

A CFT dual of General Black Holes

A long-term program

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Counting Black Holes in D=4,5.

- **BPS black holes**: very good quantitative understanding in terms of *chiral* 2D CFTs.
- **Other Extremal:** some understanding in terms of *chiral* 2D CFTs.
- Near-BPS: good qualitative understanding in terms of non-chiral 2D CFTs.
- *Generic:* clues suggest understanding in terms of *non-chiral* 2D CFTs.

The *goal of this talk*: review some of these clues.

Setting

- **Setting** (in this talk): D = 4 black holes in N = 4 string theory.
- *Note:* no further restriction, so setting includes Schwarzchild, Kerr....; and also BPS black holes.
- **The proposal:** a dual description in terms of a 2D SCFT with (0, 4) SUSY.
- Some references:
 - FL, arXiv: 9702153.
 - M. Cvetic and FL, arXiv: 9705192, 1106.3341, 1112.4846.
 - A. Castro, A. Maloney, and A. Strominger, arXiv:1004.0996.

Entropy Phenomenology (Kerr)

• The Entropy of Kerr black holes

$$S = 2\pi \left(G_4 M^2 + \sqrt{G_4^2 M^4 - J^2} \right)$$

• The form of the entropy suggests a *chiral split*

$$S = S_L + S_R \, .$$

- The two chiral halves of the CFT interact weakly, at least for the purpose of semi-classical counting.
- Model for Angular Momentum: angular momentum is identified with the R-charge of the (0,4) SCFT.
 Physics: only R-movers in the 2D CFT have the ability to carry J.

Parametric Charges

• The addition of charges in the context of N = 4 SUGRA involves the introduction of parametric mass and angular momentum m, aand parametric charges δ_i :

$$G_4 M = \frac{1}{4} m \sum_{i=1}^{4} \cosh 2\delta_i ,$$

$$G_4 Q_i = \frac{1}{4} m \sinh \delta_i , \ (i = 1, 2, 3, 4) ,$$

$$G_4 J = ma(\Pi_c - \Pi_s) .$$

Abbreviations

$$\Pi_c \equiv \prod_{i=1}^4 \cosh \delta_i , \quad \Pi_s \equiv \prod_{i=1}^4 \sinh \delta_i .$$

Chiral Entropy

• The entropy is much more complicated after general charges have been added

$$S_L = \frac{2\pi m^2}{G_4} (\Pi_c + \Pi_s) ,$$

$$S_R = 2\pi \sqrt{\frac{m^4}{G_4^2} (\Pi_c - \Pi_s)^2 - J^2}$$

• An *encouraging feature*: the dependence on angular momentum can still be accounted for by the physical assumption that only right movers are able to carry angular momentum.

Diffeomorphism Anomaly

• Define *physical temperatures for left and right movers independently* by the generalized first laws

$$dS_L = \beta_L dM$$
 , $dS_R = \beta_R dM$.

(It follows that the Hawking temperature is $\beta_H = \beta_L + \beta_R$).

• Model chiral contributions to black hole entropy as a 1D gas on a circle with radius \mathcal{R} :

$$S_{R,L} = \frac{\pi}{6} c T_{R,L} \times 2\pi \mathcal{R} \; .$$

• An *encouraging feature*: the central charges determined this way are not chiral ($c_R = c_L$).

So there is *no diffeomorphism anomaly* in the semi-classical theory, as expected.

Level Matching

• The model for the semiclassical entropy employes the Cardy form:

$$S = 2\pi\sqrt{kh_L} + 2\pi\sqrt{kh_R}$$

Notation: c=6k for (0, 4) SCFTs.

• An *encouraging feature*: there is a quantization condition

$$k(h_L - h_R) = J^2 + J_4$$
.

that can be interpreted as level matching of the 2D CFT dual to the general black holes.

Notation: the quartic invariant is the integer

$$J_4 = \vec{q}^2 \vec{p}^2 - (\vec{q} \cdot \vec{p})^2$$

• *Moduli Independence*: the charge vectors \vec{q} , \vec{p} depend on moduli but J, J_4 do not. This is a generalized *attractor mechanism*.

Strategy

- *Goal*: find further evidence for 2D conformal geometry.
- Strategy (for now): analyze details of the geometry.
- *Inspiration*: analysis of BTZ black holes (using AdS₃ geometry), analysis of Kerr black holes using near horizon symmetry....
- Strategy (medium term): precise description by tensoring of chiral CFT's that succesfully describe extremal black holes, using level matching to constrain 0-modes.

The Black Hole Geometry

The *explicit geometry* with all charges turned on:

$$ds_4^2 = -\Delta^{-\frac{1}{2}}G(dt + \mathcal{A}_{\phi}d\phi)^2 + \Delta^{\frac{1}{2}}\left(\frac{dr^2}{X} + d\theta^2 + \frac{X\sin^2\theta}{G}d\phi^2\right) ,$$

where

$$X = r^{2} - 2mr + a^{2},$$

$$G = r^{2} - 2mr + a^{2} \cos^{2} \theta,$$

$$\mathcal{A} = \frac{2ma \sin^{2} \theta}{G} [(\Pi_{c} - \Pi_{s})r + 2m\Pi_{s}] d\phi,$$

$$\Delta_{0} = \prod_{i=1}^{4} (r + 2m \sinh^{2} \delta_{i}) + 2a^{2} \cos^{2} \theta [r^{2} + mr \sum_{i} \sinh^{2} \delta_{i} + 4m^{2} (\Pi_{c} - \Pi_{s})\Pi_{s} + 2m^{2} \sum_{i < j < k} \sinh^{2} \delta_{i} \sinh^{2} \delta_{j} \sinh^{2} \delta_{k} + \frac{1}{2}a^{2}]$$

Abbreviation: $\mathcal{A}_{\mathrm{red}} = \frac{G}{a \sin^2 \theta} \mathcal{A}_{\phi}$ depends on r alone.

The Laplacian

The entropy formula suggests a dual 2D CFT. But the corresponding features in the scalar wave equation are much more precise.

The general Laplacian:

$$\Delta_0^{-\frac{1}{2}} [\partial_r X \partial_r - \frac{1}{X} (\mathcal{A}_{\mathrm{red}} \partial_t - \partial_\phi)^2 + \frac{\mathcal{A}_{\mathrm{red}}^2 - \Delta_0}{G} \partial_t^2 + \frac{1}{\sin \theta} \partial_\theta \sin \theta \partial_\theta + \frac{1}{\sin^2 \theta} \partial_\phi^2]$$

Recall: Kerr black holes allow separation of variables. The addition of charges (largely) maintains this property:

$$\frac{\mathcal{A}_{\text{red}}^2 - \Delta_0}{G} = r^2 + 2mr \left(1 + \sum_{i=1}^4 s_i^2 \right) + 8m^2 (\Pi_c - \Pi_s) \Pi_s$$
$$-4m^2 \sum_{i < j < k} s_i^2 s_j^2 s_k^2 + a^2 \cos^2 \theta \,.$$

Hypergeometric Structure?

- The resulting radial equation is "almost" hypergometric.
- If it were precisely hypergeometric we could find an $SL(2)^2$ symmetry which perhaps would be *enhanced to conformal symmetry*.
- The *obstacle* is that the effective potential

$$V_{\rm eff} = \frac{\mathcal{A}_{\rm red}^2 - \Delta_0}{G}$$

rises too quickly far from the black hole: $V_{\rm eff} \sim r^2$ because $\Delta_0 \sim r^4$ corresponds to asymptotically flat space and $G \sim r^2$.

• *Interpretation*: we should *enclose the black hole in a box* that regulate asymptotic behavior.

The Subtraction Prescription

- *Prescription:* Modify the conformal factor so that $\Delta \sim r$ instead of $\Delta_0 \sim r^4$.
- **Consequence:** the scalar wave equation in the subracted geometry is **exactly** a hypergeometric equation.
- **Technical Assumptions:** determine Δ uniquely by assuming separability is maintained, and also that Δ is analytical in r, θ .

 $\Delta = (2m)^3 (\Pi_c^2 - \Pi_s^2)r + (2m)^4 \Pi_s^2 - (2m)^2 (\Pi_c - \Pi_s)^2 a^2 \cos^2 \theta$

• Interpretation: Δ encodes the environment of the black hole, not its internal structure.

General Causal Structure

• Recall the *full metric*:

$$ds_4^2 = \Delta^{\frac{1}{2}} \left(\frac{dr^2}{X} + \frac{X\sin^2\theta}{G} d\phi^2 \right) + \left[\Delta^{\frac{1}{2}} d\theta^2 - \Delta^{-\frac{1}{2}} G (dt + \mathcal{A}_{\phi} d\phi)^2 \right]$$

Assume X = X(r), $G = X - a^2 \sin^2 \theta$ (and \mathcal{A}_{ϕ} , Δ general functions of r, θ).

- The *ergosphere* (inside static limit): G = 0.
- The *event horizon* (t, ϕ subspace changes signature): X = 0.

General Thermodynamics

- Remain in the *general context* where $X, \mathcal{A}_{red}, \Delta$, are not specified.
- The *thermodynamics* determined for these black holes:

$$\Omega_H = \frac{a}{(\mathcal{A}_{\text{red}})_{\text{hor}}}, \\ \beta_H = \frac{4\pi (\mathcal{A}_{\text{red}})_{\text{hor}}}{(\partial_r X)_{\text{hor}}}, \\ S = \frac{\pi (\mathcal{A}_{\text{red}})_{\text{hor}}}{G_4}.$$

• The *striking feature*: the conformal factor Δ does not appear in the thermodynamics.

It also did not appear in the causal structure.

• Interpretation: Δ is a feature of the embedding into an ambient spacetime.

Supporting Matter

- **Issue:** the subtracted geometry (with modified Δ) does not satisfy the equations of motion with the original matter.
- Interpretation: the sources supporting the modified Δ represents the physical matter that realizes the box enclosing the black hole.
- In situations (such as near extremality) where the offending terms can be approximated away the "additional" matter is negligible.
- In general the *matter is sensible*: it satisfies "all" energy conditions.
- *A simple realization* of the matter: 5D minimal SUGRA, KK reduced to 4D.

A 5D Lift

 It is illuminating to *lift the 4D geometry to a 5D geometry*, via an auxiliary coordinate α:

$$ds_{5}^{2} = \Delta (d\alpha + \mathcal{B})^{2} + \Delta^{-1/2} ds_{4}^{2} = -\frac{X}{\rho} dt^{2} + \frac{dr^{2}}{X} + \frac{\rho}{4m^{2}(\Pi_{c} - \Pi_{s})^{2}} (d\alpha + \frac{\mathcal{A}_{\text{red}}}{\rho} dt)^{2} + d\Omega_{2}^{'2} .$$

The linearly shifted radial coordinate:

$$\rho = 8m^3 [r(\Pi_c^2 - \Pi_s^2) - \frac{a^2}{2m} (\Pi_c - \Pi_s)^2 + 2m\Pi_s^2].$$

- In the 5D form *separability is manifest.*
- **Bonus**: massive scalar fields that couple minimally to the 5D geometry are also separable.

Virasoro

- The full 5D space is *locally* $AdS_3 \times S^2$.
- So there is Virasoro² symmetry and also an SU(2) R-symmetry.
- Except: the reduction of the 5D space to the physical 4D spacetime introduces *global issues*.
- Concretely: physical deformations are independent of α so the full Virasoro is not realized on the physical space.

Central Charge

• The *phenomenological model* determines the CFT central charge in units of the CFT radius \mathcal{R} :

$$c\mathcal{R} = \frac{24m^3}{G_4} (\Pi_c^2 - \Pi_s^2) .$$

The cubic dependence on mass generalizes the MSW formula far off extremality.

• The scale \mathcal{R} relates the physical spacetime temperature to the dimensionless CFT temperature:

$$T_{L,R}^{ ext{CFT}} = T_{L,R}^{ ext{phys}} \mathcal{R} \; .$$

• The *lift to 5D* leaves \mathcal{R} arbitrary *a priori*.

Summary

Progress towards conformal symmetry of general black holes:

- The *phenomenology of the black hole entropy formula* suggests a 2D CFT origin.
- Black hole thermodynamics in a manner that exhibits *independence of the conformal factor*.
- The *subtraction procedure* isolates a part of the geometry that has $SL(2)^2$ symmetry.
- Rudiments of a *full Virasoro symmetry* even for very general black holes.