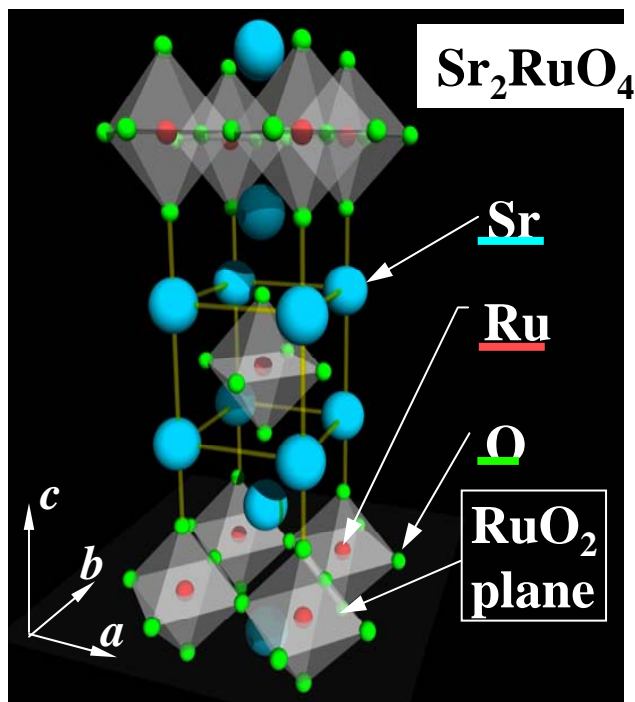




# Knight Shift Measurements on Superconducting $\text{Sr}_2\text{RuO}_4$



**Layered Perovskite structure**

**Maeno *et al.* Nature 372, 532 ('94)**

**K. Ishida<sup>A,B</sup>, H. Murakawa,<sup>A</sup>**

**H. Mukuda,<sup>B</sup> Y. Kitaoka,<sup>B</sup>**

**Z. Q. Mao,<sup>A\*</sup> Y. Maeno,<sup>A</sup>**

**A *Dept. of Physics, Kyoto Univ.***

**B *Dept. of Material Physics,***

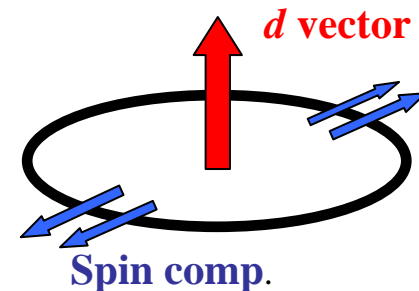
***Osaka Univ.***

***\*Tulane Univ. ( USA)***

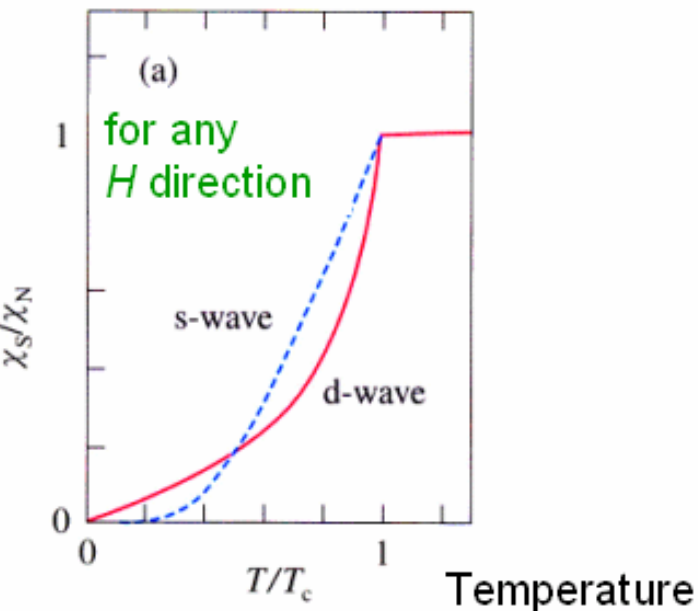
# Superconducting wave function

$$\Psi(\mathbf{r}_1, \sigma_1; \mathbf{r}_2, \sigma_2) = \chi(\sigma_1, \sigma_2) \psi(\mathbf{r}_1, \mathbf{r}_2)$$

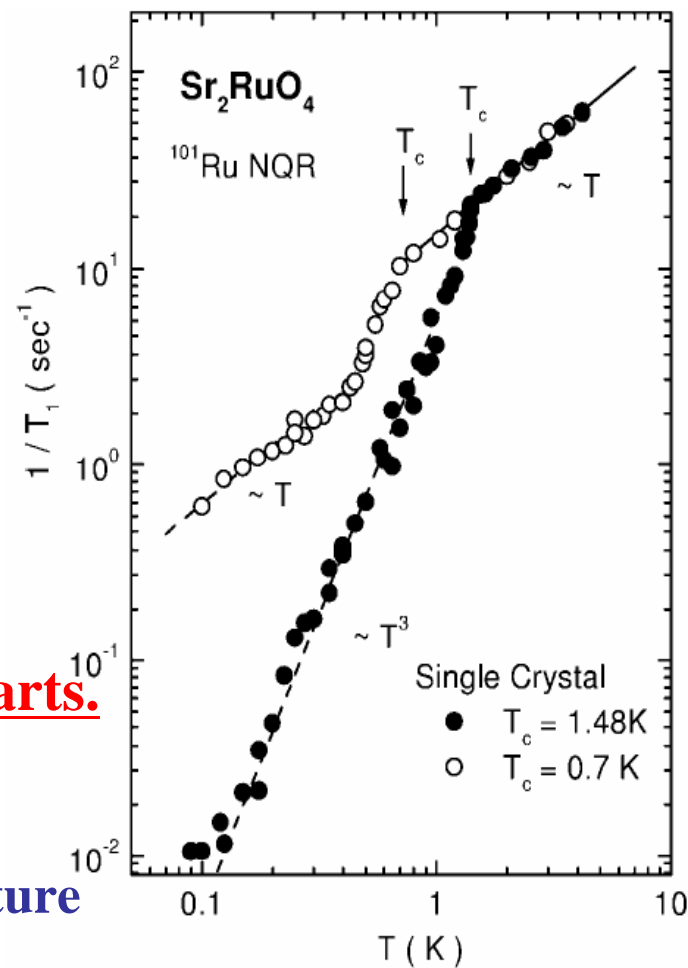
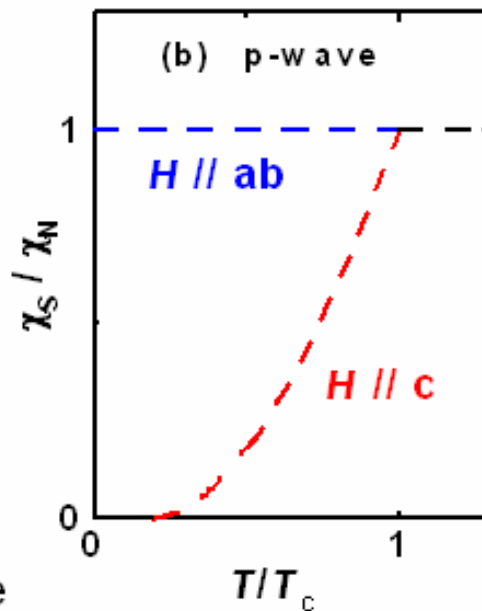
**Spin part**      **Orbital part**



Singlet



Triplet



**NMR gives important information about both parts.**

**Knight shift measurements : Spin susceptibility**

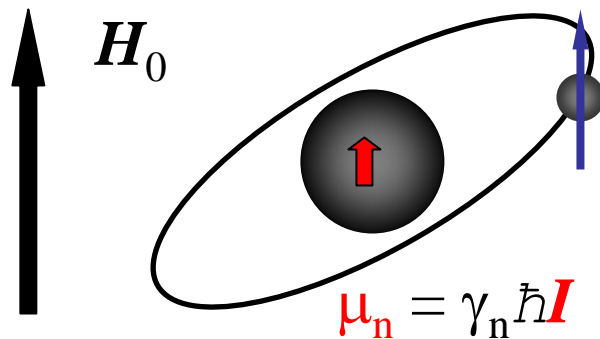
**Nuclear spin-lattice relaxation rate  $1/T_1$  : Gap structure**

# Knight-Shift Measurement

( Hyperfine Interaction )

## Interaction between nuclear spin and electronic spins

Nuclear spins ( $\mathbf{I}$ ) are coupled with electronic spins ( $\mathbf{S}$ ) in the external magnetic field ( $\mathbf{H}_0$ ).



$$\mu_e = \gamma_e \hbar \mathbf{S}$$

Hamiltonian of the nuclear spins

$$H_Z = -\gamma_n \hbar \mathbf{I} \mathbf{H}_0 + \mathbf{A} \mathbf{I} \mathbf{S} = -\gamma_n \hbar \mathbf{I} \mathbf{H}_{\text{Loc}}$$

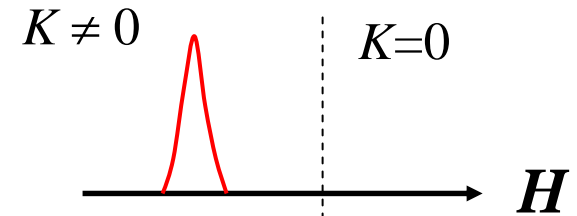
$\mathbf{A}$ : Hyperfine coupling constant

$$\mathbf{H}_{\text{Loc}} = \mathbf{H}_0 - \frac{\mathbf{A}}{\gamma_n \hbar} \mathbf{S} = \mathbf{H}_0 - \frac{\mathbf{A}}{\gamma_n \hbar} \{ \underbrace{\langle \mathbf{S} \rangle}_{\text{av. static}} + \underbrace{\delta \mathbf{S}}_{\text{dynamics}} \}$$

Shift of the spectrum (Static properties)

$$K = \frac{H_0 - \langle H_{\text{Loc}} \rangle}{H_0} = \frac{A}{\gamma_n \hbar} \cdot \frac{\langle S \rangle}{H_0} = \frac{A}{N_A \gamma_e \gamma_n \hbar^2} \chi$$

$$\left( \chi = \frac{M}{H_0} = \frac{N_A \gamma_e \hbar \langle S \rangle}{H_0} \right)$$



# Example of $K$ - $\chi$ plot

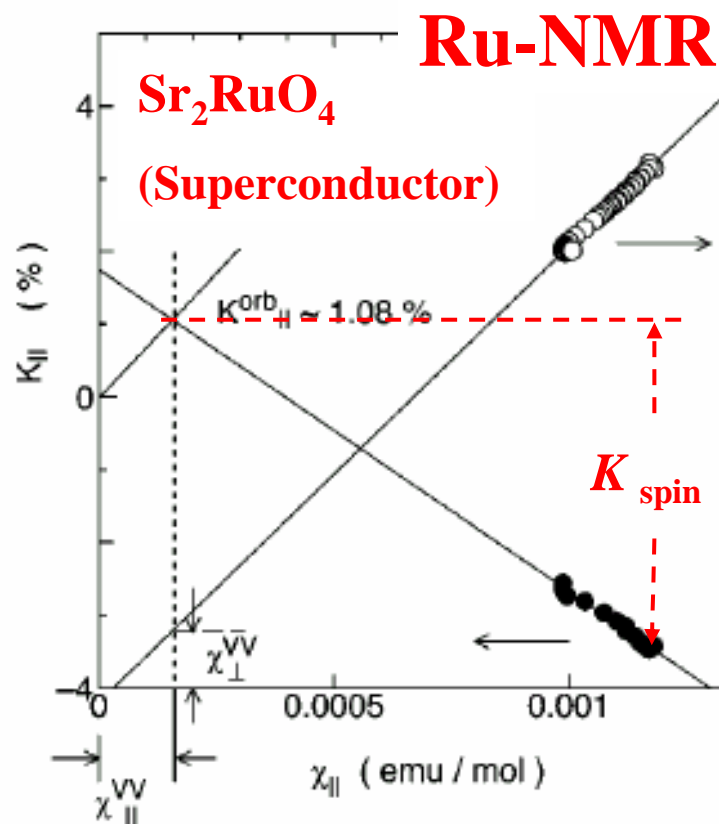


FIG. 2. Plot of  $^{99}\text{K}_{\parallel}$  (Ref. 18) against  $\chi_{\parallel}$  (left scale). The relation between  $\chi_{\perp}$  and  $\chi_{\parallel}$  are also plotted (right scale).

$$K_{\text{spin}} = \frac{A_{\text{hf}}}{\mu_{\text{B}} N_{\text{A}}} \chi_{\text{spin}} \quad \sim -4.5\%$$

$$K_{\text{orb}} = \frac{A_{\text{VV}}}{\mu_{\text{B}} N_{\text{A}}} \chi_{\text{VV}} \quad \sim 1\%$$

## $\text{CaRuO}_3$ (Nearly Ferro. Metal)

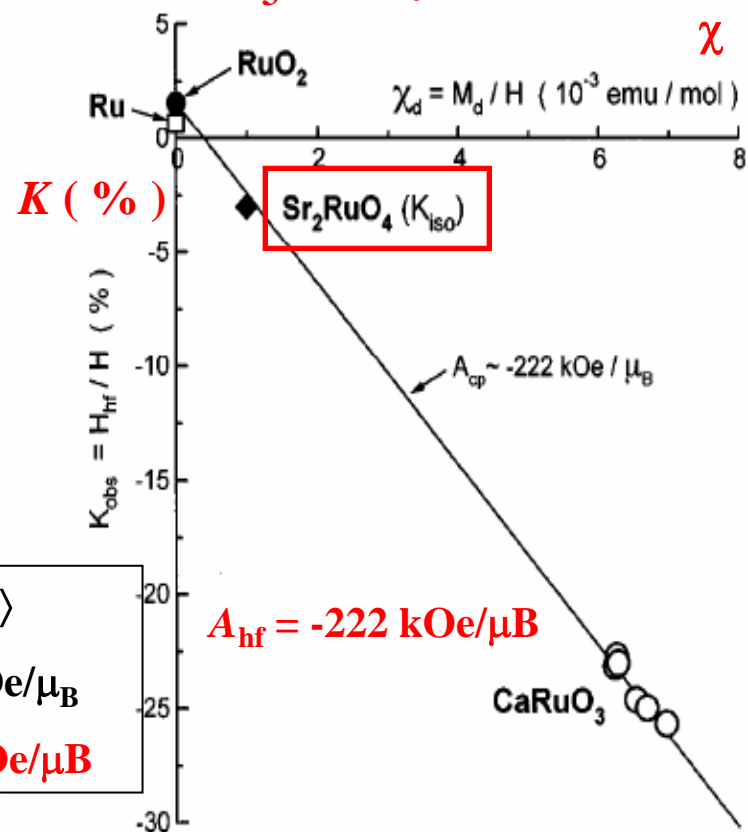
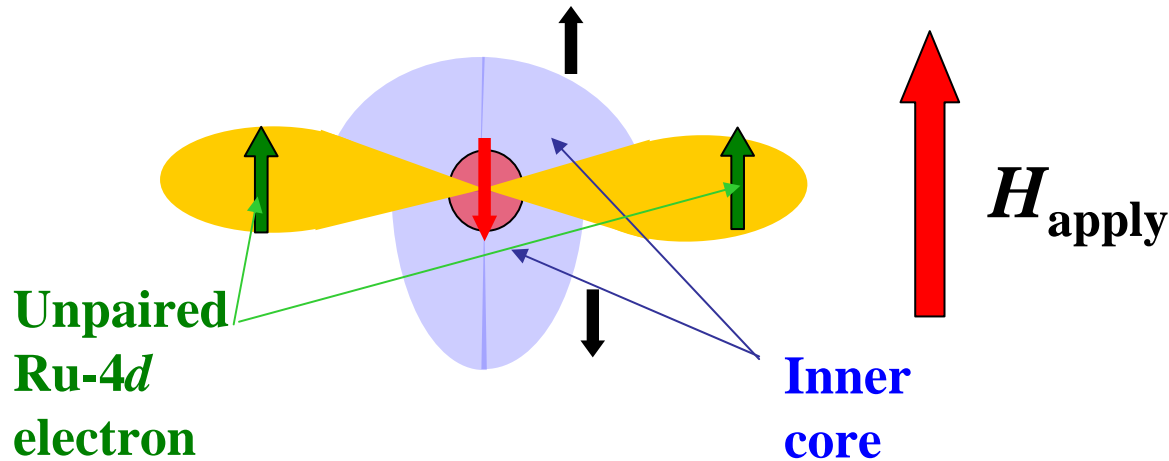


FIG. 5. Knight shift vs susceptibility  $[M_{\text{d}}(H)/H]$  with an implicit parameter of the external field  $H$ . The hyperfine coupling constant due to the inner core polarization  $A_{\text{cp}} \approx -222 \pm 50 \text{ kOe}/\mu_{\text{B}}$  is estimated from a slope of linear line when  $K_{5s} + K_{\text{orb}}$  is assumed to be the same value as  $K_{\text{obs}} = +1.59\%$  in  $\text{RuO}_2$ . This value is close to  $A_{\text{hf}} \sim -300 \pm 60 \text{ kOe}/\mu_{\text{B}}$  in the FM state of  $\text{SrRuO}_3$ . Note that the Knight shift for  $\text{Sr}_2\text{RuO}_4$  lies on the same line of  $K_{\text{obs}}$  vs  $\chi_{\text{obs}}$  plot as that for  $\text{CaRuO}_3$ .

**Core polarization effect**

# Inner Core Polarization Effect by Unpaired $d$ electrons



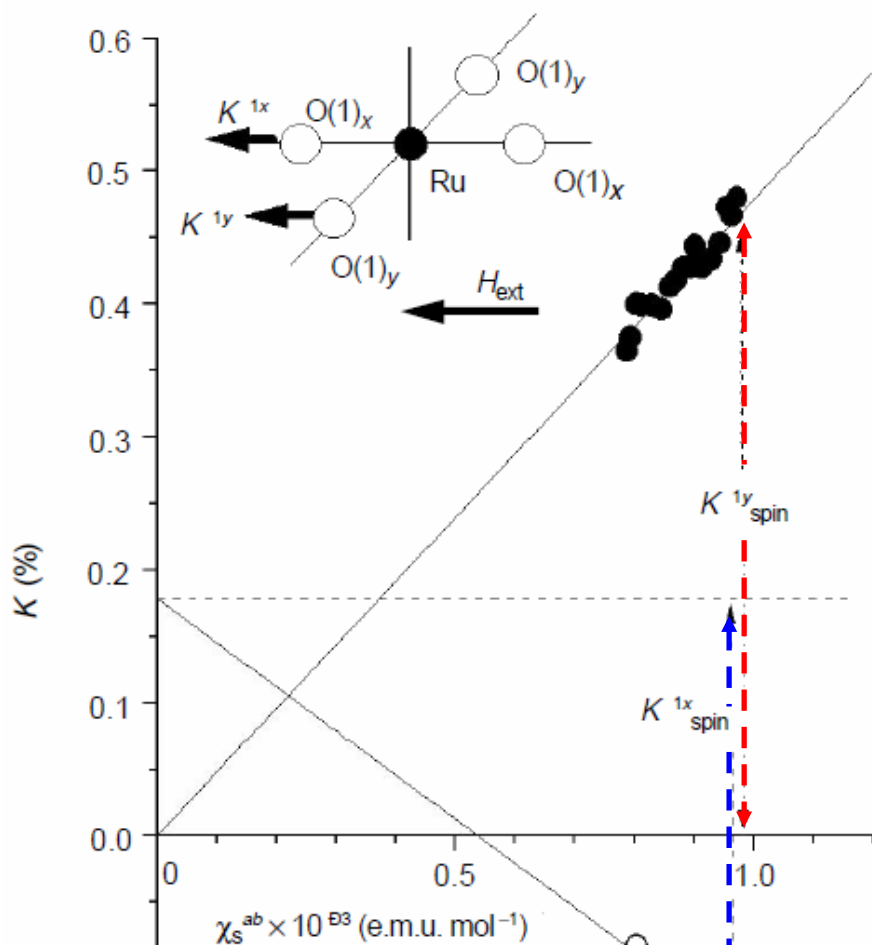
$$H_{\text{nul.}} = \sum_i \left\{ |\psi_i(0)|_{\uparrow}^2 - |\psi_i(0)|_{\downarrow}^2 \right\}$$

**Negative isotropic field is induced**

$$K_{\text{CP}} = \frac{8\pi}{3} \chi_d \sum_i \left\{ |\psi_i(0)|_{\uparrow}^2 - |\psi_i(0)|_{\downarrow}^2 \right\}$$

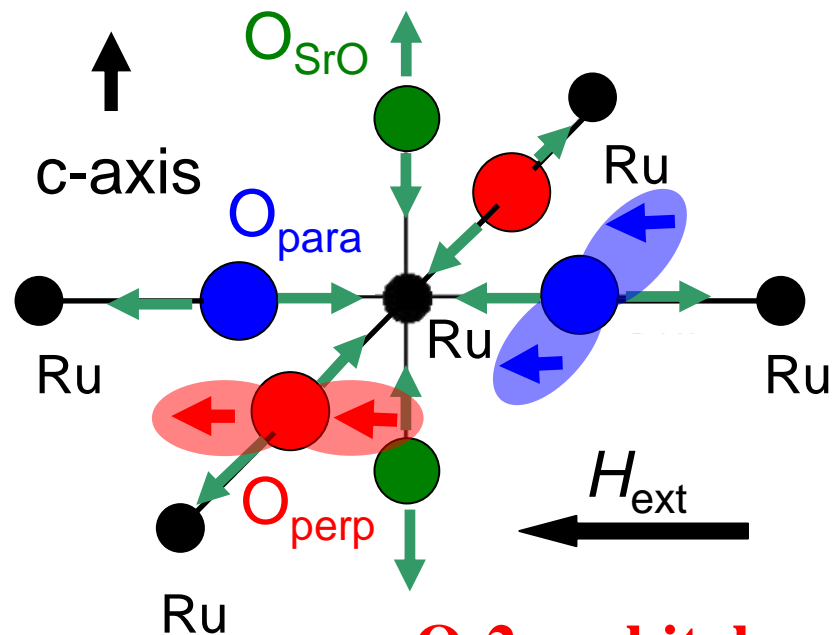
	$A_{\text{hf}}$
$\text{Sr}_2\text{RuO}_4$ (SC.)	$\sim -250 \text{ kOe}/\mu_{\text{B}}$
<b>c.f. <math>\text{SrRuO}_3</math></b> <b>(Ferro.)</b>	$H_{\text{int}} / M_0$ $\sim -300 \text{ kOe}/\mu_{\text{B}}$
$\text{CaRuO}_3$ (NearlyFerro.)	$\sim -222 \text{ kOe}/\mu_{\text{B}}$

**Ru-4d state is in the Fermi level in these Ruthenates**

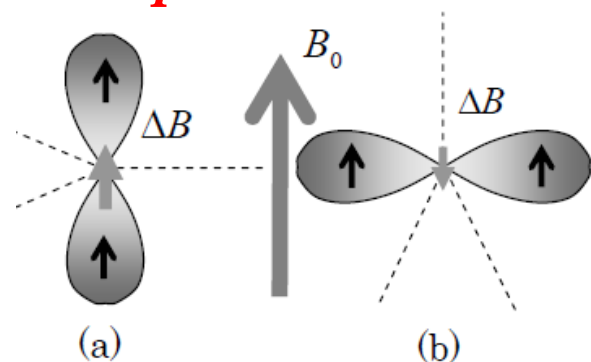


$$K_{spin}^x \sim -0.3 \%$$

$$K_{spin}^y \sim 0.5 \%$$



**O-2p orbitals**



**O-2p orbitals hybridized with Ru-4d<sub>xy</sub> orbitals are important.**

# $^{17}\text{O}$ Knight Shift measurement

K. Ishida *et al.* Nature 395, 658 (98)

Mukuda *et al.* J. Low Temp. Phys. 117, 1567 (99)

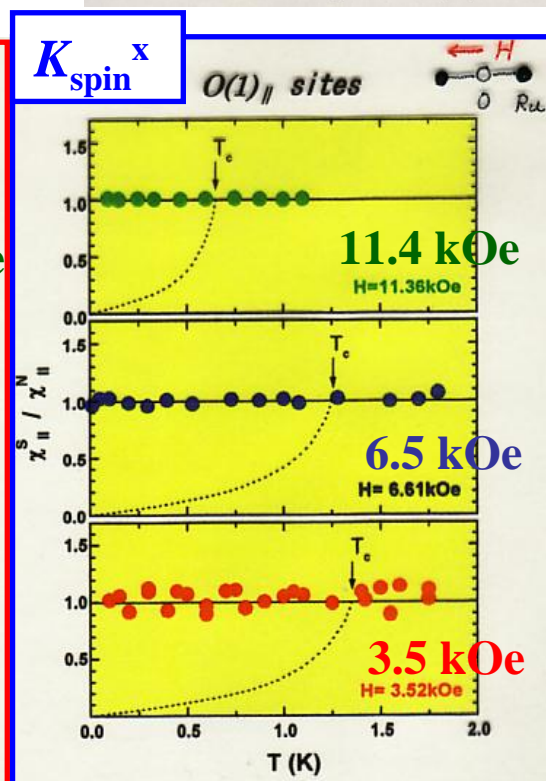
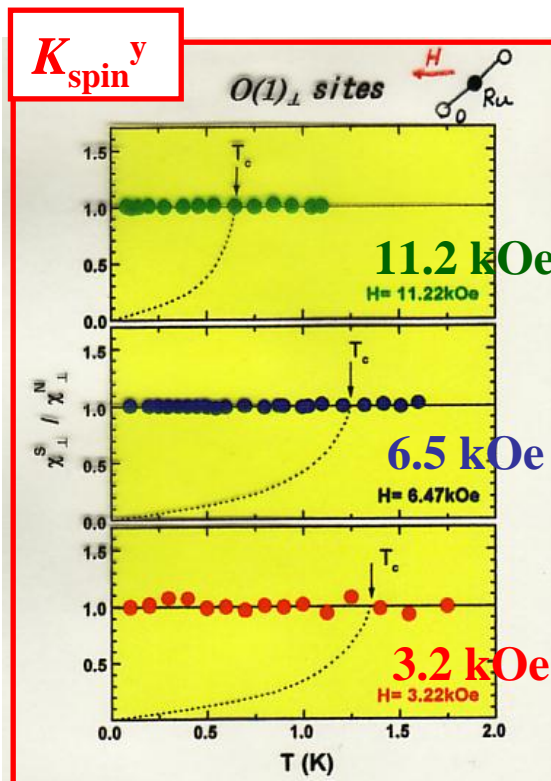
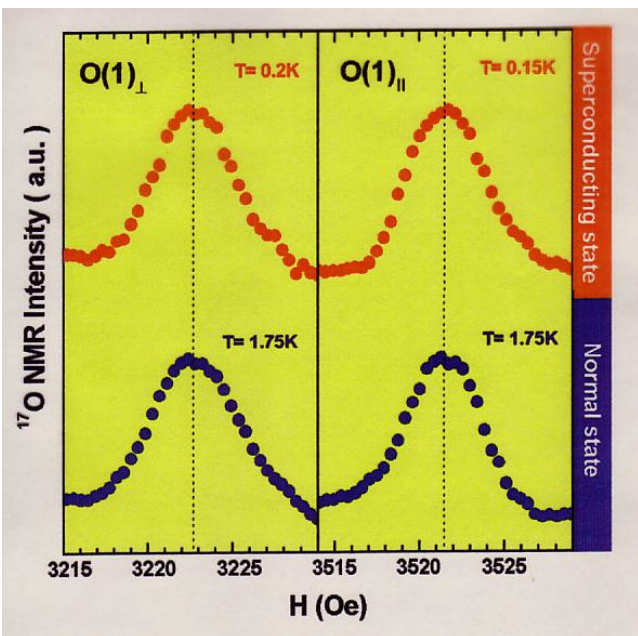
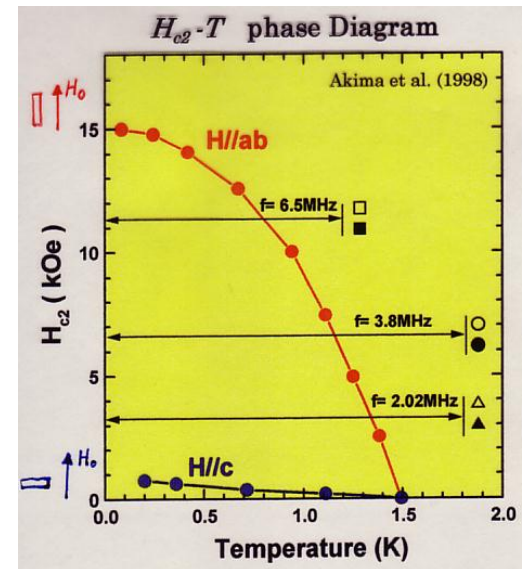
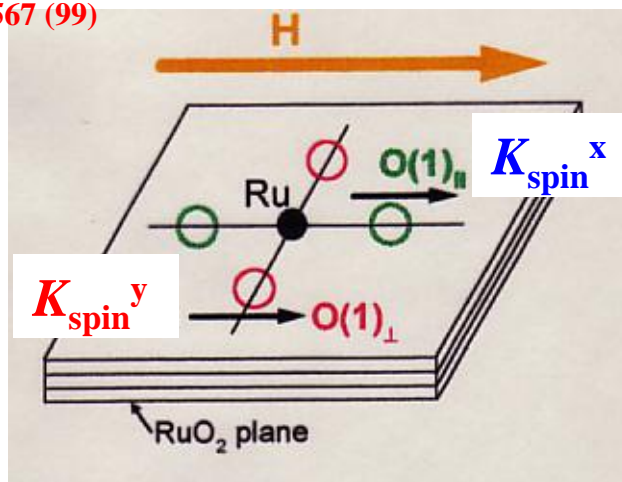
## Spin part in $K$

$$K_{\text{spin}}^x = -0.3 \%$$

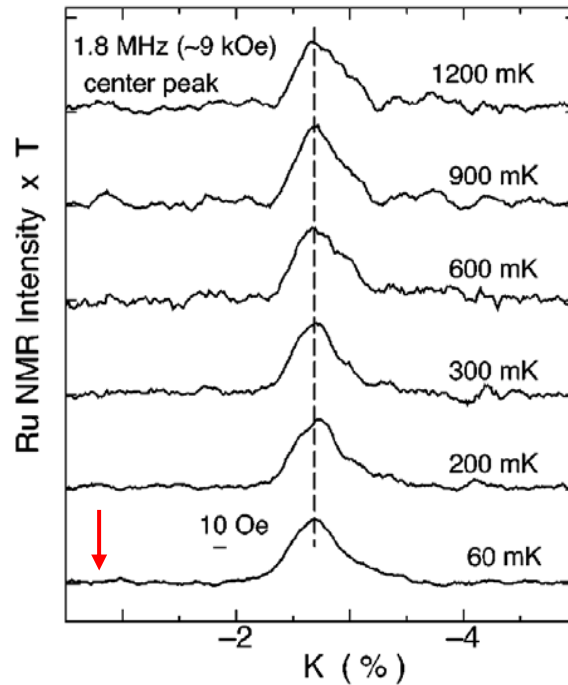
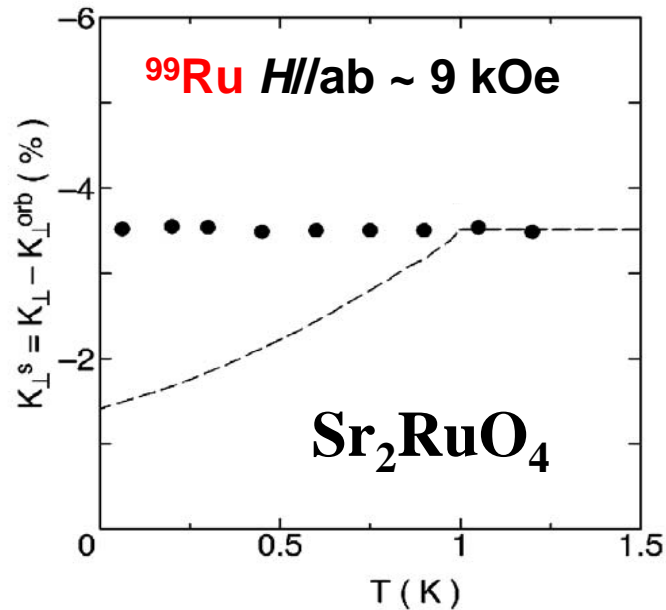
$$K_{\text{spin}}^y = 0.5 \%$$

$H // ab \sim 3.5 \text{ kOe}$

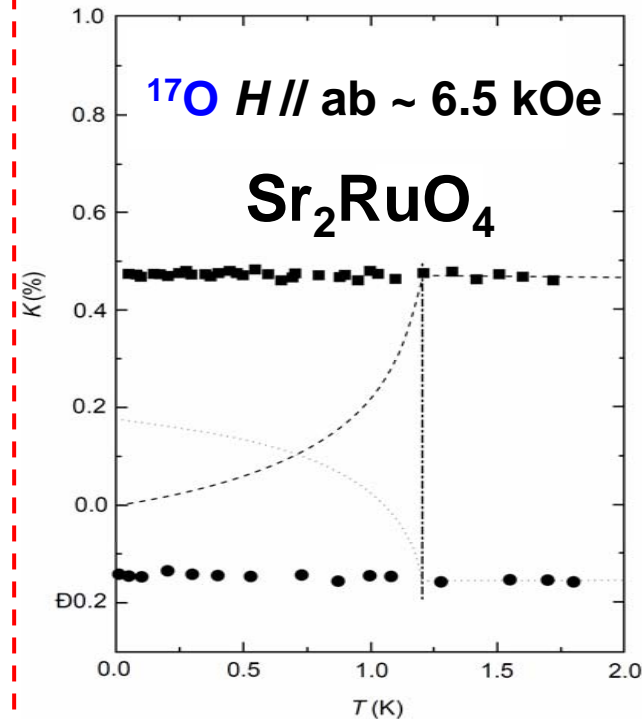
FWHM  $\sim 50\text{e}$



# Ru, $^{17}\text{O}$ Knight-shift measurements



$H > 3$  kOe  $// ab$



K. Ishida, H. Mukuda, Y. Kitaoka *et al.*  
*Phys. Rev. B* 63 (2001) 060507(R).

K. Ishida, H. Mukuda, Y. Kitaoka  
*et al.* *Nature*. 396 (1998) 658.

**Knight-shift is unchanged in the SC state**

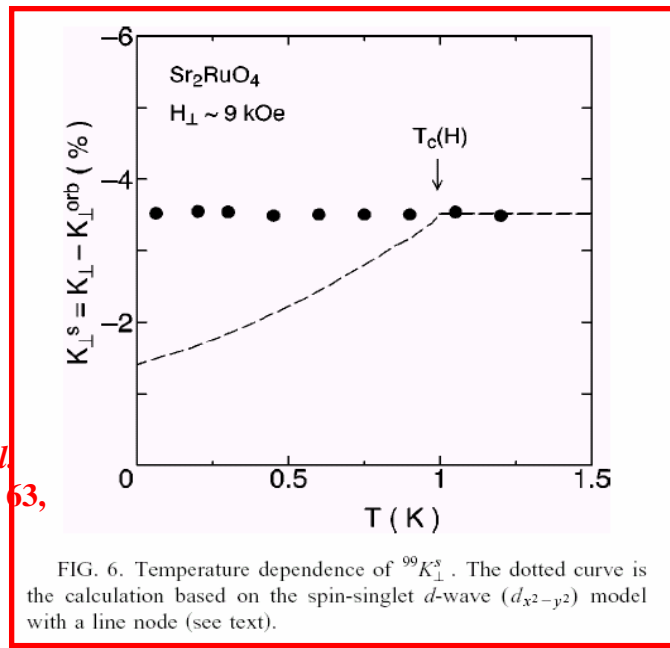


**Spin-triplet superconductivity**



# KS behavior in the SC state

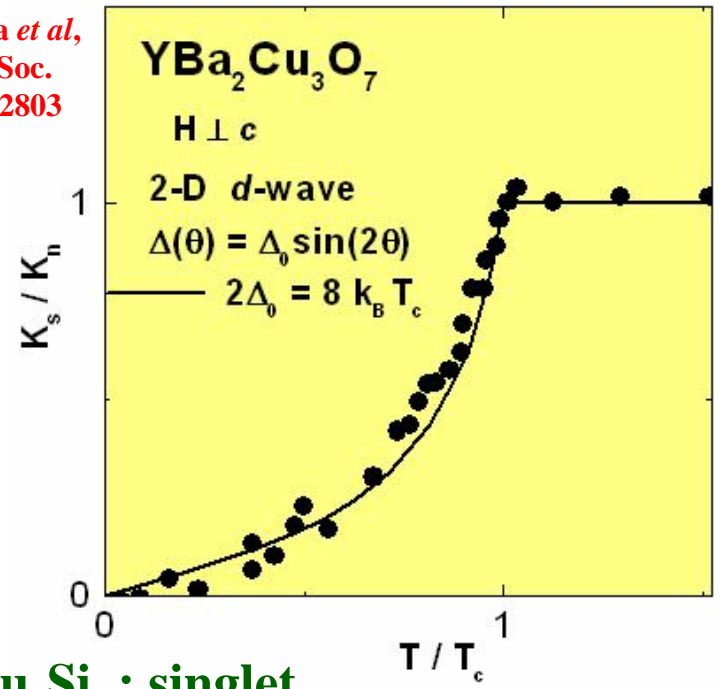
## T-dep. KS in Sr<sub>2</sub>RuO<sub>4</sub>



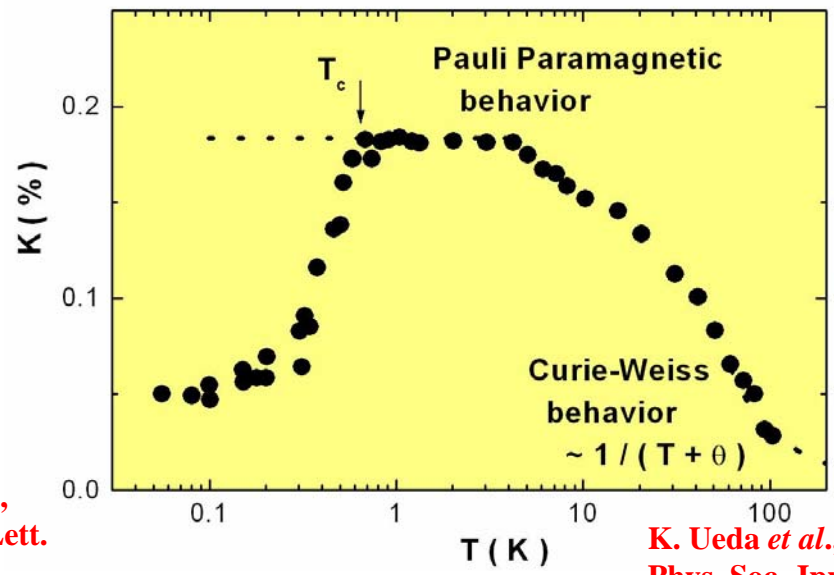
K. Ishida *et al.*  
 Phys. Rev. B **63**,  
 060507(R)

## High- $T_c$ : singlet

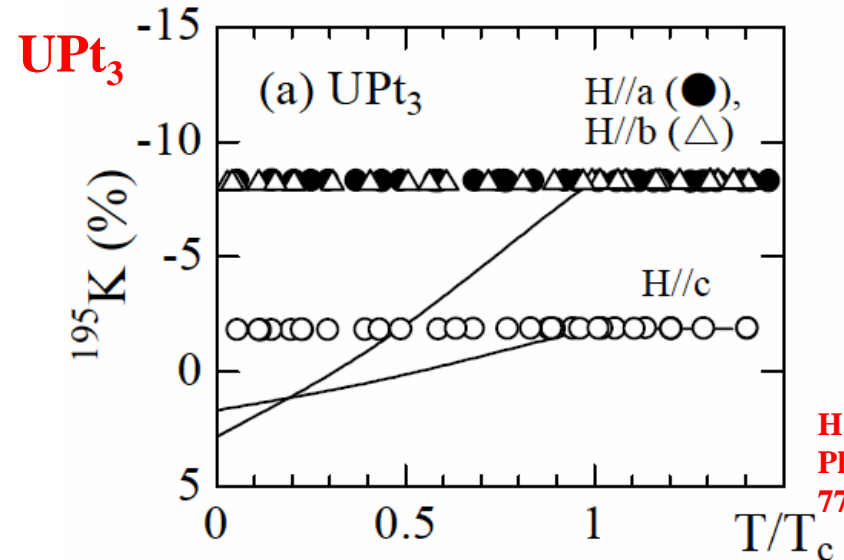
K. Ishida *et al.*,  
 J. Phys. Soc.  
 Jpn. **63**, 2803  
 (93)



## CeCu<sub>2</sub>Si<sub>2</sub> : singlet



K. Ueda *et al.*, J.  
 Phys. Soc. Jpn  
**56**, 867 (87)

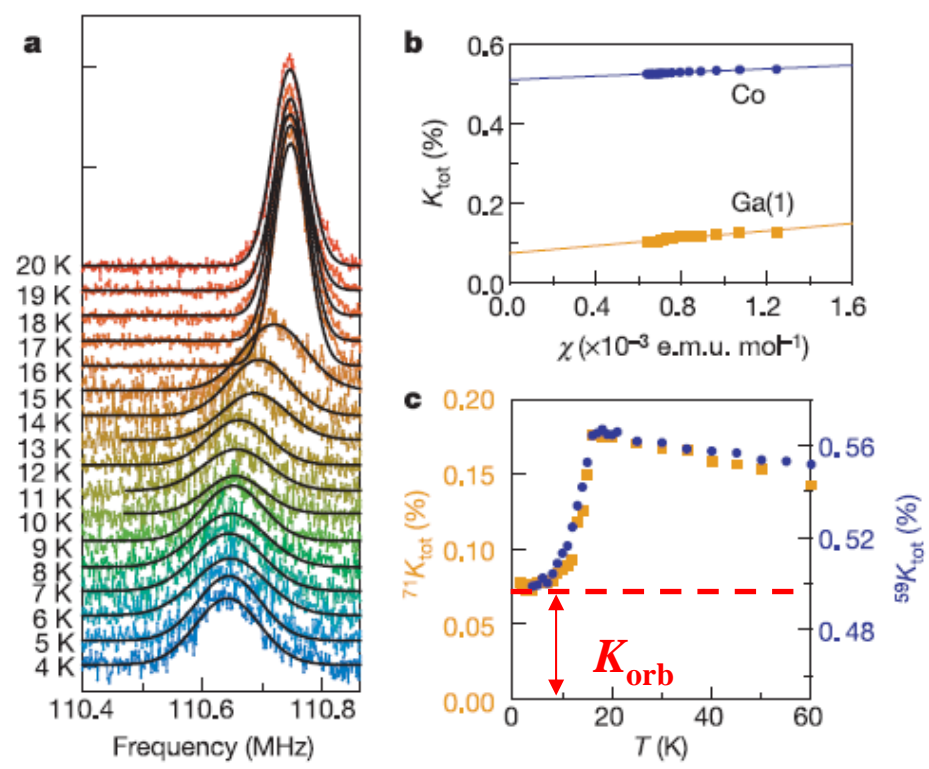


H. Tou *et al.*,  
 Phys. Rev. Lett.  
**77** 1374 (96)

# Other example (Spin-singlet Superconductor)

**PuCoIn<sub>5</sub> ( $T_c \sim 20$  K)**

**N. Curro *et al.* Nature 434, 622 (05)**

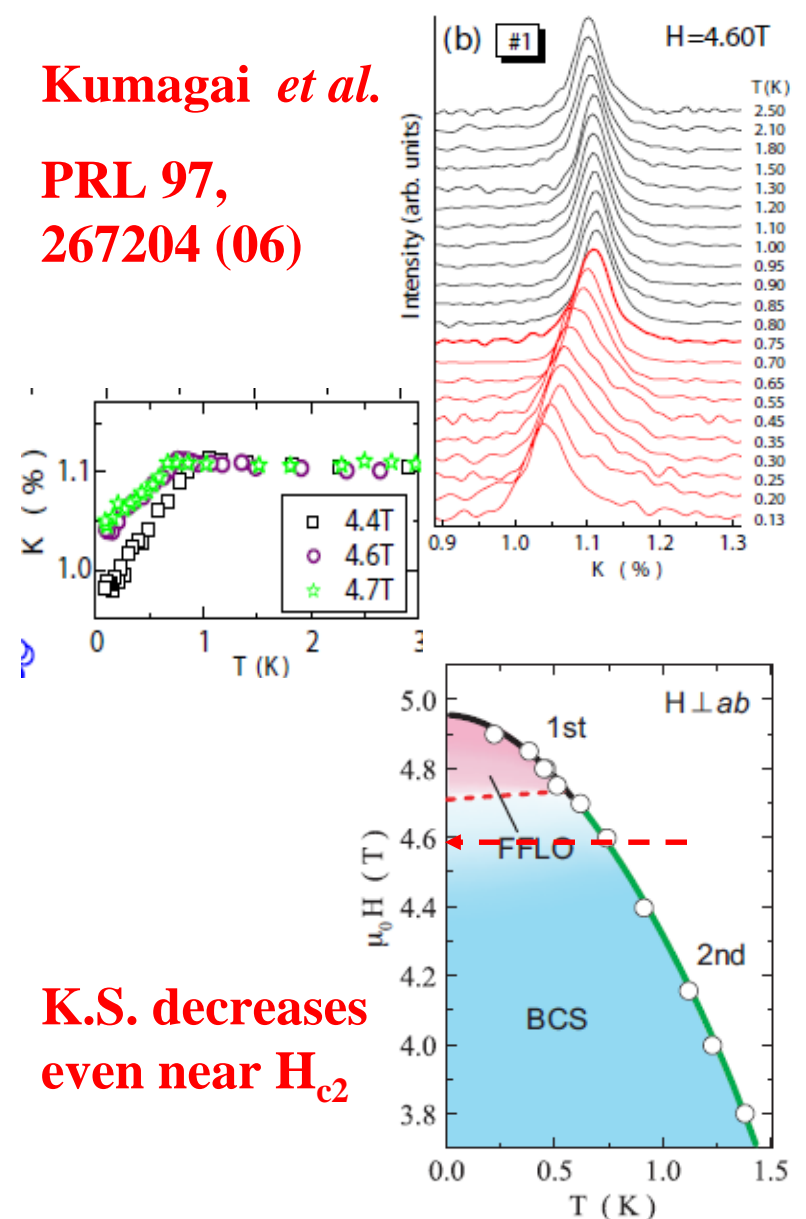


**Figure 1** Knight shift measurements in PuCoGa<sub>5</sub>. **a**, NMR spectra of  $^{71}\text{Ga}$  in 8 T at a series of temperatures through  $T_c$ . The spectra have been offset vertically for clarity. Solid lines are gaussian fits. **b**, The normal-state magnetic shift  $K_{\text{tot}}$  of the  $^{59}\text{Co}$  and  $^{71}\text{Ga}(1)$  versus bulk susceptibility  $\chi$ . The intercepts and hyperfine constants are given by  $^{59}K_{\text{orb}} = 0.53\%$ ,  $^{71}K_{\text{orb}} = 0.088\%$ , and  $^{59}A = 1.5\text{kOe } \mu_B^{-1}$ ,  $^{71}A = 4.1\text{kOe } \mu_B^{-1}$ . Solid lines are fits to the low-temperature data. **c**, The total magnetic shift  $K_{\text{tot}}$  of the  $^{59}\text{Co}$  and  $^{71}\text{Ga}(1)$  versus temperature.

**CeCoIn<sub>5</sub> ( $T_c \sim 2.3$  K)**

**Kumagai *et al.***

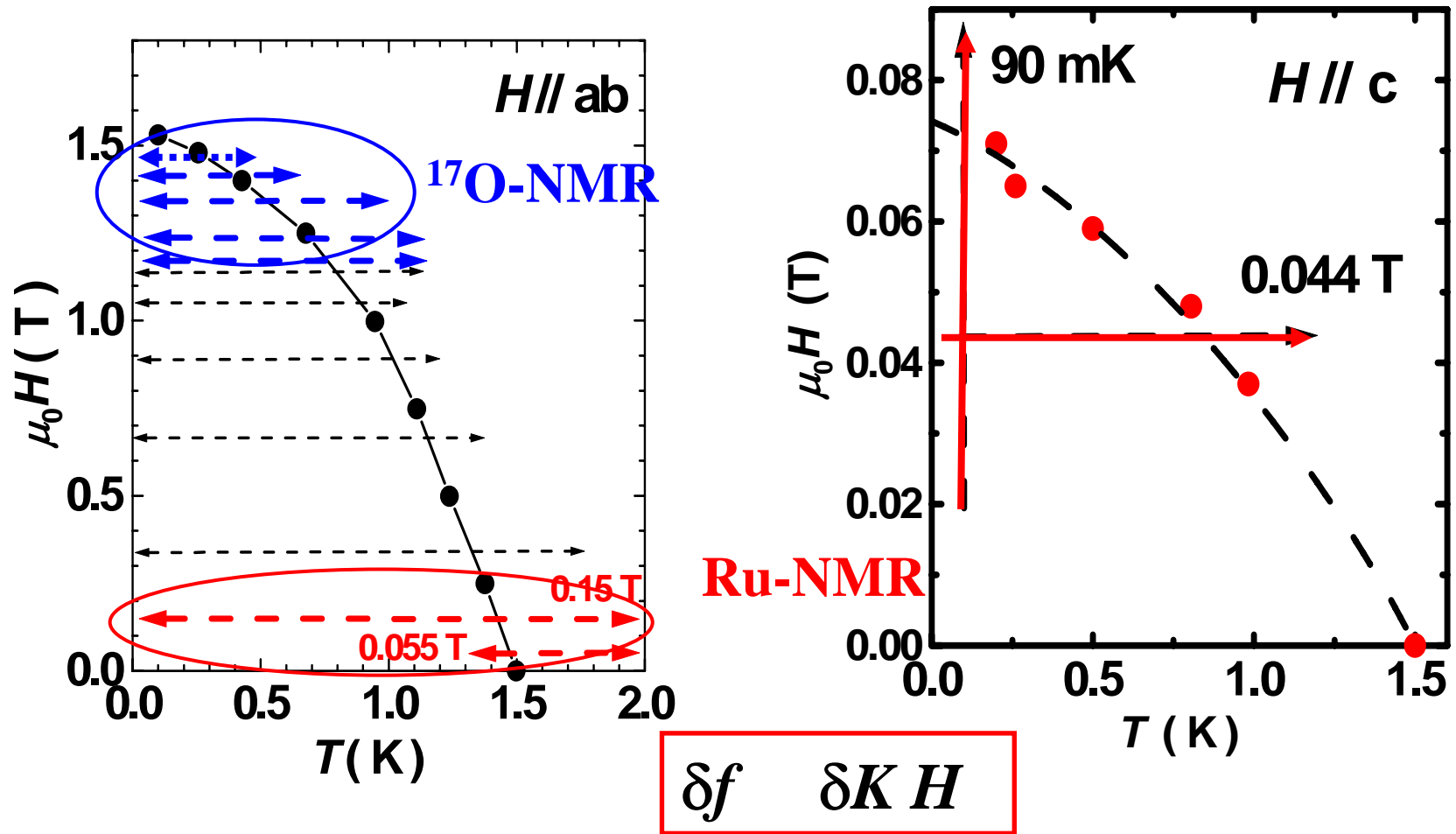
**PRL 97, 267204 (06)**



**K.S. decreases even near  $H_{c2}$**

# Summary of the Knight-Shift Measurements so far

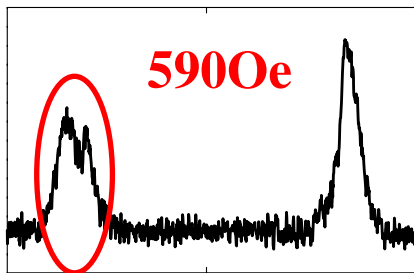
Applied fields were controlled with an accuracy less than 0.5 degree.



Knight shift is unchanged in the field of  $\mu_0 H_{ab} > 0.055\text{ T}$ ,  $\mu_0 H_c > 0.02\text{ T}$ .

# $T$ dependence of $^{101}\text{K}$ at $H = 440$ Oe

$H // c$

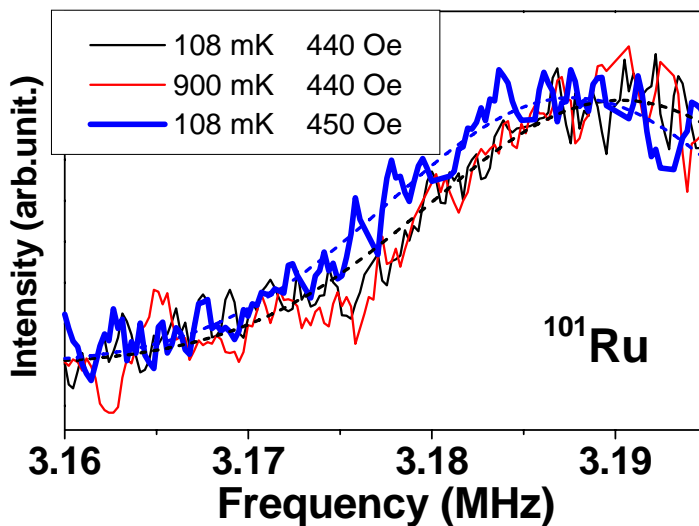
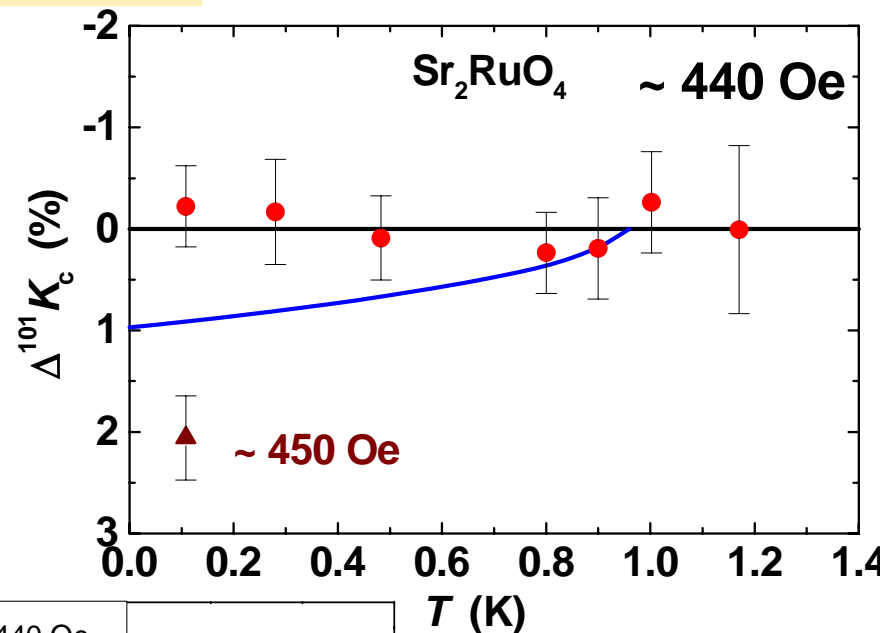
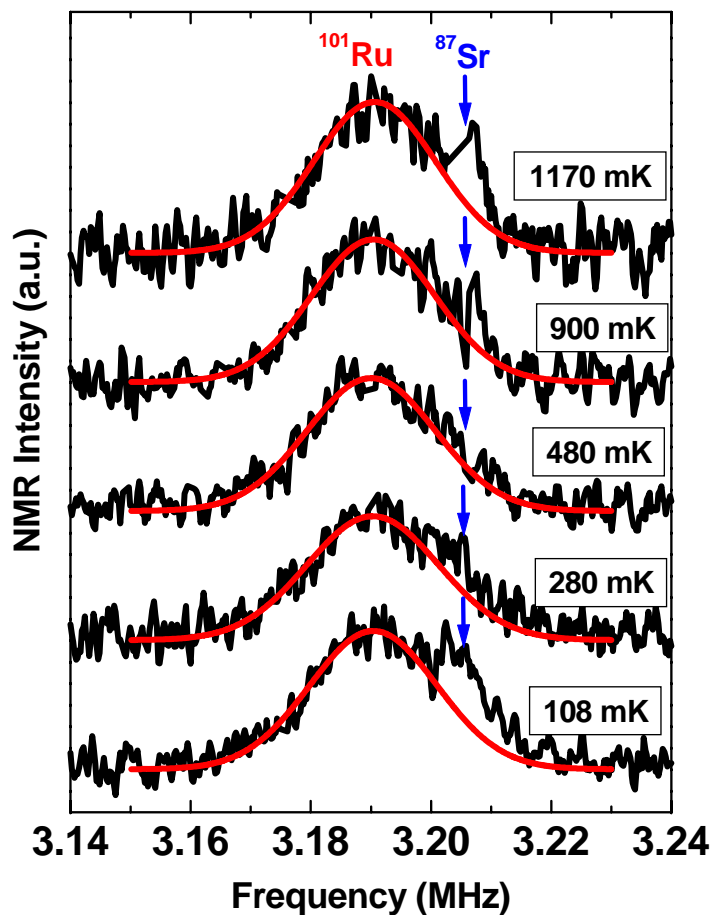


## Ru-NMR

$H // c$

$\kappa \sim 2.3$

$^{101}\text{Ru}$  NMR at 440 Oe at 3.19 MHz



$\delta K \sim 1\%$

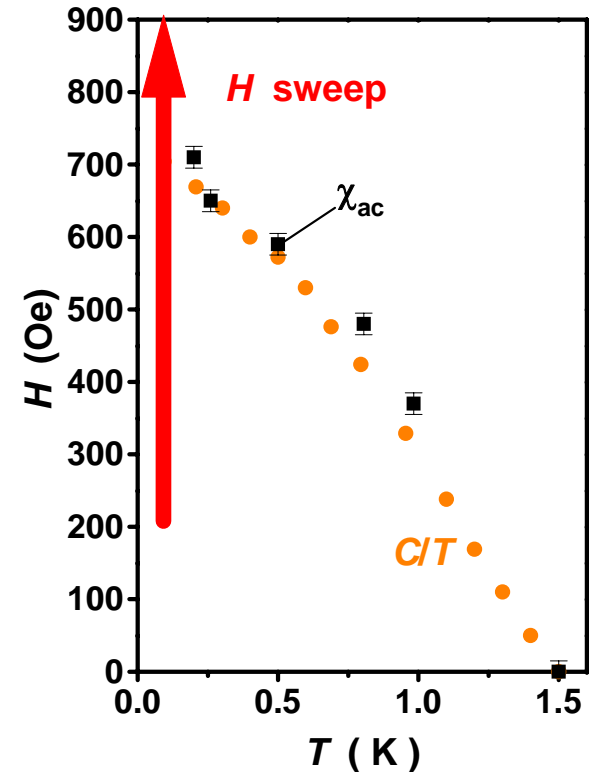
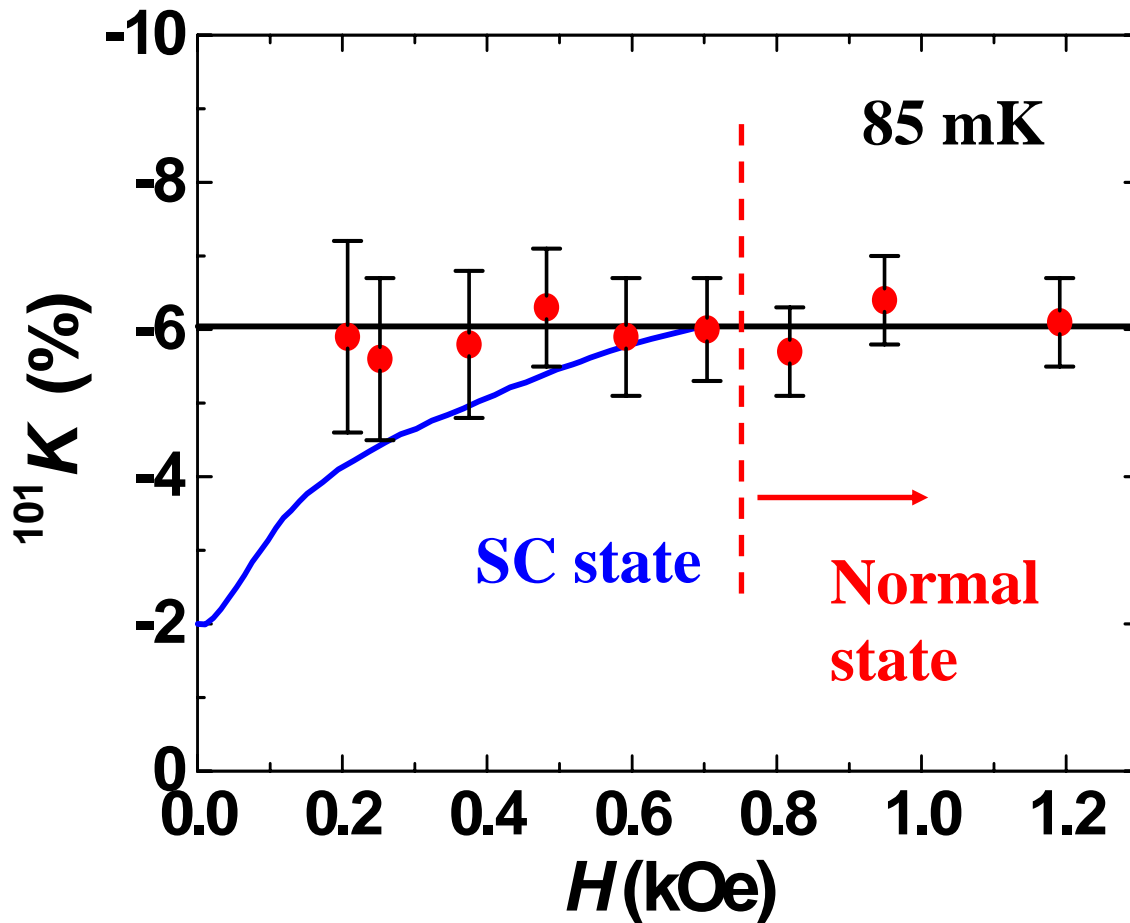
$\delta H = 4.4$  Oe

$\delta f = \gamma \times \delta H$

$\sim 1$  kHz

Within the experimental accuracy,  
 $^{101}\text{K}$  is invariant with  $T$ .

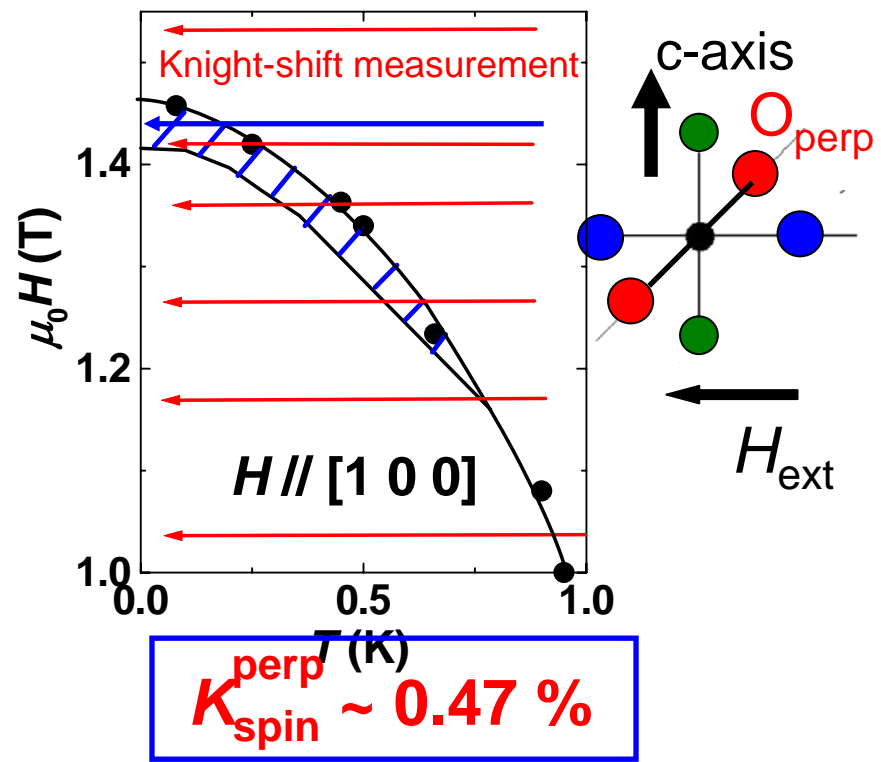
# $H$ dependence of the Ru Knight Shift at Low $T$



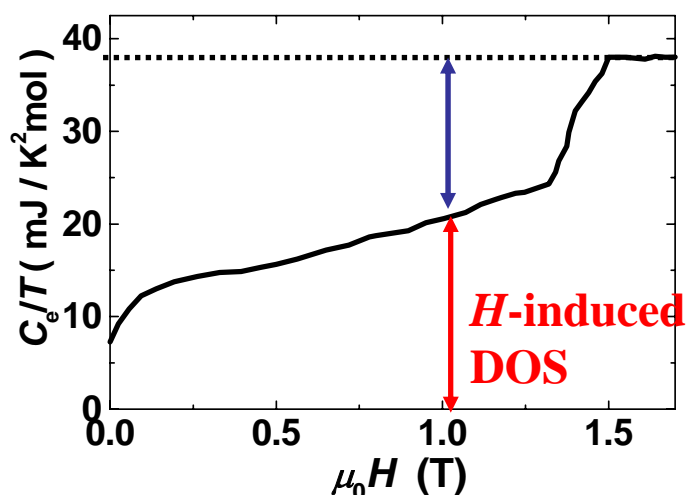
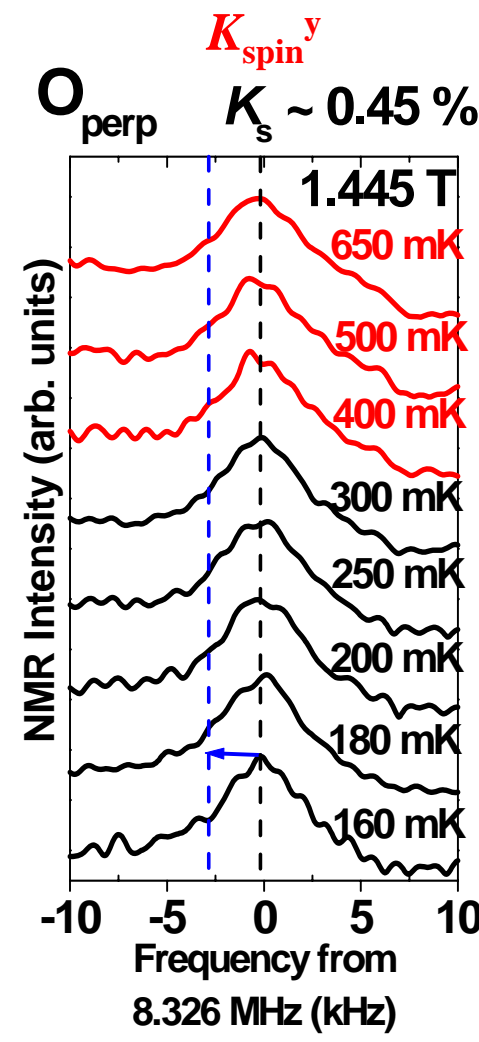
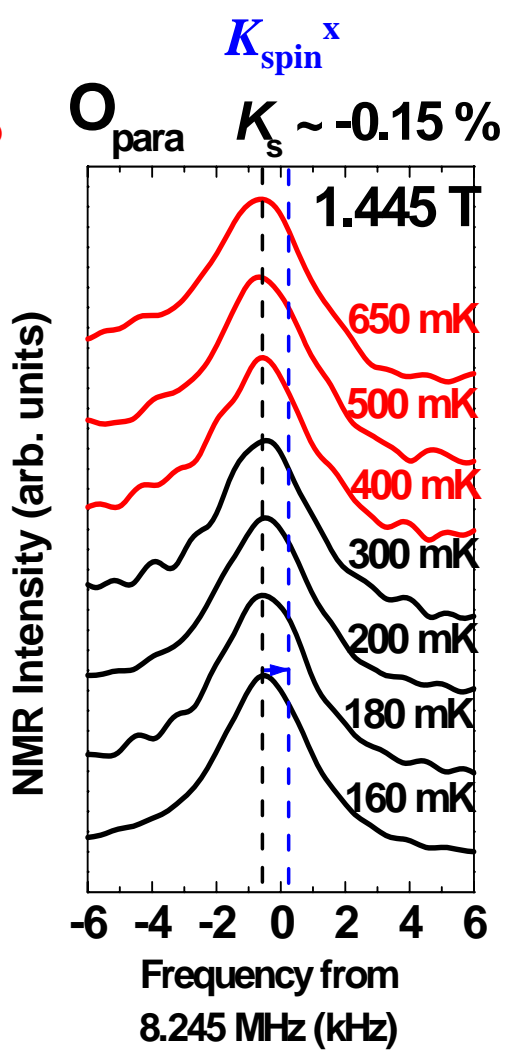
Within the experimental accuracy,

$^{101}\text{K}_S$  ( $\sim -2\%$ ) is unchanged across  $T_c$ .

# $^{17}\text{O}$ NMR at "D" phase



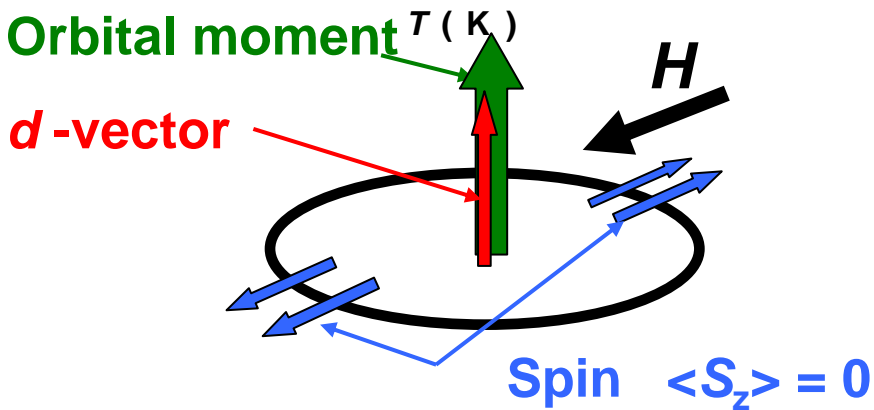
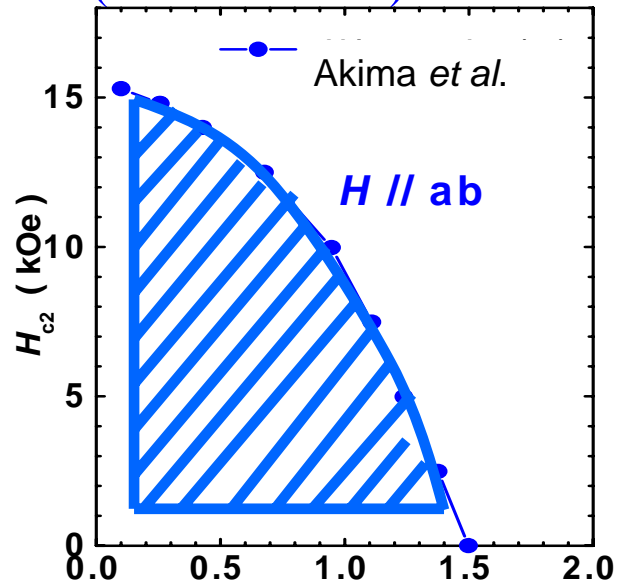
# $^{17}\text{O}$ NMR



# Summary

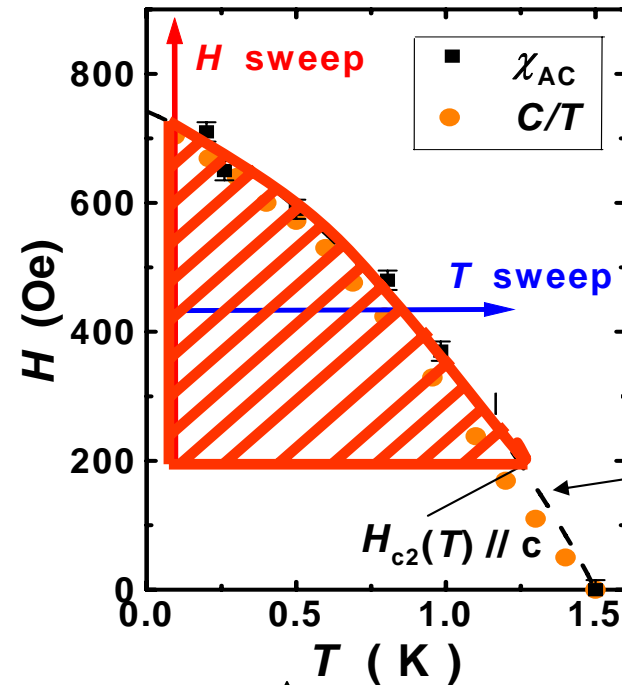
KS was measured in the wider-field range shown below, but the decrease of KS was not observed in the range.

$H // ab$  ( $H > 550$  Oe)

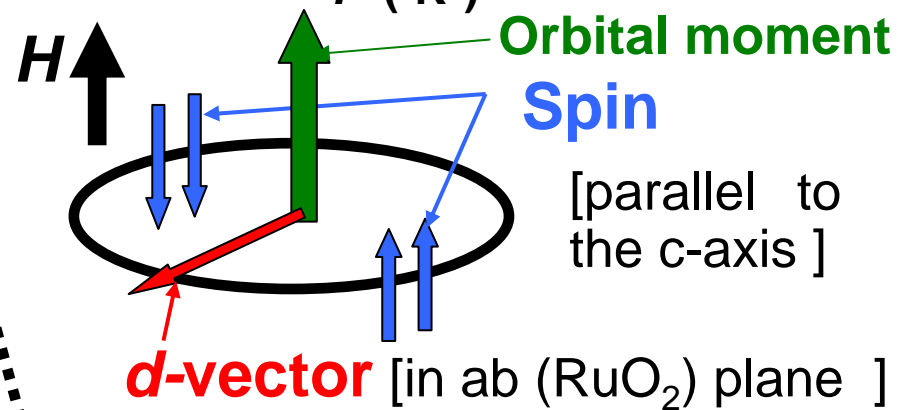


$$\mathbf{d} = \Delta_0 (\sin k_x + i \sin k_y) \text{ (A state)}$$

$H // c$  ( $H > 200$  Oe)



Due to Meissner effect, precise measurement of KS was impossible



# The effect of the spin-orbit interaction at Ru-4d orbitals ( $\lambda$ -s)

Y. Yanase and M. Ogata JPSJ 72, 673 (2003).

The symmetry breaking interaction in the  $d$ -vector space is **in the second order with respect to  $\lambda$** .

**Ikeda &  
Nomura**

$$F_{d\perp c} - F_{d\parallel c} \sim UJ_H \lambda^2 / W^3$$

Anisotropic energy is reduced to  $\sim 10^{-4}$ .

$$\Delta T_c \sim 0.01 T_c \quad H_{\text{flip}} \sim 1 \text{ kOe}$$

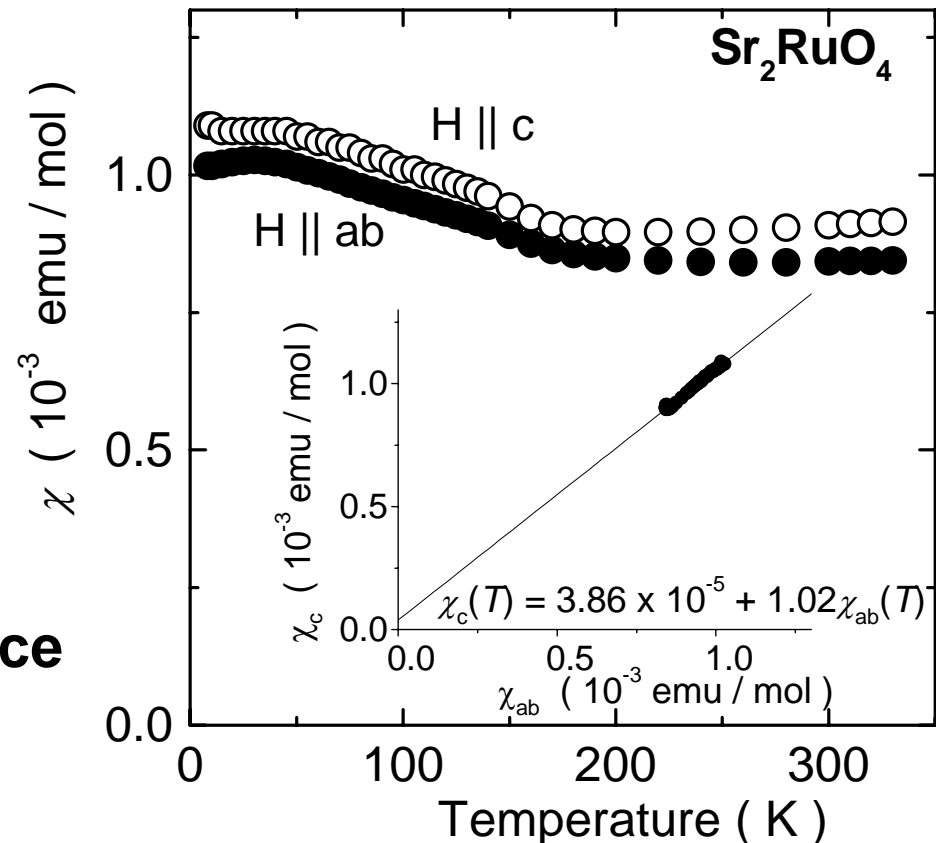
This is the order estimation.

## DC susceptibility

The normal state DC susceptibility is almost isotropic.

$$\chi_c / \chi_{ab} \sim 1.02$$

The anisotropy in the spin space seems to be small.





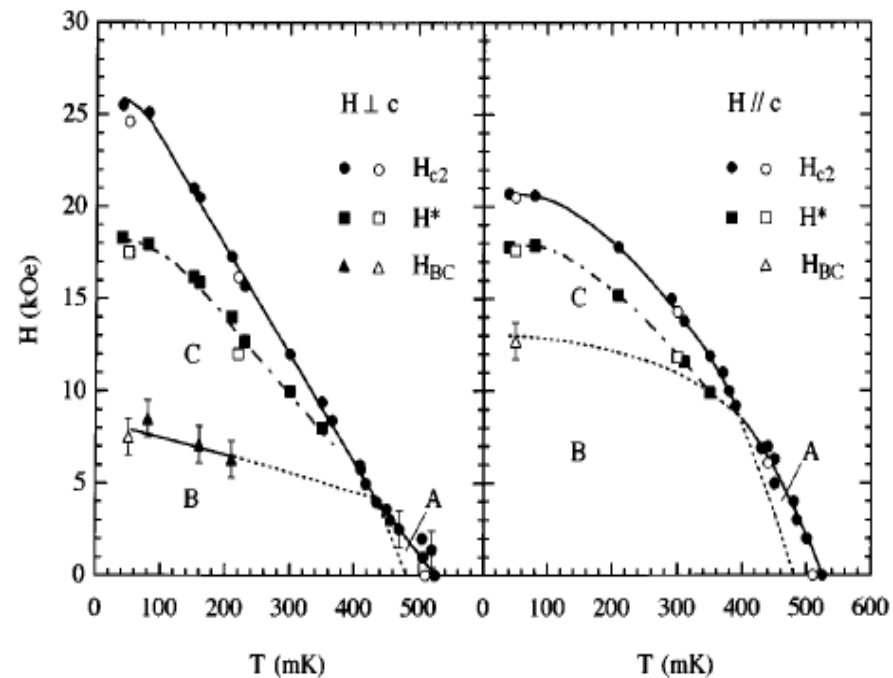
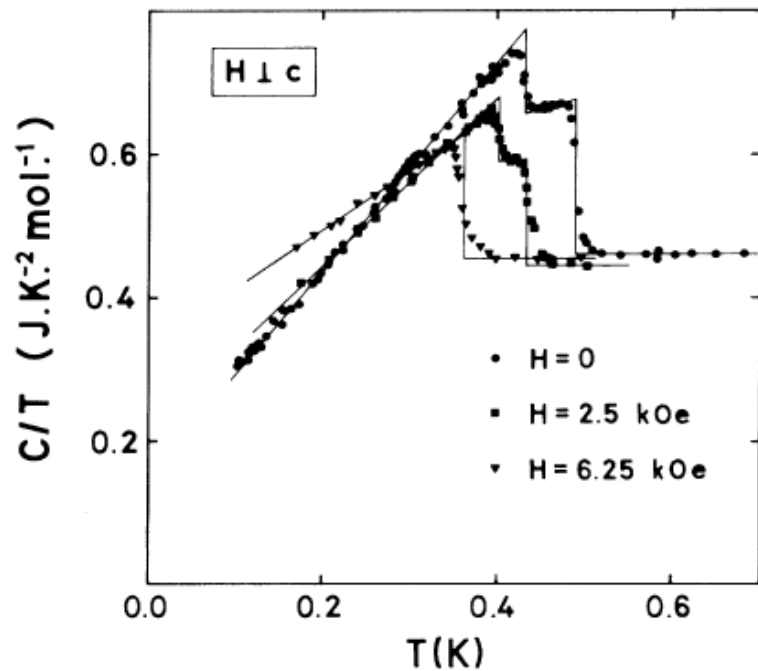
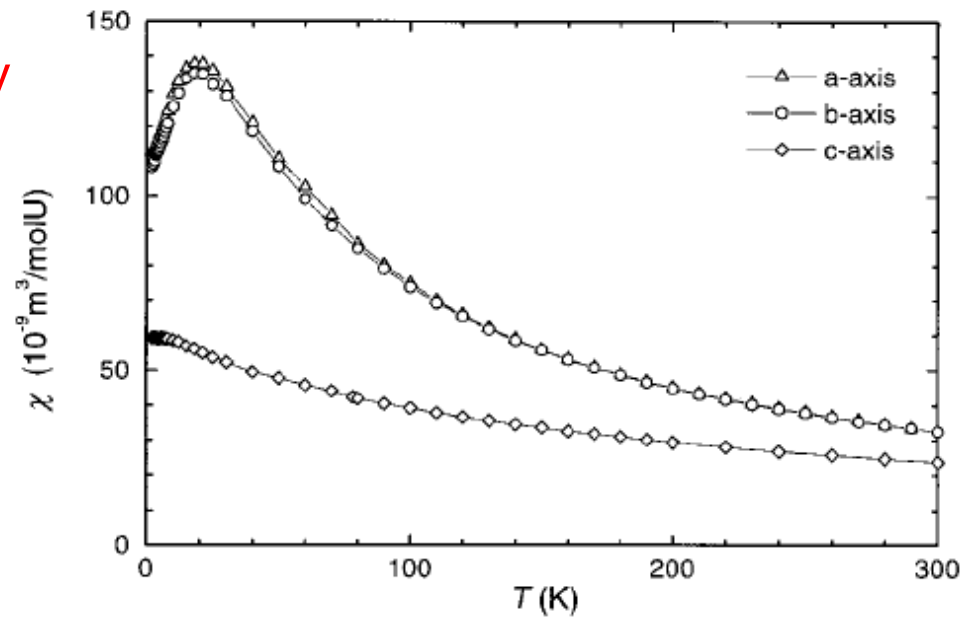
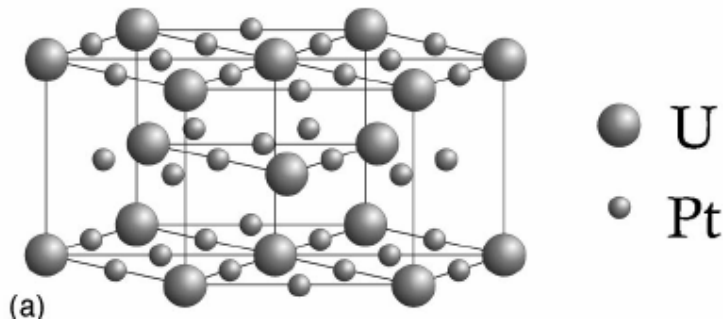
# UPt<sub>3</sub> : Spin-triplet Multi-phase superconductor

$T_{c1} \sim 0.58$  K,

$T_{c2} \sim 0.53$  K

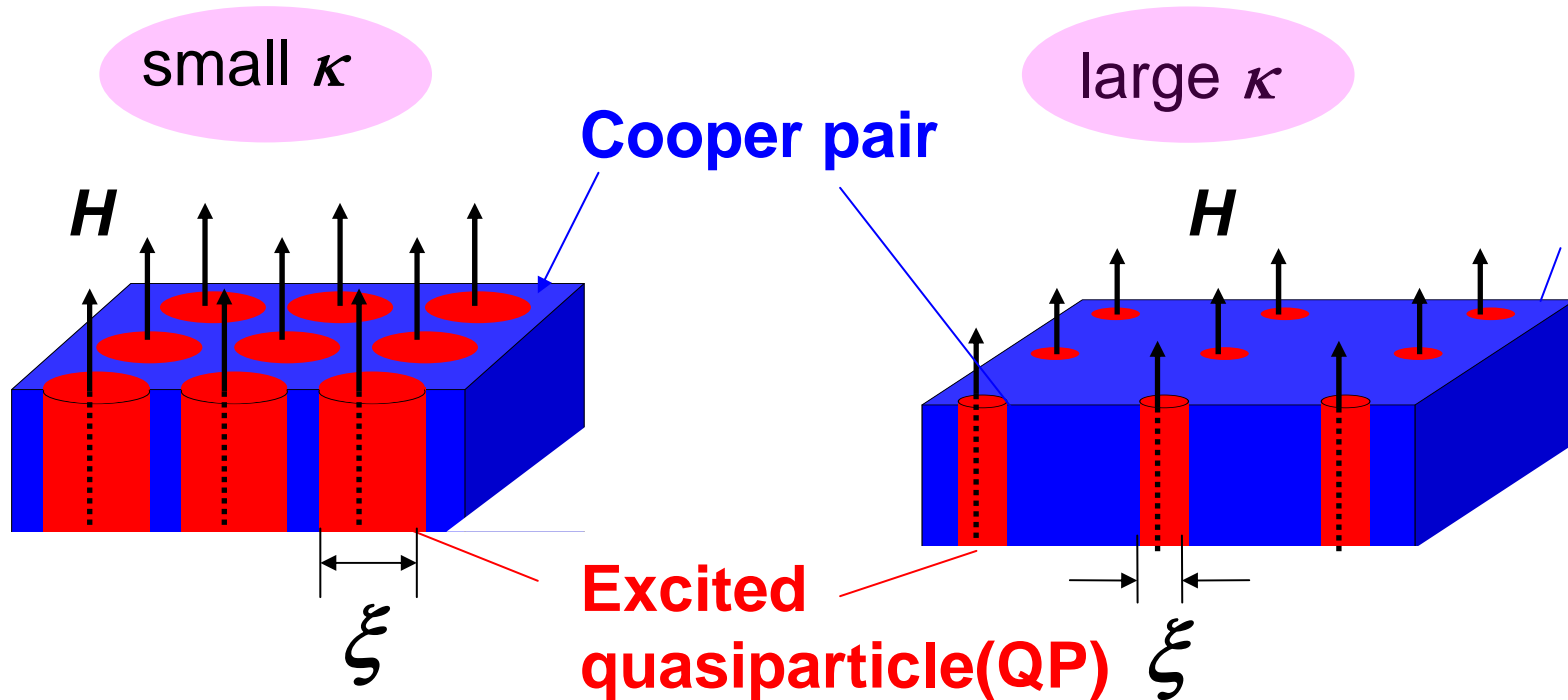
Magnetic anomaly

$T_M \sim 5$  K



# Ginzburg Landau Parameter $\kappa$

$\kappa \sim 2.7$  in  $H // c$  ( $\text{Sr}_2\text{RuO}_4$ )



Estimation the ratio of the spin susceptibility from the excited QP is needed.



Specific heat data (Deguchi *et al.*) is used for estimating the QP's contribution.

# Superconducting Parameter in Sr<sub>2</sub>RuO<sub>4</sub>

Parameter	<i>ab</i>	<i>c</i>
$\mu_0 H_{c2\parallel c}(0)$ (T)	0.075	
$\mu_0 H_{c2\parallel ab}(0)$ (T)	1.50	
$\mu_0 H_c(0)$ (T)	0.023	
$\xi(0)$ (Å)	660	33
$\lambda(0)$ (Å)	1520	$3.0 \times 10^4$
$\kappa(0)$	2.3	46
$\gamma_s = \xi_{ab}(0) / \xi_c(0)$	20	