

Migration in selfgravitating discs

Evolution of supermassive black hole binaries at centi pc scale

Retrograde versus prograde

Constanze Roedig

with Alberto Sesana

The Johns Hopkins University, Baltimore

Albert Einstein Insitut, Potsdam

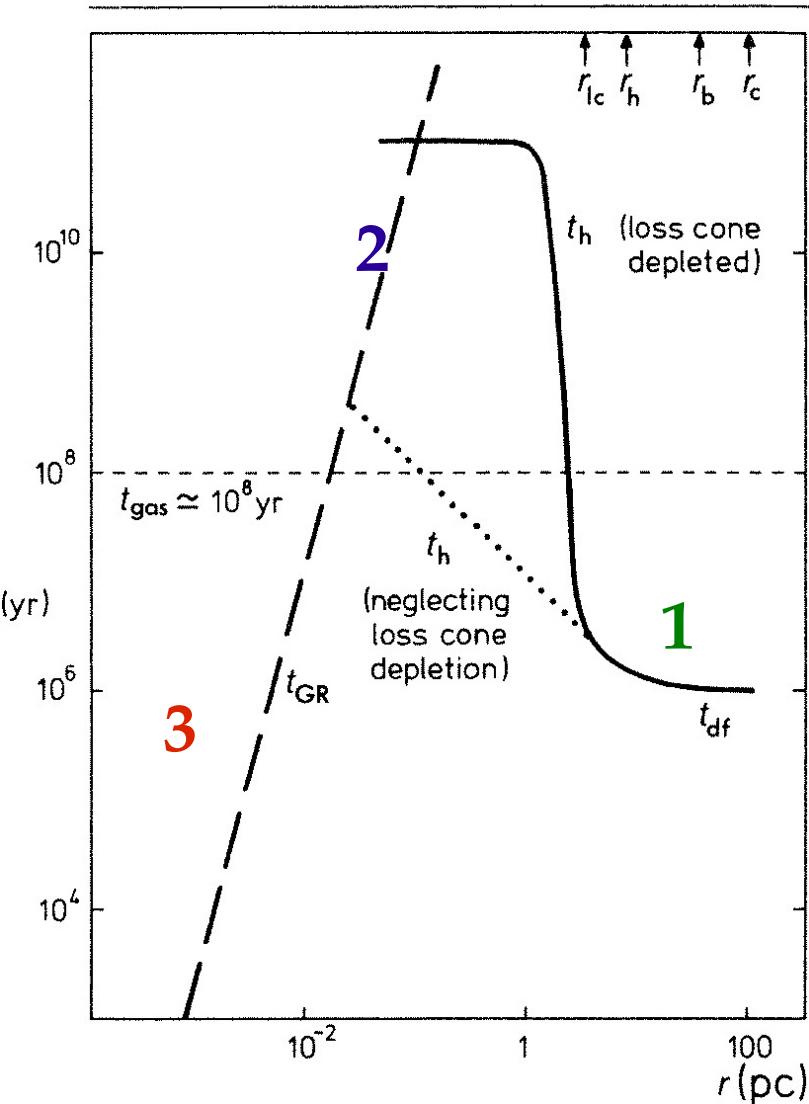
Outline:

- ✓ Introducing...
 - ...the physical problem: SMBHB in circumbinary discs
 - ...what we've learnt previously
- ✓ Flipping the Sign: retrograde migration in a selfgravitating disc
 - insights from comparisons && analytical explanations
- ✓ The final fate: Introducing the tilting instability for high e
- ✓ Summary

the BH merger paradigm

(Begelman, Blandford, Rees 1980)

Nature Vol. 287 25 September



(1) dynamical friction regime

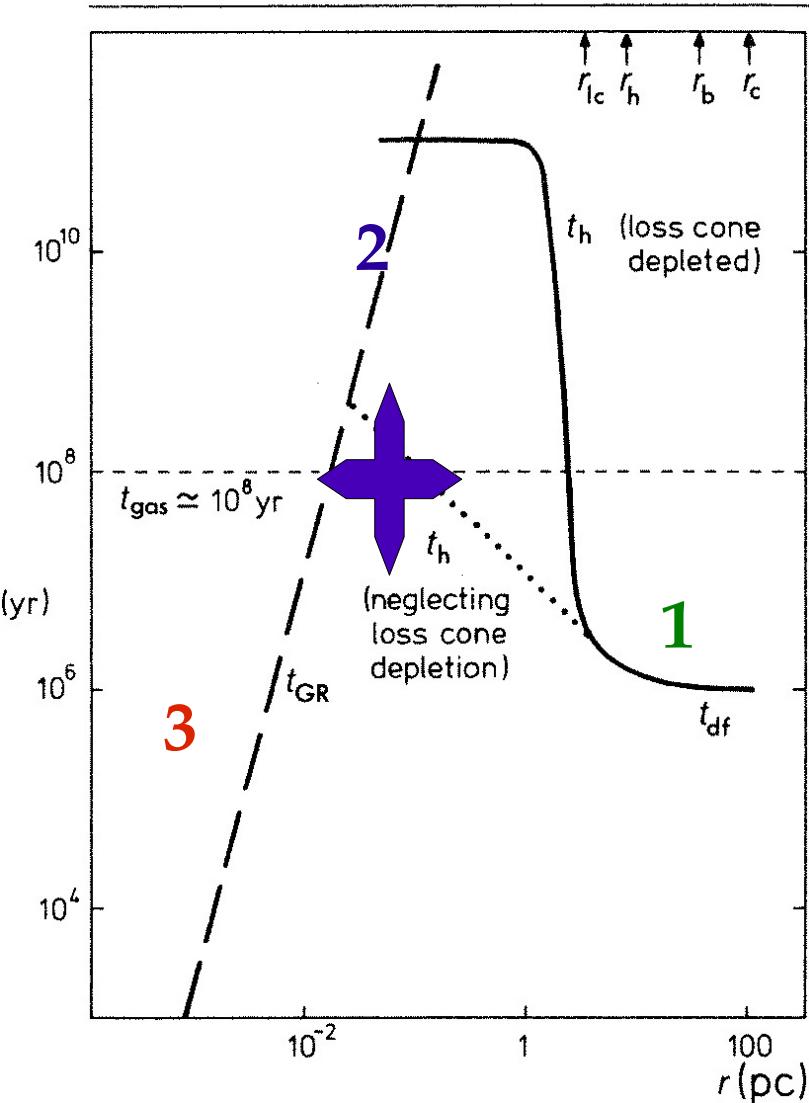
(2) binary hardening regime
by interaction with
gas and/or stars

(3) emission of GWs
inspiral/merger/ringdown

Focussing on centi-parsec scale

(Begelman, Blandford, Rees 1980)

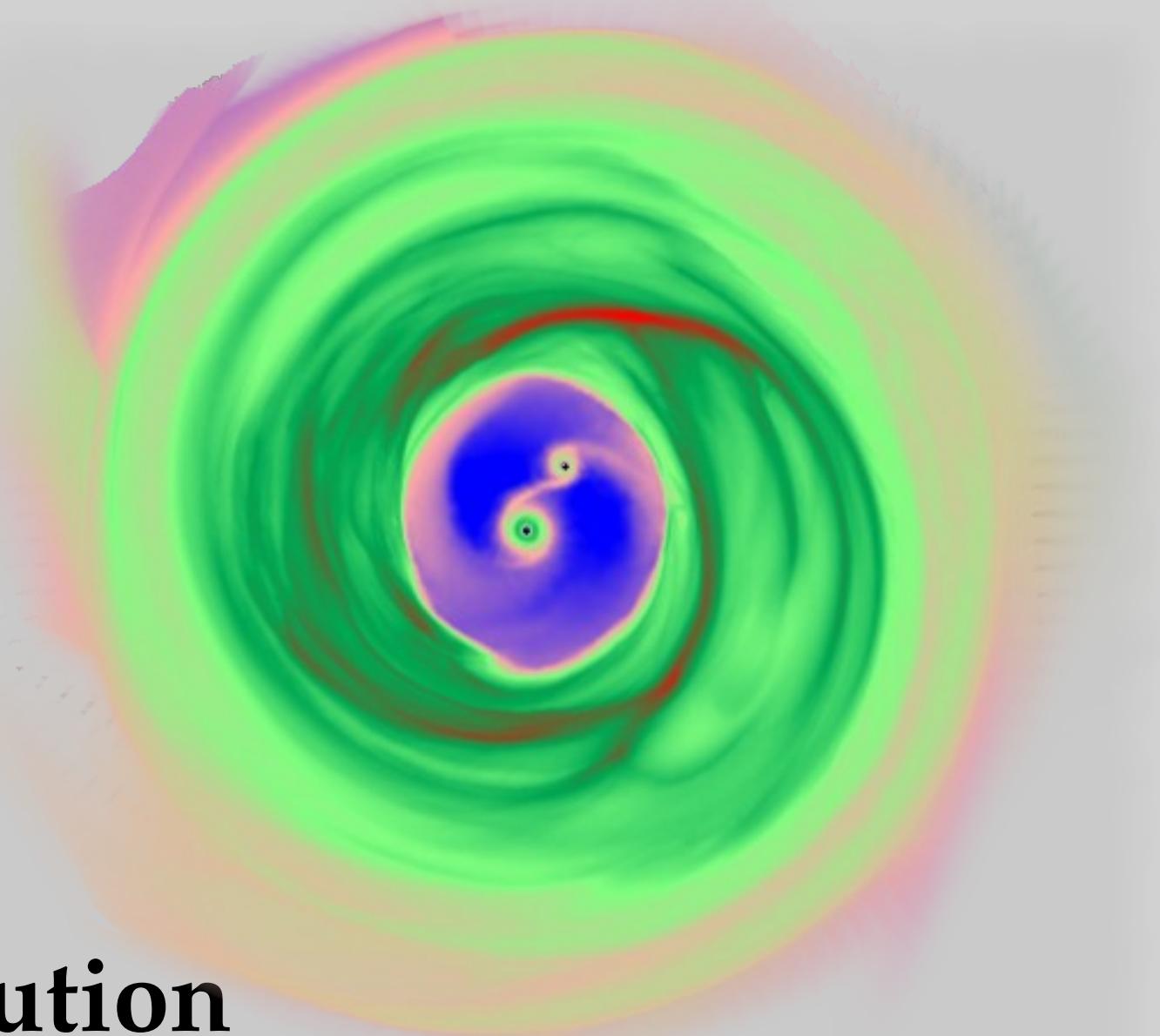
Nature Vol. 287 25 September



(1) dynamical friction regime

(2) binary hardening regime
by interaction with
gas and/or stars

(3) emission of GWs
inspiral/merger/ringdown

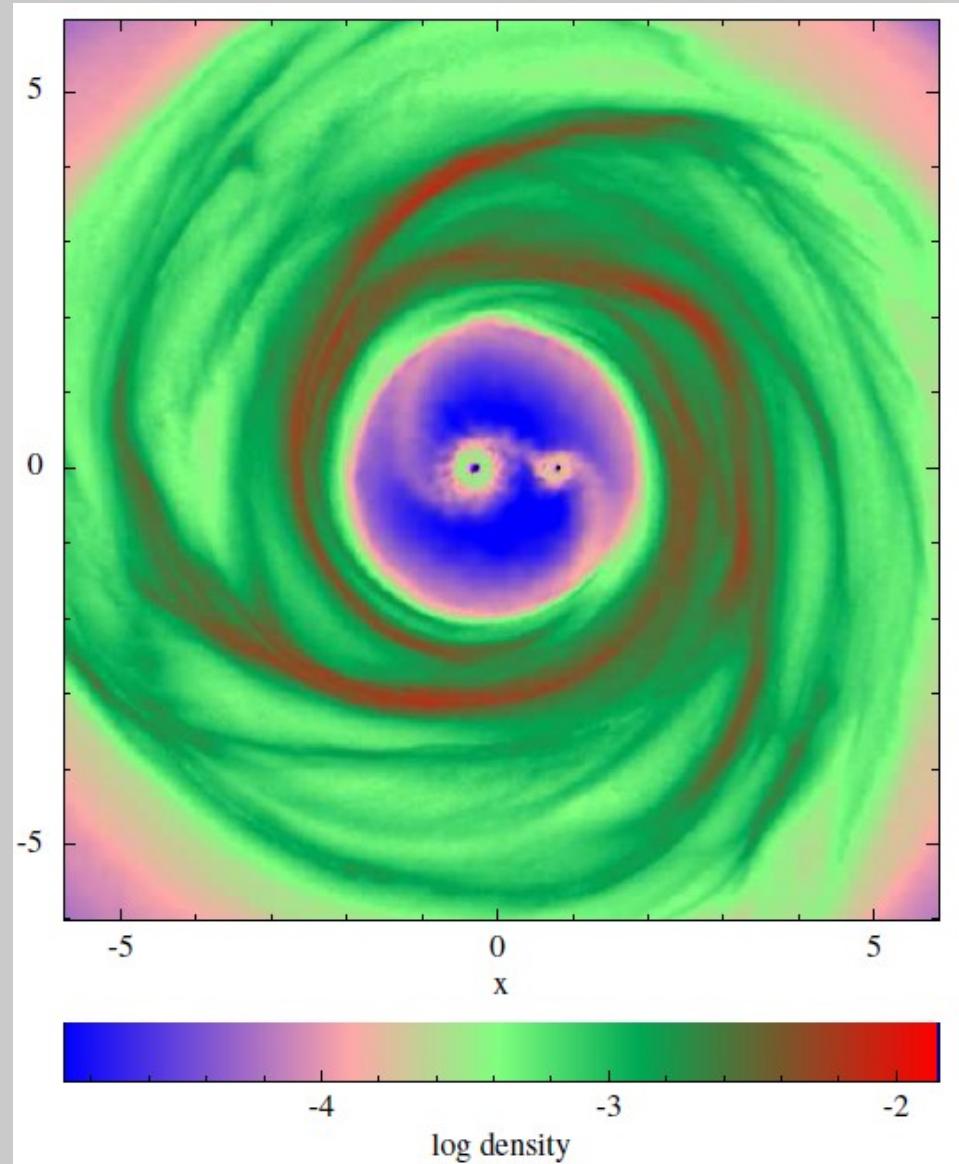


Secular evolution of binaries in self-gravitating discs

Secular evolution using GADGET-2

Disc needed to redistribute Energy E and Angular momentum L

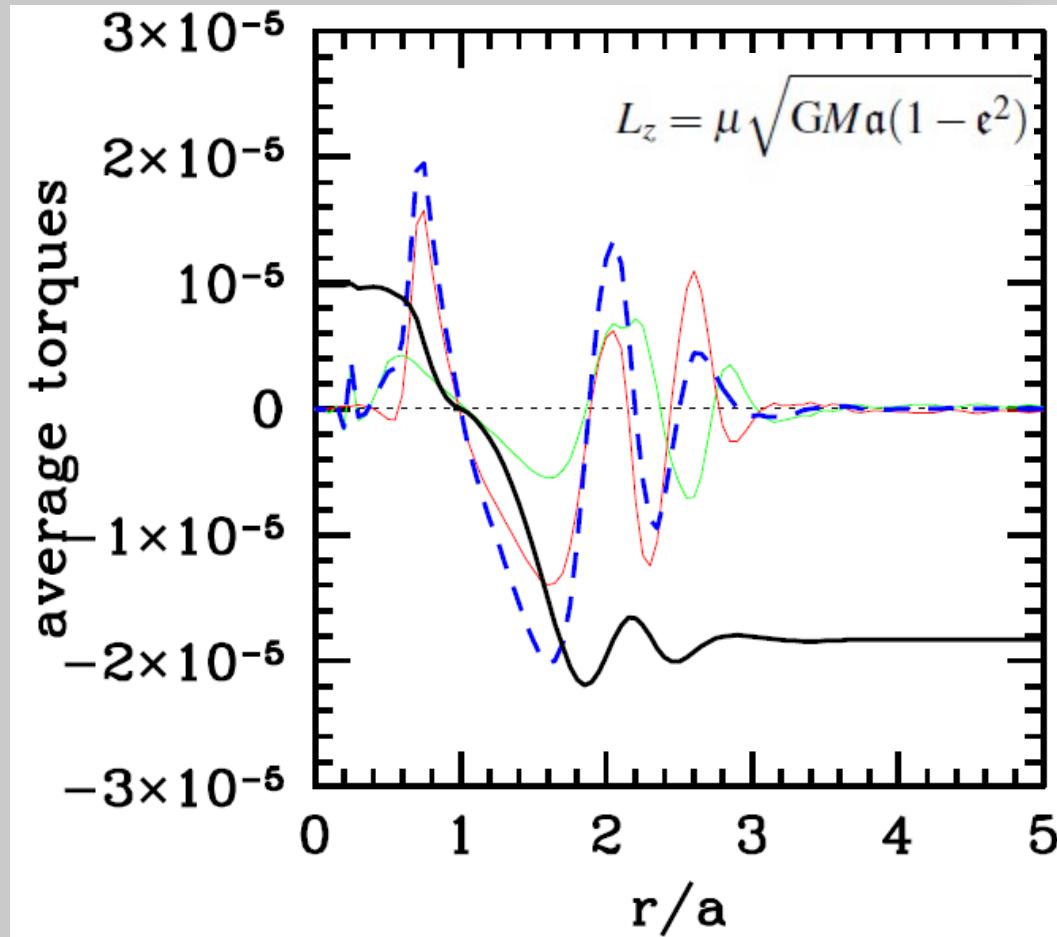
- Selfgravity as "viscosity"
- $\beta = t_{\text{cool}}/t_{\text{dyn}} = 10$
- "life" binary
- Accretion transfers mass + momentum
- Cavity of size $2a$ sustained
- Formation of minidisks on t_{dyn}
- Mean motion resonances not applicable
kinematic streams & "stream-bending"
dynamically important



Secular evolution using GADGET-2

Gravitational and accretion torques change a and e

- Prograde disc:
- Strongest torque inside cavity
- Opposite sign torque at minidisc
- Total torque of disc onto binary:
Negative --> a shrinks
--> e grows
 $\text{cavity expands to } 2a(1+e)$
- Balance reached at $e \sim 0.6$
- Shrinkage $\dot{a} \sim \dot{M}$ (!!)

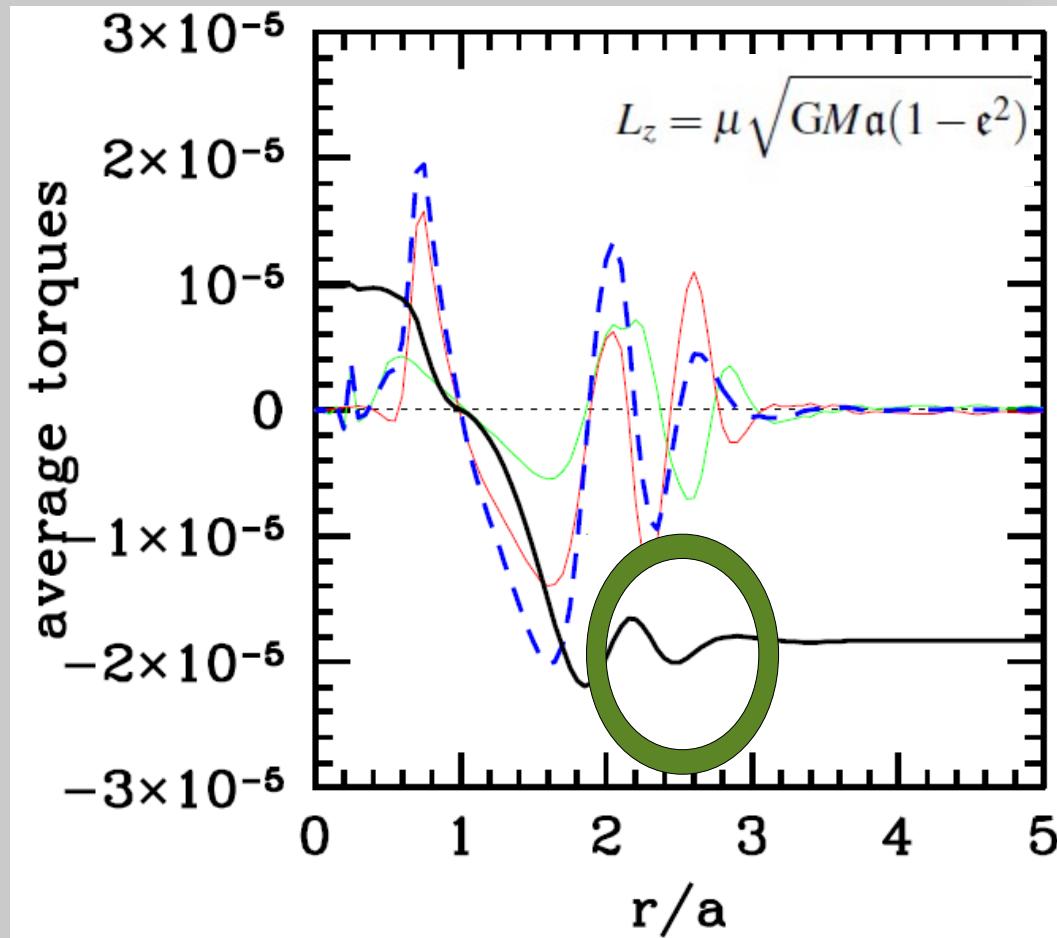


(Roedig et al, 2011, 2012, 2013)

Secular evolution using GADGET-2

Gravitational and accretion torques change a and e

- Prograde disc:
- Strongest torque inside cavity
- Opposite sign torque at minidisc
- Total torque of disc onto binary:
Negative --> a shrinks
--> e grows
cavity expands to $2a(1+e)$
- Balance reached at $e \sim 0.6$
- Shrinkage $\dot{a} \sim \dot{M}$ (given by Σ)

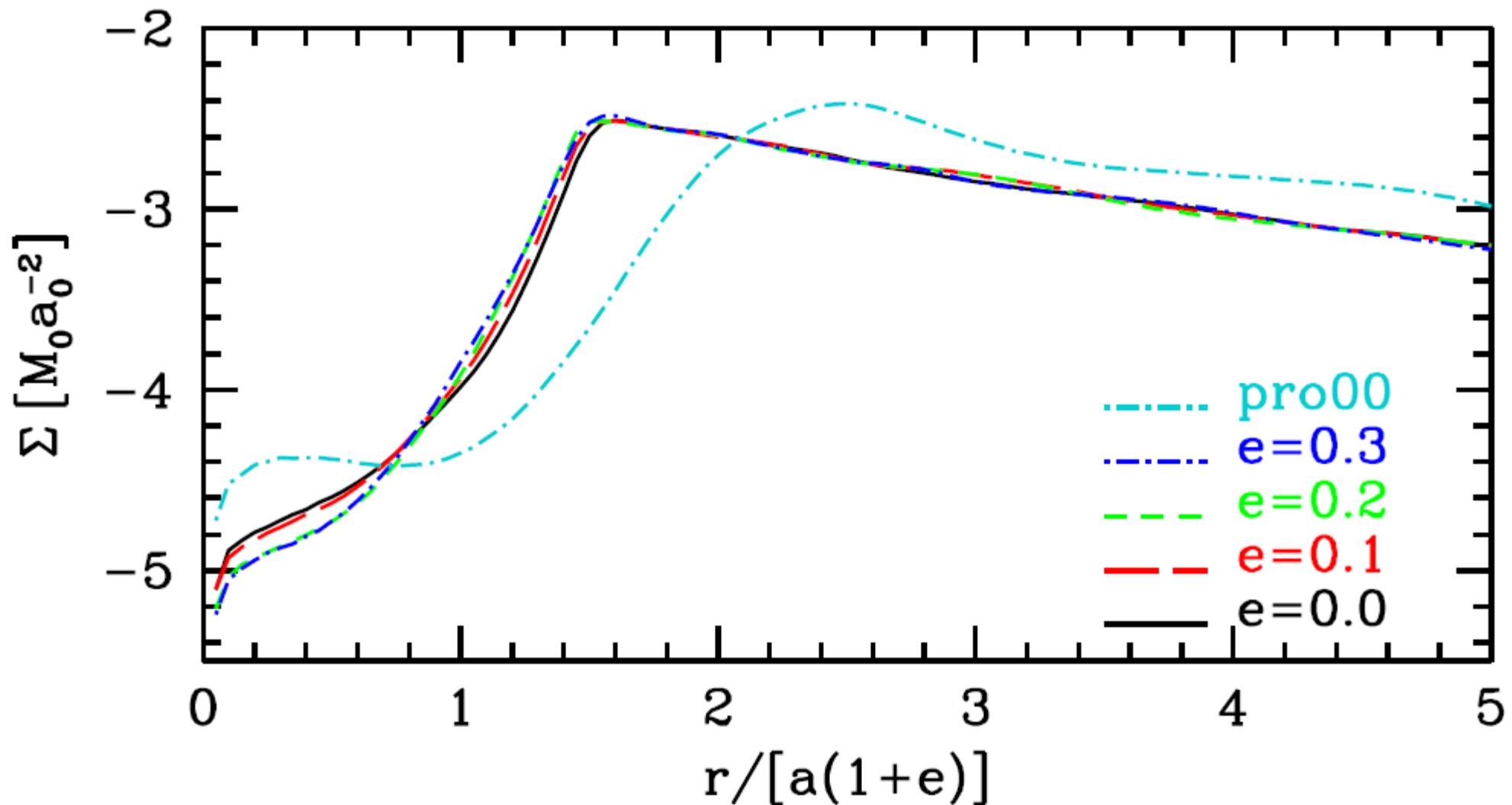


(Roedig et al, 2011, 2012, 2013)

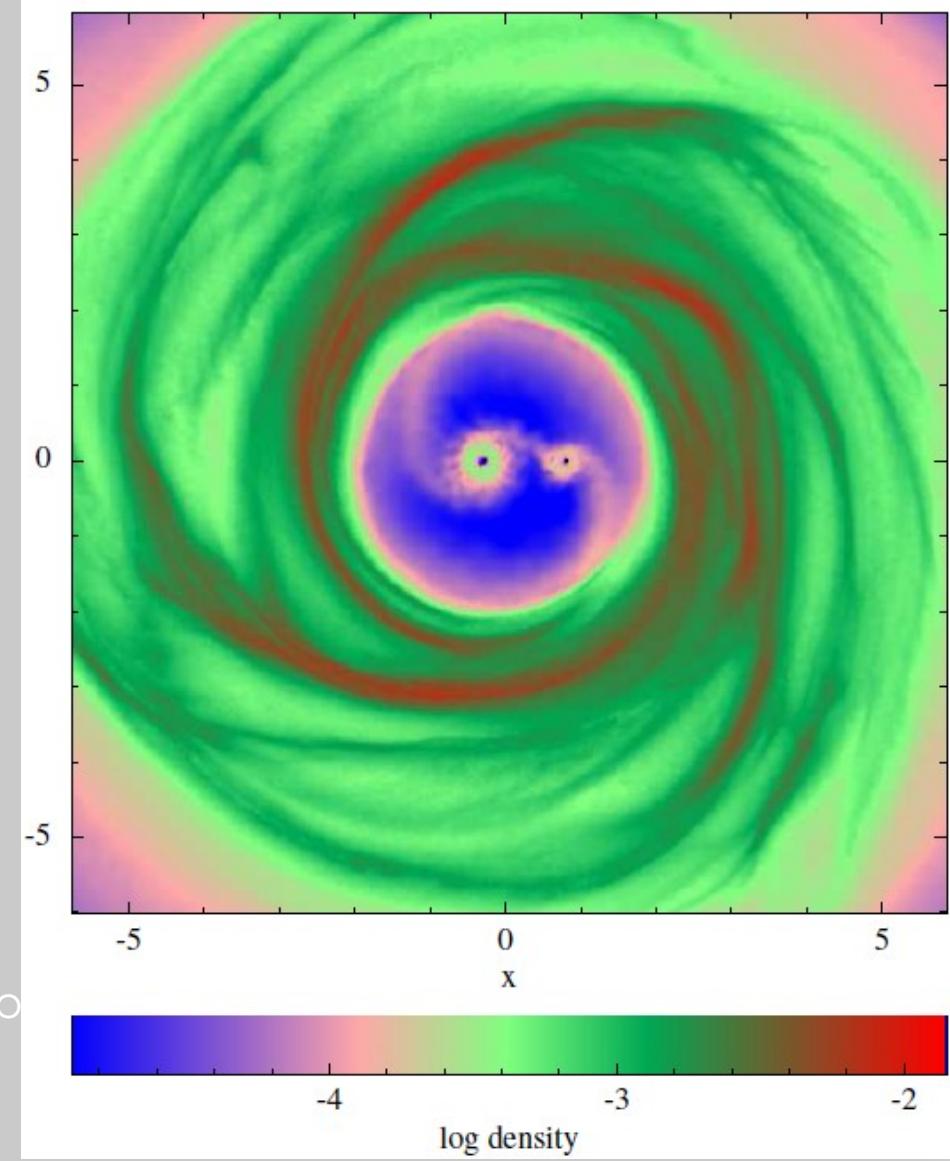
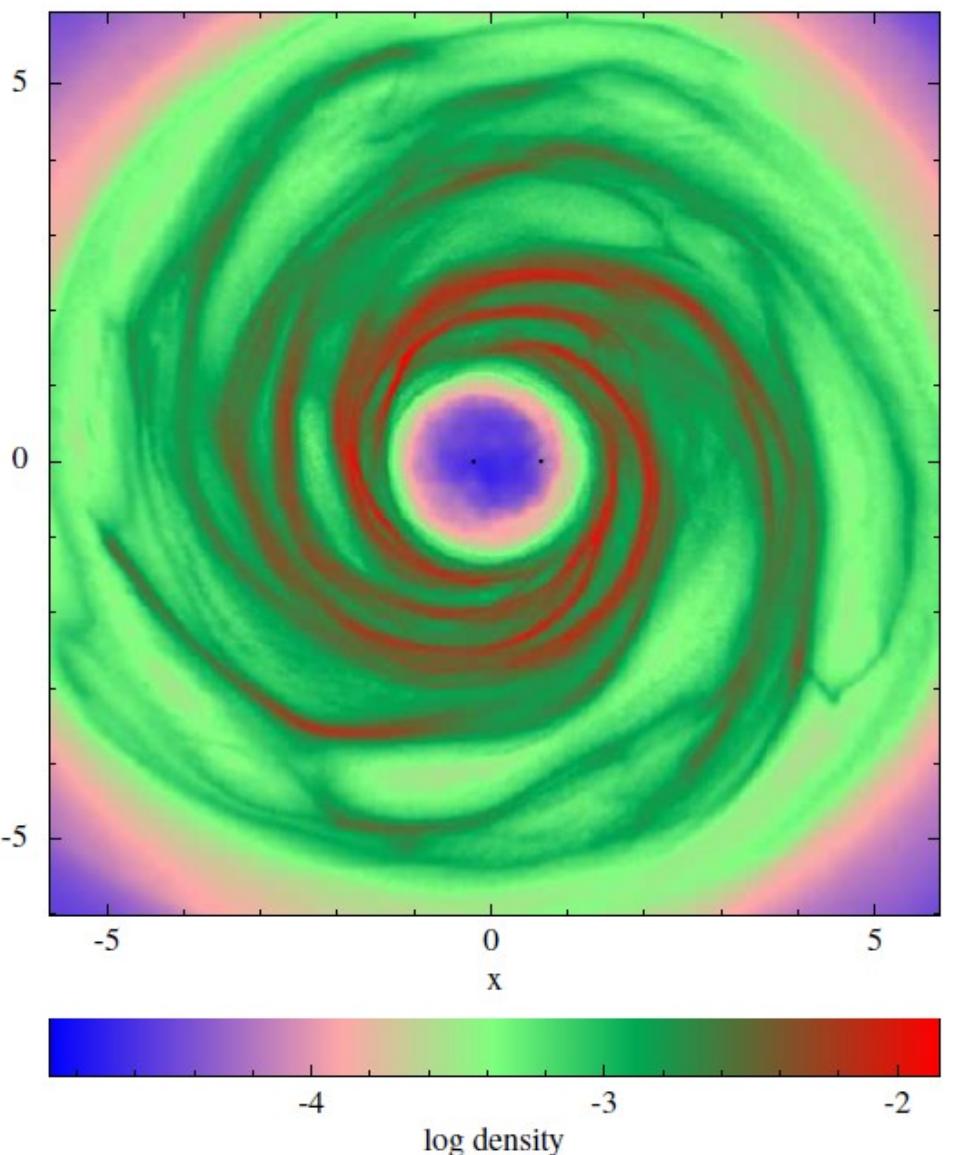
Retrograde versus prograde

Surface density self similar to a $(1+e)$

$$\Sigma \sim r^{-4/3}$$



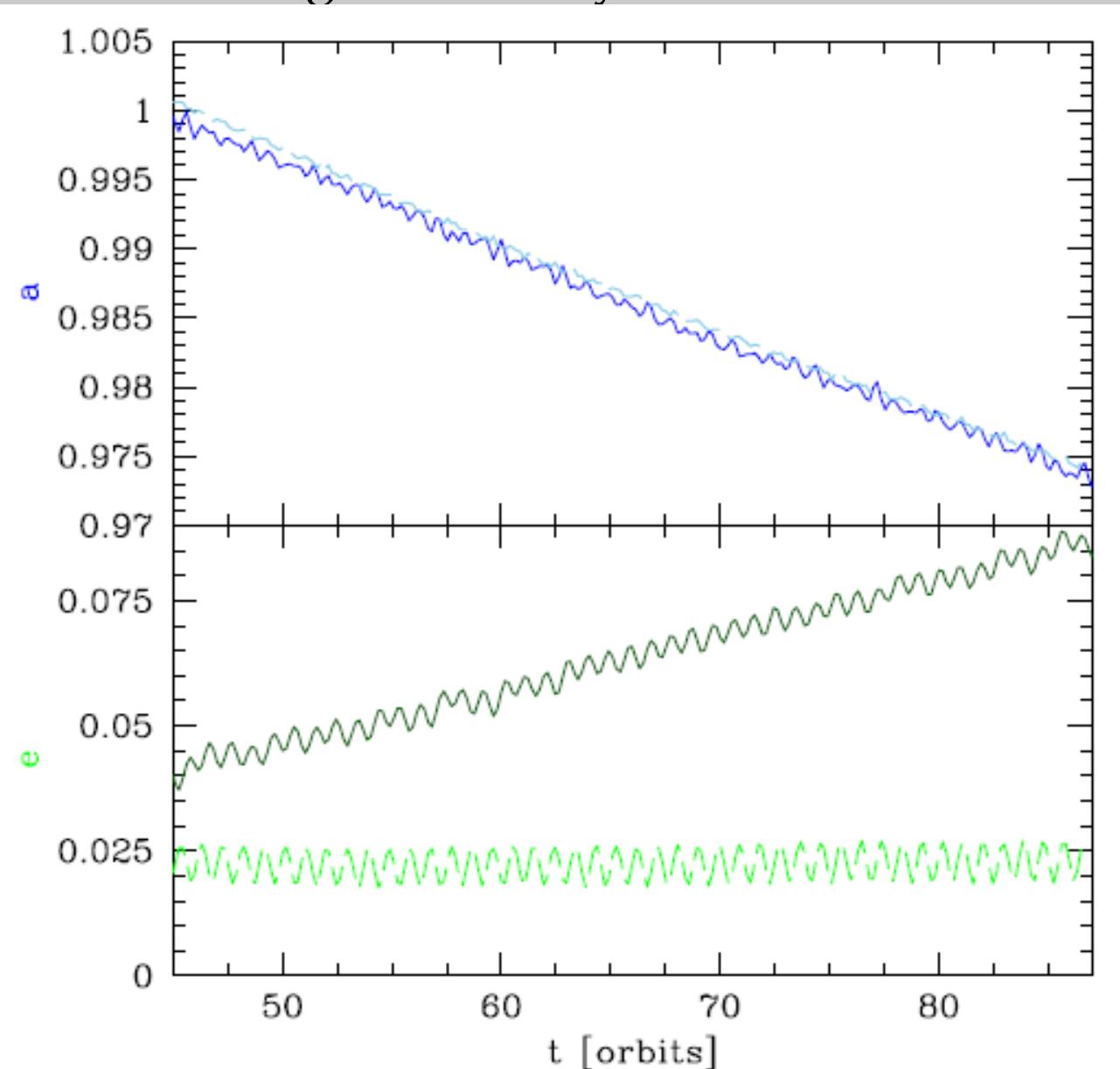
Retrograde versus prograde



Ro

Retrograde versus prograde: 1. Circular BHs

Shrinkage identical, retrograde binary remains circular

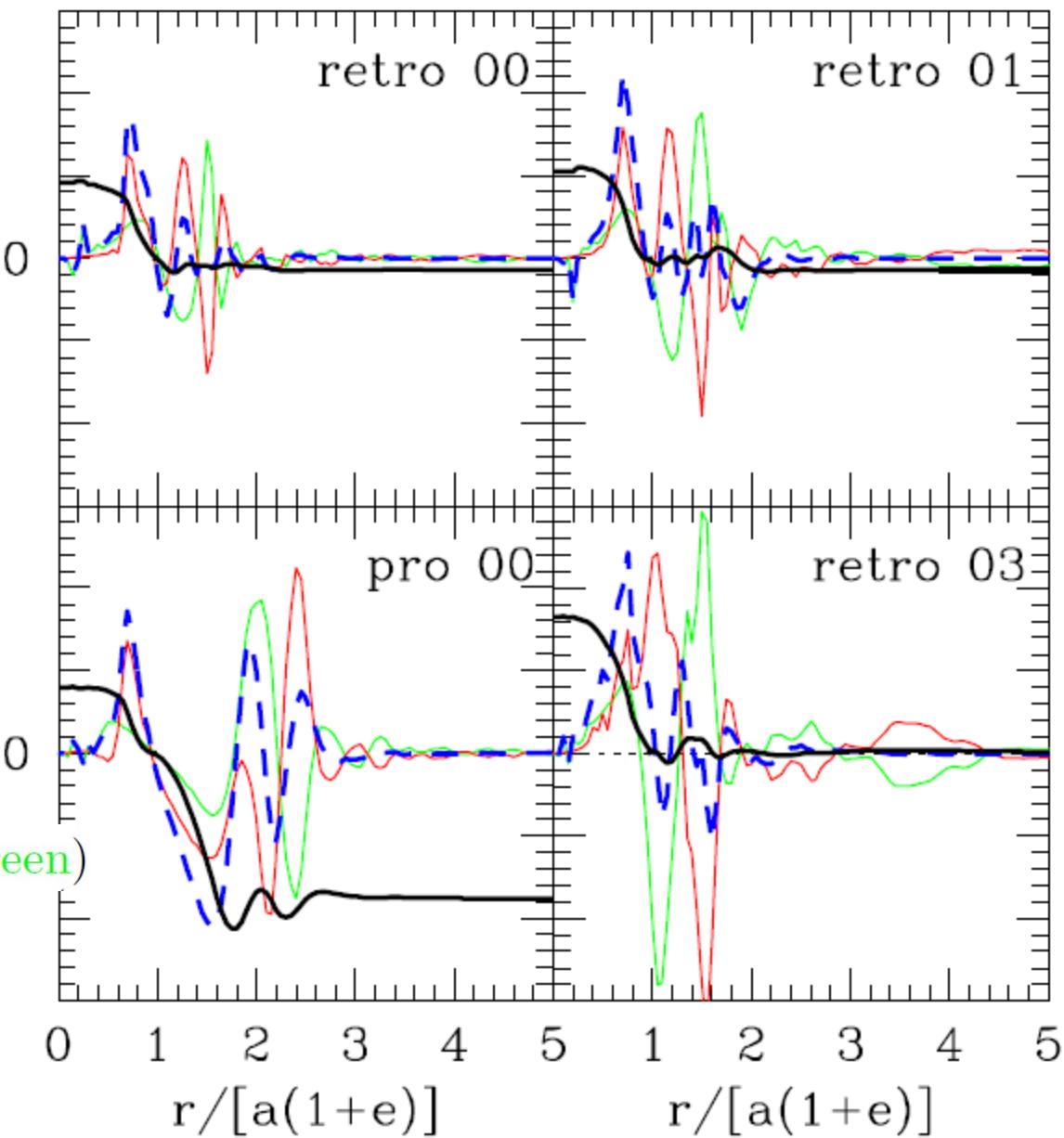


(Roedig &
Sesana submitted)

Retrograde torques: no stream-component

$$T_{\text{in}}(\epsilon) \approx T_{\text{in}}(\epsilon = 0) + 10 \cdot T_{\text{in}}(\epsilon = 0) \epsilon^2$$

differential torque on the primary (red)
differential torque on the secondary (green)
the sum of the two (dashed-blue)
integrated torque (black)

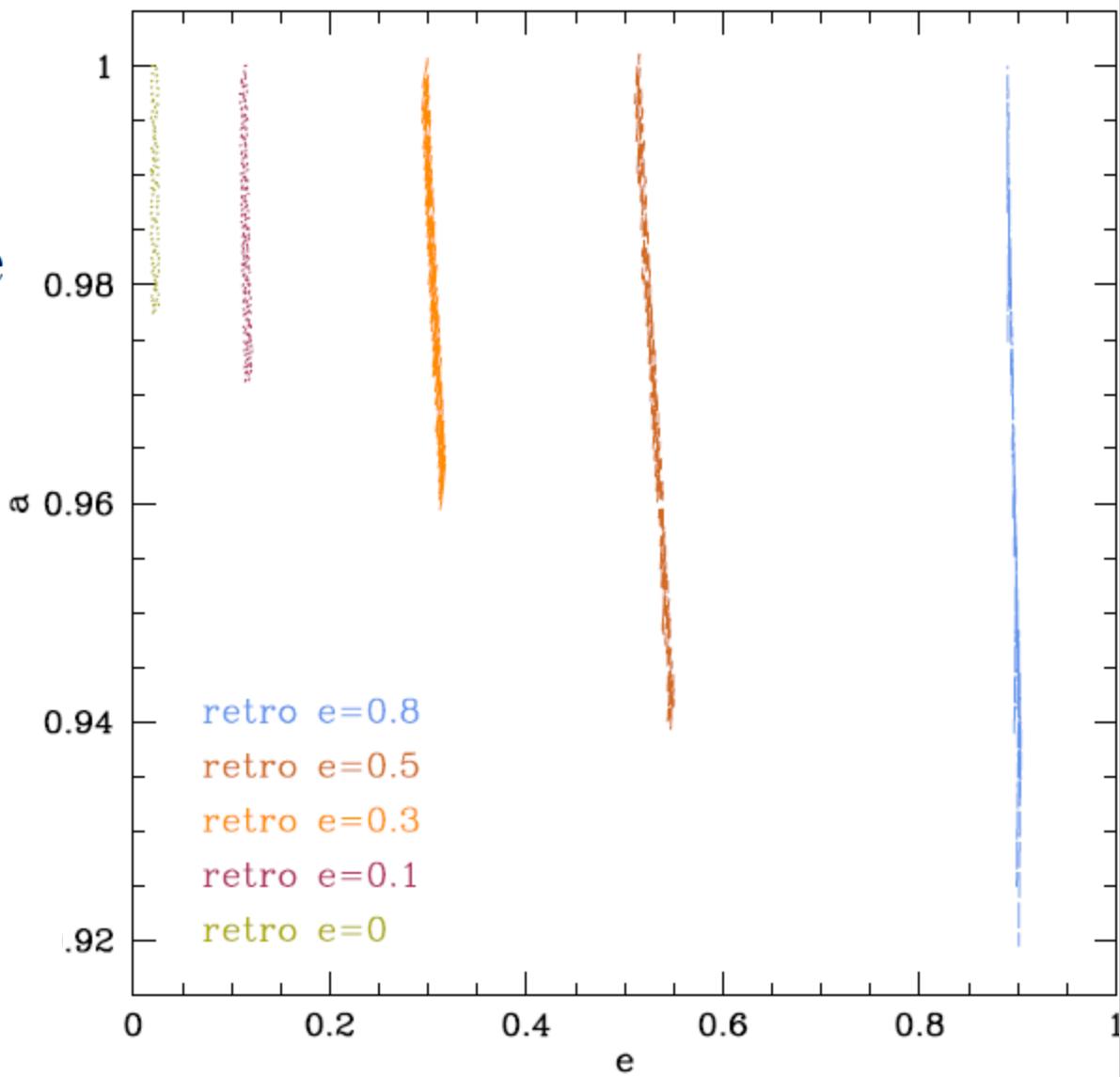


Retrograde decay exponential

$$\dot{e}(e) \approx -0.0034 + 0.09 e$$

$\dot{a}(e)$ is linear and negative

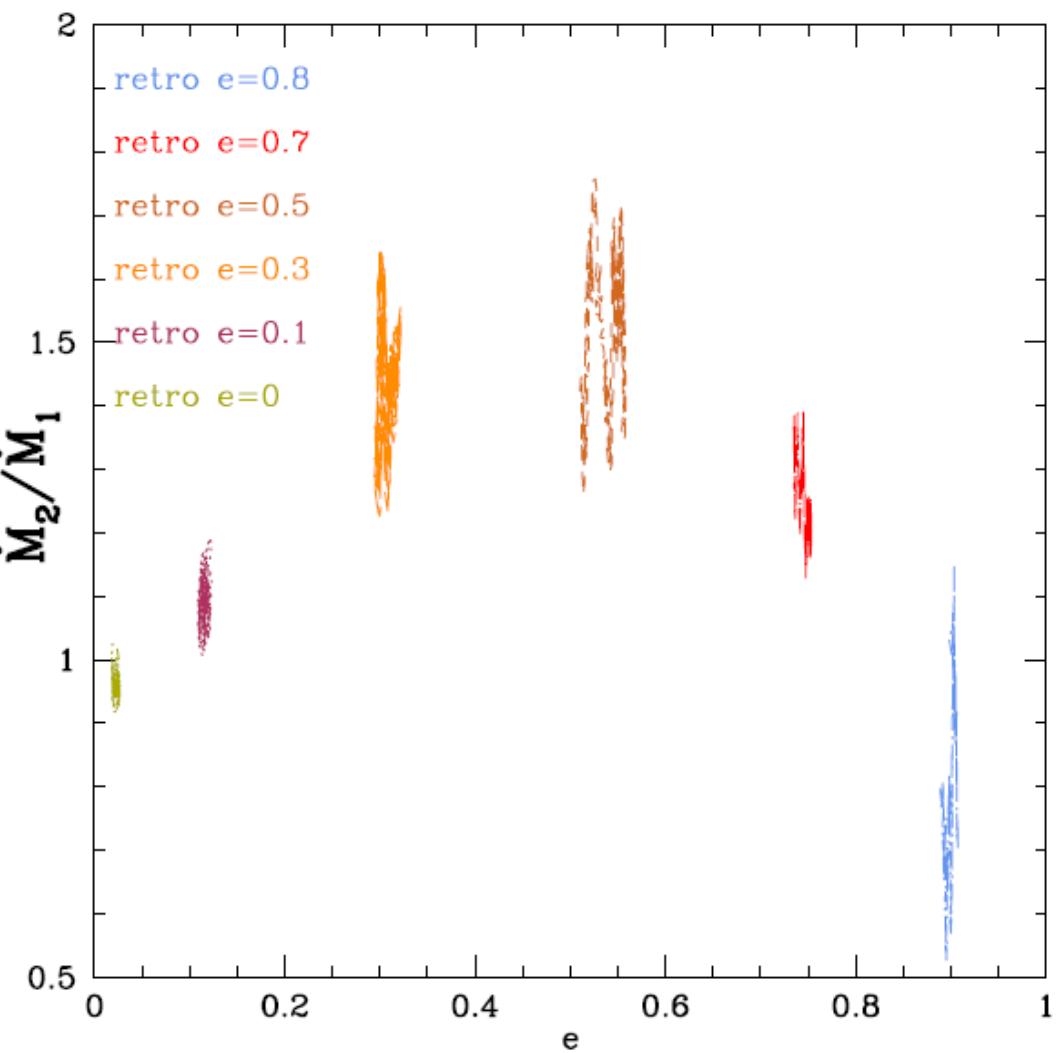
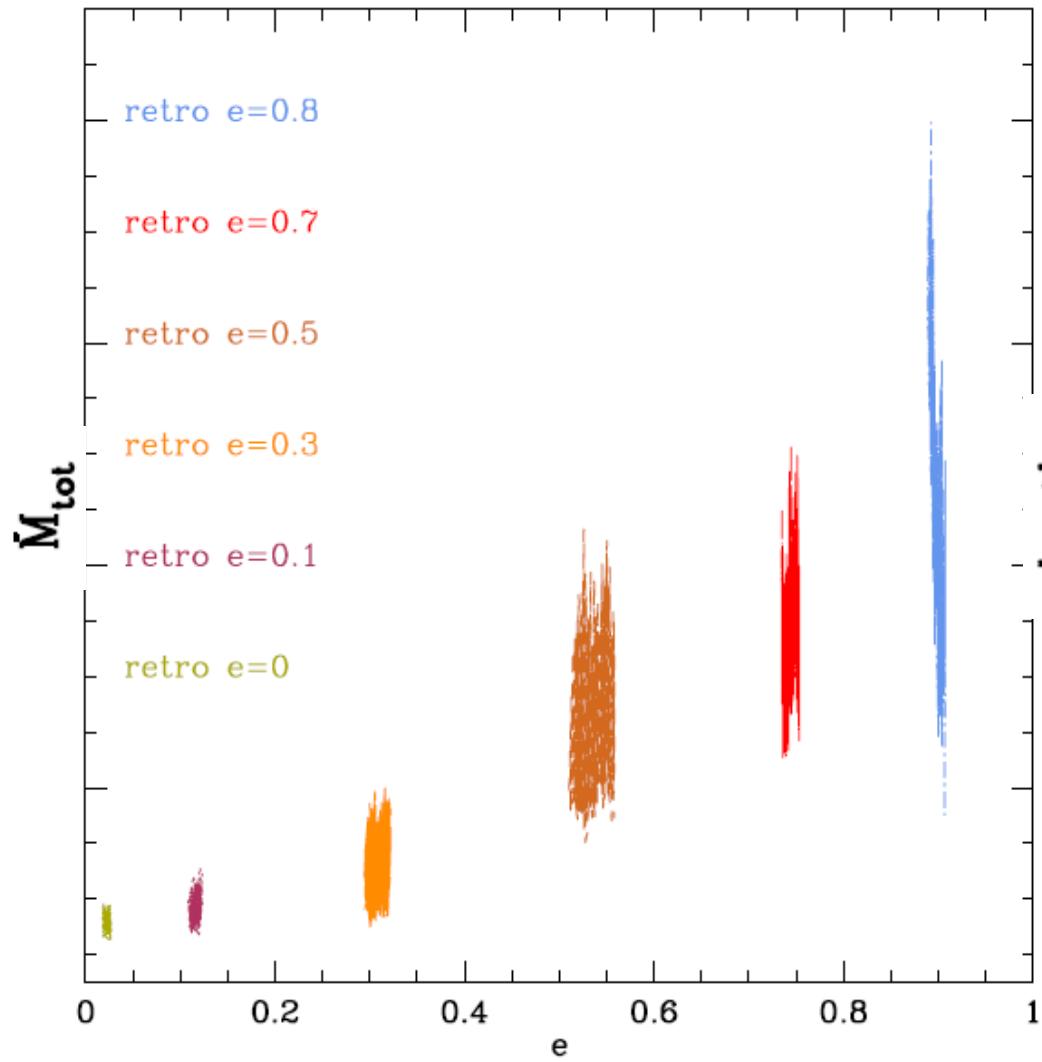
$$\frac{\dot{a}(t)}{a} = -\frac{2}{M_2}(\dot{\mu} + \dot{M}_{\text{BHB}})$$



Impact on both BHs once eccentricity high

Accretion rates and ratio depends on eccentricity

Turnover when primary Bondi-radius large at apocentre



Analytical understanding: M2-impact model

Purely kinematic interaction between secondary BH and gas at apocenter
with fraction α of accreted gas

Ignoring CoM motion

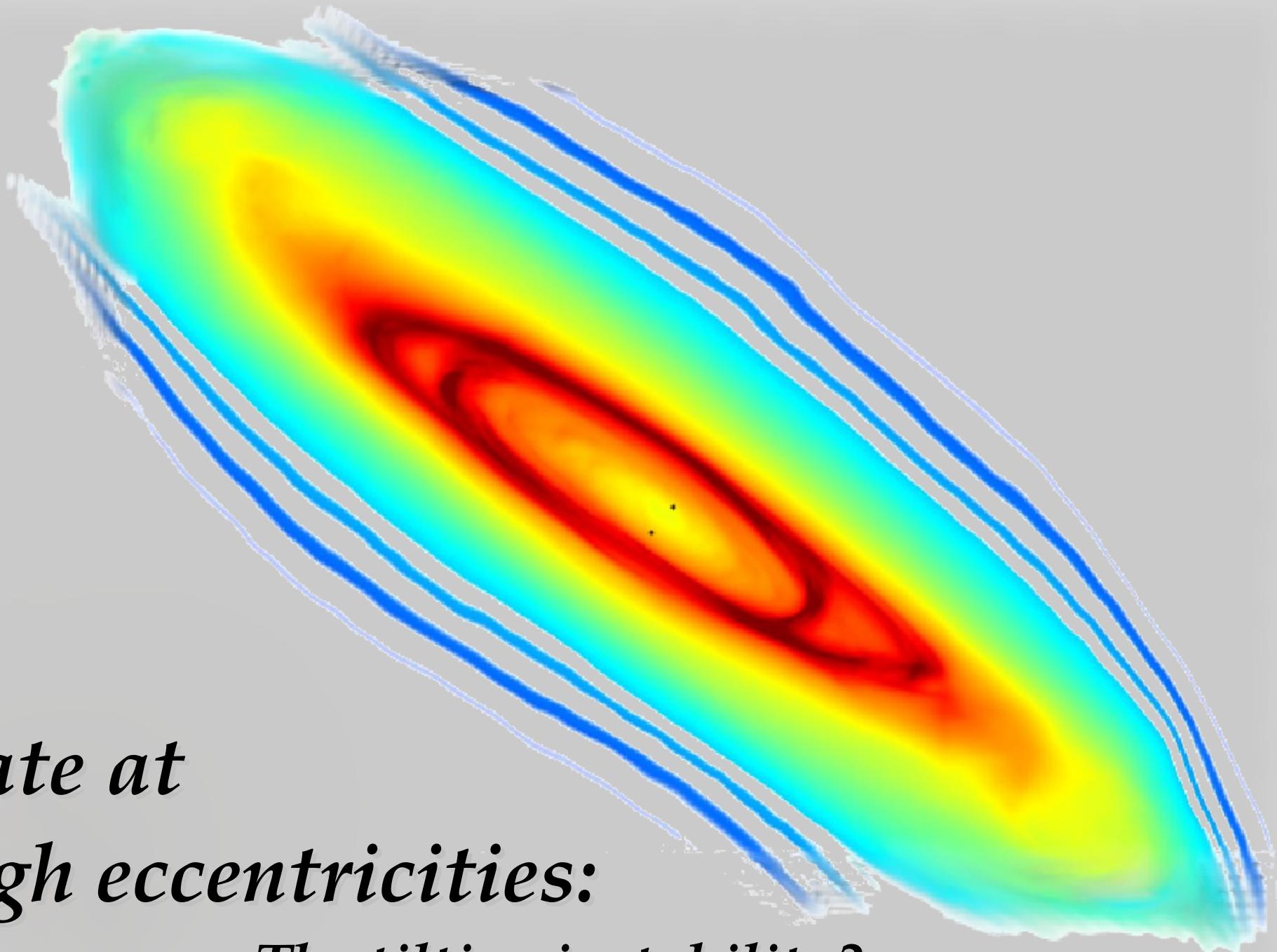
(compare with Nixon et al 2011)

$$\frac{\Delta \mathfrak{a}}{a} = -\frac{2\Delta M}{M_2(1 + \mathfrak{e})} \mathcal{F}(e, q, \alpha),$$

$$\Delta e = \frac{2\Delta M}{M_2} \mathcal{F}(e, q, \alpha),$$

$$\mathcal{F}(e, q, \alpha) = (1 - e)^{1/2} (1 + q)^{1/2} + \alpha \frac{1 - \mathfrak{e}}{1 + \mathfrak{q}}$$

*fate at
high eccentricities:
The tilting instability?*



Out-of-plane perturbations at high eccentricity

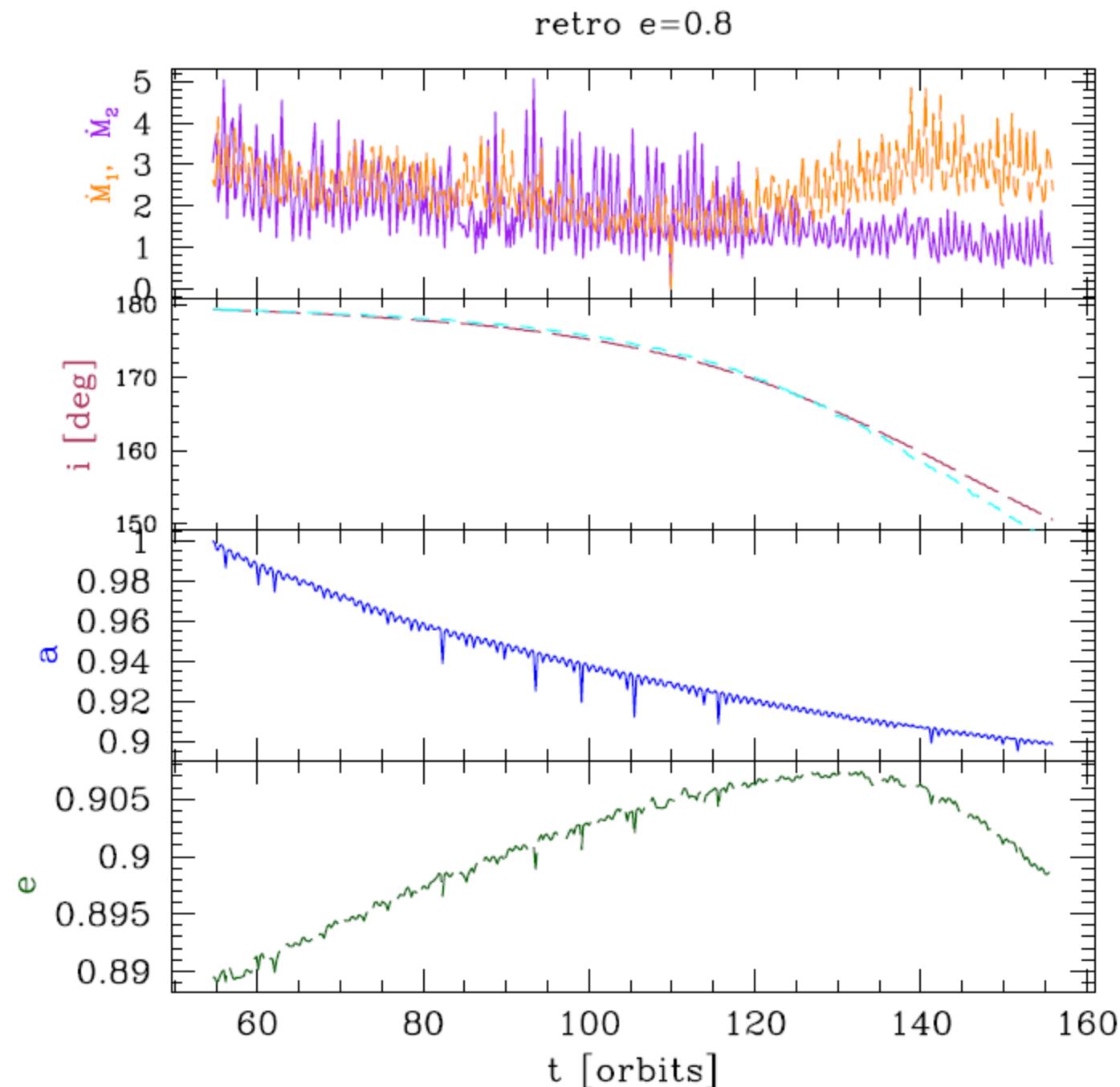
mass accretion rates
M1 and M2

$$i = \arccos \left(\frac{\langle \mathbf{L}_{\text{BHB}} | \mathbf{L}_{\text{Disc}} \rangle}{|\mathbf{L}_{\text{BHB}}| |\mathbf{L}_{\text{Disc}}|} \right)$$

Semi-major axis a

Eccentricity e

Constanze I



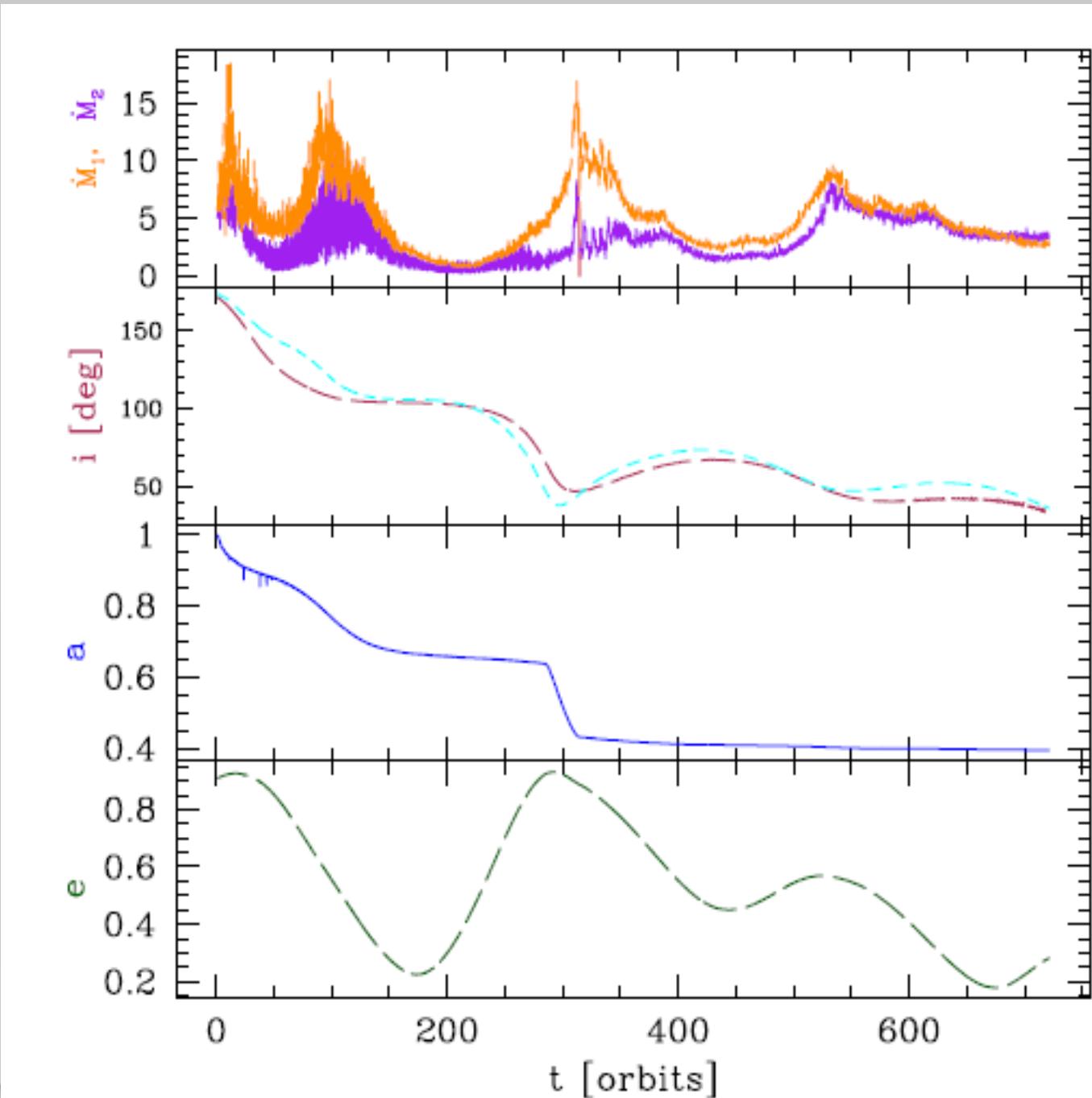
Kozai-like exchange of inclination & eccentricity

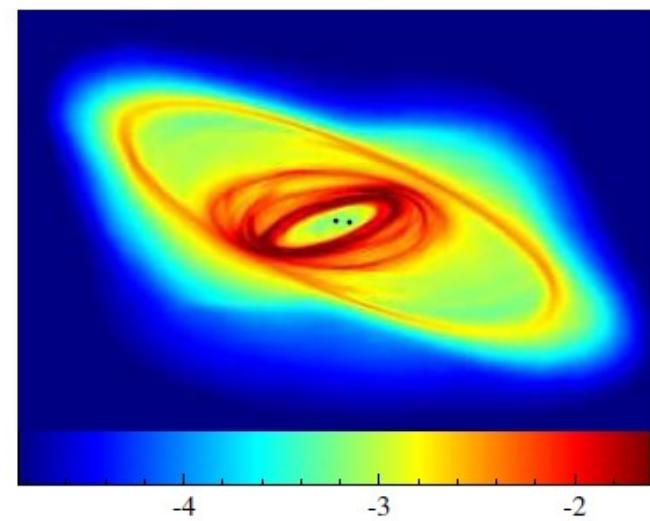
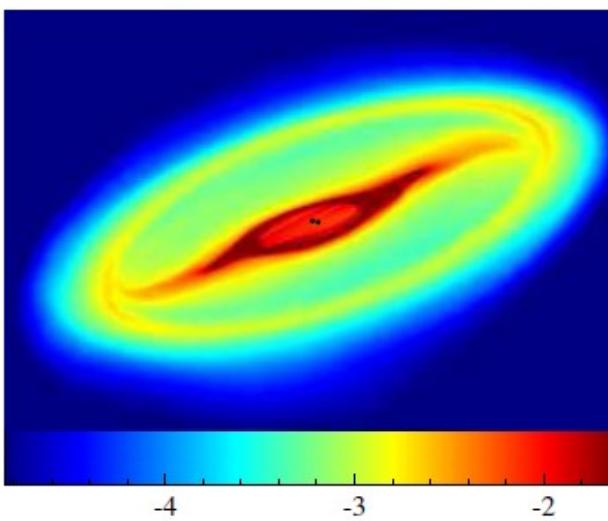
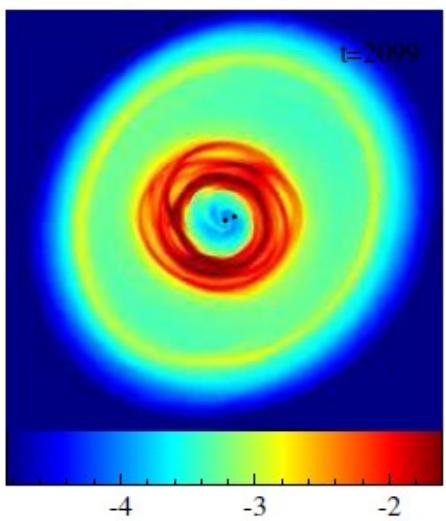
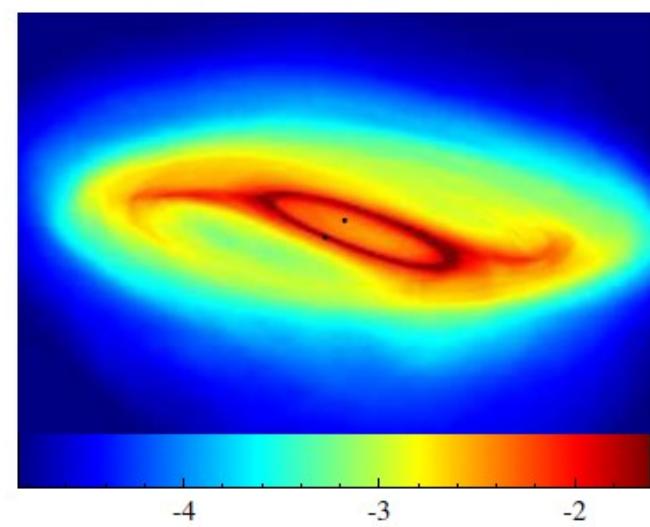
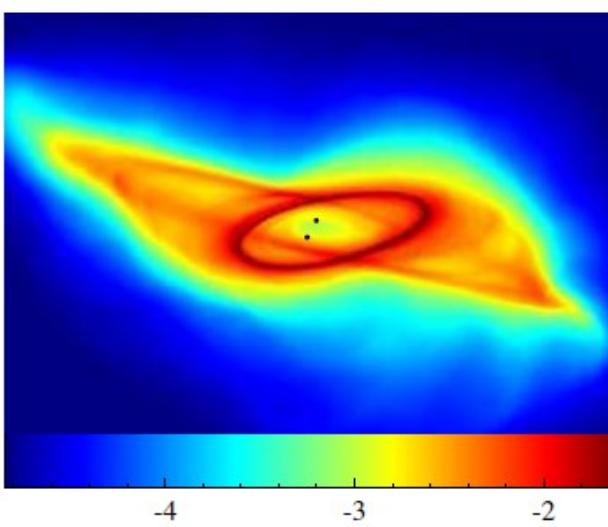
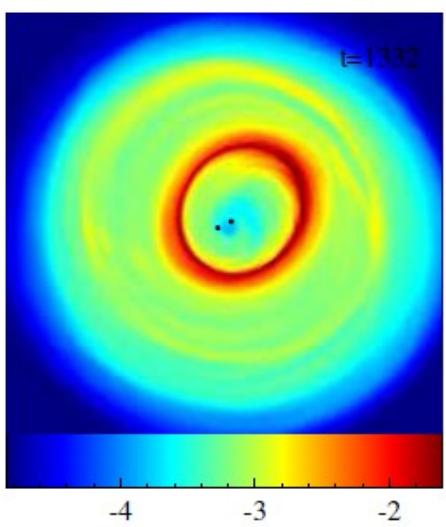
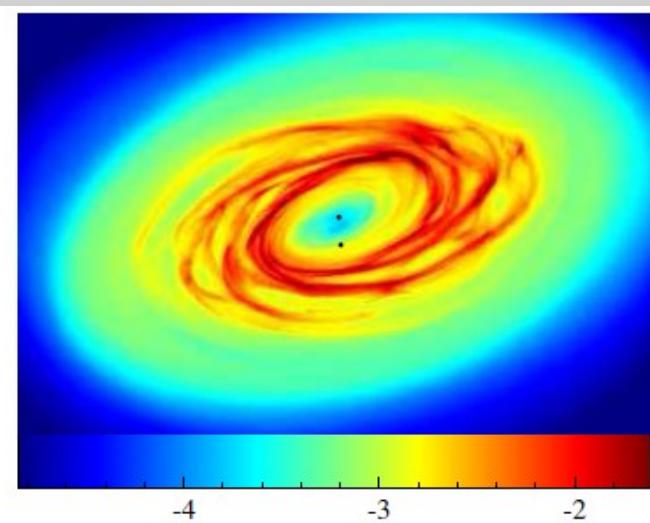
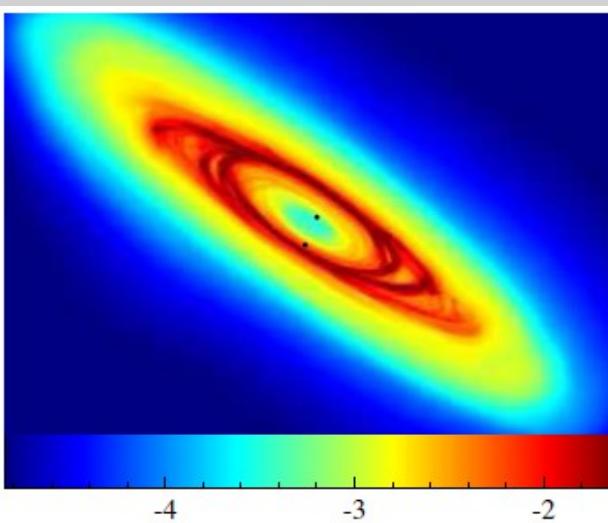
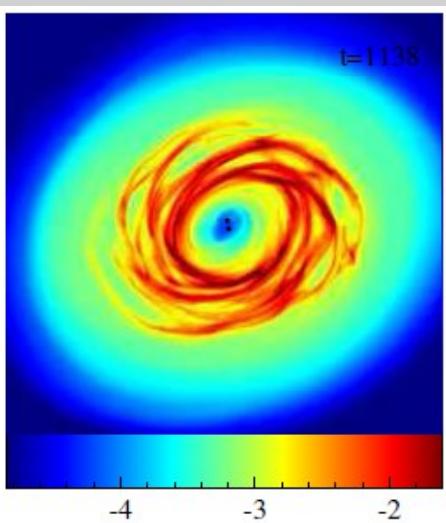
mass accretion rates
 M_1 and M_2

Inclination i

Semi-major axis a

Eccentricity e





Summary:

- Circular retrograde discs shrink similar to prograde ones
- They remain circular
- If eccentricity excited, retrograde BHB shrinks exponentially
- Torques \sim Accretion rates \sim eccentricity
- Ballistic impact-model
- Ratio of accretion rates M_2/M_1 turnover at $e \sim 0.45$
- Increasing movement of BHB-Centre of Mass
- At high ecc: stochastic perturbations from clumpy disc destabilize binary : it leaves the plane
- Kozai like exchange of inclination and eccentricity
- Final stage: prograde and eccentric
OR: plunge-like merger if GW important

Thank you for your attention !

Special thanks to:

Kareem Sorathia, Sarah Buchanan & Rowdy

