

# Spin-singlet nature of various dimerized cuprates. From *ab-initio* to experiment.

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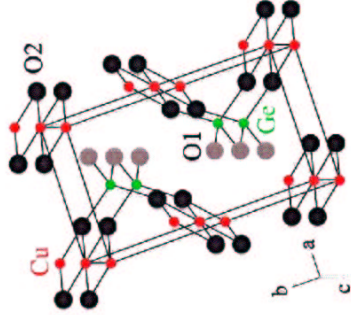
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## Outline

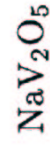
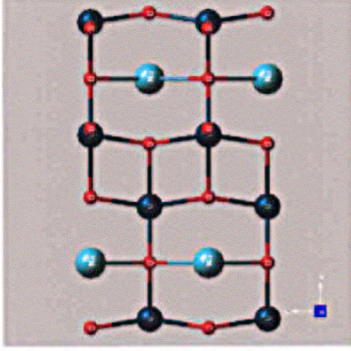
- Motivation - Systems
- Dimerized-cuprates, why are they interesting?
  - CaCuGe<sub>2</sub>O<sub>6</sub> -related to CuGeO<sub>3</sub>-
  - KCuCl<sub>3</sub> and TiCuCl<sub>3</sub>
- Microscopic description: *ab-initio* + QMC
- Summary and Conclusions

Low-dimensional transition-metal compounds

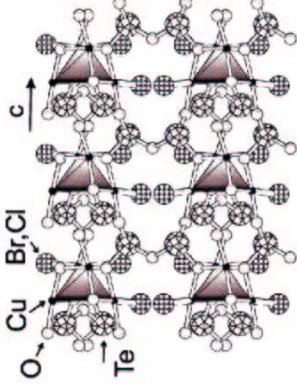
• Cuprates:



• Vanadates:



• Tellurates:



**Properties:** Spin-Peierls phase transition, Frustration, Charge-ordering,  
Quantum-critical behavior, superconductivity, ...

– **Electron correlation important** –

Methods

• **Density Functional Theory**.- Electronic Structure

↓ + Electron correlation

*Effective Model*

Hubbard, Heisenberg, t-J, ...

↓

• **Many-Body techniques**.- Bond-Operator Theory, Variational Ansätze,  
DMRG, QMC, ...

↓

*Physical Picture*

## Dimerized Cuprates, why are they interesting?

Quantum spin systems with a singlet-triplet gap in the spin excitation spectrum

- close gap by a strong enough **H**. Some dimers are promoted to triplets. These triplets behave as a gas of quantum particles. They can undergo different kinds of **phase transitions** as a function of **H** or **T** depending on topology, dimension, ...

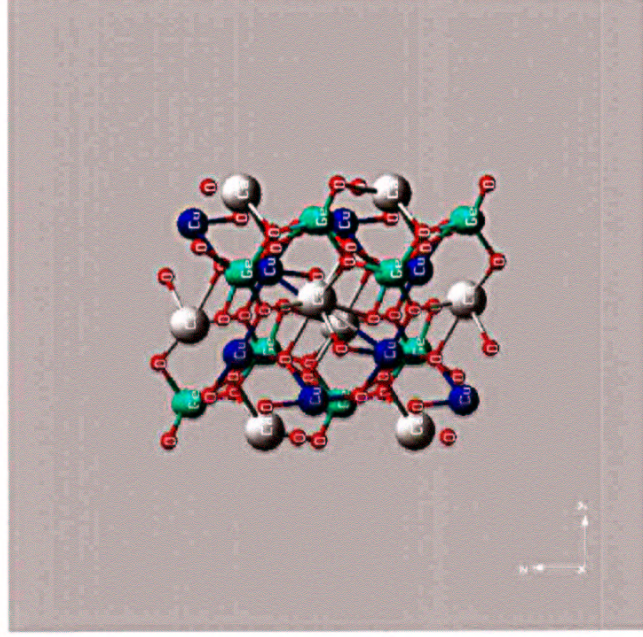
→ **rich physics!**

... often true nature of the exchange networks, not expected from the crystal structure...

→ **microscopic investigation *ab-initio* + QMC**

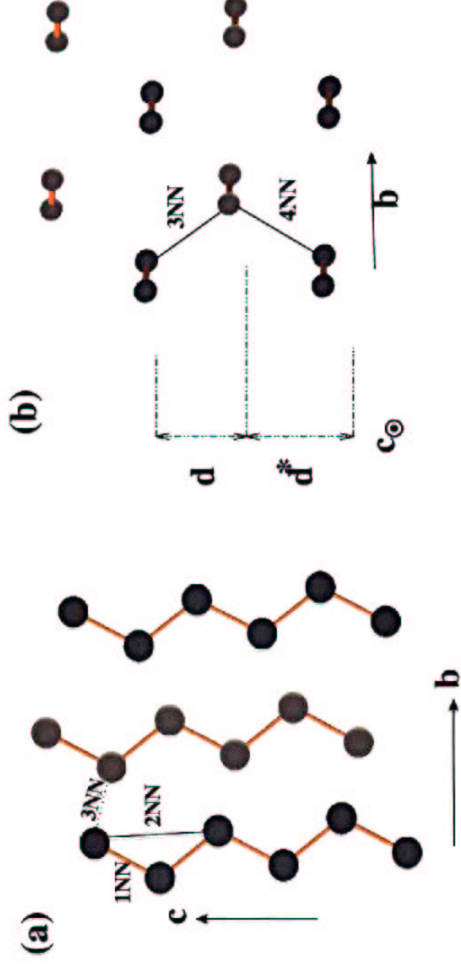
Examples:  $\text{CaCuGe}_2\text{O}_6$ ,  $\text{TlCuCl}_3$  and  $\text{KCuCl}_3$

## $\text{CaCuGe}_2\text{O}_6$ - Structure



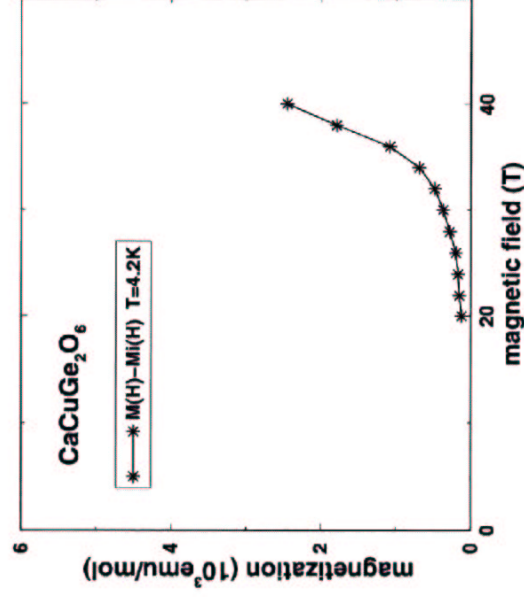
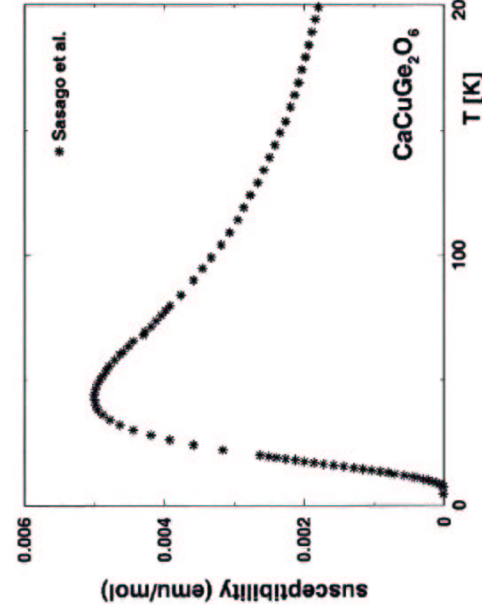
- space group  $P2_1/c$
- $\text{CuO}_6$ -zigzag chains along  $c$

## CaCuGe<sub>2</sub>O<sub>6</sub> - Structure



- $\text{Cu}^{2+}$  -chains alternate between two neighboring  $bc$  planes (black and grey)
- $\text{dist}(1\text{NN})=3.07\text{\AA}$ ,  $d=4.46\text{\AA}$ ,  $d^*=5.35\text{\AA}$

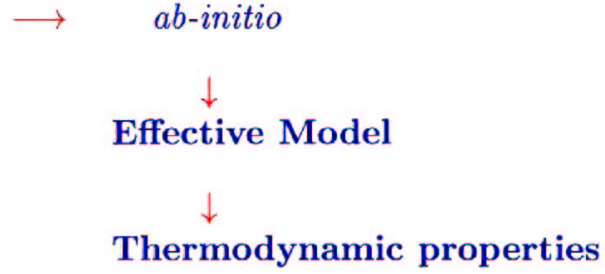
## CaCuGe<sub>2</sub>O<sub>6</sub>- Magnetism



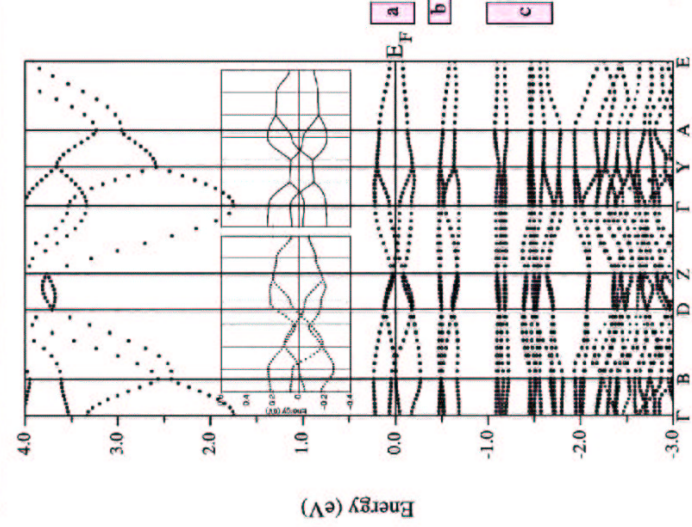
- spin-singlet ground state with an energy gap
- INS (Zheludev *et al.*) energy gap  $\sim 6$  meV
- no spin-Peierls phase transition

## Microscopic model?

Which Cu-pairs can constitute AF dimers?



## CaCuGe<sub>2</sub>O<sub>6</sub>- Band Structure



LAPW and LMTO scheme (GGA)

- a= Cu  $3d_{x^2-y^2}$
- b= Cu  $3d_{3z^2-1}$
- c=Cu  $3d_{xy}, 3d_{xz}, 3d_{yz}$

• *downfolded-TB Model:*

$$H_{TB} = \sum_{(i,j)} t_{ij} (c_i^\dagger c_j + H.c.)$$

$$t_{1NN} = 68\text{meV}, \quad t_{2NN} = 8\text{meV}$$

$$t_{3NN} = 88\text{meV}, \quad t_{4NN} = 4\text{meV}$$

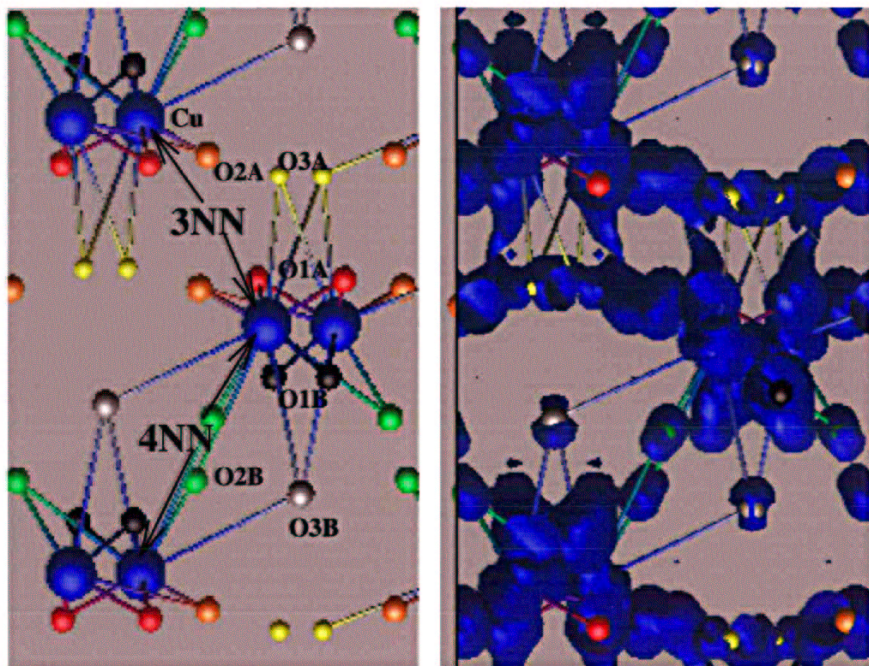
-system of interacting dimers -

$$\text{dimers: } J_3 \sim 4 \frac{t_{3NN}^2}{U}, \quad U \sim 4.2\text{eV},$$

$$J_3 \sim 85\text{K} \sim 7\text{meV}$$

- longer-ranged magnetic interactions dominate over *short-ranged* interactions -

### CaCuGe<sub>2</sub>O<sub>6</sub>- Electron Density

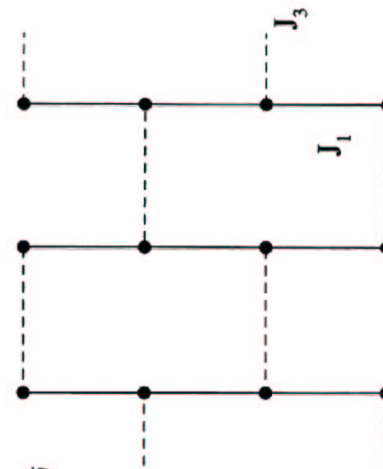


### CaCuGe<sub>2</sub>O<sub>6</sub>- Spin-Model

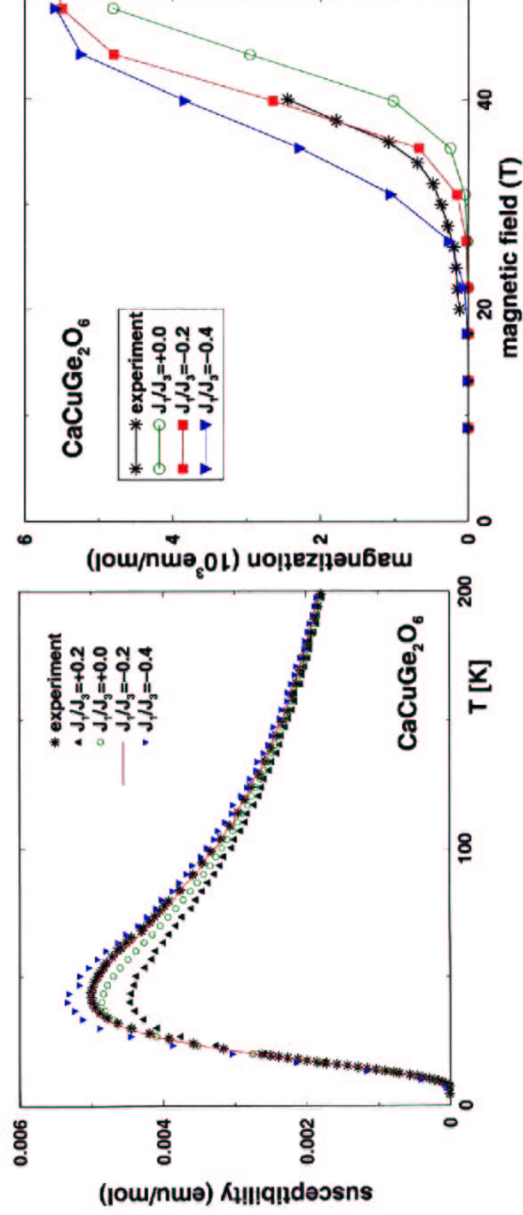
System of interacting dimers

$$H_{eff} = J_3 \sum_{\langle i,j \rangle} \mathbf{S}_i \mathbf{S}_j + J_1 \sum_{(i,j)} \mathbf{S}_i \mathbf{S}_j$$

→ calculation of thermodynamic properties  
by QMC (stochastic series expansion)



## CaCuGe<sub>2</sub>O<sub>6</sub>- Magnetism



- Optimal value  $J_1/J_3 = -0.2$ ,  $J_3 = 67\text{K} = 5.8\text{meV} \Rightarrow$   
- weak ferromagnetic interdimer coupling  $J_1 < 0$  -
- The model has two quantum critical points  $J_1 \simeq 0.55J_3$ ,  $J_1 \simeq -0.9J_3$

## Order of Magnitude - Comparison with CuGeO<sub>3</sub>

### CaCuGe<sub>2</sub>O<sub>6</sub>:

- $ab-initio$  + Examination of  $\chi$  and  $M$  by QMC  $\Rightarrow$
- system of dimers formed by  $3\text{NN}$   $s = \frac{1}{2}$   $\text{Cu}^{2+}$   
with FM 1NN interdimer couplings  $J_1/J_3 \sim -0.2$ .

$J_3(ab-initio) \sim 7\text{meV}$ ,  $J_3(\text{QMC}) = 5.8\text{meV}$

FM  $J_1$   $\alpha_{Cu-O-Cu} = 92^\circ$ ,  $\beta_{Cu-O-Cu} = 98^\circ$

- no spin-Peierls phase transition

### CuGeO<sub>3</sub>:

- Frustrated  $J_1$ - $J_2$  system + spin-Peierls phase transition

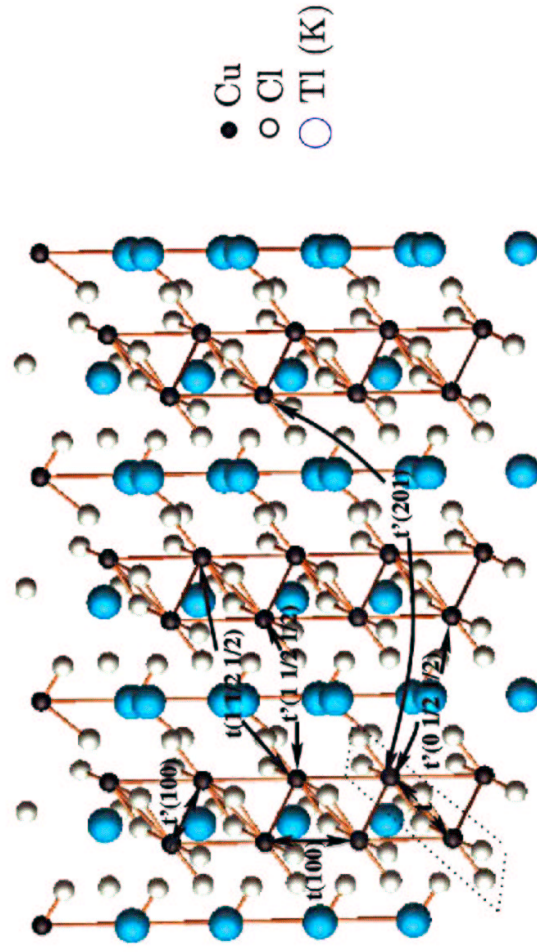
$J_1 \sim 10.3\text{meV}$ ,  $J_2/J_1 \sim 0.24$ - $0.34$

### Spin-gap systems $\text{KCuCl}_3$ and $\text{TlCuCl}_3$

- Isostructural  $P2_1/c$
  - **INS** (Cavadini *et al.* ('00) Tanaka *et al.* ('01) Oosawa *et al.* ('02) ):
    - spin excitation gap  $\Delta \sim 0.65\text{meV}$
    - *strongly*-coupled spin-dimer system
    - Néel ordering for  $H > H_g = \Delta/g\mu_B \sim 6\text{T}$  **BEC of magnons**
  - **$\text{TlCuCl}_3$ :**
    - spin excitation gap  $\Delta \sim 2.6\text{meV}$
    - *weakly*-coupled spin-dimer system
    - Néel ordering for  $H > H_g = \Delta/g\mu_B \sim 22\text{T}$
  - **$\text{KCuCl}_3$ :**
    - spin excitation gap  $\Delta \sim 2.6\text{meV}$
    - *weakly*-coupled spin-dimer system
    - Néel ordering for  $H > H_g = \Delta/g\mu_B \sim 22\text{T}$
- *ab-initio*
- **What determines their different behavior?**
    - structural differences?
    - role  $\text{Tl} \leftrightarrow \text{K}$  ?

### $\text{KCuCl}_3$ and $\text{TlCuCl}_3$ - Structure

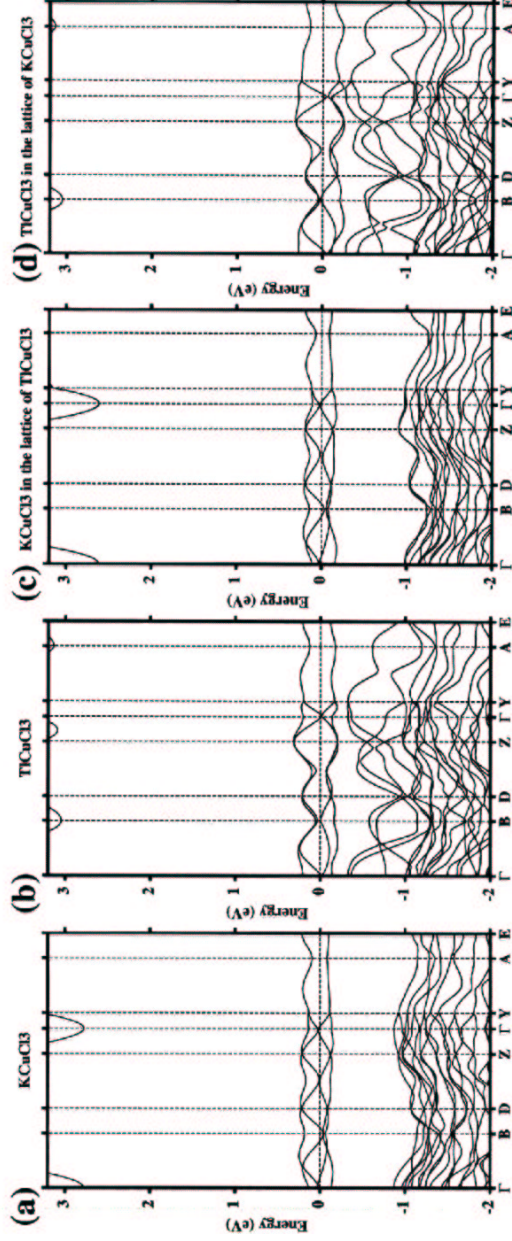
Coupled two-leg  $s=1/2$   $\text{Cu}^{2+}$  ladders separated by  $\text{Tl}^+$  ( $\text{K}^+$ ) ions





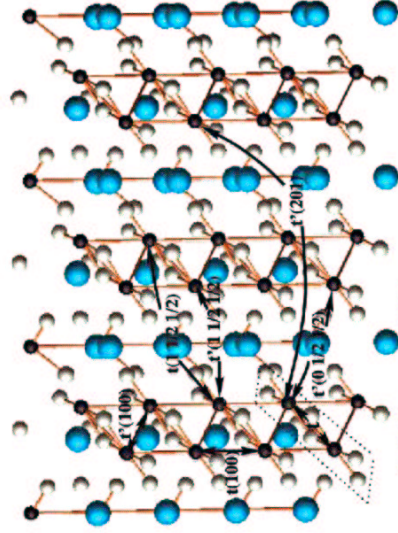
### KCuCl<sub>3</sub> and TiCuCl<sub>3</sub> - Band Structure

- Four narrow bands close to EF formed by  $\text{Cu } d_{x^2-y^2}$  orbitals admixed with Cl p.



- Change of lattice parameters has a **minor** effect on the band structure ((a) ↔ (c))
- Substitution of  $\text{K}^+$  by  $\text{Ti}^+$  in the same lattice has a **significant** effect ((a) ↔ (d))

### KCuCl<sub>3</sub> and TiCuCl<sub>3</sub> - Effective Model



Downfolding + TB:

**TiCuCl<sub>3</sub>**:  $t=134$  meV,  $t'_{2,0,1}=65$  meV,  $t_{1,-1/2,1/2}=32$  meV,  $t_{0,1/2,1/2}=8$  meV

**KCuCl<sub>3</sub>**:  $t=116$  meV,  $t'_{2,0,1}=37$  meV,  $t_{1,-1/2,1/2}=4$  meV,  $t_{0,1/2,1/2}=5$  meV

Comparison with INS analysis (Cavadini et al. ('01), Oosawa et al. ('02)):

**TiCuCl<sub>3</sub>**:  $J=5.7$  meV,  $J_{2,0,1}=-1.5$  meV,  $J_{1,-1/2,1/2}=0.5$  meV,  $J_{0,1/2,1/2}=-0.06$  meV

**KCuCl<sub>3</sub>**:  $J=4.3$  meV,  $J_{2,0,1}=-0.4$  meV,  $J_{1,-1/2,1/2}=0.3$  meV,  $J_{0,1/2,1/2}=-0.003$  meV

## KCuCl<sub>3</sub> and TiCuCl<sub>3</sub> - Comparison

*ab-initio*:

- Structural differences **minor** effect.  
TiCuCl<sub>3</sub> lattice is compressed along *a* and enlarged in the *bc* plane with respect to KCuCl<sub>3</sub>.
- Evident role of TI<sup>+</sup> in Cu<sup>2+</sup> - Cu<sup>2+</sup> path.  
i.e. larger hybridization effect between Cu 3d<sub>x<sup>2</sup>-y<sup>2</sup></sub> and TI 6s and 6p energy states
- TiCuCl<sub>3</sub> stronger interdimer couplings than in KCuCl<sub>3</sub> in agreement with predictions from INS experiments.

## Summary and Conclusions

- Microscopic investigation (*ab-initio* + QMC) of CaCuGe<sub>2</sub>O<sub>6</sub>:
  - system with *longer-ranged* dominated exchange couplings (3NN) with ferromagnetic (1NN) interdimer couplings.
  - other examples: CaV<sub>4</sub>O<sub>9</sub>, (VO)<sub>2</sub>P<sub>2</sub>O<sub>7</sub>
- Microscopic investigation (*ab-initio*) of TiCuCl<sub>3</sub> and KCuCl<sub>3</sub>:
  - corroborates the coupled dimer behavior predicted by INS.
  - differences in the interdimer coupling related to contribution of TI<sup>+</sup> to the chemical bonding between two Cu<sup>2+</sup> ions.