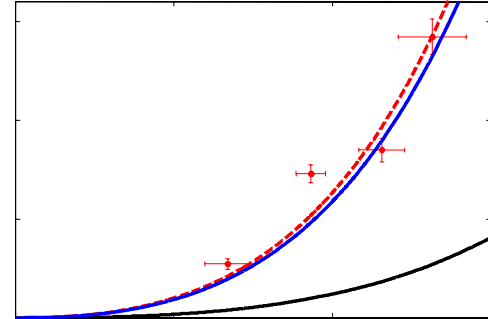
The onset of vortex proliferation coincides with the loss of quasi-LRO

Z. Hadzibabic *et al.*, Nature **441**, 1118 (2006)

see also Schweikhard, Tung and Cornell, cond-mat/0704.0289 for KT in a lattice

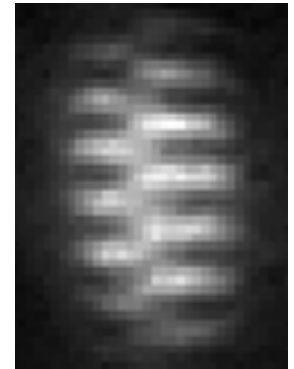
So far in atomic Flatland...

- Phase transition with a critical point (N_C, T_C) :
- eliminates conventional BEC
 - agrees quantitatively with BKT + LDA



Direct visualization of free vortices:

- coincides with loss of quasi-long-range order
- supports the microscopic basis of the theory



Open questions/future:

Equation of state?

Tune the interactions from $g \sim 1$ to $g \sim 10^{-4}$

Superfluidity – transport, dissipation?

Resolve tightly bound vortex pairs in the superfluid state?

THE END