

Plasma Instabilities in Relativistic Shocks

Antoine Bret

Universidad Castilla-La Mancha – Ciudad Real – Spain

Nonlinear Processes in Astrophysical Plasmas: Particle Acceleration, Magnetic Field Amplification, and Radiation Signatures

September 28, 2009 - October 2, 2009 Kavli Institute for Theoretical Physics

Outline

- Kind of systems considered
- Formalism
- Electron beam/plasma interaction No B₀
- Proton beam, B₀ //
- Conclusions

The kind of system considered

- Infinite beam/plasma system, homogenous, collisionless, current and charge neutral, magnetized or not.
- LINEAR evolution.
- Settings are numerous:



Formalism

Many ways to perturbate the system:

k Oblique – because real word perturbations don't choose "easy" alignment - Σ all orientations



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- Vlasov + Maxwell Linearized
- Perturbation $\propto \exp(i\mathbf{k}.\mathbf{r}-i\boldsymbol{\Theta}t)$ with $\mathbf{k}=(k_{\prime\prime},k_{\perp})$
- Dielectric tensor T

• Det(**T**)=0

$$\mathcal{T} = \begin{pmatrix} T_{xx} & T_{xy} & T_{xz} \\ T_{xy}^{*} & T_{yy} & T_{yz} \\ T_{xz}^{*} & T_{yz}^{*} & T_{zz} \end{pmatrix}$$

Every single unstable mode is recovered

About the dispersion equation

General form:

$$\det \left| \frac{\omega^2}{c^2} \varepsilon_{ij} + k_i k_j - k^2 \delta_{ij} \right| = 0$$

Where,

can't

$$\varepsilon_{kl}(\omega, \mathbf{k}) = \delta_{kl} + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \int \int \frac{p_k \partial f_{\alpha}^{(0)}}{\partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \int \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)} / \partial \mathbf{p}}{\partial p_l \partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \int \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)} / \partial \mathbf{p}}{\partial p_l \partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \int \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)} / \partial \mathbf{p}}{\partial p_l \partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \int \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)} / \partial \mathbf{p}}{\partial p_l \partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \int \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)} / \partial \mathbf{p}}{\partial p_l \partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \int \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)} / \partial \mathbf{p}}{\partial p_l \partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \int \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)} / \partial \mathbf{p}}{\partial p_l \partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \int \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)} / \partial \mathbf{p}}{\partial p_l \partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \int \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)} / \partial \mathbf{p}}{\partial p_l \partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \int \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)} / \partial \mathbf{p}}{\partial p_l \partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \int \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)} / \partial \mathbf{p}}{\partial p_l \partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \int \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)} / \partial \mathbf{p}}{\partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \int \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)} / \partial \mathbf{p}}{\partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)} / \partial \mathbf{p}}{\partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)} / \partial \mathbf{p}}{\partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)} / \partial \mathbf{p}}{\partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)} / \partial \mathbf{p}}{\partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)} / \partial \mathbf{p}}{\partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)} / \partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)} / \partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)} / \partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)} / \partial p_l} d^3 p + \sum_{\alpha=b,p} \frac{n_{\alpha}}{\omega^2} \int \frac{p_k p_l \mathbf{k} \cdot \partial f_{\alpha}^{(0)$$

Cold regime (DF Dirac δ 's) and no B0



Fainberg, et al, Sov. Phys. JETP 30, 528 (1970). Califano, et al, Phys. Rev. E 58, 7837 (1998). Bret, PoP 12, 082704, (2005)

Kinetic effects (still no B_0)

• Waterbag calculation (10 s CPU instead of 3 weeks). $N_b/N_p = 0.1$, $\gamma_b = 4$





Beam Temp. first stabilizes Weibel & modes with $k_{\perp} >> k_{\prime\prime}$

L.O. Silva et al., Bull. Am. Phys. Soc. 46, 205 (2001)

Waterbag model: PIC's confirmations



Dieckmann, PoP, 2006



Gremillet, PoP, 2007



Kong, PoP, 2009

Protons

Electrons

return curren

Kinetic regime (still no B₀)

Typical spectrum and generated patterns



IF dominant...

Kinetic regime (still no B₀) – Modes hierarchy



A. Bret et al., PRL 100, 205008 (2008).

Shock physics context



Oblique modes: Growth rate scalings

Oblique modes govern in the following limits:



Limits are relevant as Oblique modes keep governing



Oblique modes: Growth rate scalings

Oblique modes similar to Two-stream ones:



Proton beam – Magnetized, cold



Proton beam – Magnetized, cold B0, Ma=1/20 All \mathbf{k} 's – largest GR(\mathbf{k}) Oblique Protons (Mass ratio 1/100) Protons, γ=20, density 1/50 -1.5 **Electrons** -2.5 -3 $\log(\gamma/\omega_{pe})$ -3.5 Buneman Weibel 2-Stream too small) -5.5 -6 Bell -6.5 1 0.001 0.5 Alfven 0 -0.5 Bell -1 -1.5 -3 -2 $\log(k_{\perp}V_{CR}/\omega_{pe})$ 10^{-4} 0.001 0.01 0.1 100 $\log(k_{//}V_{CR}/\omega_{pe})$

Conclusions

- Beam plasma instabilities are relevant for:
 - Shock Formation (the trigger),
 - Particle acceleration.
- The simplest system (no B₀, fixed ions, e-beam+RC) has been analyzed thoroughly, including in the kinetic regime
- Shock Formation: comparable density shells. Weibel/Oblique dominant.
- Particle Acceleration: Diluted beam regime, Oblique modes.
- System with CR, B_0 : rich, multi-scale spectrum.

THANK YOU

Shock Physics: Formation

- Typical scenario: colliding shells.
- PIC scheme: Plasma interact with itself
 -> equal density shells (Fila/Weibel)
- What if different density shells?
- If un-magnetized, pair plasmas, or e-protons: Nb/Np







Shock Physics: Particle acceleration

- Relativistic, diluted, e-beam/plasma interaction, B₀ // Flow.
- Typical sprectrum:



For a thorough analysis of $B_0 \perp$ Flow case, see: Lemoine & Pelletier, MNRAS, *In Press* - arXiv:0904.2657

Shock Physics: Particle acceleration

- Consider a relativistic, diluted, e-beam/plasma interaction.
- Dominant mode:



Outline of the Talk

- System considered
- Plasma instabilities and relativistic shocks
 - Shock Formation
 - Particle acceleration
- Oblique Modes Cold and Waterbag results
- Kinetic theory of Oblique Modes
 - Scaling laws for the Growth Rate
- Non-parallel flows
- Conclusions