

Modeling Ca cycling in cardiac cells: from ion channels to whole cell dynamics

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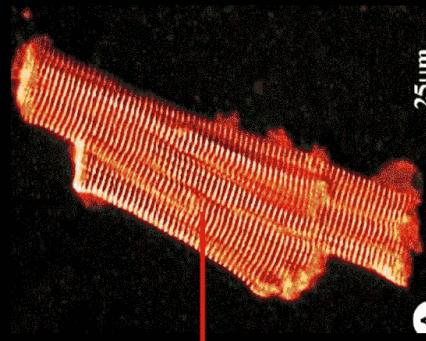
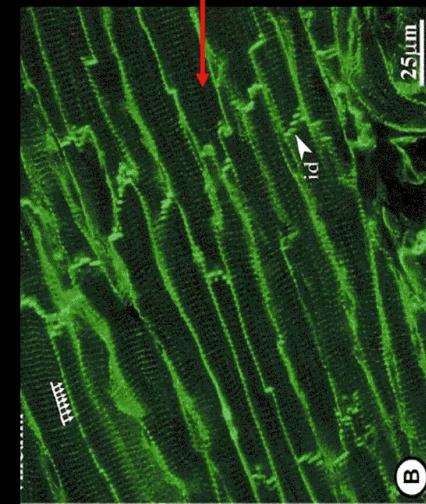
Collaborators:

Alain Karma, Daisuke Sato, Northeastern Univ.
Robert Rowett, Alan Garfinkel, Zhilin Qu, and James Weiss, UCLA

Outline

- Overview of some important features of Ca cycling in cardiac cells.
- Developing a phenomenological model of Ca cycling
- Some features of the nonlinear dynamics of paced cells.
- Relating whole cell properties to Ca signalling at the ion channel level.
- Summary

Cardiac tissue and cells

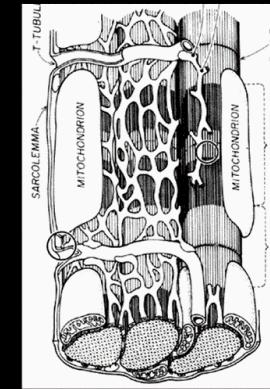
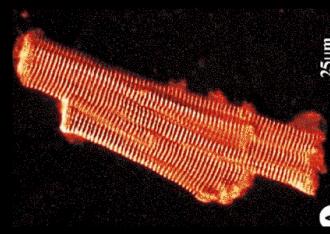
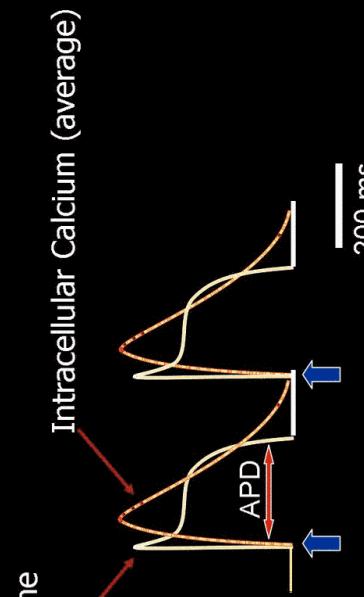


Function

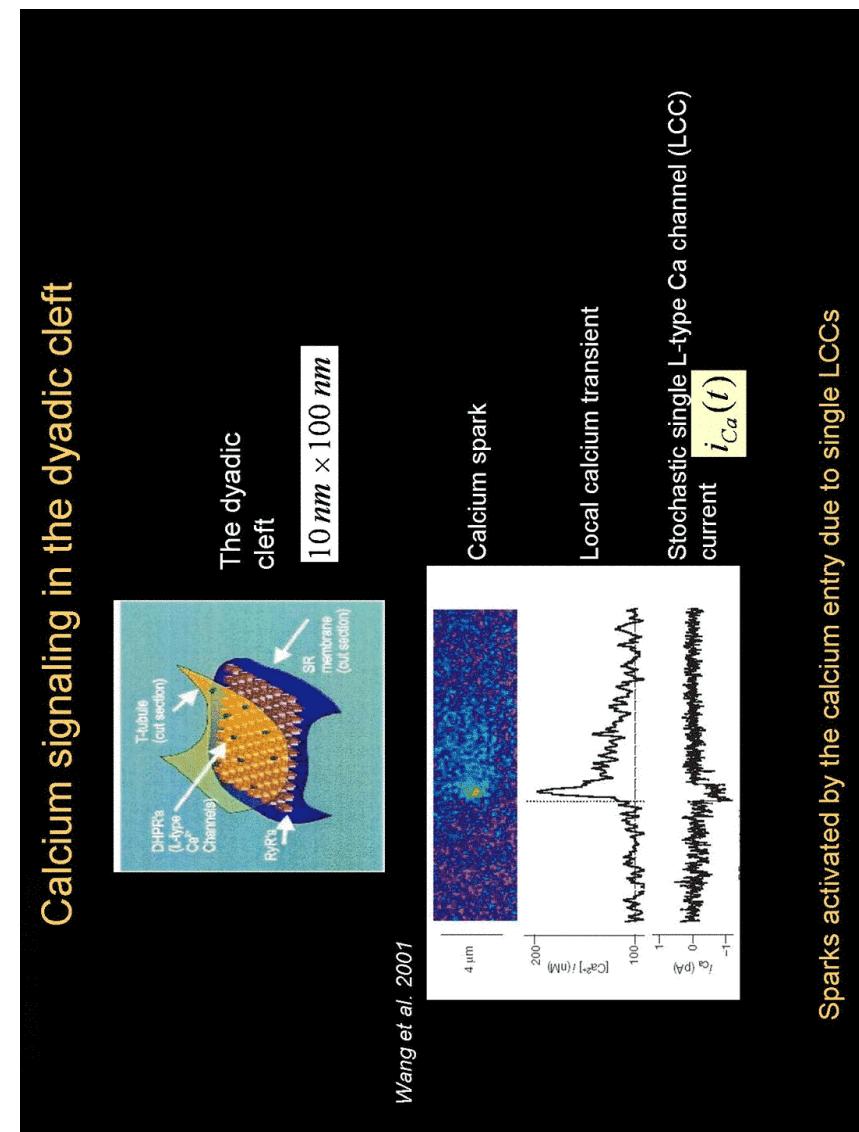
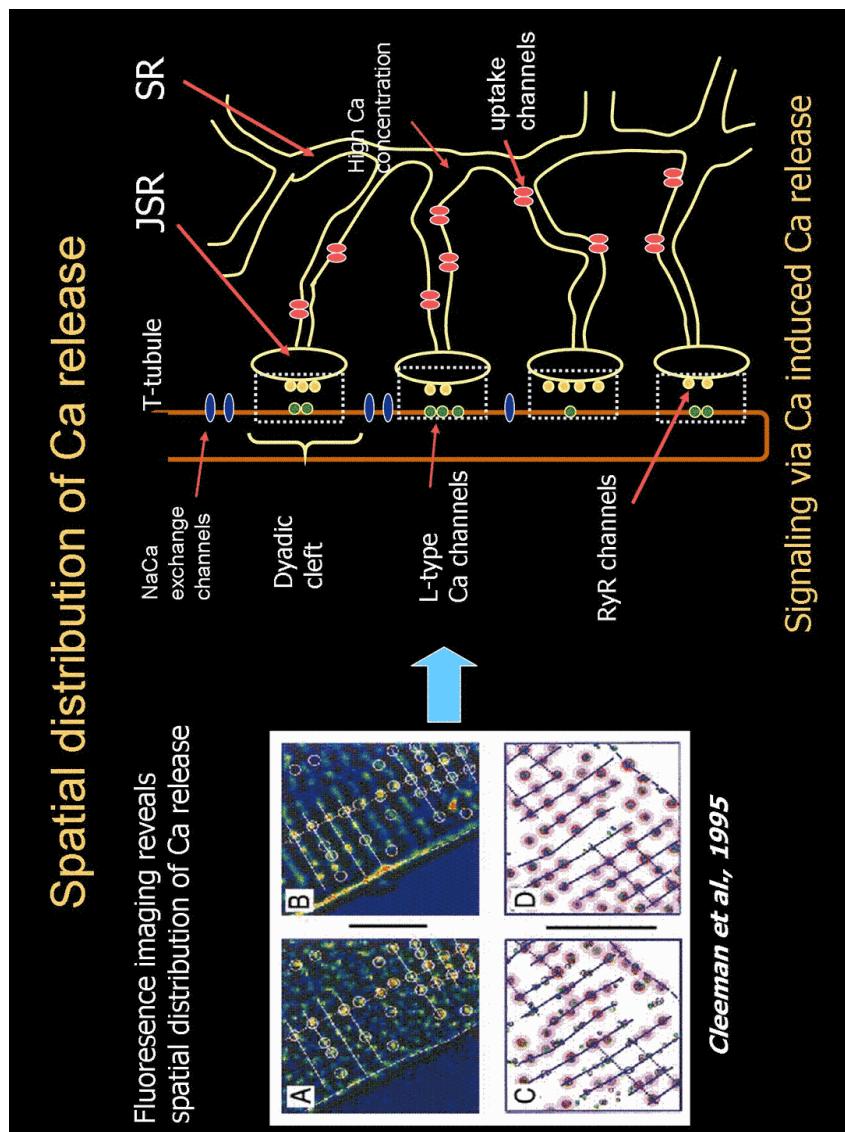
- Conduct electrical waves
- Contract in response to an electrical stimulus

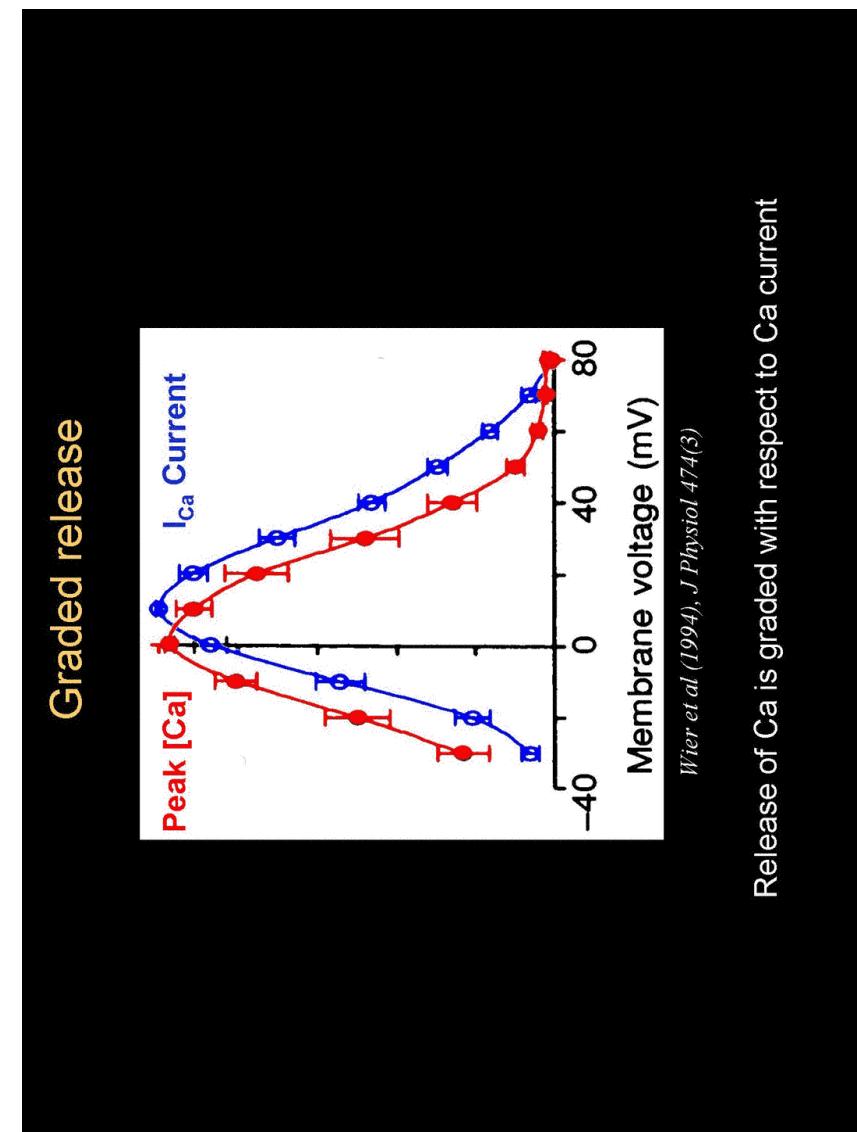
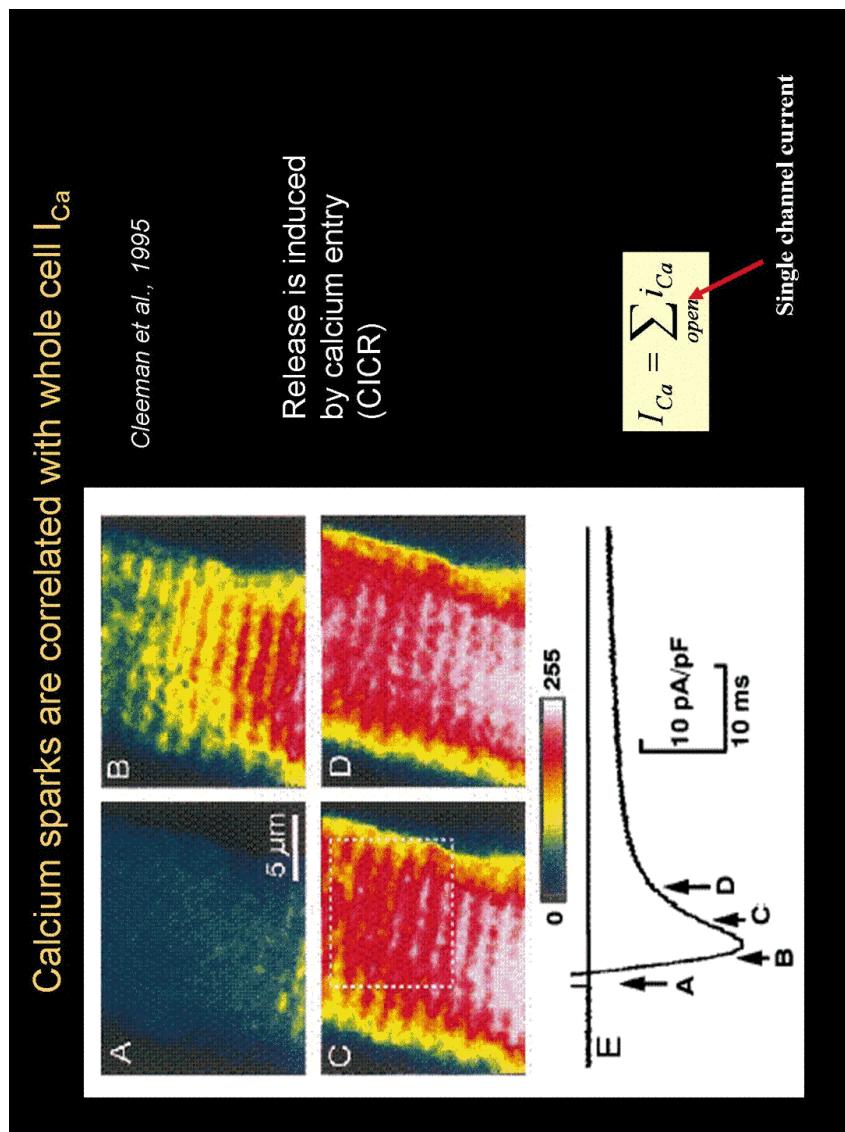
Basic Cell physiology

Action potential regulates intracellular calcium



Calcium → Contraction





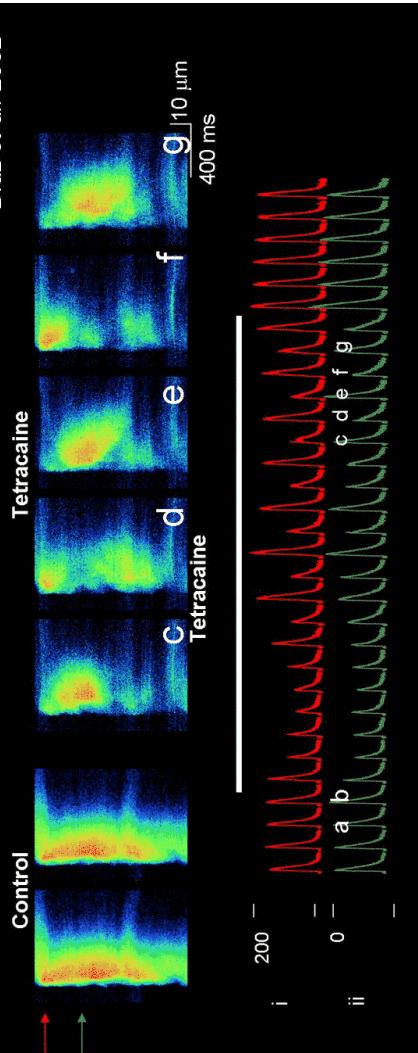
Paced cardiac cells exhibit nonlinear dynamical properties

Ca transient alternans observed in rabbit Ventricular, when cell is paced with periodic voltage clamp.

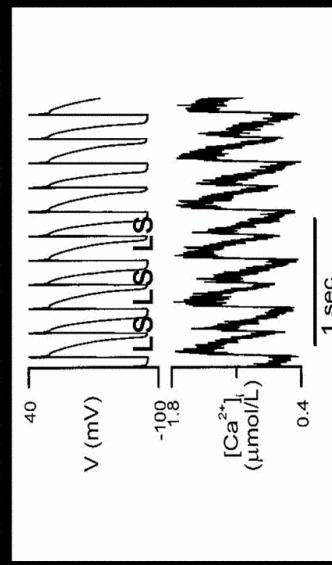
From lab of J. Weiss, UCLA



Diaz et al. 2002

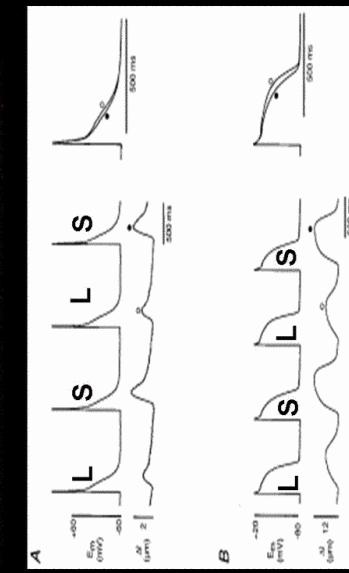


Concordant electro-mechanical alternans



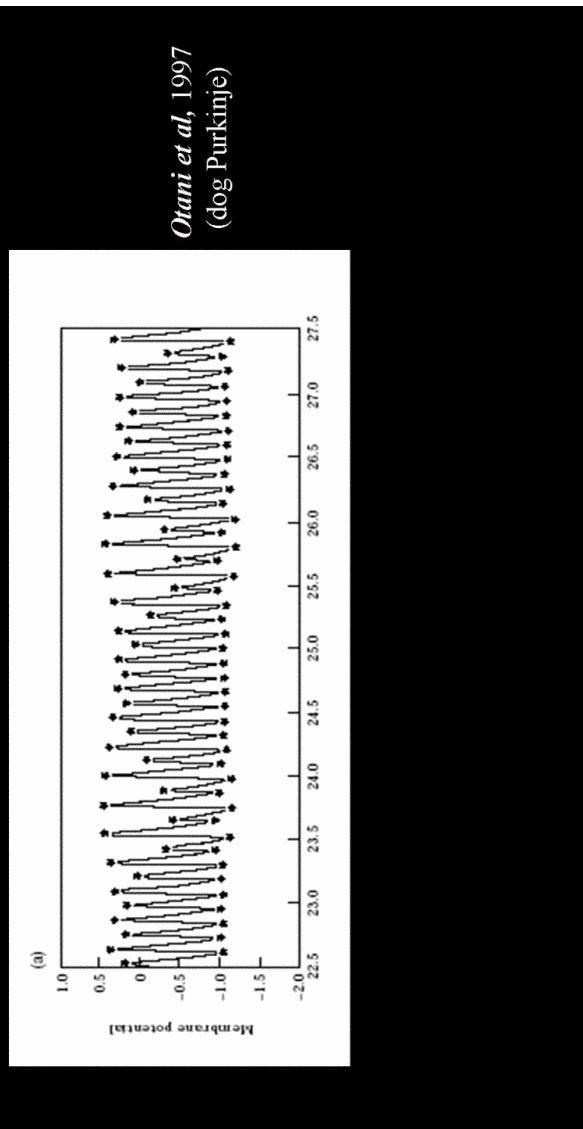
(Rabbit)
Chudin et al. 1999

Discordant electro-mechanical alternans

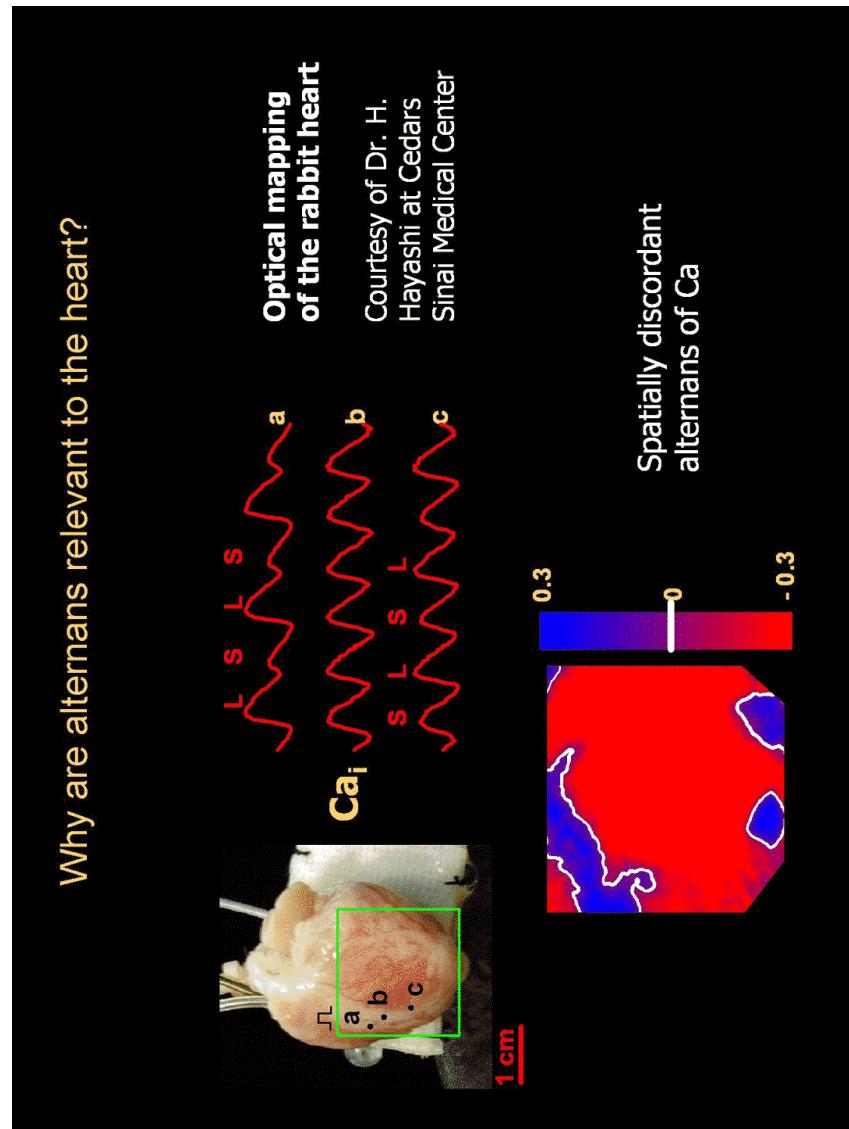


(guinea pig)
J. Houser, et al. 2000

Quasiperiodic dynamics



Why are alternans relevant to the heart?



Can we construct a model of Ca cycling where:

1. Release is triggered by Ca entry via a ClCR mechanism operative within $\sim 10^4$ dyadic cleft.
2. Release occurs via discrete release events called Ca sparks.
3. Release mirrors the whole cell Ca entry (graded release).
4. Can be used to understand the rich dynamical behavior of Ca cycling.

A Phenomenological approach:

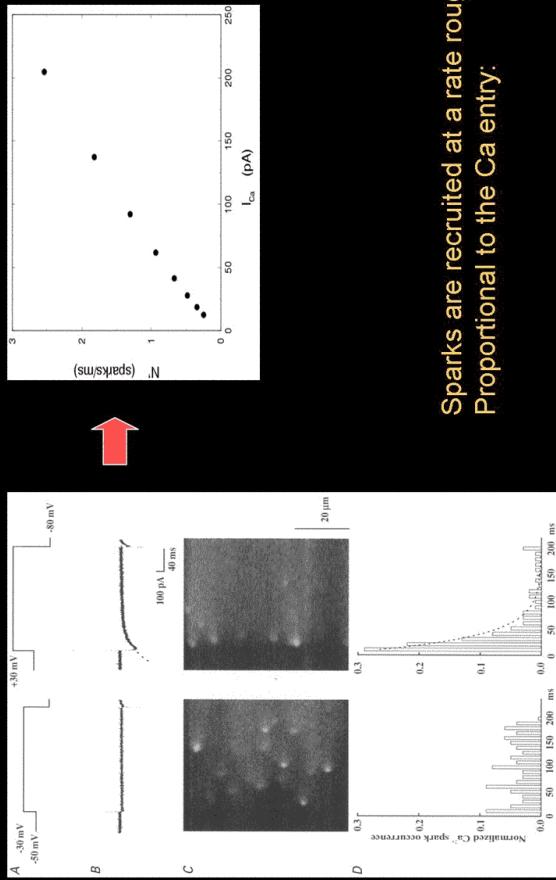
Phenomenological model of Ca release at the dyadic cleft (spark):

$$I_{\text{spark}}(t) = J(c_j') e^{-(t-t')/\tau_s}$$

Concentration of JSR

Duration of spark $\approx 10 - 20 \text{ ms}$

Spark recruitment rate



Sparks are recruited at a rate roughly proportional to the Ca entry:

$$\frac{dN_s(t)}{dt} = gI_{Ca}(t)A(c_j'(t))$$

Summing discrete release events

Whole cell current due to summation of discrete Ca fluxes:

$$I_r(t) = \int_{t_0}^t \frac{dN(t')}{dt'} J(c_j'(t')) e^{-(t-t')/\tau_s} dt'$$

Satisfies

$$\frac{dI_r(t)}{dt} = gI_{Ca}(t)\mathcal{Q}(c_j'(t)) - \frac{I_r(t)}{\tau_s}$$

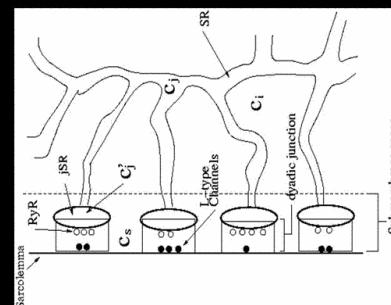
$$\mathcal{Q}(c_j'(t)) = A(c_j'(t))J(c_j'(t))$$

Full Model Equations for Ca cycling

$$\frac{dc_s}{dt} = \frac{\alpha_i}{\alpha_s} \left(I_r(t) - \frac{c_s - c_i}{\tau_d} - I_{Ca}(c_s, V(t)) + I_{NaCa}(c_s, V(t)) \right)$$

$$\frac{dc_i}{dt} = \frac{c_s - c_i}{\tau_d} - I_{up}(c_i)$$

$$\frac{dc_j}{dt} = \frac{\alpha_i}{\alpha_j} \left(-I_r(t) + I_{up}(c_i) \right)$$



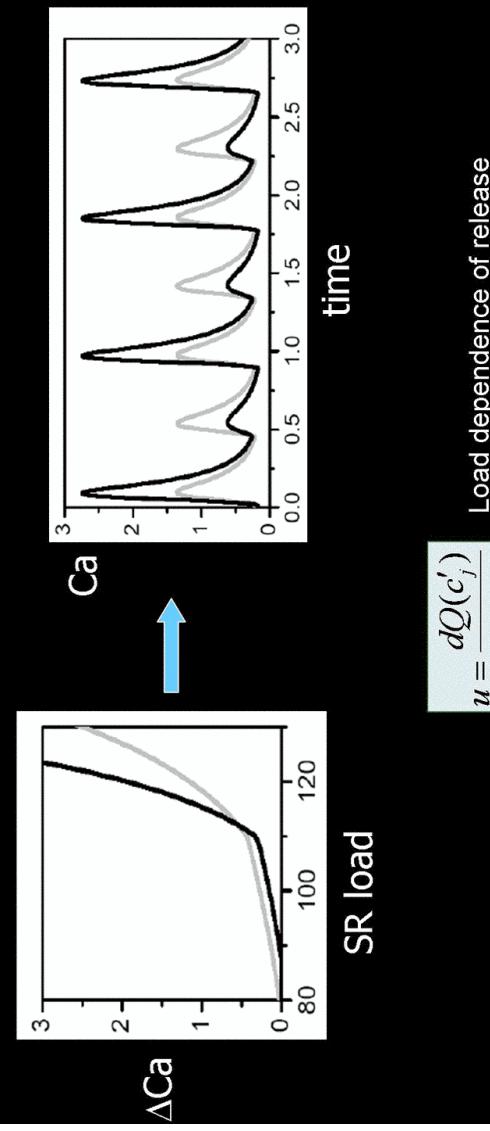
Biophys J., 2003

$$\frac{dI_r(t)}{dt} = gI_{Ca}(t)\mathcal{Q}(c_j'(t)) - \frac{I_r(t)}{\tau_s}$$

$$\frac{dc_j'(t)}{dt} = \frac{c_j(t) - c_j'(t)}{\tau_j}$$

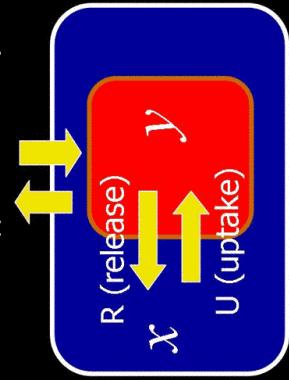
Instability Mechanism

Steep SR-release vs. SR-load relationship induces alternans



Nonlinear map reduction

Δ (NaCa exchanger and L-type Ca current)



$$\begin{aligned}x_{n+1} &= x_n + R_n - U_n + \Delta_n \\y_{n+1} &= y_n - R_n + U_n\end{aligned}$$

Calcium cycling dynamics

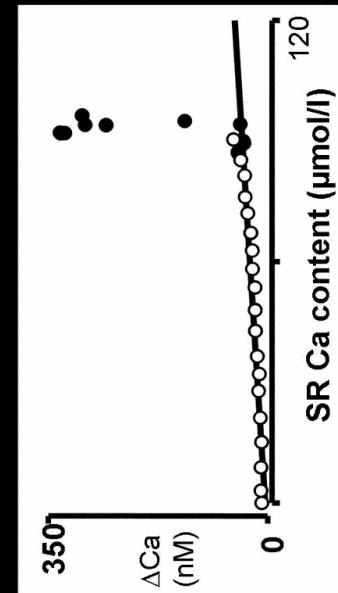
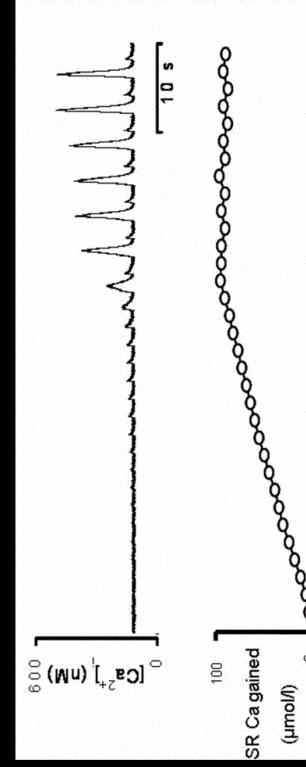
Shiferaw et al., 2005

Alternans occurs when:

$$(1-\varphi) \left(-1 - \frac{\partial R}{\partial x_n} + \frac{\partial R}{\partial y_n} \right) > 1$$

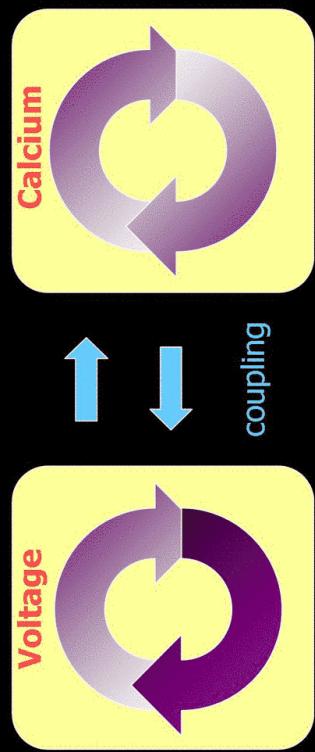
Experimental Validation

Diaz et al., 2004



Nonlinear dynamics of voltage and Ca cycling

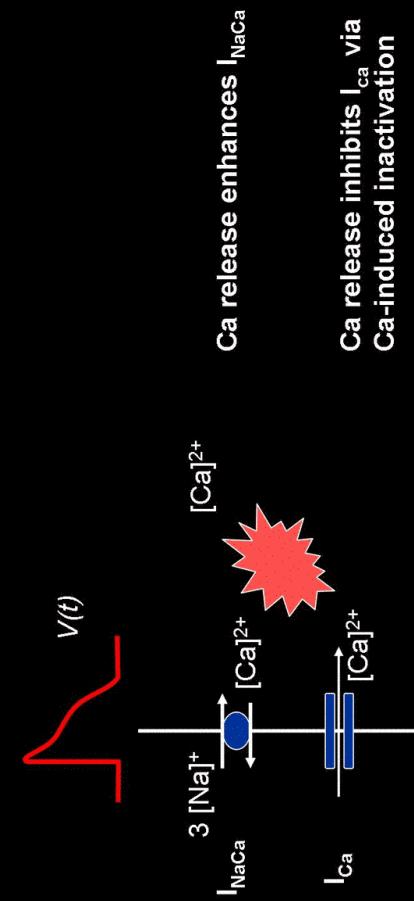
Two coupled nonlinear systems

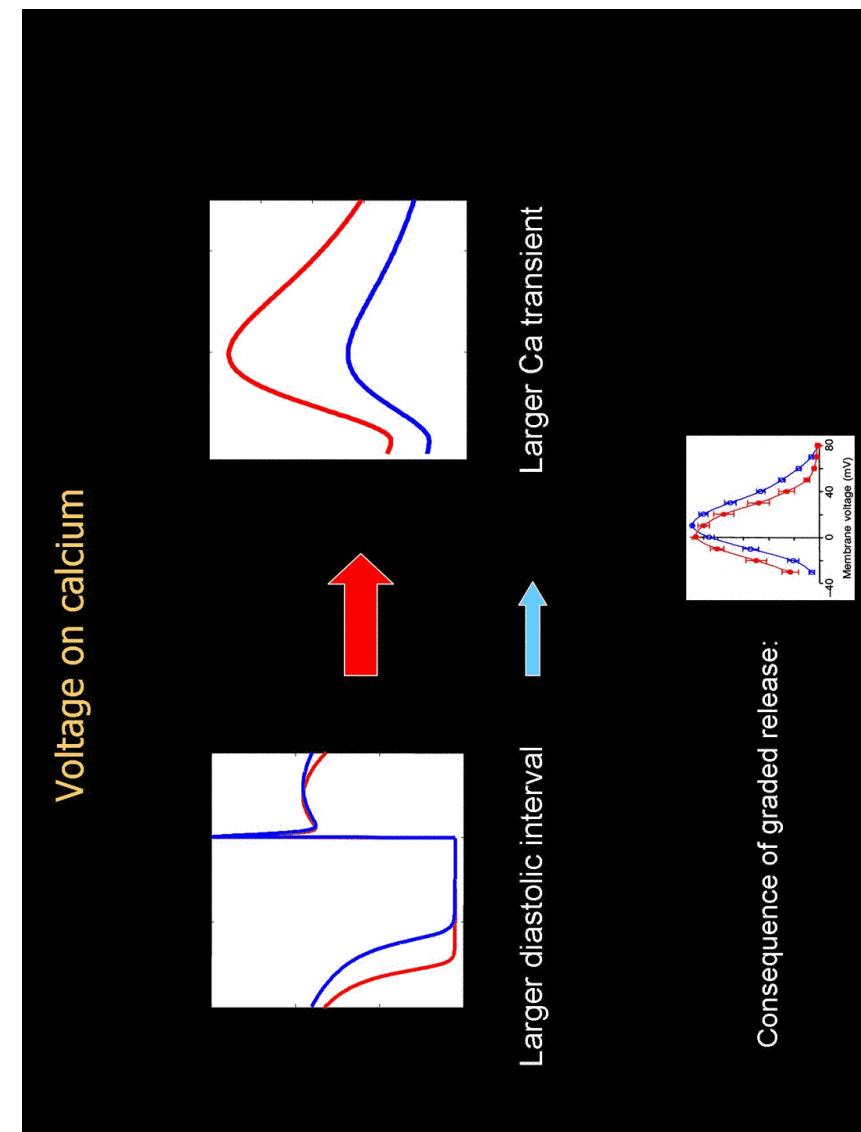
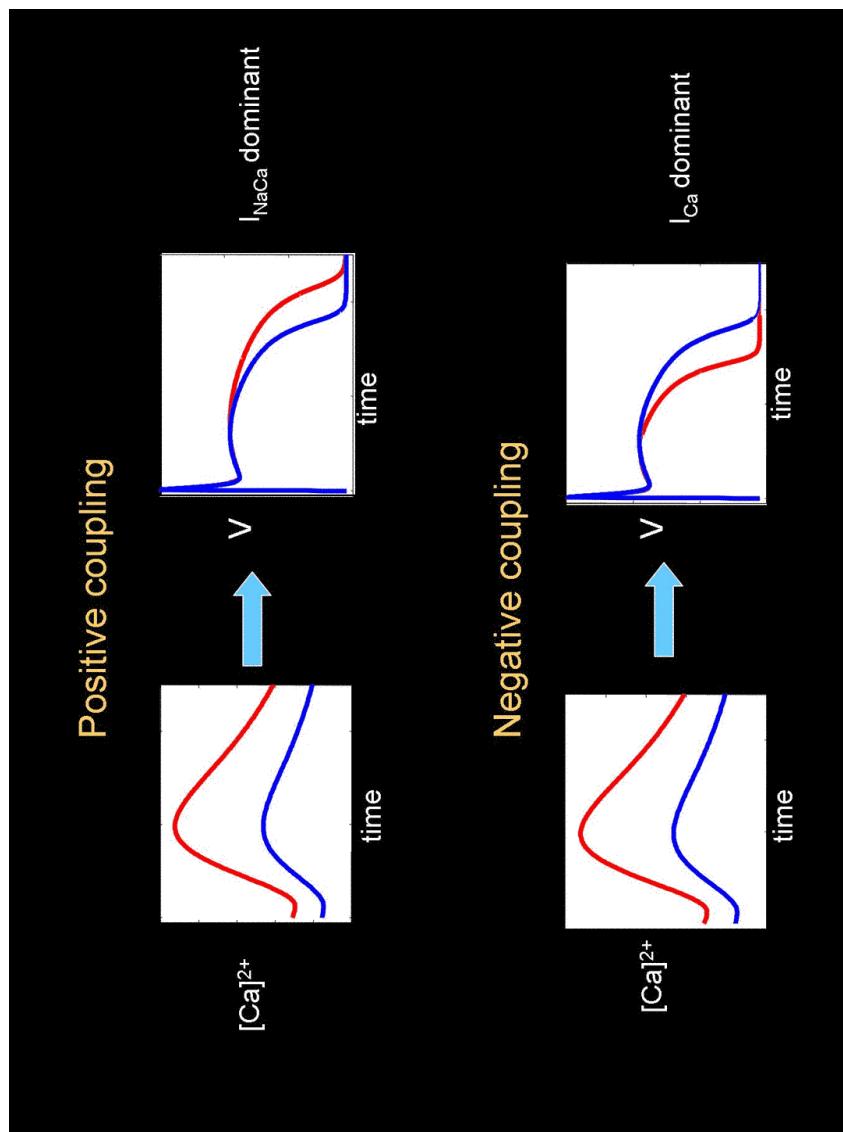


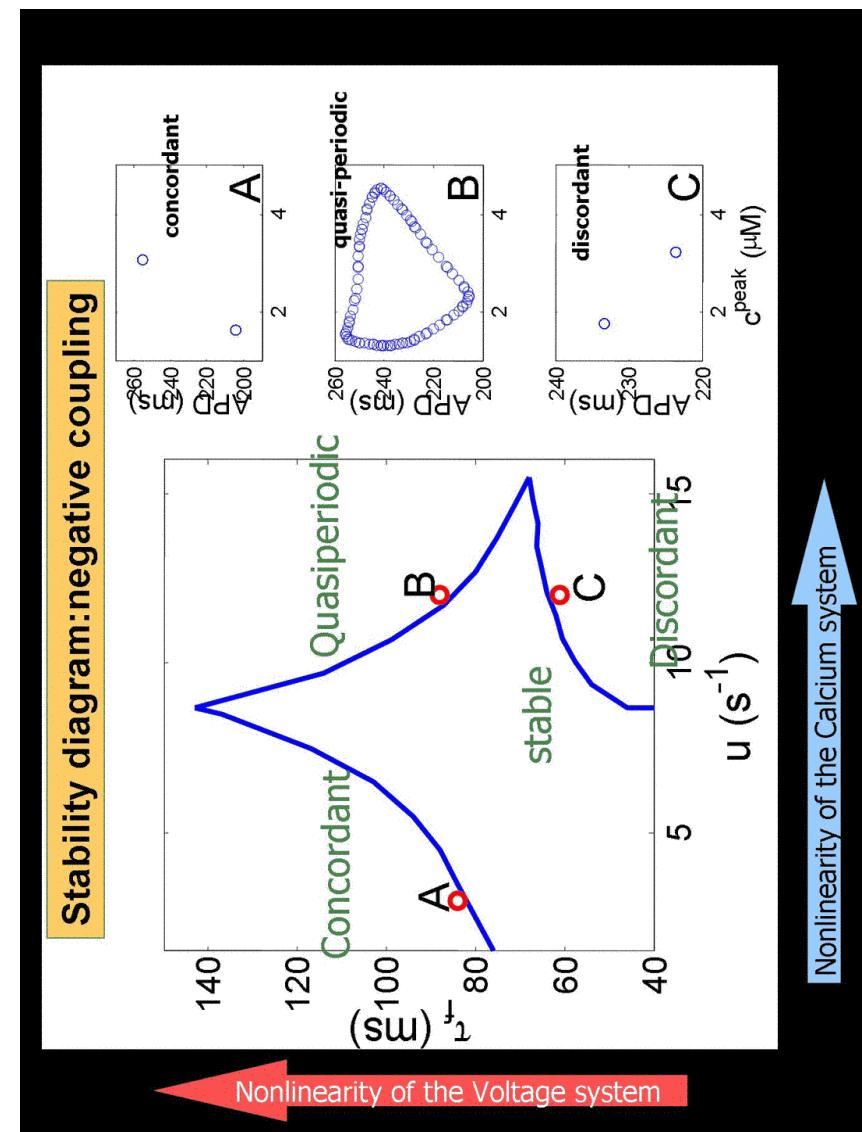
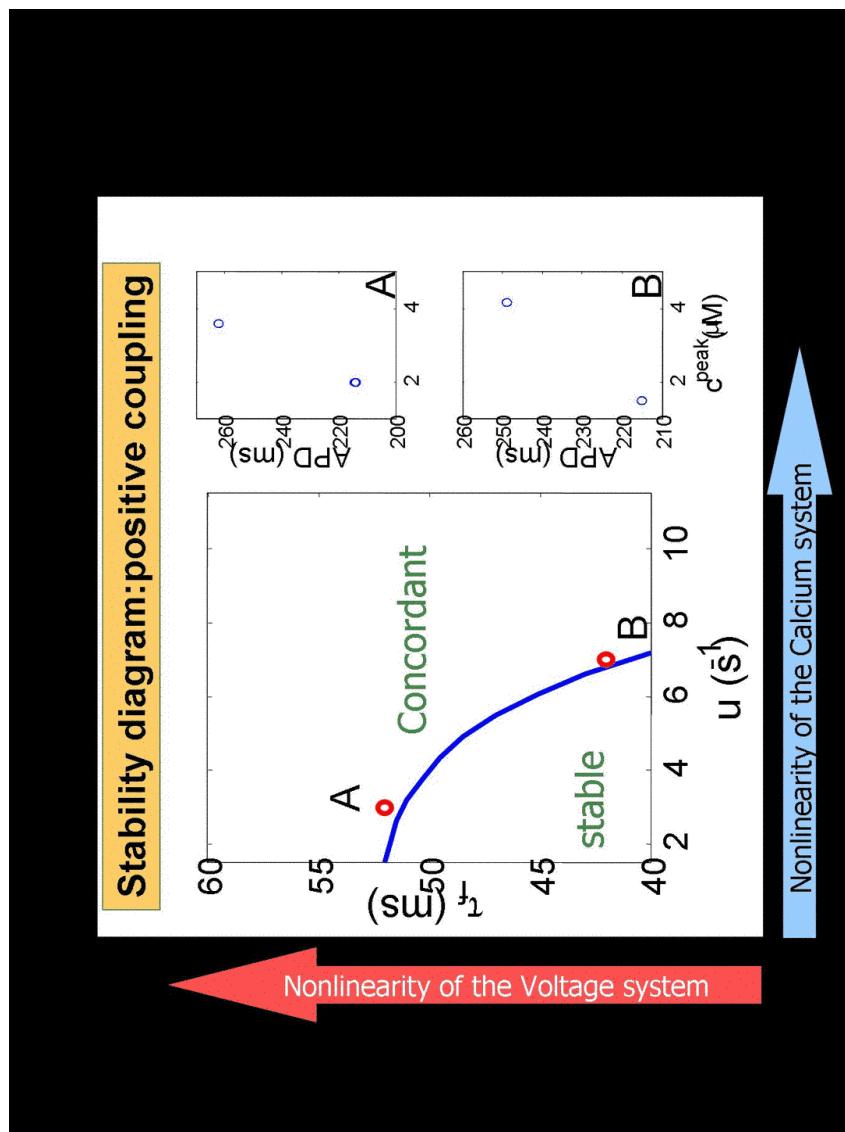
Bi-directional coupling

$$\dot{V} = -(I_{Na} + I_{K1} + I_{Ks} + I_{Kr} + I_{Kp} + I_{to} + I_{NaCa} + I_{Ca}) / C_m$$

Fox et al., 2002



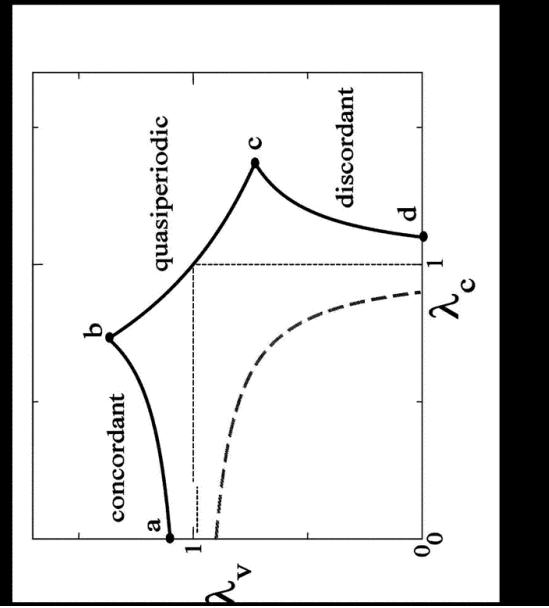




Discrete maps

$$\begin{aligned}x_{n+1} &= x_n + R_n - U_n + \Delta_n \\y_{n+1} &= y_n - R_n + U_n\end{aligned}$$

$$A_{n+1} = F(D_n, x_n, y_n)$$



The restitution slope.

$$\lambda_v = \frac{\partial F}{\partial D_n}$$

$$\lambda_c = -1 - \frac{\partial(R_n - U_n)}{\partial x_n} + \frac{\partial(R_n - U_n)}{\partial y_n}$$

Degree of instability of Ca cycling.

$$C = \frac{\partial(R_n - U_n)}{\partial D_n} \left(\frac{\partial F}{\partial y_n} - \frac{\partial F}{\partial x_n} \right)$$

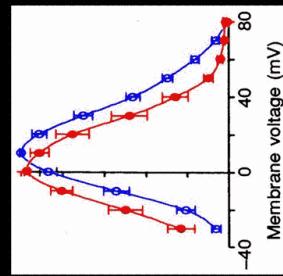
Degree of coupling between voltage and Ca systems.

Limitations

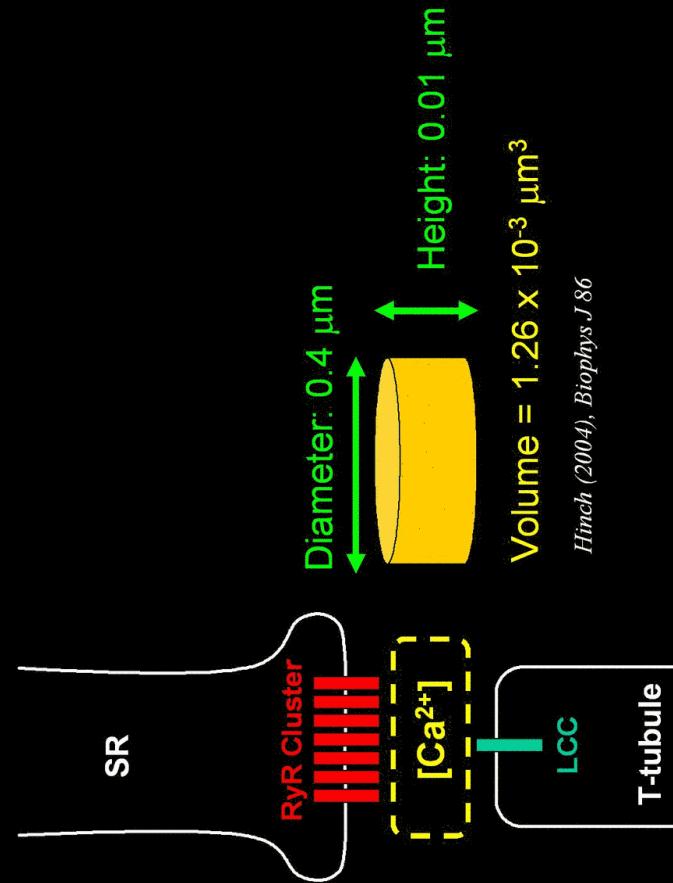
Model does not yield insight on the role of ion channels.

Can we understand the basic whole cell properties in terms of the kinetics and interaction of ion channels?

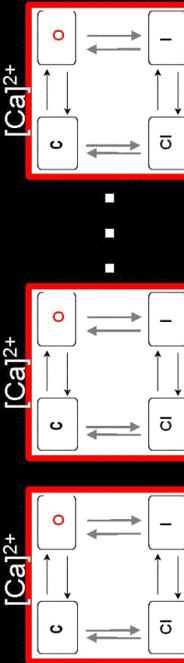
For example, can we explain graded release in terms of Ca signaling between ion channels in the dyadic cleft??



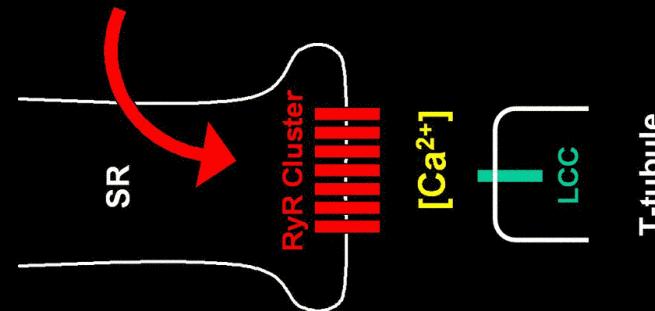
Computational model of a single dyad

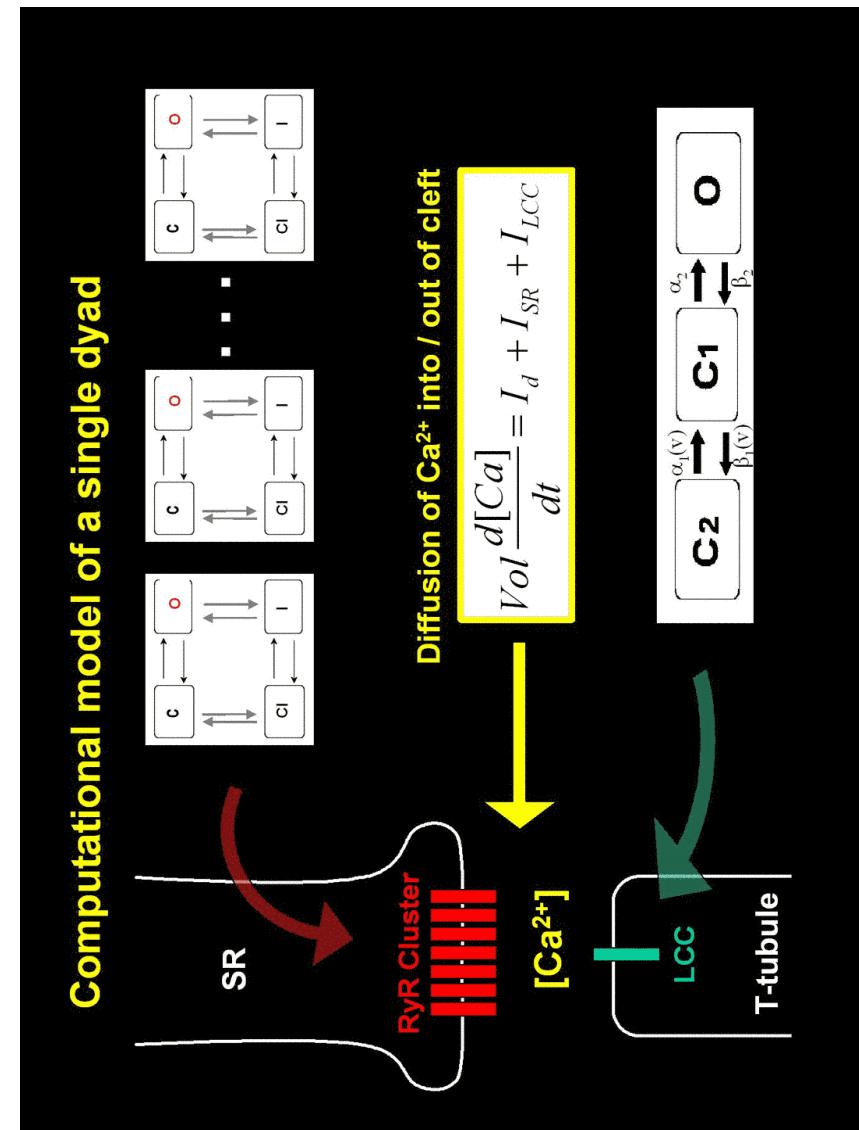
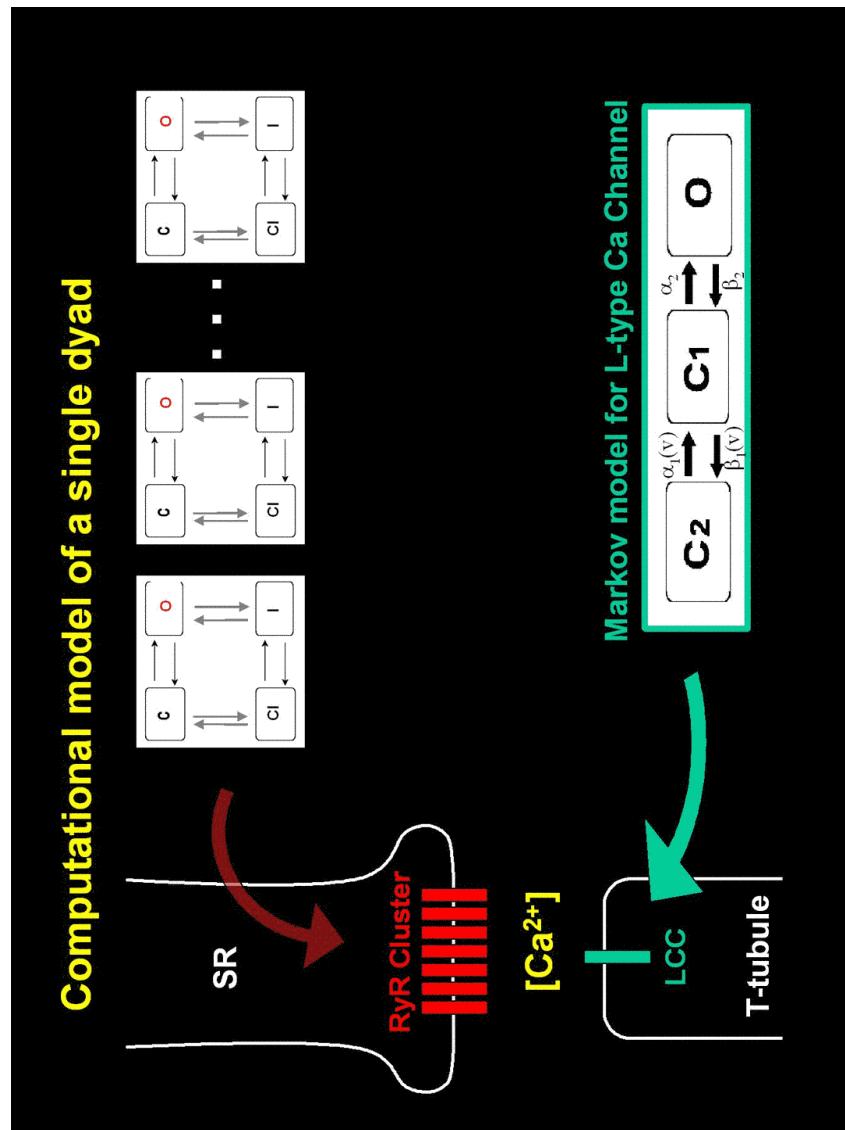


Modified from Stern (1999), *J Gen Physiol* 113(3)

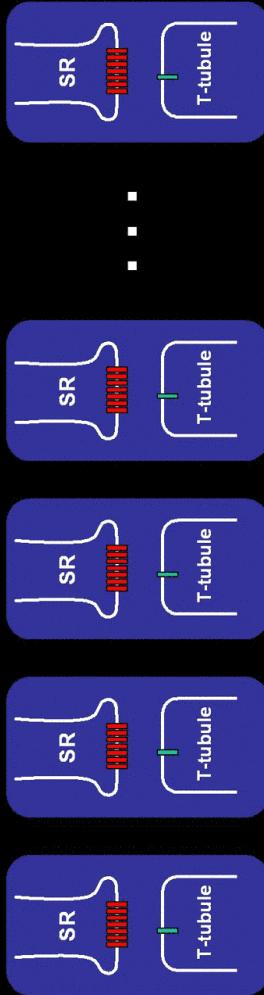


Multiple channels modeled simultaneously





Whole-cell data is composed of multiple independent simulations

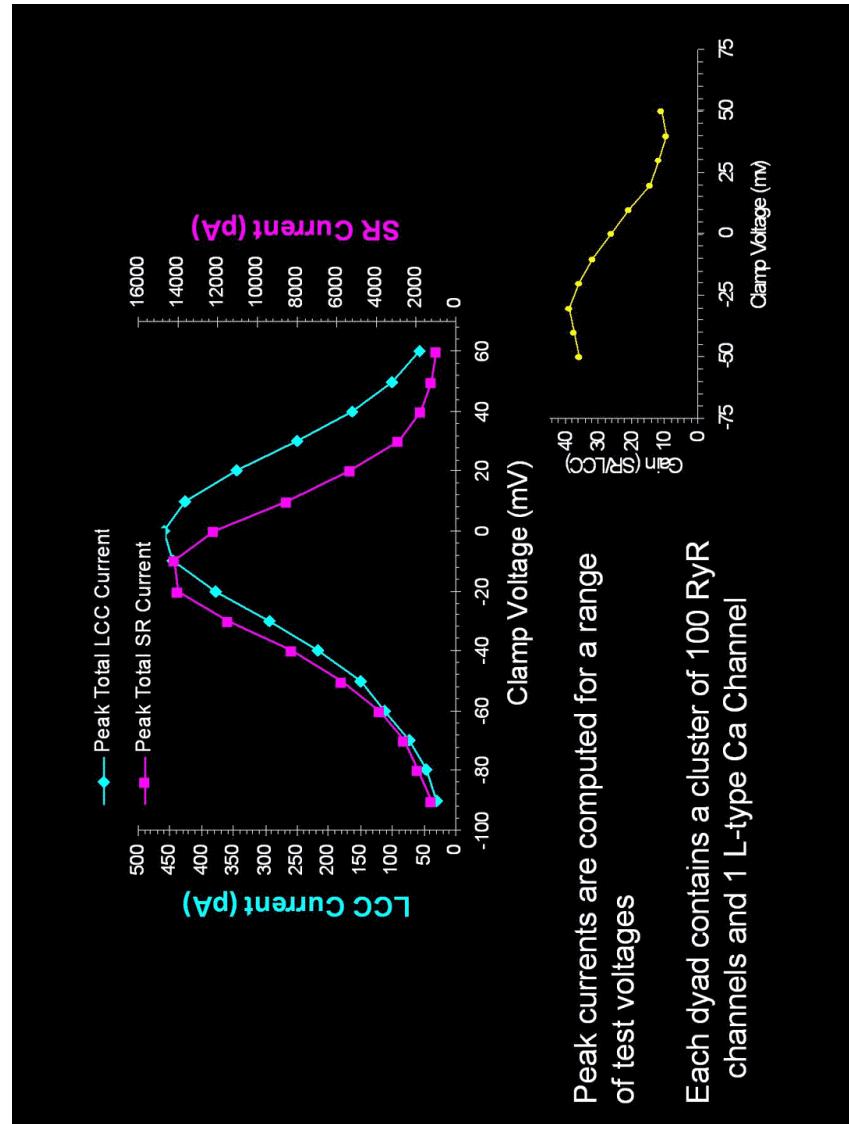


~ 100 RyR channels per dyadic junction
~ 50,000 independent dyadic junctions per cell

About 5 million computations are done at each time step (1 μ s)

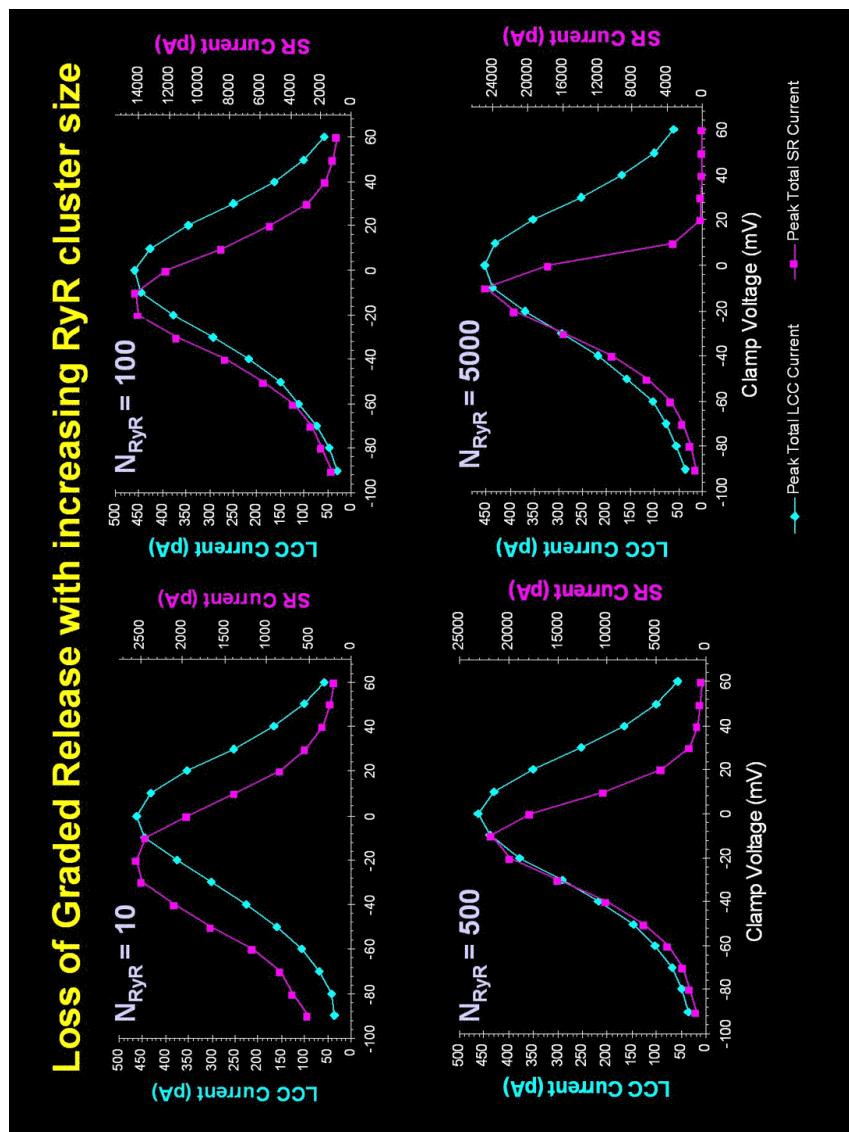
Whole-cell data is computed after summing up currents from each individual dyadic junction:

Peak LCC current Peak SR release

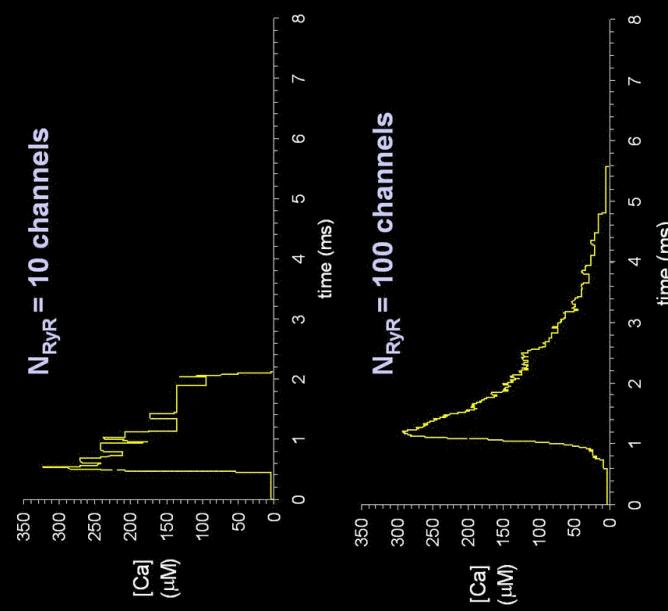


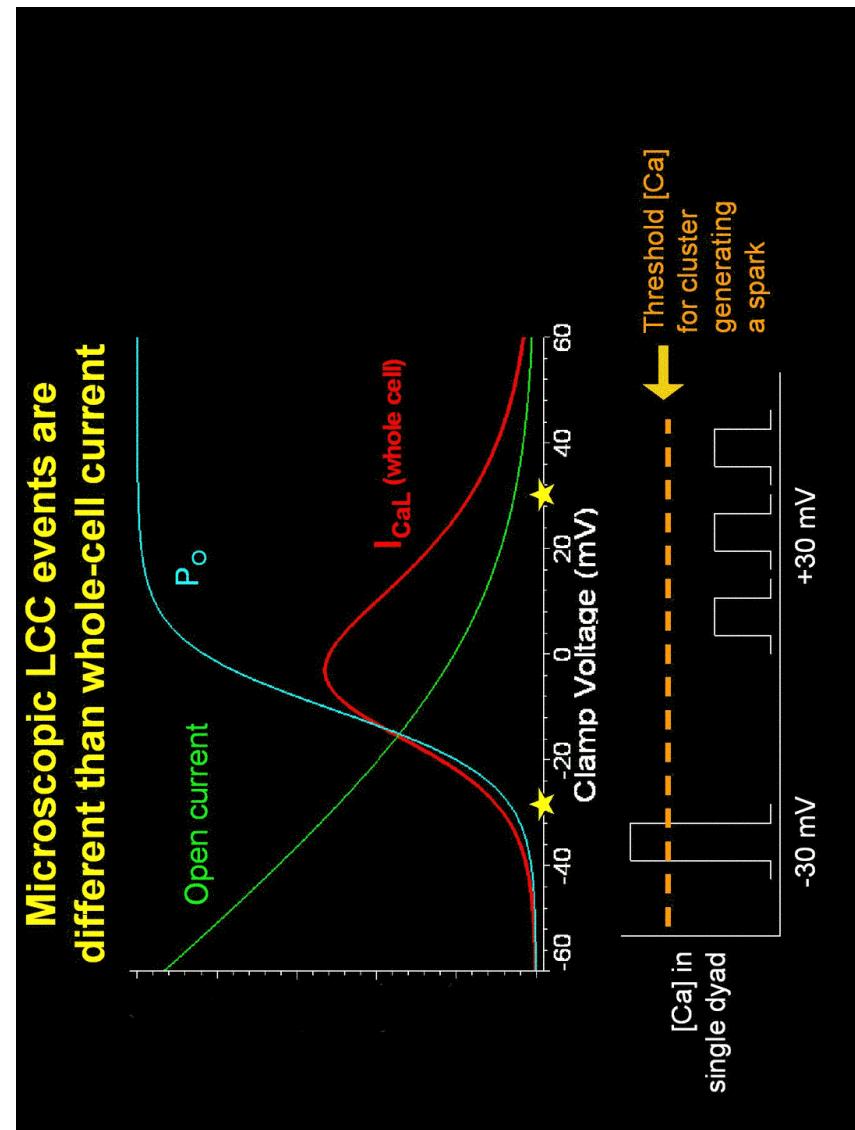
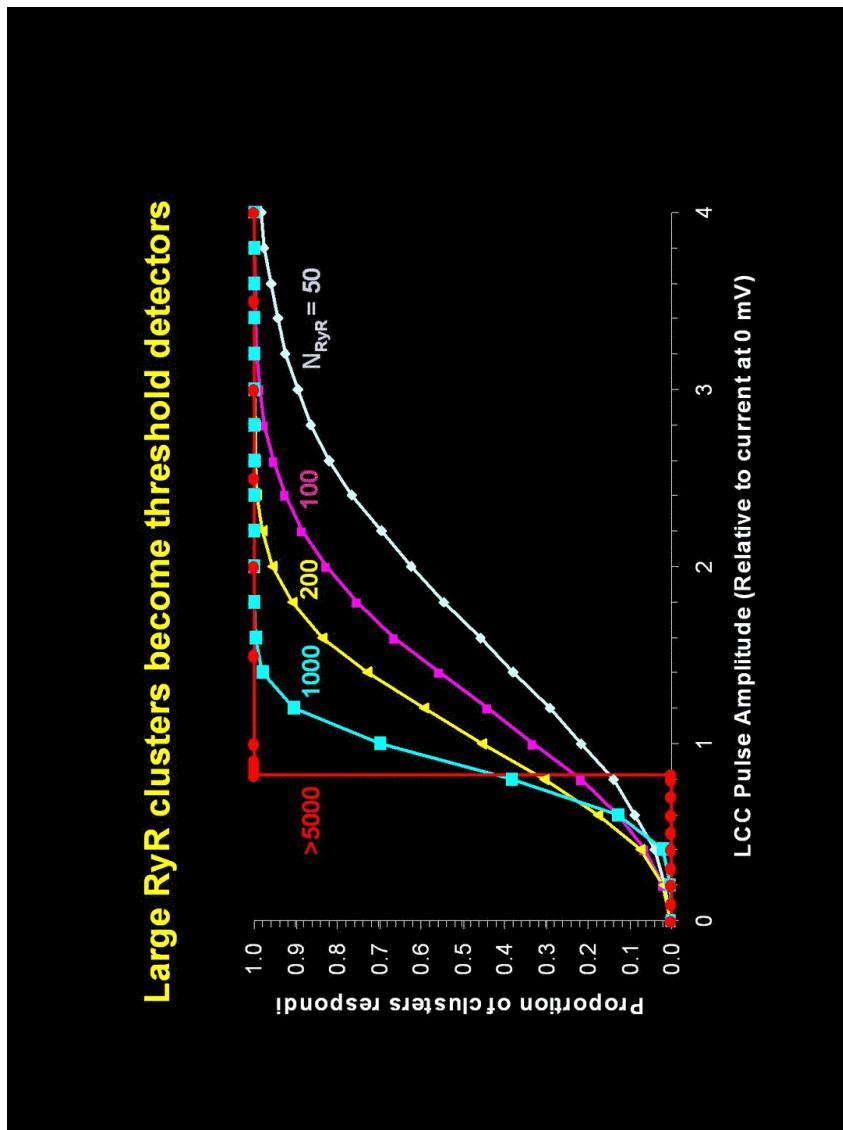
Peak currents are computed for a range of test voltages

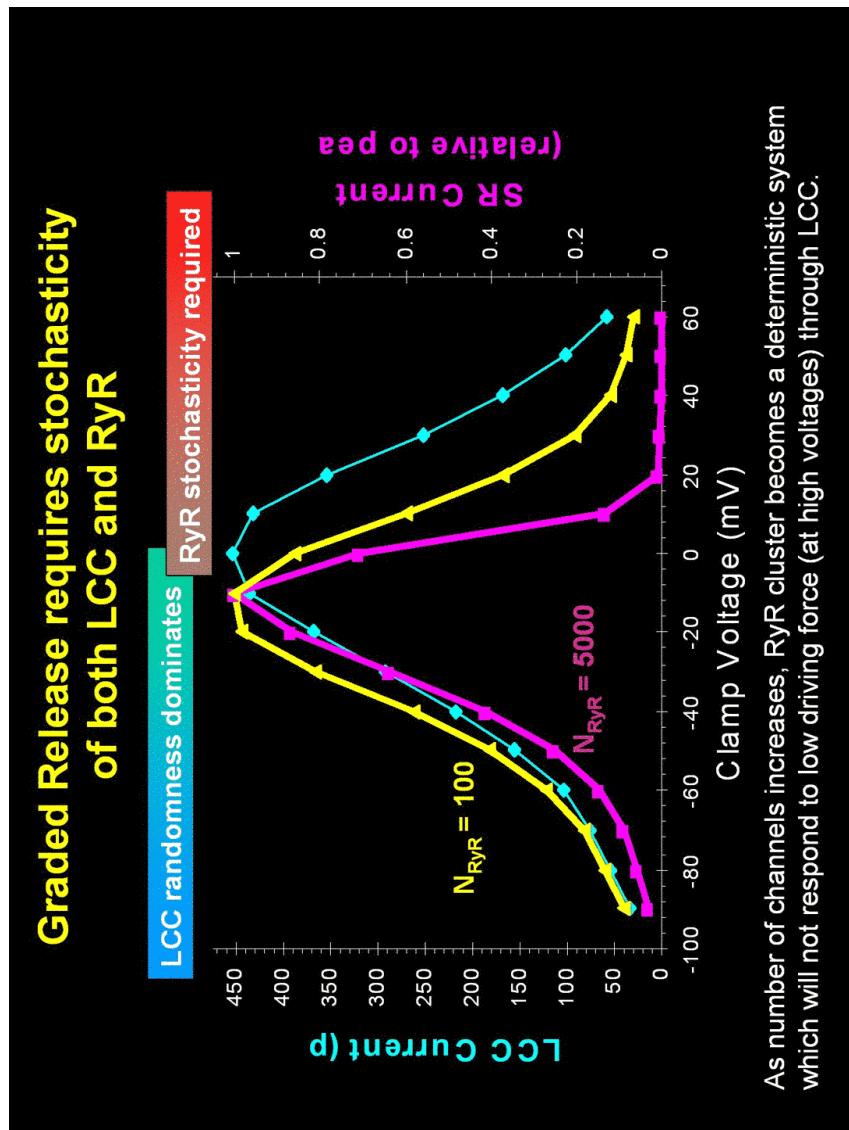
Each dyad contains a cluster of 100 RyR channels and 1 L-type Ca Channel



**Small RyR clusters do not reproduce
Ca sparks in junctions**







Main points

1. Rich nonlinear dynamics of voltage-calcium of cardiac cells. Dynamics can be understood in terms of basic features of Ca release, and bi-directional coupling.
2. Whole cell properties dependent on local ion channel signaling in the dyadic cleft.
3. Stochasticity at the ion channel level is critical to understanding whole cell properties.