

## Confinement and Bechgaard salts

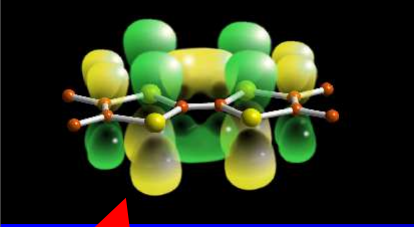
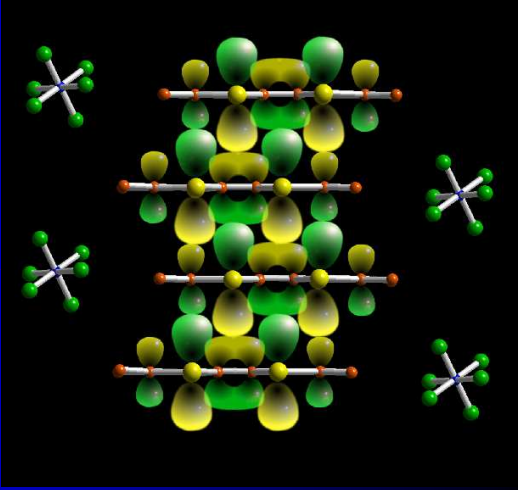
G. Gruner, L. DeGiorgi  
A. Schwarz, V. Vescoli, M. Dressel

A. Georges, N. Sandler

S. Biermann, A. Lichtenstein

- Organic Superconductors: quasi-one dimensional systems (Bechgaard salts)
- TMTTF and TMTSF molecules
- Remarkable properties :
  - Non Fermi liquid behavior
  - Quantum Hall effect
  - Superconductivity and Frohlich conductivity

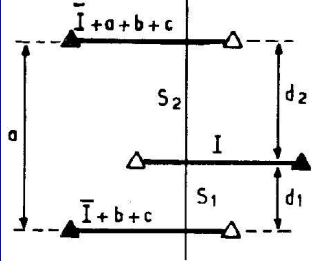
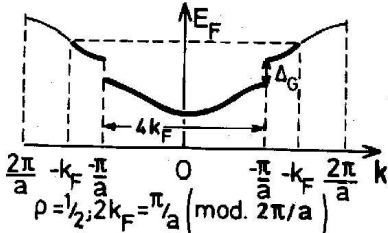
## Organic (super-) conductors

TMTSF<sub>2</sub>(X)

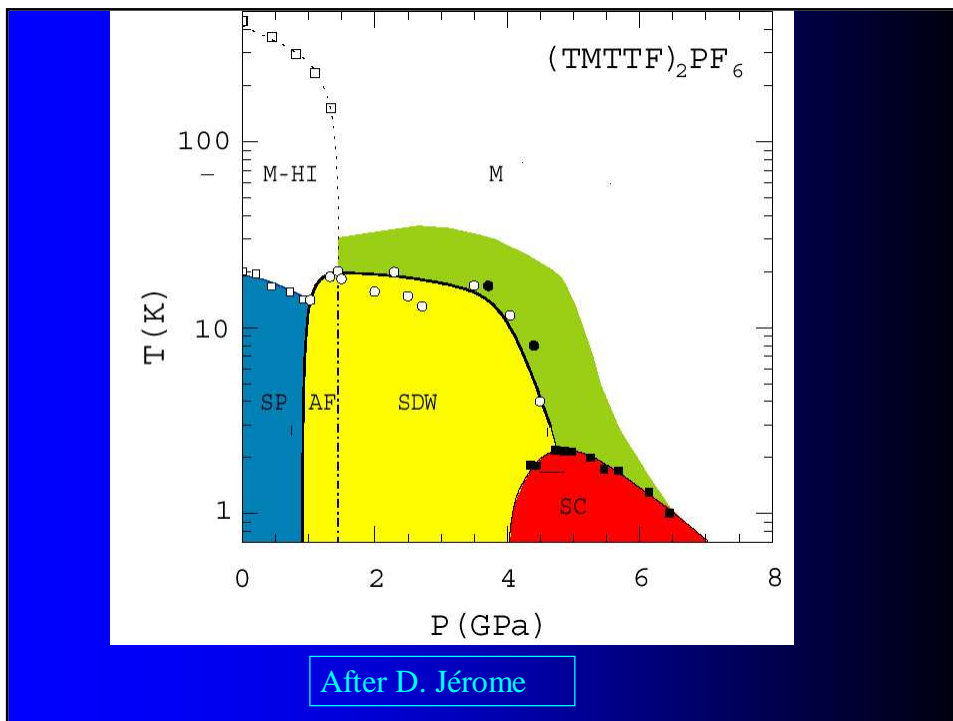
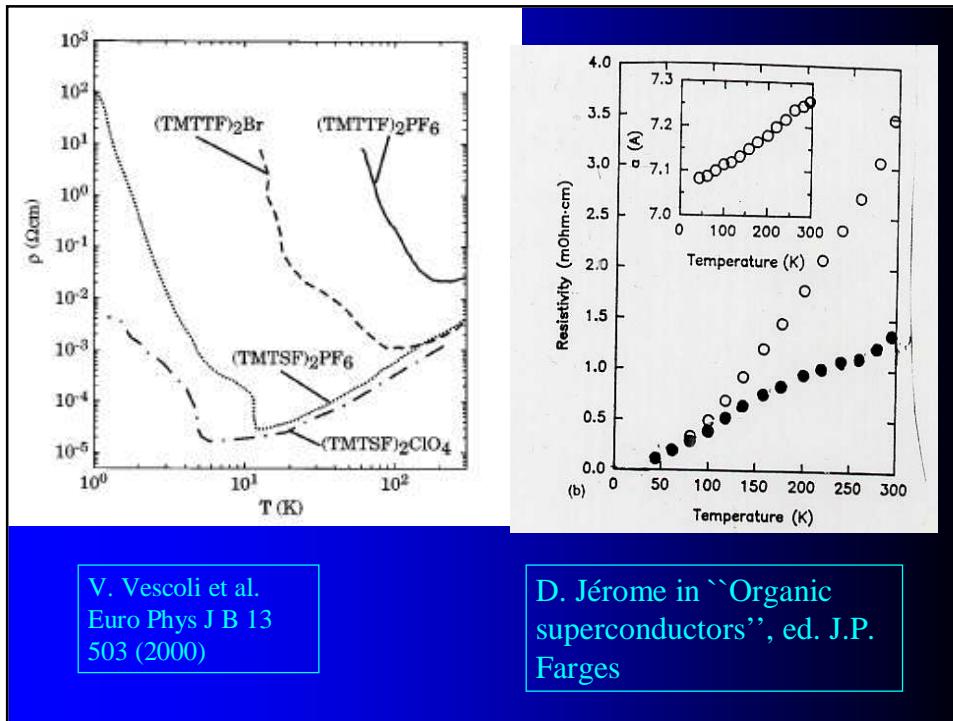
- Propagation of electrons along the chains

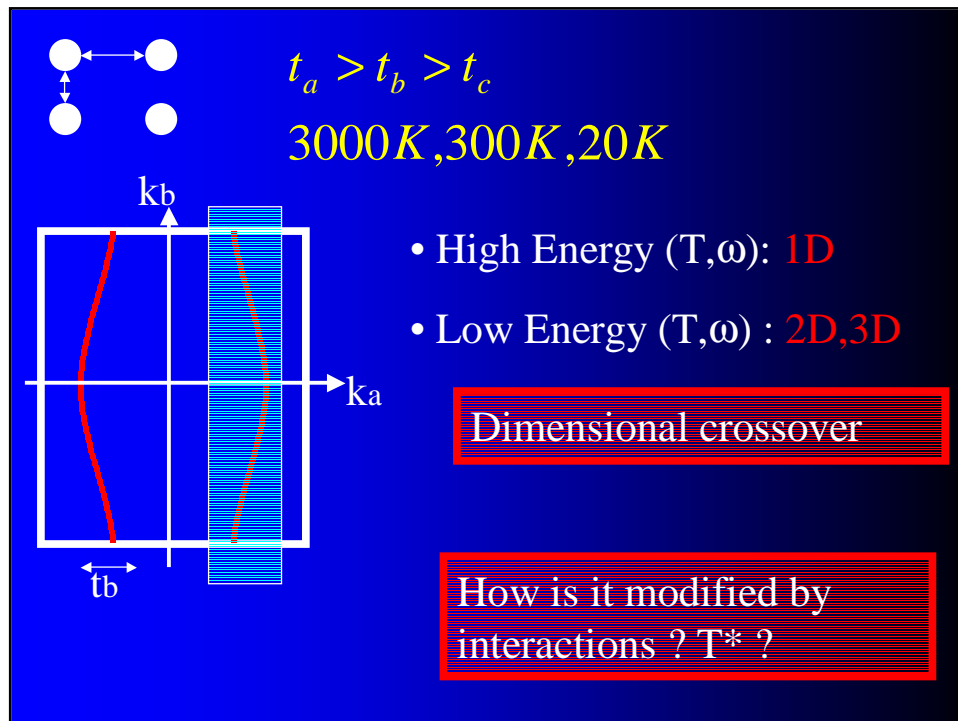
- Chemistry : quarter filled
- Small dimerization: half filled
- Mott insulator
- Is dimerization important ?

	TMTTF			TMTSF		
	PF6	ClO4	Br	PF6	ClO4	NO3
$\Delta = d_1 - d_2$	0.1	0.04	0.03	0.03	0.01	0.01

Deconfinement and Bechgaard salts





## Questions

- Is the high temperature phase a Luttinger liquid/Mott insulator?
- What is the strength of interactions?
- How to describe the dimensional crossover? (scale?)
- What is the resulting 'Fermi liquid'?

## Luttinger liquid

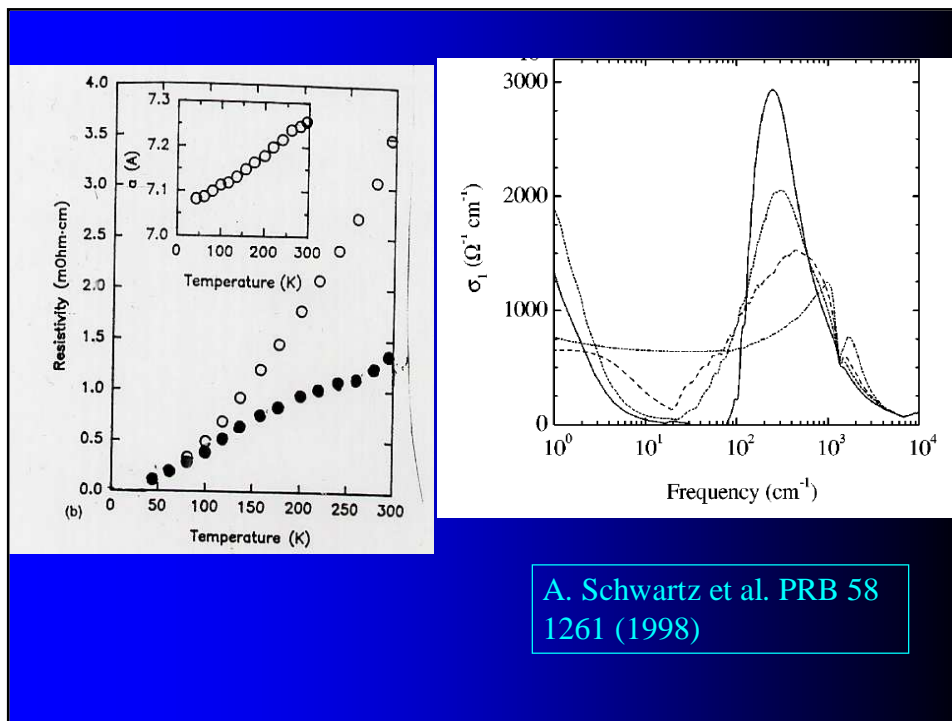
- Spin charge separation



- No fermionic quasiparticles
- Power law decay of correlation functions

$$\langle S(x)S(0) \rangle = \frac{1}{x^2} + \cos(2k_F x) \left(\frac{1}{x}\right)^{1+K_\rho} + \dots$$

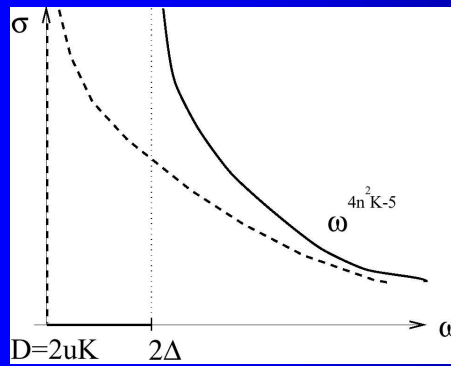
$K_\rho$  contains all information about interactions



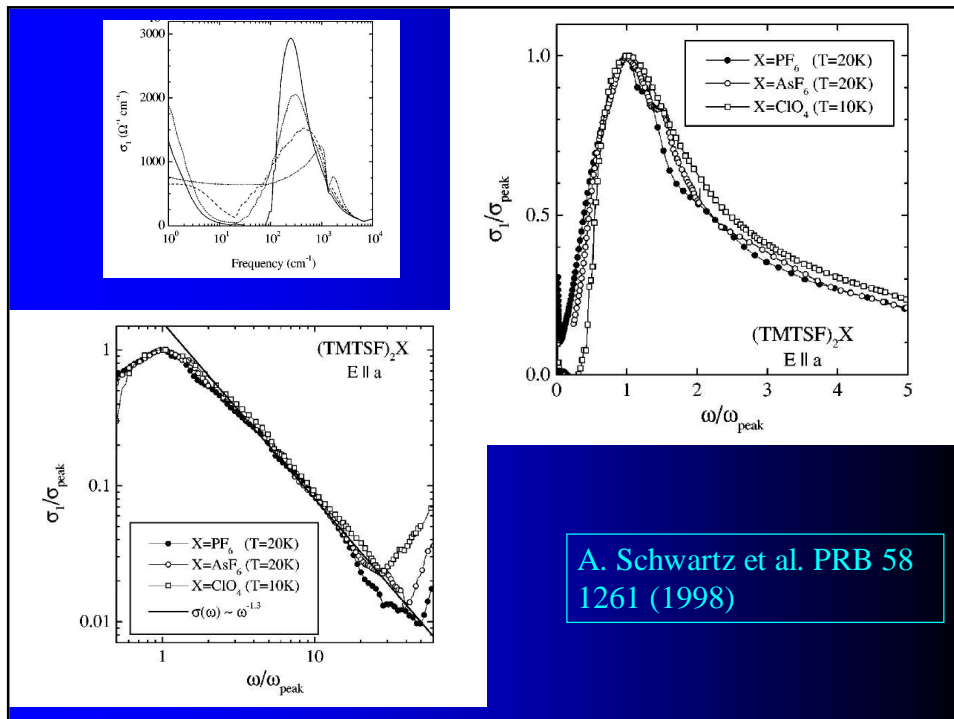
A. Schwartz et al. PRB 58  
1261 (1998)

# Transport in a LL

(TG PRB 44 2905 (91); Physica B 230 975 (97))



- Mott insulator for  $1/4$  filling also !
- Power law in  $\sigma(\omega)$  determines  $K\rho$
- Deviations from 1D law gives  $E_c$

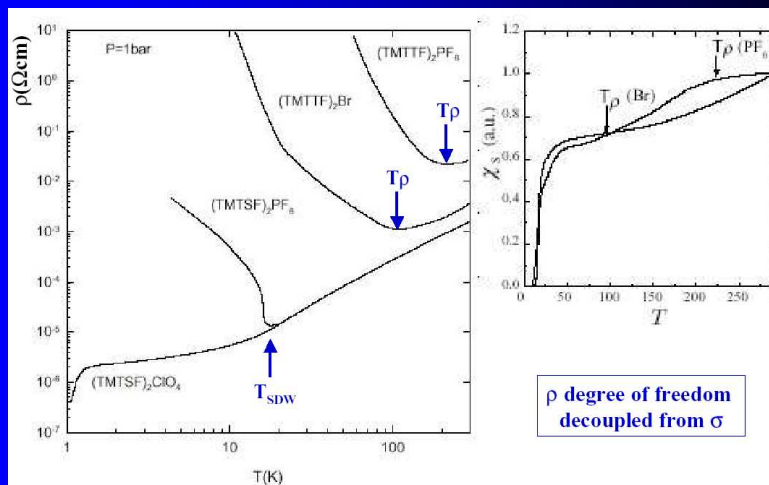


A. Schwartz et al. PRB 58 1261 (1998)

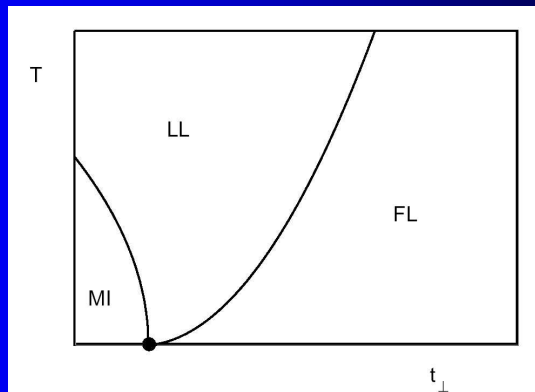
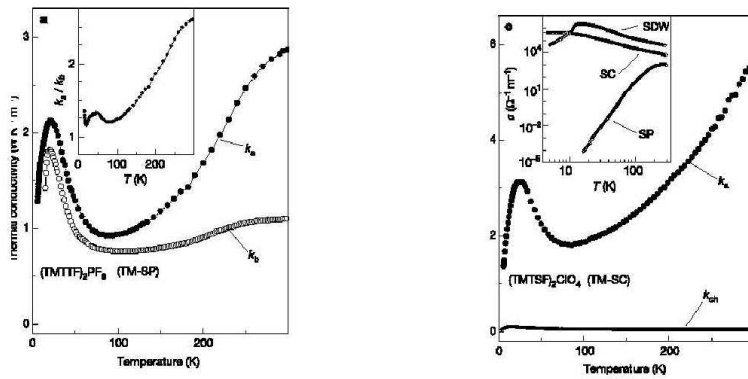
## Consequences

- Luttinger Liquid !
- $K\rho = 0.23$  very strong interactions
- Dimerization not important (1/4 filled Mott insulator) [confirmed in new non-dimerized compound]
- TF vs SF likely to be due to change of interactions (quantum critical point)
- **Dimensional crossover at  $E=100\text{K}$  !**

## Charge spin separation ?



Lorenz et-al, Nature 418, 614, 2002



- Deconfinement : How to study ?
- Difficult (RG, RPA, etc.)

(S. Biermann, A. Georges, TG, A. Lichtenstein, cond-mat 0201542)



- Renormalization arguments

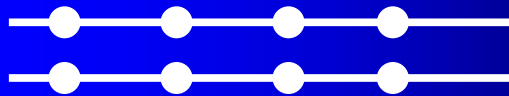
(Bourbonnais, Brazovskii+Yakovenko, Schulz)

$$E^* \sim t_{\perp} (t_{\perp}/t)^{\alpha/(1-\alpha)}$$

$$\alpha = \frac{1}{4}(K_{\rho} + 1/K_{\rho}) - \frac{1}{2}$$

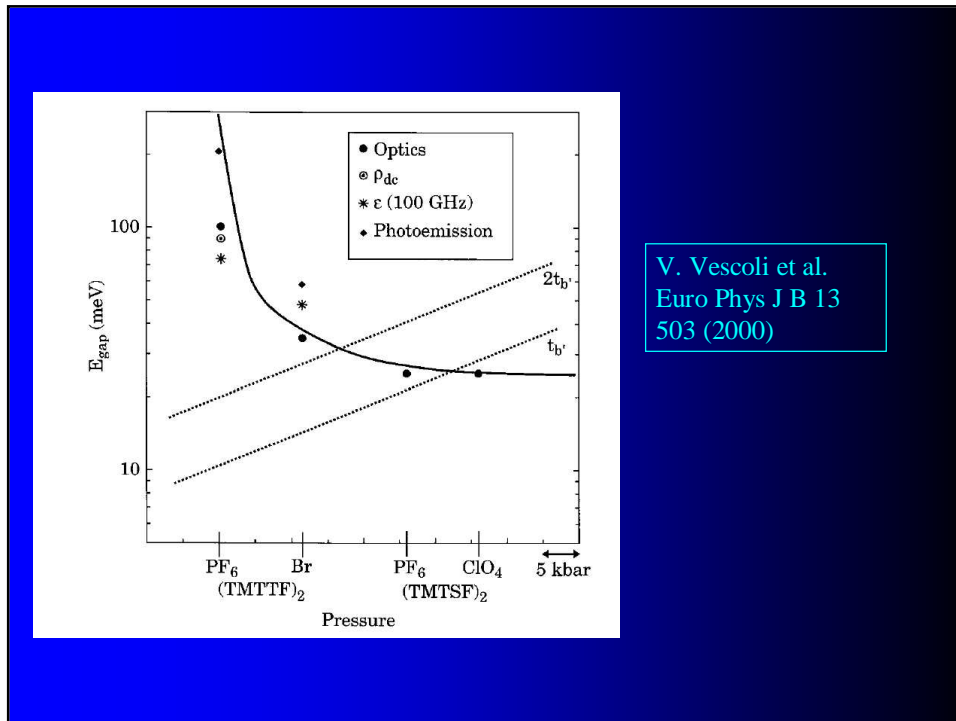
Strong reduction of crossover temperature. But hopping still relevant !

## Mott insulators: confinement



- 1 chain : Mott insulator  $U > 0$
- 3d : Mott insulator  $U > U_c$

Competition Mott insulator/Interchain hopping



## Dimensional Crossover

(A. Georges, TG, N. Sandler PRB 61 16393 (00))

(also E. Arrigoni 1999; 2000)



- $d = \infty + 1$  (chDMFT)

- Effective one dimensional theory

$$S_{\text{eff}} = - \int \int_0^\beta d\tau d\tau' \sum_{ij,\sigma} c_{i\sigma}^+(\tau) \mathcal{G}_0^{-1}(i-j, \tau - \tau') c_{j\sigma}(\tau') + \int_0^\beta d\tau H_{1D}^{\text{int}}[\{c_{i\sigma}, c_{i\sigma}^+\}], \quad (4)$$

$$G(k, i\omega_n) = \int d\epsilon_\perp \frac{D(\epsilon_\perp)}{i\omega_n + \mu - \epsilon_k - \Sigma(i\omega_n, k) - \epsilon_\perp}.$$

$$\text{Re } \sigma_\perp(\omega, T) \propto t_\perp^2 \int d\epsilon_\perp D(\epsilon_\perp) \int \frac{dk}{2\pi} \int d\omega' A(\epsilon_\perp, k, \omega') \times A(\epsilon_\perp, k, \omega + \omega') \frac{f(\omega') - f(\omega' + \omega)}{\omega}.$$

- Self consistent theory for  $\Sigma$
- Feedback of  $t_\perp$  in  $\Sigma$  (a priori important for deconfinement)
- Different from RPA

$$G(k, k_\perp, i\omega_n) = \frac{1}{i\omega_n - \epsilon_k - \epsilon_\perp(k_\perp) - \Sigma_{1D}(k, i\omega_n)}$$

- Difficult to solve the equations analytically

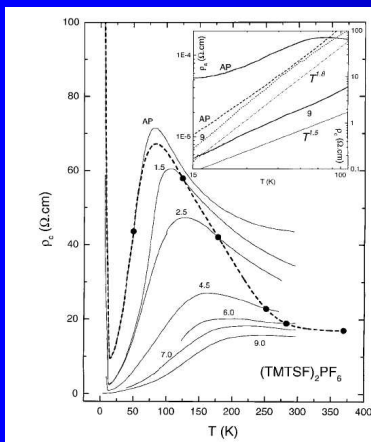
## Transverse transport

$$\sigma(\omega, T) \propto (\omega, T)^{2\alpha-1}$$

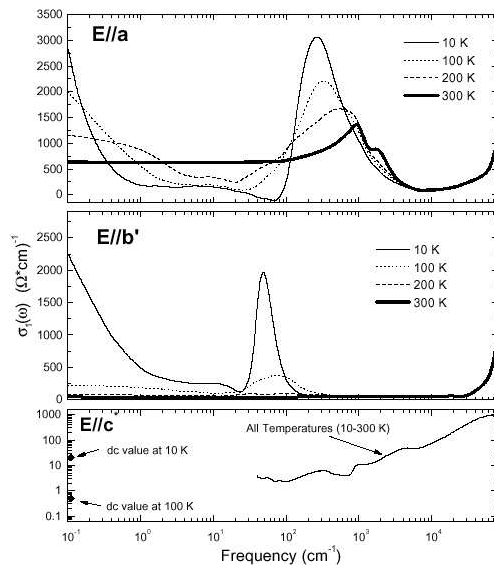
$$\alpha = \frac{1}{4}(K + K^{-1}) - \frac{1}{2}$$

$$\alpha \approx 0.6$$

$$\omega_{D\perp}^2 / \omega_{P\perp}^2 \propto (t_{\perp} / t)^{2\alpha/(1-\alpha)} = Z^2$$



J. Moser et al. Euro Phys. J. B 1 39 (1998)



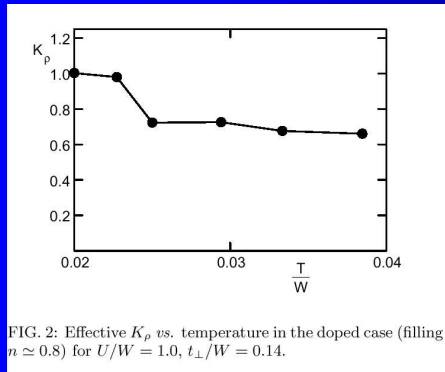
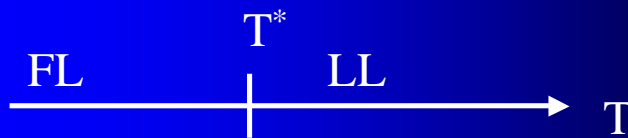
V. Vescoli et al. Euro Phys J B 11 365 (1999)

## Numerical Solution

(S. Biermann, A. Georges, A. Lichtenstein, TG, PRL 87 276405 (2001))

- Hubbard Model
- QMC (16 – 32 sites)
- 32 time slices ( $T/W = 1/50$ )

## Incommensurate case



$$T^* \approx \frac{t_\perp}{\pi} \left( \frac{t_\perp}{t} \right)^{\frac{\theta}{1-\theta}}$$

$$T^* \approx 0.5 \frac{t_\perp}{\pi}$$

## Commensurate case

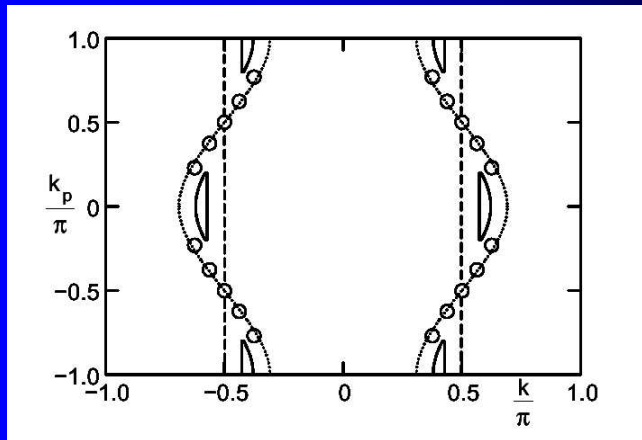


TABLE II: Effective  $K_\rho$  at half-filling, as a function of  $t_\perp/W$  for  $U/W = 0.65$  and  $T/W = 1/40$ .

$t_\perp/W$	0.00	0.04	0.07	0.11	0.14	0.16	0.18
$K_\rho$	0.00	0.02	1.01	1.09	1.07	1.06	1.04

$$t_\perp^* \approx \Delta_{1D}$$

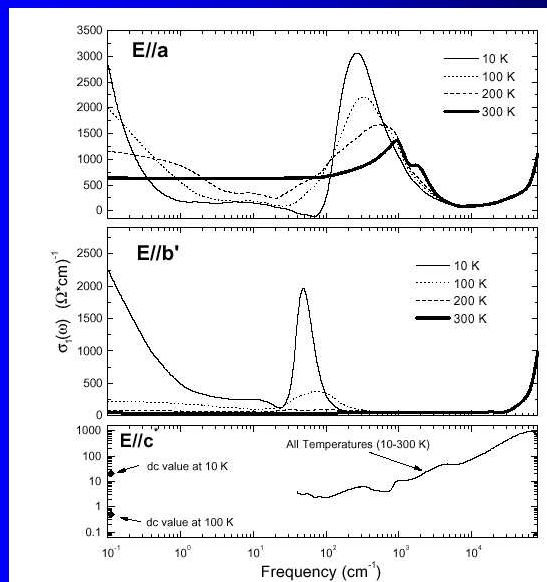
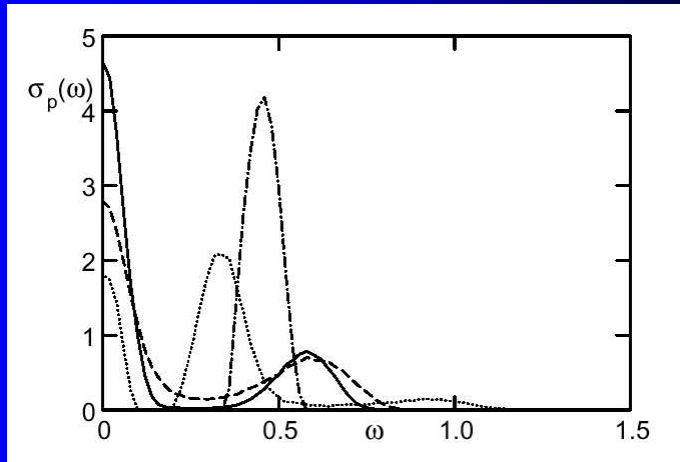
## Fermi Surface



Z

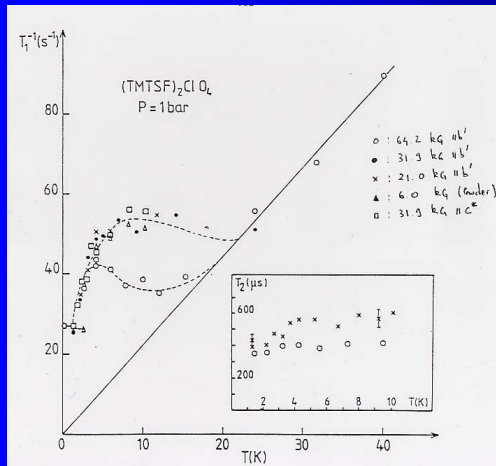
$k_\perp/\pi$	0.23	0.38	0.50	0.62	0.77
$Z(k_\perp)$	0.78	0.77	0.76	0.77	0.79

## Transverse conductivity



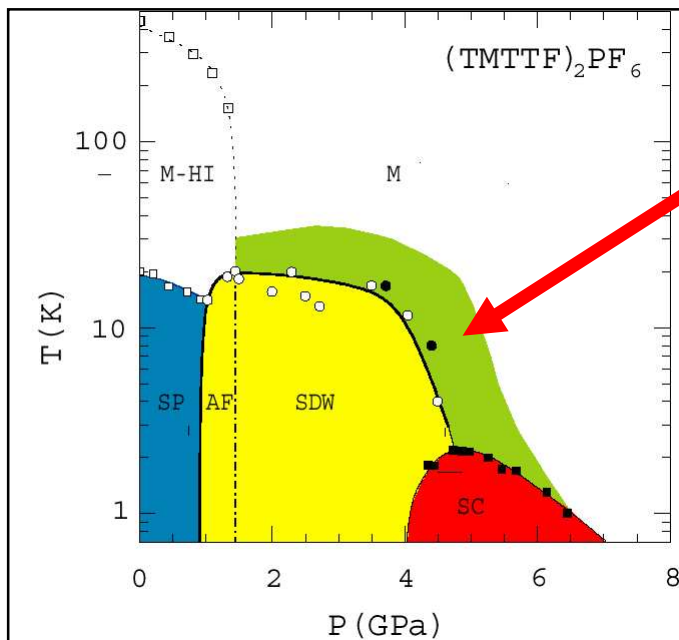
V. Vescoli et al. Euro Phys J B 11  
365 (1999)

## Puzzle from NMR



F. Creuzet et al., J. Phys. Lett 45 L755 (84)

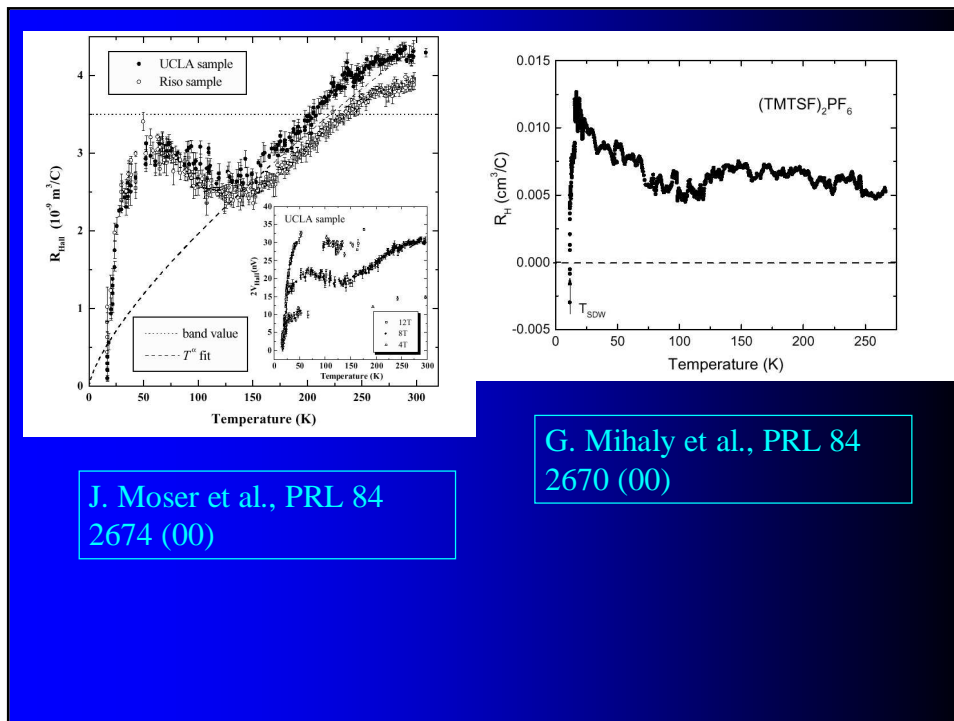
Non fermi liquid below E<sub>c</sub> ?



?

After D. Jérôme





## Hall Effect

(A. Lopatin, A. Georges, TG PRB 62 (00))

If no momentum relaxation (no umklapp) along the chains:

$$\rho_{yx} = \frac{H}{ne c} \frac{2\alpha k_F}{v_F}$$

With umklapp : scaling expression for Hall

$$R_H(T) = R_H^0 \left( 1 + a_1 \frac{g_3}{T^{1/x}} + a_2 \frac{g_3^2}{T^{2/x}} + \dots \right)$$

$$2x = 1/(1 - n^2 K_\rho)$$

Plot of  $R_H$  vs  $\rho_{xx}/T$

## Conclusions

- Transport proves LL nature of high energy phase
- $K=0.23$ ,  $\frac{1}{4}$  filled Mott insulators
- $E_c = 100K$  much higher than expected
- Good method to tackle the dimensional crossover

- $\frac{1}{4}$  Filling with chDMFT
- Other physical quantities : Hall effect
- Nature of the 2D phase ? 2D non Fermi liquid ?
- Ordered phases (superconductivity)