

Investigations of rapidly rotating, stably stratified and non-hydrostatic flows

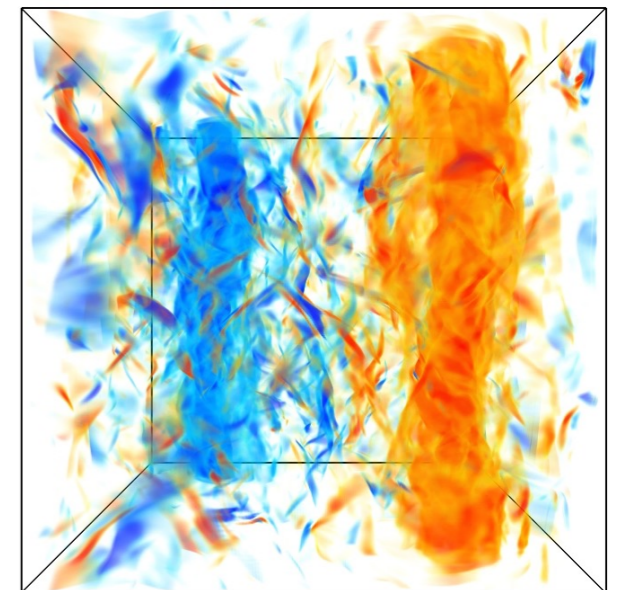
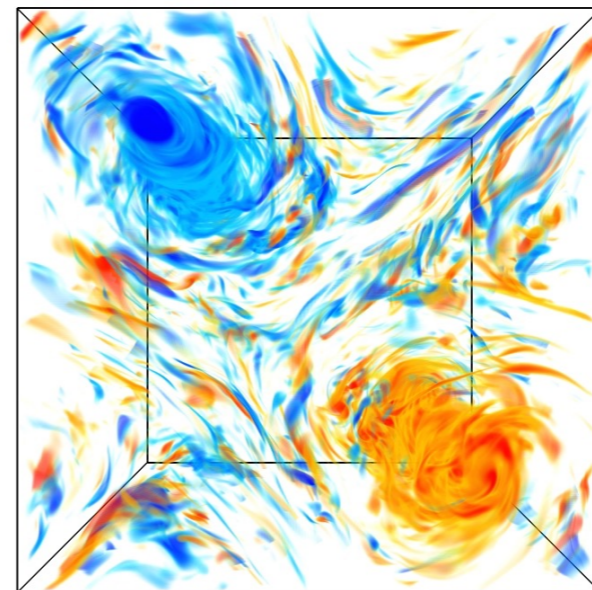
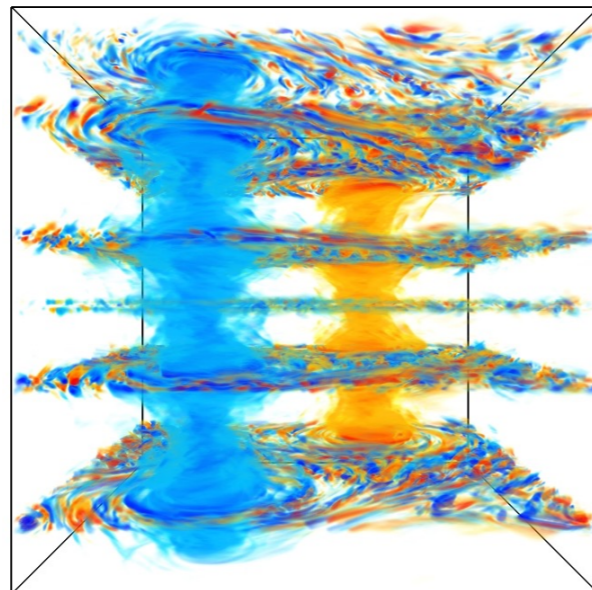
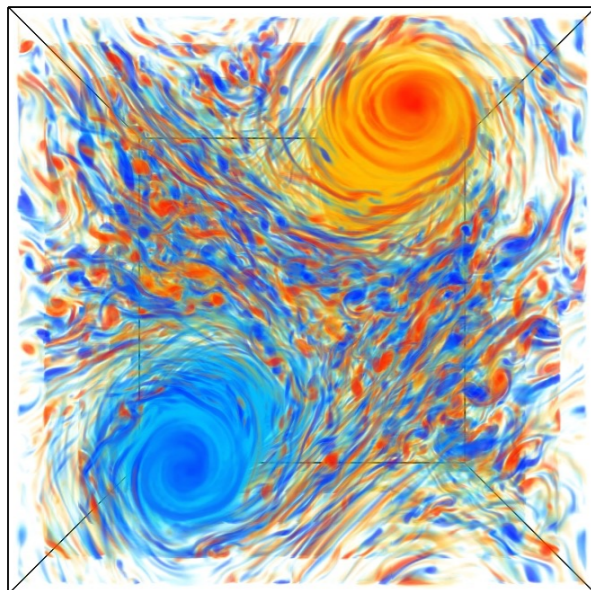
Keith Julien¹

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University of Colorado Boulder
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NSF EAR CSEDI-1067944

Nieves, Grooms, J. & Weiss (2016) *Journal of Fluid Mechanics*, vol. 801



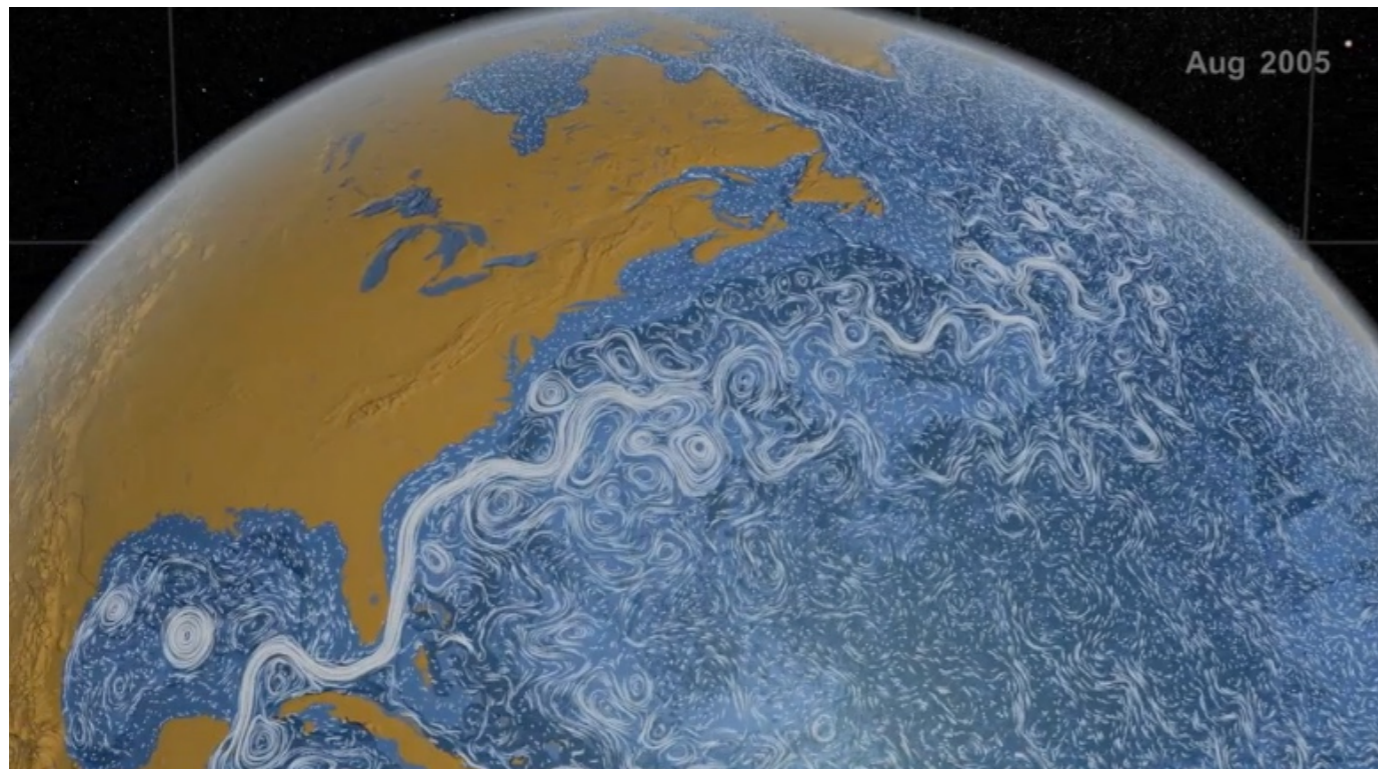
Motivation:

- **Strong system rotation** & **stable density stratification** characterize many geophysical flows

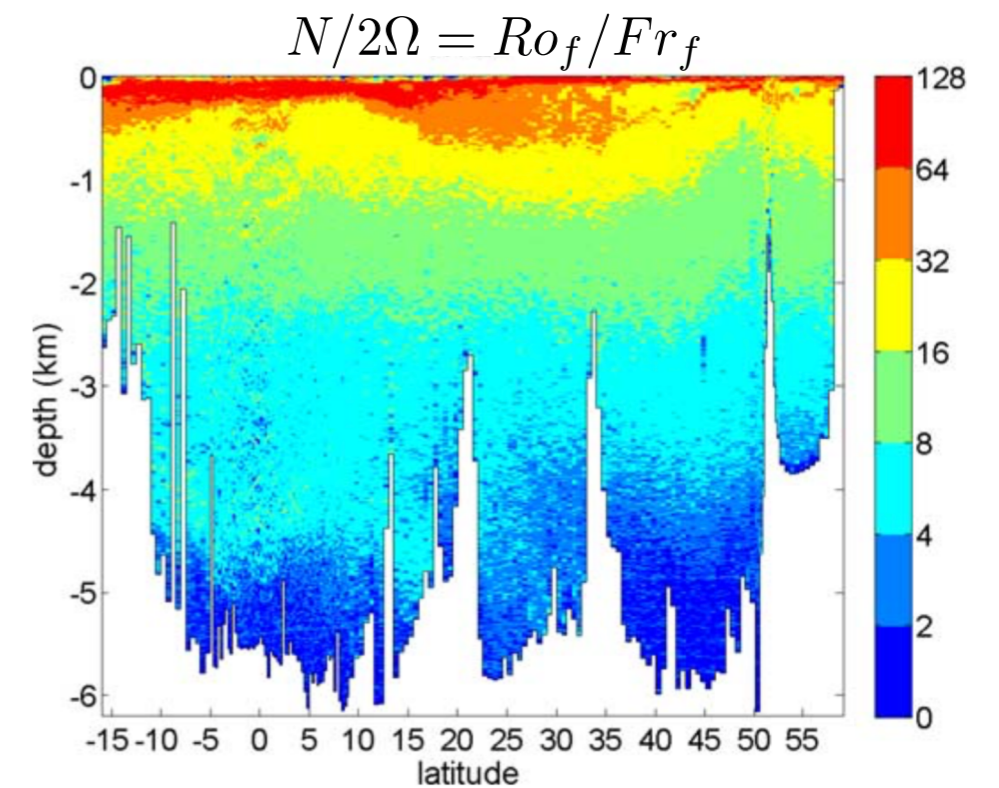
- Primary balance: **geostrophy** where $Ro_f \ll 1$
- Primary balance: **hydrostacy** where $Fr_f \ll 1$
- Strong spatial anisotropy $L_f \gg H$

$$Ro_f = \frac{U_f^*}{2\Omega L_f^*} \quad \text{Rossby}$$
$$Fr_f = \frac{U_f^*}{N_0^* L_f^*} \quad \text{Froude}$$
$$Re_f = \frac{U_f^* L_f^*}{\nu} \quad \text{Reynolds}$$

$$Fr_f \ll Ro_f \ll 1$$



Source: NASA (GCM synthesizing satellite and in-situ data)



North-South section of Pacific ocean at 179° E
Gerkema et al. Rev of Geophysics 2004

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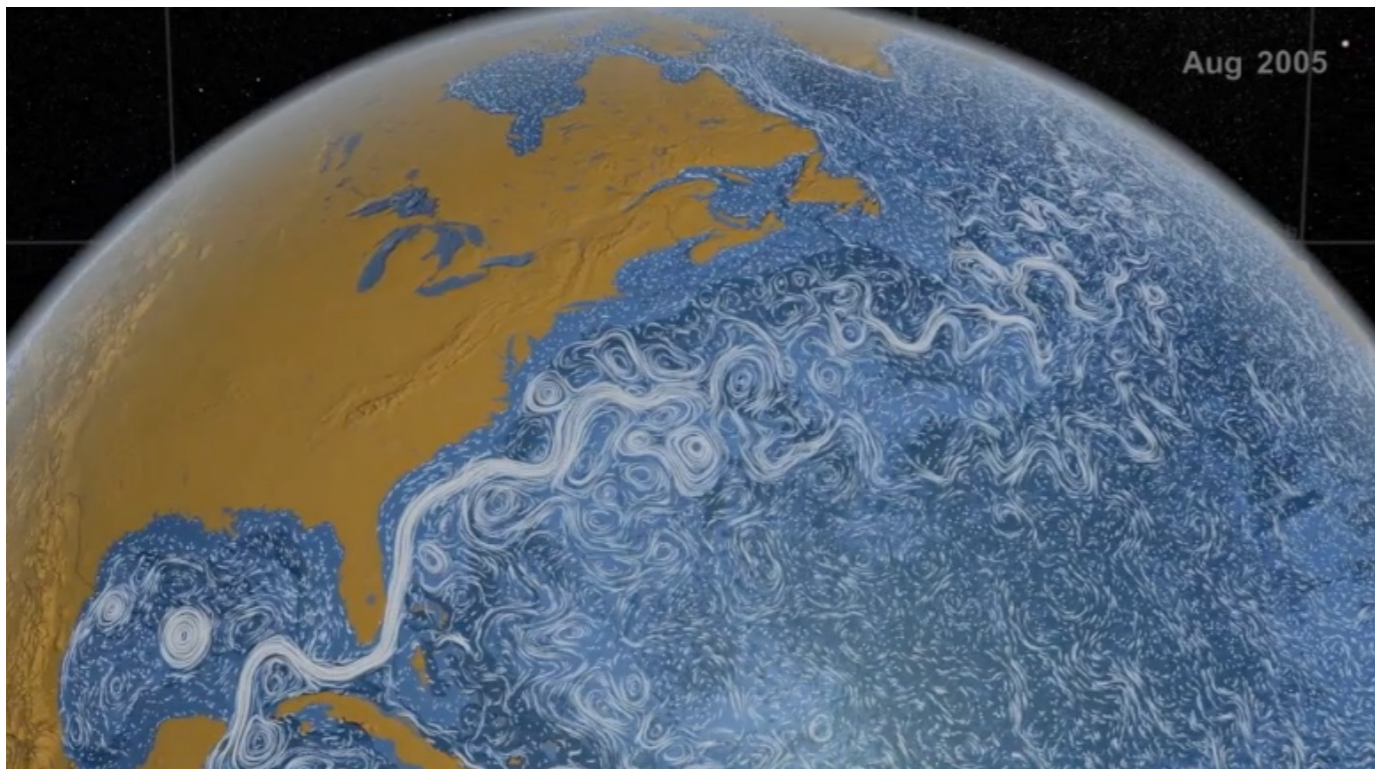
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Consider normal modes:
temporal separation btw wave & eddies

$$\propto \exp[i(\omega t + \mathbf{k}_\perp \cdot \mathbf{x}_\perp + k_z z)]$$

$$\omega_{\text{wave}}^2 = \frac{1}{Fr_f^2} \sin^2 \theta_{\mathbf{k}} + \frac{1}{Ro_f^2} \cos^2 \theta_{\mathbf{k}},$$

$$\omega_{\text{eddy}}^2 = 0$$

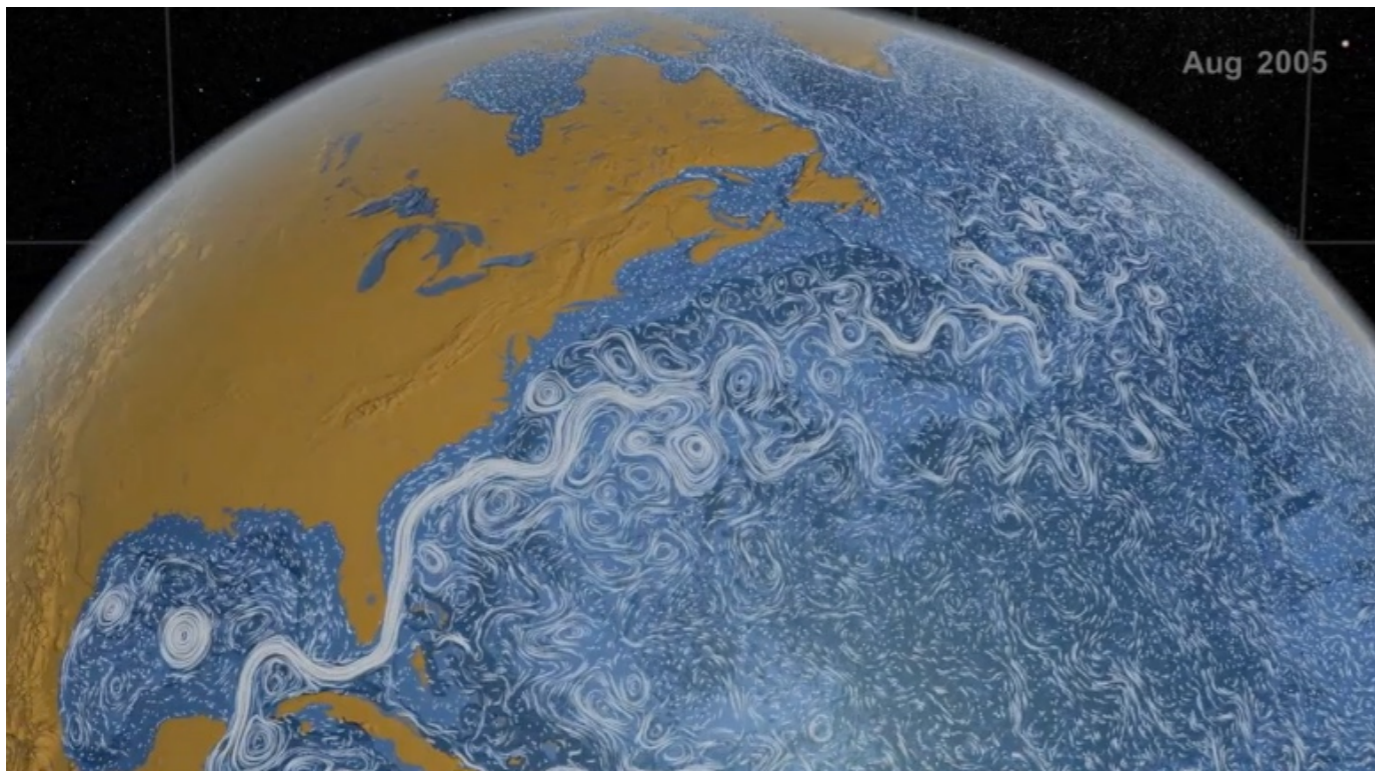
$$\tan \theta_{\mathbf{k}} = \frac{k_\perp}{k_z}$$

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$$\omega_{\text{wave}} \sim \mathcal{O} \left(\min \left\{ \frac{1}{Fr_f}, \frac{1}{Ro_f} \right\} \right) \gg \omega_{\text{eddy}}$$

Motivation: Hydrostatic-QG

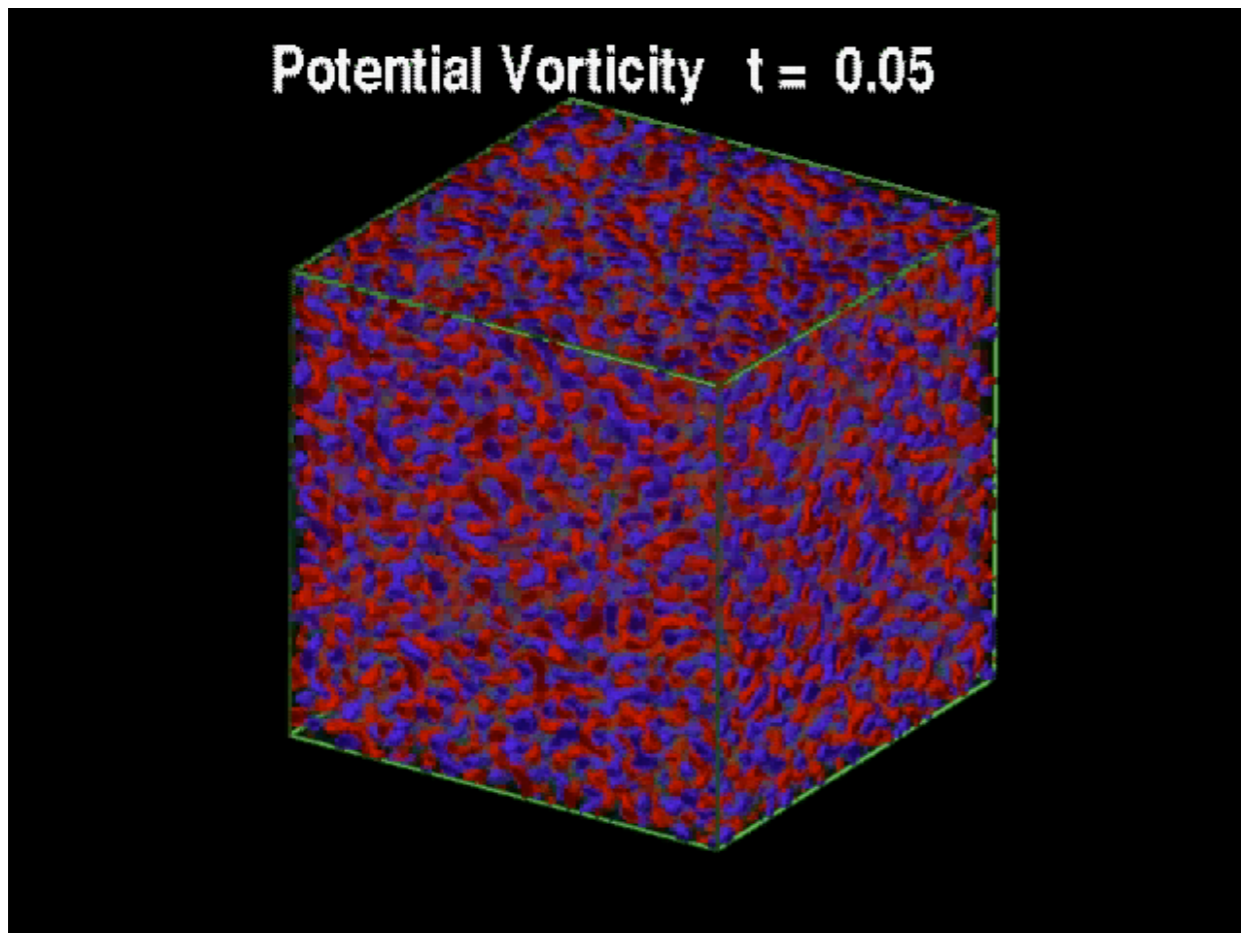
• **Strong system rotation** & **stable density stratification** characterize many geophysical flows

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- Result classical QG balance **filtering fast wave dynamics**

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McWilliams, Weiss & Yavneh JFM 1999

Classical QG (Charney '48,71): **reduction to PV**

$$p = \Psi, \quad \mathbf{u}_\perp = \nabla_\perp \times \Psi \hat{\mathbf{z}}, \quad b = \partial_z \Psi \quad w = \mathcal{O}\left(\frac{H}{L_f} Ro_f\right)$$

$$D_t^\perp \left(\nabla_\perp^2 \Psi + \frac{4\Omega^2}{N_0^{*2}} \partial_{zz} \Psi \right) = f + d$$

Inverse energy cascade $E(k) \approx k^{-5/3}, k < k_f$

Direct enstrophy cascade $\approx k^{-3}, k > k_f$

$$\omega_{\text{wave}} \sim \mathcal{O}\left(\min\left\{\frac{1}{Fr_f}, \frac{1}{Ro_f}\right\}\right) \gg \omega_{\text{eddy}}$$

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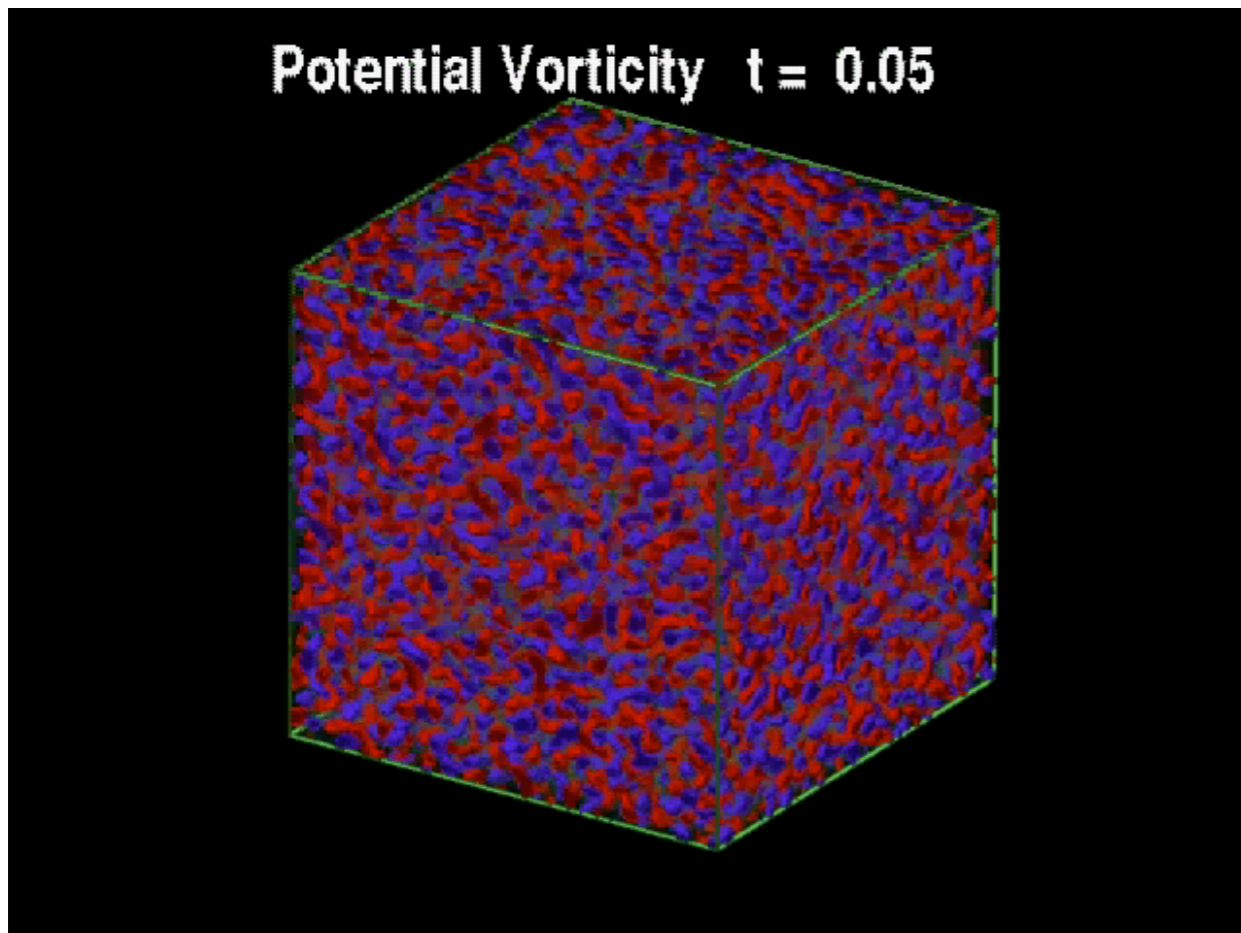
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Validity of QG - Embid & Majda GAFD 1998
 - Temam & Wirosoetisno JAS '10,'11

Evidenced in DNS - Smith & Waleffe JFM '02

Marino et al EPL '13

Whitehead & Wingate JFM '14

Motivation:

- **Strong system rotation** & **weak stable density stratification** also prevalent in geophysical flows

Current focus

- Primary balance: geostrophy where $Ro_f < 1$
- weak stable stratification where $Fr_f = \mathcal{O}(1)$

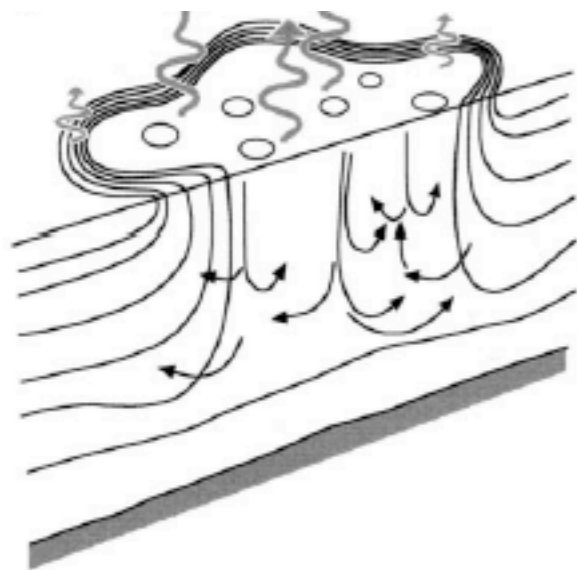
Examples

- Open ocean deep convection, convective overshoot
- High latitude abyssal oceans (not well studied)

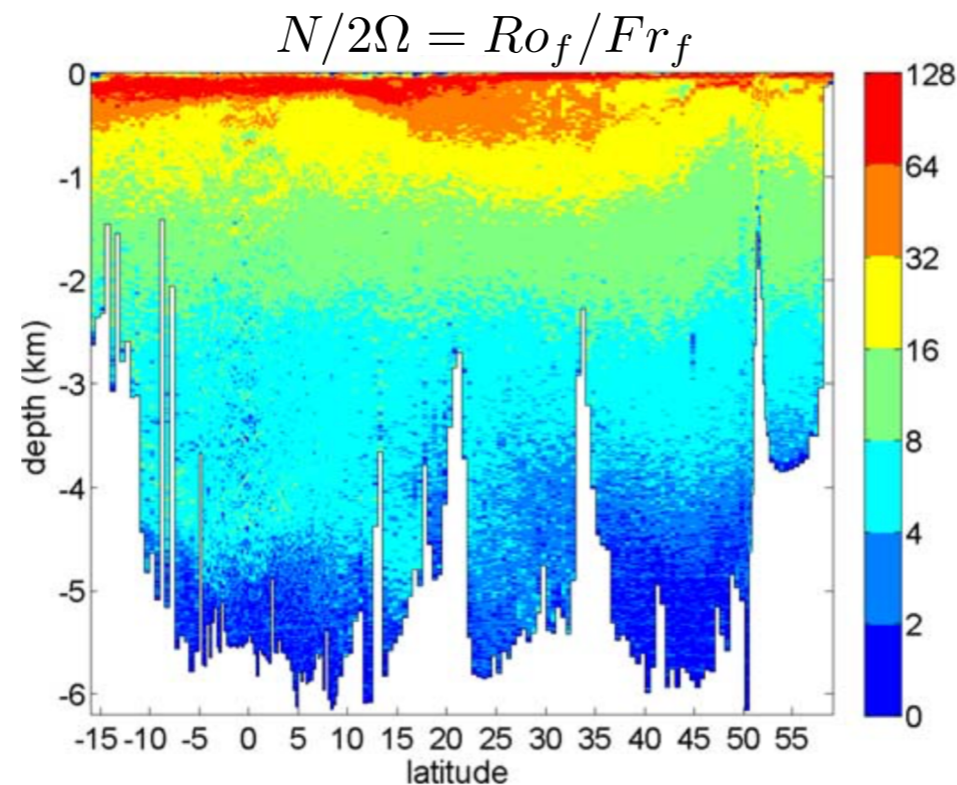
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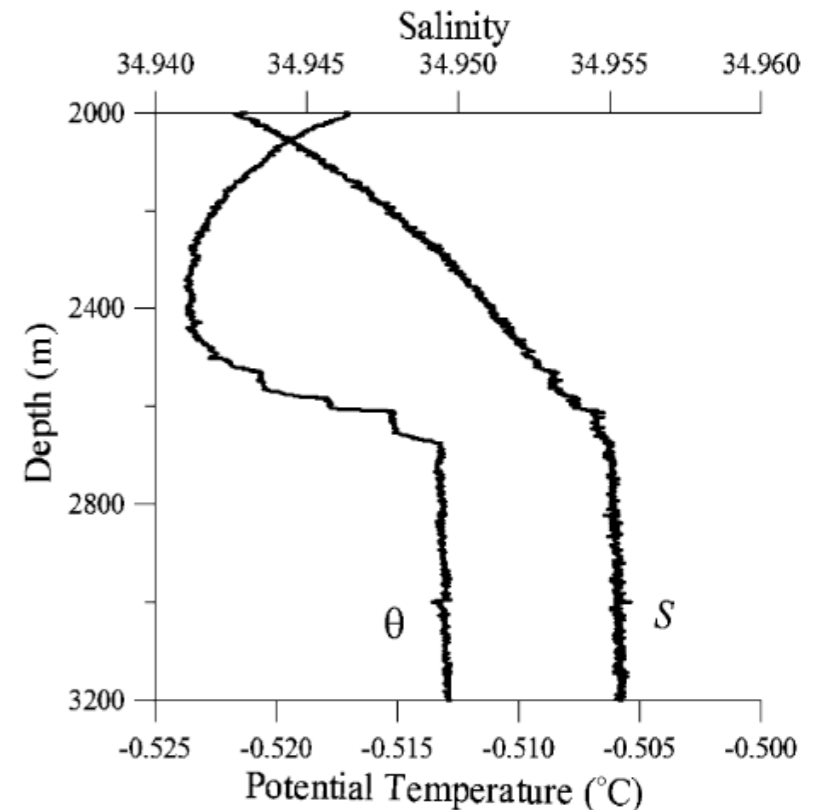
$$Re_f = \frac{U_f^* L_f^*}{\nu} \quad \text{Reynolds}$$



Marshall & Schott RG 1999



North-South section of Pacific ocean at 179° E
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Timmermans et al. DSR 2003

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- Primary balance: geostrophy where $Ro_f < 1$
- **weak stable stratification** where $Fr_f = \mathcal{O}(1)$

Examples

- Open ocean deep convection, convective overshoot
- High latitude abyssal oceans (not well studied)
- **Result: no separation between eddy and wave dynamics**

$$\omega_{\text{wave}}^2 = \frac{1}{Fr_f^2} \sin^2 \theta_{\mathbf{k}} + \frac{1}{Ro_f^2} \cos^2 \theta_{\mathbf{k}},$$

$$\omega_{\text{eddy}}^2 = 0$$

$$\tan \theta_{\mathbf{k}} = \frac{k_{\perp}}{k_z}$$

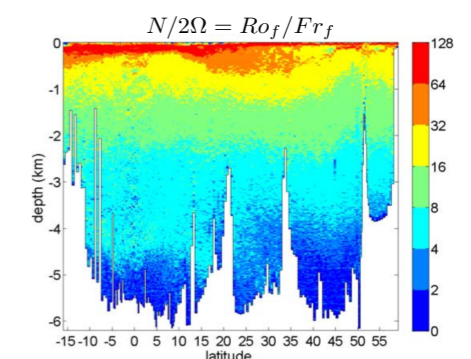
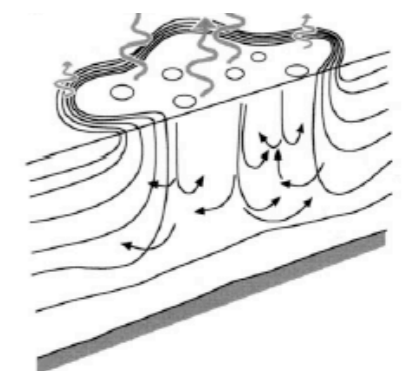
$$A = \frac{k_{\perp}}{k_z} = \mathcal{O}(Ro_f^{-1}) \implies \theta = \pm \left(\frac{\pi}{2} - \mathcal{O}(Ro_f) \right)$$

Considering anisotropy: columnar dynamics is slow

$$Ro_f = \frac{U_f^*}{2\Omega L_f^*} \quad \text{Rossby}$$

$$Fr_f = \frac{U_f^*}{N_0^* L_f^*} \quad \text{Froude}$$

$$Re_f = \frac{U_f^* L_f^*}{\nu} \quad \text{Reynolds}$$



Governing equations (*non-dimensional* Boussinesq equations):

$$D_t \mathbf{u} + \left(\frac{1}{Ro_f} \hat{\mathbf{z}} \times \mathbf{u} = -Eu_f \nabla p \right) + b \hat{\mathbf{z}} + \frac{1}{Re_f} \nabla^2 \mathbf{u}$$

imposed system rotation

$$D_t b + \frac{1}{Fr_f^2} S(z) w = \frac{1}{\sigma Re_f} \nabla^2 b$$

$$\boldsymbol{\Omega} = \Omega \hat{\mathbf{z}}$$

imposed ambient stratification

$$\nabla \cdot \mathbf{u} = 0$$

non-dimensional numbers

$$N_0^* \equiv |g \rho_0^{*-1} (\partial_{z^*} \delta \hat{\rho}^*(z^*))_{max}|$$

$$S(z) \equiv -\partial_z \delta \hat{\rho}$$

$$Ro_f = \frac{U_f^*}{2\Omega L_f^*} \quad \text{Rossby number}$$

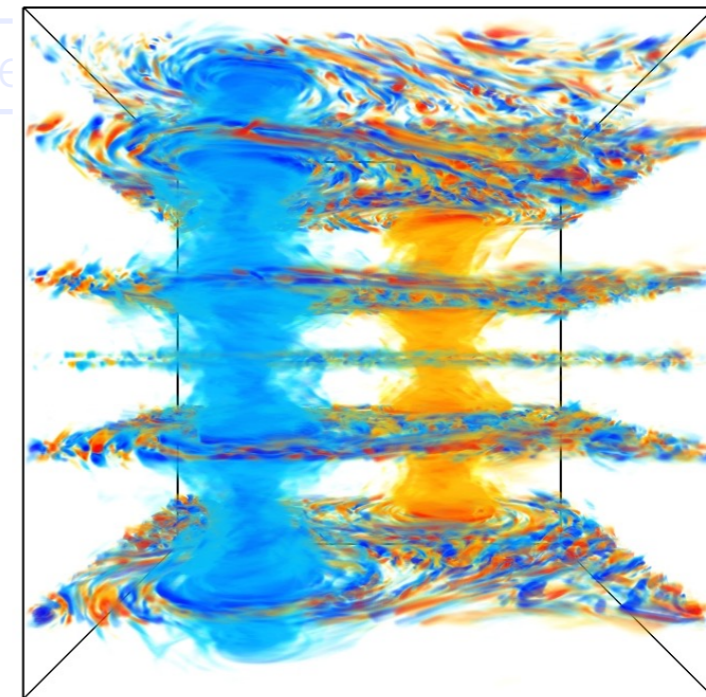
$$Fr_f = \frac{U_f^*}{N_0^* L_f^*} \quad \text{Froude number}$$

$$Eu_f = \frac{\delta p_0}{\rho_0 U_f^{*2}} \quad \text{Euler number}$$

$$Re_f = \frac{U_f^* L_f^*}{\nu} \quad \text{Reynolds number}$$

$$\sigma = \frac{\nu}{\kappa} \quad \text{Prandtl number}$$

energy injection



motions

3D periodic rotating box stratification - weak to strong

Governing equations (*non-dimensional* Boussinesq equations):

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imposed system rotation

$$D_t b + \frac{1}{Fr_f^2} S(z) w = \frac{1}{\sigma Re_f} \nabla^2 b$$

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non-dimensional numbers

$$N_0^* \equiv |g \rho_0^{*-1} (\partial_{z^*} \delta \hat{\rho}^*(z^*))_{max}|$$

$$S(z) \equiv -\partial z \delta \hat{\rho}$$

energy injection: required to excite fluid motions
vertical motions only

$$Ro_f = \frac{U_f^*}{2\Omega L_f^*} \quad \text{Rossby number}$$

$$Fr_f = \frac{U_f^*}{N_0^* L_f^*} \quad \text{Froude number}$$

$$Eu_f = \frac{\delta p_0}{\rho_0 U_f^{*2}} \quad \text{Euler number}$$

ϵ_f^* injection rate

L_f^* injection scale

characteristic scales

$$U_f^* = (\epsilon_f^* L_f^*)^{1/3}$$

$$T_f^* = (L_f^{*2} \epsilon_f^{*-1})^{1/3}$$

$$Re_f = \frac{U_f^* L_f^*}{\nu} \quad \text{Reynolds number}$$

$$\sigma = \frac{\nu}{\kappa} \quad \text{Prandtl number}$$

Internal length scales (partitioning of parameter space)

$L_K \equiv Re_f^{-3/4}$ Kolmogorov scale - equivalence of viscous and eddy turnover times

$L_D \equiv \frac{ARo_f}{Fr_f}$ Deformation radius - PE to KE conversion by baroclinic instability

$L_\Omega \equiv Ro_f^{3/2}$ Zeman scale* - Coriolis (buoyancy) force negligible

$L_N \equiv Fr_f^{3/2}$ Ozmidov scale

aspect ratio

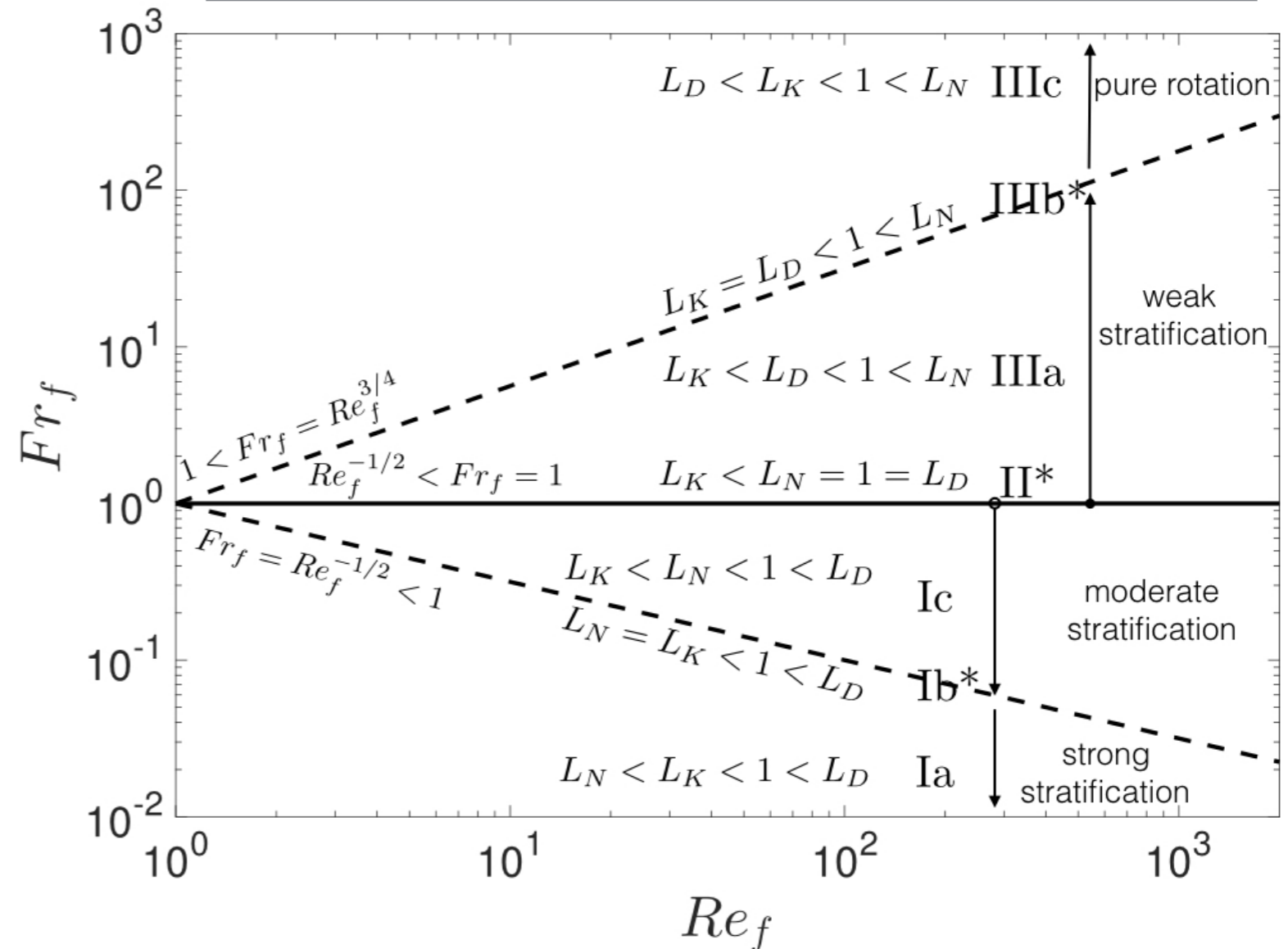
$$A = \frac{H^*}{L_f^*} \sim Ro_f^{-1}$$

We focus on flows with

$$Ro_f = o\left(Re_f^{-1/2}\right)$$

$$L_\Omega = o(L_K)$$

partitioning of parameter space for $Ro_f \ll 1$



Governing equations (*non-dimensional* Boussinesq equations):

$$D_t \mathbf{u} + \left(\frac{1}{Ro_f} \hat{\mathbf{z}} \times \mathbf{u} - Eu_f \nabla p \right) + b \hat{\mathbf{z}} + \frac{1}{Re_f} \nabla^2 \mathbf{u}$$

$$D_t b + \frac{1}{Fr_f^2} S(z) w = \frac{1}{\sigma Re_f} \nabla^2 b$$

$$\nabla \cdot \mathbf{u} = 0$$



$$A = Ro_f^{-1} \gg 1$$

$$w = \mathcal{O}(ARo_f) \sim 1$$

leading order properties:

$$\hat{\mathbf{z}} \times \mathbf{u}_0 + \nabla_{\perp} p'_0 = 0 \quad \mathbf{u}'_0 = -\nabla_{\perp} \times \psi'_0 \hat{\mathbf{z}} + w'_0 \hat{\mathbf{z}} \quad (\text{geostrophic solution})$$

$$\nabla_{\perp} \cdot \mathbf{u}_0 = 0 \quad p'_0 = \psi'_0 \quad (\text{geostrophic streamfunction})$$

Non-hydrostatic quasi-geostrophic (NH-QG) equations



$$\partial_t \zeta_0 + J[\psi_0, \zeta_0] - \partial_Z w_0 = \frac{1}{Re_f} \nabla_{\perp}^2 \zeta_0 \quad \text{vertical vorticity}$$

$$\partial_t w_0 + J[\psi_0, w_0] + \partial_Z \psi_0 = b_0 + \frac{1}{Re_f} \nabla_{\perp}^2 w_0 + f_{w_0} \quad \text{vertical velocity}$$

$$A = Ro_f^{-1} \gg 1 \\ w = \mathcal{O}(ARo_f) \sim 1$$

$$\left. \begin{aligned} \partial_t b_0 + J[\psi_0, b_0] + w_0 \left(\partial_Z \bar{b}_{-1} + \frac{1}{Fr_f^2} S(Z) \right) &= \frac{1}{\sigma Re_f} \nabla_{\perp}^2 b_0 \\ \partial_T \bar{b}_{-1} + \partial_Z (\overline{w_0 b_0}) &= \frac{1}{\sigma Re_f} \partial_Z^2 \bar{b}_{-1} \end{aligned} \right\} \text{buoyancy}$$

leading order properties:

$$\begin{aligned} \hat{\mathbf{z}} \times \mathbf{u}_0 + \nabla_{\perp} p'_0 &= 0 & \mathbf{u}'_0 &= -\nabla_{\perp} \times \psi'_0 \hat{\mathbf{z}} + w'_0 \hat{\mathbf{z}} \quad (\text{geostrophic solution}) \\ \nabla_{\perp} \cdot \mathbf{u}_0 &= 0 & p'_0 &= \psi'_0 \quad (\text{geostrophic streamfunction}) \end{aligned}$$

$$J[\psi_0, \cdot] \equiv \mathbf{u}'_0 \cdot \nabla_{\perp} \quad (\text{vertical advection is a subdominant effect})$$

Linearization for inviscid NH-QG equations about $S(Z) = 1$ gives

$$\omega_{\text{wave}}^2 = \frac{1}{Fr_f^2} + \frac{k_Z^2}{k_{\perp}^2} \quad k_Z = \mathcal{O}(1) \quad \begin{aligned} &(\text{wave-eddy interactions occur on order-one timescale}) \\ &(\text{inviscid conserved quantity - energy not enstrophy}) \end{aligned}$$

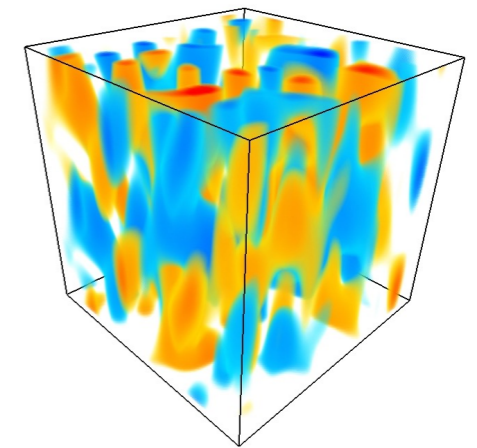
Numerical simulation of NH-QG equations

$$\partial_t \zeta_0 + J[\psi_0, \zeta_0] - \partial_Z w_0 = \frac{1}{Re_f} \nabla_{\perp}^2 \zeta_0$$

$$\partial_t w_0 + J[\psi_0, w_0] + \partial_Z \psi_0 = b_0 + \frac{1}{Re_f} \nabla_{\perp}^2 w_0 + \boxed{f_{w_0}} \left. \vphantom{\frac{1}{Re_f} \nabla_{\perp}^2 w_0} \right\} \text{stochastic forcing of vertical velocity used to excite fluid motions (waves)}$$

$$\partial_t b_0 + J[\psi_0, b_0] + w_0 \left(\partial_Z \bar{b}_{-1} + \frac{1}{Fr_f^2} S(Z) \right) = \frac{1}{\sigma Re_f} \nabla_{\perp}^2 b_0$$

$$\partial_T \bar{b}_{-1} + \partial_Z (\overline{w_0 b_0}) = \frac{1}{\sigma Re_f} \partial_Z^2 \bar{b}_{-1}$$



time-stepping:

- deterministic dynamics evolved using 3rd-order Runge-Kutta
- white-noise forcing is nowhere differentiable in time
stochastic dynamics evolved using Euler-Maruyama

$$u_{n+1}^* = u_{n+1} + \sqrt{\epsilon_f \Delta t} \chi_n$$

$$\mathbb{E} [u_{n+1}^{*2}] = \mathbb{E} [u_{n+1}^2] + \epsilon_f \Delta t$$

numerical resolution:

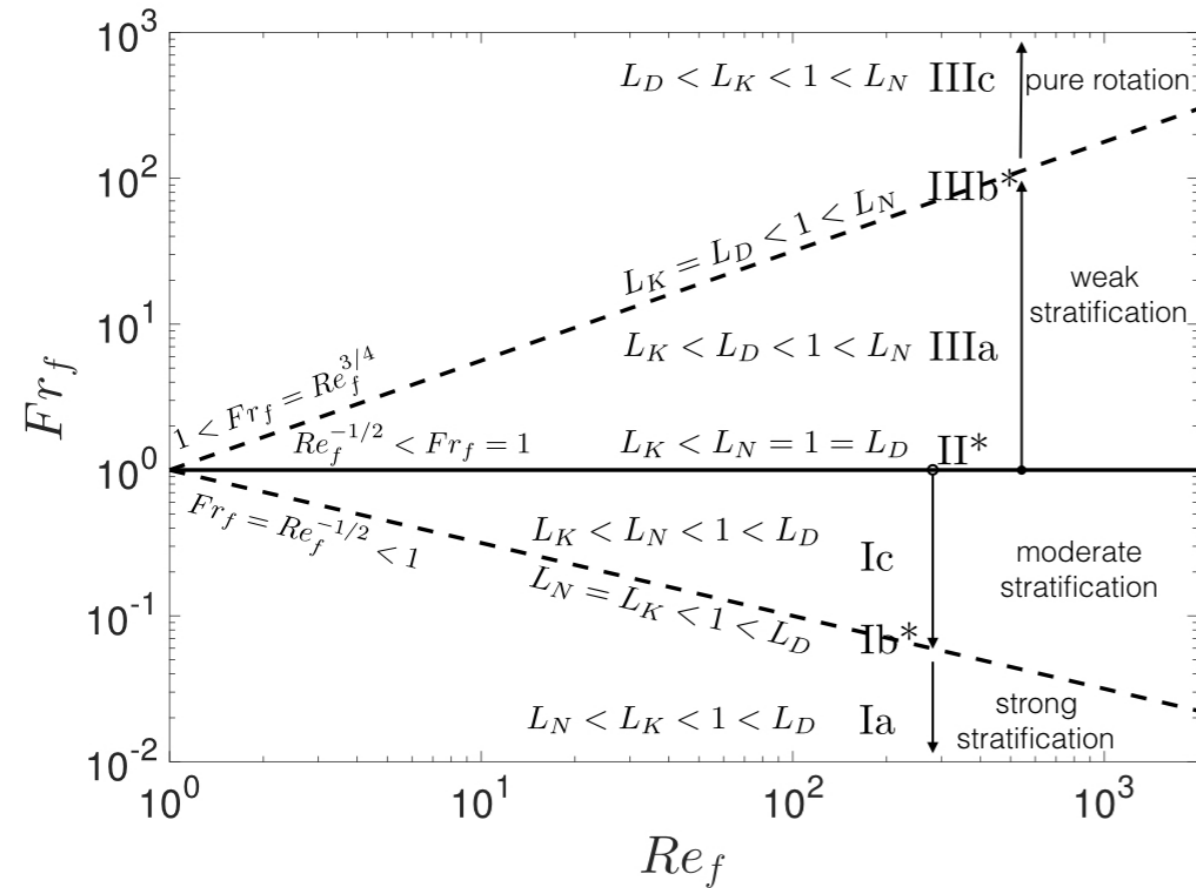
$$\Delta x = 2L_K, \quad L_K = Re^{-3/4}$$

$$N_{x,y,z} = \frac{L_b}{2} Re^{3/4}, \quad L_b = 10$$

Domain size: $L_b \times L_b \times 1$

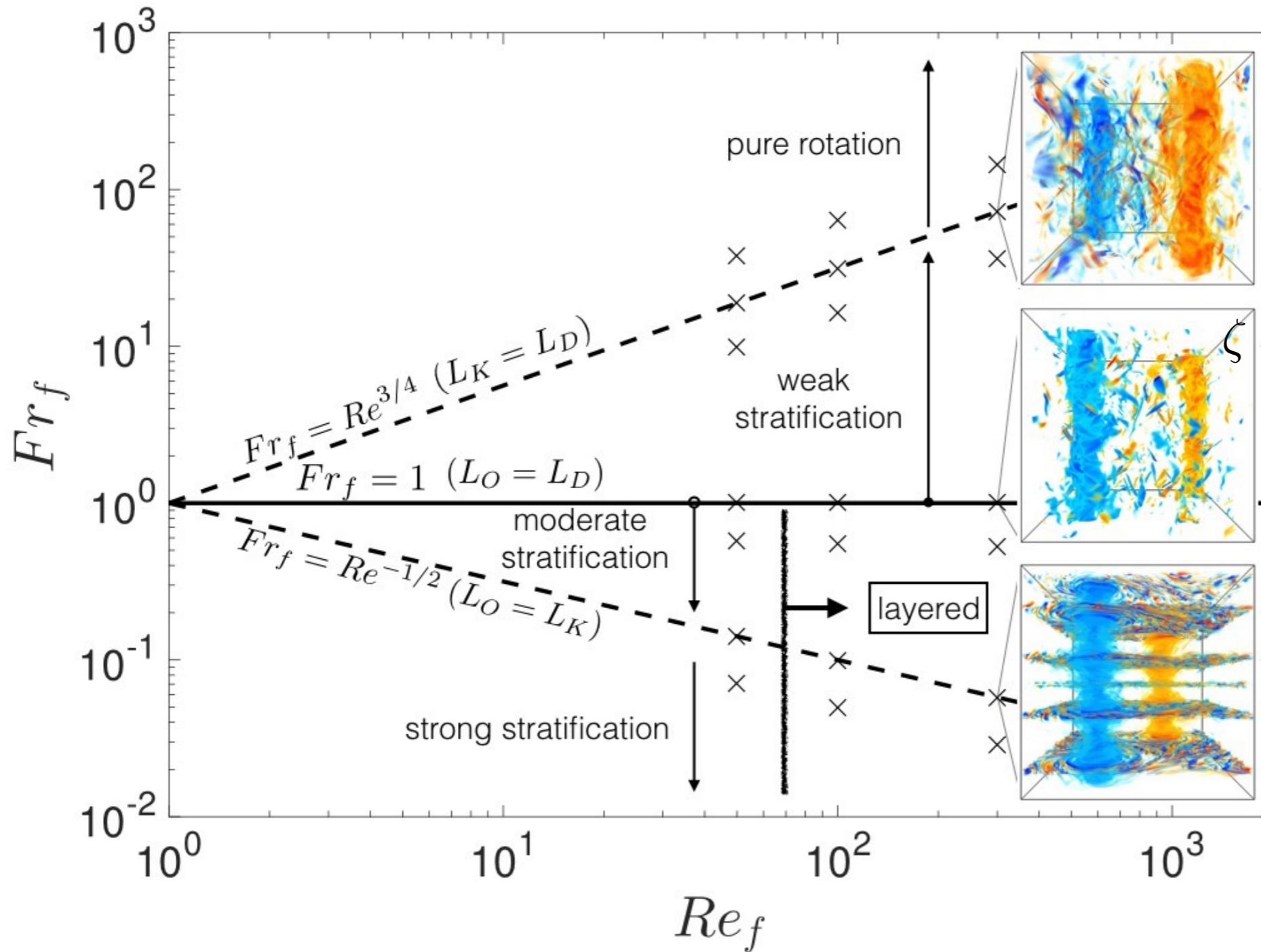
Numerical Survey

partitioning of parameter space for $Ro_f \ll Re_f^{1/2} < 1$



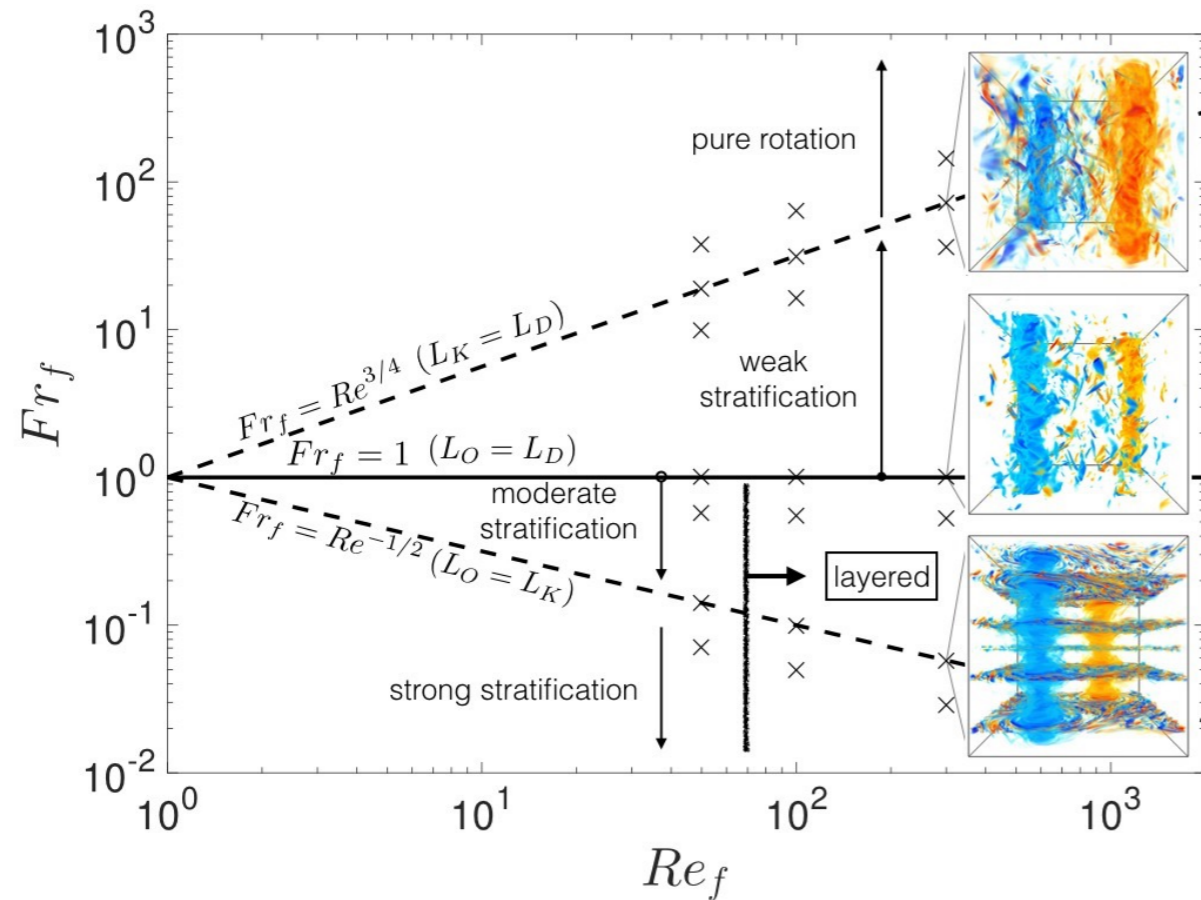
Regime	$Fr_f(Re_f)$	$Fr_f(Re_f = 50)$	$Fr_f(Re_f = 100)$	$Fr_f(Re_f = 300)$
Ia	$\frac{1}{2}Re_f^{-1/2}$	0.0707	0.0500	0.0289
Ib*	$Re_f^{-1/2}$	0.1414	0.1000	0.0577
Ic	$\frac{1}{2}(1 + Re_f^{-1/2})$	0.5707	0.5500	0.5289
II*	1	1	1	1
IIIa	$\frac{1}{2}(1 + Re_f^{3/4})$	9.9015	16.311	36.542
IIIb*	$Re_f^{3/4}$	18.803	31.623	72.084
IIIc	$2Re_f^{3/4}$	37.606	63.246	144.17
Grid resolution	$N_x \times N_y \times N_z$	$96 \times 96 \times 96$	$192 \times 192 \times 192$	$384 \times 384 \times 384$

Results (qualitative): large scale barotropic mode



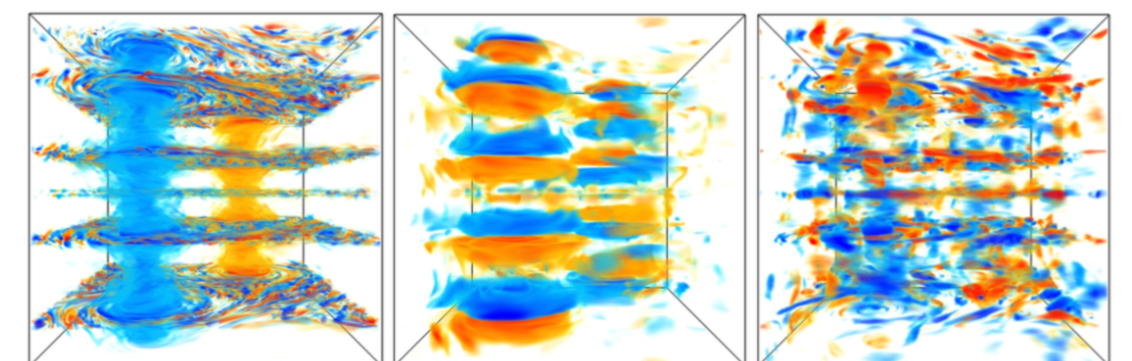
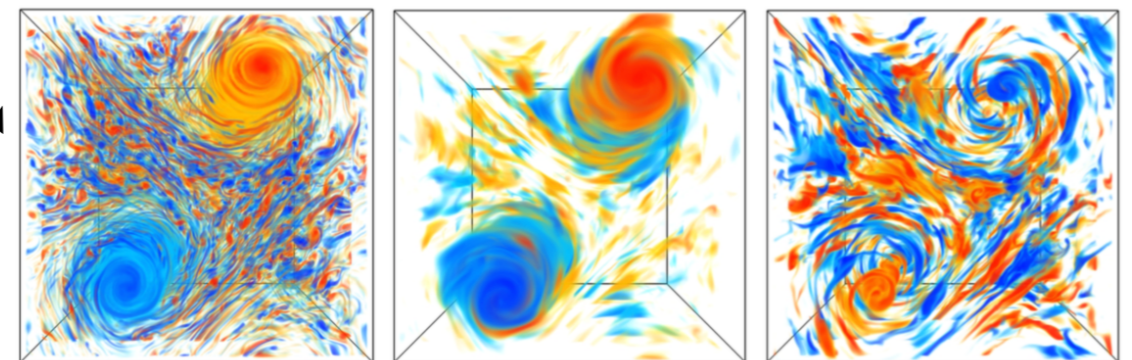
Flow organizes into a barotropic dipole surrounded by small scale geostrophic turbulence

Results (qualitative): large scale barotropic mode



- tendency to form well-defined layers
- turbulence significantly enhanced within layers
- provided Re sufficiently large
- layering not observed in classical QG (coupling to waves absent)

$$Re_f = 300, \quad Re_c = 4290 \quad Fr_f < 1$$

(a) Vertical vorticity, ζ (b) Buoyancy, b (c) Vertical velocity, w

Results: Layering

mixing leads to locally reduced stratification

$$\partial_Z \left(\overline{w_0 b_0} - \frac{1}{Pe} \partial_Z \bar{b}_{-1} \right) = 0$$

vorticity RMS profiles:

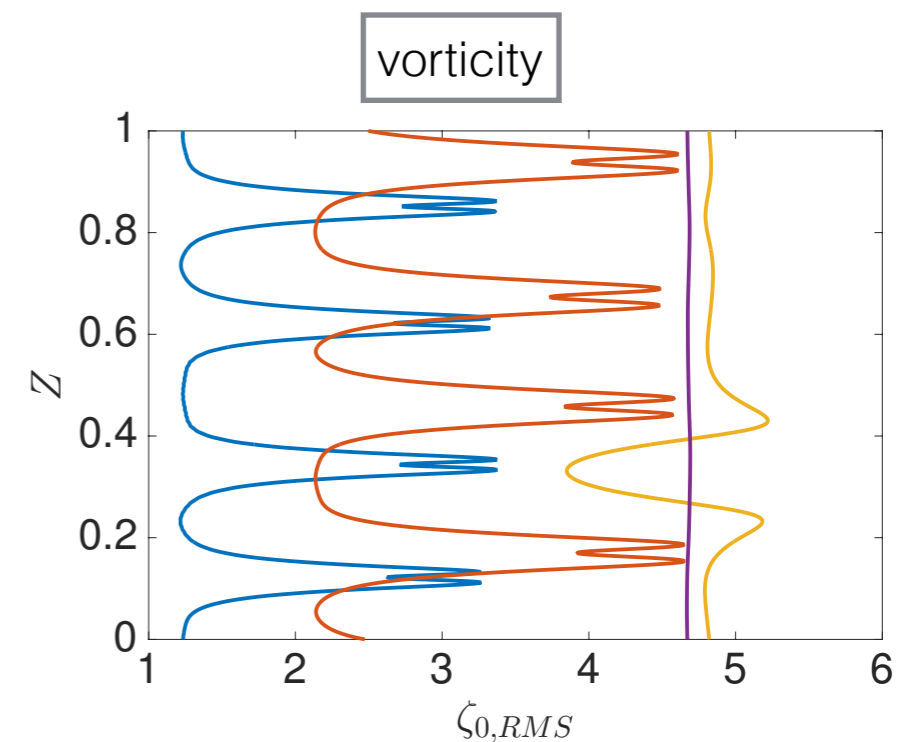
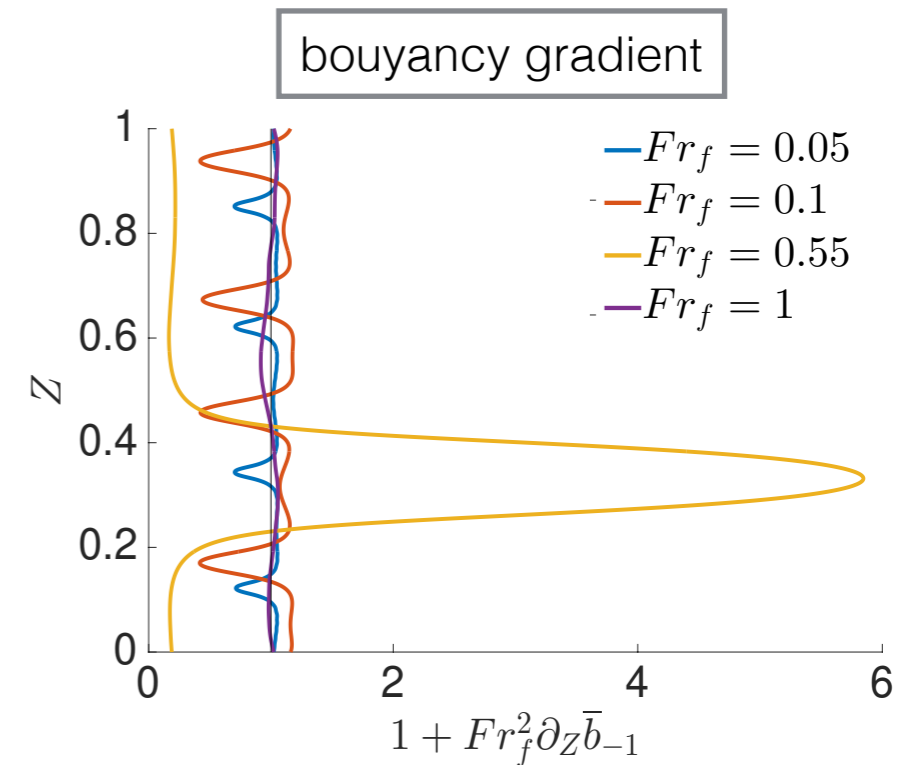
- layering: regions of enhanced vortex stretching
- evidence of sublayers (thickness: distance b/w maxima)
- locally increased kinetic dissipation rate $Re_f^{-1} \overline{\zeta^2}^A$

layer location:

- minima of mean-buoyancy gradient
- local minima in vorticity RMS profile

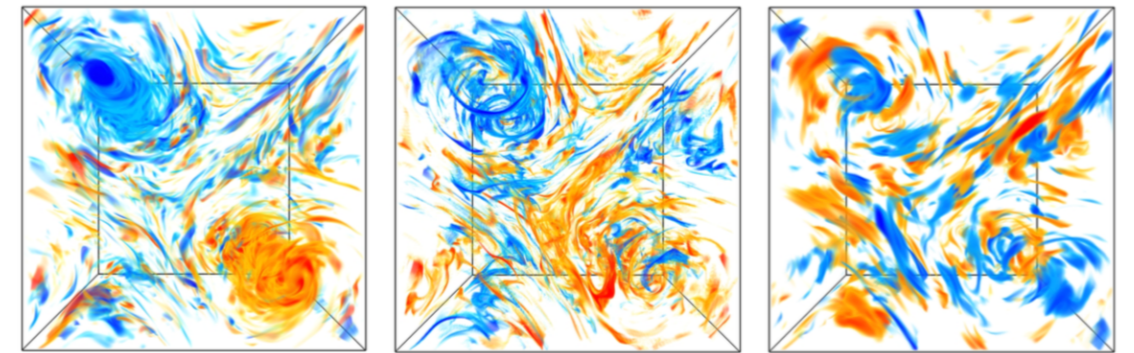
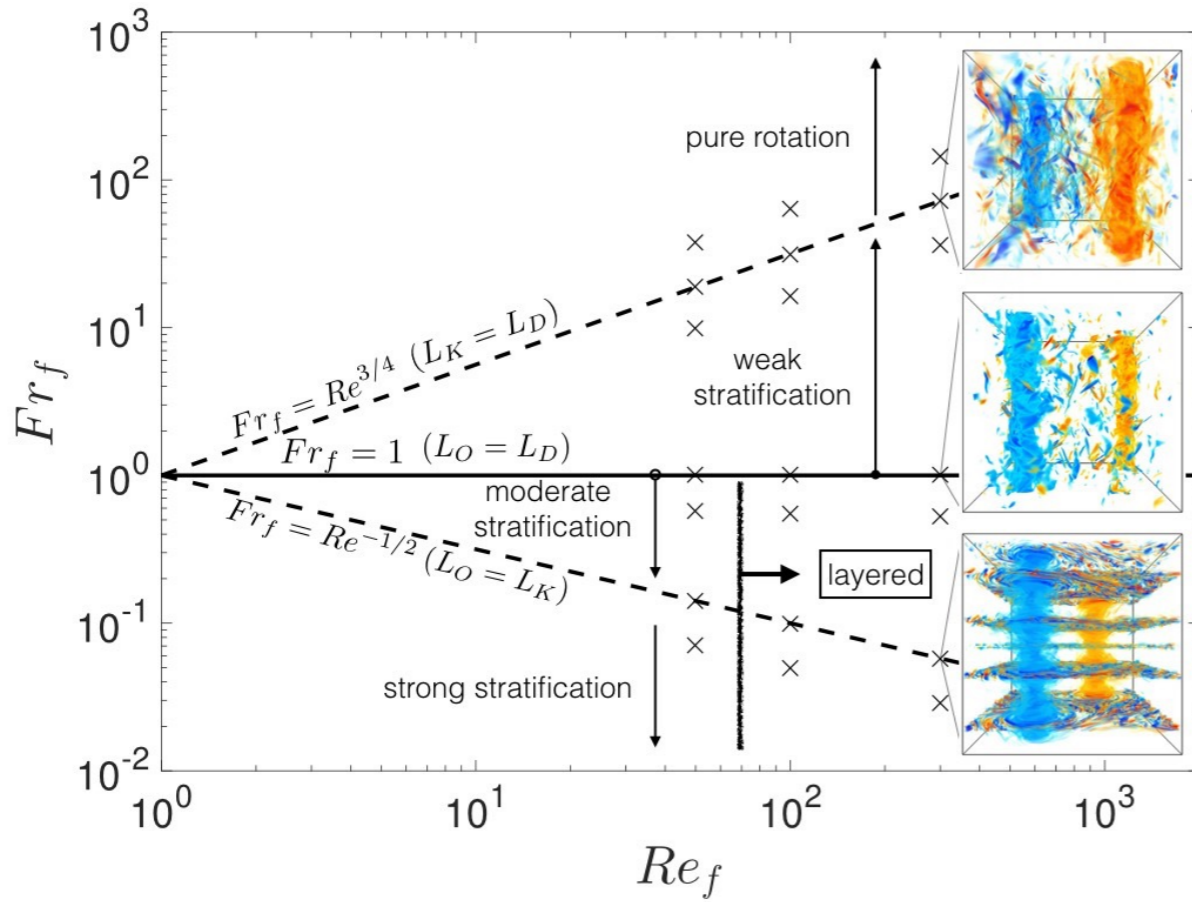
layering here is dissimilar to

- layering in purely stratified turbulence (Kimura & Herring)
- doubly-diffusive systems with steep buoyancy gradients
- K-H shear layer



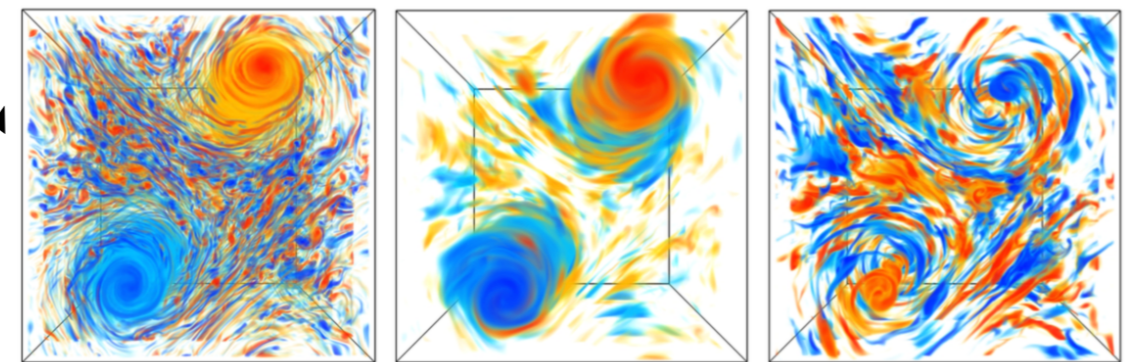
Results (qualitative): large scale barotropic mode

$Re_f = 300, \quad Re_c = 5742 \quad Fr_f > 1$



(a) Vertical vorticity, ζ (b) Buoyancy, b (c) Vertical velocity, w

$Re_f = 300, \quad Re_c = 4290 \quad Fr_f < 1$

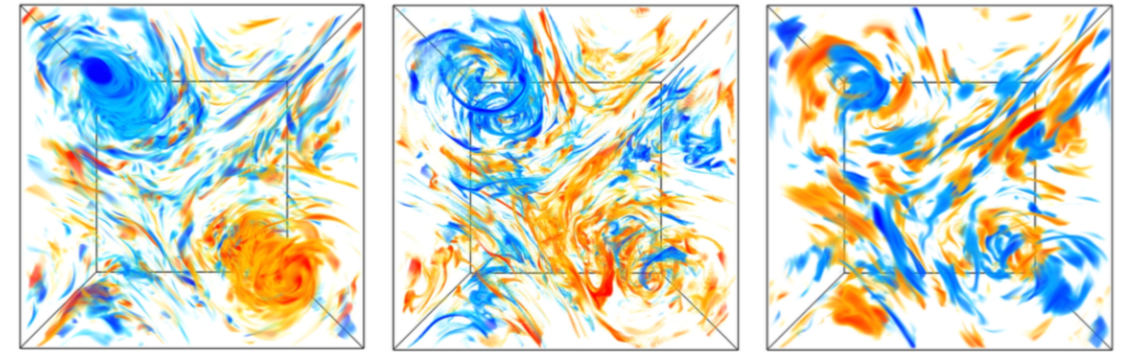
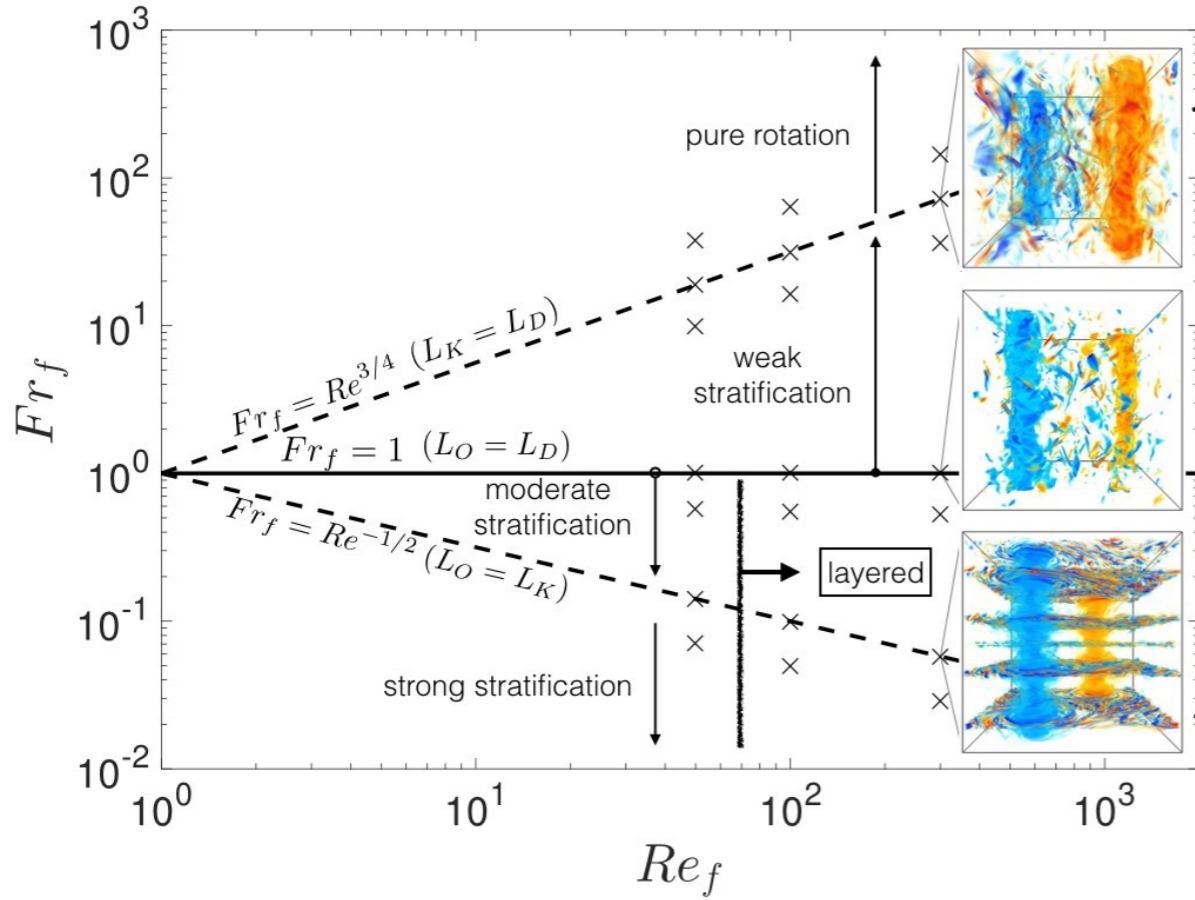


(a) Vertical vorticity, ζ (b) Buoyancy, b (c) Vertical velocity, w

- *Turbulence unobstructed by layers*
- *Does final state preserve a priori classification?*

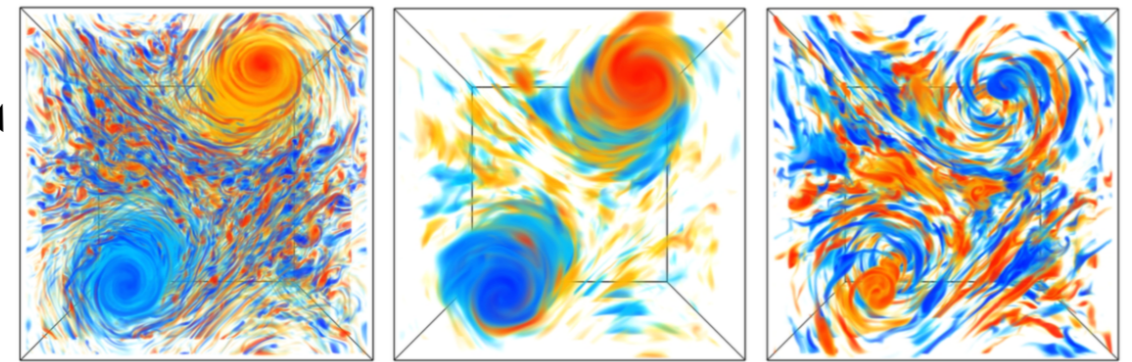
Results (qualitative): large scale barotropic mode

$Re_f = 300, \quad Re_c = 5742 \quad Fr_f > 1$



(a) Vertical vorticity, ζ (b) Buoyancy, b (c) Vertical velocity, w

$Re_f = 300, \quad Re_c = 4290 \quad Fr_f < 1$



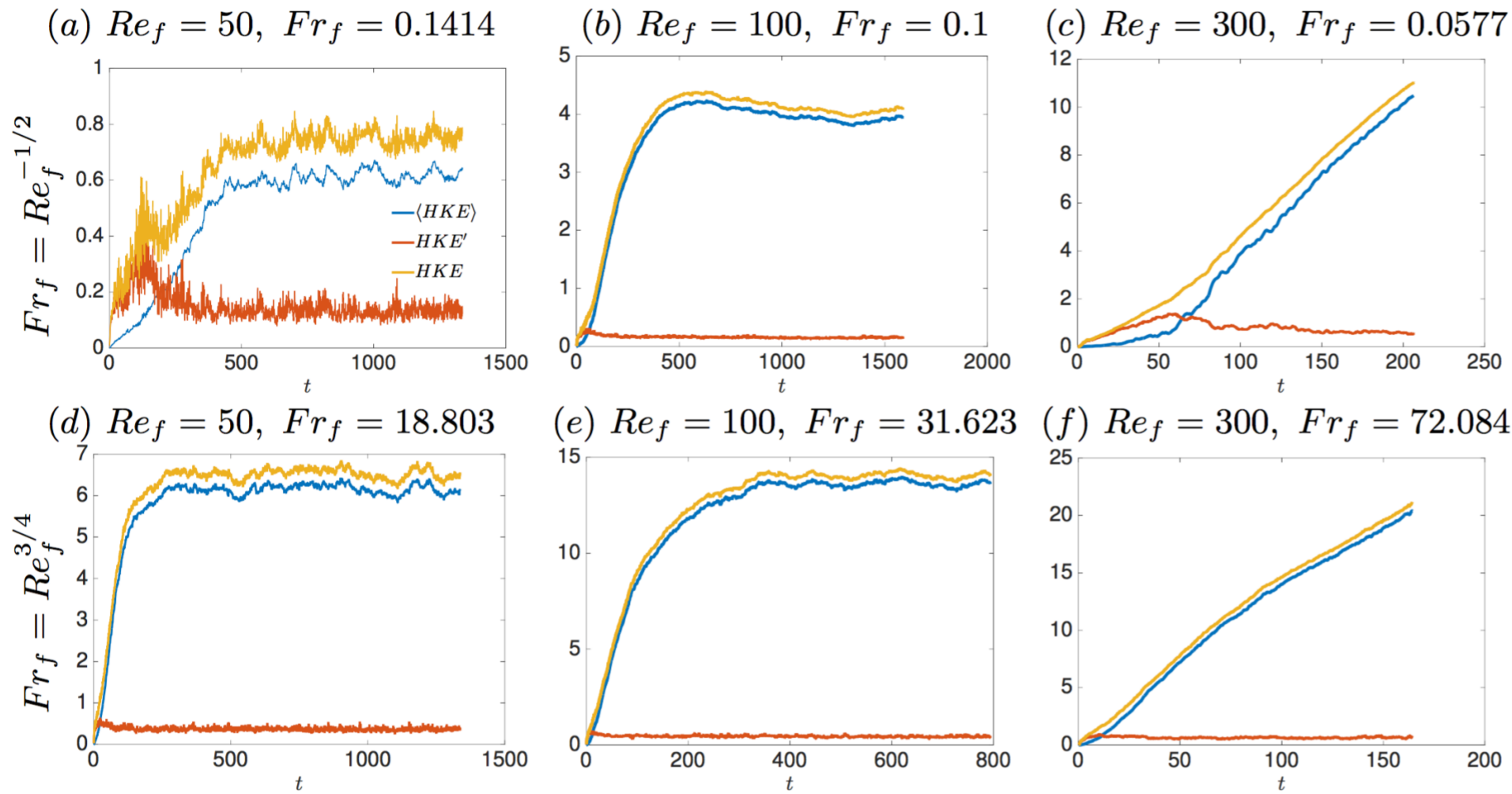
(a) Vertical vorticity, ζ (b) Buoyancy, b (c) Vertical velocity, w

[measured quantities based on energy centroid](#)

$$k_c = \frac{\int k E(k) dk}{\int E(k) dk} \quad Re_c = Re_f U_c L_c, \quad Fr_c = Fr_f \frac{U_c}{L_c}$$

Re_f	Re_c	Fr_f	Fr_c	L_c	U_c
50	150	0.1414	0.0118	6.0	0.5
100	980	0.1000	0.0200	7.0	1.4
300	4290	0.0577	0.0195	6.5	2.2
50	604	18.80	4.50	7.1	1.7
100	2190	31.62	12.99	7.3	3.0
300	5742	72.08	31.67	6.6	2.9

Results: Energetics and equilibrium



- baroclinic & barotropic decomposition:

$$\psi = \langle \psi \rangle + \psi'$$

$$\mathbf{u}_{bt} = \nabla^\perp \langle \psi \rangle \quad (w_{bt}, b_{bt}) = 0$$

$$\mathbf{u}_{bc} = \nabla^\perp \psi' \quad (w_{bc}, b_{bc}) = (w, b)$$

- horizontal kinetic energy (HKE)

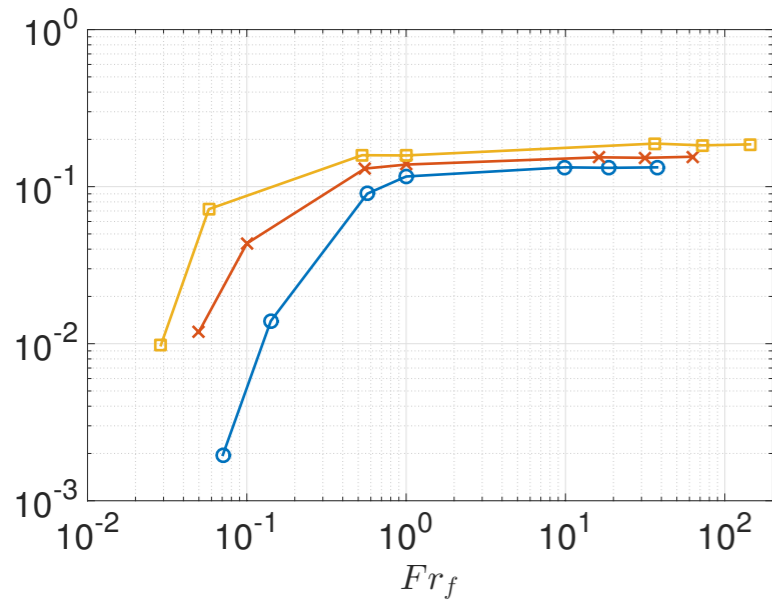
$$\partial_t \langle HKE \rangle := \partial_t \left[\frac{1}{2} \overline{|\nabla_\perp \langle \psi_0 \rangle|^2}^{\mathcal{A}} \right] = \overline{\langle \psi_0 \rangle \langle J[\psi'_0, \zeta'_0] \rangle}^{\mathcal{A}} \dots \dots \dots \text{(baroclinic forcing of barotropic mode, c.f. 2D BVE)}$$

$$\partial_t HKE' := \partial_t \left[\frac{1}{2} \overline{|\nabla_\perp \psi'_0|^2}^{\mathcal{A}} \right] = -\overline{\langle \psi_0 \rangle \langle J[\psi'_0, \zeta'_0] \rangle}^{\mathcal{A}} + \overline{\langle w'_0 \partial_Z \psi'_0 \rangle}^{\mathcal{A}} \dots \dots \text{(vortex stretching: strictly baroclinic)}$$

Results: energy conversion

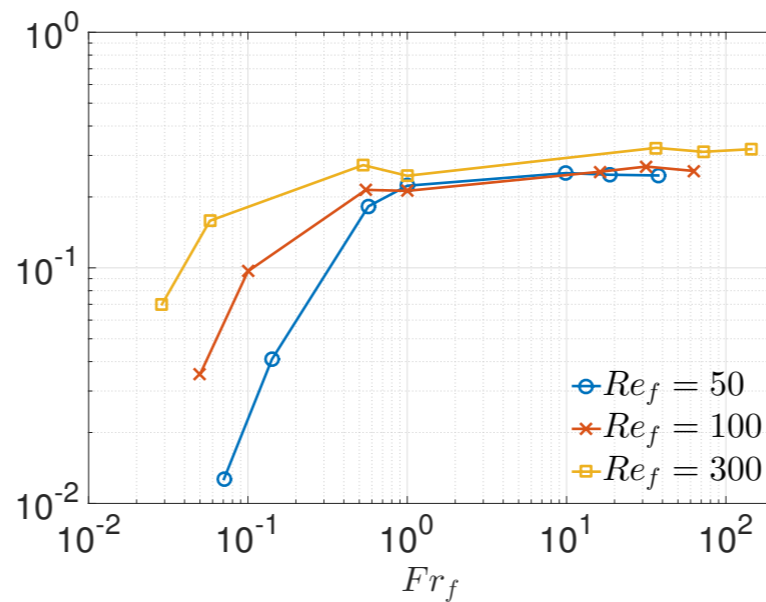
baroclinic forcing

$$\overline{\langle \psi_0 \rangle \langle J[\psi'_0, \zeta'_0] \rangle}^{\mathcal{A}}$$



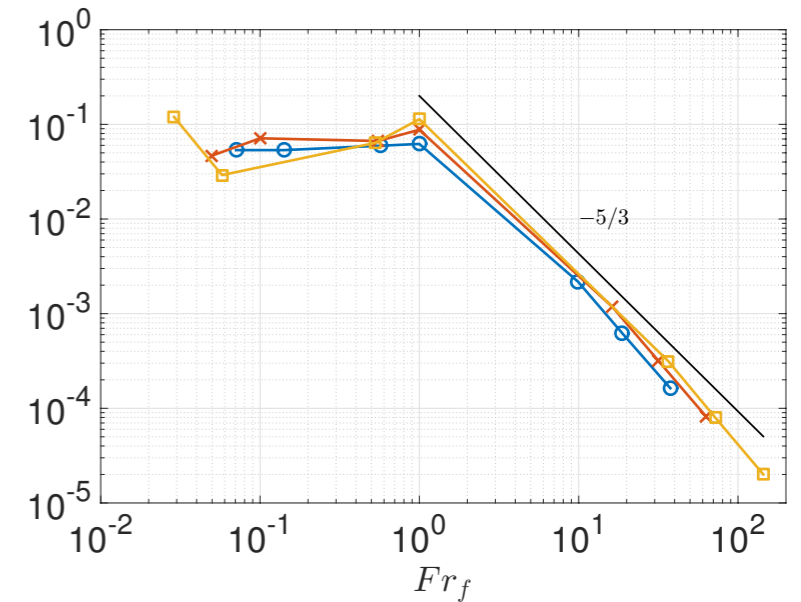
vortex stretching

$$\overline{\langle w_0 \partial_Z \psi_0 \rangle}^{\mathcal{A}}$$

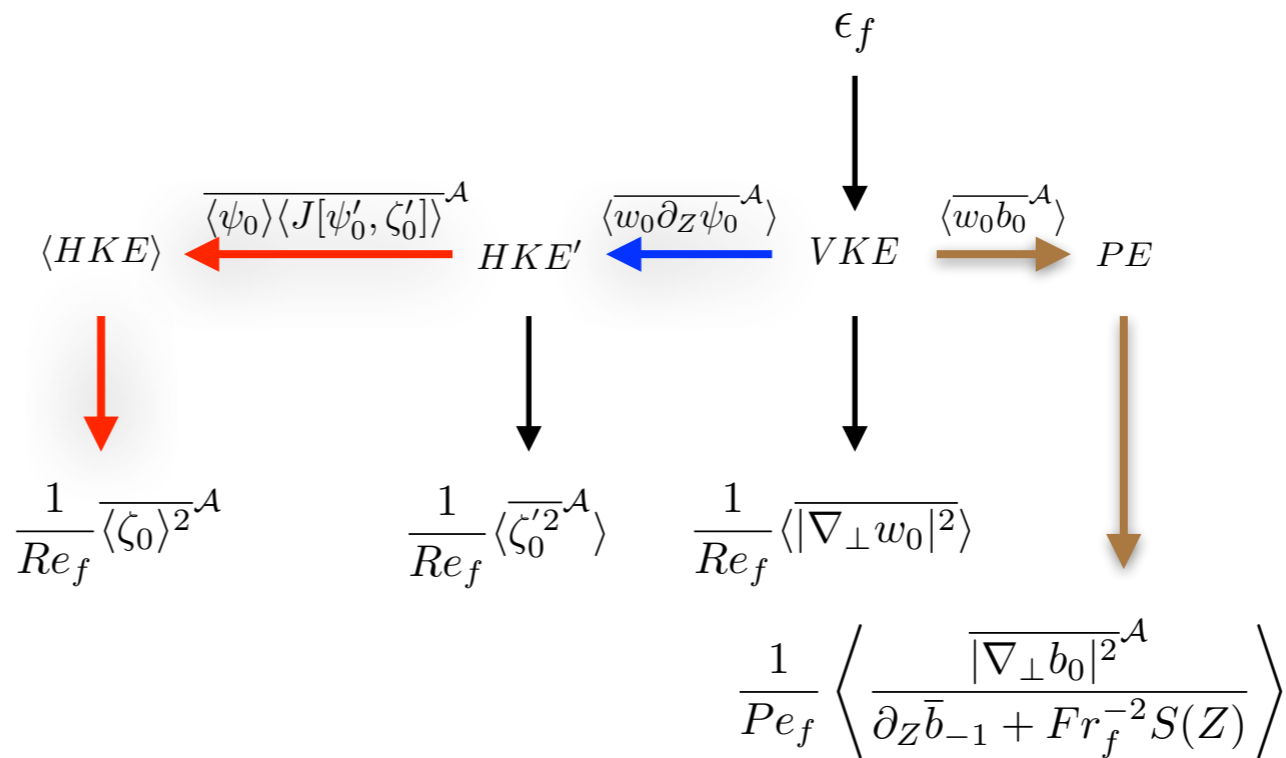


vertical buoyancy flux

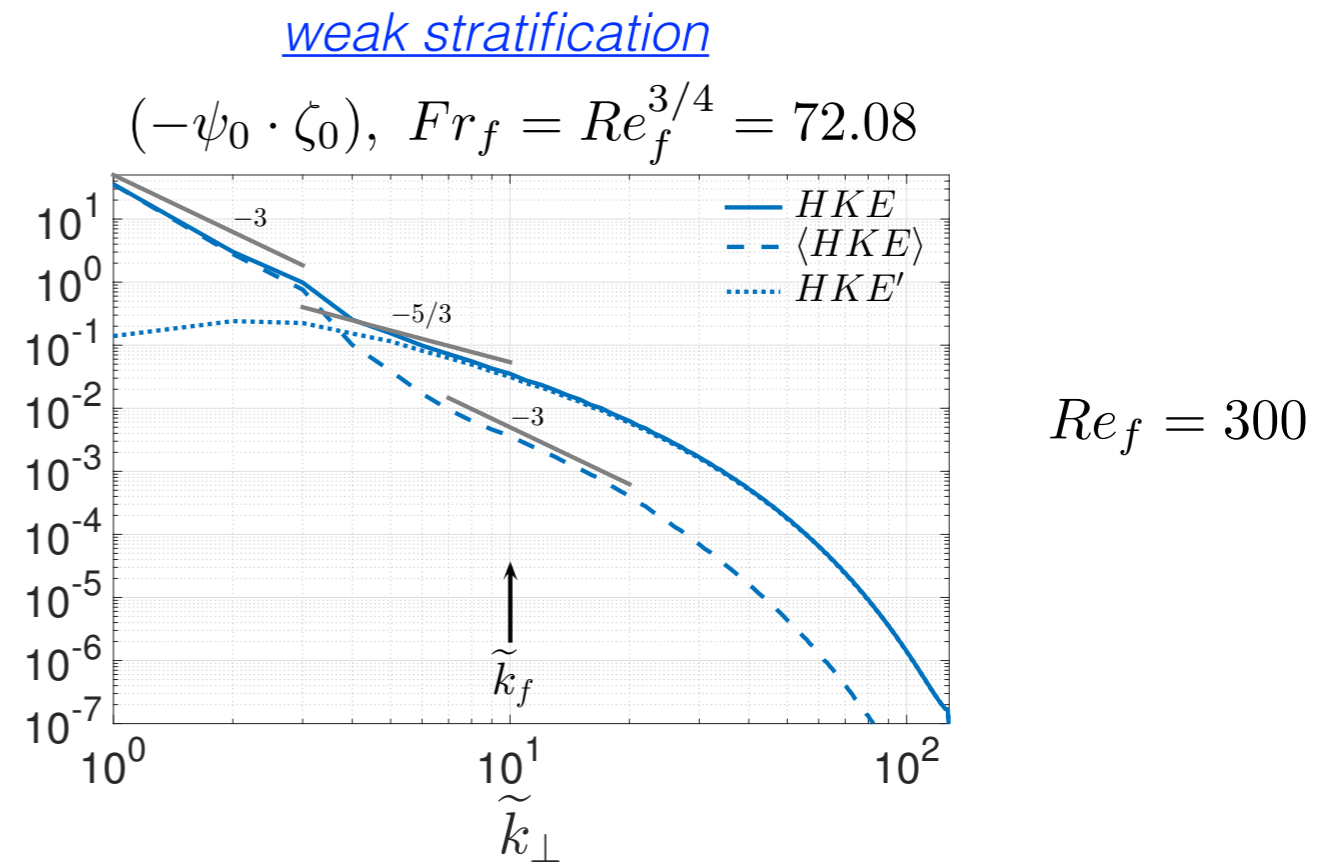
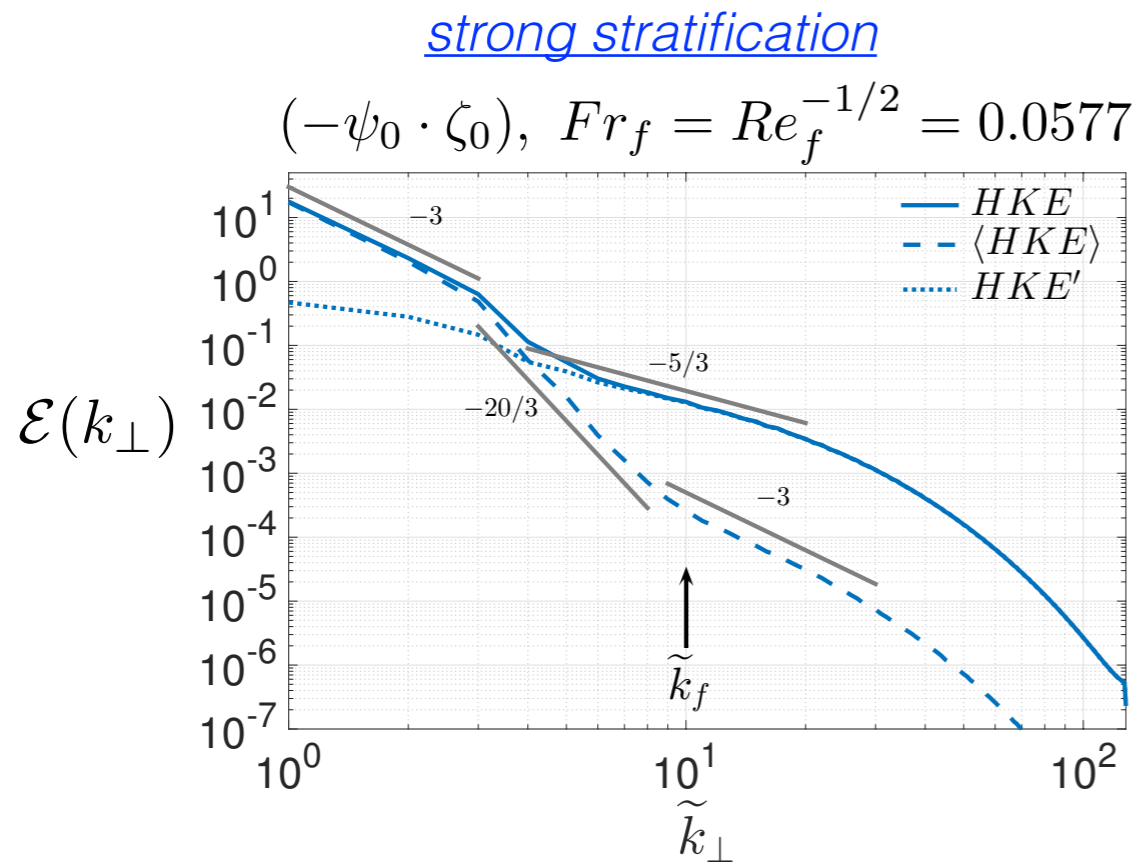
$$-\overline{\langle w_0 b_0 \rangle}^{\mathcal{A}}$$



energy conversions



Results: Energy spectra for energy containing scales (inverse cascade)



- baroclinic & barotropic decomposition:

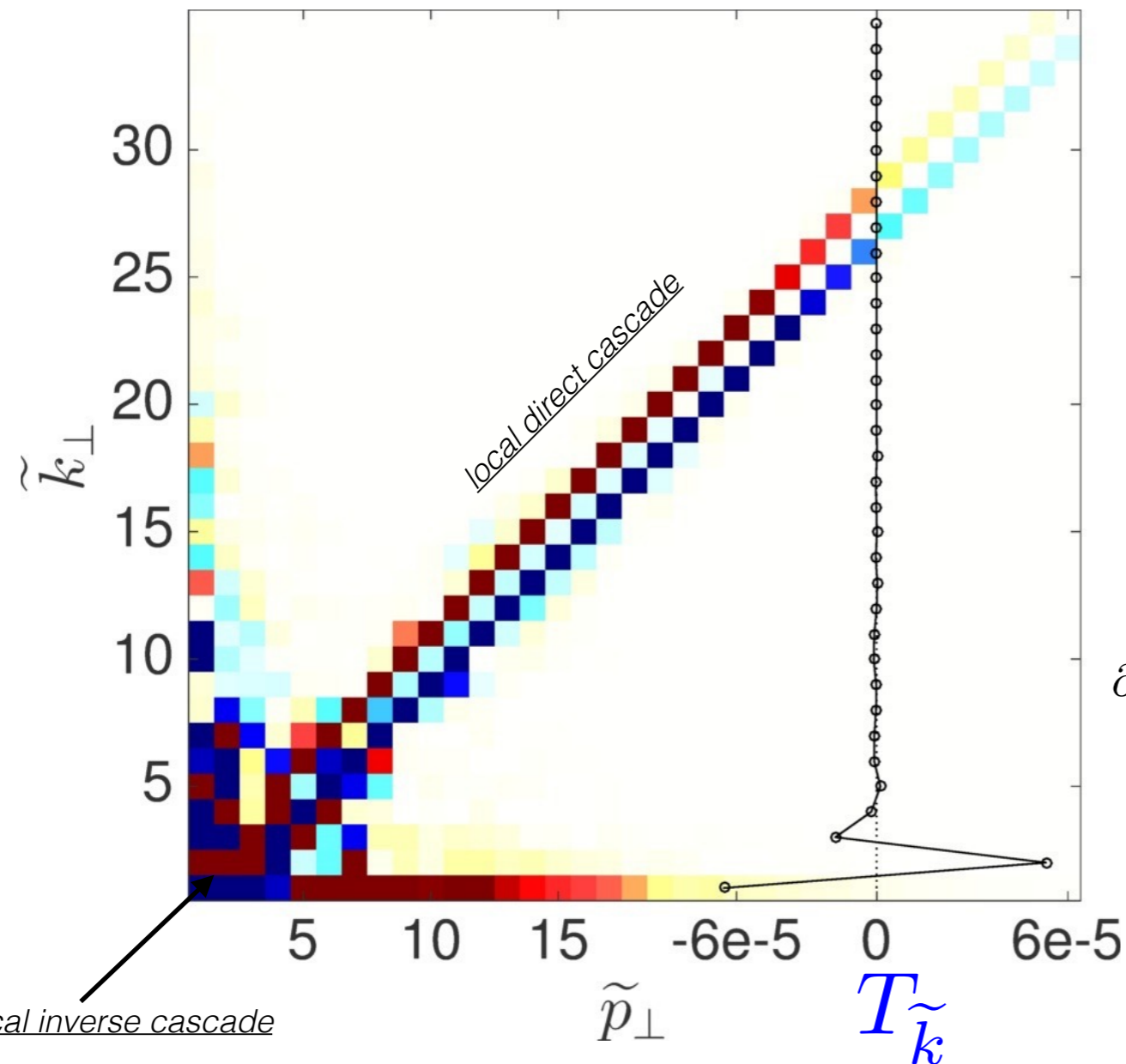
$$\psi = \langle \psi \rangle + \psi'$$

$$\mathbf{u}_{bt} = \nabla^{\perp} \langle \psi \rangle \quad (w_{bt}, b_{bt}) = 0$$

$$\mathbf{u}_{bc} = \nabla^{\perp} \psi' \quad (w_{bc}, b_{bc}) = (w, b)$$

- Steep transition in strong stratification regime result of reduced VKE-HKE conversion
- -3 regime at high wavenumber indicative of forward enstrophy cascade
- -5/3 regime observed in baroclinic HKE - indicative of forward energy cascade

Results: Nonlocal (barotropic) inverse cascade

$$T_{\tilde{k}\tilde{p}}$$


$$\partial_t \langle \zeta_0 \rangle + J[\langle \psi_0 \rangle, \langle \zeta_0 \rangle] = -\langle J[\psi'_0, \zeta'_0] \rangle + \frac{1}{Re_f} \nabla_\perp^2 \langle \zeta_0 \rangle,$$

barotropic triad interactions

$$\partial_t \langle HKE \rangle_k = \sum_{k=p+q} T_{kpq} + \sum_{k=p+q} F_{kpq} + D_k$$

- Transfer map illustrating energy conversion to mode k from mode p
- Local forward cascade co-exist with a nonlocal inverse cascade (through baroclinic interactions)

Summary

- Exploration of *rapidly rotating* and *weakly* stratified geophysical flows
- Two regimes identified:
 1. layered regime (strong stratification $Fr_f < 1$)
 2. columnar regime (strong stratification $Fr_f \geq 1$)
- Distinct from classical QG (where inertia-gravity waves are absent)
- Quantified regimes by energetics
 - observed energy equilibration for moderate Reynolds
 - diagnosed scales active in energy conversion
- ?? DNS verification