

# THE EFFECTS OF DIFFUSIVE SHOCK ACCELERATION ON THE EMITTED THERMAL X-RAY SPECTRUM IN SNR SHOCKS

DAN PATNAUDE  
SAO AND KITP

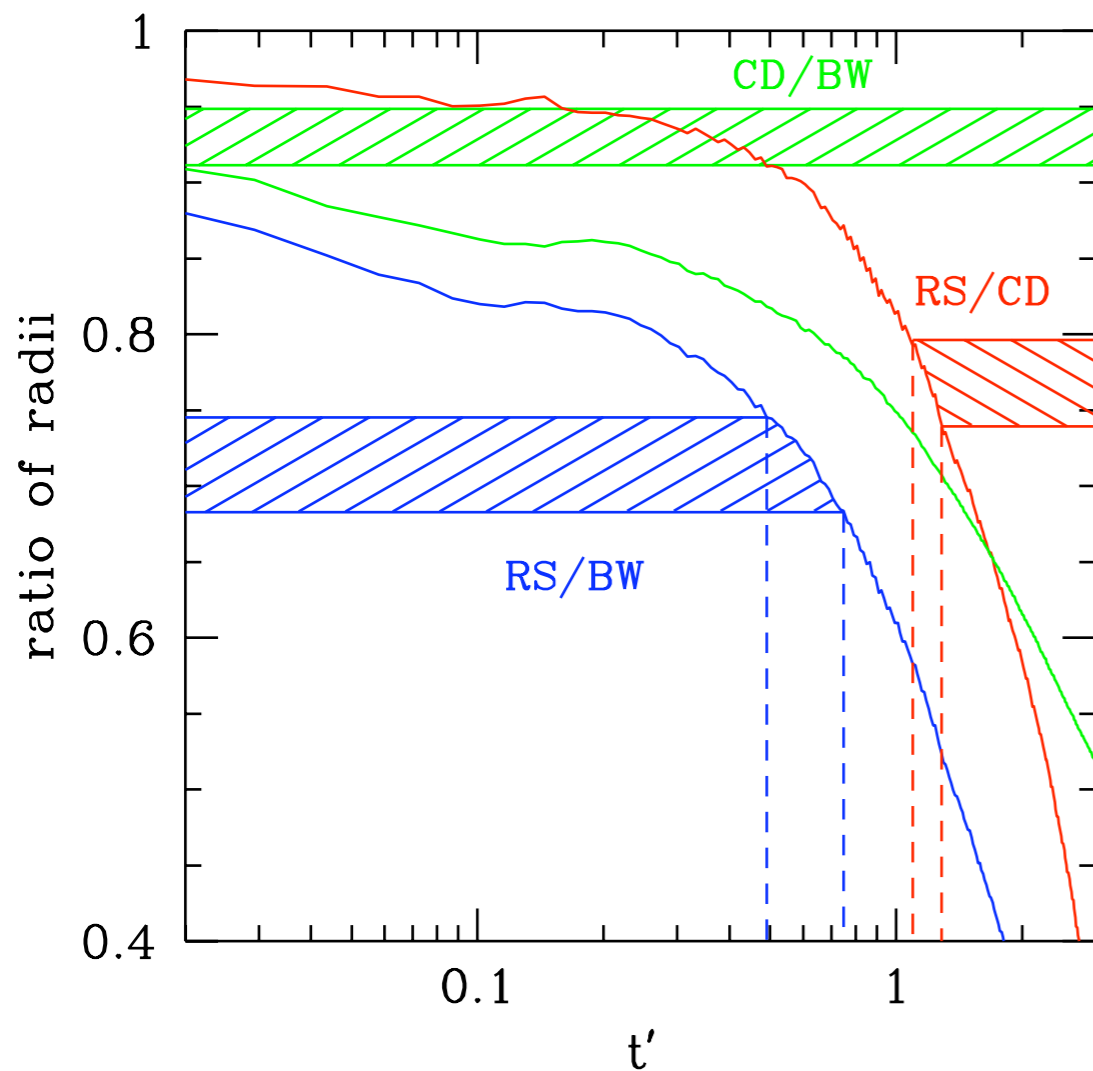
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# DSA: A BRIEF OVERVIEW

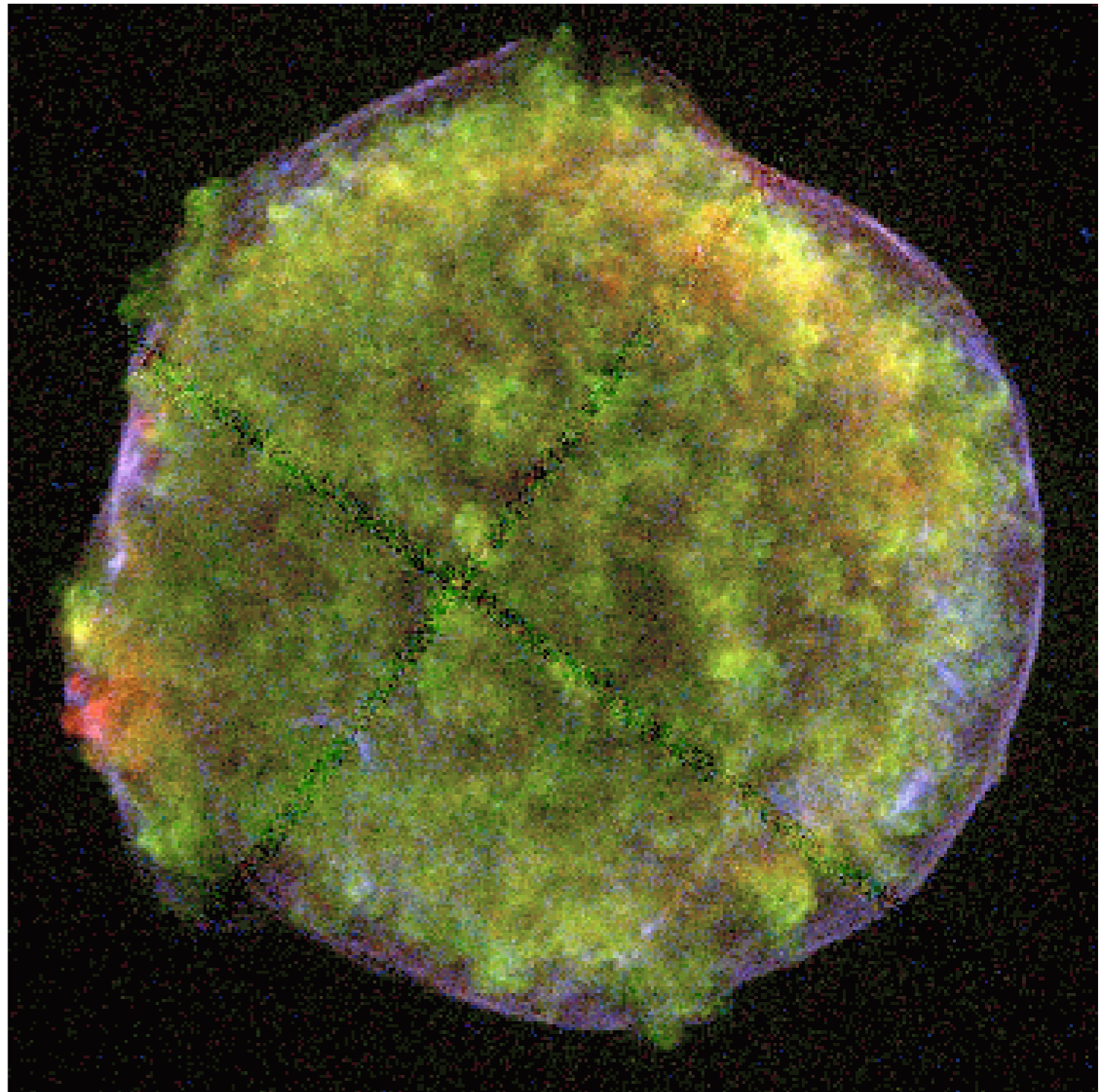
- Efficient diffusive shock acceleration lowers the shock temperature and raises the postshock density (Jones & Ellison, 1991; Berezhko & Ellison, 1999)
- The nonequilibrium ionization is dependent upon both the shock temperature and the shock density through their relation to the electron temperature,  $T_e$ , and electron density,  $n_e$
- A number of Galactic SNRs show both thermal and nonthermal emission behind the forward shock, including SN1006 (Vink et al. 2003; Bamba et al. 2008) and Tycho (Hwang et al. 2002; Cassam-Chenaï et al. 2007)
- In SNR RX J 1713.7-3946, the lack of thermal X-ray emission is an important constraint on the ambient density and significantly impacts models for TeV emission (Slane et al. 1999; Ellison et al. 2001, Aharonian et al. 2007; Katz & Waxman 2008)

# EVIDENCE FOR EFFICIENT DSA IN SNRs

Efficient shock acceleration softens the EOS in the shocked gas. In Tycho, the location of the blastwave suggests that it has been modified considerably by cosmic ray acceleration.



Time evolution of Tycho's SNR (Warren et al. 2005)



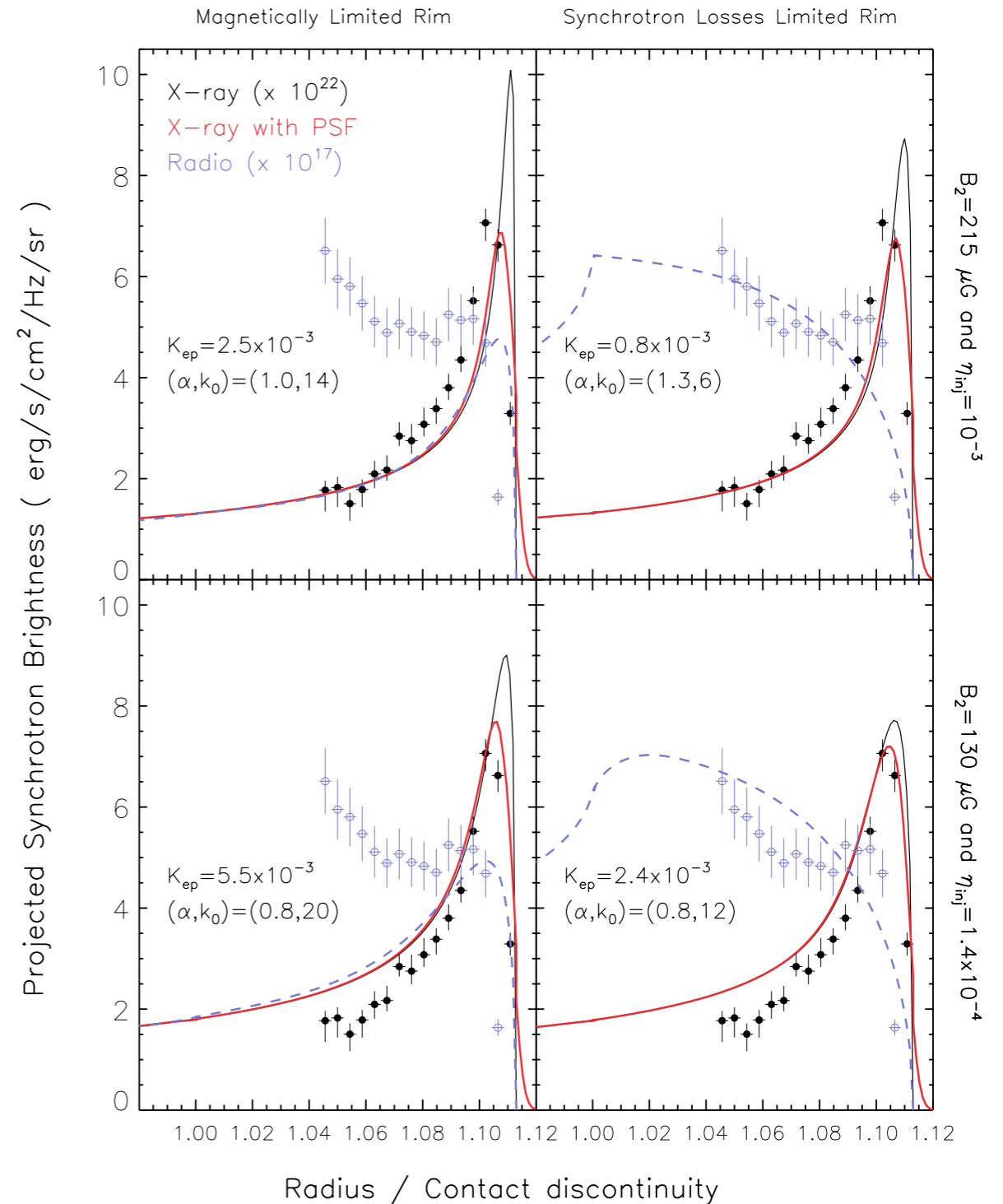
Tycho viewed in X-rays (Warren et al. 2005)

# EVIDENCE FOR EFFICIENT DSA IN SNRs

- Thin synchrotron rims:

The radial profile of the X-ray bright synchrotron rims in Tycho can be explained by models for amplified magnetic fields at the shock front and acceleration of electrons to TeV energies.

Additionally, the synchrotron dominated rims can be used to constrain the ambient medium density to be  $0.6 \text{ cm}^{-3}$ .

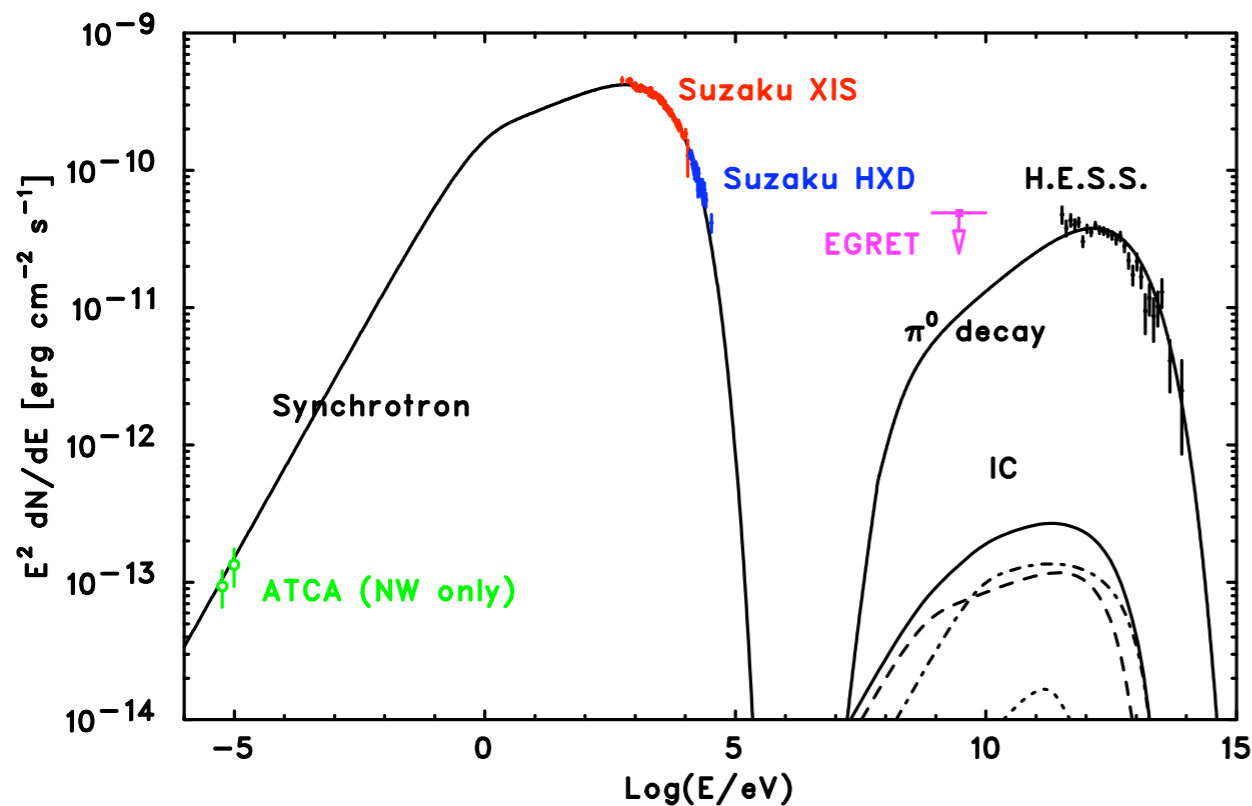
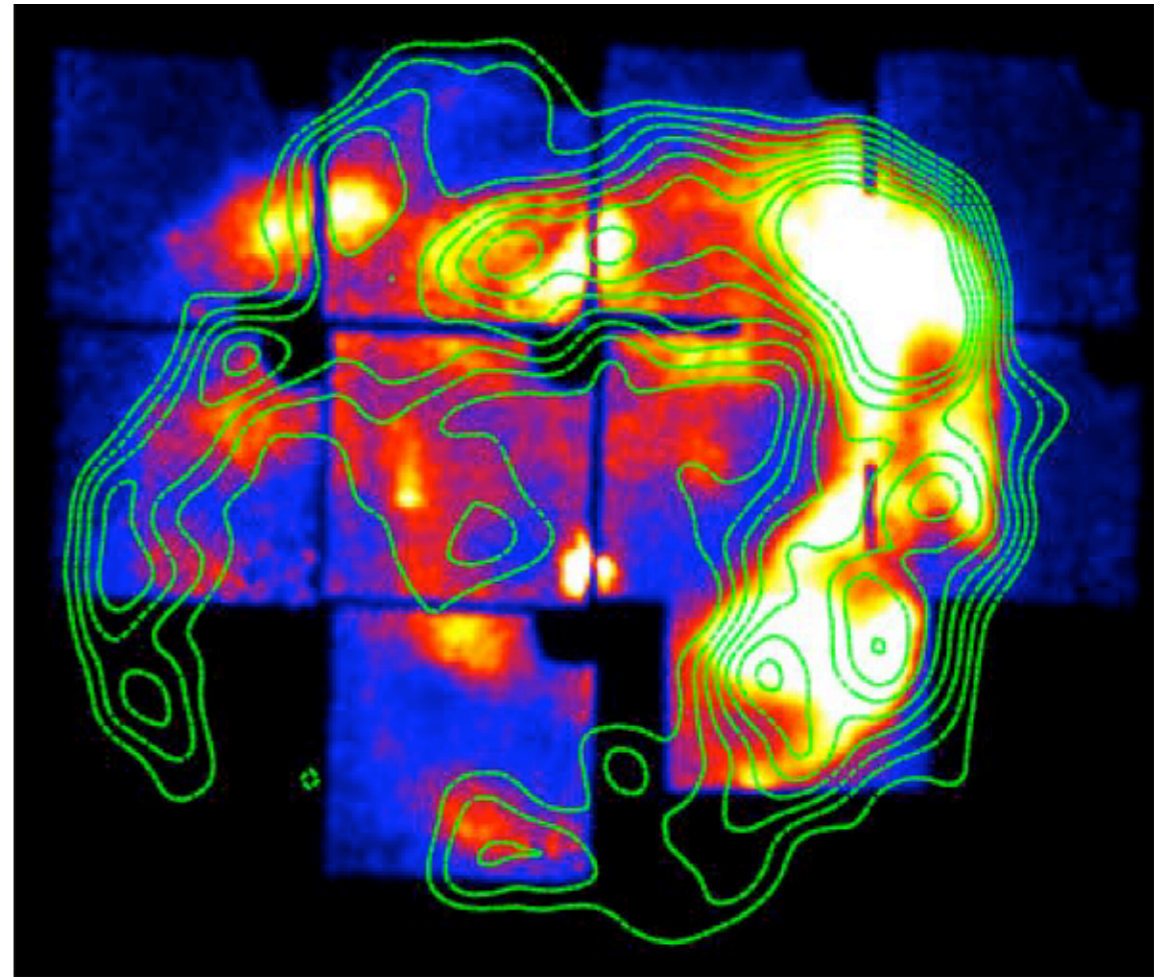


Line of sight projections of radio and X-ray rims with a varying spectral cutoff ( $\alpha$ ) (Cassam-Chenaï et al. 2007).

# EVIDENCE FOR EFFICIENT DSA IN SNRs

TeV  $\gamma$ -ray emission:

HESS detections of TeV gamma rays provides direct evidence for efficient acceleration of particles. However, the origin remains an open question

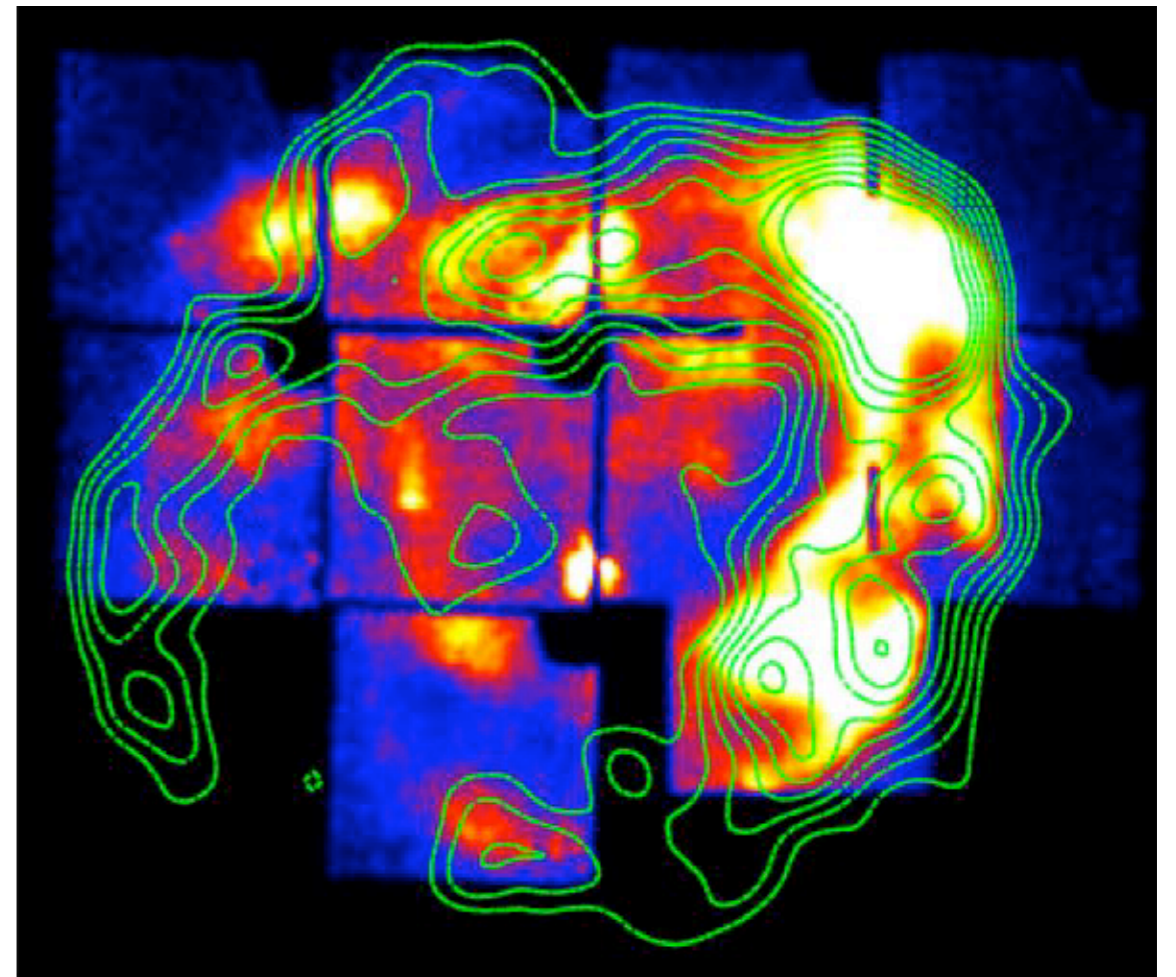
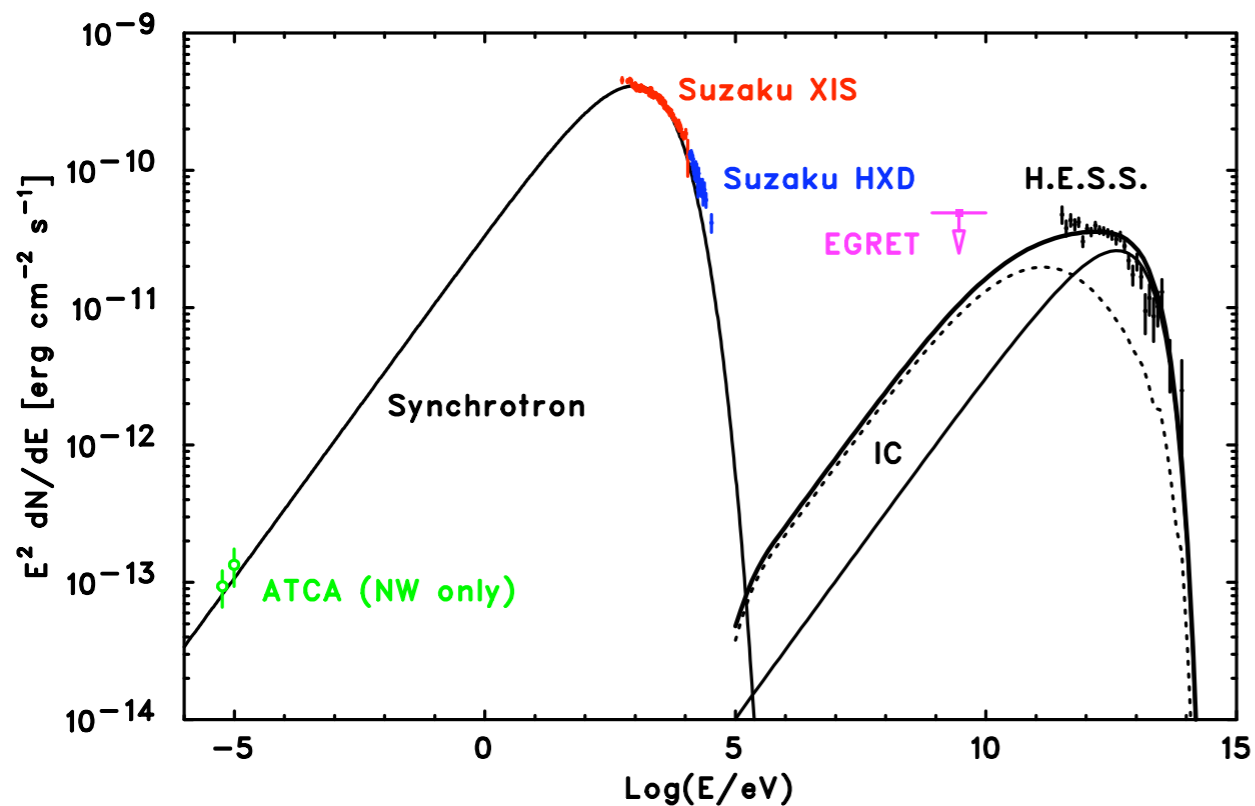


Suzaku XIS image and H.E.S.S. gamma-ray image (contours) of RX J1713-3946. (Left): Broadband SED assuming a hadronic or leptonic origin to the TeV emission (Tanaka et al. 2009).

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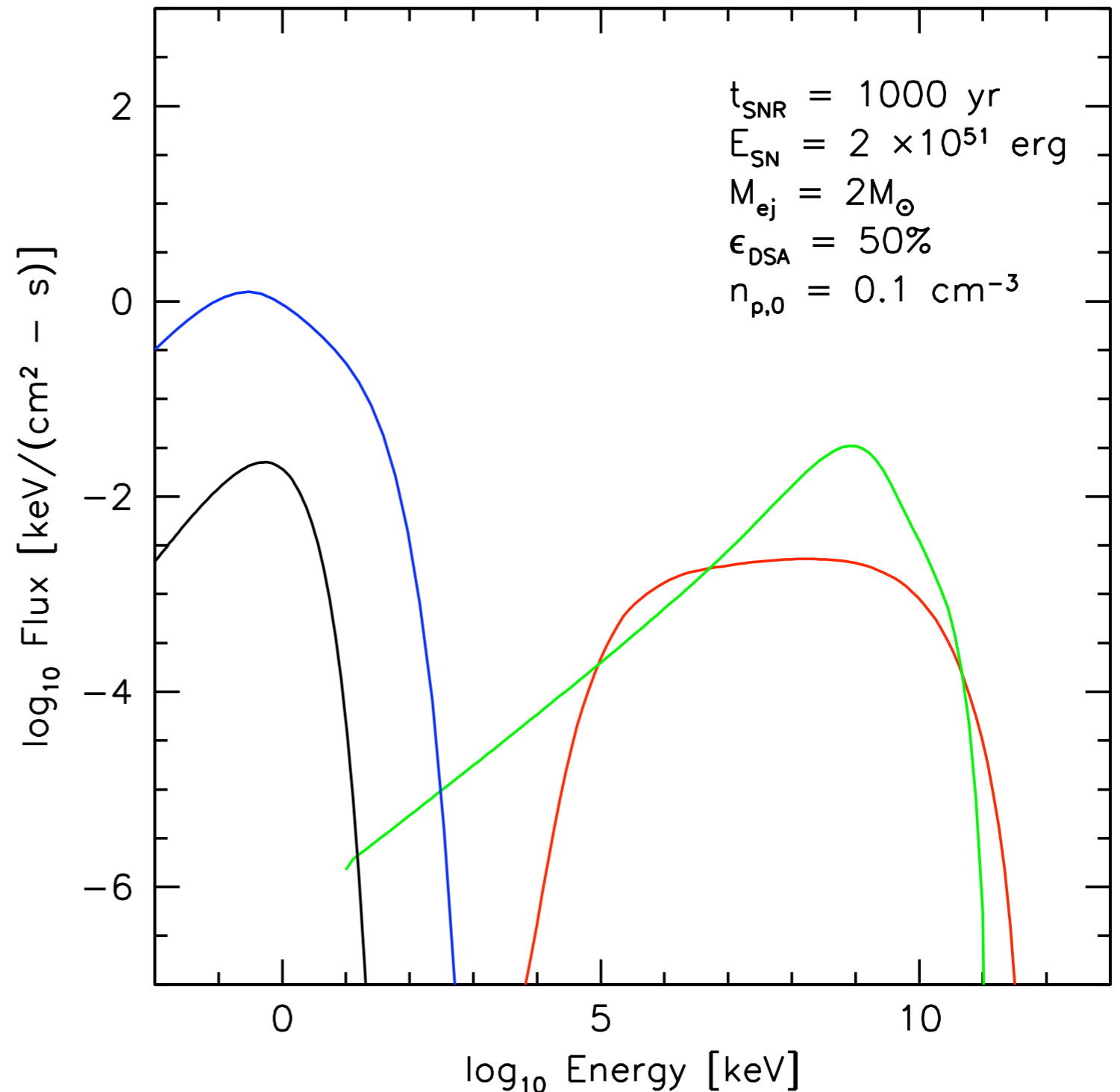
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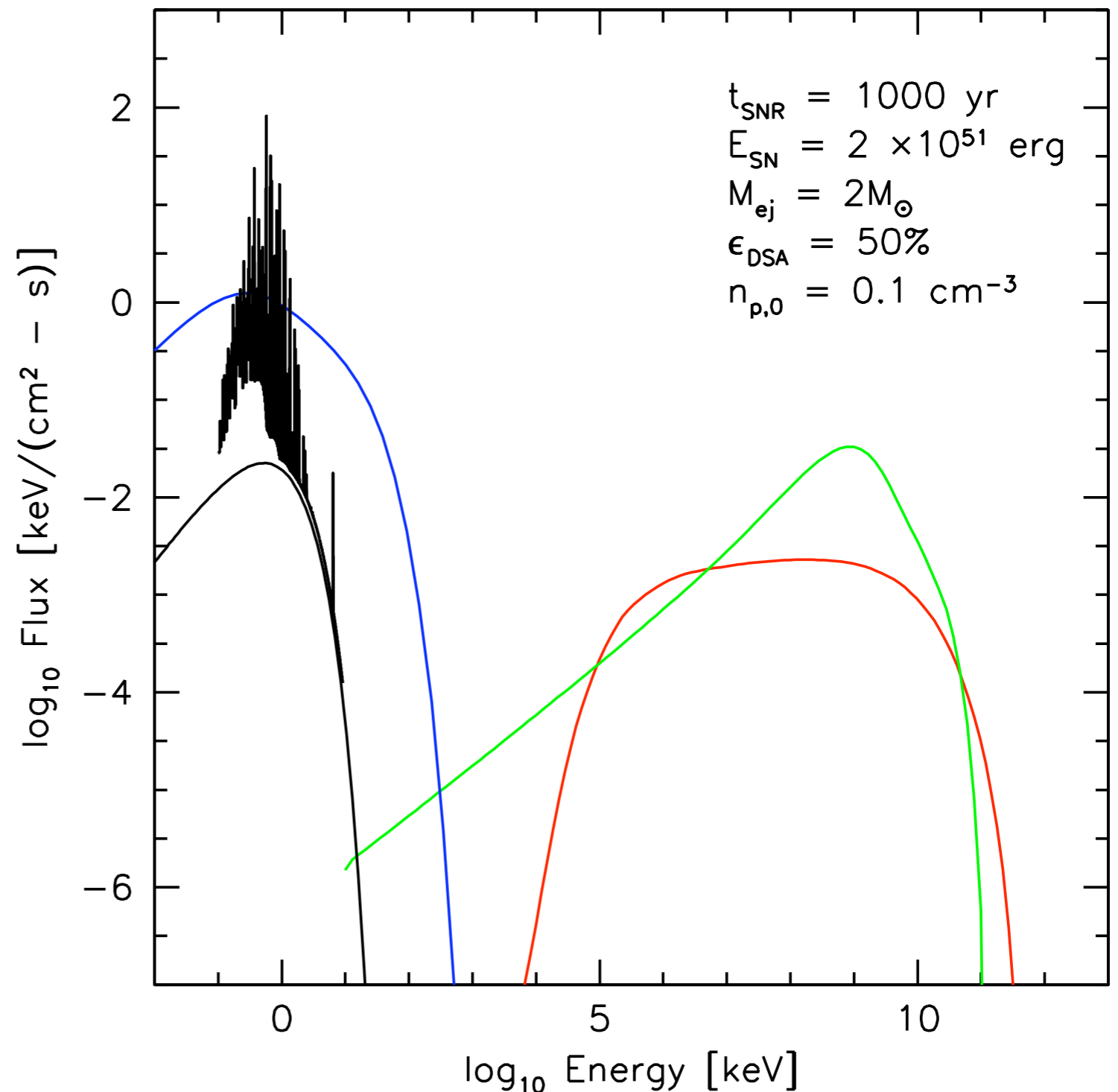
# The importance of including X-ray line emission in our simulations

Often, the lack of thermal emission is used to set the ambient medium density and place constraints upon the source of the TeV emission.



# The importance of including X-ray line emission in our simulations

the problem of course is that the X-ray line emission can extend well above the continuum!

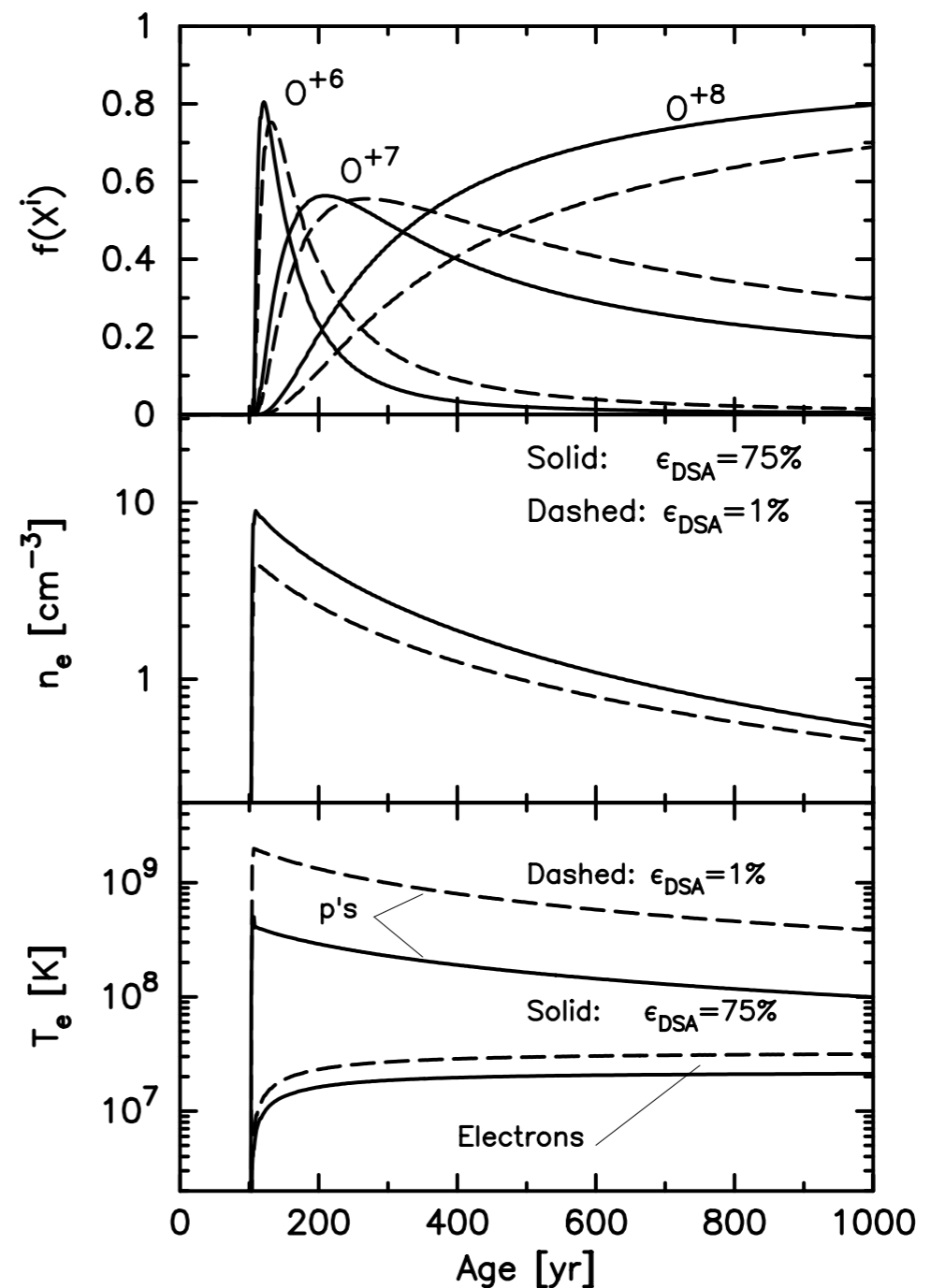




Existing evidence suggests that in order to understand the morphology of supernova remnants and other exotic astrophysical objects, we need models that self-consistently calculate the hydrodynamics, diffusive shock acceleration, and nonequilibrium ionization in shocks

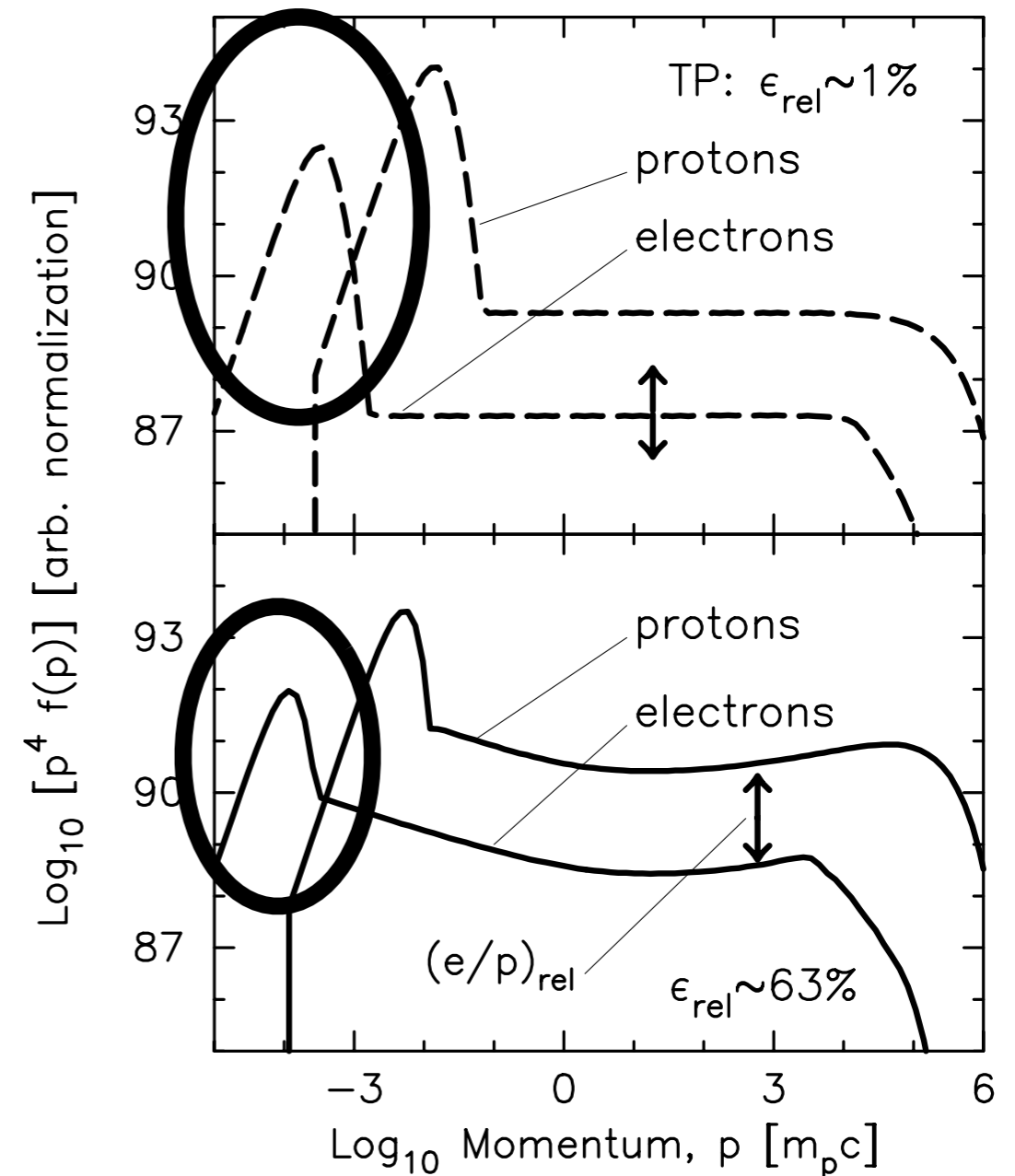
# OUR NEW MODEL: CR+HYDRO+NEI

- Uses semi-analytic model for DSA of Amato & Blasi (2005) and Blasi (2005)
- The ionization of the shock heated gas at a distance behind the shock is determined by  $n_e$  and  $T_e$
- Ionization structure is determined by solving the collisional ionization equations in a Lagrangian gas element
- $T_e$  is calculated by assuming heating via Coulomb collisions, but more efficient heating is considered

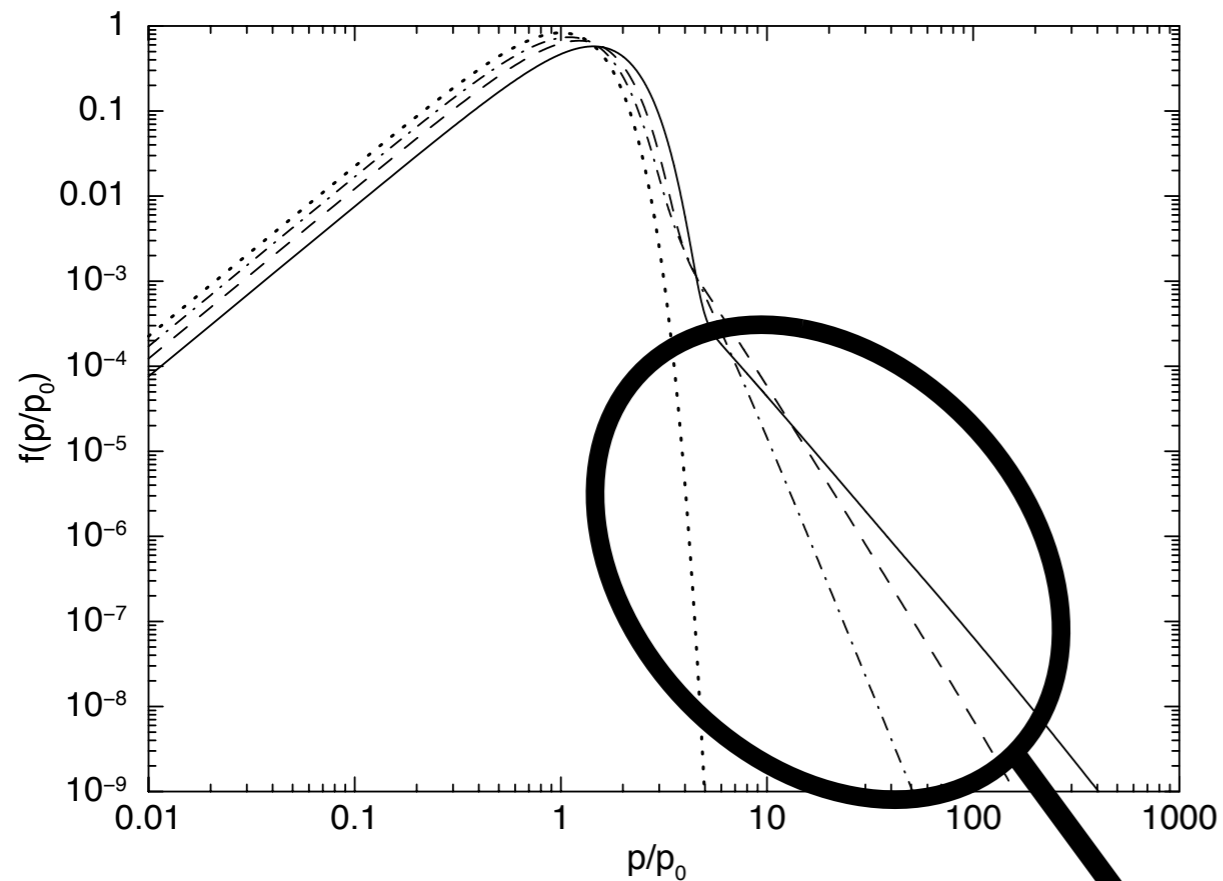


Time evolution of a spherically symmetric Lagrangian mass shell which is crossed by the forward shock at 100 yr. The CSM proton number density for this example is  $n_{p,0} = 1 \text{ cm}^{-3}$ . Here, and in all other examples, the unshocked CSM temperature is  $T_0 = 10^4 \text{ K}$ , and the unshocked magnetic field is  $B = 15 \mu\text{G}$ .

- Efficient shock acceleration produces a significant nonthermal particle population
- Our model, however, only treats ionization from the thermal population
- Using a Hybrid model (thermal + powerlaw tail) for the electron distribution, Porquet et al (2001) showed that nonthermal effects can alter the ionization balance
- The effect is much less pronounced in ionizing plasmas

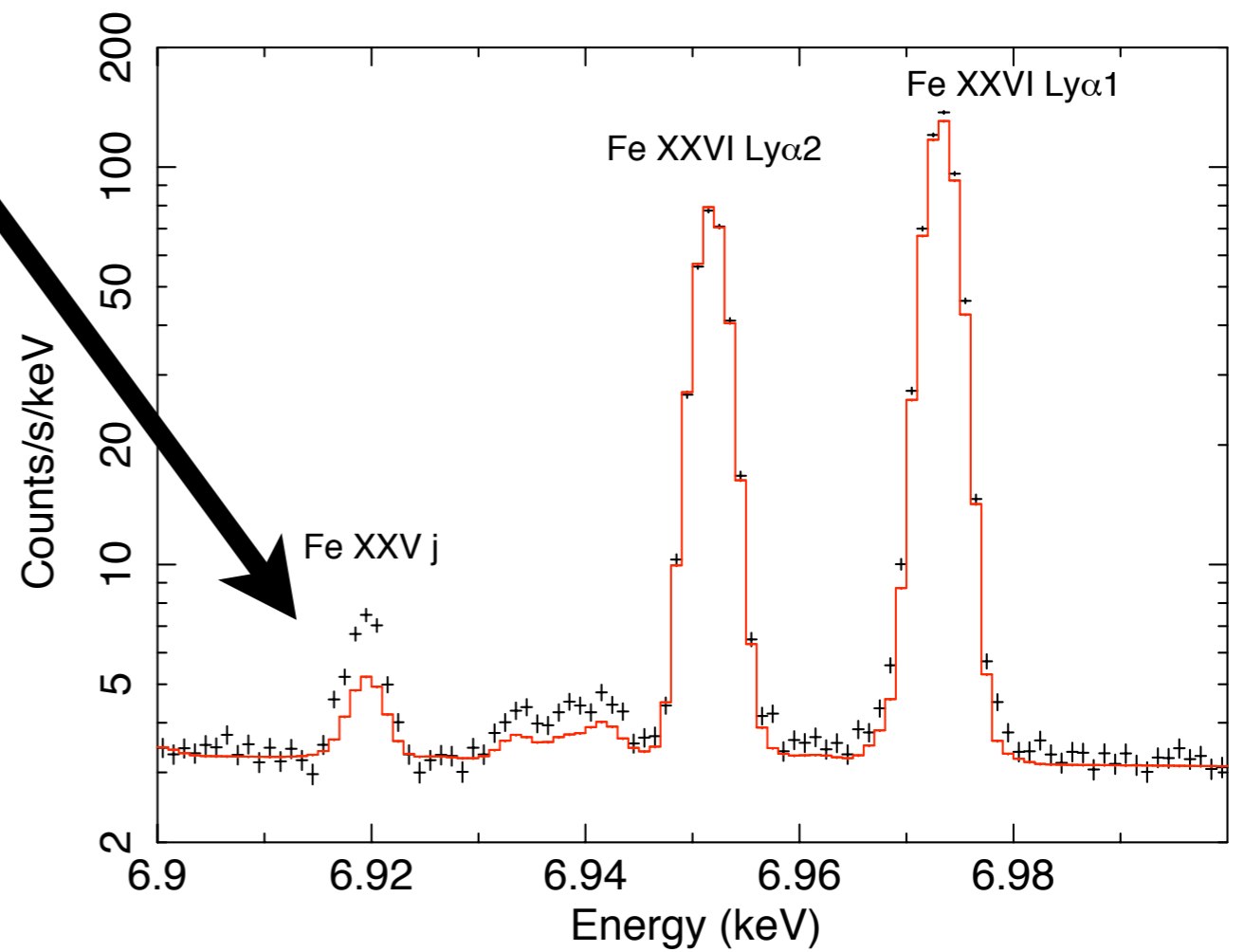


Electron and proton spectra,  $p^4 f(p)$ , for TP and efficient DSA. The up/down arrow indicates the normalization of the powerlaw component of the distribution (Ellison et al. 2007).



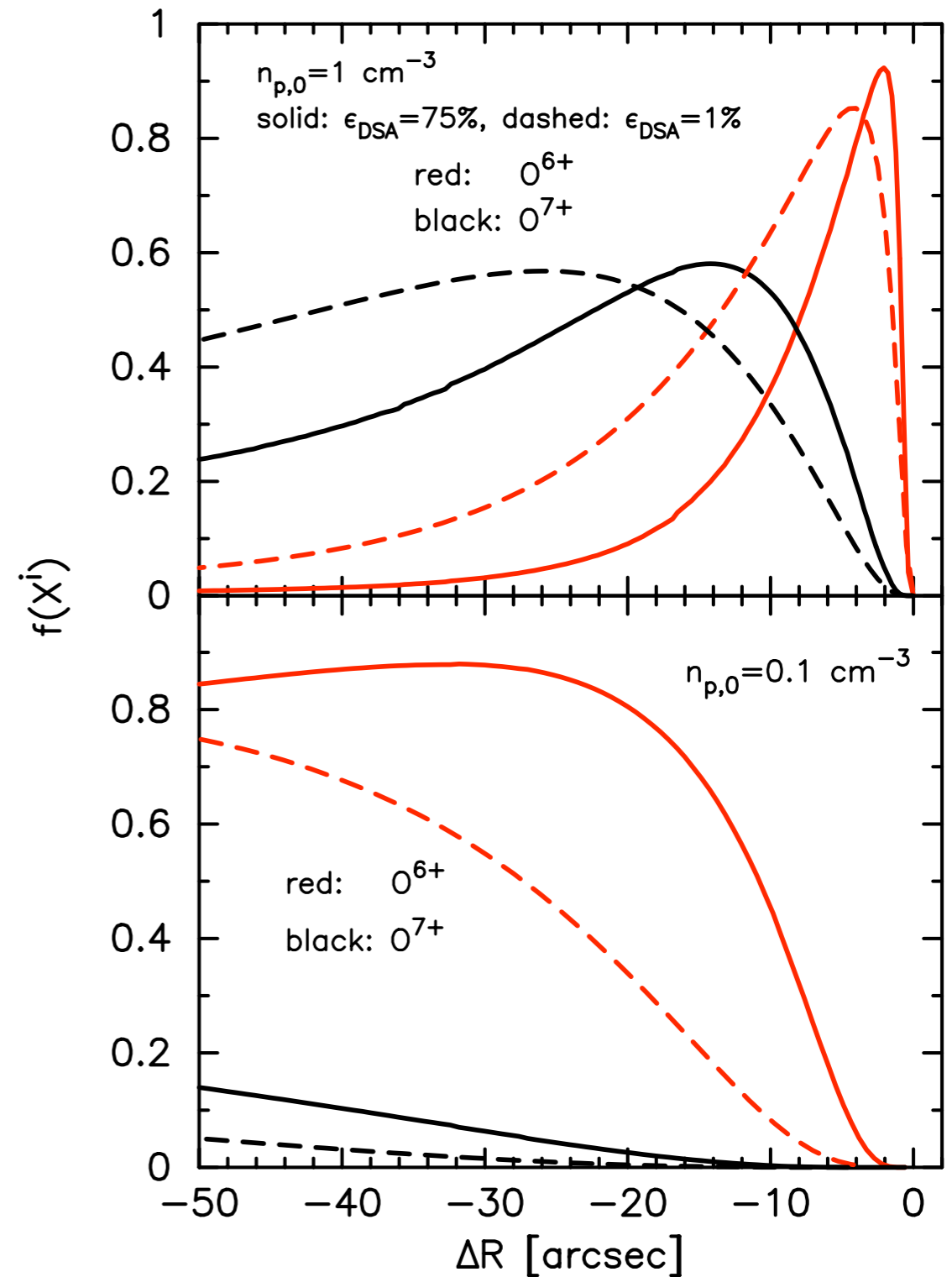
Using a sum of Maxwellians to approximate the particle distribution, Kaastra et al. (2009) showed that the presence of nonthermal particles can alter the relative intensities of satellite lines.

Above: summed particle distribution function behind a cluster shock. Right: simulated IXO spectrum (crosses) with the best fit Maxwellian-plasma model, showing an excess at Fe XXV j line.



# EXAMPLE RESULTS:

- In the efficient models, the charge state for a particular ion peaks closer to the shock front
- For instance, in the  $n_{p,0} = 1.0 \text{ cm}^{-3}$  models,  $\text{O}^{6+} \sim 2''$  behind the FS for efficient models, and  $\sim 4''$  behind for test particle models
- The resolved spatial and spectral structure could provide useful diagnostics for Galactic SNRs undergoing efficient shock acceleration

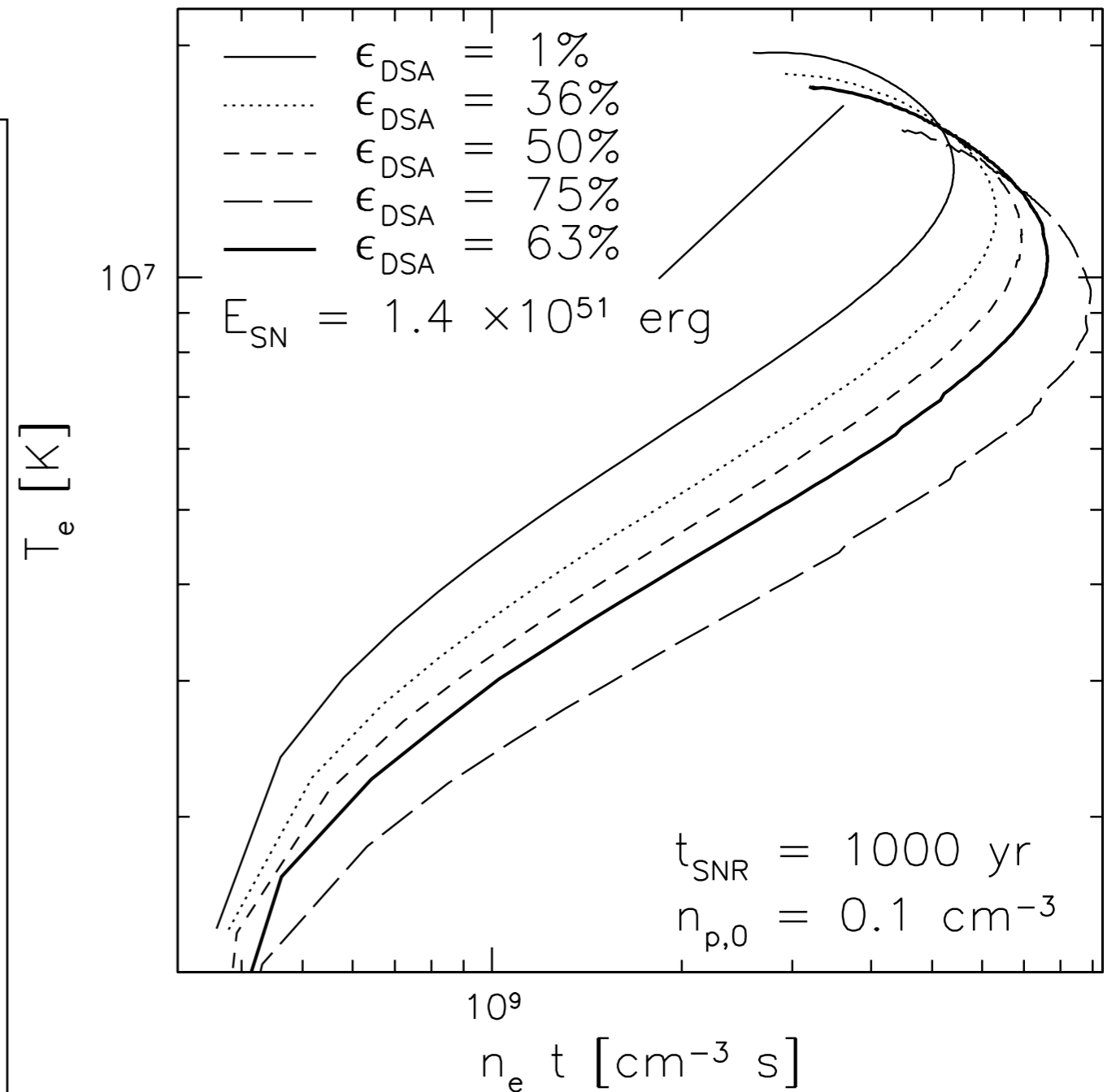


Top: Ionization fraction as a function of distance behind the forward shock for  $\text{O}^{6+}$  and  $\text{O}^{7+}$  with  $n_{p,0} = 1.0 \text{ cm}^{-3}$ . Bottom: Ionization fractions of  $\text{O}^{6+}$  and  $\text{O}^{7+}$  with  $n_{p,0} = 0.1 \text{ cm}^{-3}$ . In both panels, the solid curves are for  $\epsilon_{\text{DSA}} = 75\%$  and the dashed curves are for  $\epsilon_{\text{DSA}} = 1\%$ . The angular scale is determined assuming the SNR is at a distance of 1 kpc and the results are calculated at  $t_{\text{SNR}} = 1000 \text{ yr}$ .

# IONIZATION AGE VS ACCELERATION

## EFFICIENCY:

- Simulation also tracks ionization age ( $n_e t$ )
- For increasing acceleration efficiency, SNRs appear to have a higher ionization age
- Additionally, models with higher  $E_{\text{SN}}$  can but lower acceleration efficiency appear spectrally similar to models with lower  $E_{\text{SN}}$  but differing acceleration efficiencies

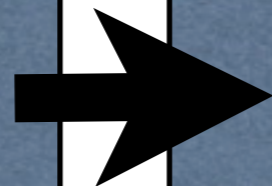


$n_e t$  vs  $T_e$  for varying acceleration efficiency. In these curves, the forward shock is in the lower left, and the contact discontinuity is at the upper right.

# Thermal X-ray Emission

Specify:

$E_{\text{sn}}, M_{\text{ej}},$   
 $n_{\text{amb}},$   
 $\epsilon_{\text{DSA}}, \dots$



Calculate:

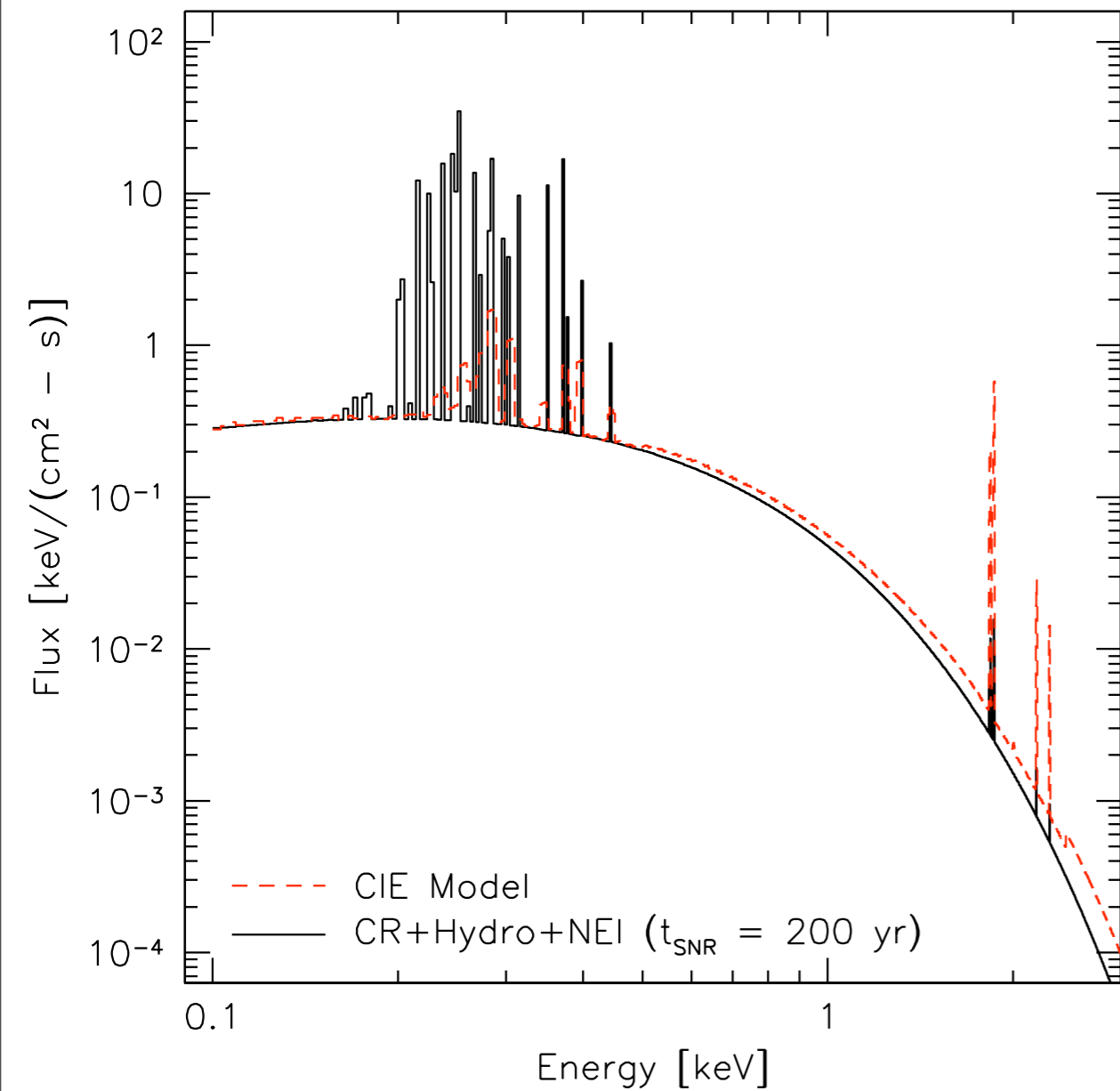
NEI, hydro,  
DSA, etc



Observables:

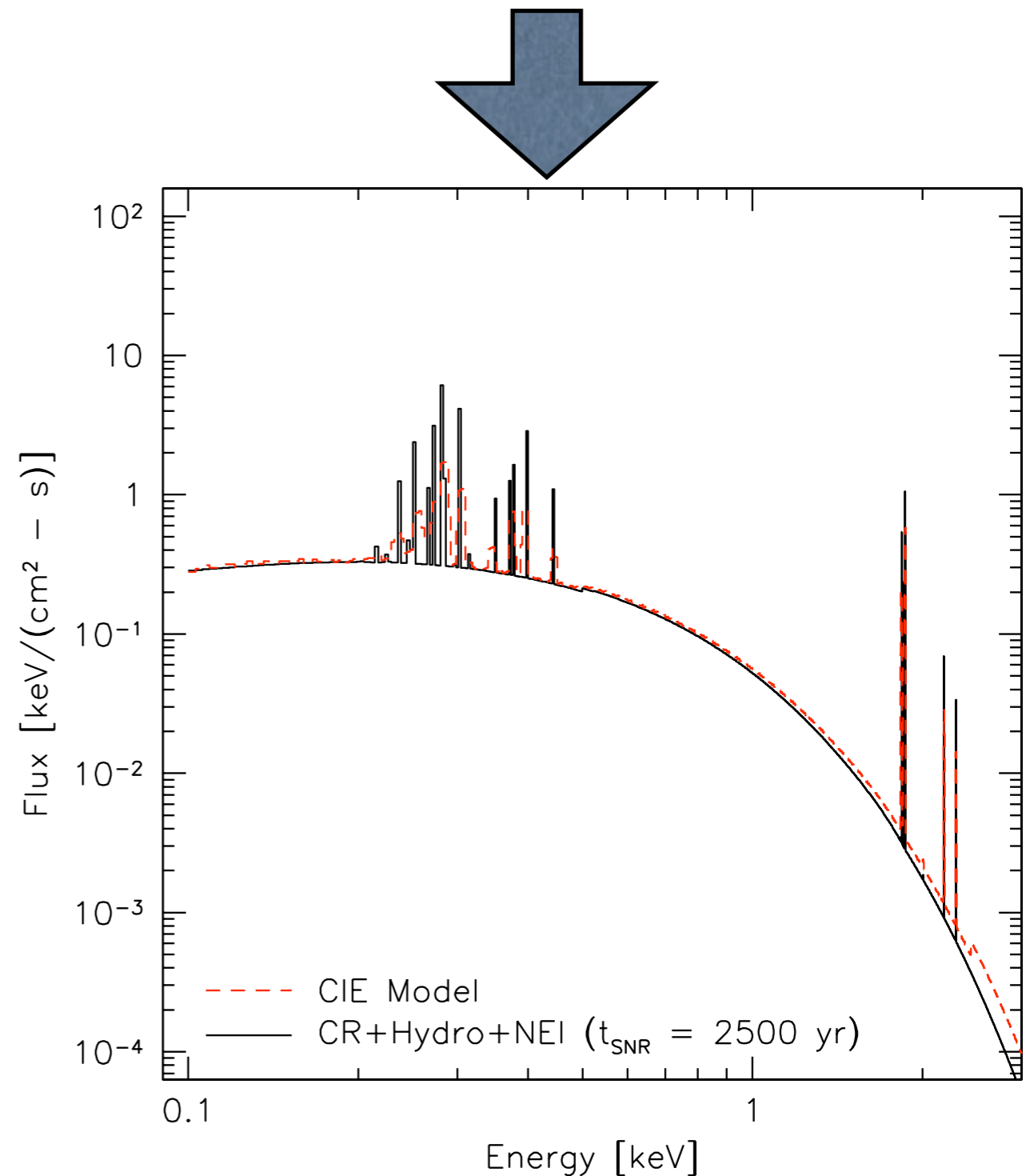
SNR  
dynamics,  
emitted  
spectrum

# testing the emissivity code:

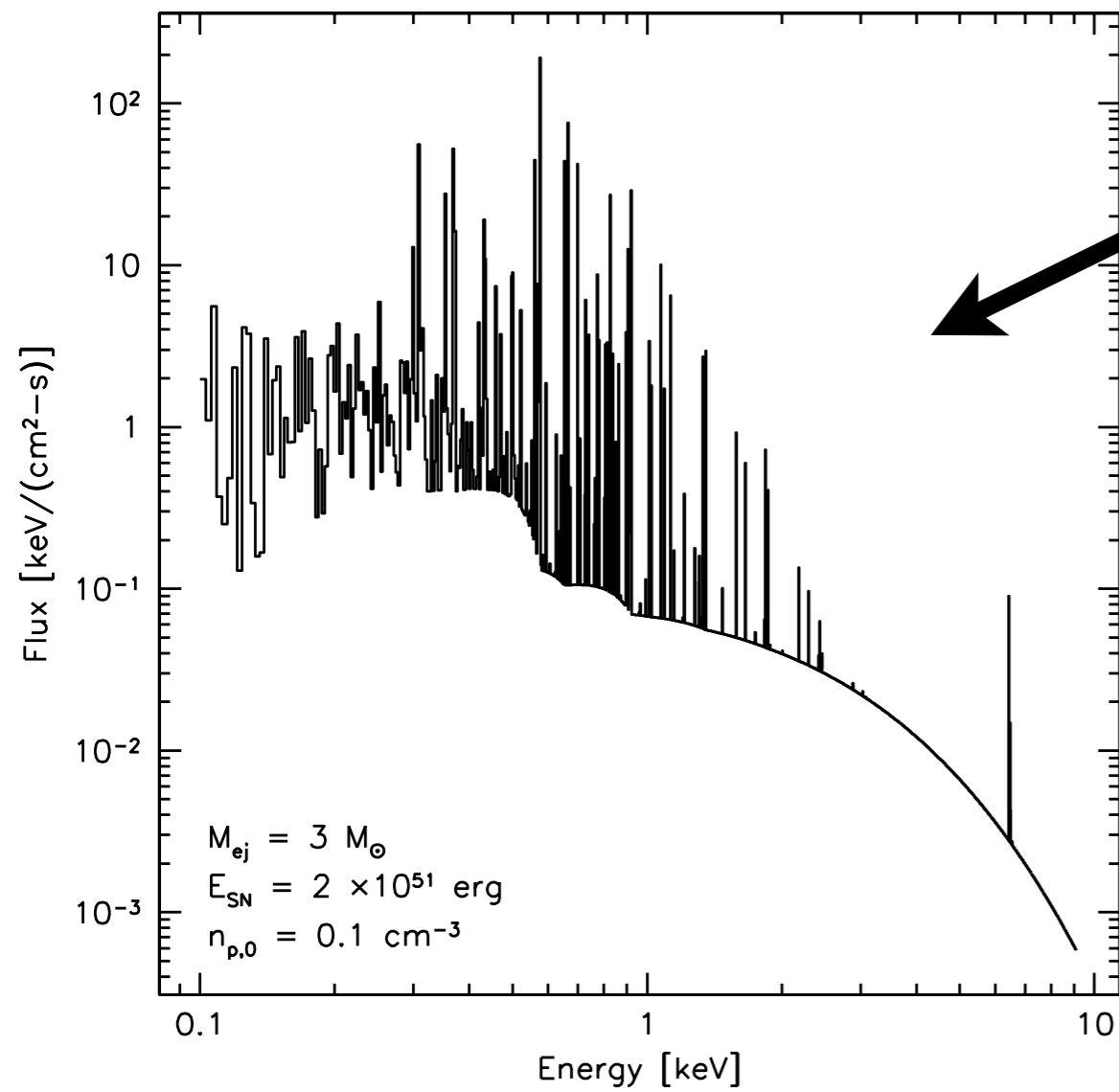


H and Si emission line spectrum  
at 200 yrs ( $n_{\text{e}t} = 1.9\text{E}+10$ )

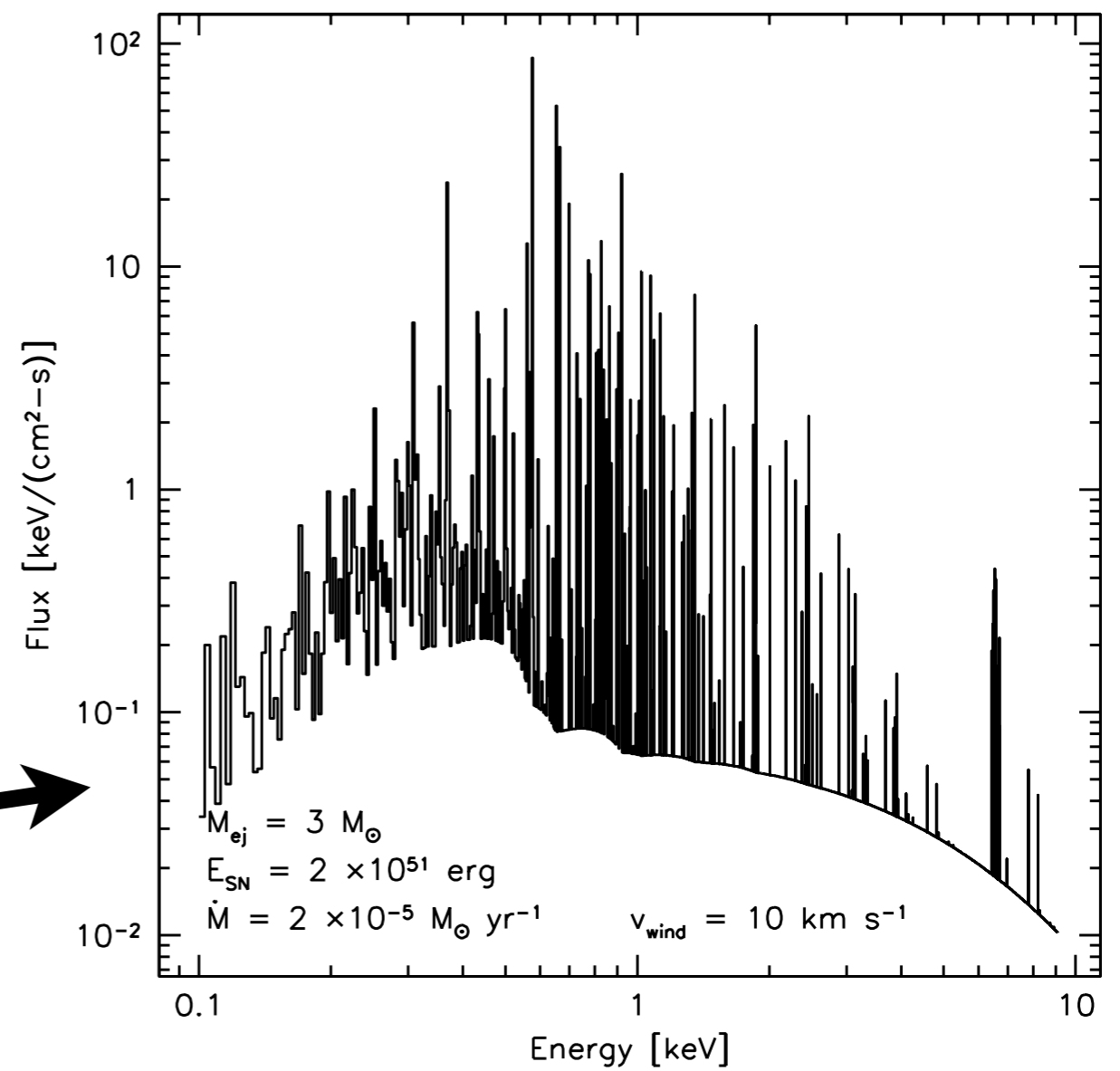
H and Si emission line spectrum  
at 2500 yr ( $n_{\text{e}t} = 2.4\text{E}+11$ )







Type Ia



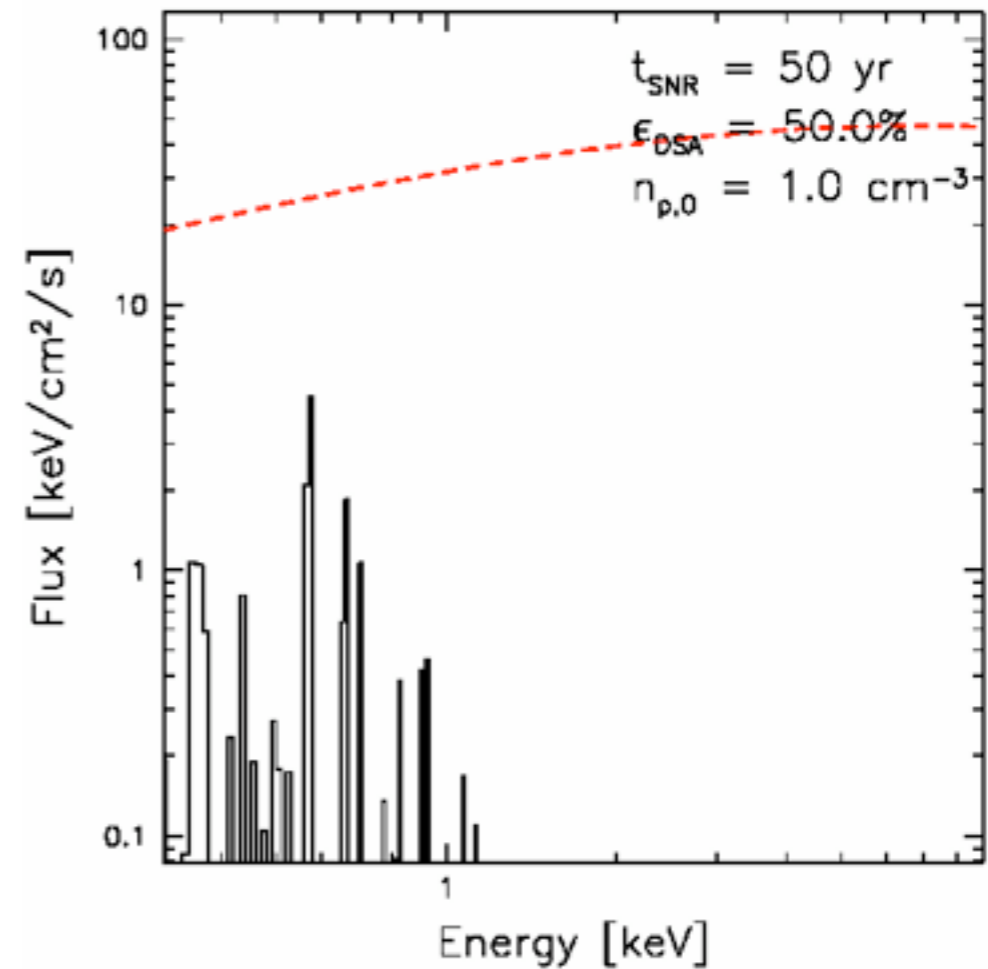
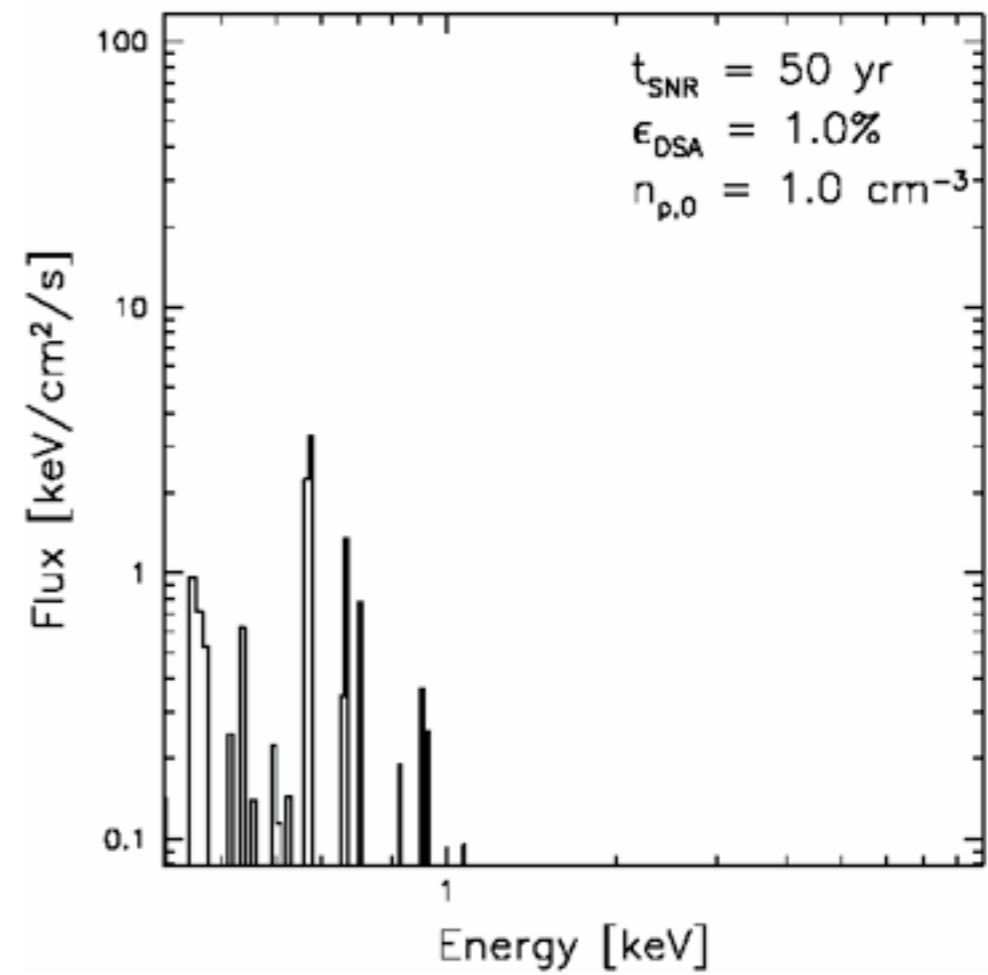
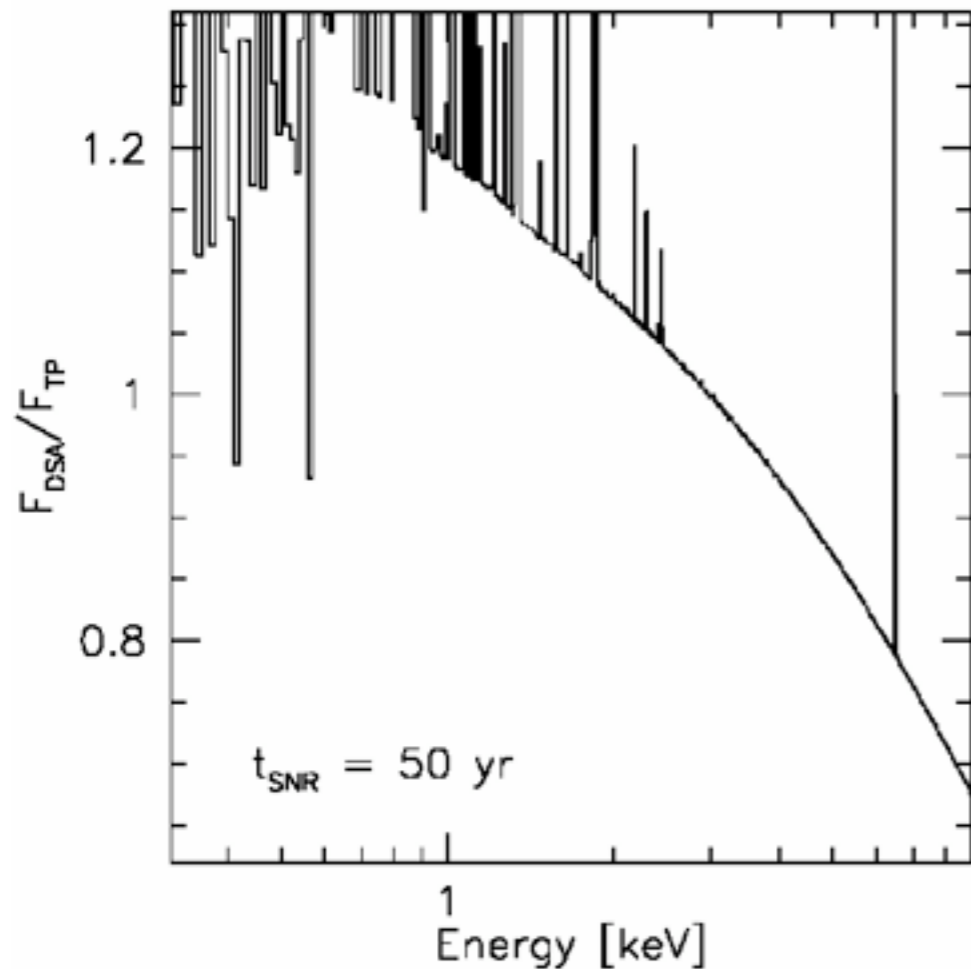
Core Collapse

# Example Results

$$E_{\text{SN}} = 10^{51} \text{ erg}$$

$$M_{\text{ej}} = 1.4 M_{\text{sun}}$$

$$n_{\text{p},0} = 1 \text{ cm}^{-3}$$



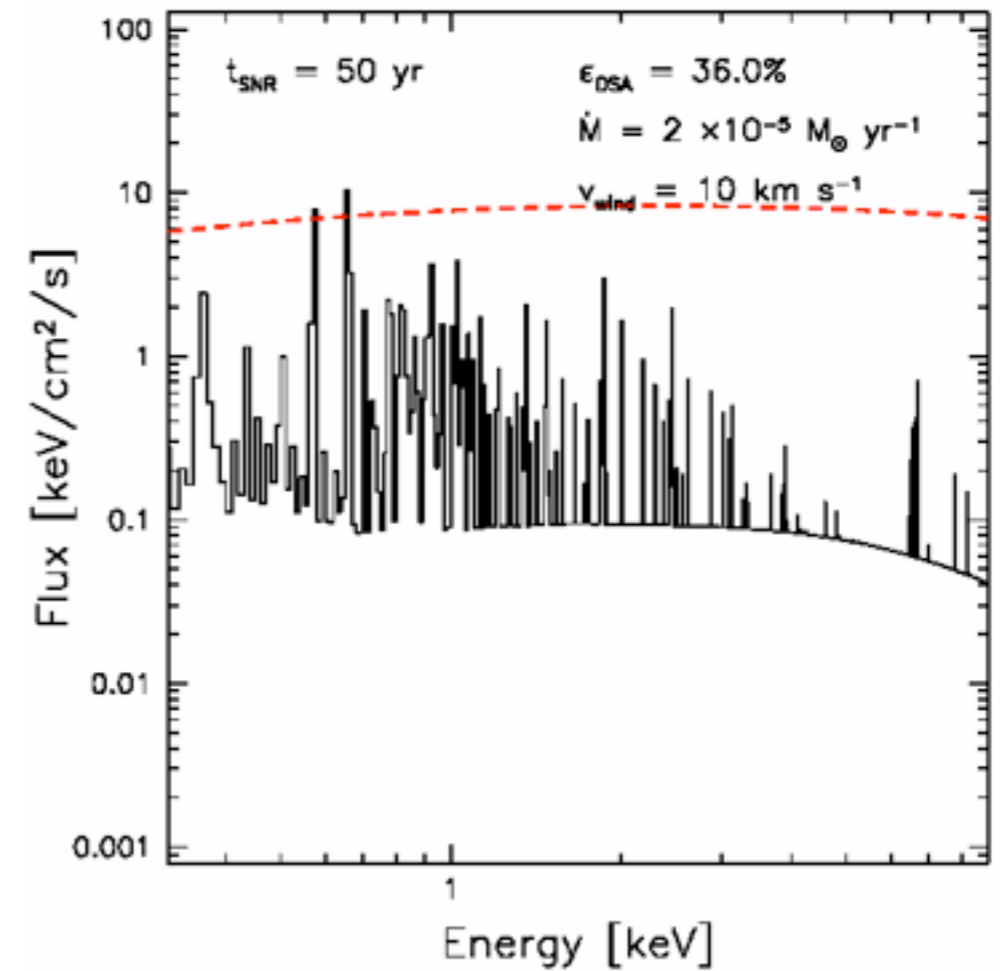
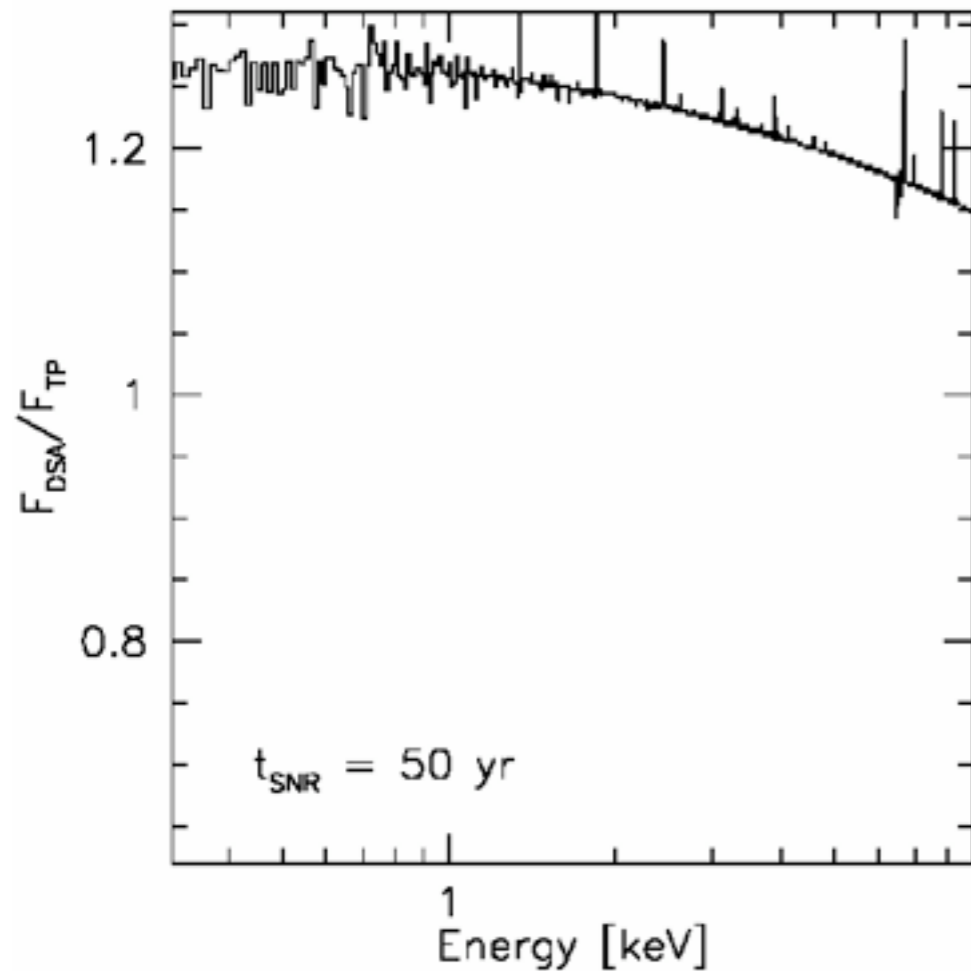
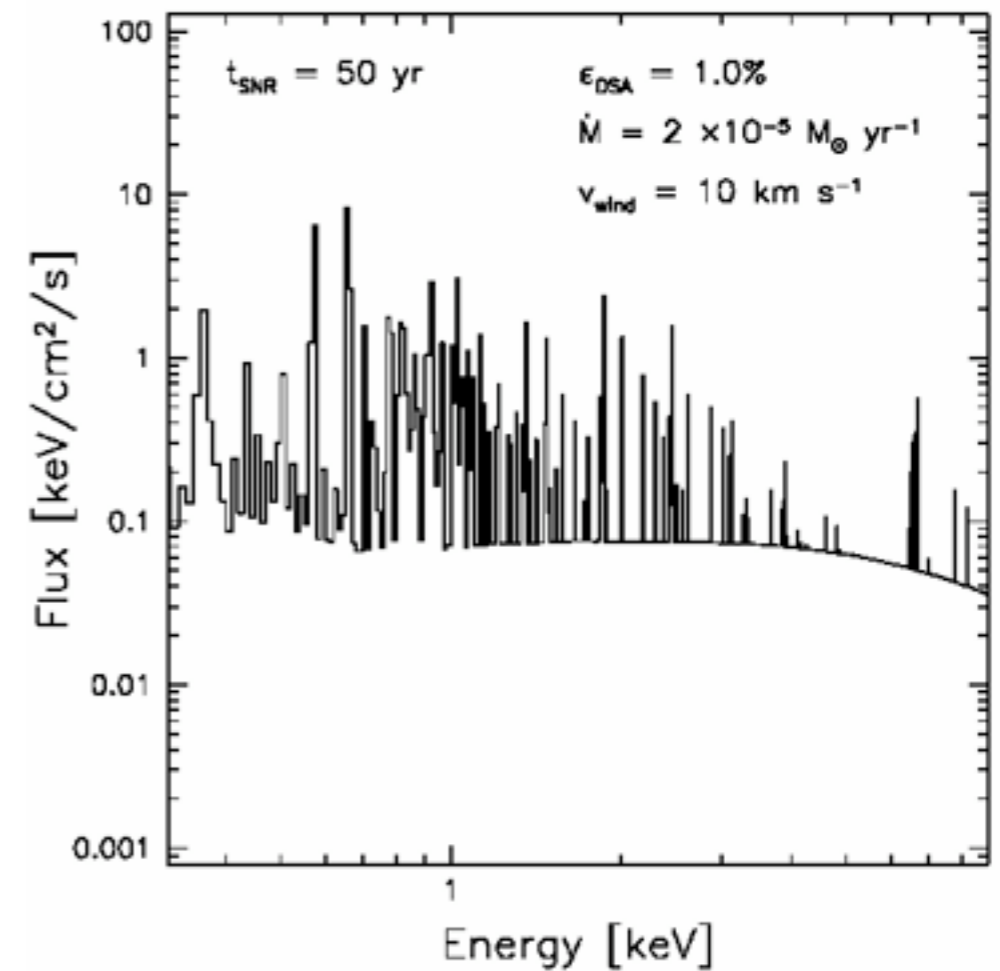
# Example Results

$$E_{\text{SN}} = 2 \times 10^{51} \text{ erg}$$

$$M_{\text{ej}} = 2 M_{\text{sun}}$$

$$\dot{M} = 2 \times 10^{-5} M_{\text{sun}} \text{ yr}^{-1}$$

$$v_{\text{wind}} = 10 \text{ km s}^{-1}$$



# CONCLUSIONS:

- Efficient cosmic-ray production via DSA significantly alters the NEI of material behind the SNR blastwave
- Because of higher postshock densities, higher charge states are reached at lower electron temperatures
- The spatial structure of the ionization varies with acceleration efficiency
- A general characteristic of the modification of the thermal emission via diffusive shock acceleration is that the line emission is brighter, but the underlying continuum is softer