



Leibniz-Institut für
Astrophysik Potsdam

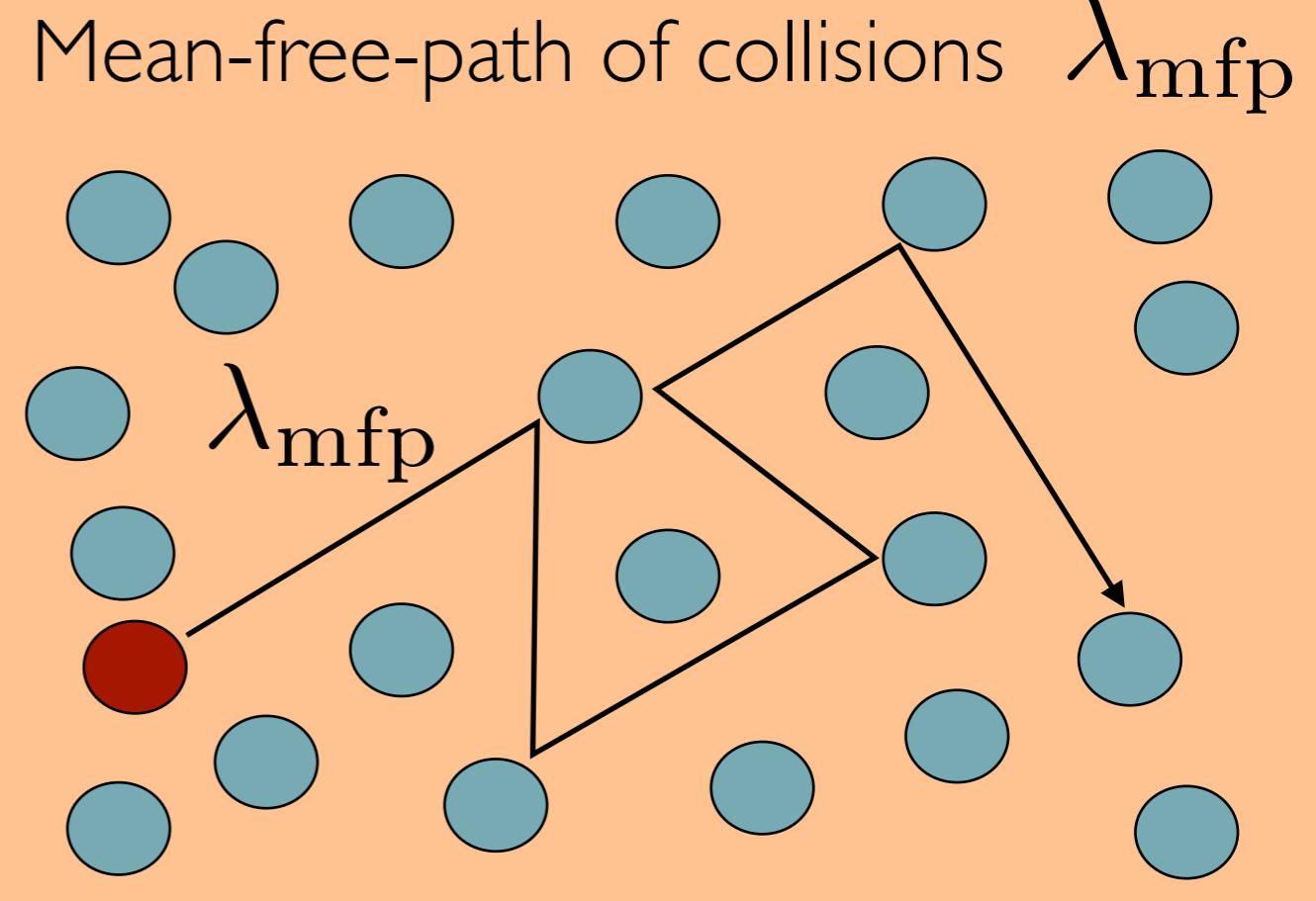
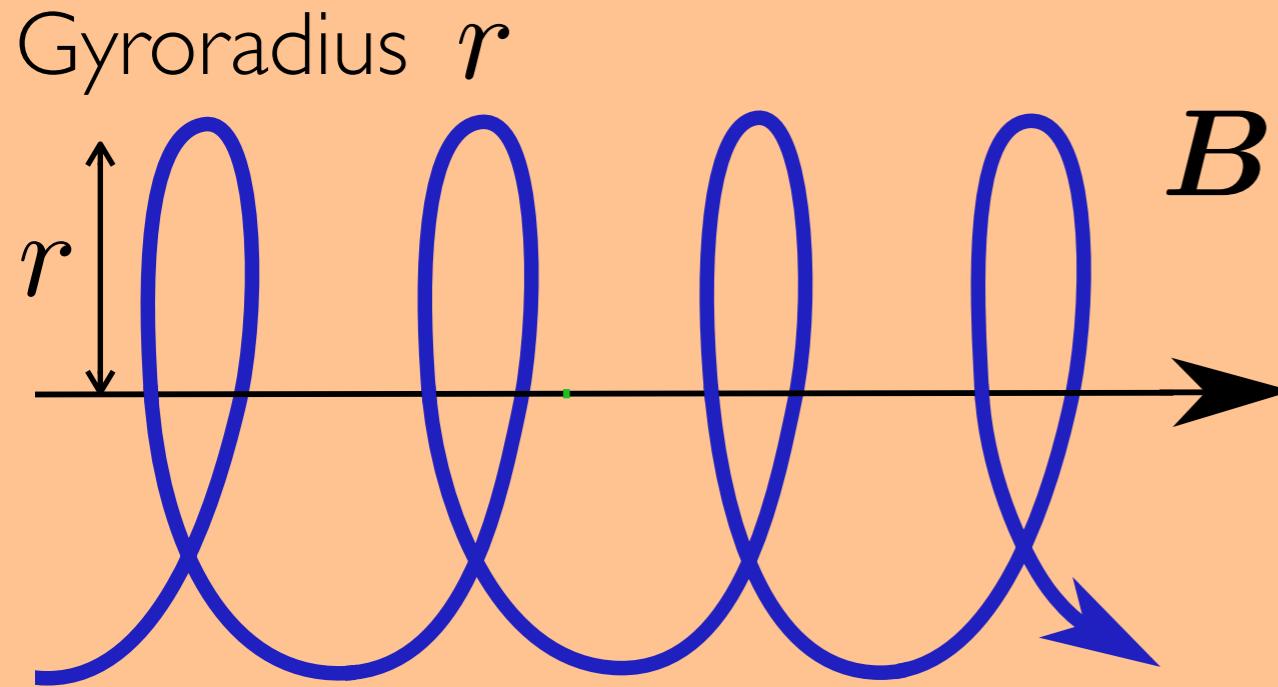
BRAGINSKII VISCOSITY ON AN UNSTRUCTURED, MOVING MESH ACCELERATED WITH SUPER-TIME-STEPPING

THOMAS BERLOK

Rüdiger Pakmor, MPA

Christoph Pfrommer, AIP

MULTISCALE PHENOMENA IN PLASMA ASTROPHYSICS, KITP 17/9/2019



Weakly collisional and magnetized

$$r \ll \lambda_{\text{mfp}} \ll L \longrightarrow$$

Transport of heat and momentum is along magnetic field lines.

Heat conduction

$$\mathbf{Q} = -\chi_{\parallel} \mathbf{b} (\mathbf{b} \cdot \nabla T)$$

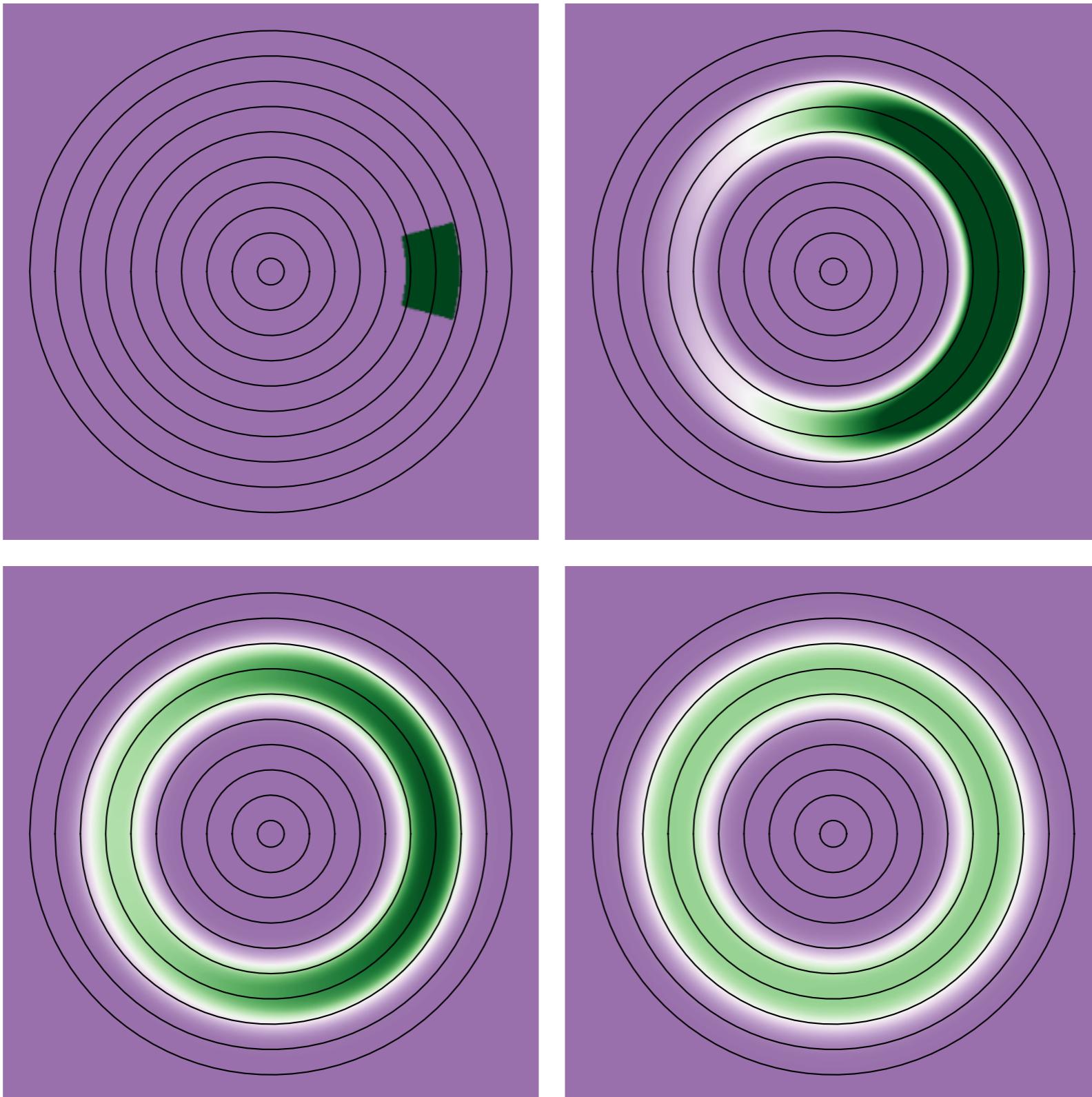
Kannan+ 2016 in Arepo
Sharma & Hammett 2007

Braginskii viscosity

$$\Pi = -\Delta p \left(\mathbf{b} \mathbf{b} - \frac{1}{3} \right),$$

$$\Delta p = \rho \nu_{\parallel} (3 \mathbf{b} \mathbf{b} : \nabla \mathbf{v} - \nabla \cdot \mathbf{v}) .$$

ANISOTROPIC DIFFUSION



ATHENA (Stone et al. 2008)

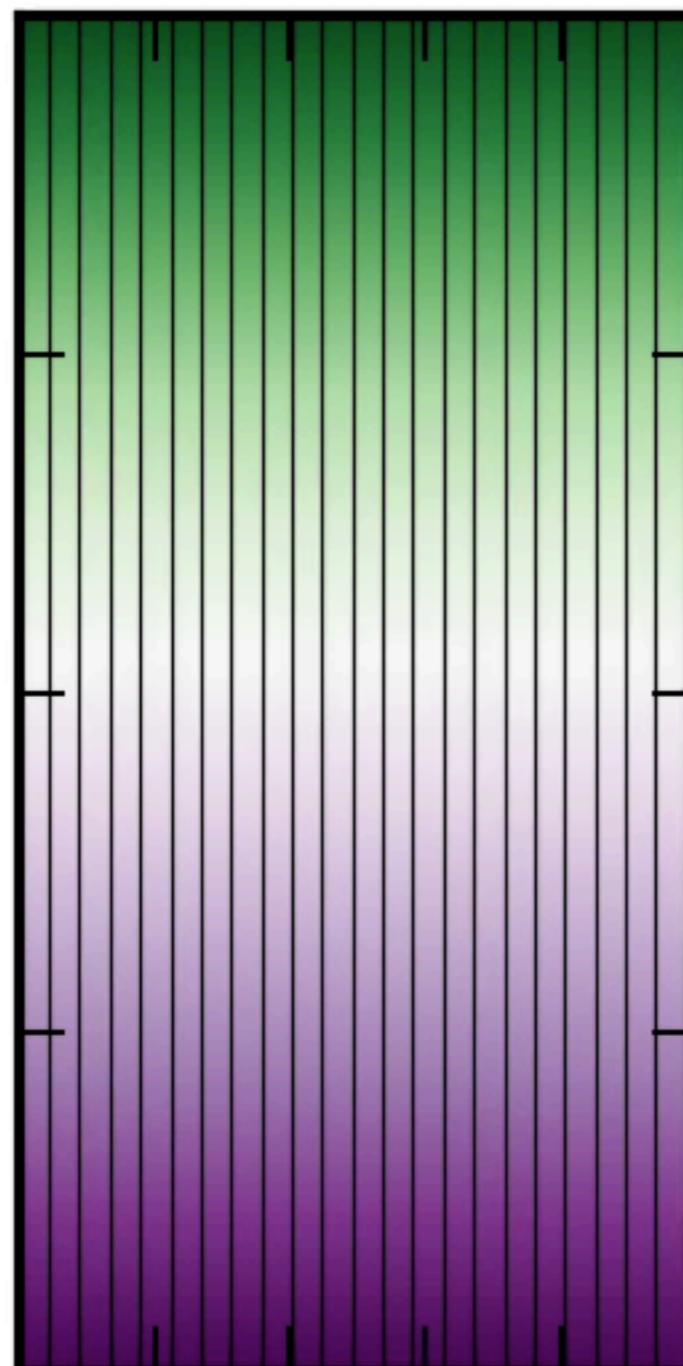
(Parrish & Stone 2005, Sharma & Hammett 2007, Berlok & Pessah 2016a)

QUASI-GLOBAL SIMULATIONS

Berlok & Pessah 2016b, ApJ

$t/t_0 = 0.0$

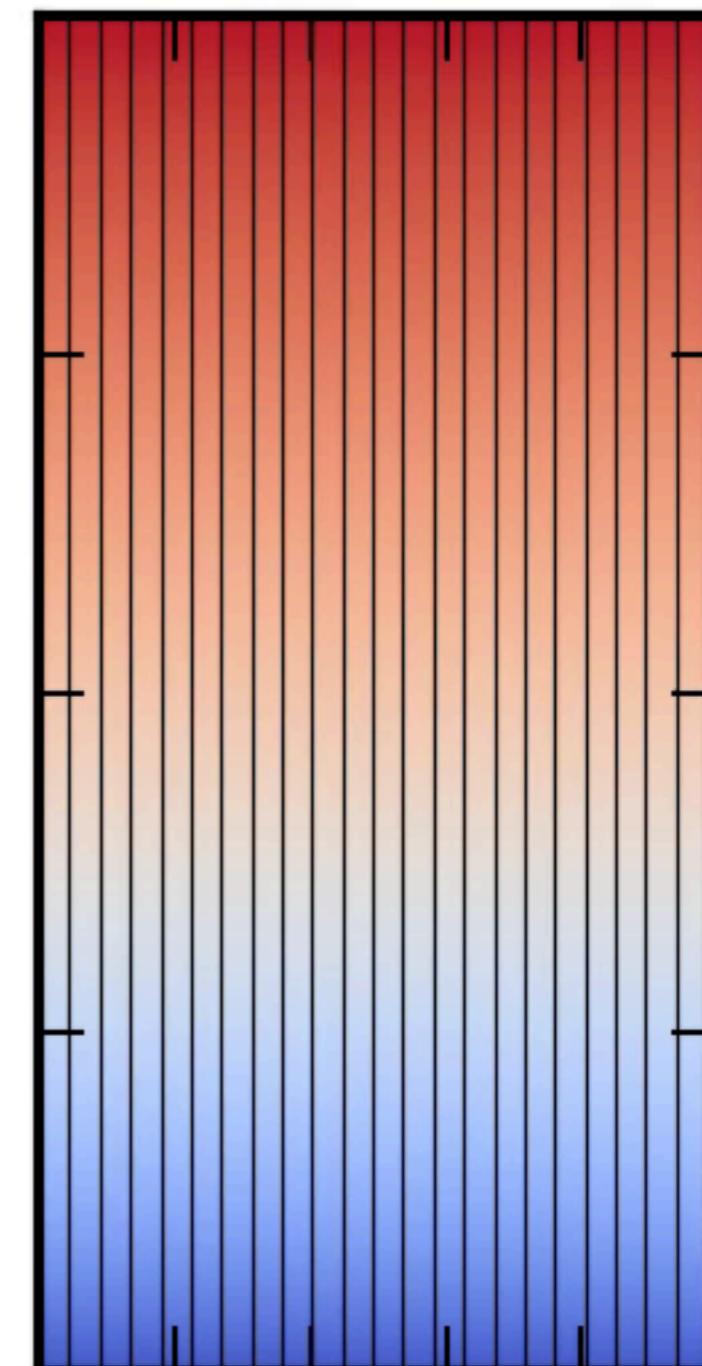
helium



Peng & Nagai 2009

hydrogen

H_0



hot

cold

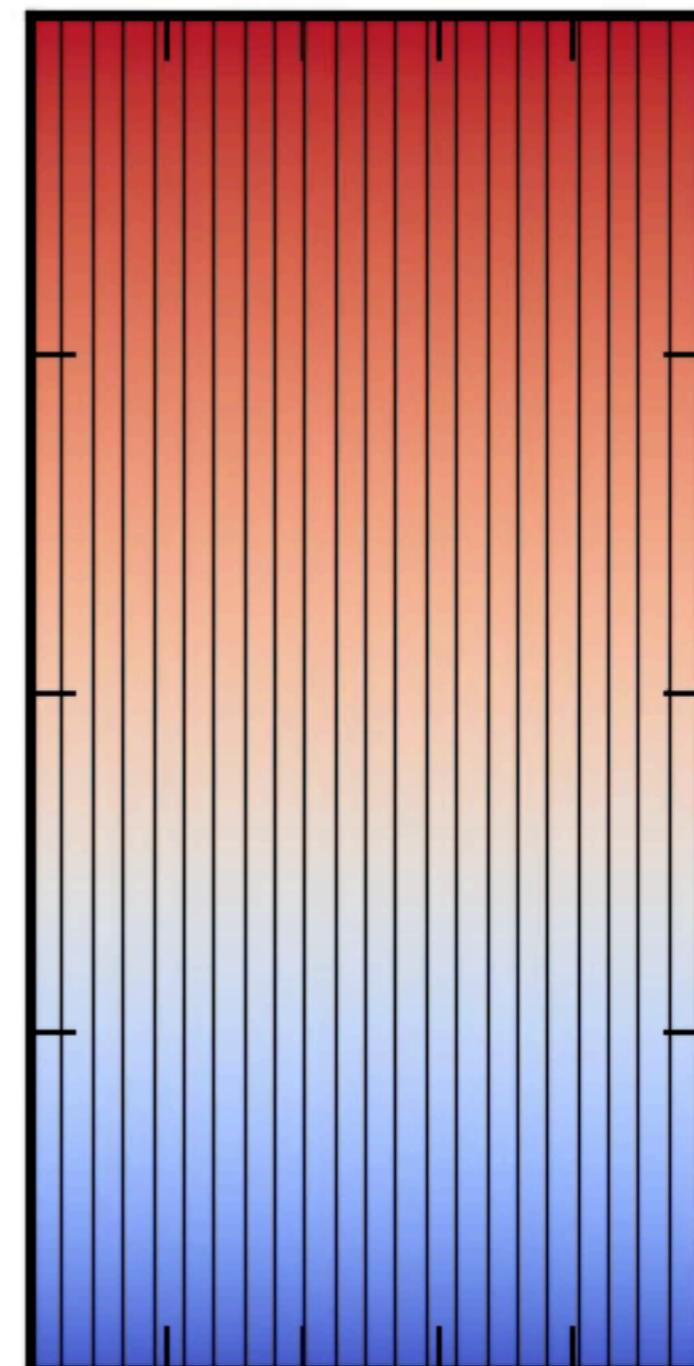
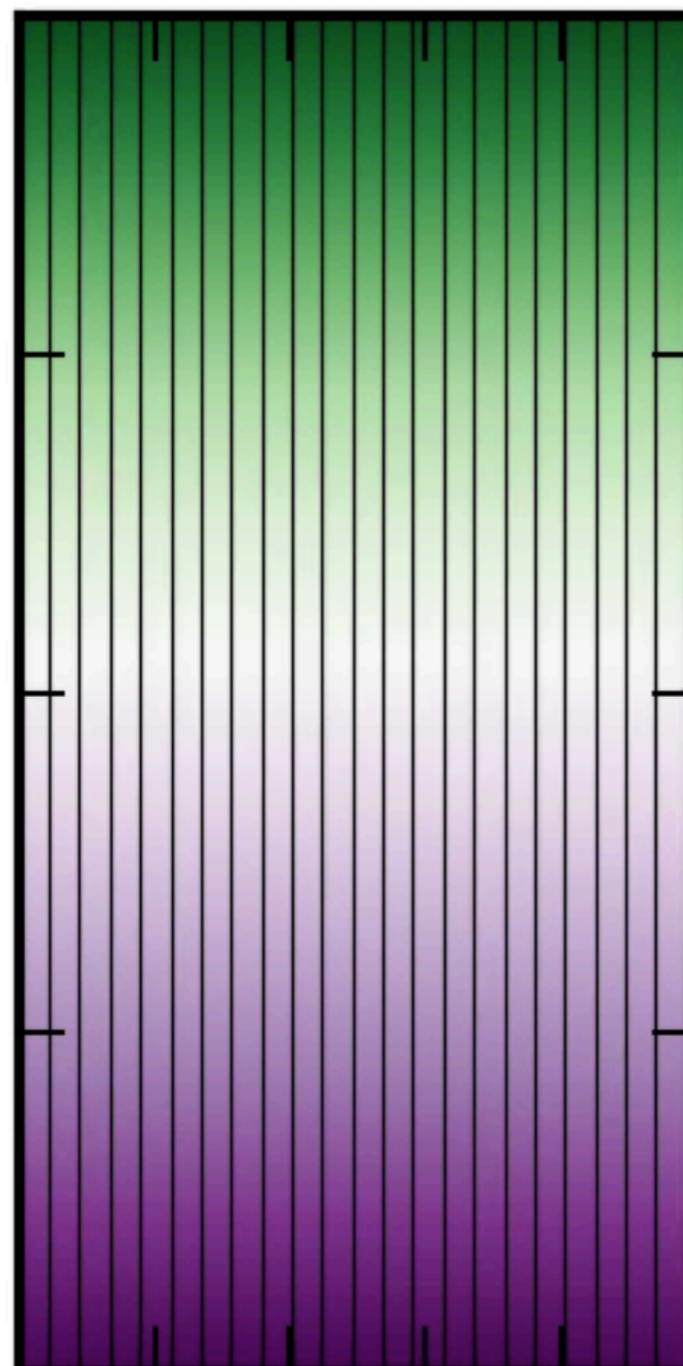
See Balbus 2000, 2001; Parrish & Stone 2005, 2007; Quataert 2008; Parrish & Quataert 2008; Parrish et al. 2008, 2009; Bogdanovic et al. 2009; Parrish et al. 2010; Ruszkowski & Oh 2010; McCourt et al. 2011, 2012; Latter & Kunz 2012, Kunz et al. 2012; Parrish et al. 2012a,b

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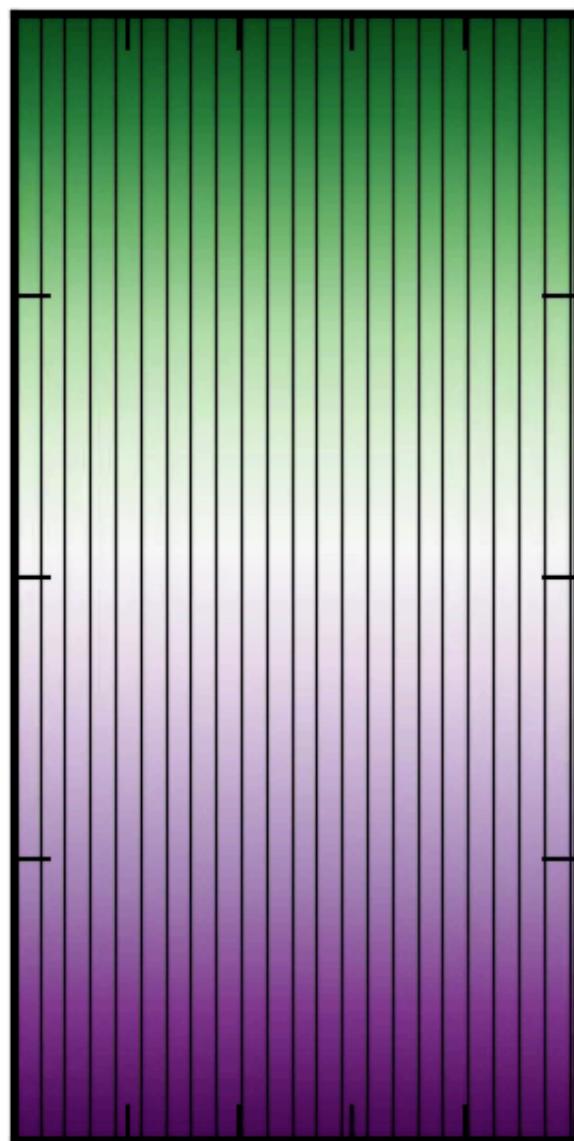
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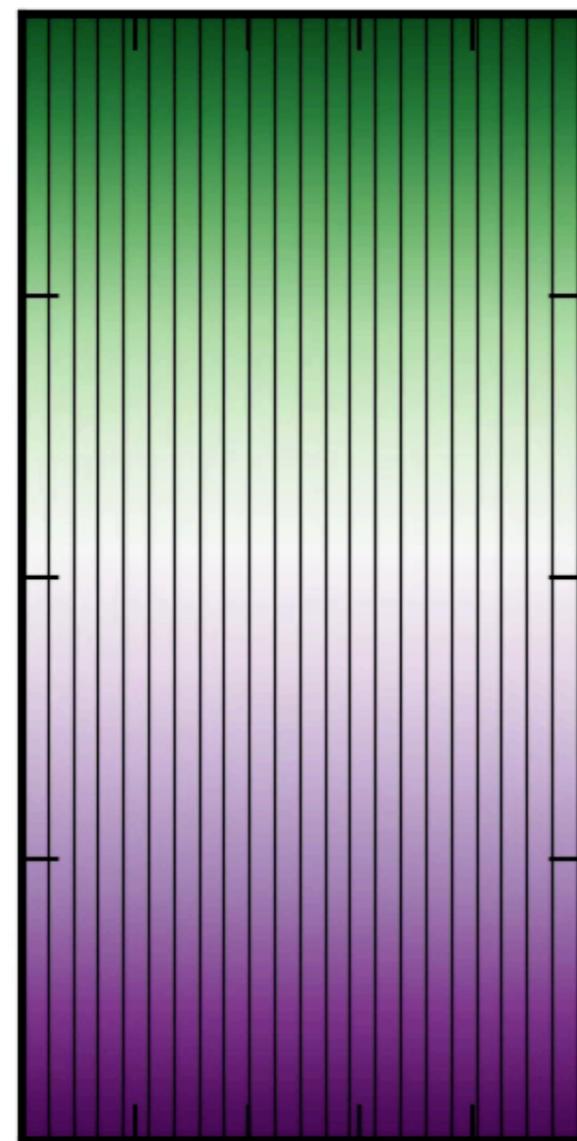
isoP



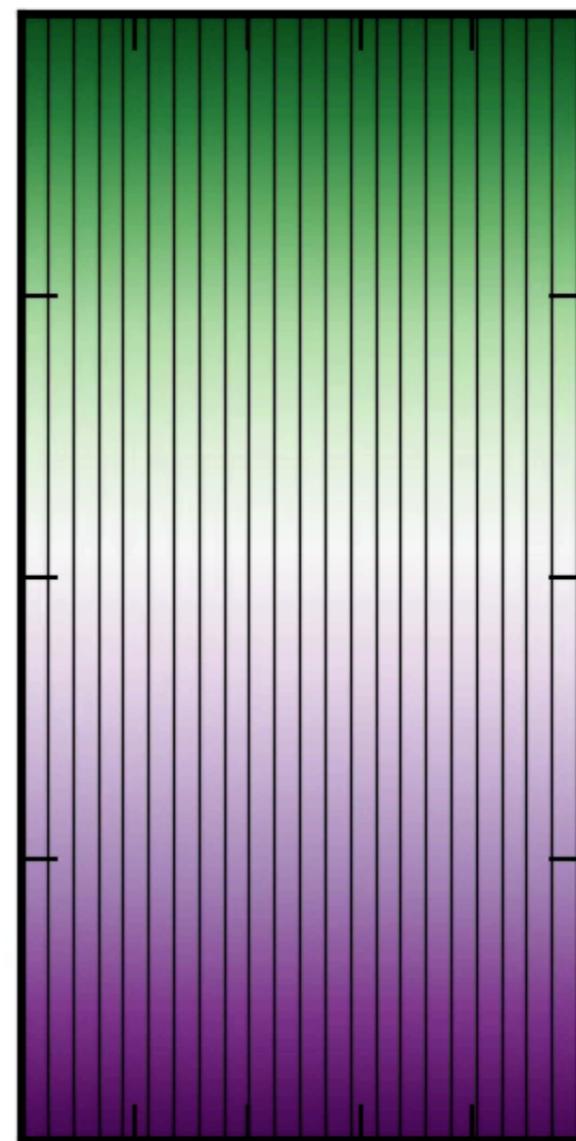
H_0

Berlok & Pessah 2016b
See also Kunz+ 2012

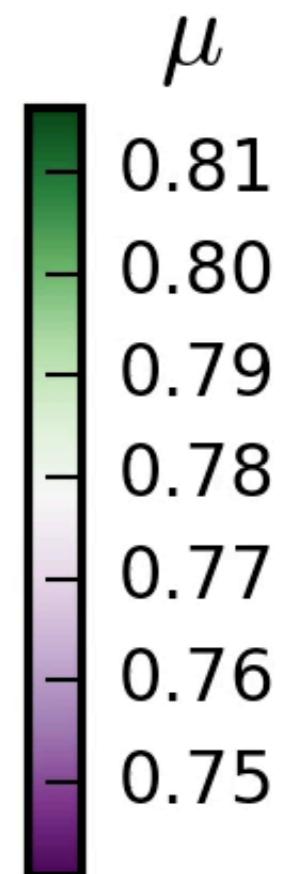
Blim



Brag



$t/t_0 = 0.0$

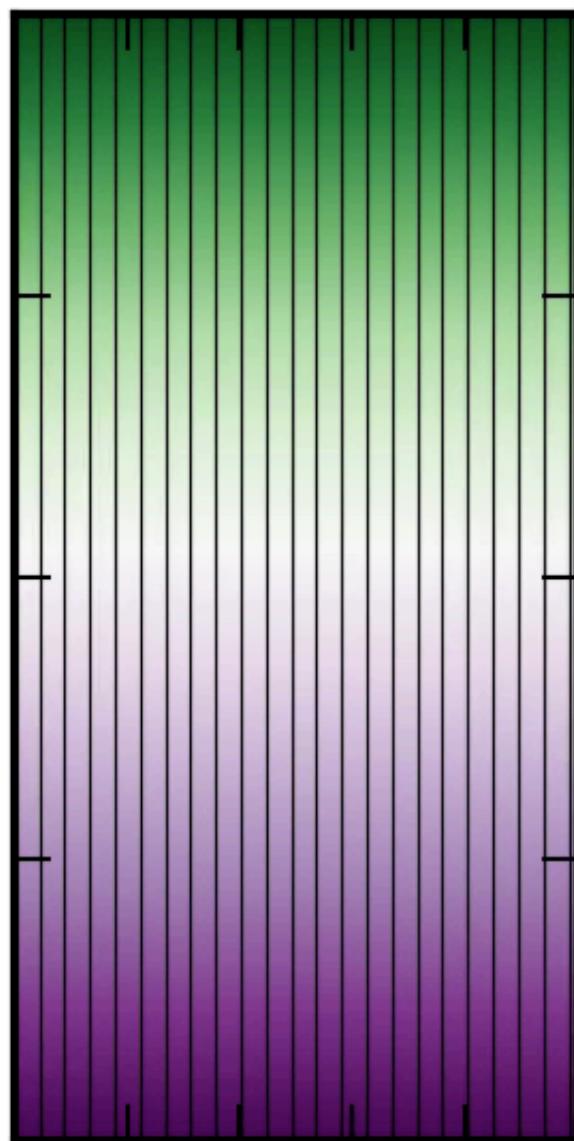


$$-\frac{B^2}{\mu_0} < \Delta p < \frac{B^2}{2\mu_0}$$

See Sharma+ 2006,
Schekochihin+ 2008, Kunz+ 2014

QUASI-GLOBAL SIMULATIONS

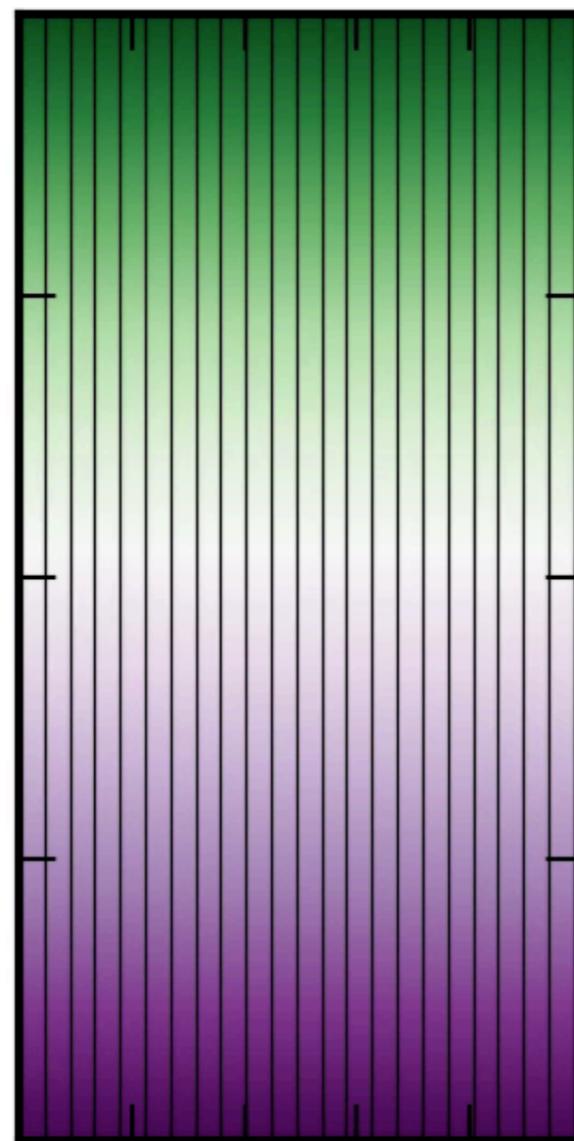
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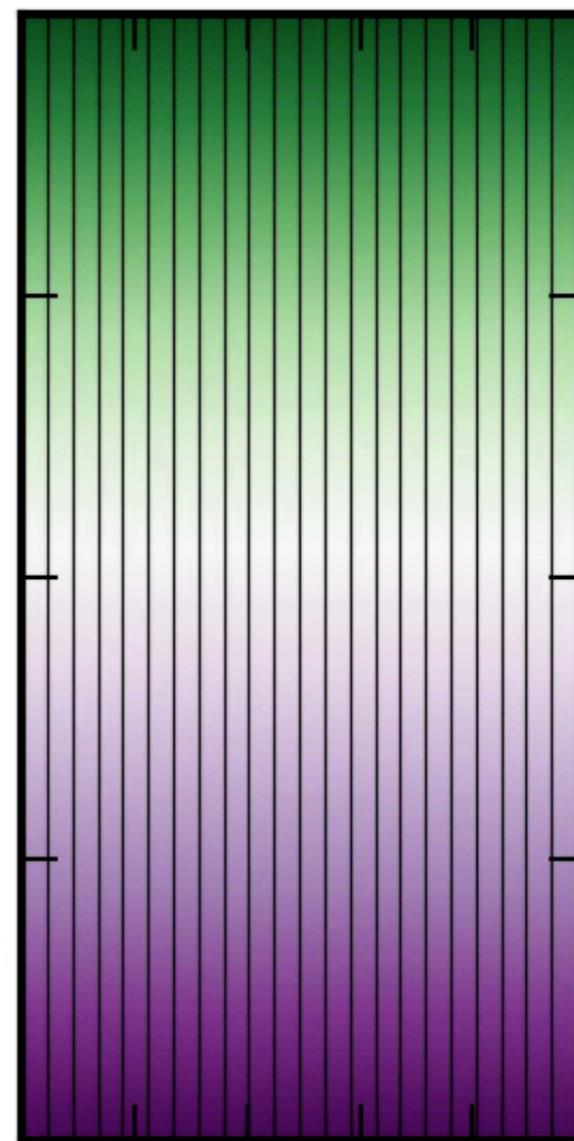
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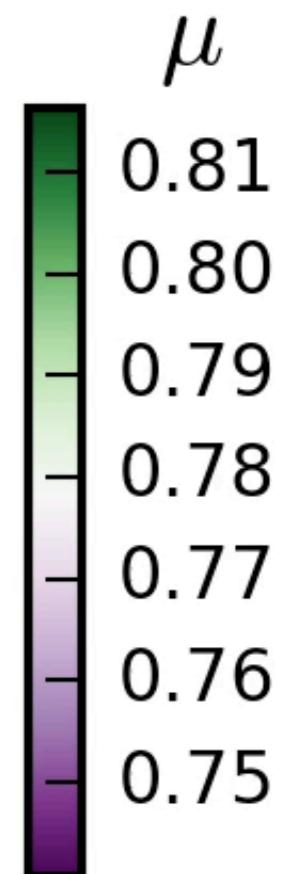
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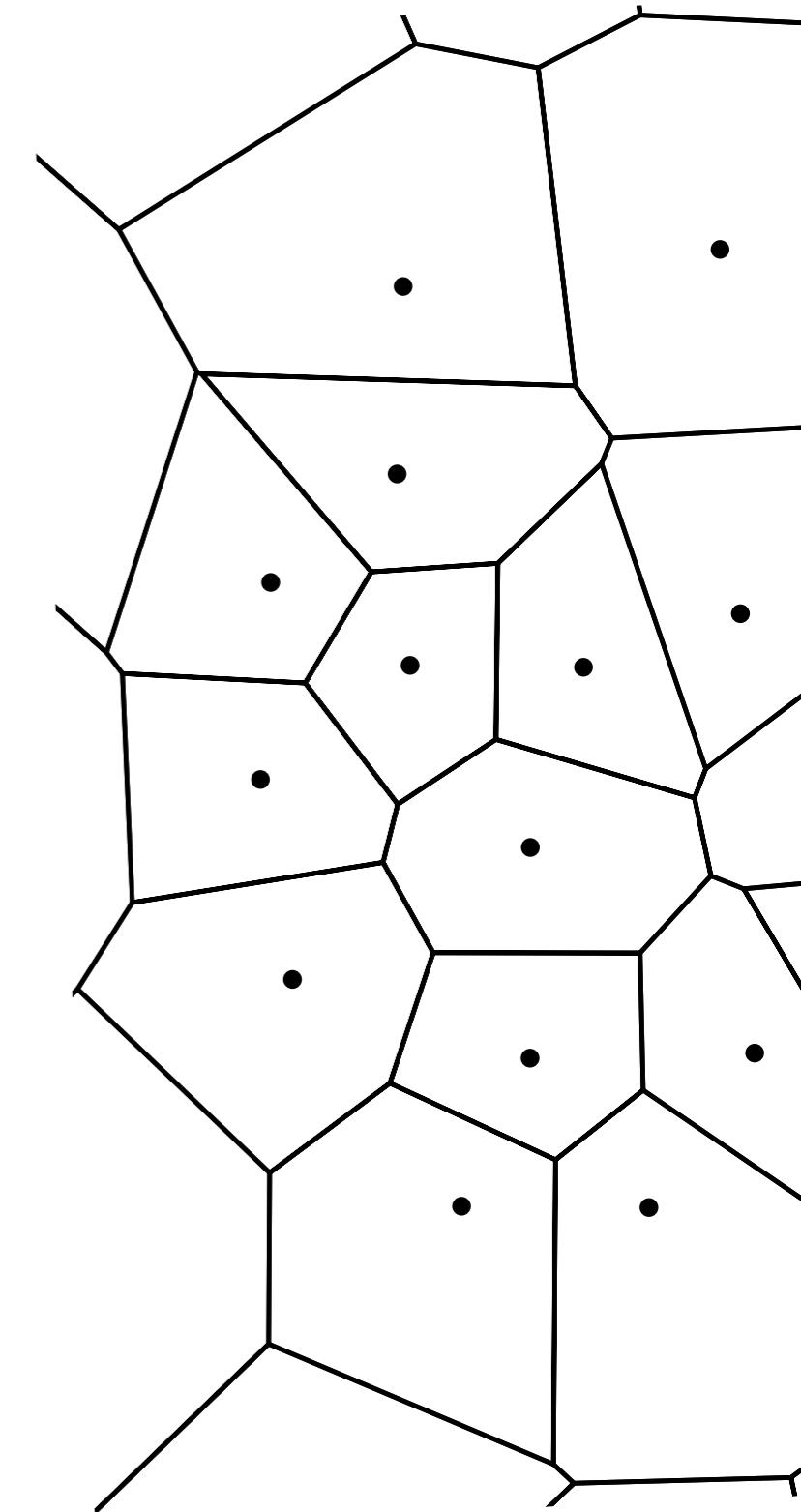
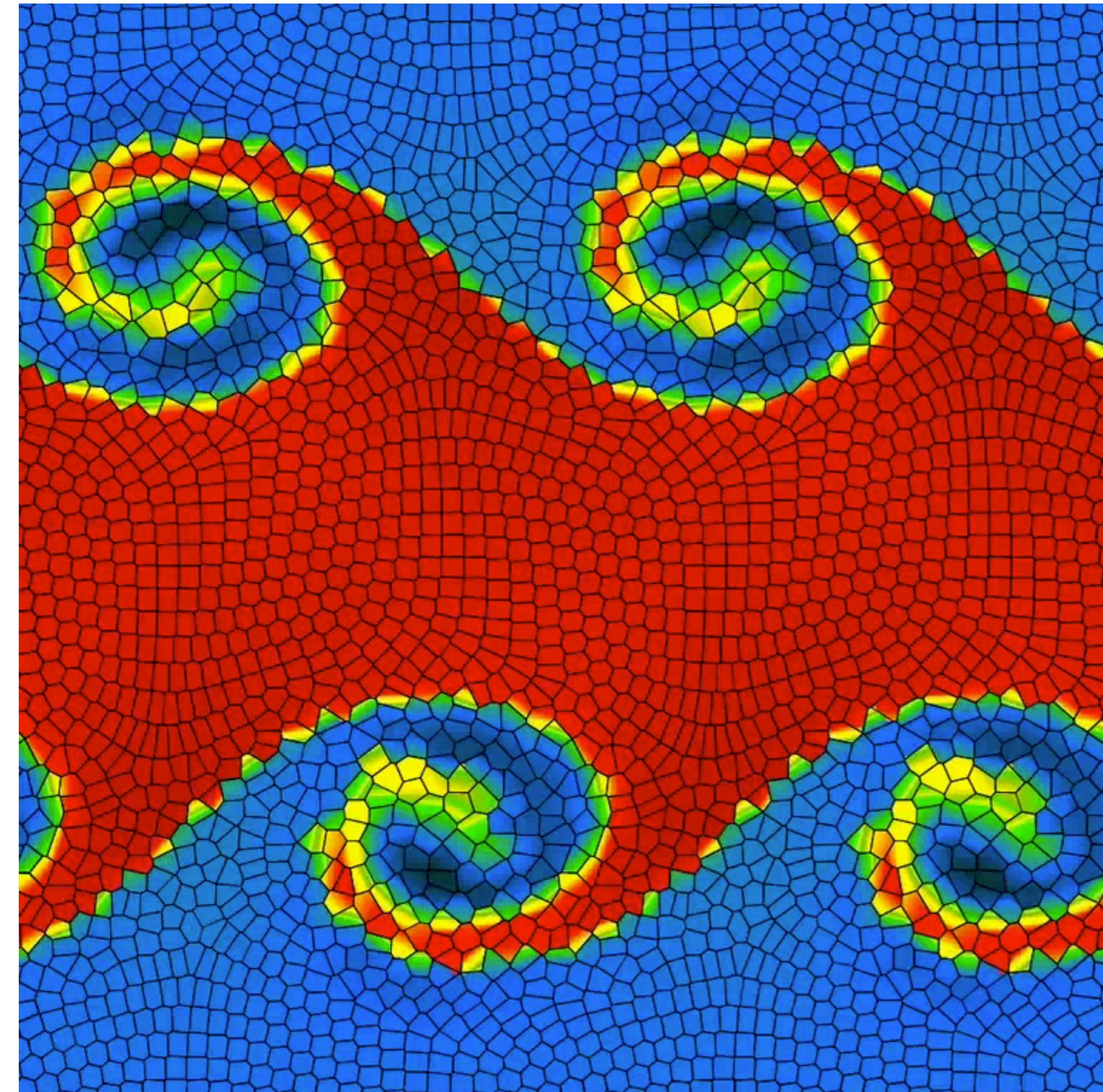


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THE MOVING MESH CODE AREPO

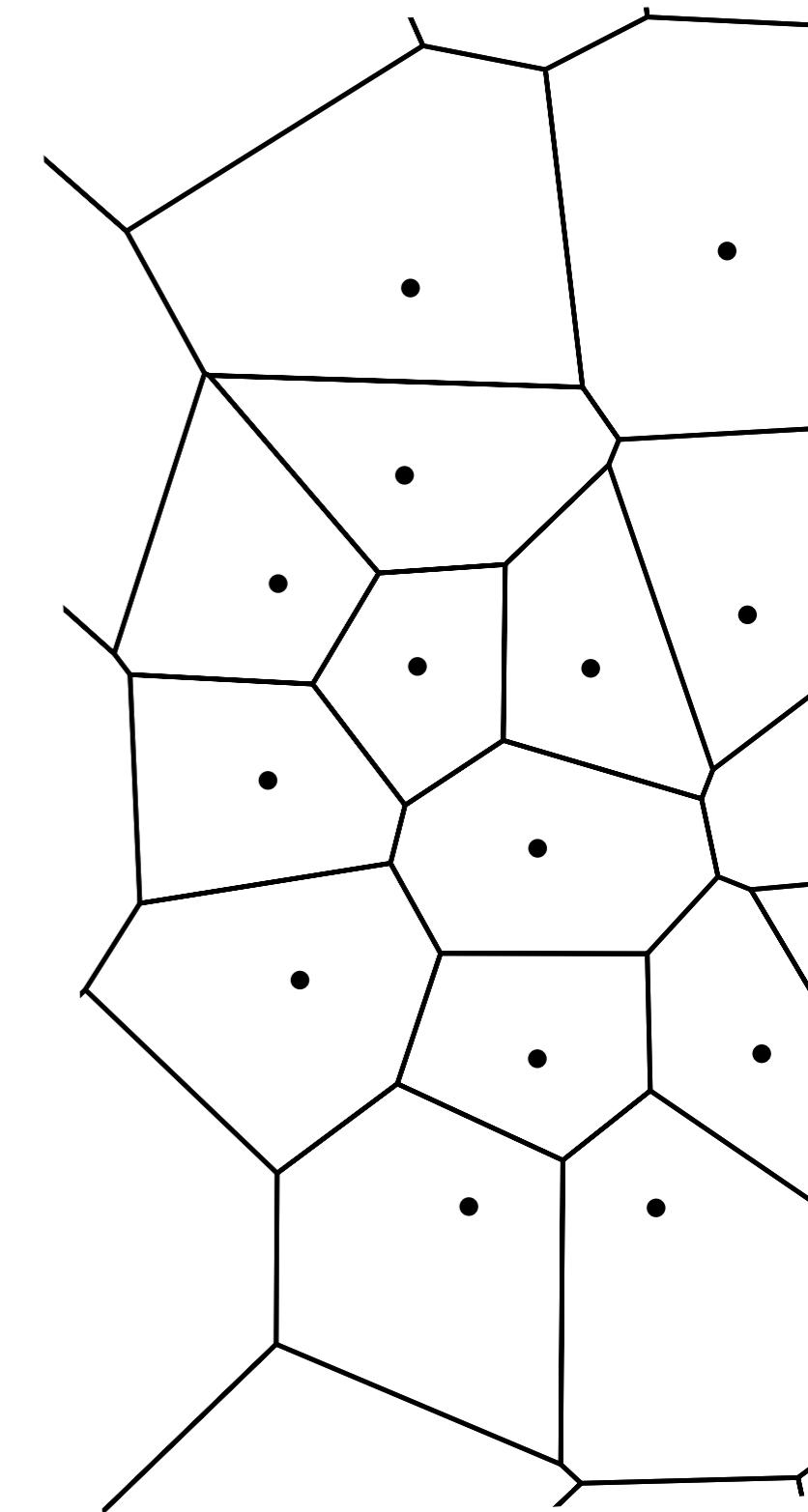
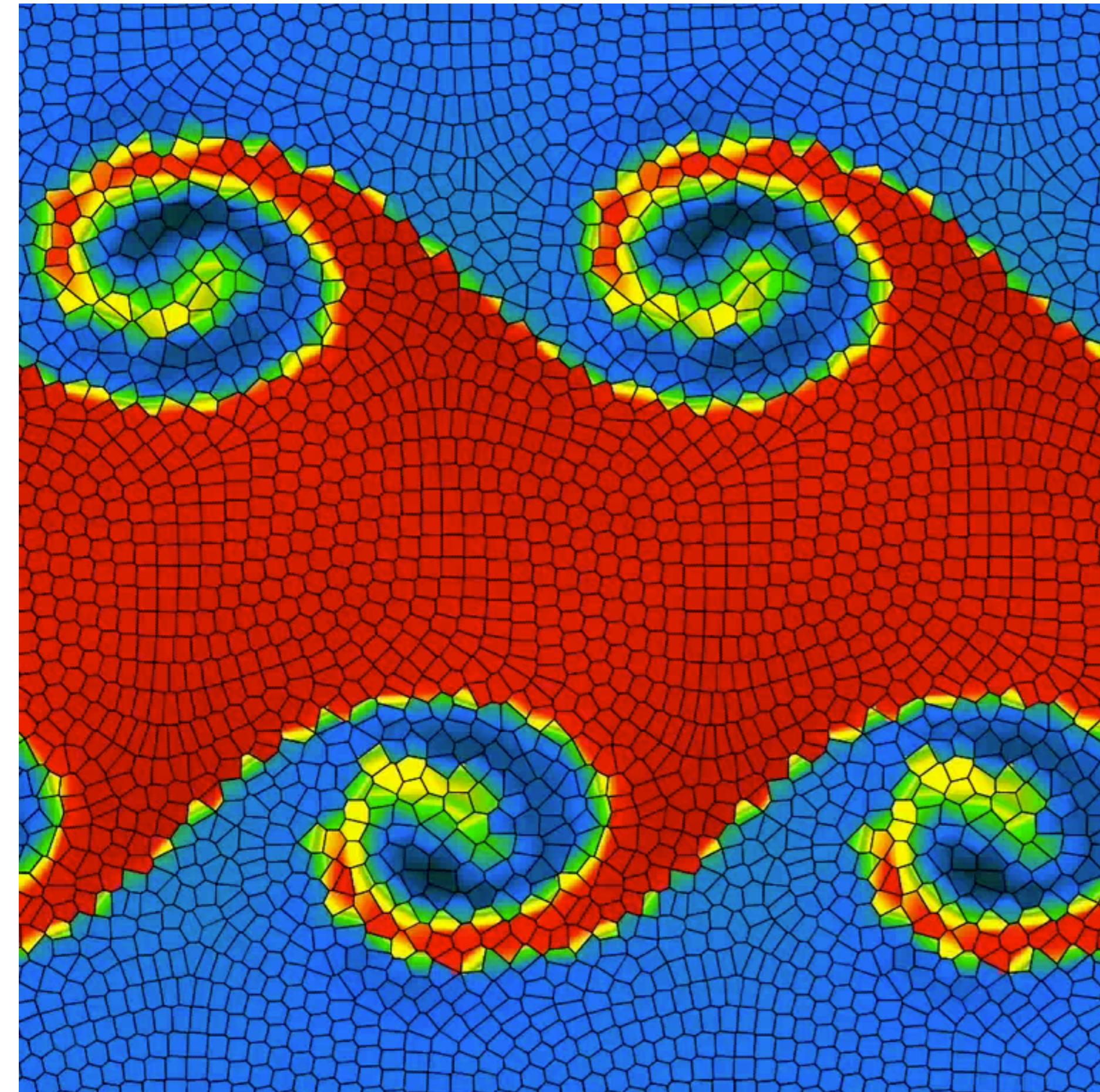
Volker Springel (2010)



Pakmor+ 2011, 2013, 2016
6

THE MOVING MESH CODE AREPO

Volker Springel (2010)



Pakmor+ 2011, 2013, 2016
6

BRAGINSKII VISCOSITY IN AREPO

$$\frac{\partial \rho \mathbf{v}}{\partial t} = -\nabla \cdot \boldsymbol{\Pi}, \quad \boldsymbol{\Pi} = -\Delta p \left(\mathbf{b}\mathbf{b} - \frac{1}{3} \right),$$

$$\frac{\partial E}{\partial t} = -\nabla \cdot (\boldsymbol{\Pi} \cdot \mathbf{v}), \quad \Delta p = \rho \nu_{\parallel} (3\mathbf{b}\mathbf{b} : \nabla \mathbf{v} - \nabla \cdot \mathbf{v}).$$

SECOND ORDER ACCURATE SUPER Timestepping

$$\Delta t \leq C \frac{(\Delta x)^2}{2d\nu_{\parallel}}$$

$$\tau = \frac{\Delta t}{4} (s^2 + s - 2)$$

$$\frac{\partial \mathbf{v}}{\partial t} = L(T, \mathbf{v})$$

$$\frac{\partial E}{\partial t} = \nabla \cdot \mathbf{F}_E(T, \mathbf{v})$$

Velocity update

$$\mathbf{Y}_0 = \mathbf{v}^n,$$

$$\mathbf{Y}_1 = \mathbf{Y}_0 + \tilde{\mu}_1 \tau \mathbf{L}(T^n, \mathbf{Y}_0),$$

$$\begin{aligned} \mathbf{Y}_j &= \mu_j \mathbf{Y}_{j-1} + \nu_j \mathbf{Y}_{j-2} + (1 - \mu_j - \nu_j) \mathbf{Y}_0 \\ &\quad + \tilde{\mu}_j \tau \mathbf{L}(T^n, \mathbf{Y}_{j-1}) + \tilde{\gamma}_j \tau \mathbf{L}(T^n, \mathbf{Y}_0) \quad \text{for } 2 \leq j \leq s \end{aligned}$$

$$\mathbf{v}^{n+1} = \mathbf{Y}_s$$

Energy update

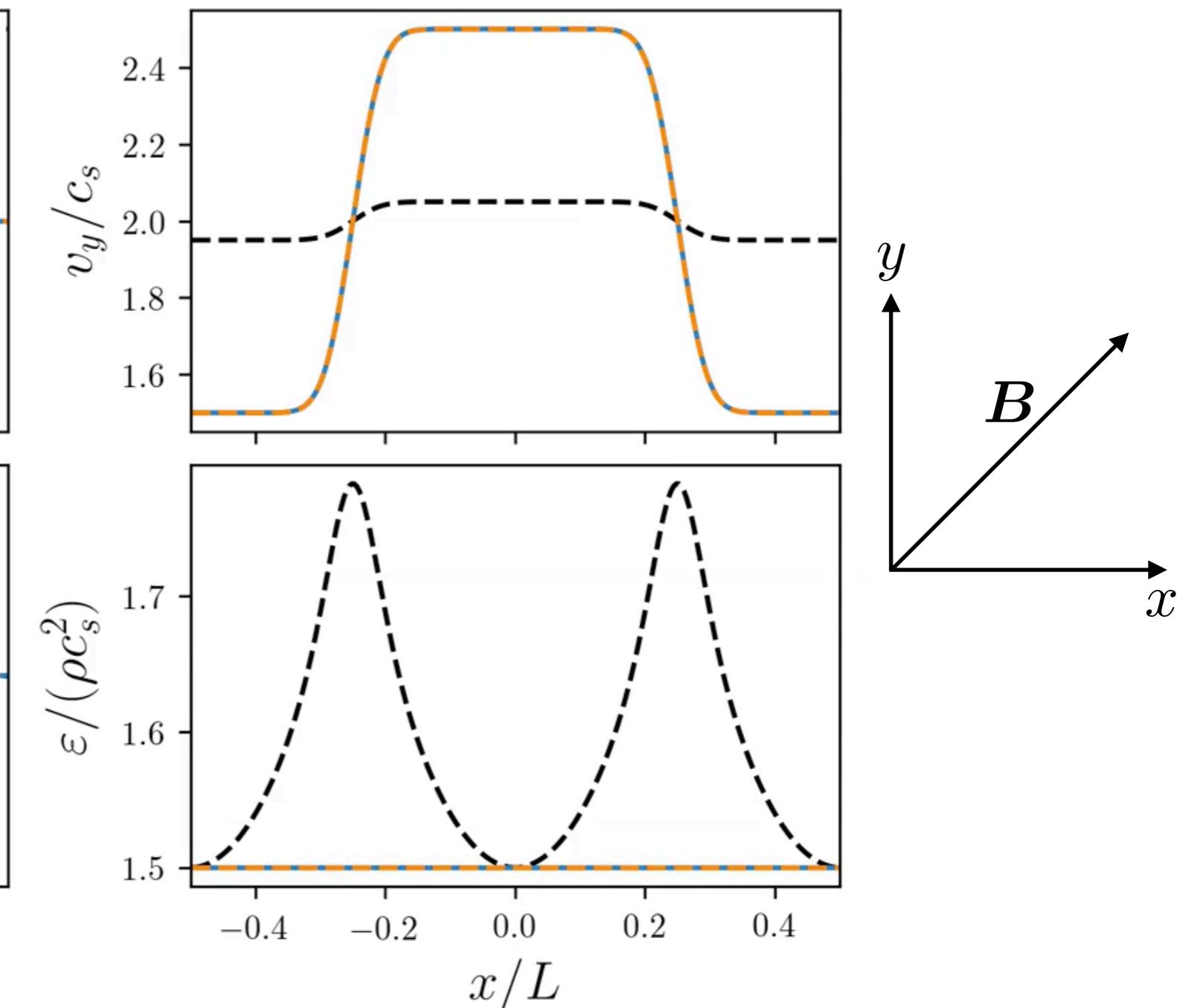
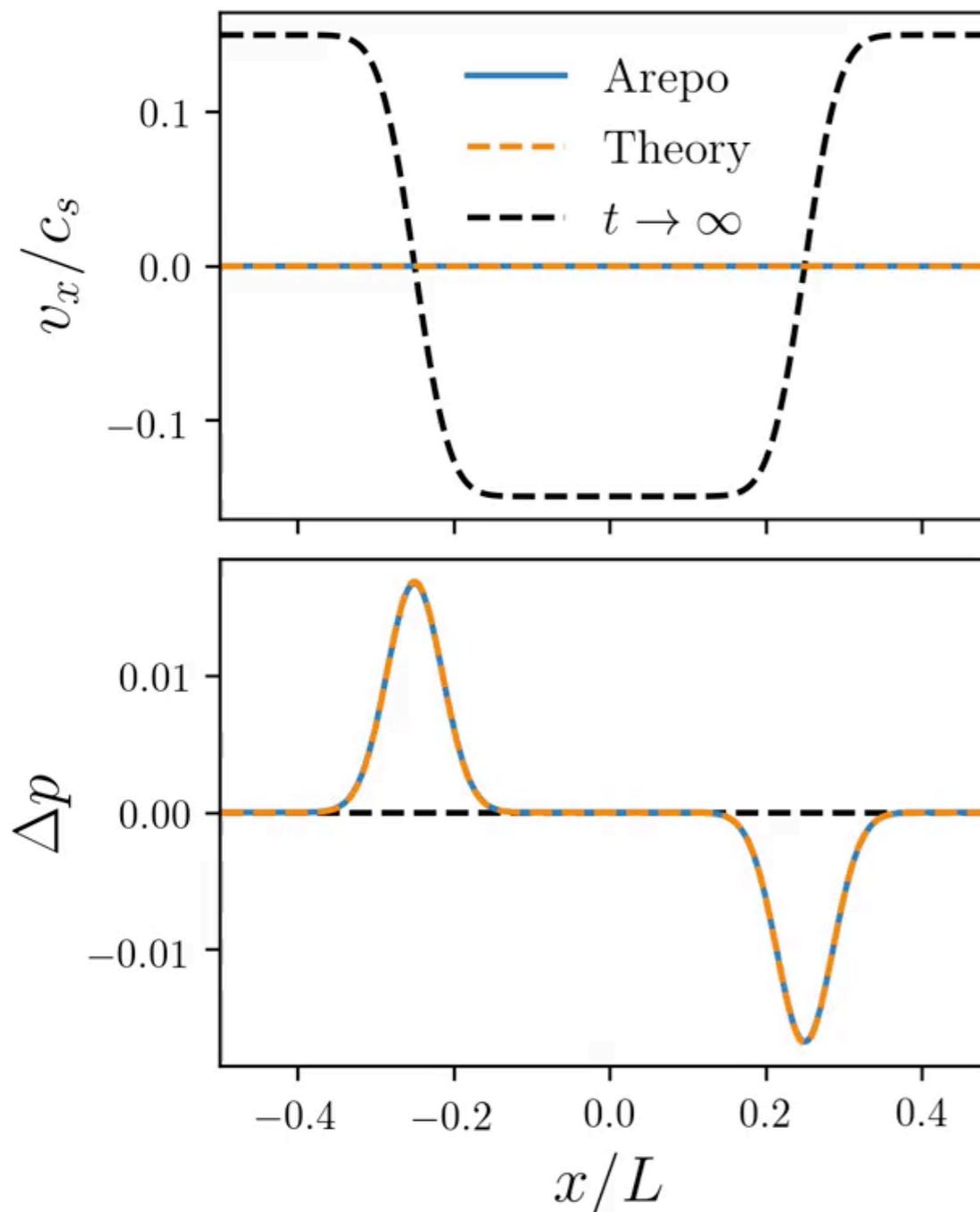
$$E^{n+1} = \frac{\tau}{2} [\nabla \cdot \mathbf{F}_E(T^n, \mathbf{v}^n) + \nabla \cdot \mathbf{F}_E(T^n, \mathbf{v}^{n+1})]$$

RKL2 STS theory in Meyer+ 2012. See also Vaidya+ 2017

$$v_x(x, t) = -c_s \sum_{n=0}^{\infty} \frac{3a_n}{10} \cos(k_n x) (1 - e^{-\gamma_n t})$$

$$v_y(x, t) = c_s \sum_{n=0}^{\infty} \frac{a_n}{10} \cos(k_n x) (1 + 9e^{-\gamma_n t})$$

$$c_s t / L = 0.0$$



$$\Delta p(x, t) = -\frac{3\rho c_s \nu_{\parallel}}{2} \sum_{n=1}^{\infty} k_n a_n \sin(k_n x)$$

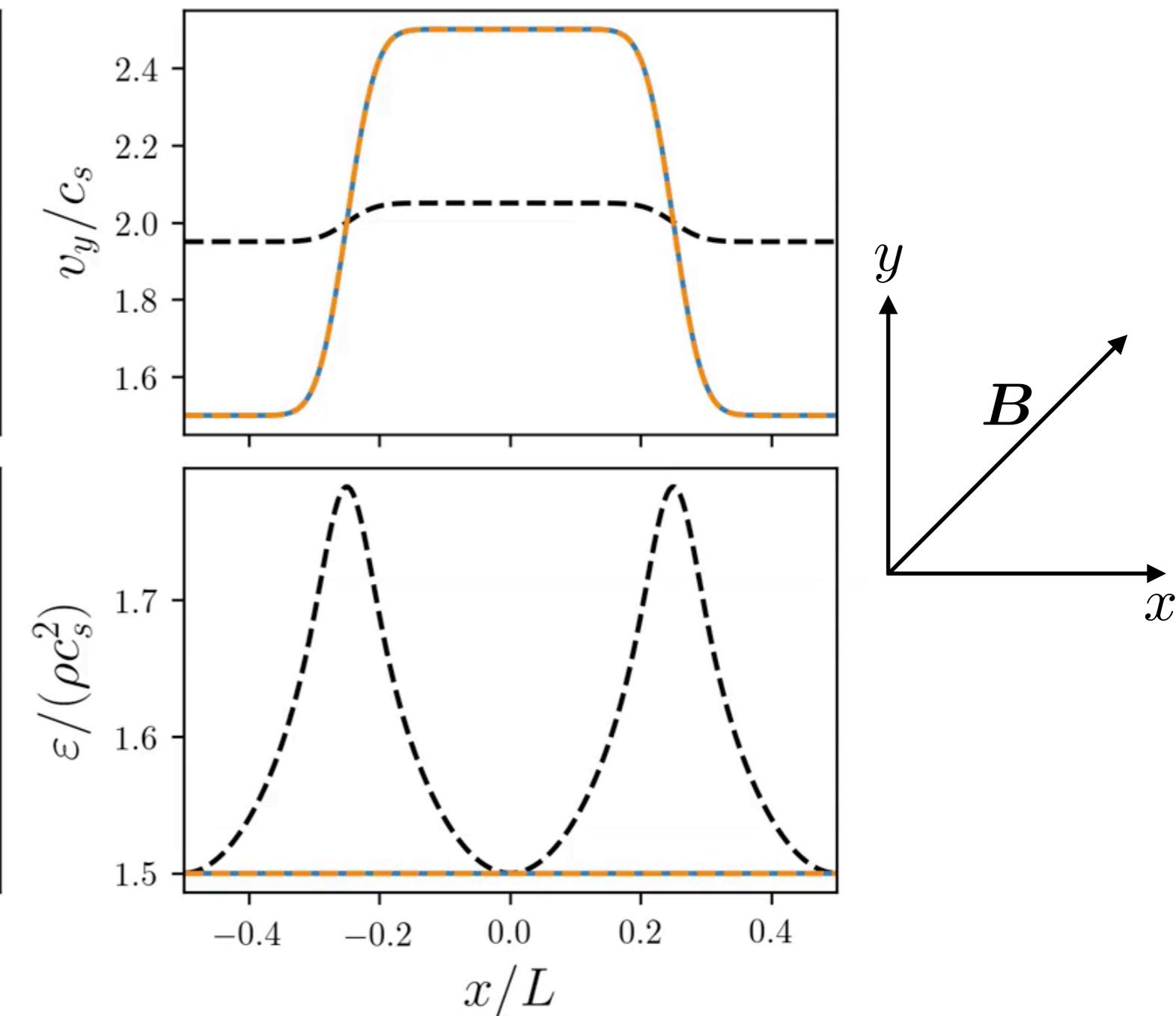
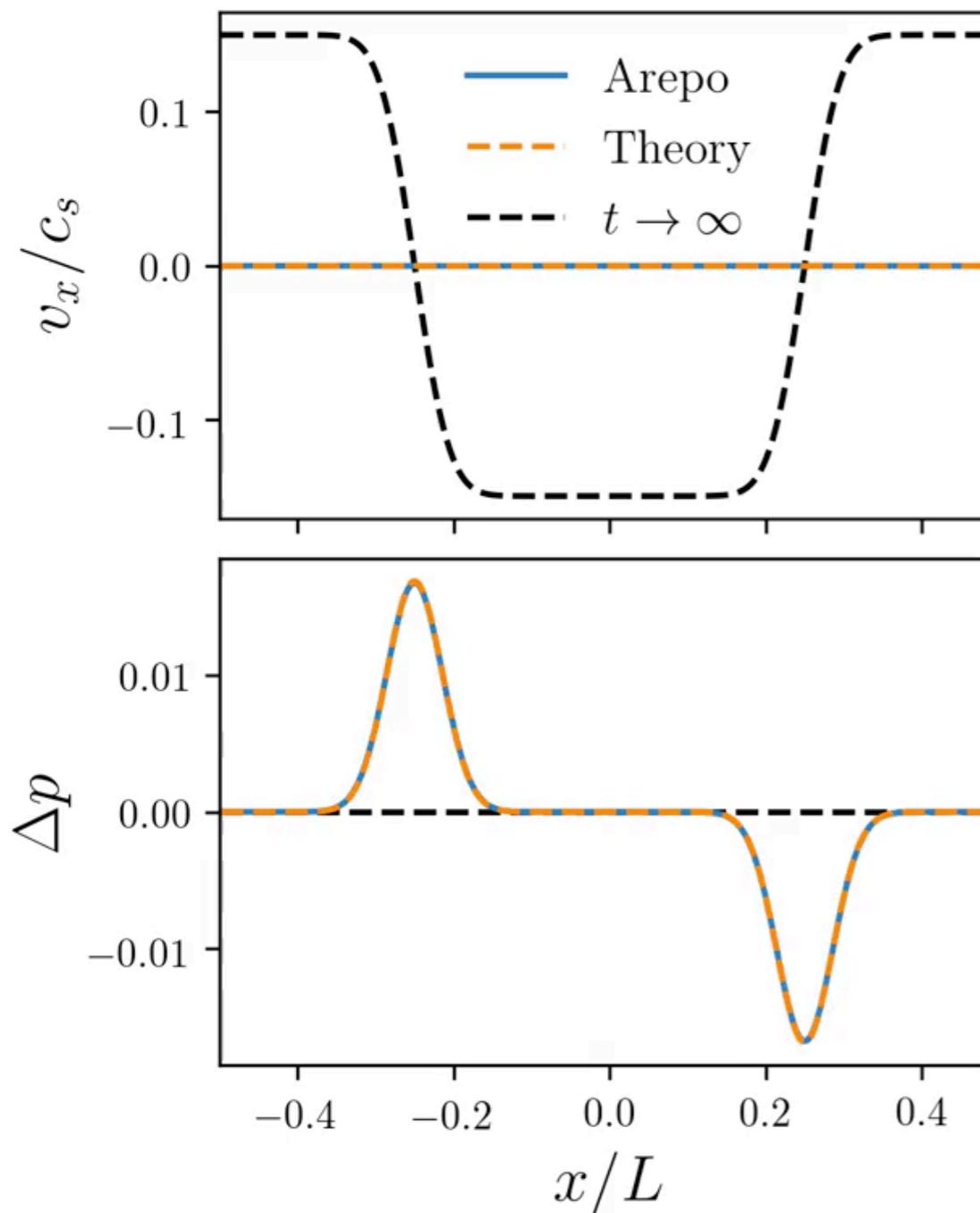
$$\gamma_n = \frac{5\nu_{\parallel}}{6} k_n^2$$

$$\begin{aligned} \varepsilon(t) = \varepsilon_0 + \frac{9\rho c_s^2}{10} \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} a_n a_m \frac{\sqrt{\gamma_n \gamma_m}}{\gamma_n + \gamma_m} \times \\ \sin(k_n x) \sin(k_m x) \left(1 - e^{-(\gamma_n + \gamma_m)t} \right), \end{aligned}$$

$$v_x(x, t) = -c_s \sum_{n=0}^{\infty} \frac{3a_n}{10} \cos(k_n x) (1 - e^{-\gamma_n t})$$

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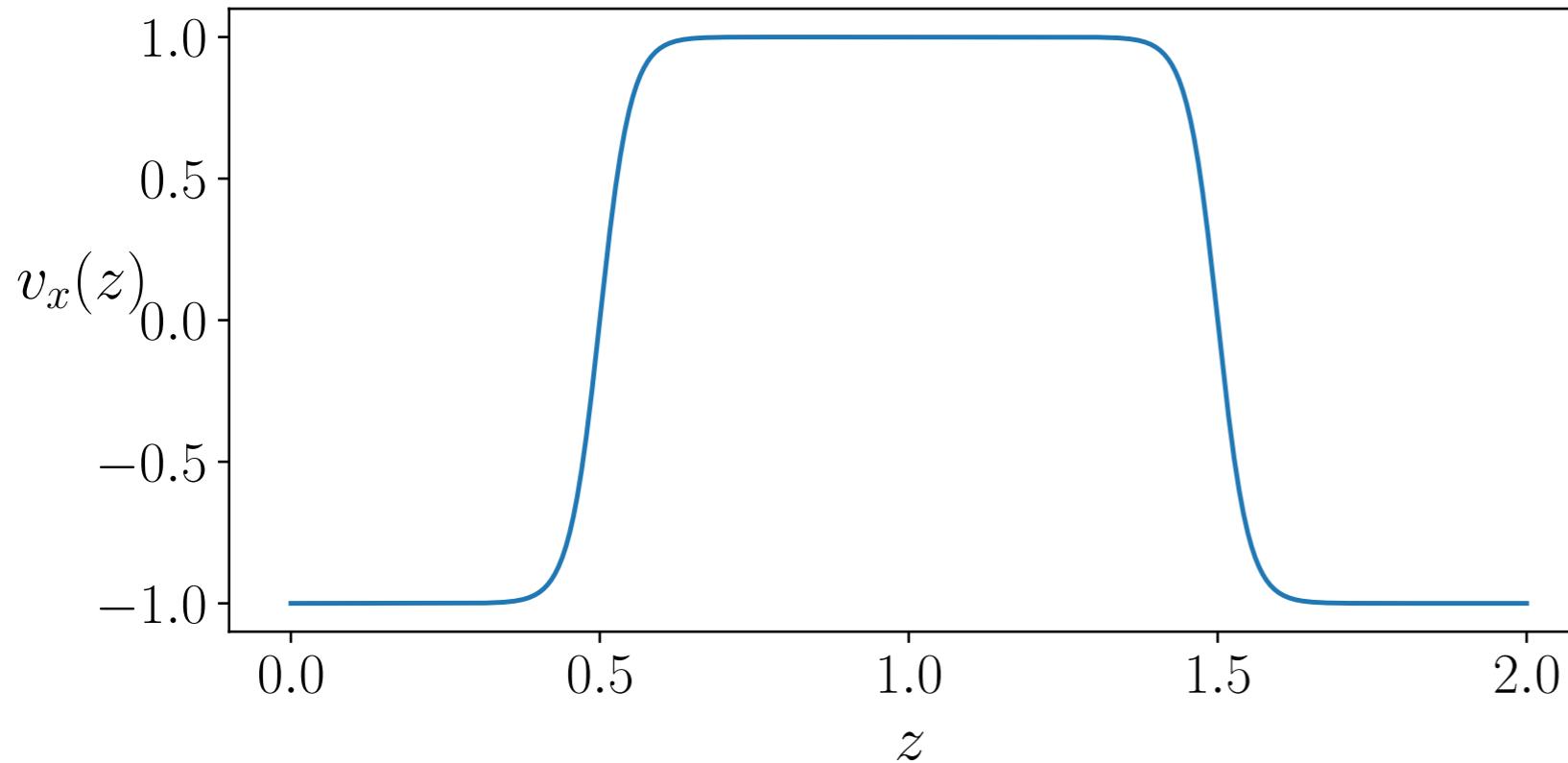


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KELVIN-HELMHOLTZ INSTABILITY WITH BRAGINSKII VISCOSITY



Smooth equilibrium necessary for convergence of KHI.
See e.g. McNally+ 2012 and Lecoanet+ 2016.

LINEAR THEORY FOR VISCOUS KELVIN-HELMHOLTZ INSTABILITY

$$-i\omega \frac{\delta\rho}{\rho} = -ik \left(v_0 \frac{\delta\rho}{\rho} + \delta v_x \right) - \frac{\partial \delta v_z}{\partial z} .$$

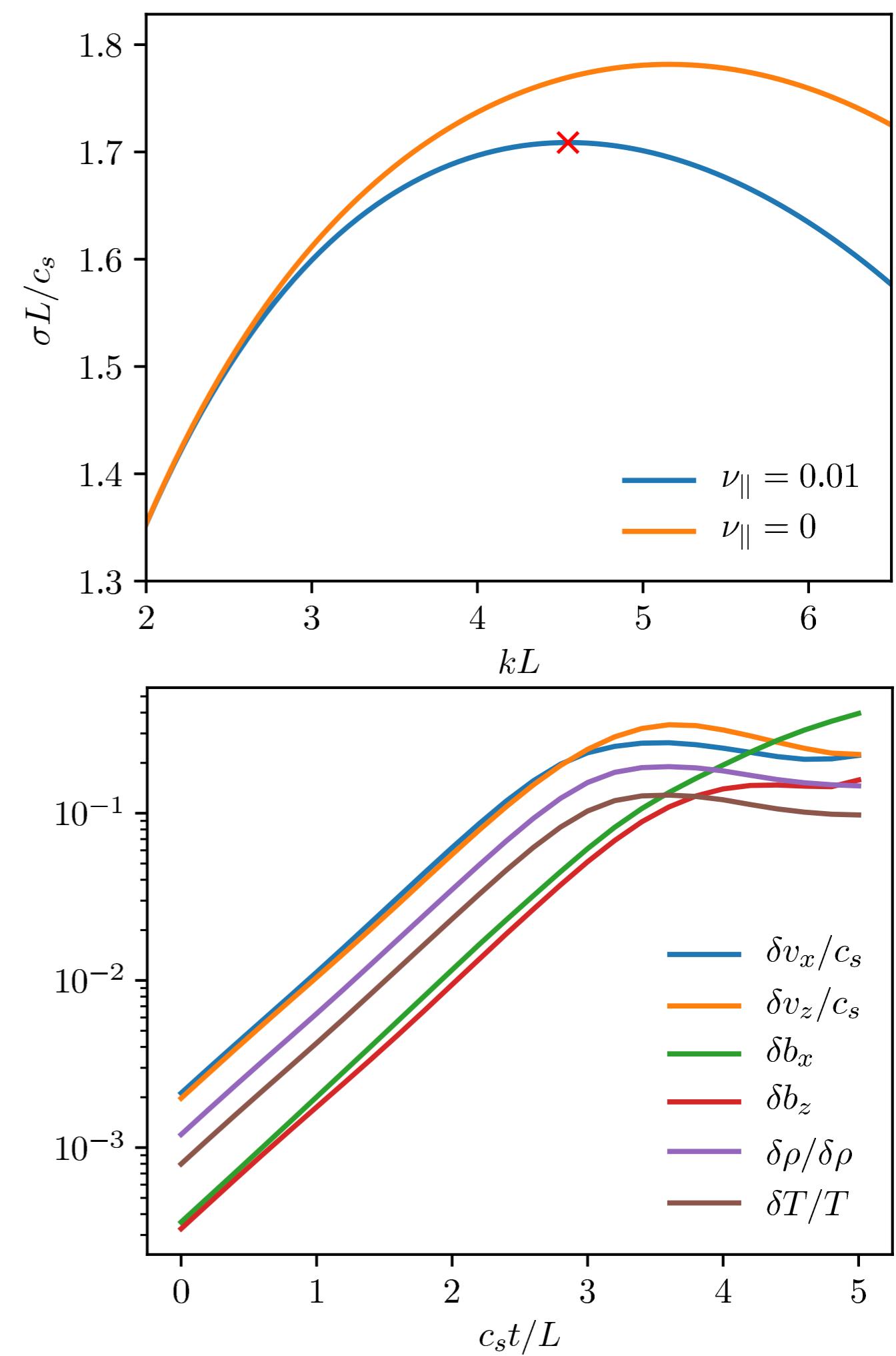
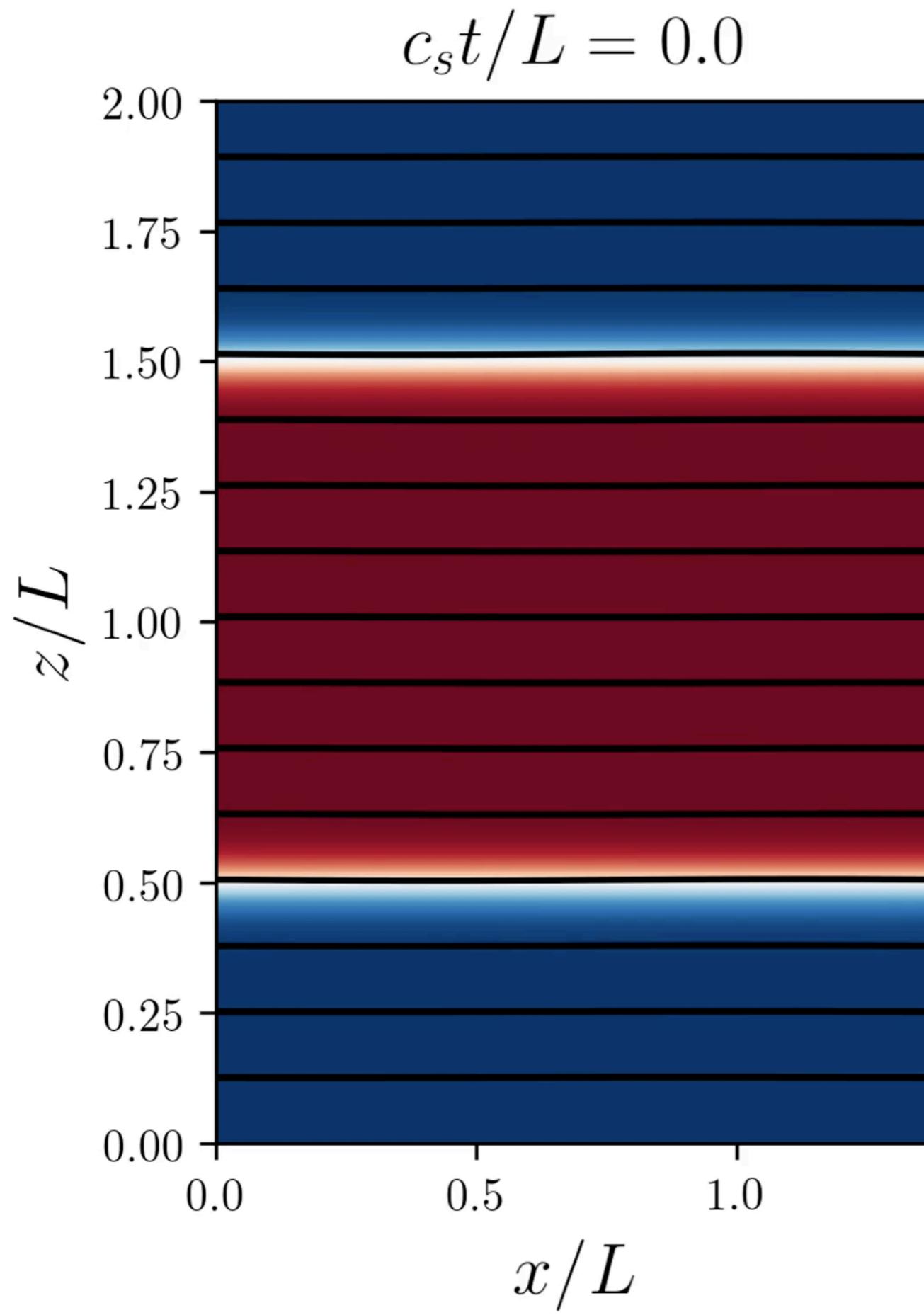
$$-i\omega \frac{\delta A}{B} = -ik v_0 \frac{\delta A}{B} + \delta v_z ,$$

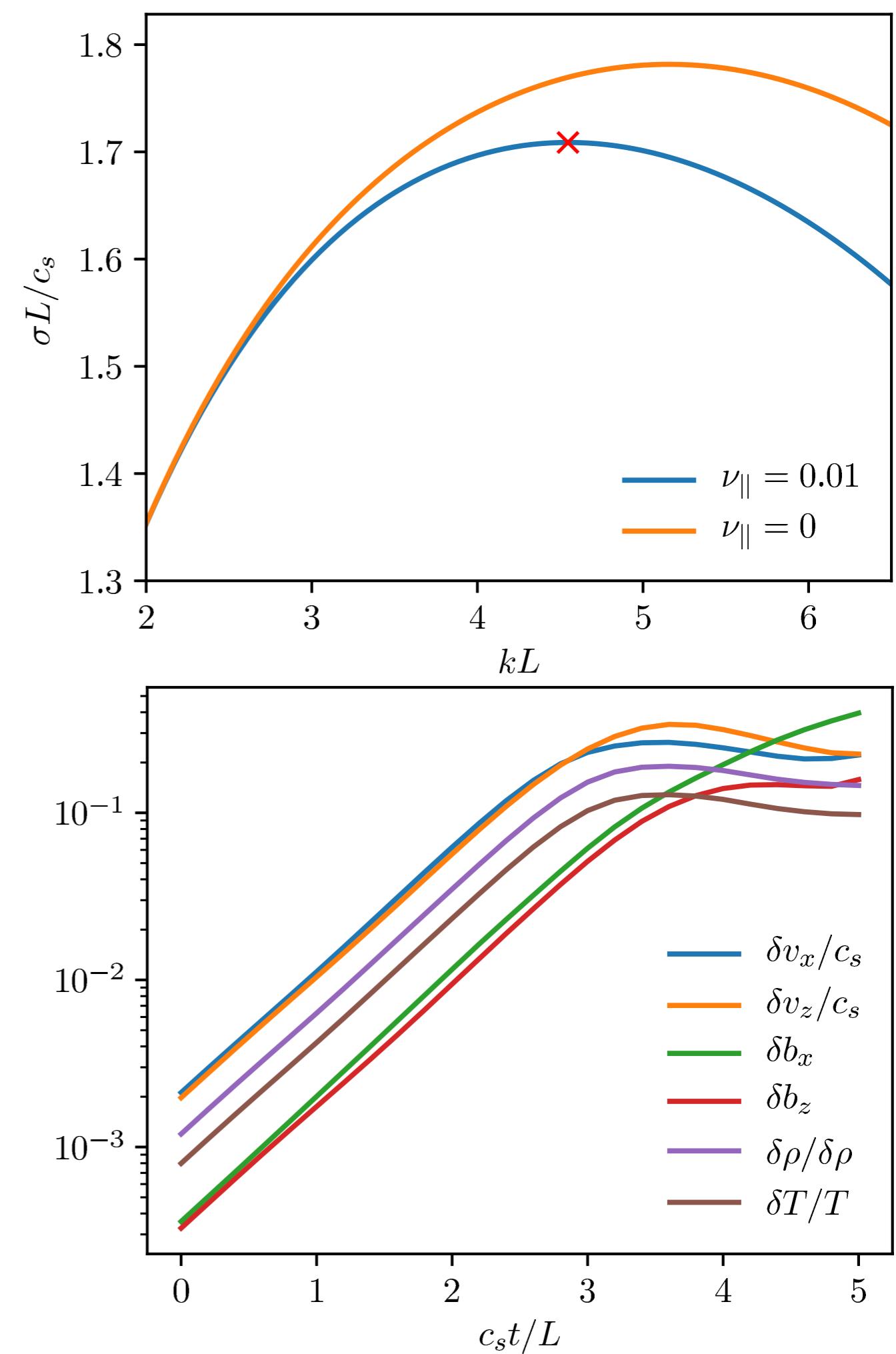
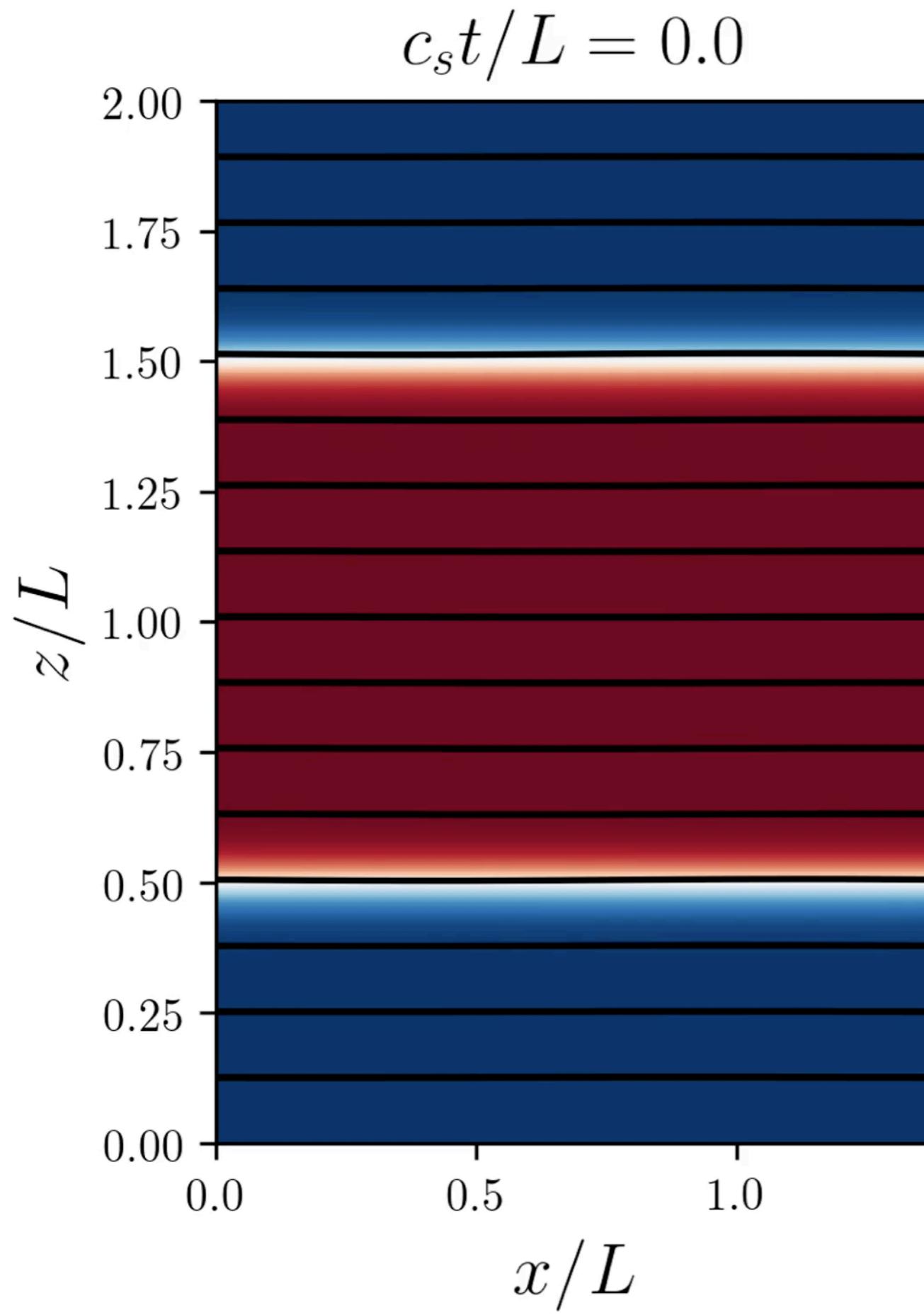
$$-i\omega \delta v_x = -ik v_0 \delta v_x - \frac{\partial v_0}{\partial z} \delta v_z - ik \frac{\delta p}{\rho} - \nu_{\parallel} \left(\frac{4}{3} k^2 \delta v_x + 2k^2 \frac{\partial v_0}{\partial z} \frac{\delta A}{B} + ik \frac{2}{3} \frac{\partial \delta v_z}{\partial z} \right) ,$$

$$-i\omega \delta v_z = -ik v_0 \delta v_z - \frac{1}{\rho} \frac{\partial \delta p}{\partial z} + v_a^2 \left(\frac{\partial^2}{\partial z^2} - k^2 \right) \frac{\delta A}{B} - \nu_{\parallel} \left(ik \frac{2}{3} \frac{\partial \delta v_x}{\partial z} + ik \frac{\partial^2 v_0}{\partial z^2} \frac{\delta A}{B} + ik \frac{\partial v_0}{\partial z} \frac{\partial}{\partial z} \frac{\delta A}{B} - \frac{1}{3} \frac{\partial^2 \delta v_z}{\partial z^2} \right) ,$$

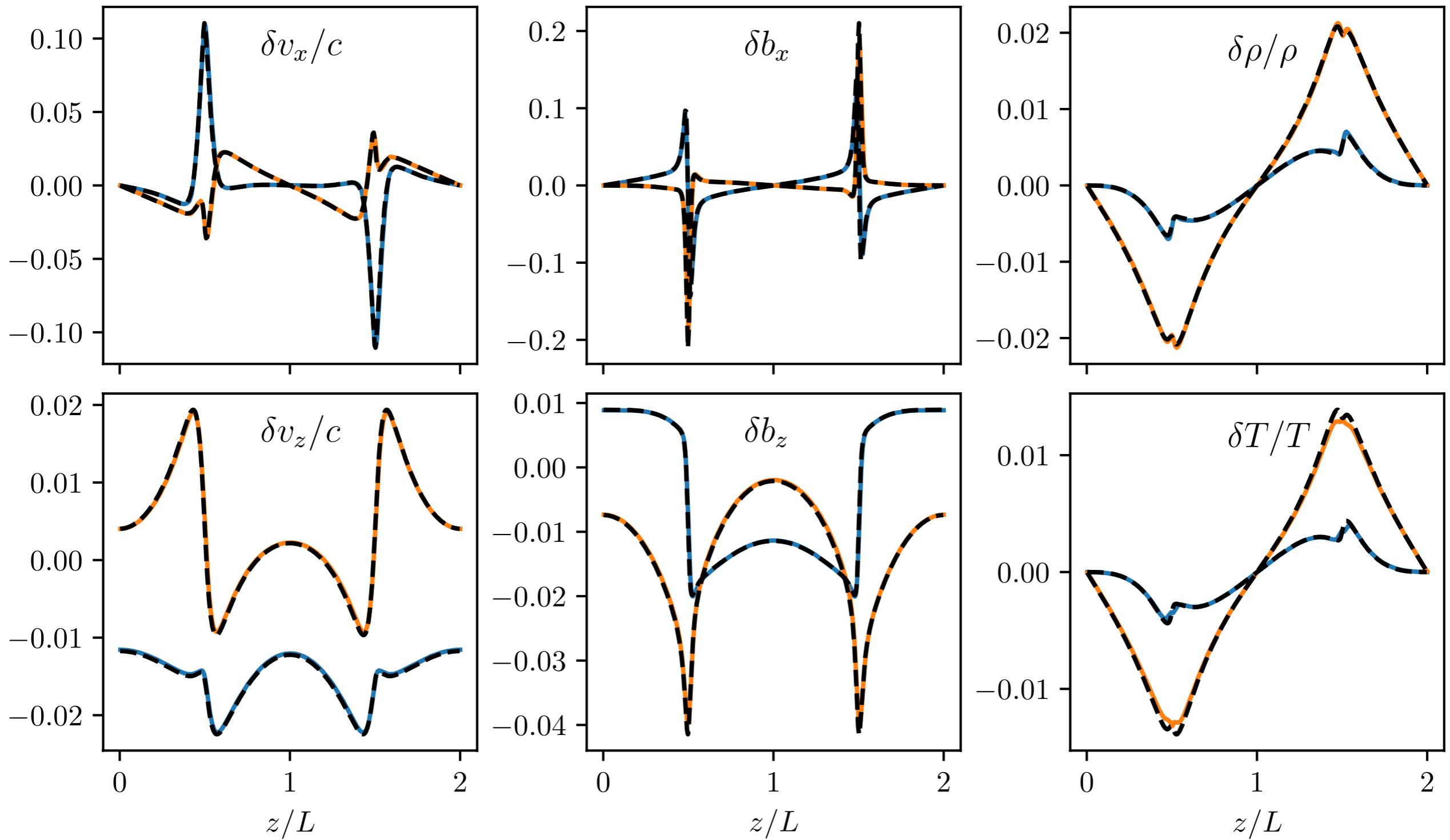
$$-i\omega \frac{\delta T}{T} = -ik \left(v_0 \frac{\delta T}{T} + \frac{2}{3} \delta v_x \right) - \frac{2}{3} \frac{\partial \delta v_z}{\partial z}$$

Berlok & Pfrommer, 2019a, MNRAS
See also Suzuki+ (2013)

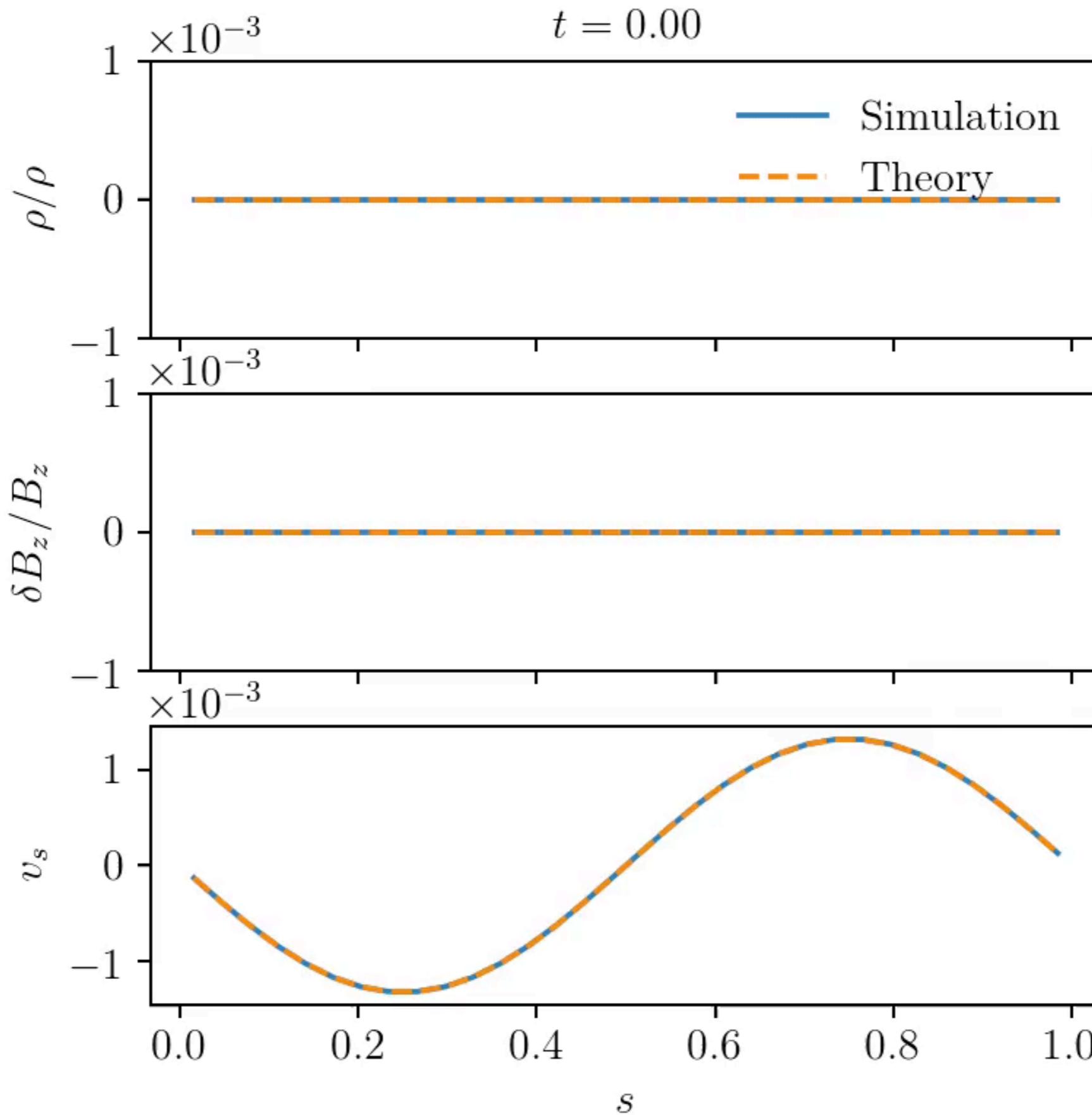




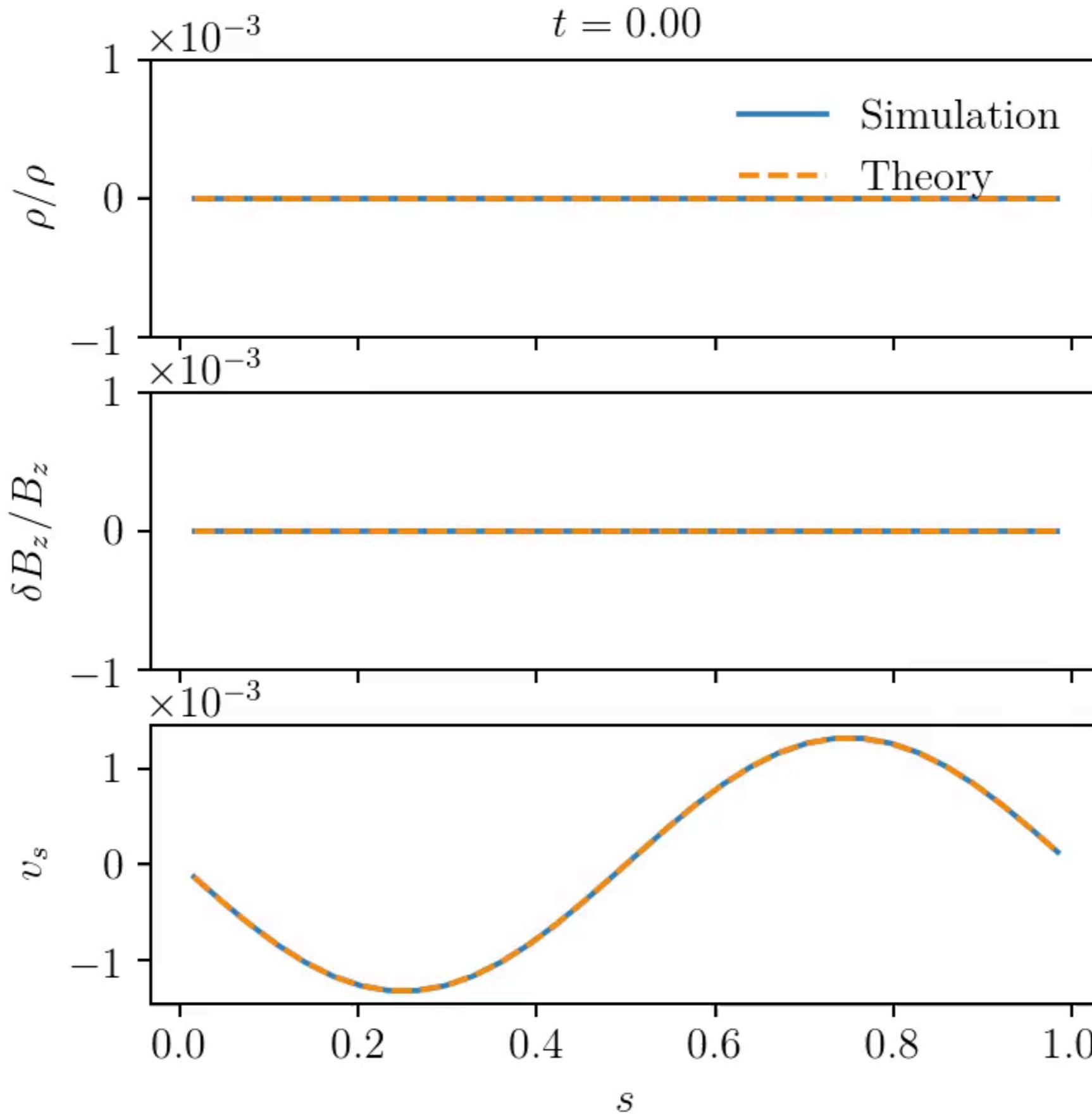
EIGENMODES OF THE INSTABILITY



DECAY OF 2D MAGNETO-SONIC WAVE



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DECAY OF 2D MAGNETO-SONIC WAVE

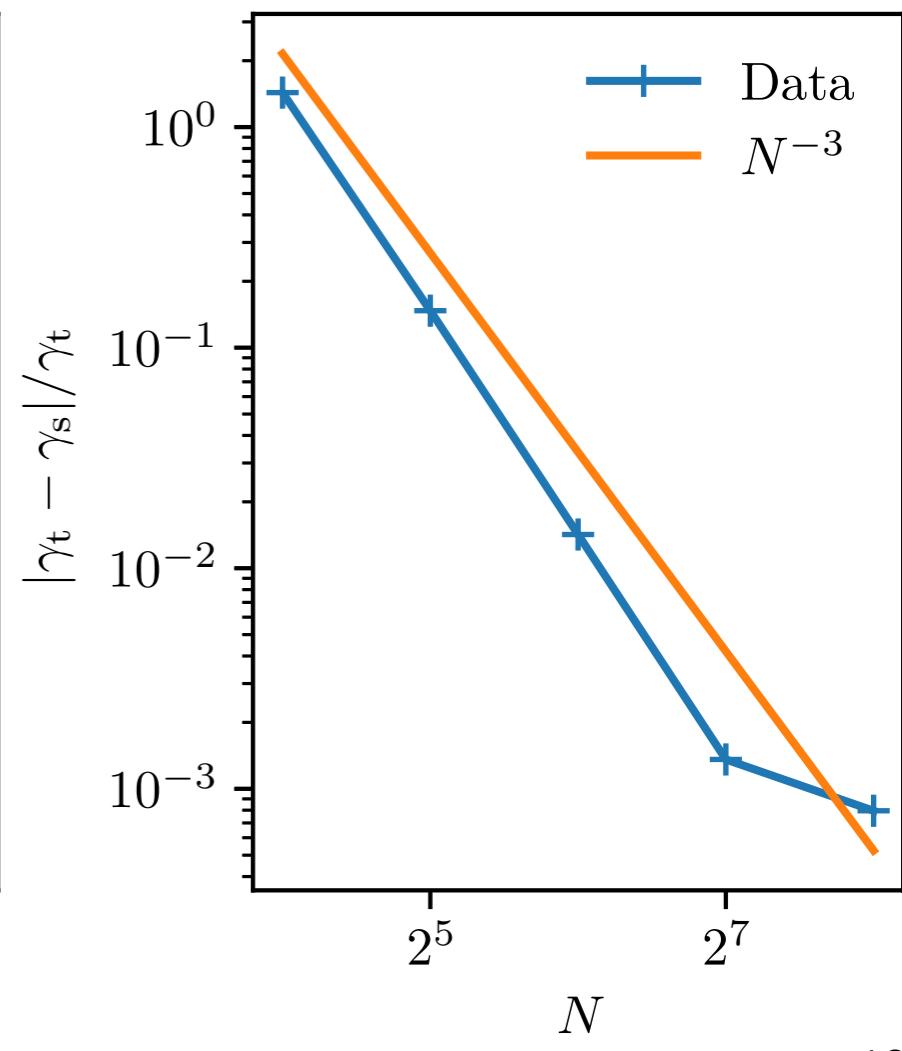
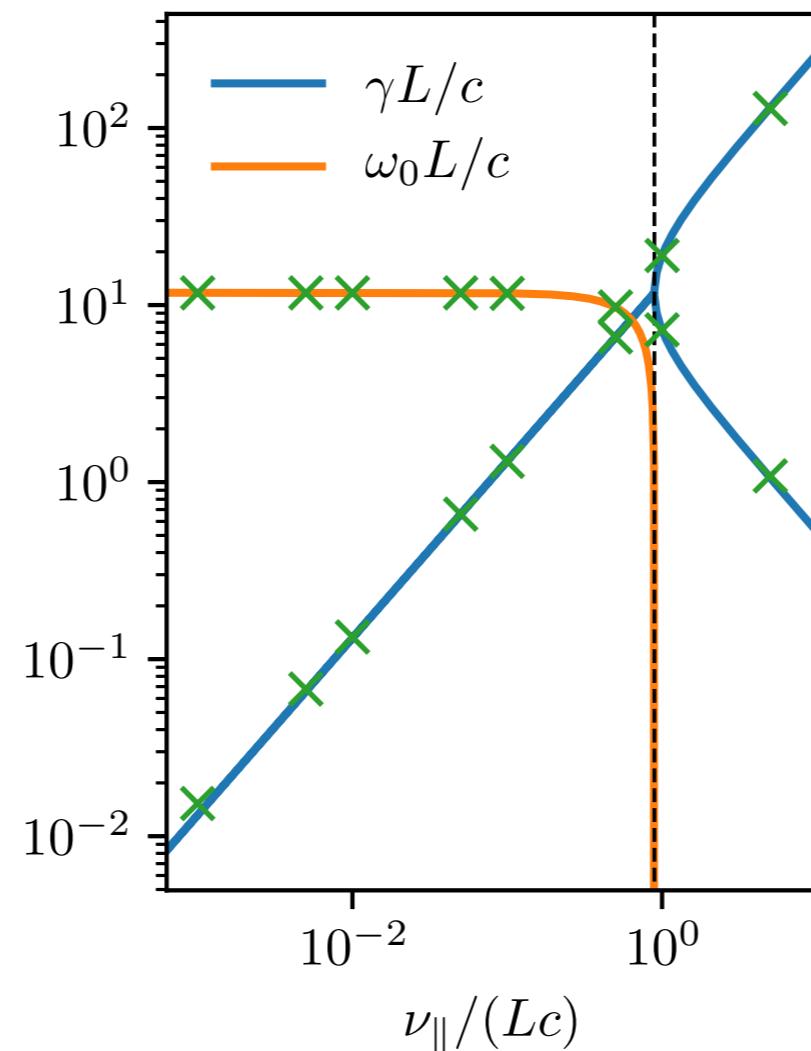
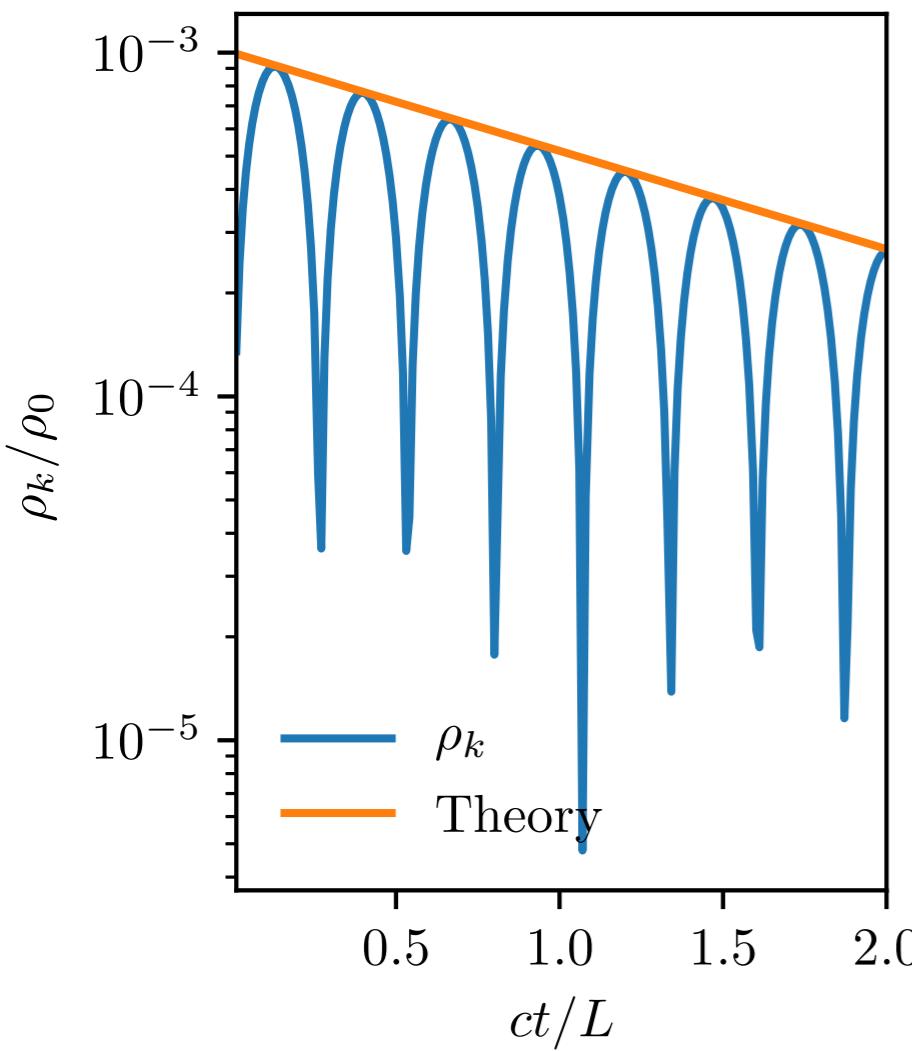
Alfvén waves

$$(\omega^2 - k_{\parallel}^2 v_a^2) = 0$$

Fast and slow magnetosonic waves

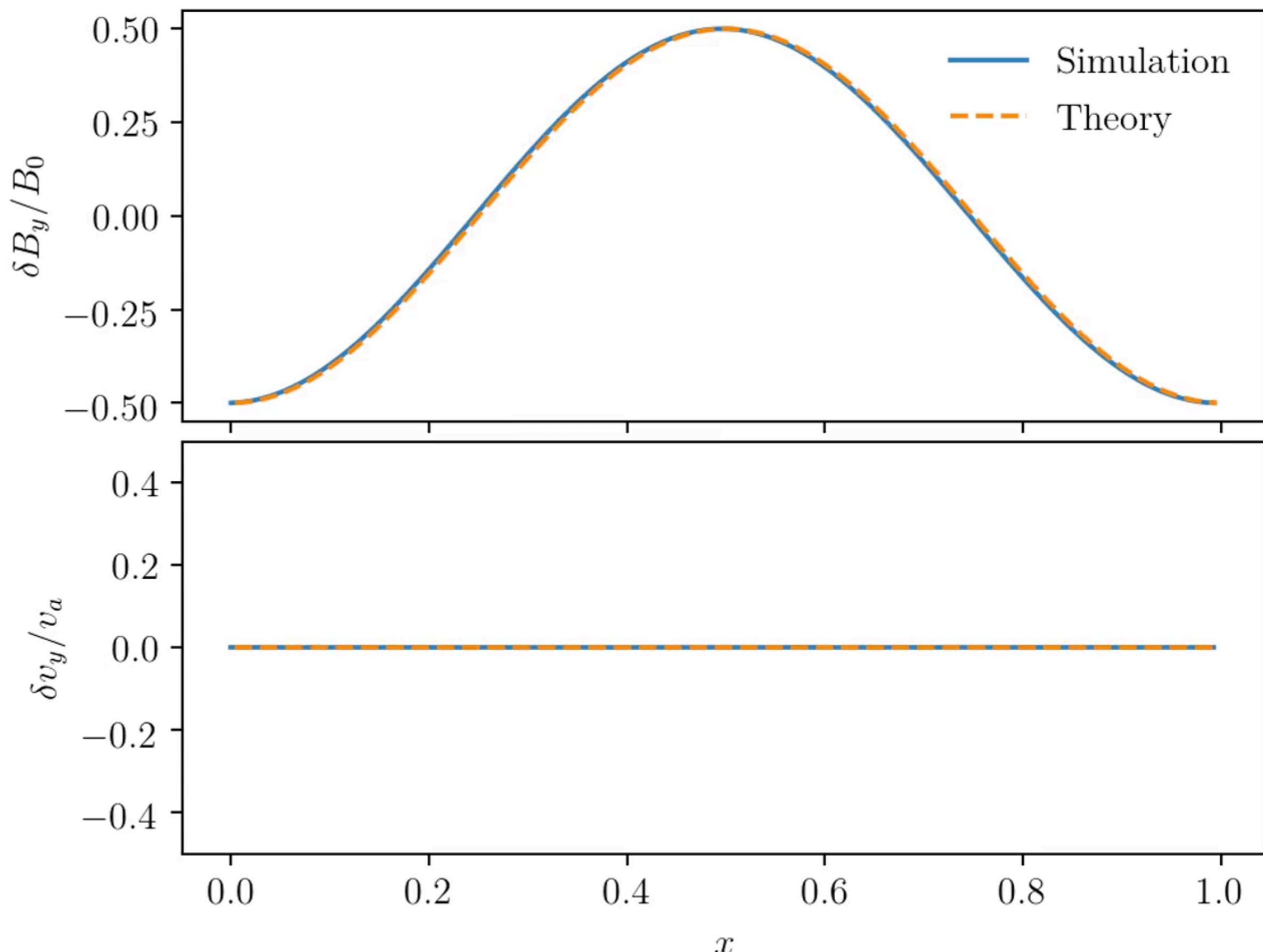
$$\begin{aligned} \omega^4 + i\omega^3 \frac{\nu_{\parallel}}{3} (4k_{\parallel}^2 + k_{\perp}^2) - \omega^2 k^2 (v_a^2 + \tilde{c}^2) \\ - i\omega \frac{\nu_{\parallel}}{3} k_{\parallel}^2 (9k_{\perp}^2 \tilde{c}^2 + 4k^2 v_a^2) + k_{\parallel}^2 k^2 \tilde{c}^2 v_a^2 = 0 \end{aligned}$$

if $k_{\parallel} = 0$ then $\omega = \pm k_{\perp} \sqrt{v_a^2 + \tilde{c}^2 - \left(\frac{k_{\perp} \nu_{\parallel}}{6}\right)^2 - i \frac{\nu_{\parallel}}{6} k_{\perp}^2}$



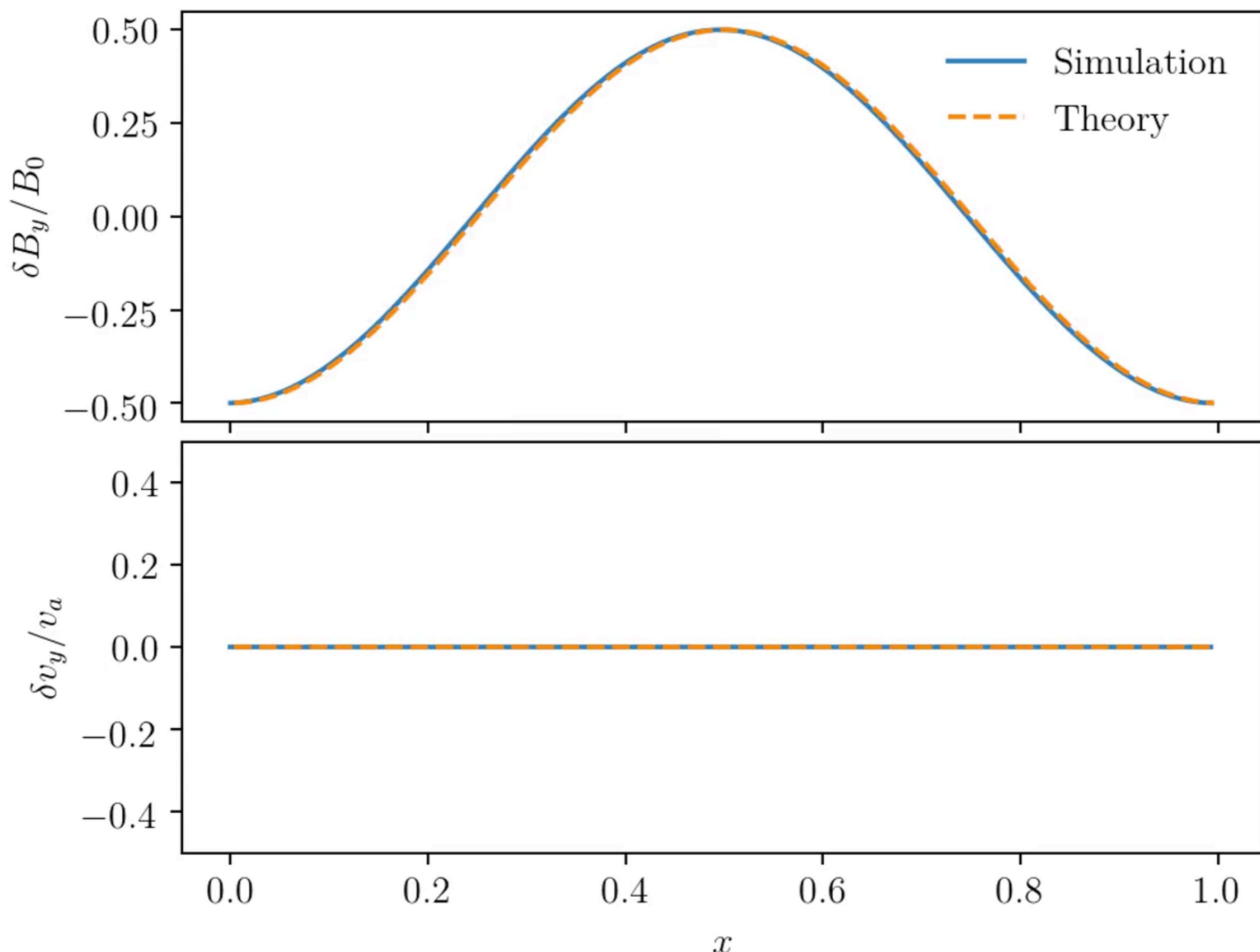
LINEARLY POLARIZED ALFVEN WAVE

$\omega_a t = 0.00$



LINEARLY POLARIZED ALFVEN WAVE

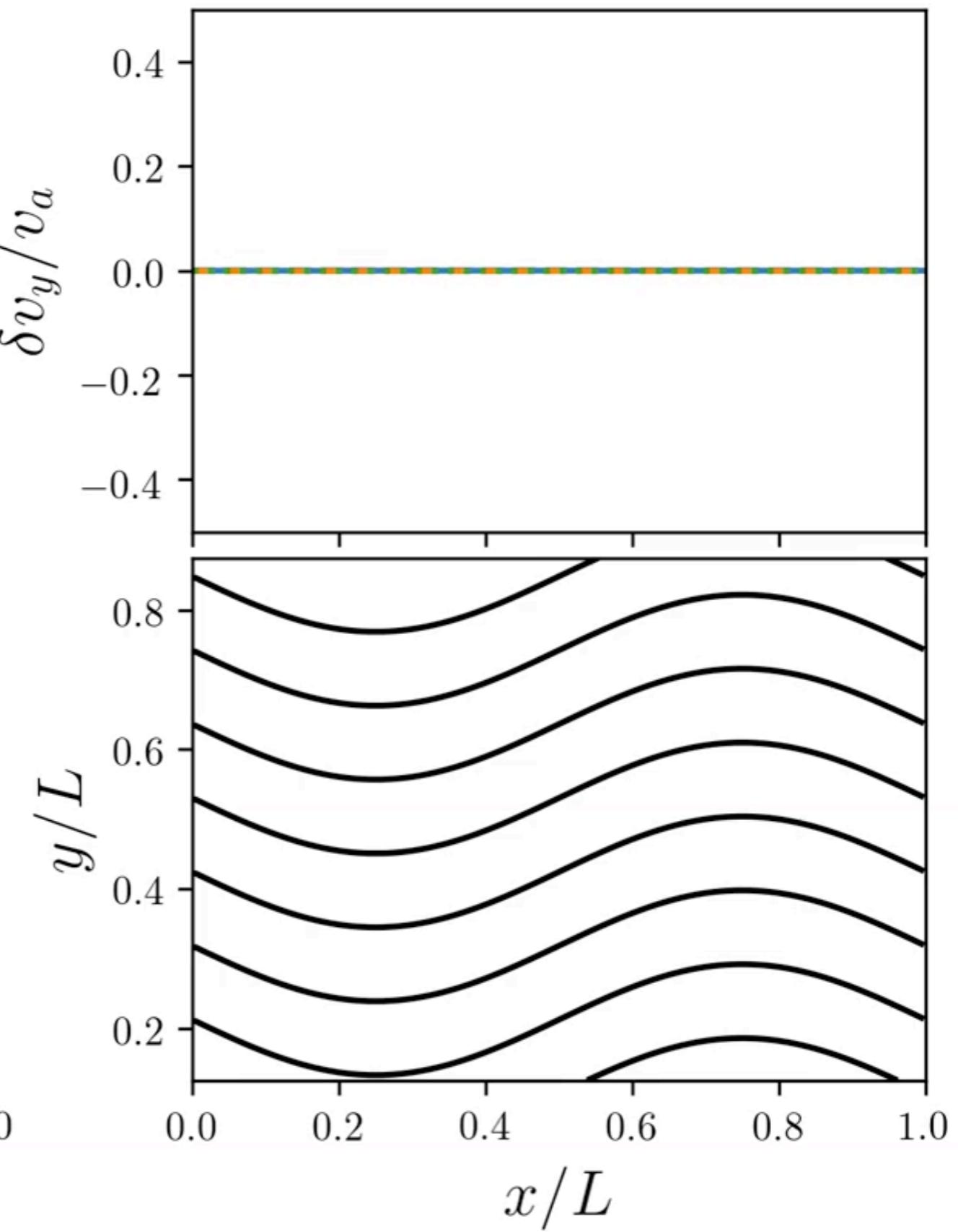
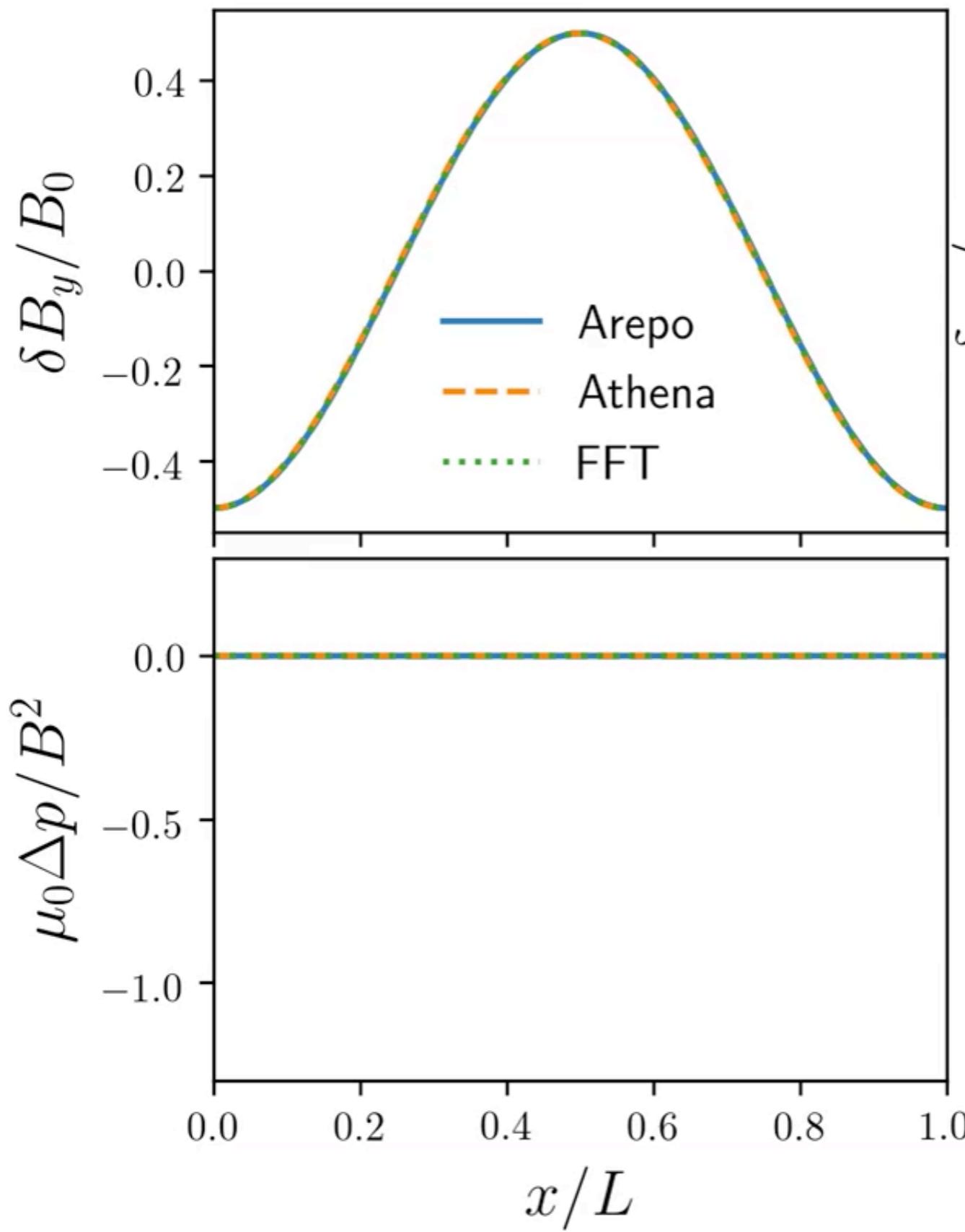
$\omega_a t = 0.00$



INTERRUPTION BY THE FIREHOSE INSTABILITY

$\omega_a t = 0.00$

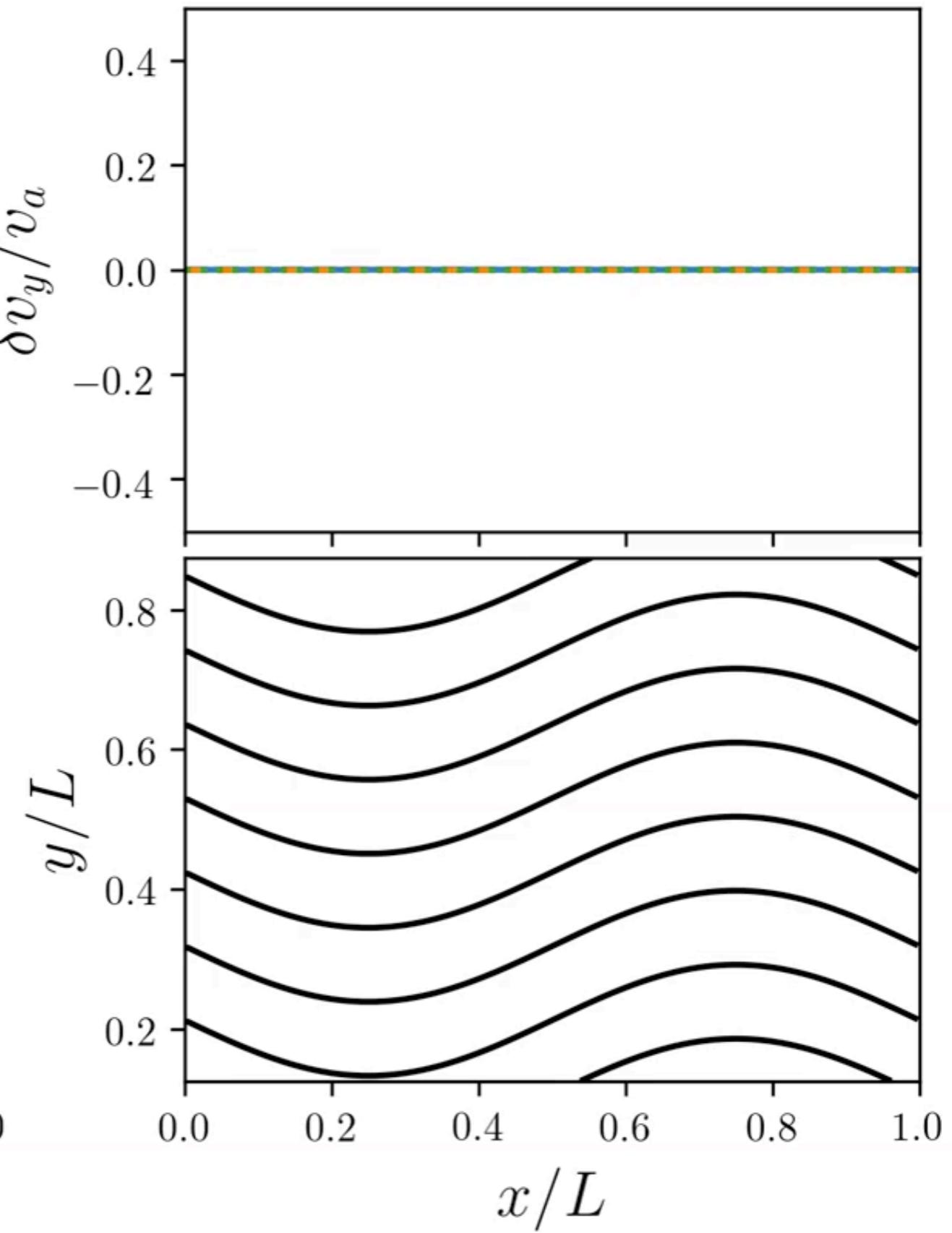
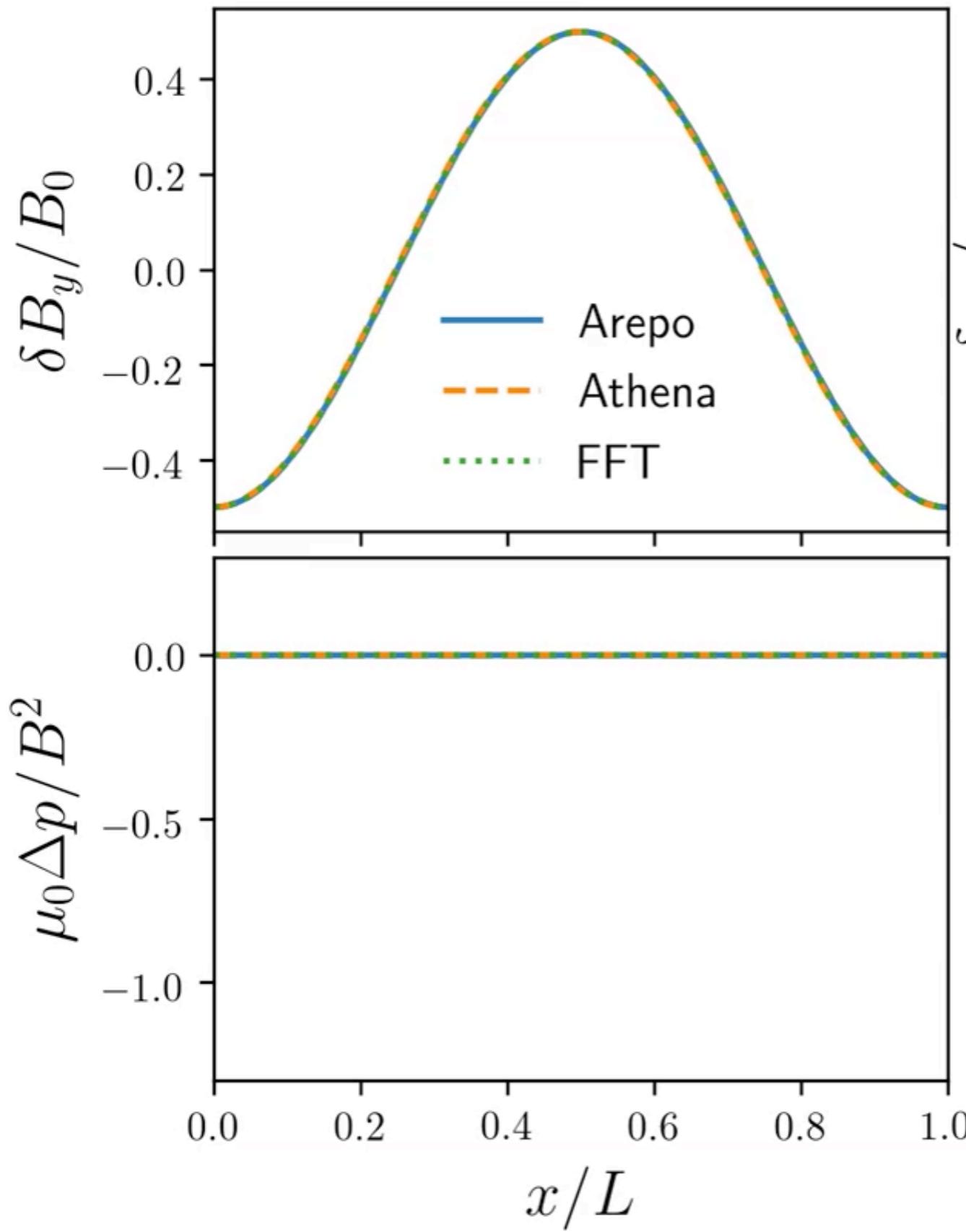
Squire+ 2016, 2017, 2019



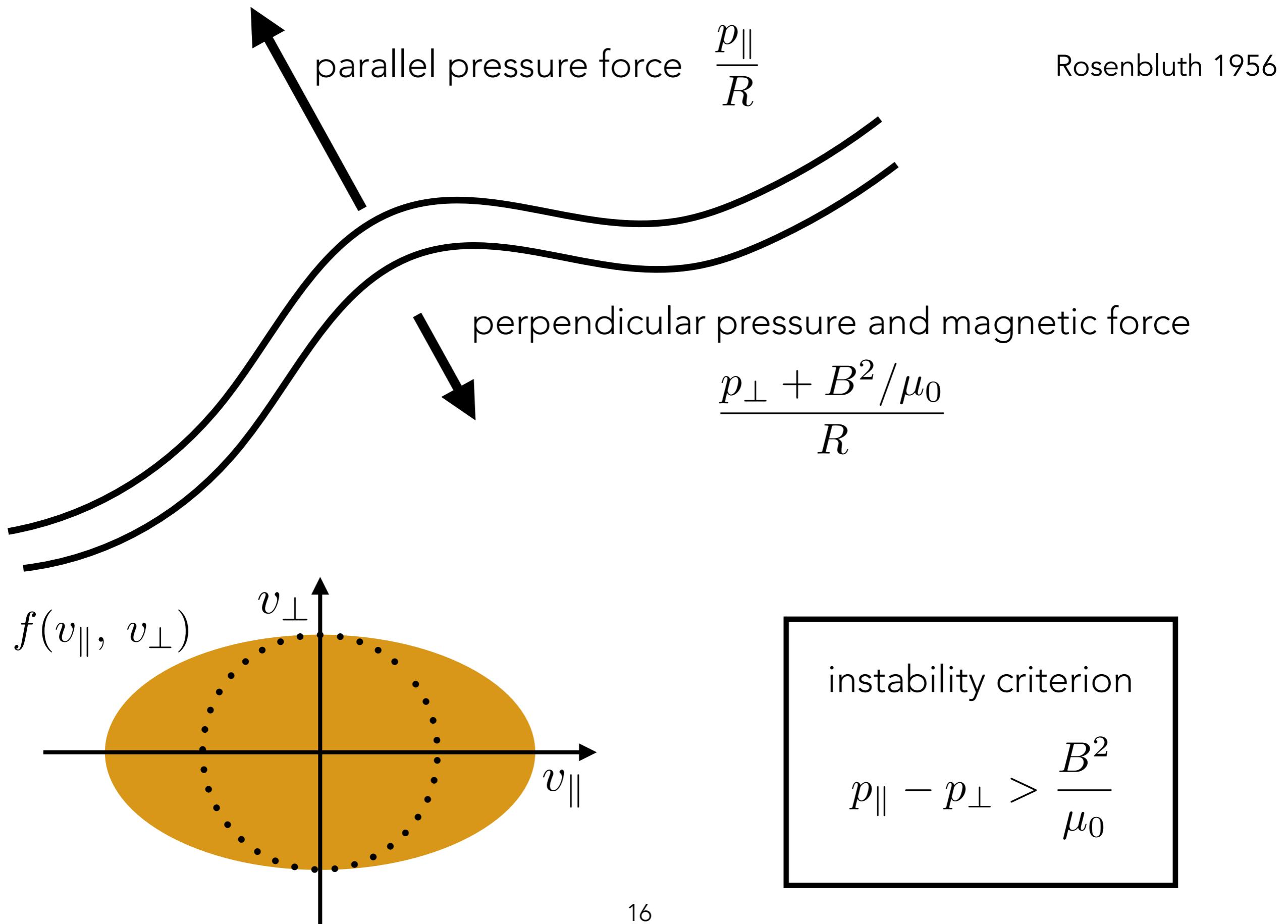
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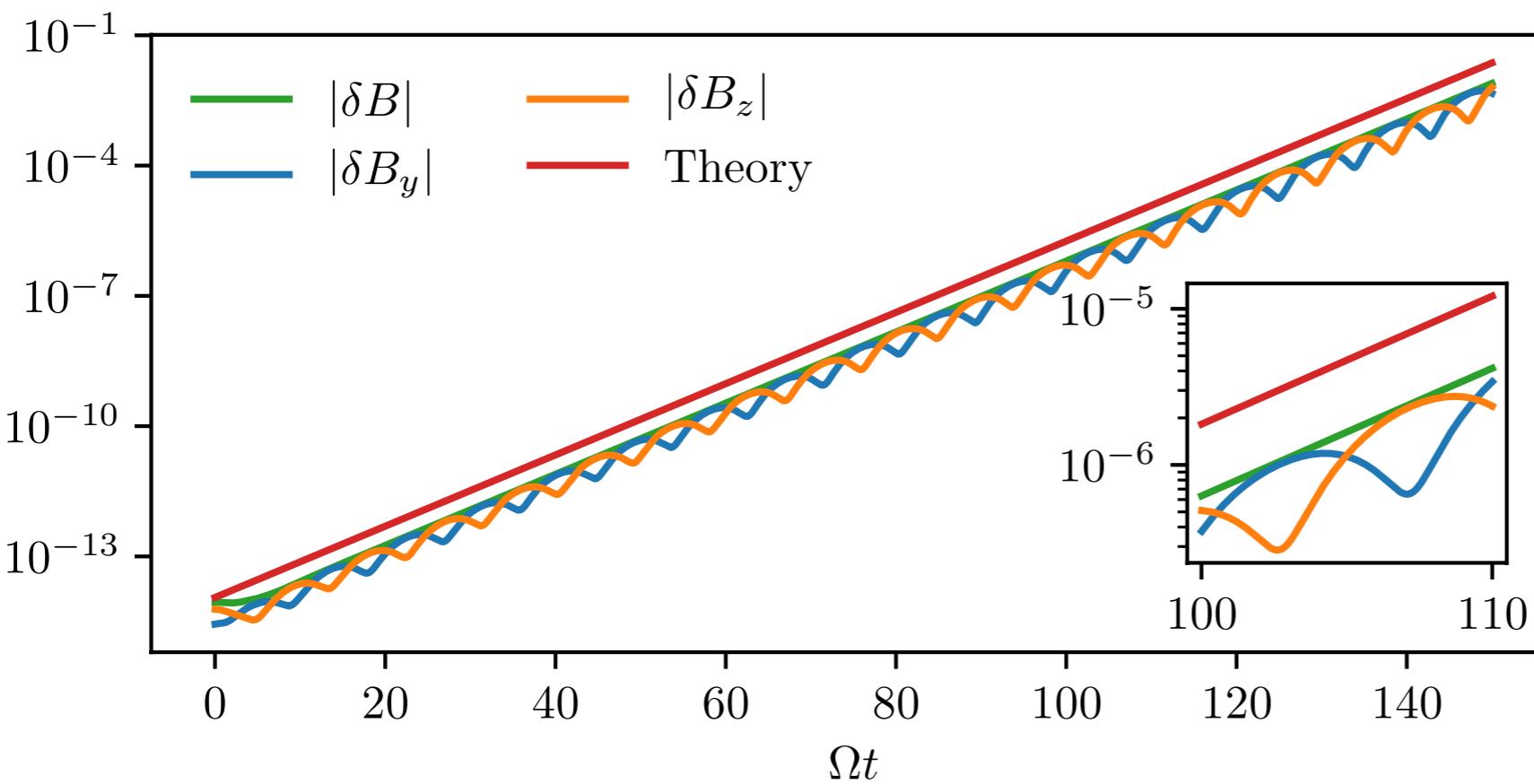
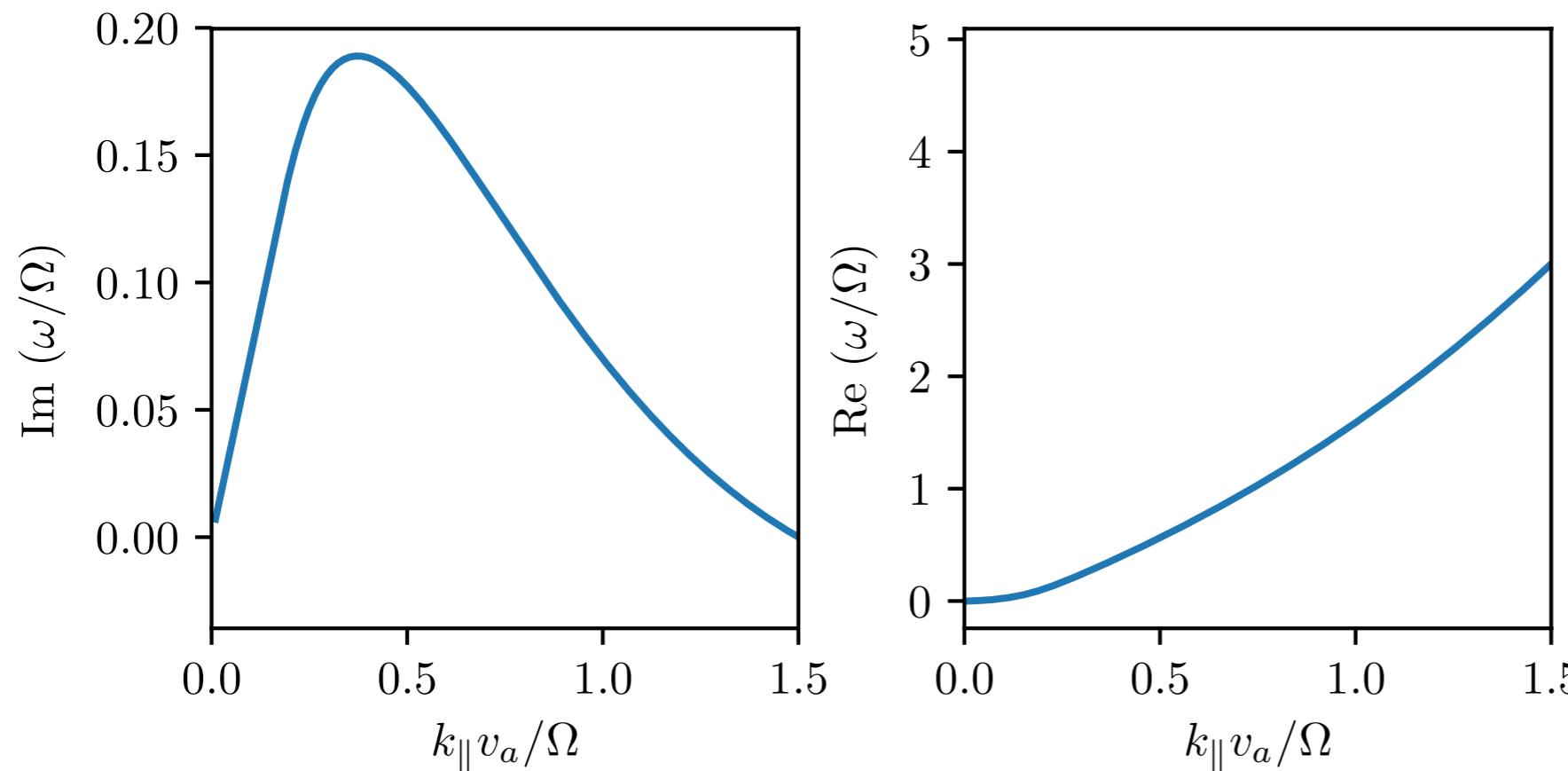


FIREHOSE INSTABILITY



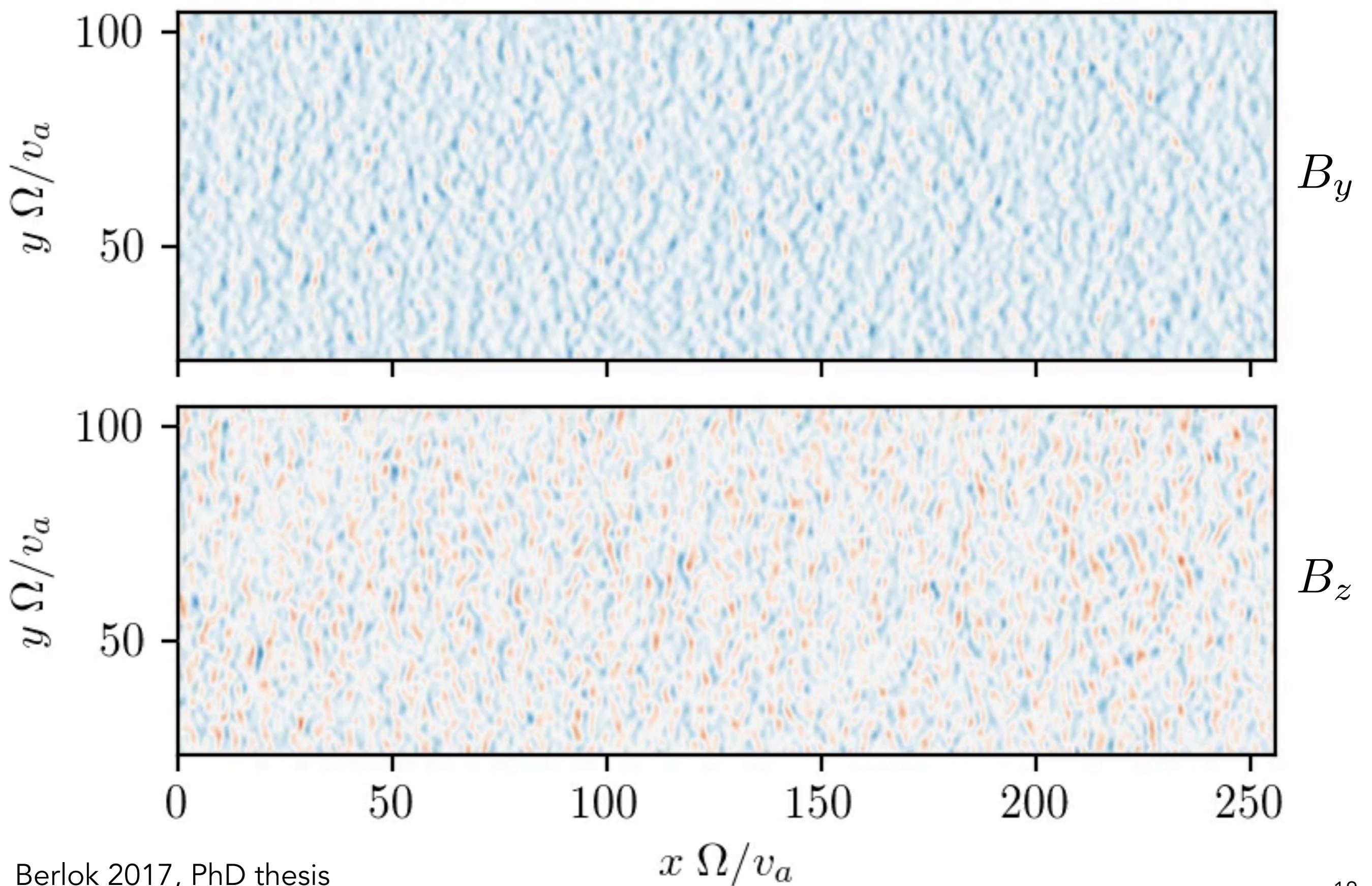
PARALLEL FIREHOSE INSTABILITY

$$\beta_{\parallel} = 4, \beta_{\perp} = 1, T_e = 0$$

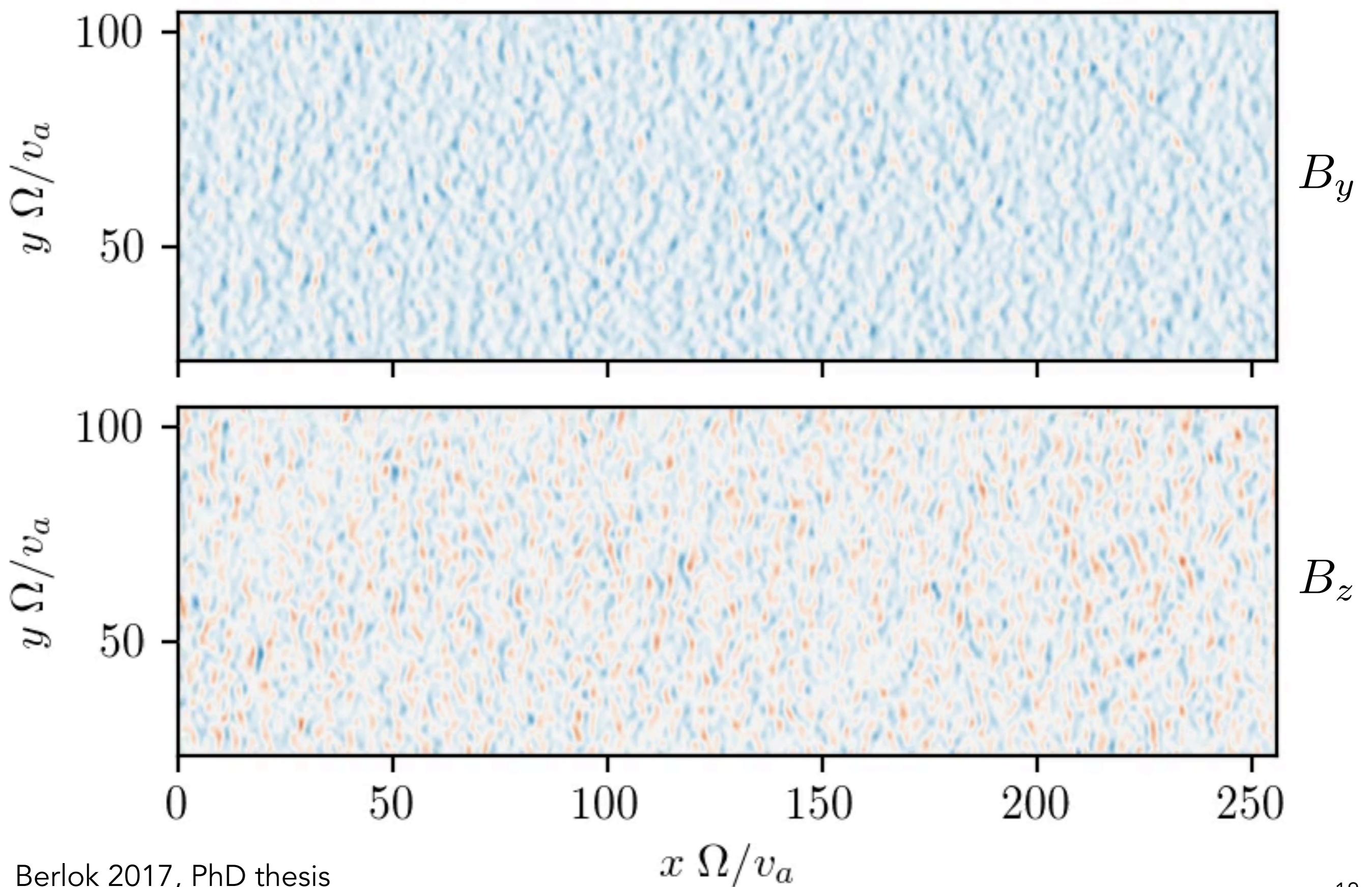


Berlok 2017, PhD thesis
 Advisors:
 Martin Pessah, Troels
 Haugbølle and Tobias
 Heinemann
<http://www.nbi.dk/~berlok/>

2D FIREHOSE INSTABILITY WITH 2D-3V HYBRID-KINETIC CODE

 $\Omega t = 0$ 

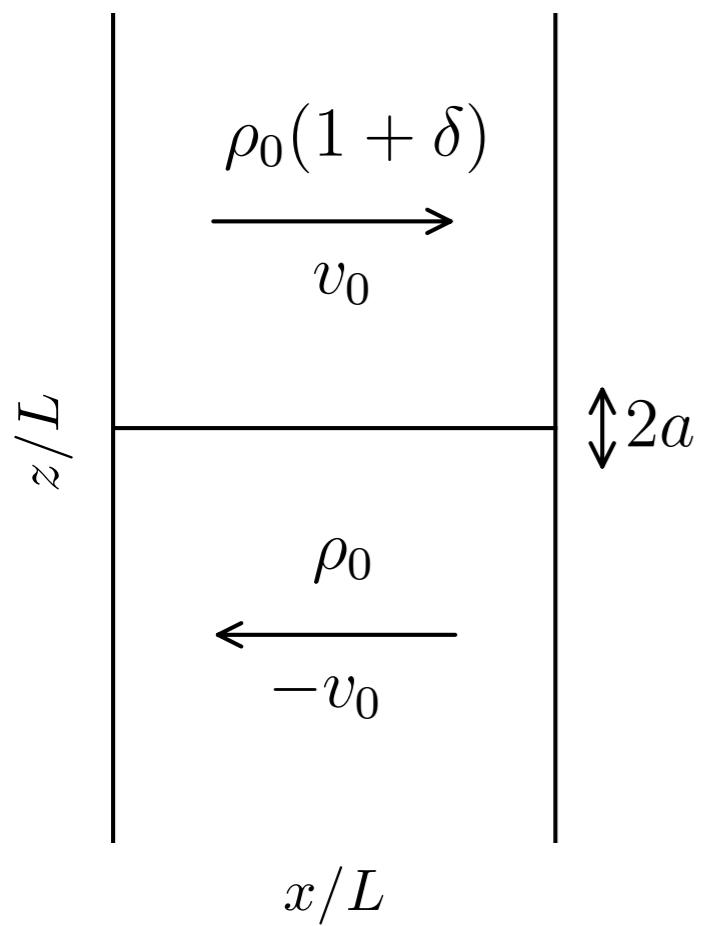
2D FIREHOSE INSTABILITY WITH 2D-3V HYBRID-KINETIC CODE

 $\Omega t = 0$ 

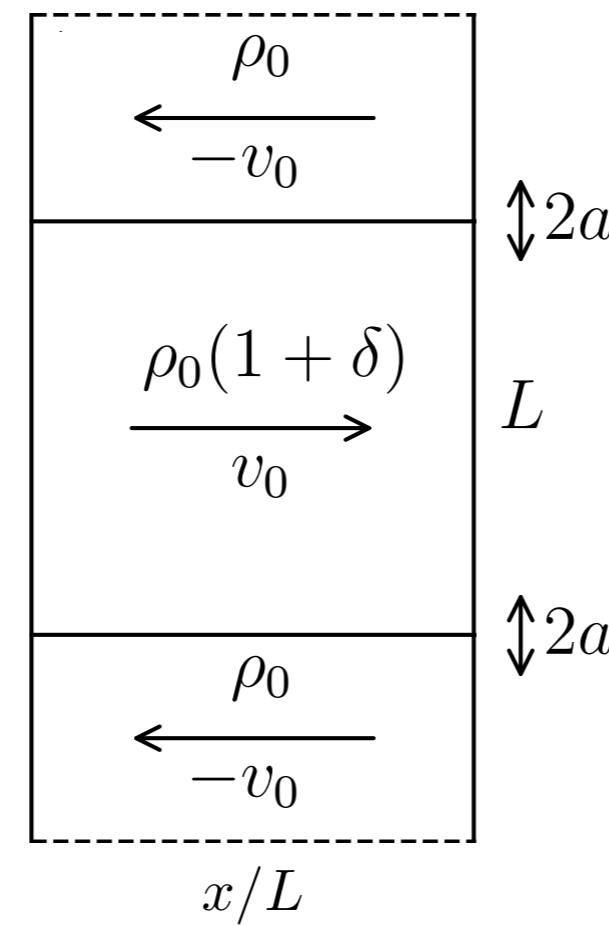
KELVIN-HELMHOLTZ INSTABILITY



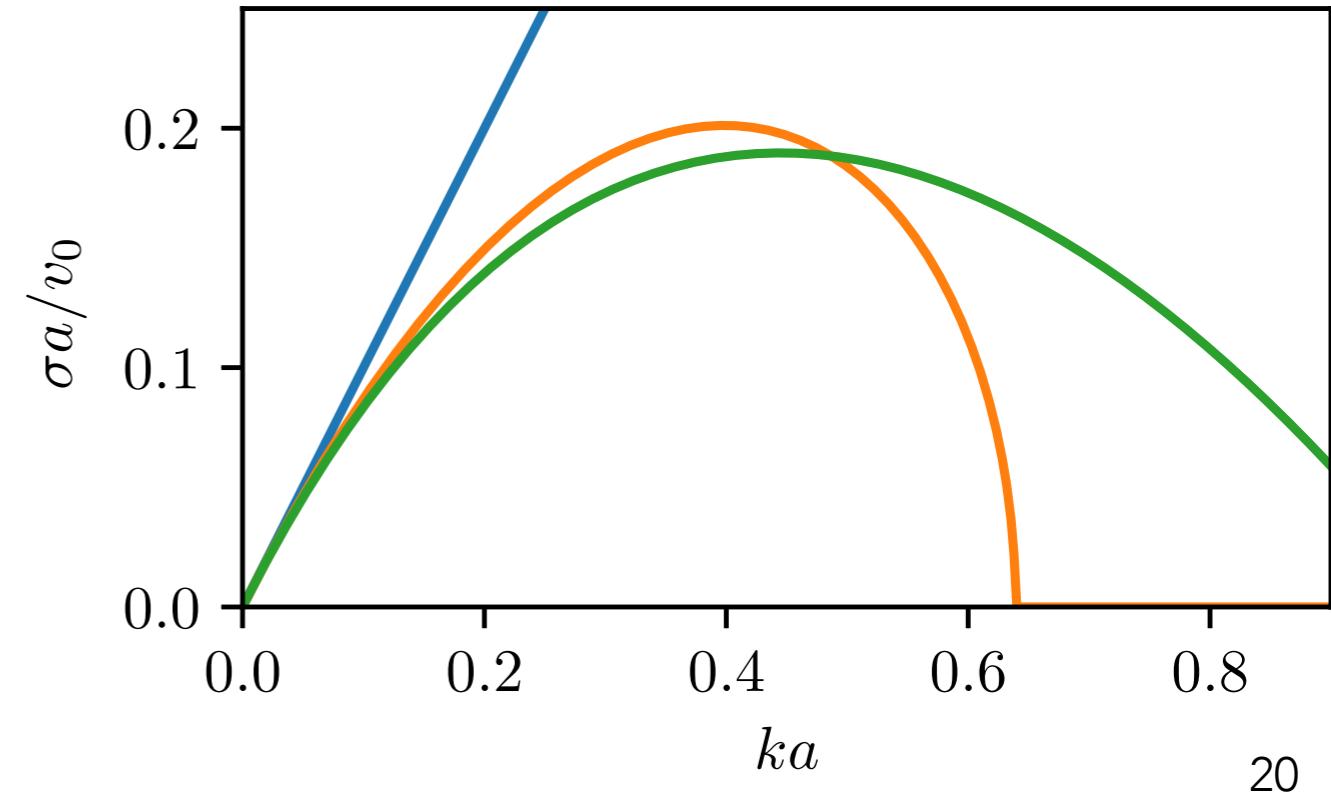
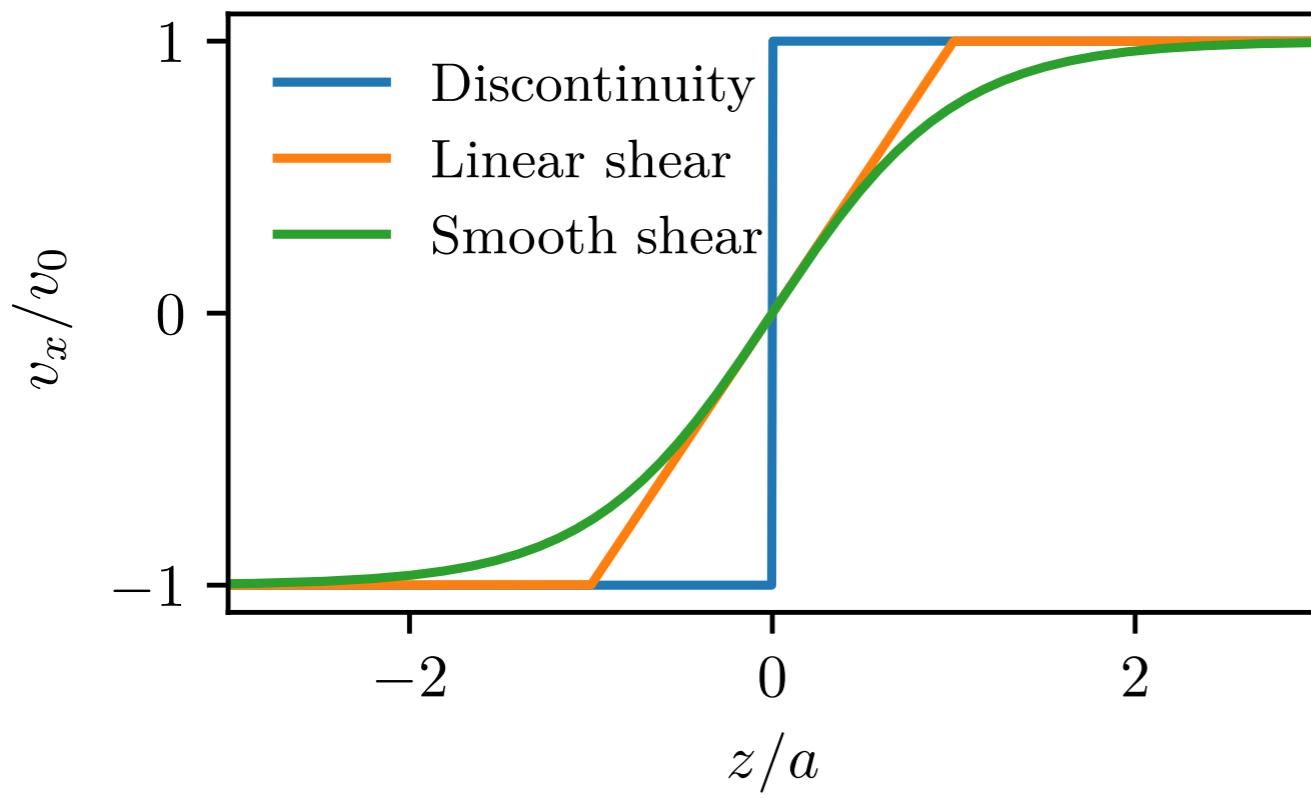
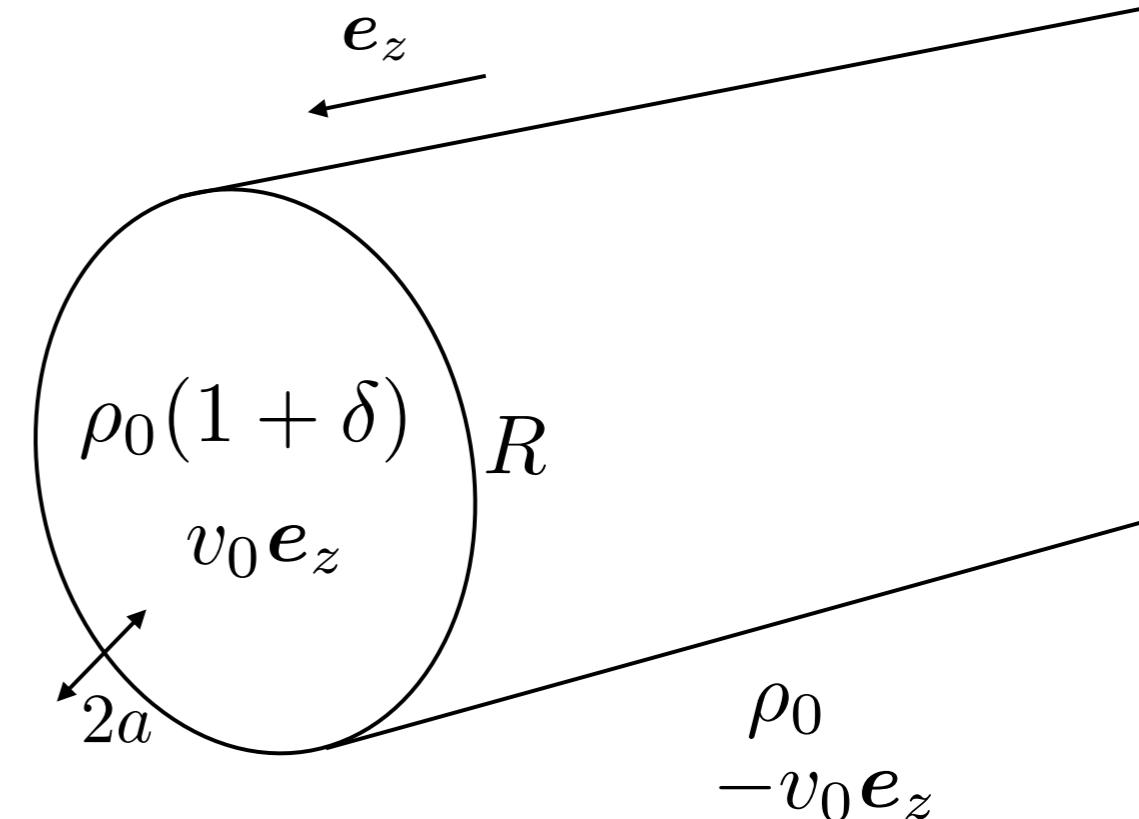
Planar sheet

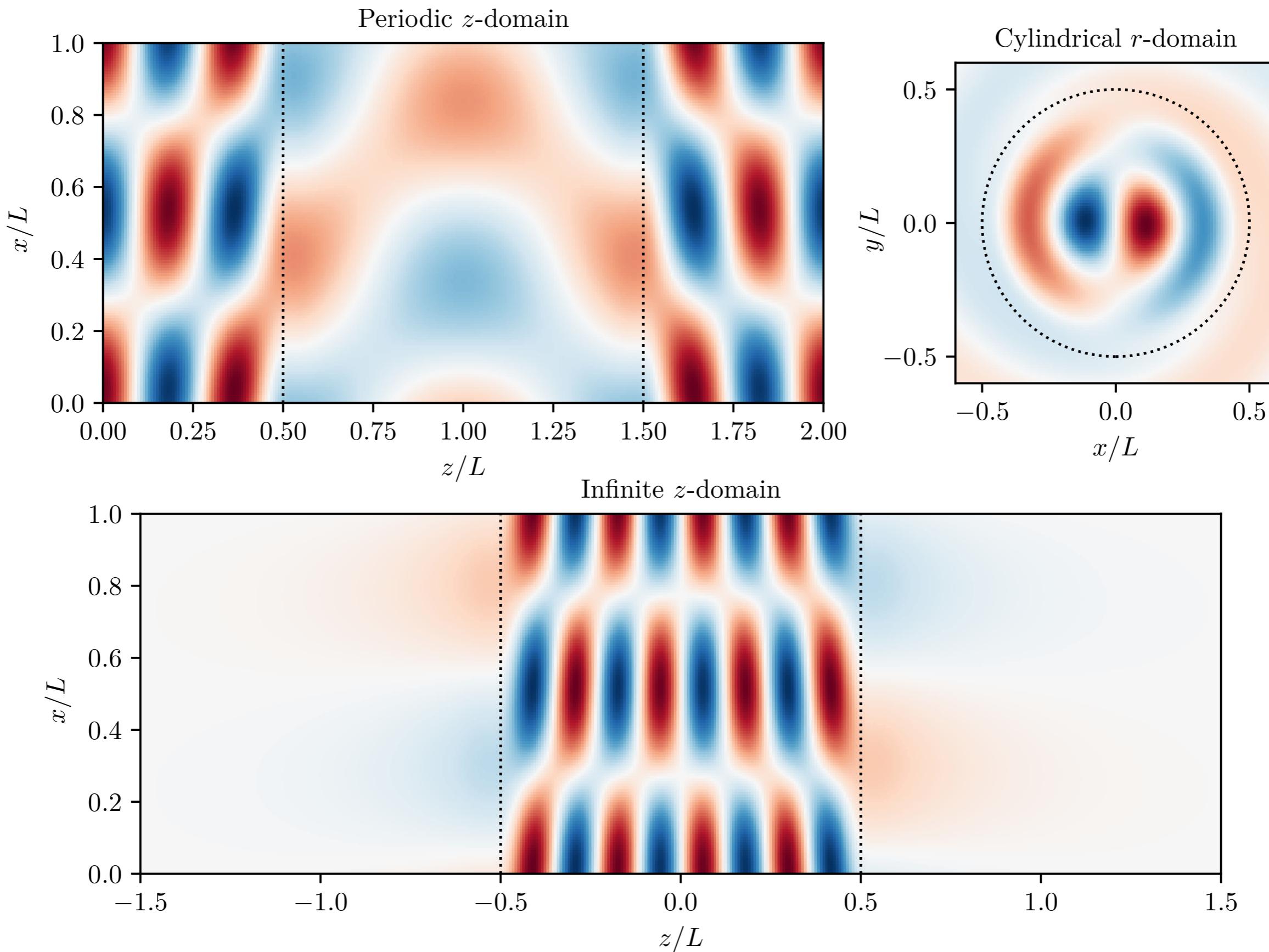


Planar slab(s)

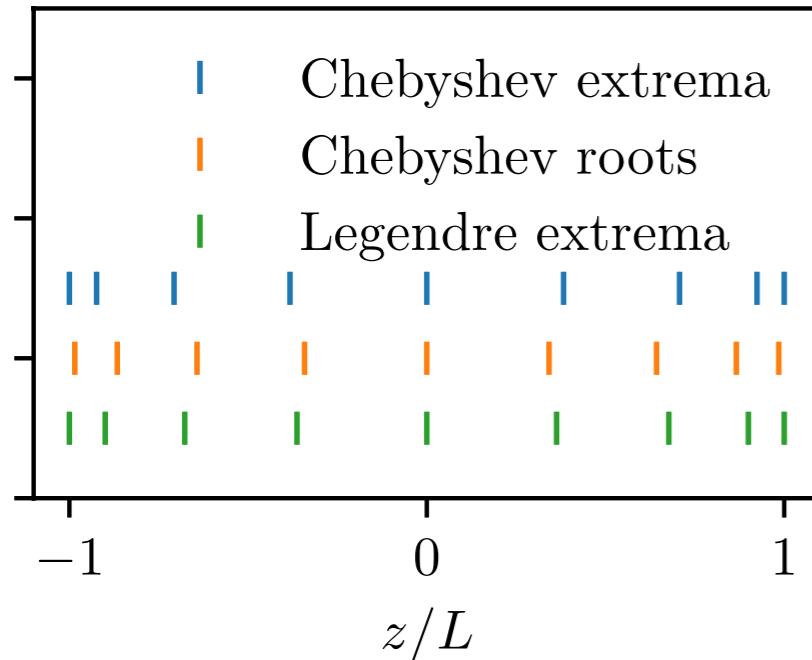


Cylindrical stream

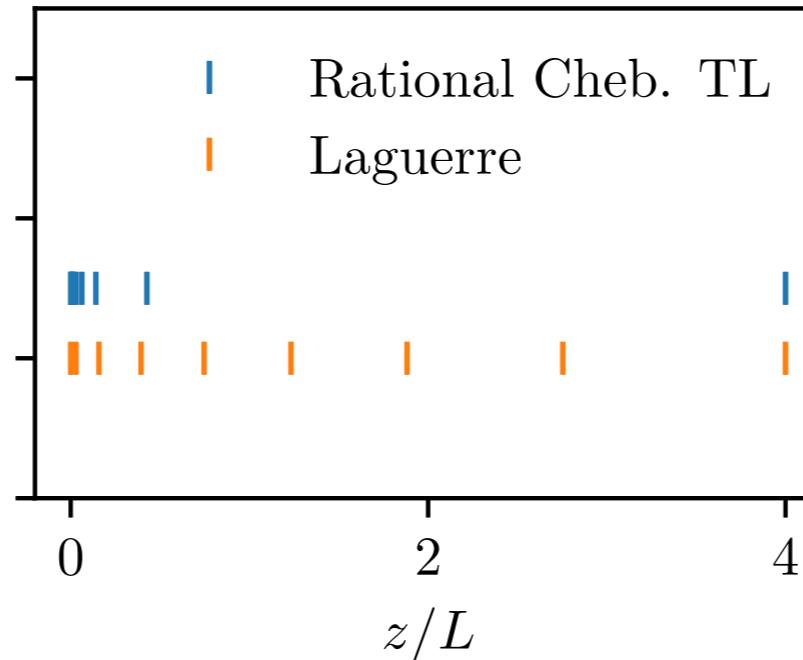




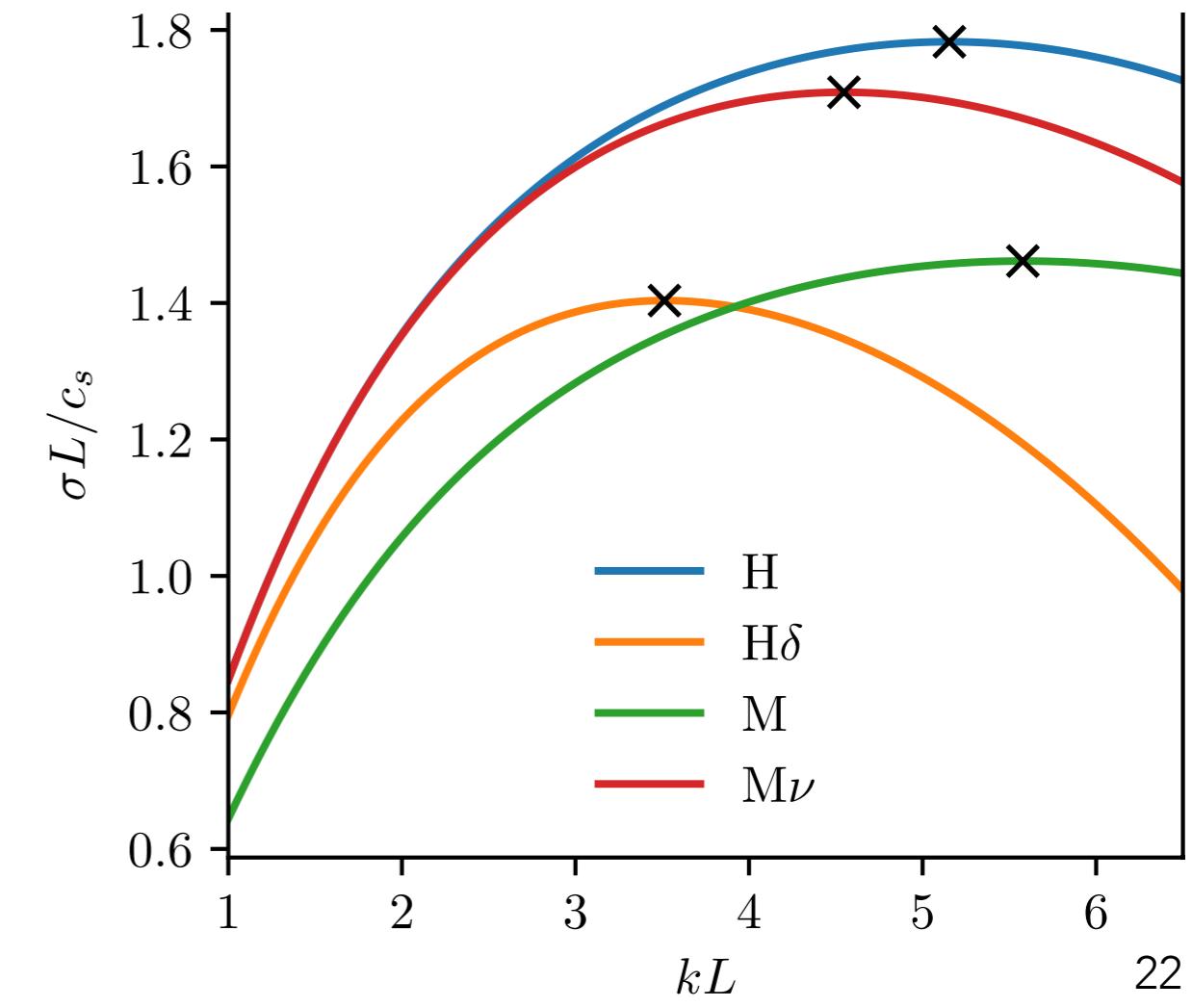
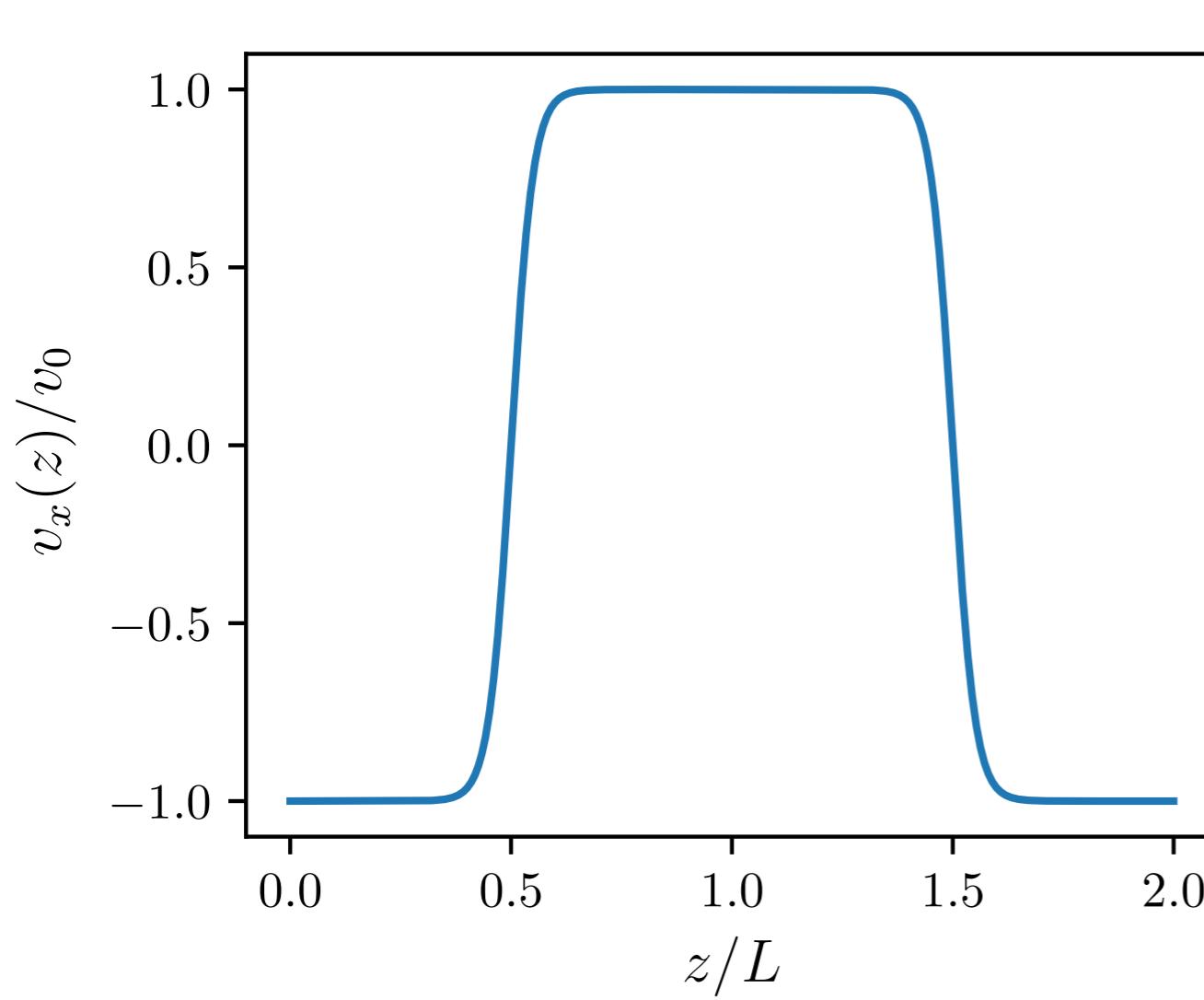
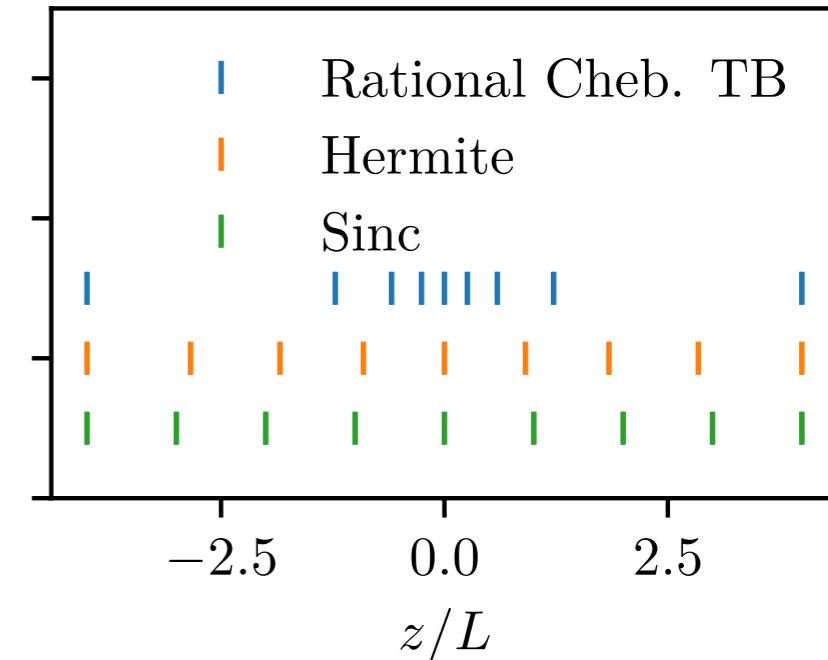
Finite domain



Semi-infinite domain



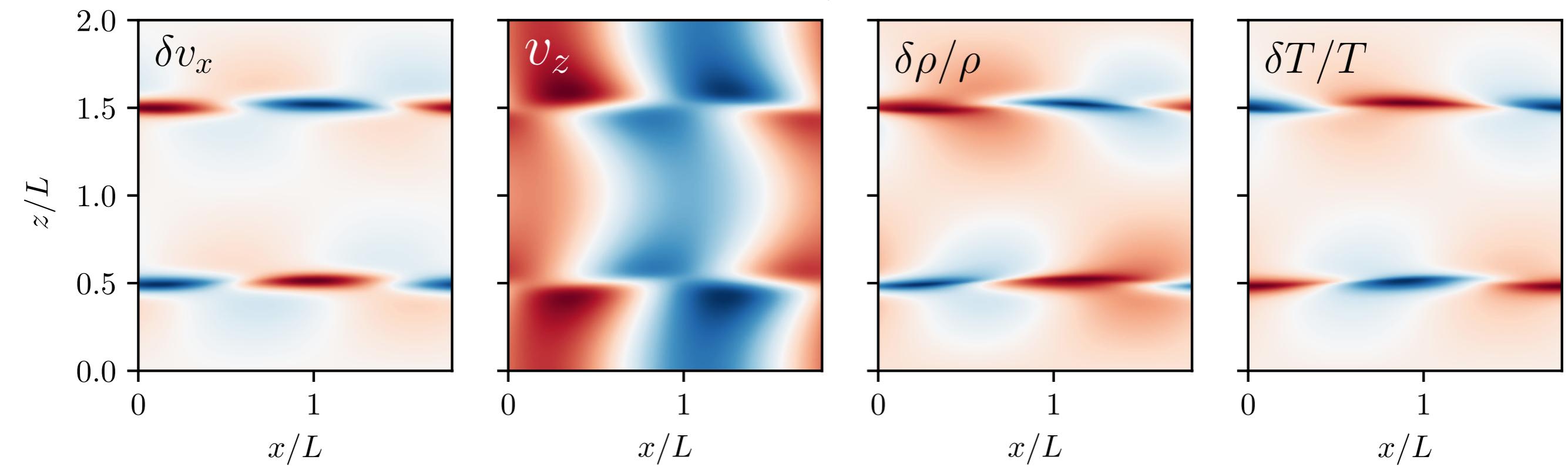
Infinite domain



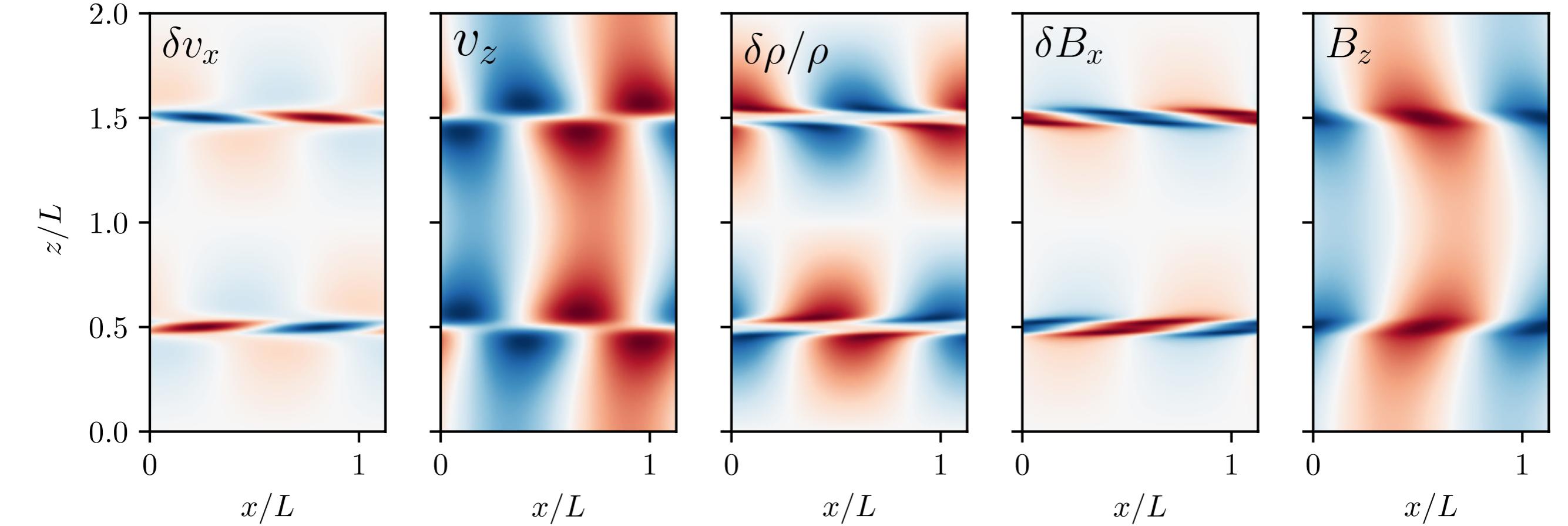
SURFACE MODE SIMULATIONS

with Athena

Hydro, $\rho_s/\rho_0 = 2$

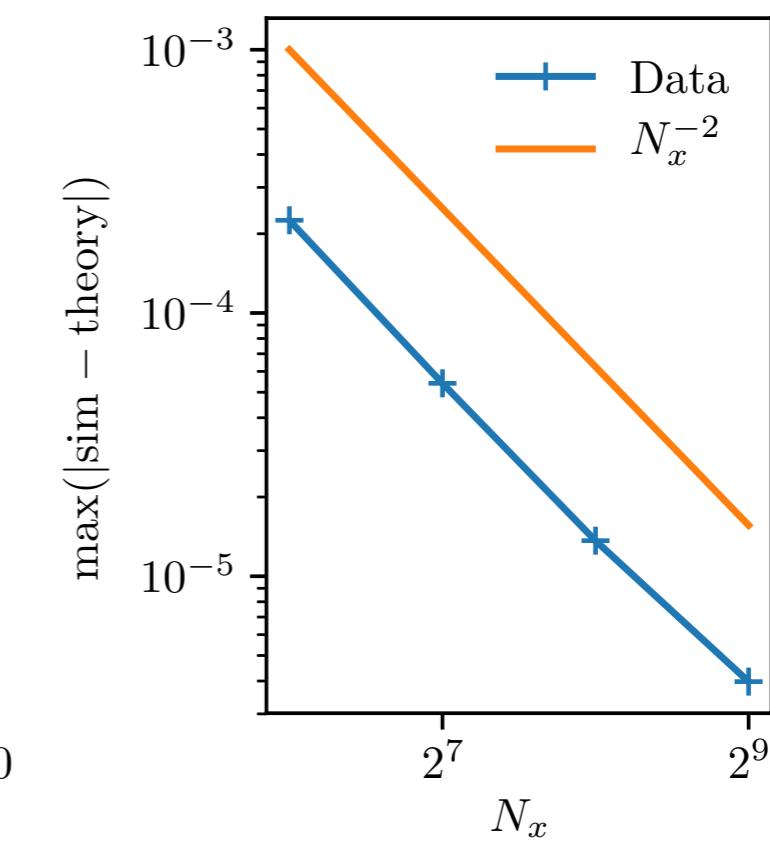
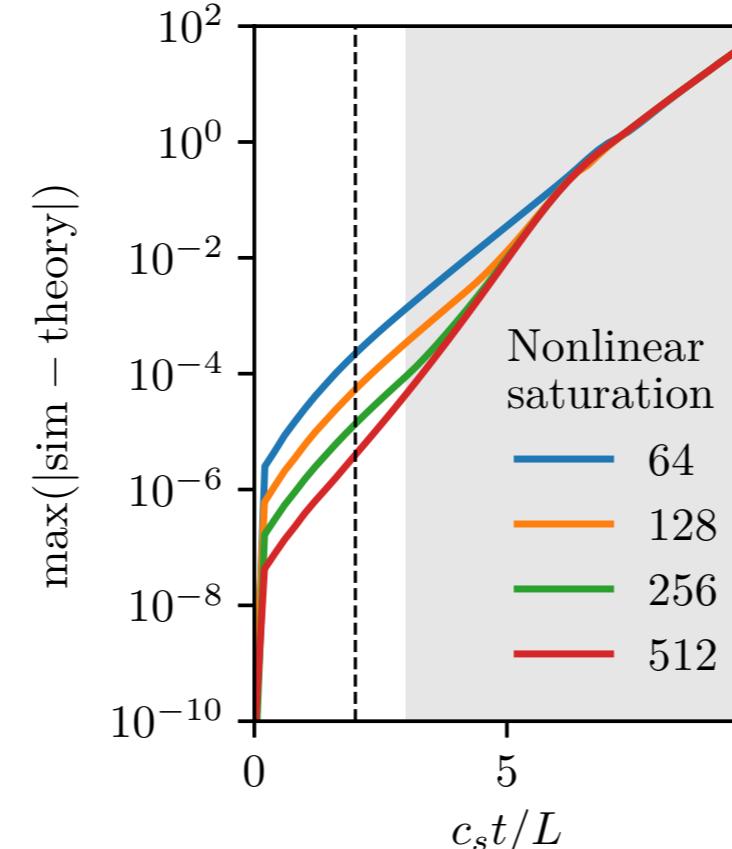
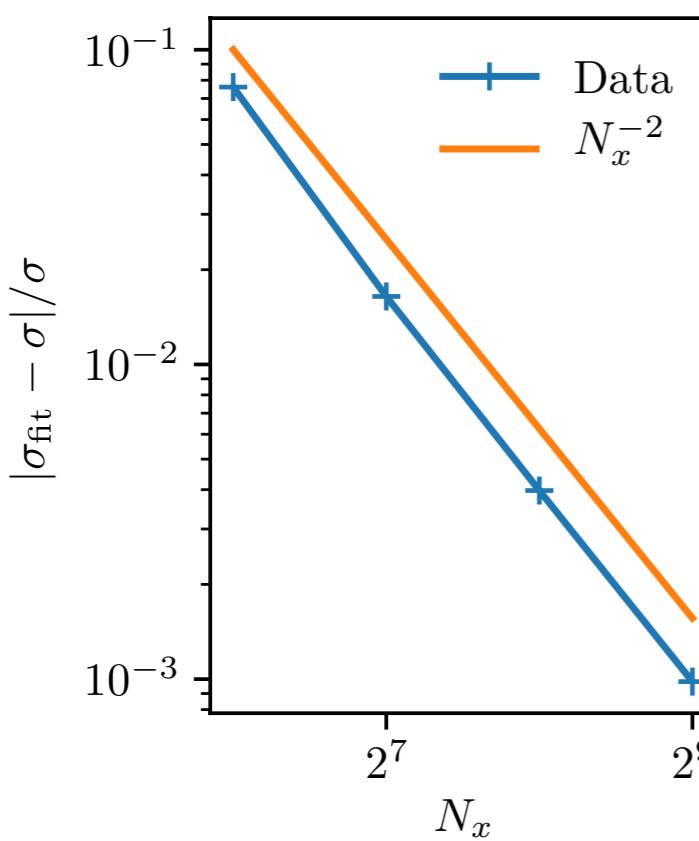
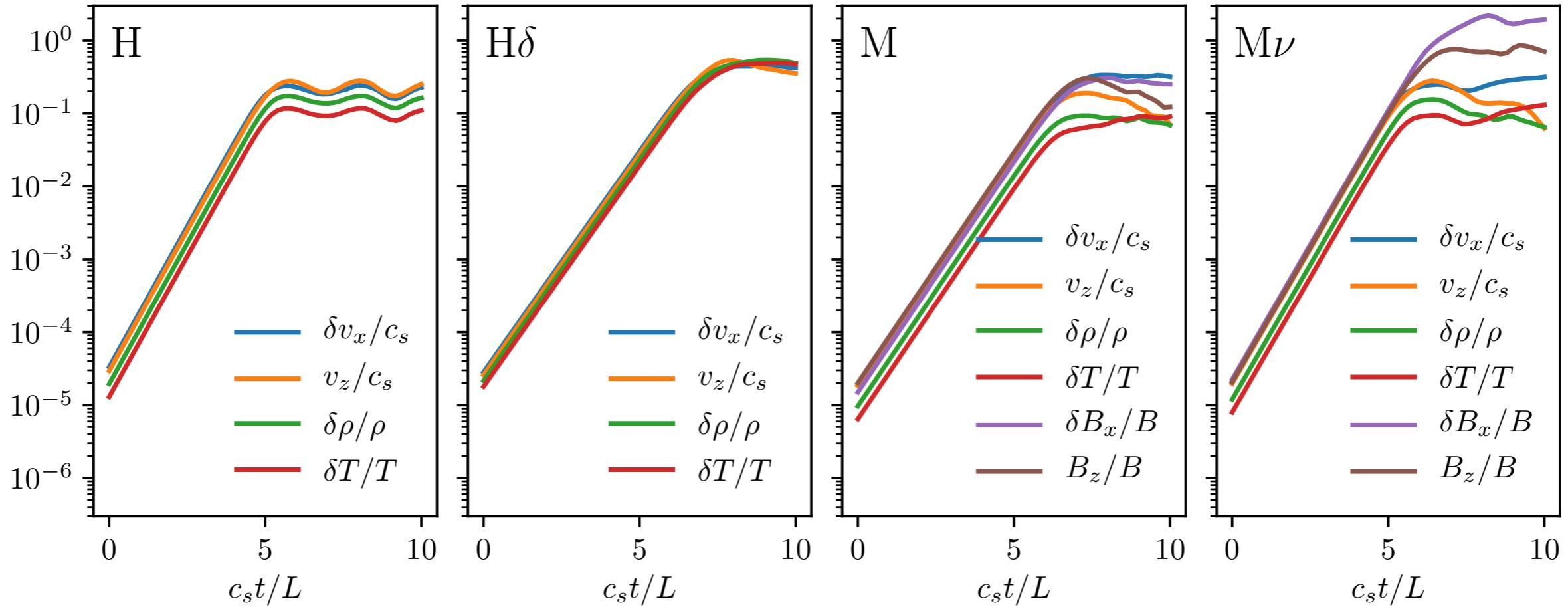


MHD, $\beta = 5$



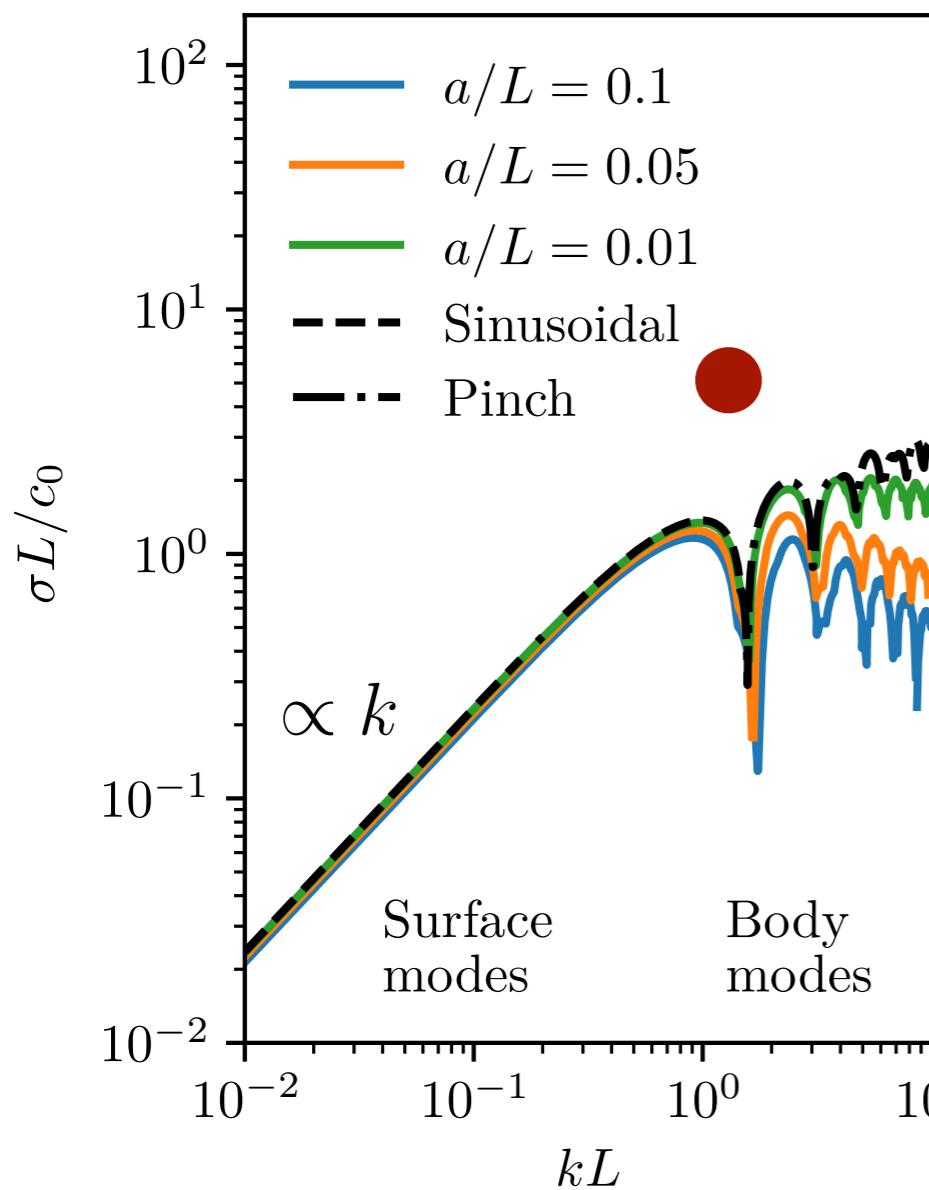
KELVIN-HELMHOLTZ INSTABILITY

with Athena

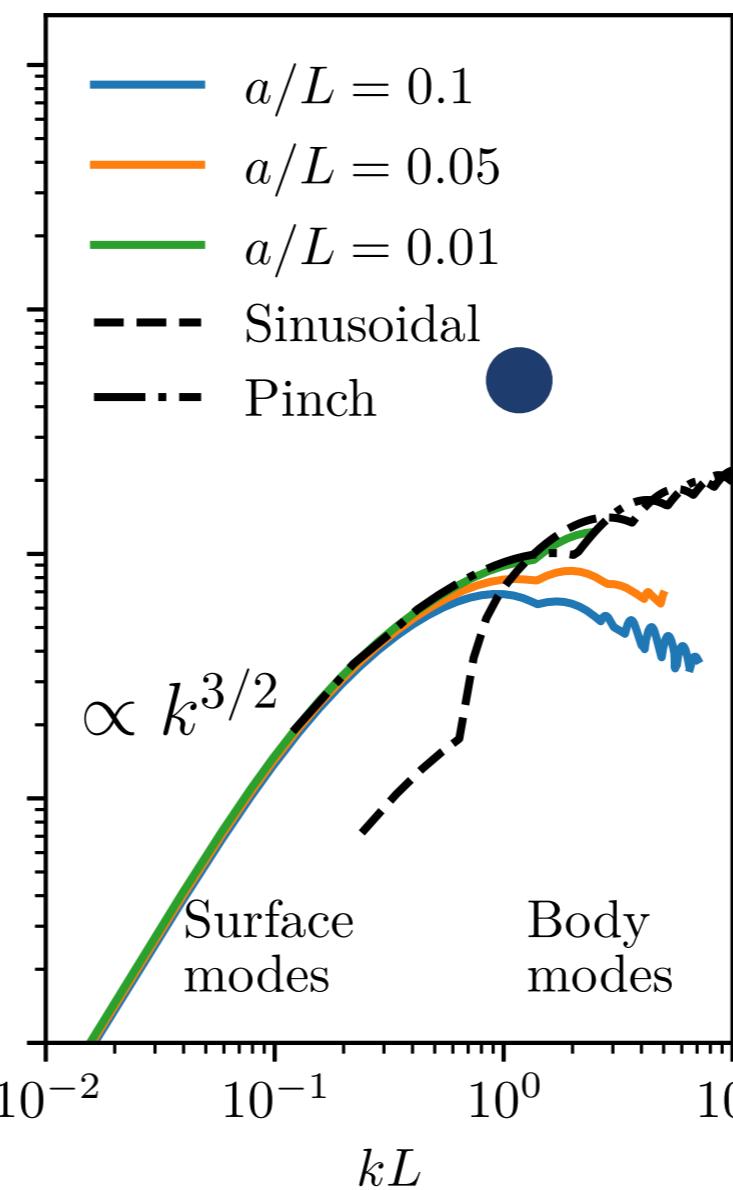


SUPersonic BODYMODE DISPERSION RELATION

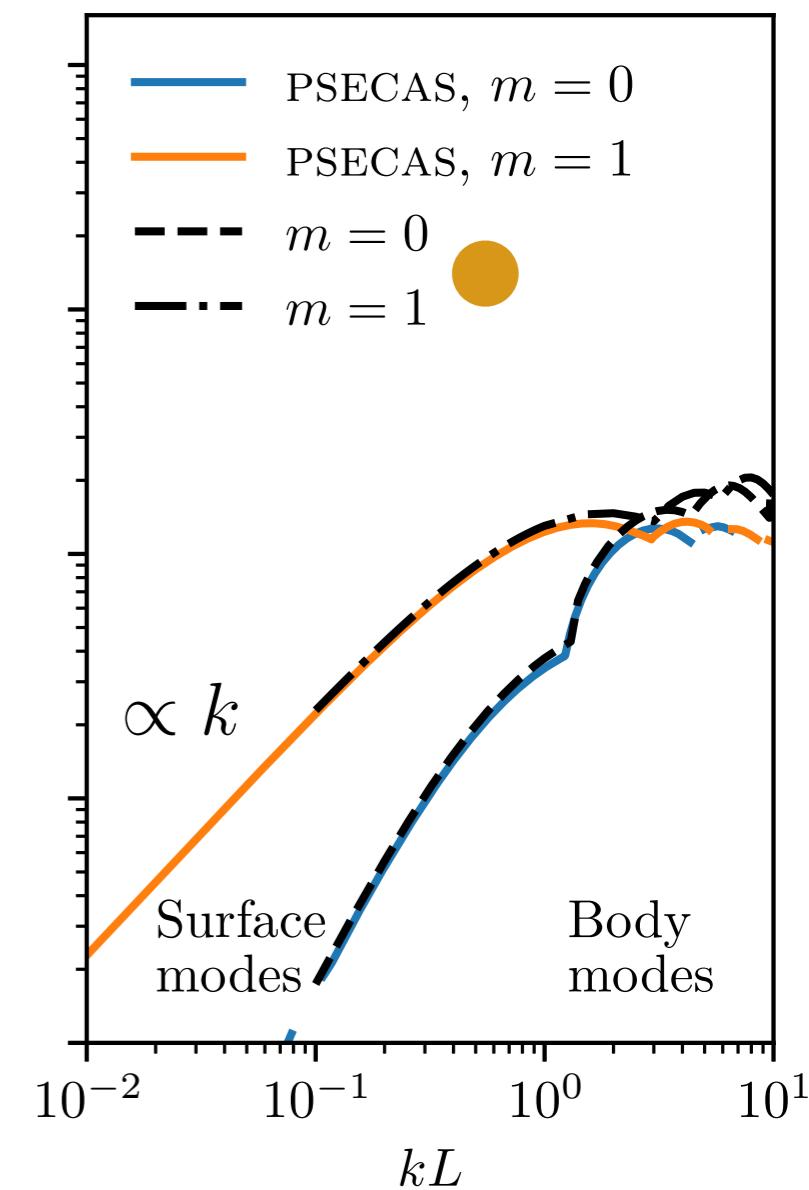
Periodic z -domain



Infinite z -domain



Cylindrical domain



Berlok & Pfrommer 2019a

$$\frac{\mathcal{T}(q_2 L/2)}{\mathcal{T}(q_1 L/2)}$$

$$\left(\frac{\omega - kv_2}{\omega - kv_1} \right)^2 \frac{\rho_2}{\rho_1} \frac{q_1}{q_2} =$$

● $-\mathcal{T}(q_2 L/2)$

Mandelker+ (2016)

$$\frac{\mathcal{I}'_m(q_1 R)}{\mathcal{I}_m(q_1 R)} \frac{\mathcal{K}_m(q_2 R)}{\mathcal{K}'_m(q_2 R)}$$

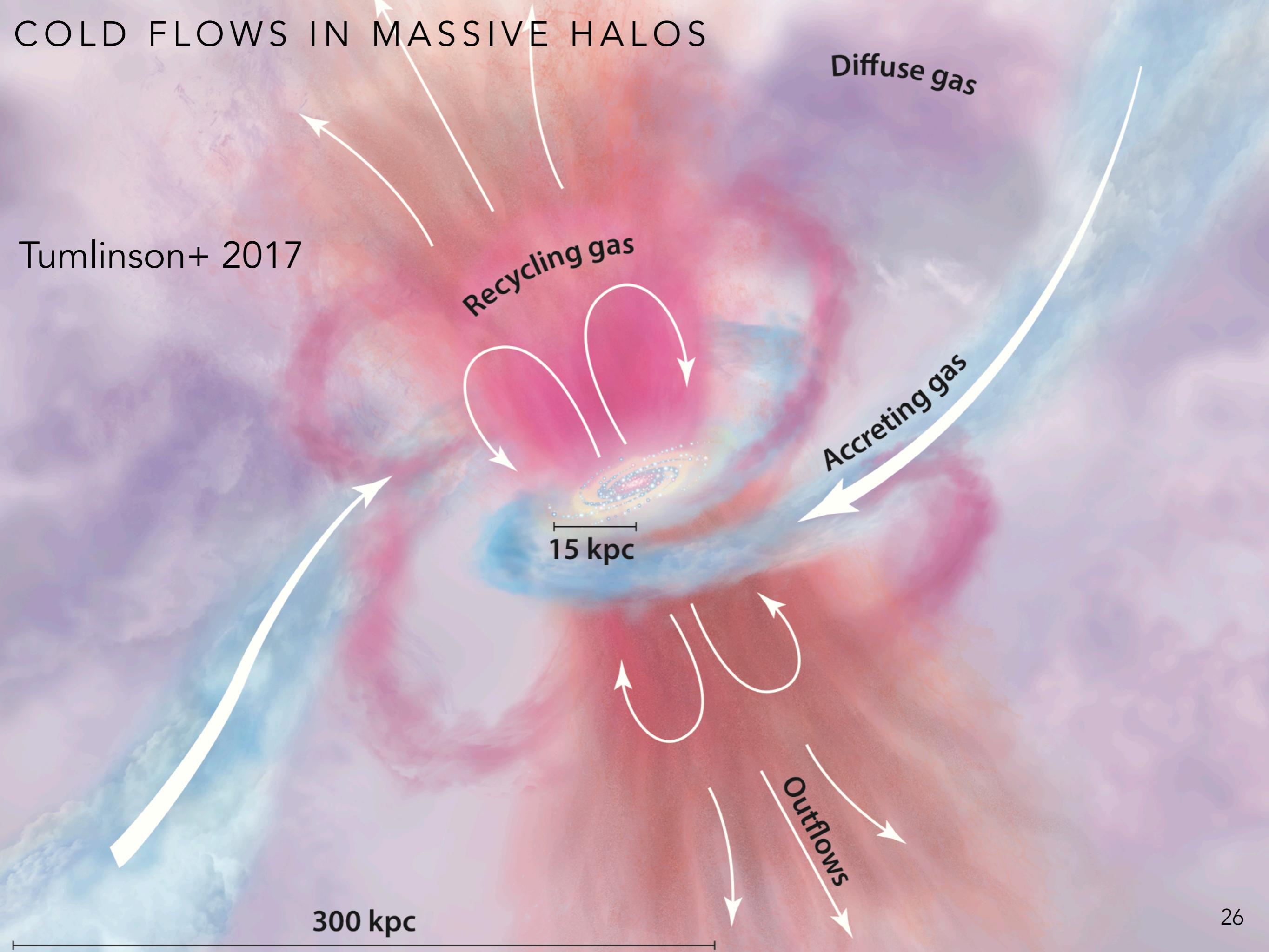
$$\mathcal{T}(z) = \tanh(z) \text{ or } \cosh(z)$$

$$q = k \sqrt{1 - \left(\frac{\omega - kv}{kc} \right)^2}$$

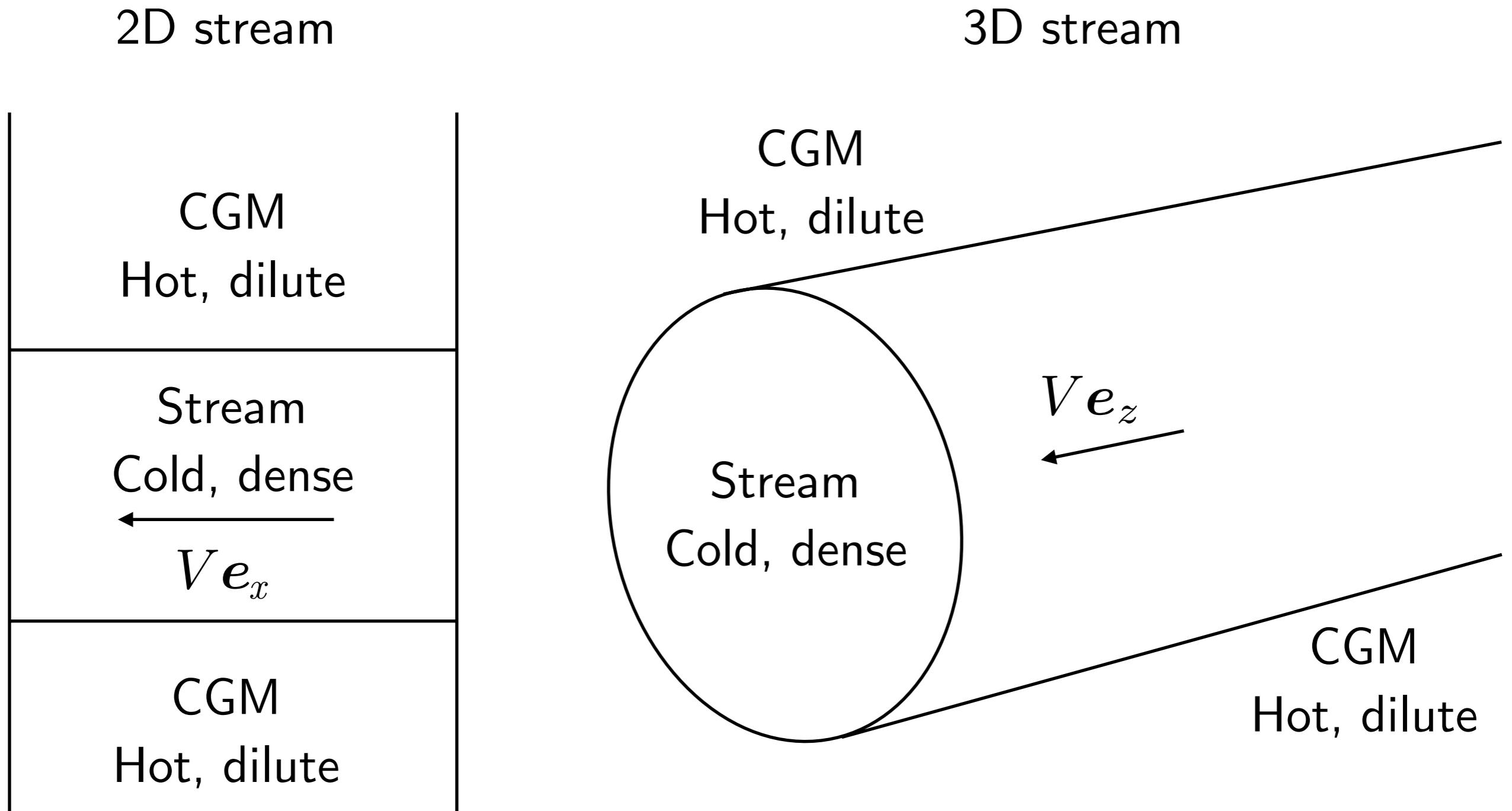
$\mathcal{I}_m, \mathcal{K}_m$ are modified Bessel functions

COLD FLOWS IN MASSIVE HALOS

Tumlinson+ 2017



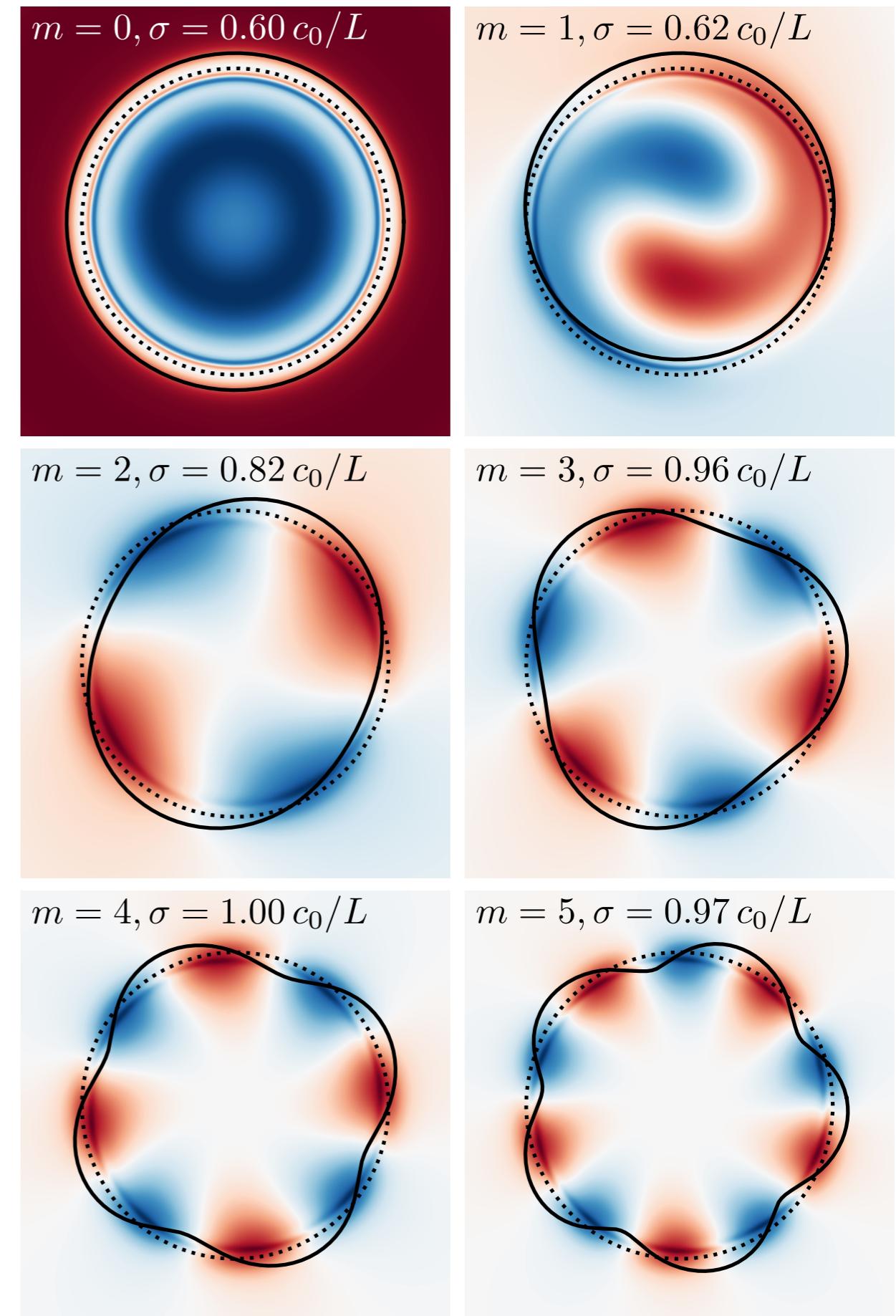
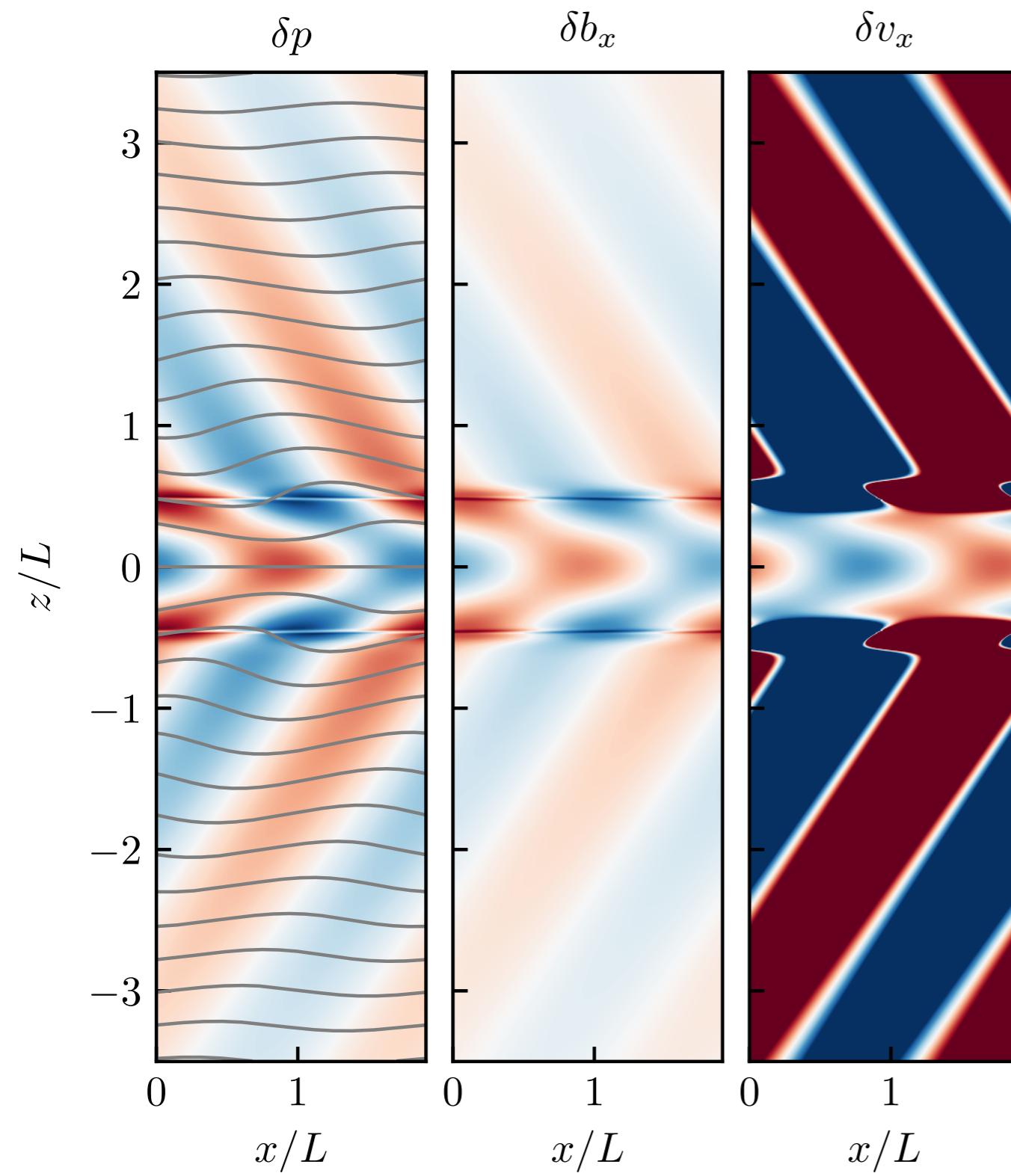
COLD FLOWS IN MASSIVE HALOS



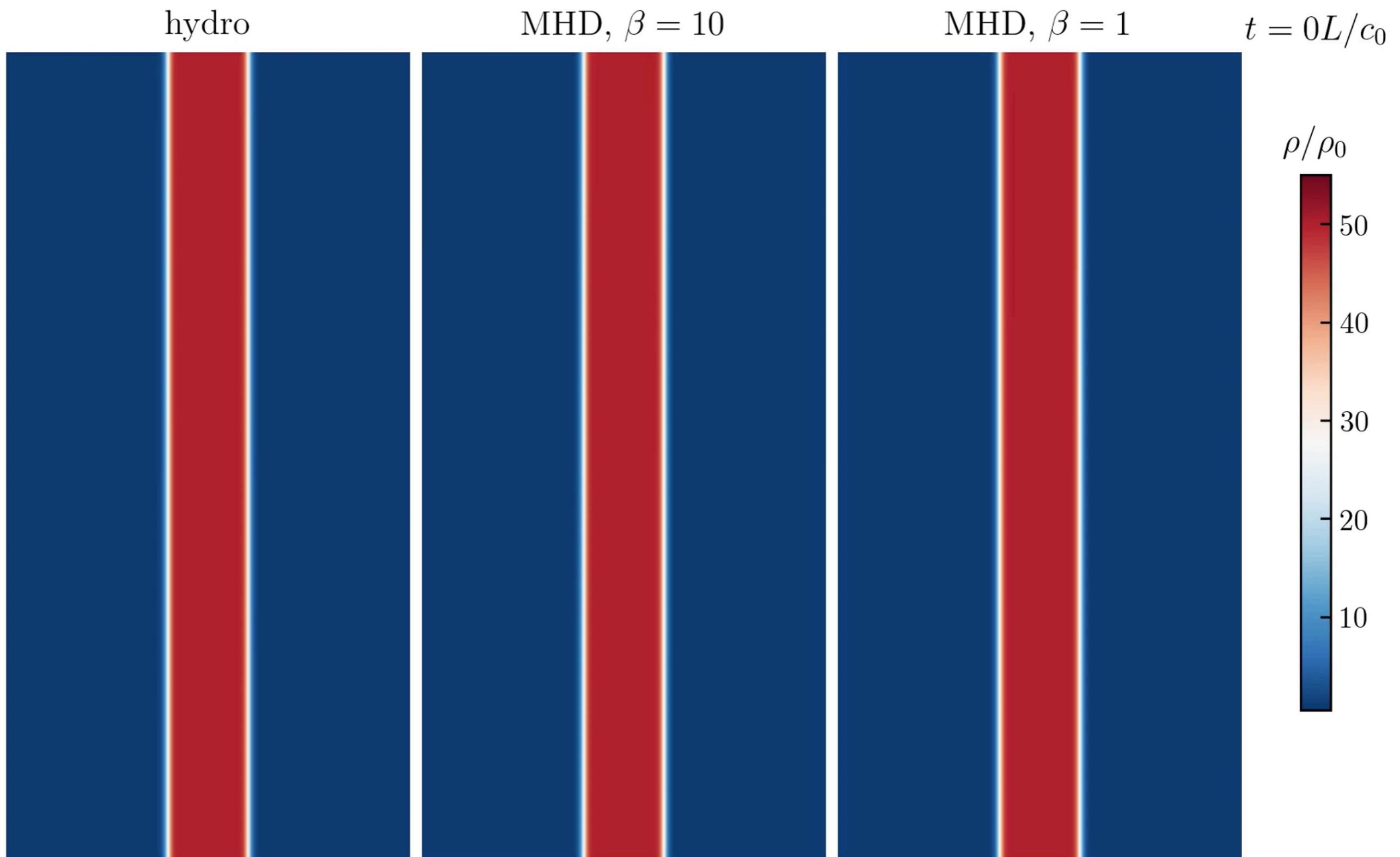
Hydrodynamic studies in Mandelker+ 2016, 2019 and Padnos+ 2019

MHD study in Berlok & Pfrommer 2019b

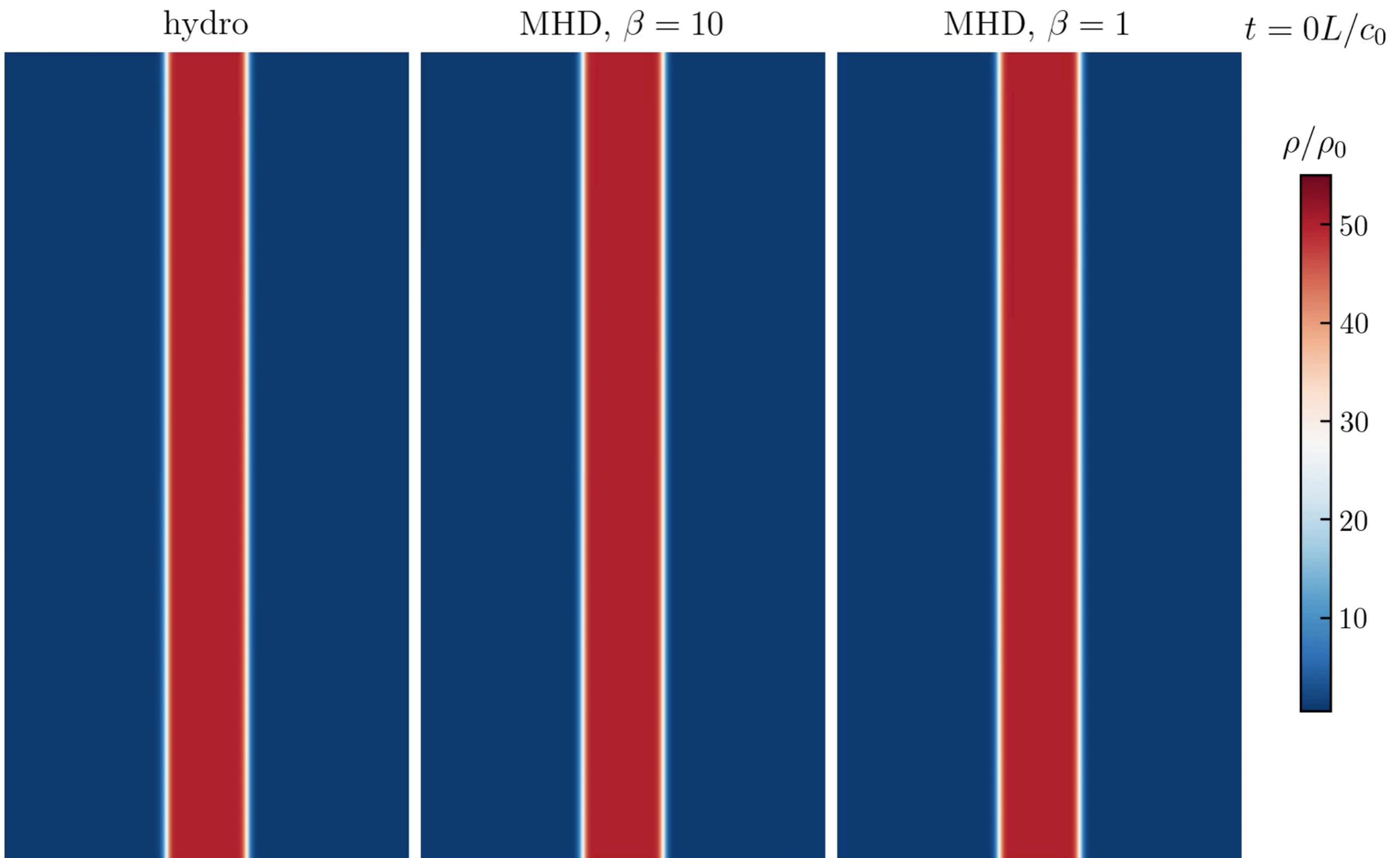
LINEAR THEORY, IDEAL MHD



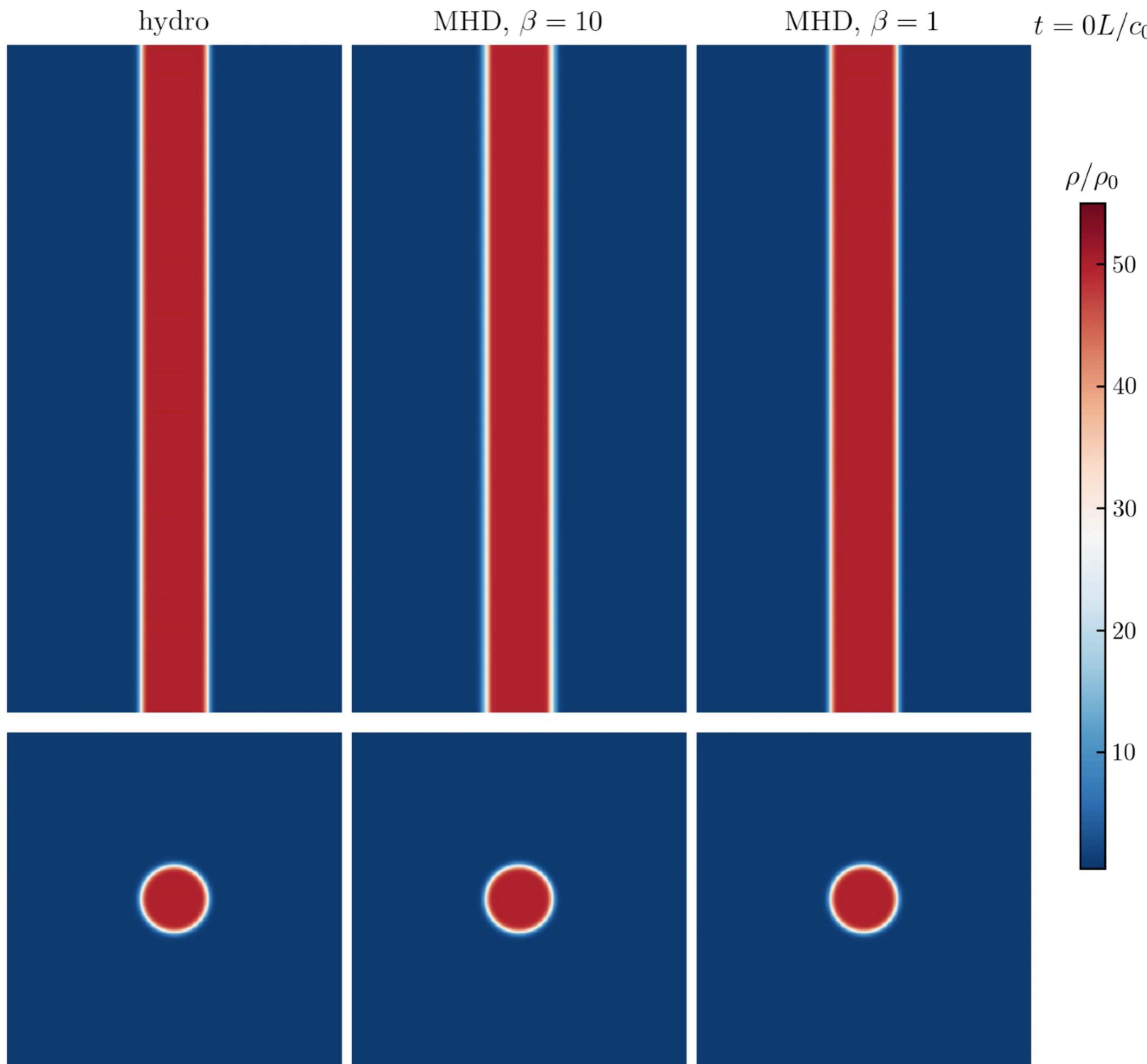
2D MAGNETIZED COLD STREAMS



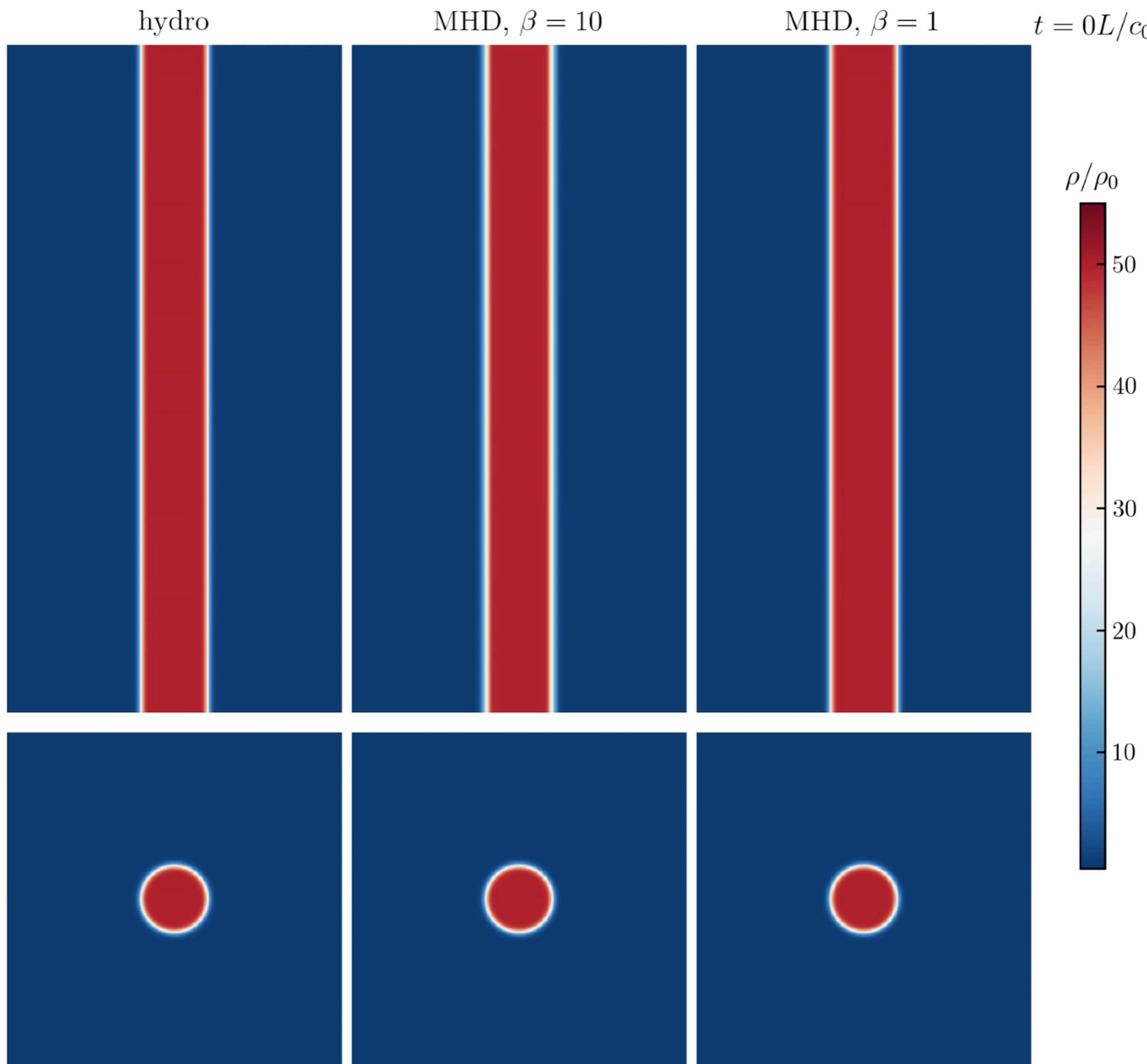
2D MAGNETIZED COLD STREAMS



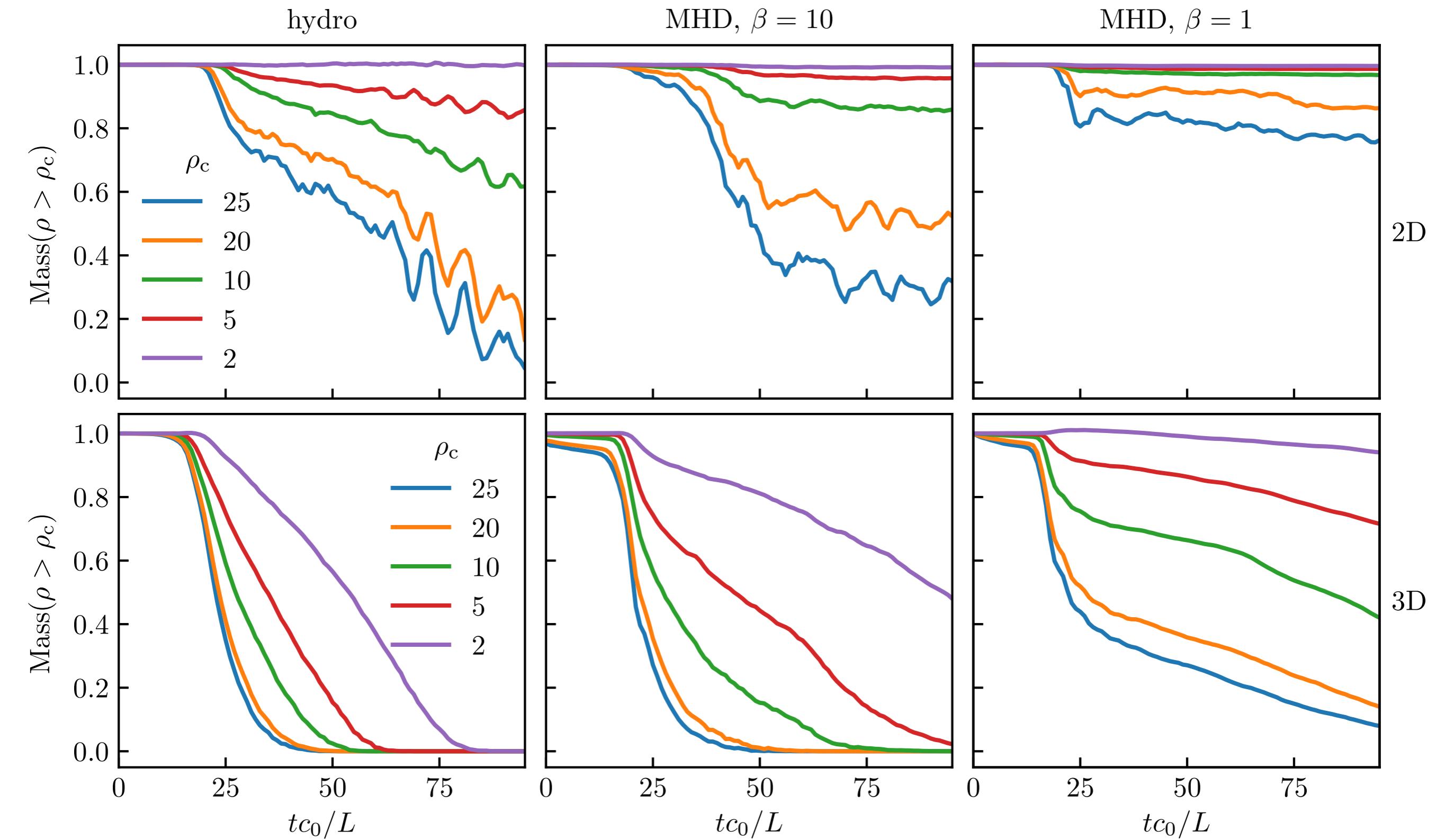
3 D



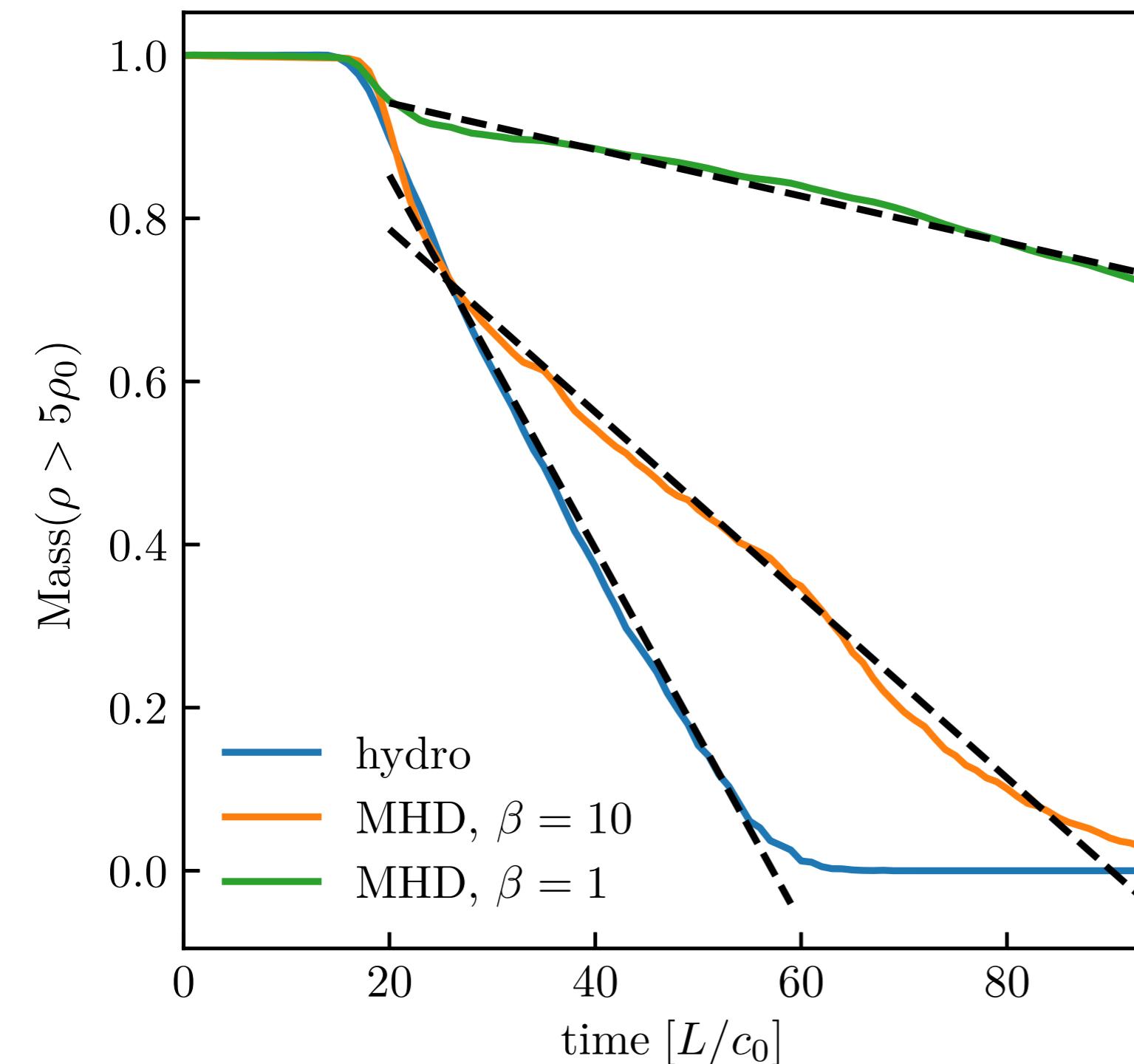
3 D



MIXING OF COLD STREAMS WITH CGM



MAGNETIC FIELDS SUPPRESS MIXING



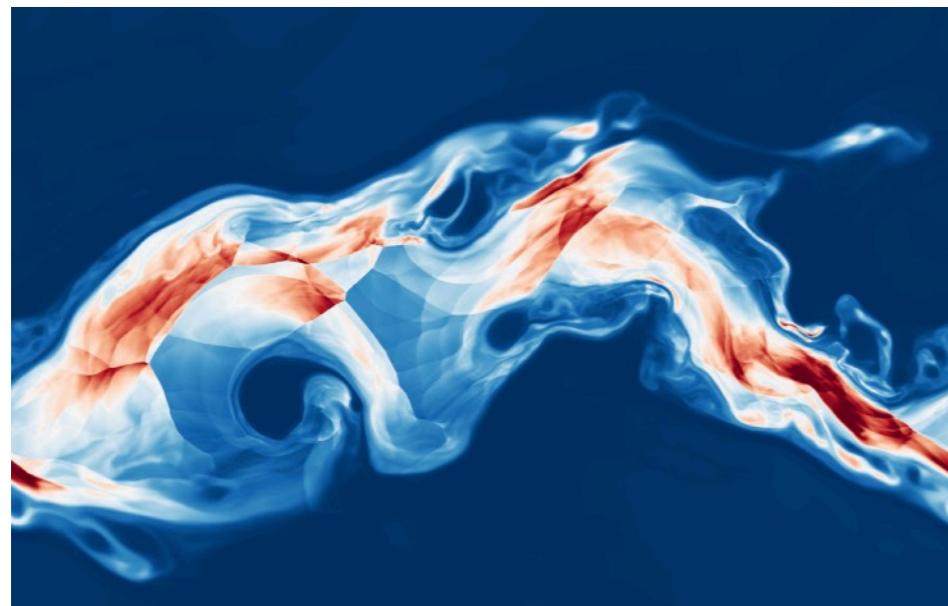
$$\frac{L}{R_v} \lesssim \left\{ \begin{array}{ll} 2 \times 10^{-2} & \text{hydro} \\ 9 \times 10^{-3} & \beta = 10 \\ 2 \times 10^{-4} & \beta = 1 \end{array} \right.$$

$$B = \sqrt{\frac{2p\mu_0}{\beta}} = 0.8 \left(\frac{n_{\text{H}}}{10^{-4} \text{cm}^{-3}} \right)^{1/2} \left(\frac{T}{10^6 \text{K}} \right)^{1/2} \beta^{-1/2} \mu\text{G}$$

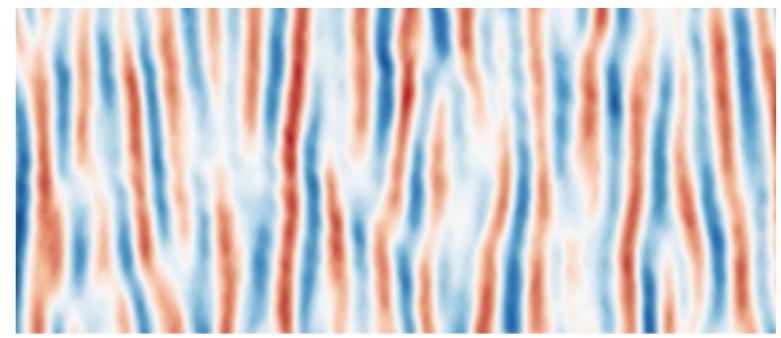
SUMMARY

- Weakly collisional and collisionless plasmas: Small scales with hybrid-kinetic codes, intermediate scales with Athena
- Large scales with Braginskii viscosity in Arepo
- Supersonic, magnetized Kelvin-Helmholtz instability in cold streams at high redshift

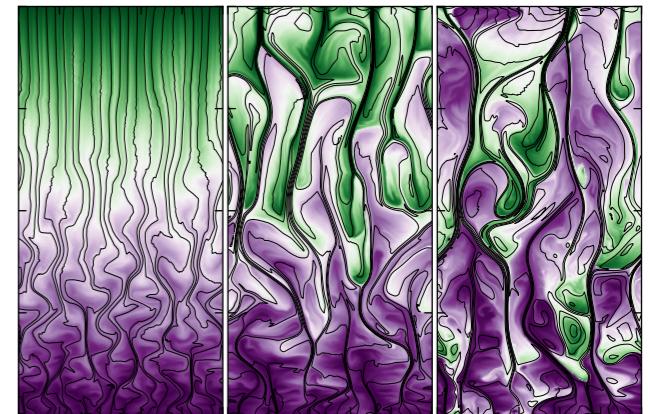
Magnetized cold streams



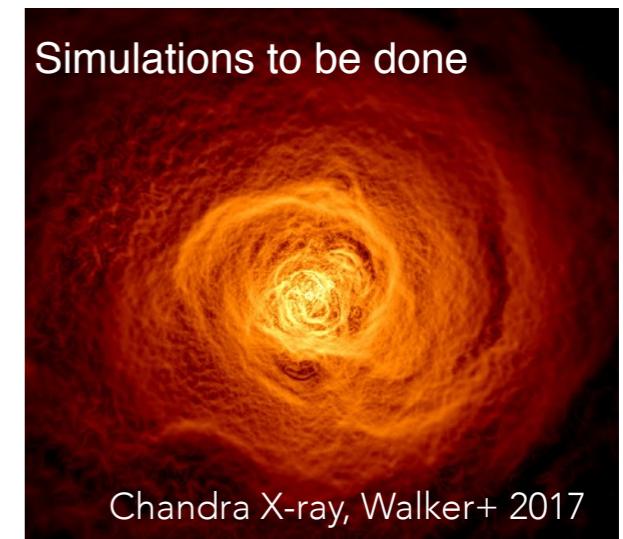
$$r_i \sim 10^{-9} \text{ pc}$$



$$H \sim 10^2 \text{ kpc}$$

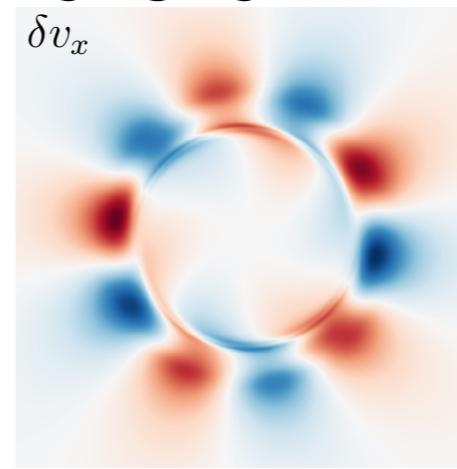


$$L \sim \text{Mpc}$$

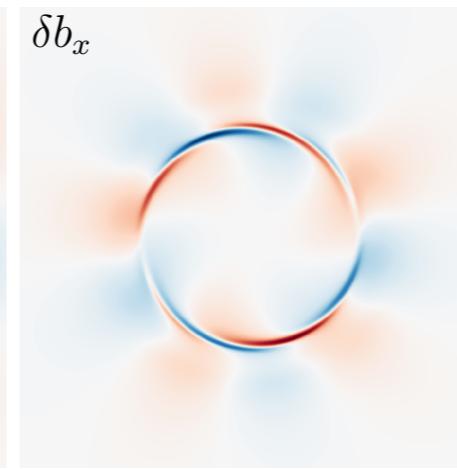


PSECAS

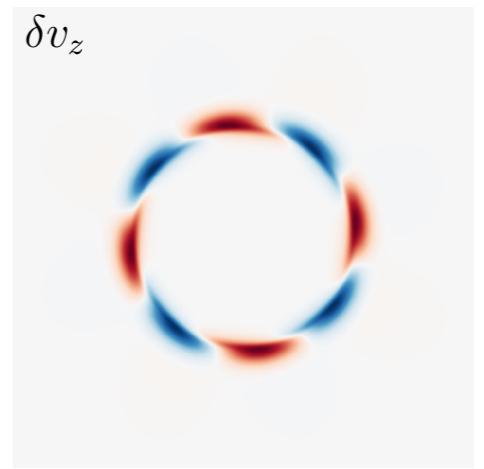
δv_x



δb_x



δv_z



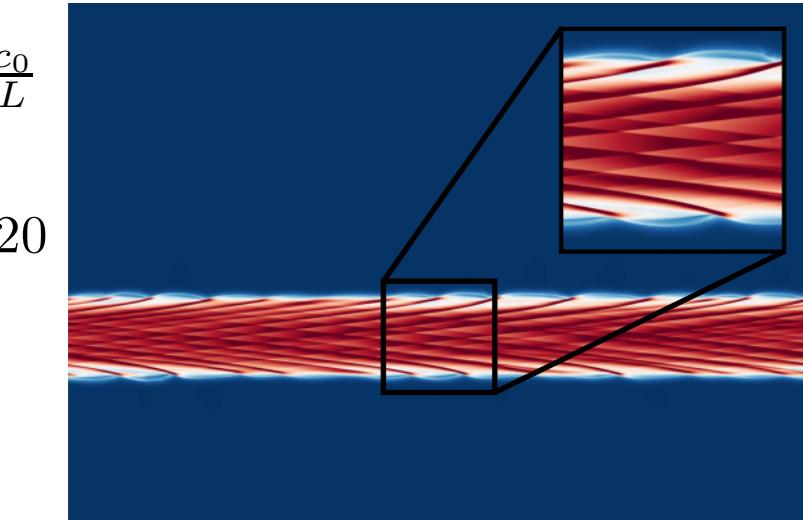
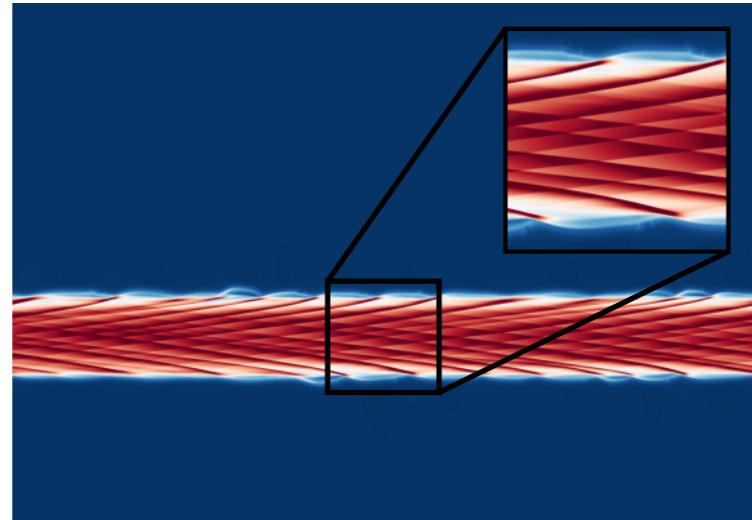
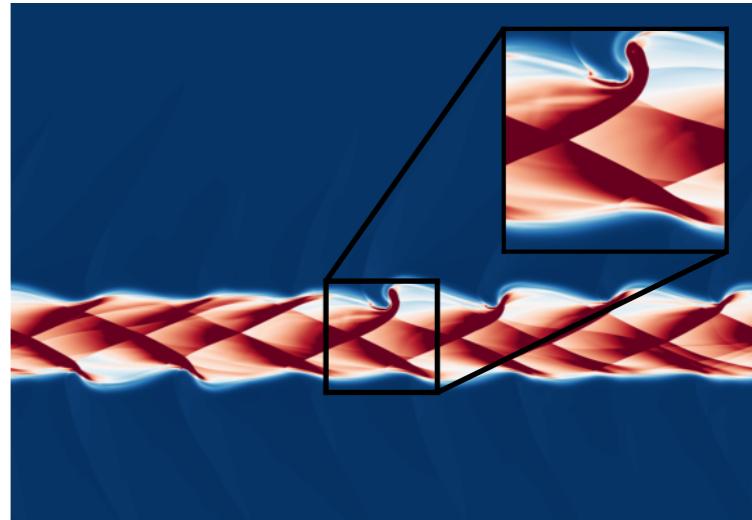


European Research Council
Established by the European Commission

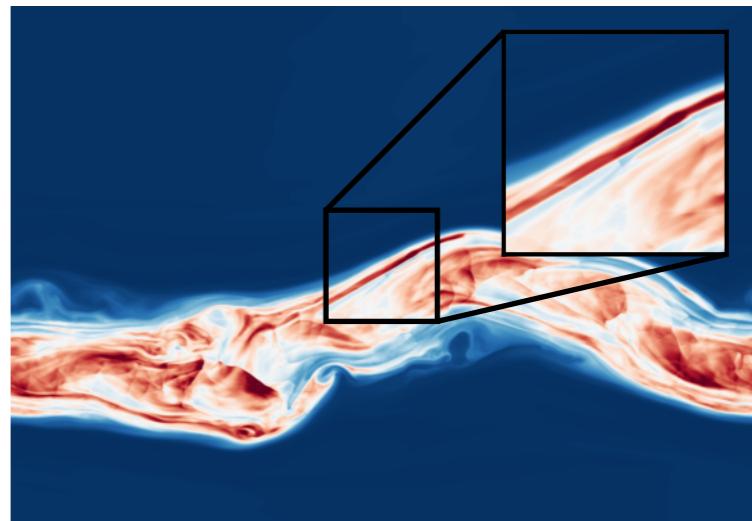
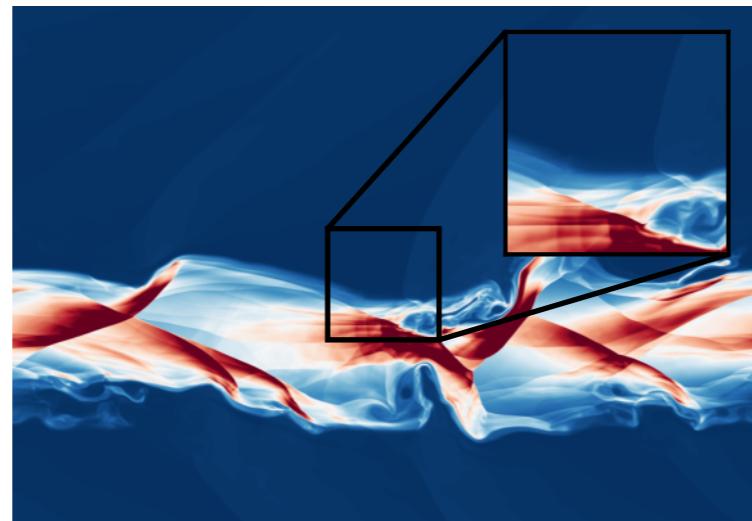
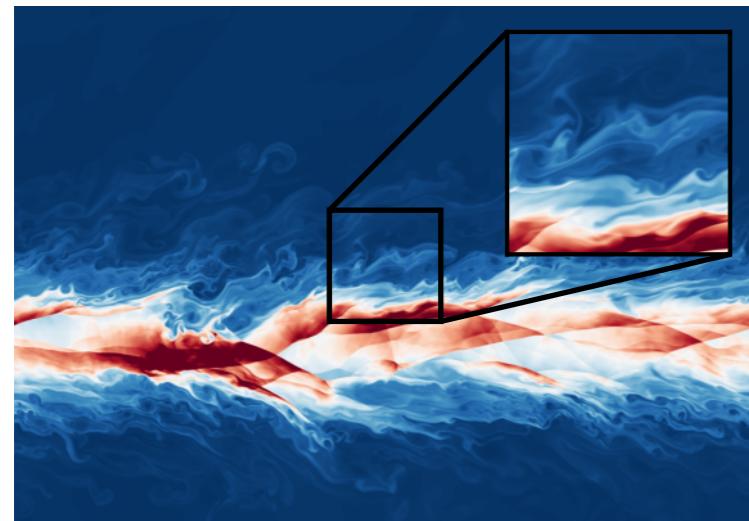
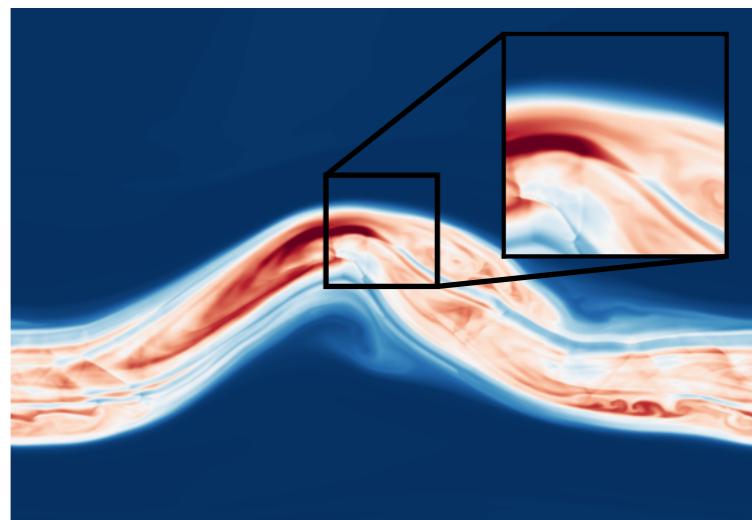
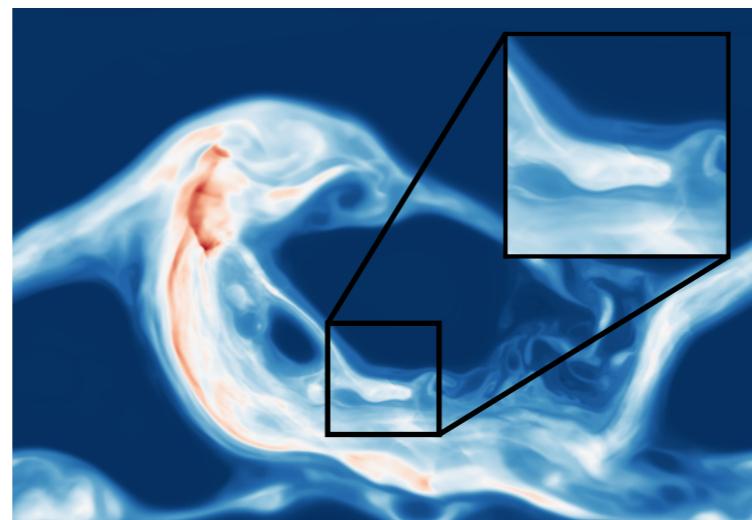
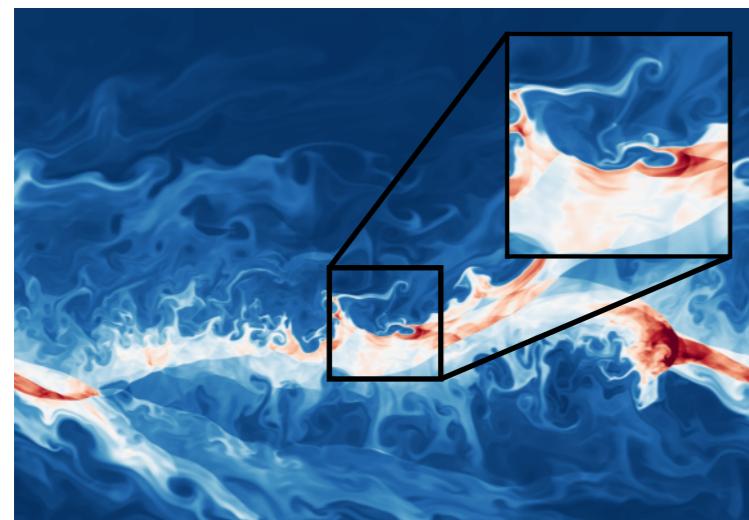
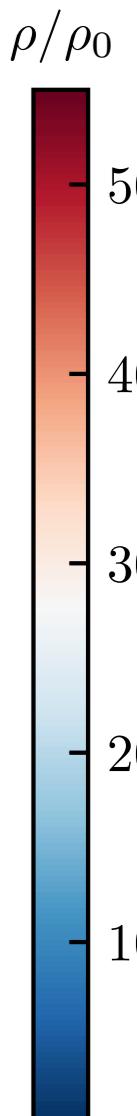


This project has received funding from the European Research Council (ERC) under the European Union’s Horizon 2020 research and innovation program (grant agreement No CRAGSMAN–646955).

hydro

 $t \frac{c_0}{L}$ MHD, $\beta = 10$ $t \frac{c_0}{L}$ MHD, $\beta = 1$ $t \frac{c_0}{L}$ 

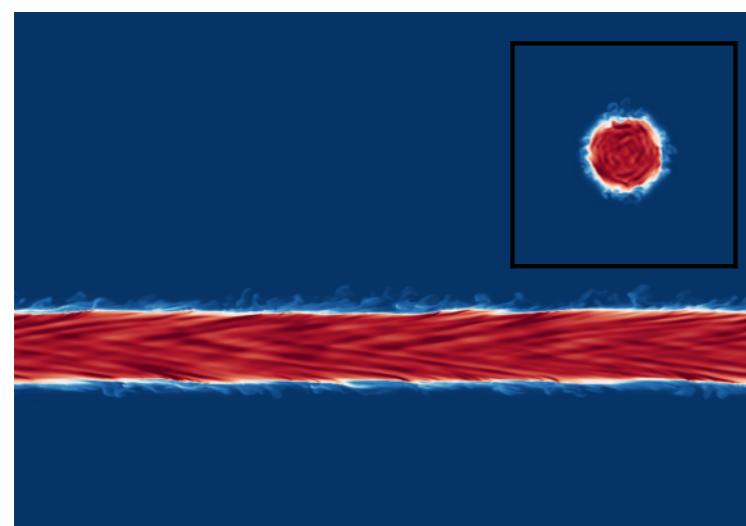
2D

 $t \frac{c_0}{L}$  $t \frac{c_0}{L}$  L |

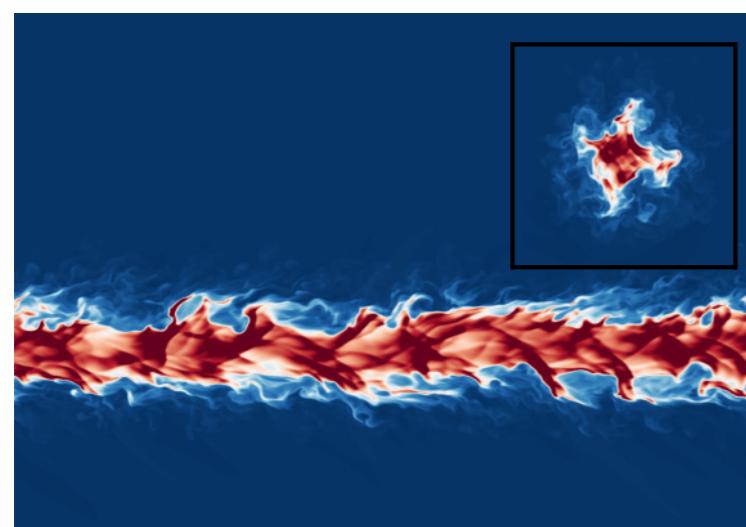
hydro

 $t \frac{c_0}{L}$

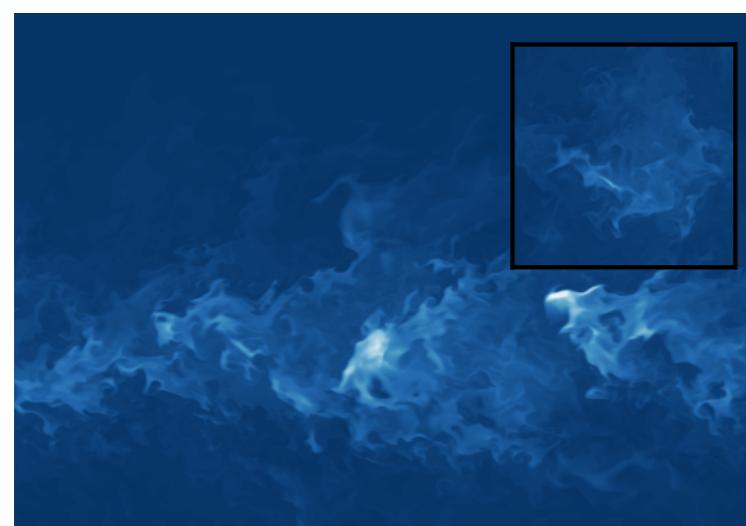
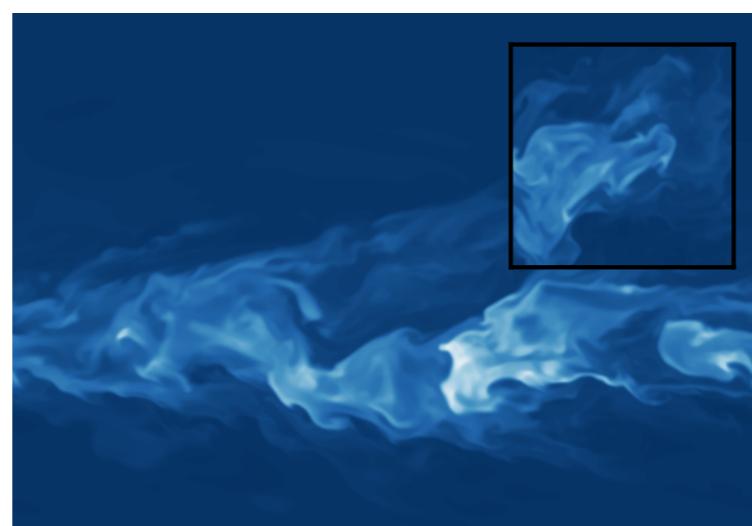
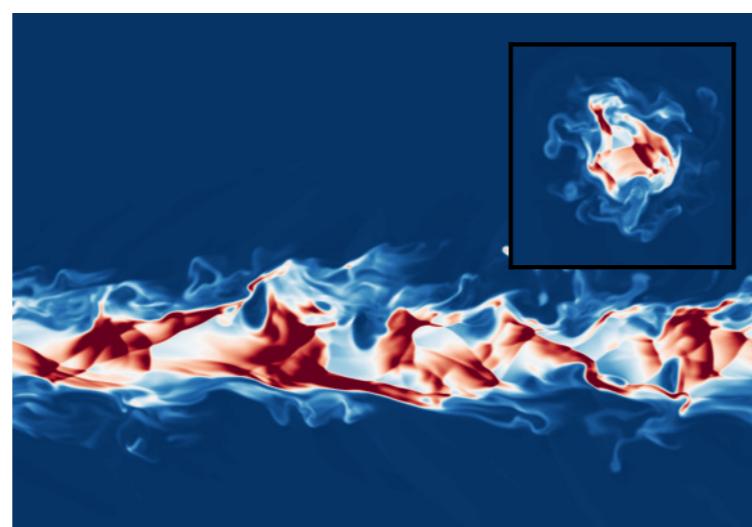
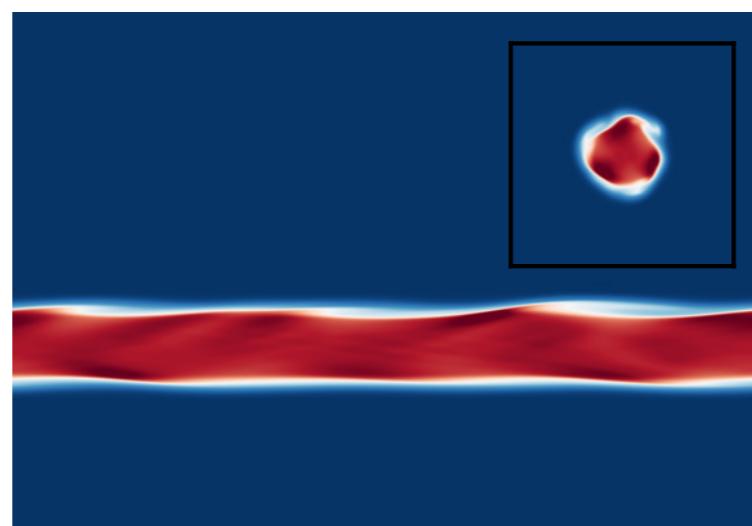
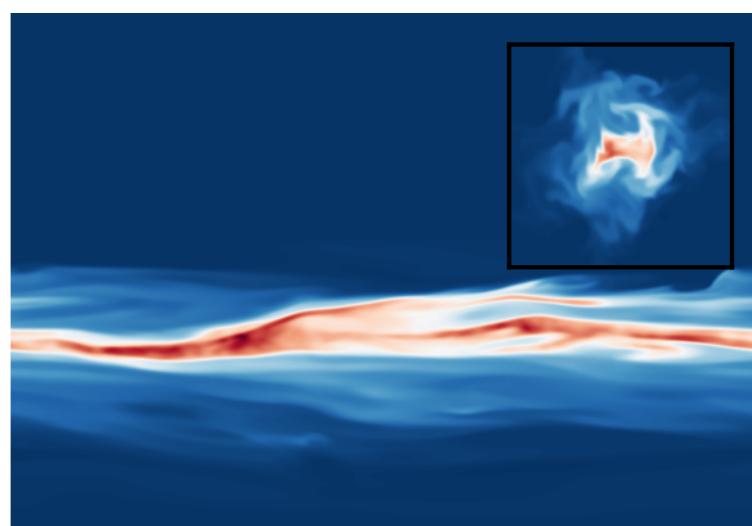
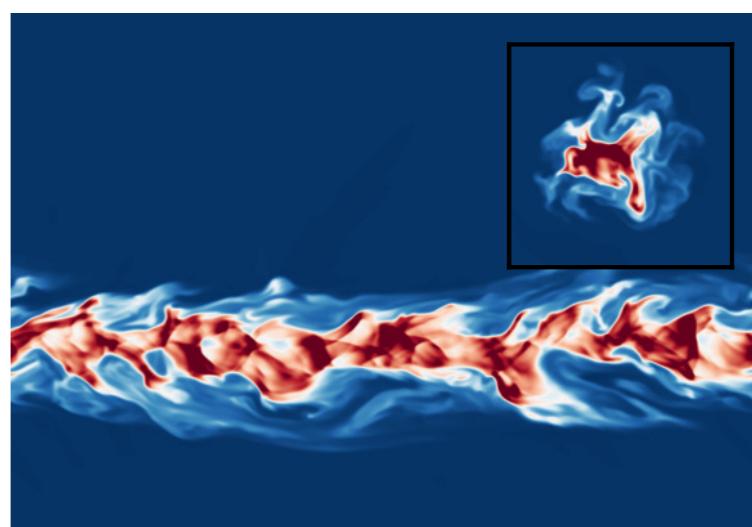
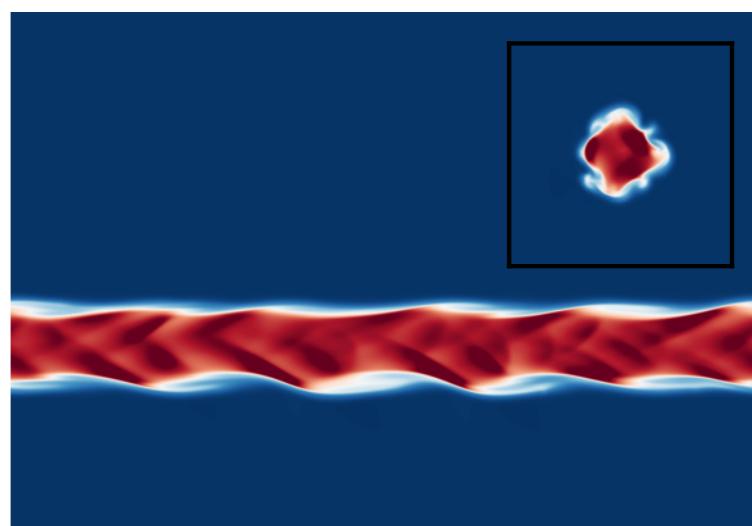
15



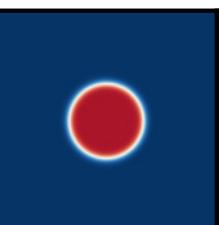
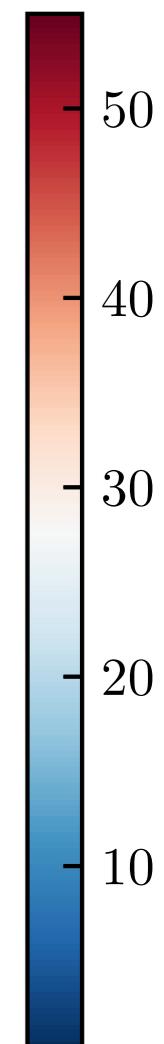
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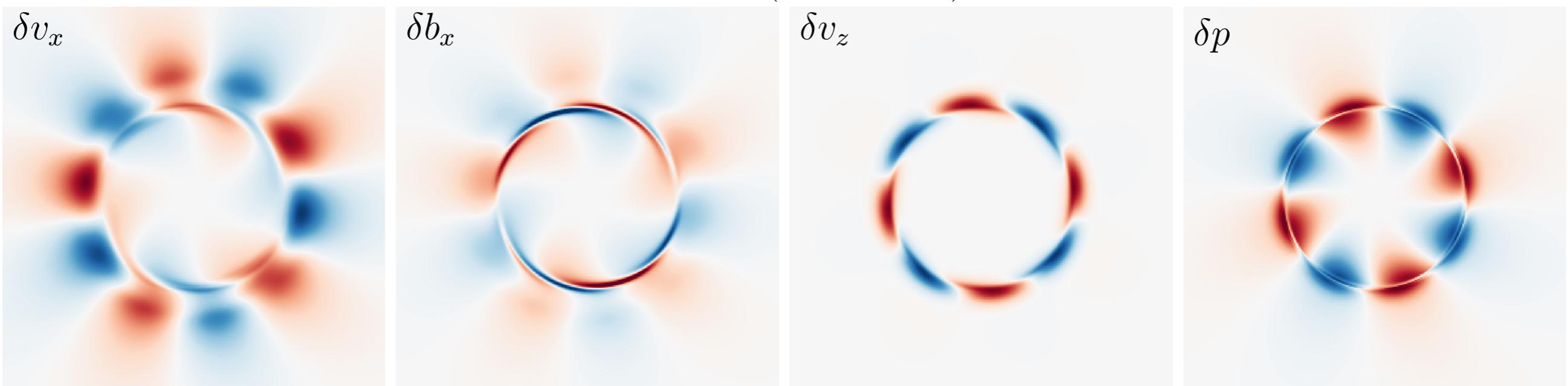
50

MHD, $\beta = 10$ MHD, $\beta = 1$ 

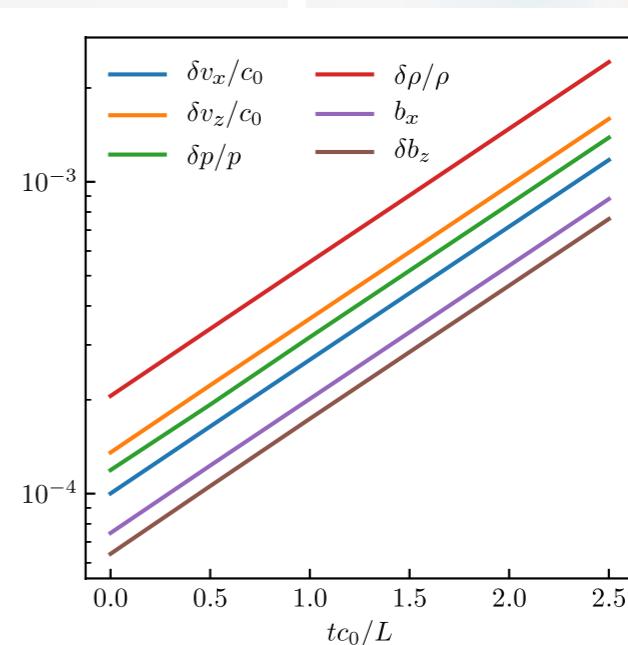
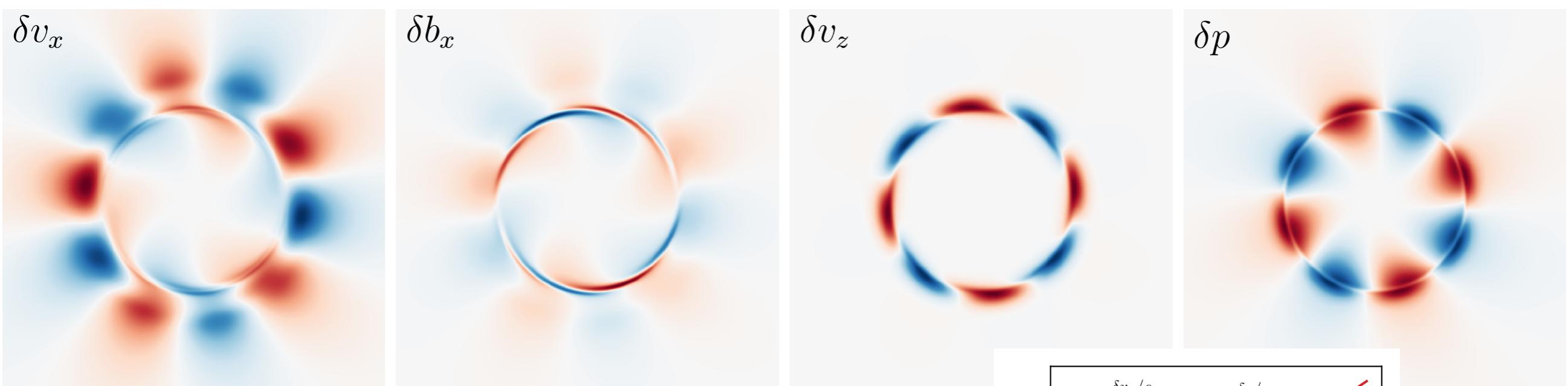
3D

 ρ / ρ_0 

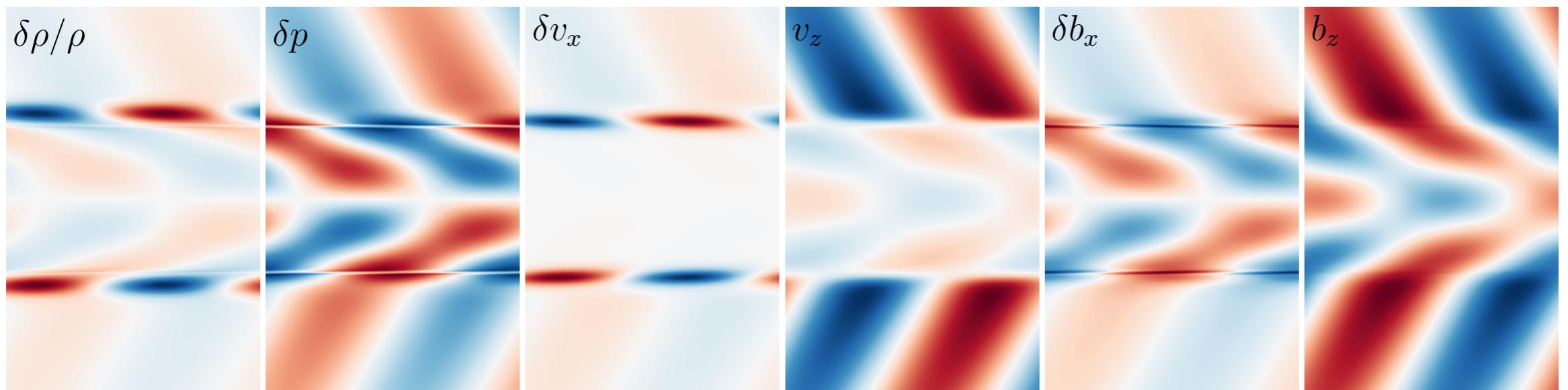
Simulation (ATHENA++)



Theory (PSECAS)



Simulation (ATHENA++)



Theory (PSECAS)

