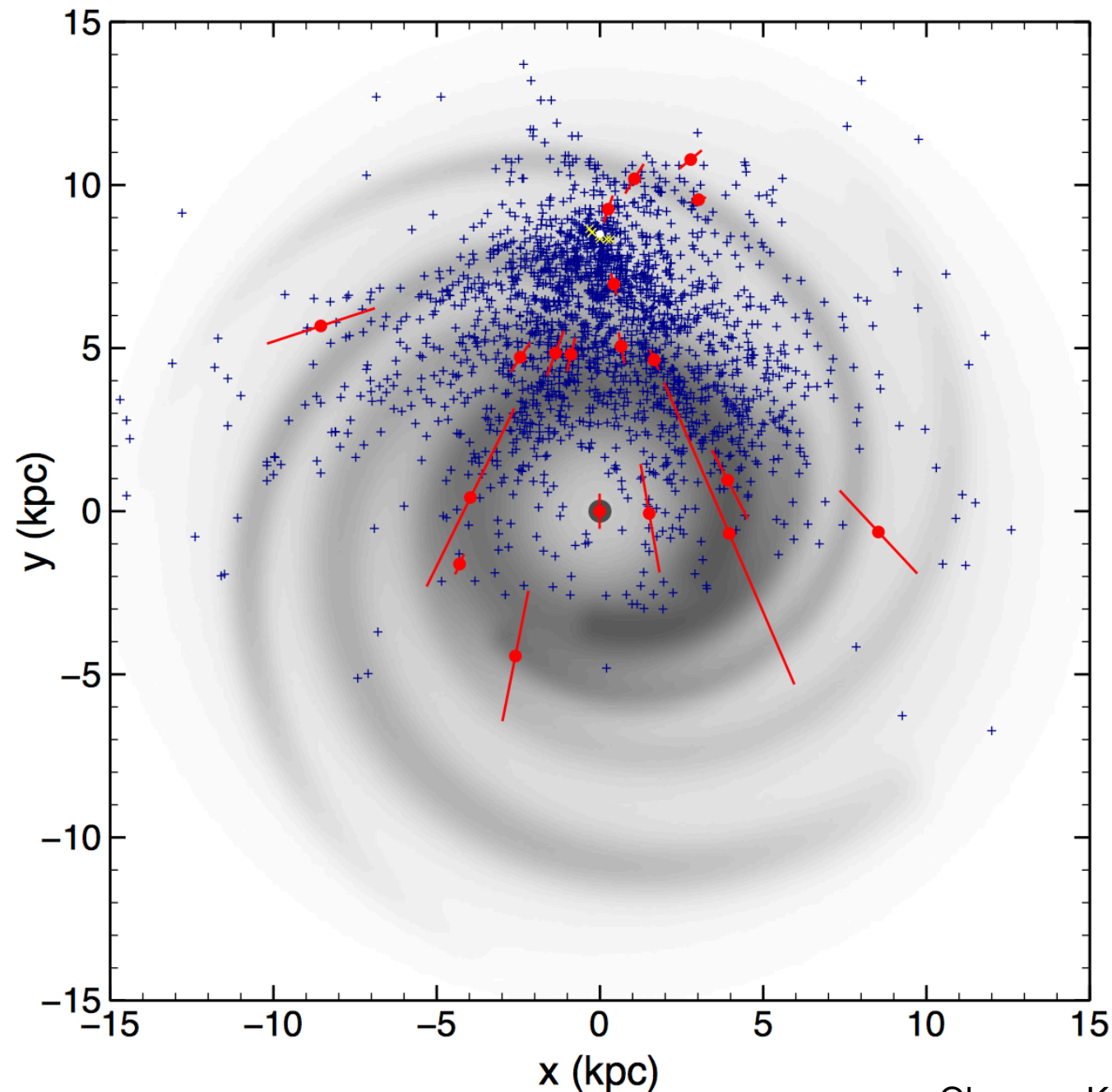


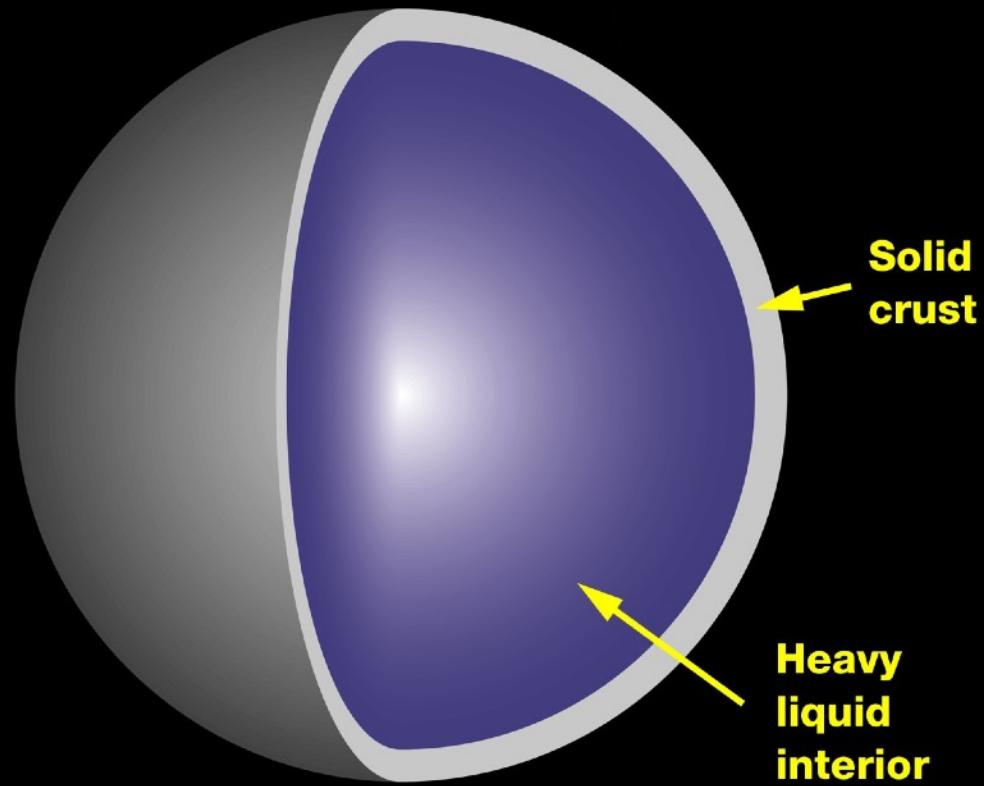
# **Magnetar flares and Fast Radio Bursts**

Andrei Beloborodov  
Columbia University and MPA

## Magnetar population (30 objects)

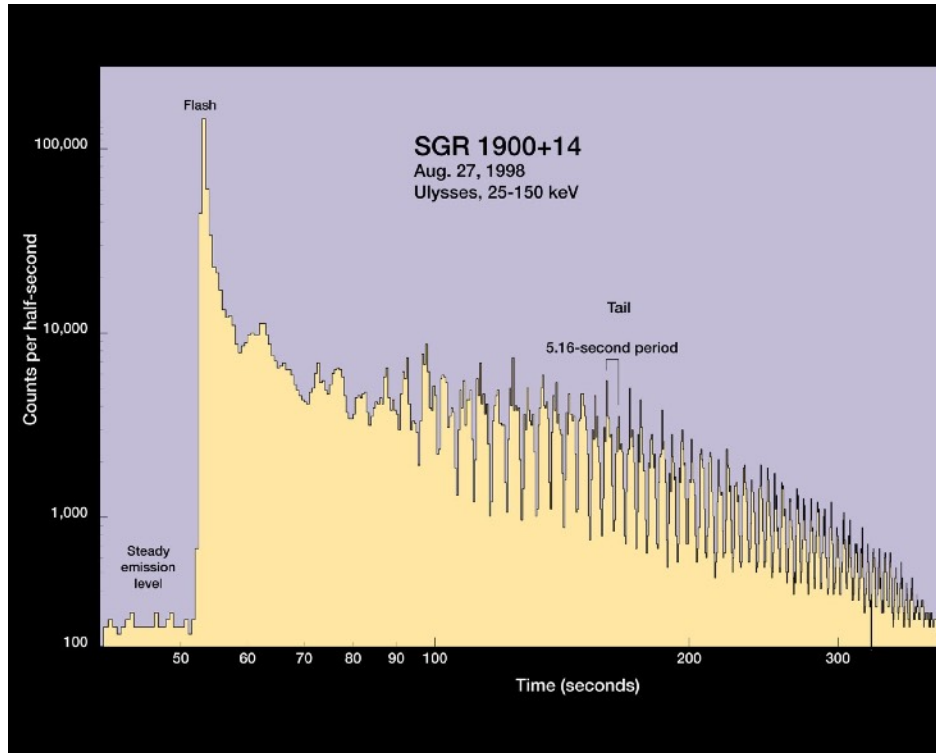


- concentrated near disk plane  $h \sim 20\text{-}30$  pc
- kick velocities  $\sim 200$  km/s
- $\sim 1/3$  found in SN remnants
- no binaries
- age 1-10 kyr
- period 2-12 s



**Development of unbalanced magnetic stresses:  
ambipolar diffusion and Hall drift**

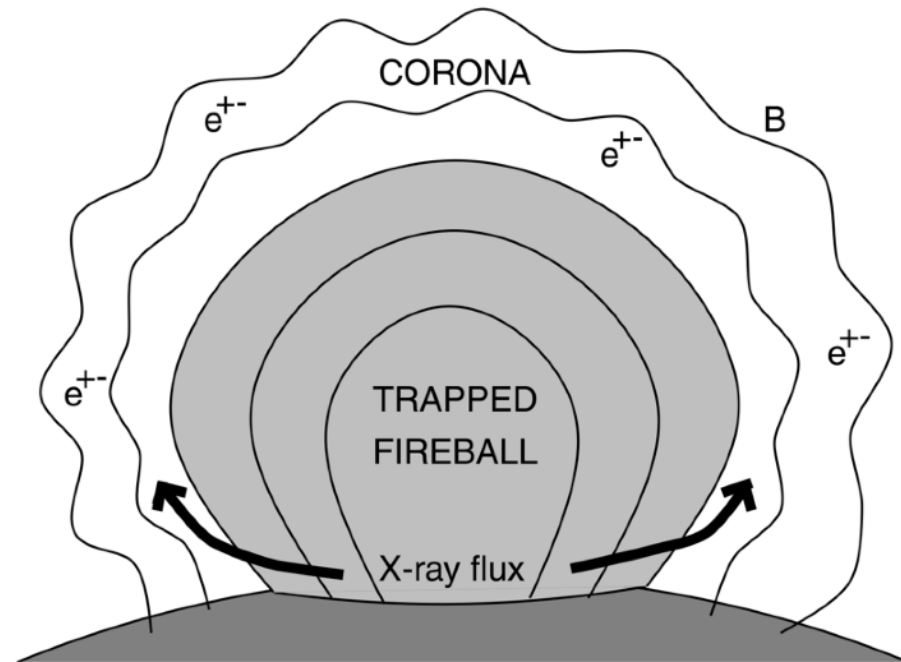
# Giant flares



Tail of the flare:

6 minutes (many rotations)

$$E_{\text{tail}} \sim 10^{44} \text{ erg}$$



Thompson, Duncan 1996, 2001  
Thompson, Blaes 1998

Is the fireball heated  
by turbulence cascade?

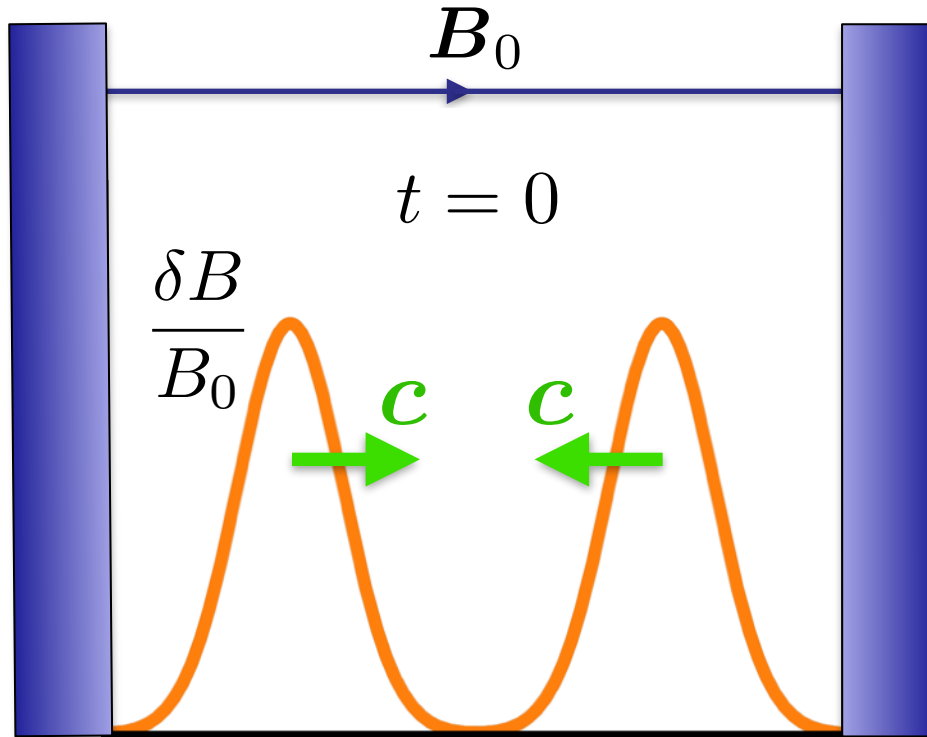
# Heating source for giant flares:

## I. Turbulence

Sudden deformation of the magnetosphere leads  
to an Alfvén wave cascade

How long does it take to dissipate the waves?

# Simulation setup: colliding Alfvén wave packets



uniform background field

two symmetric wave packets

$$\delta \mathbf{B} \perp \delta \mathbf{E} \quad \delta \mathbf{B}, \delta \mathbf{E} \perp \mathbf{B}_0$$

reflecting boundaries  
( $\Leftrightarrow$  periodic box)

force-free electrodynamics

3D

Parameter:

$$\xi \equiv \frac{\delta B_{\max}}{B_0}$$

# Force-free electrodynamics (FFE)

$$\partial_t \mathbf{B} = -\nabla \times \mathbf{E}$$

$$\partial_t \mathbf{E} = \nabla \times \mathbf{B} - \mathbf{J}$$

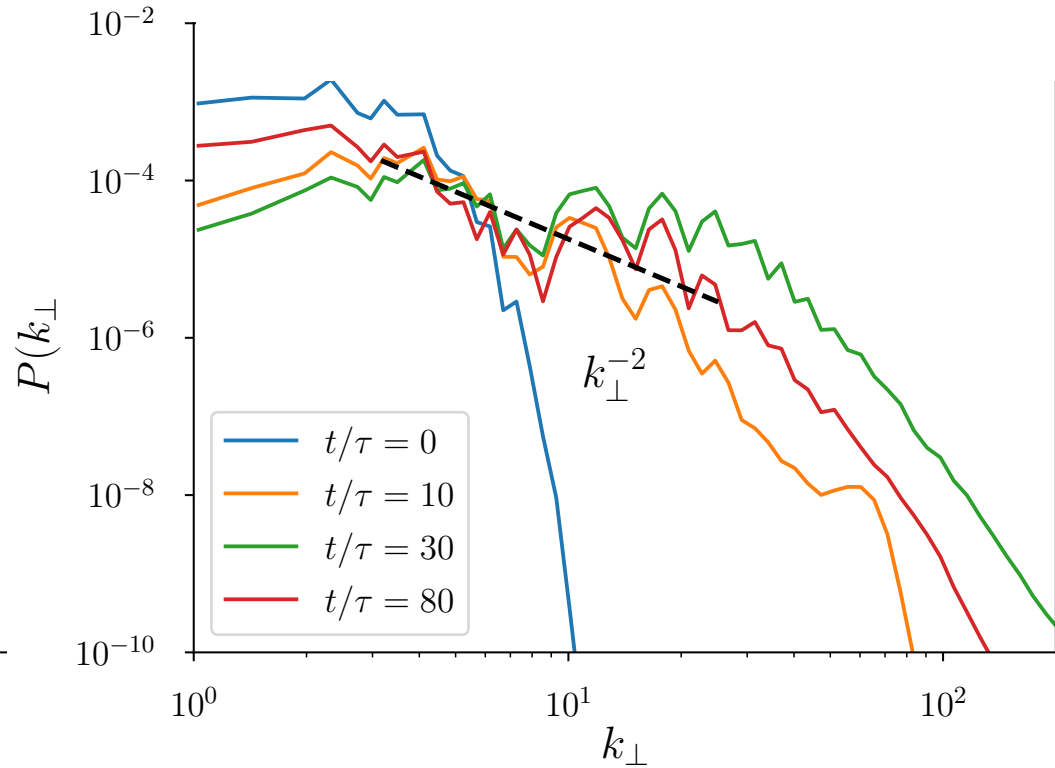
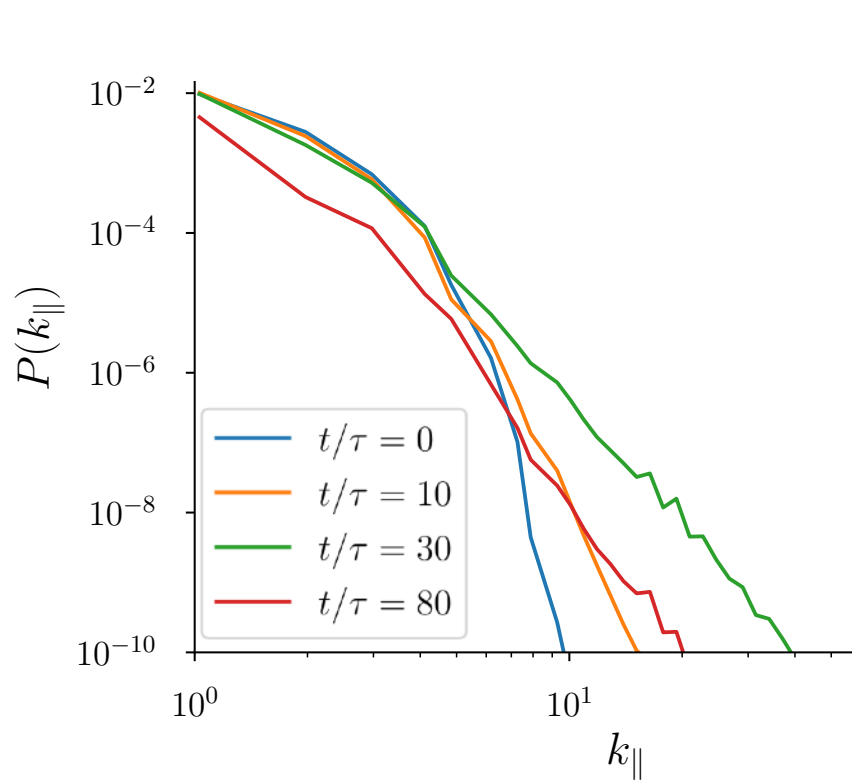
$$\rho \mathbf{E} + \frac{\mathbf{J} \times \mathbf{B}}{c} = 0 \quad \Rightarrow \quad \mathbf{J} = \frac{\mathbf{B} \cdot \nabla \times \mathbf{B} - \mathbf{E} \cdot \nabla \times \mathbf{E}}{B^2} \mathbf{B} + \nabla \cdot \mathbf{E} \frac{\mathbf{E} \times \mathbf{B}}{B^2}$$

Eigen modes in FFE:

- Alfvén wave
- “Fast” wave (no electric current, as if in vacuum)

Both modes propagate with the speed of light

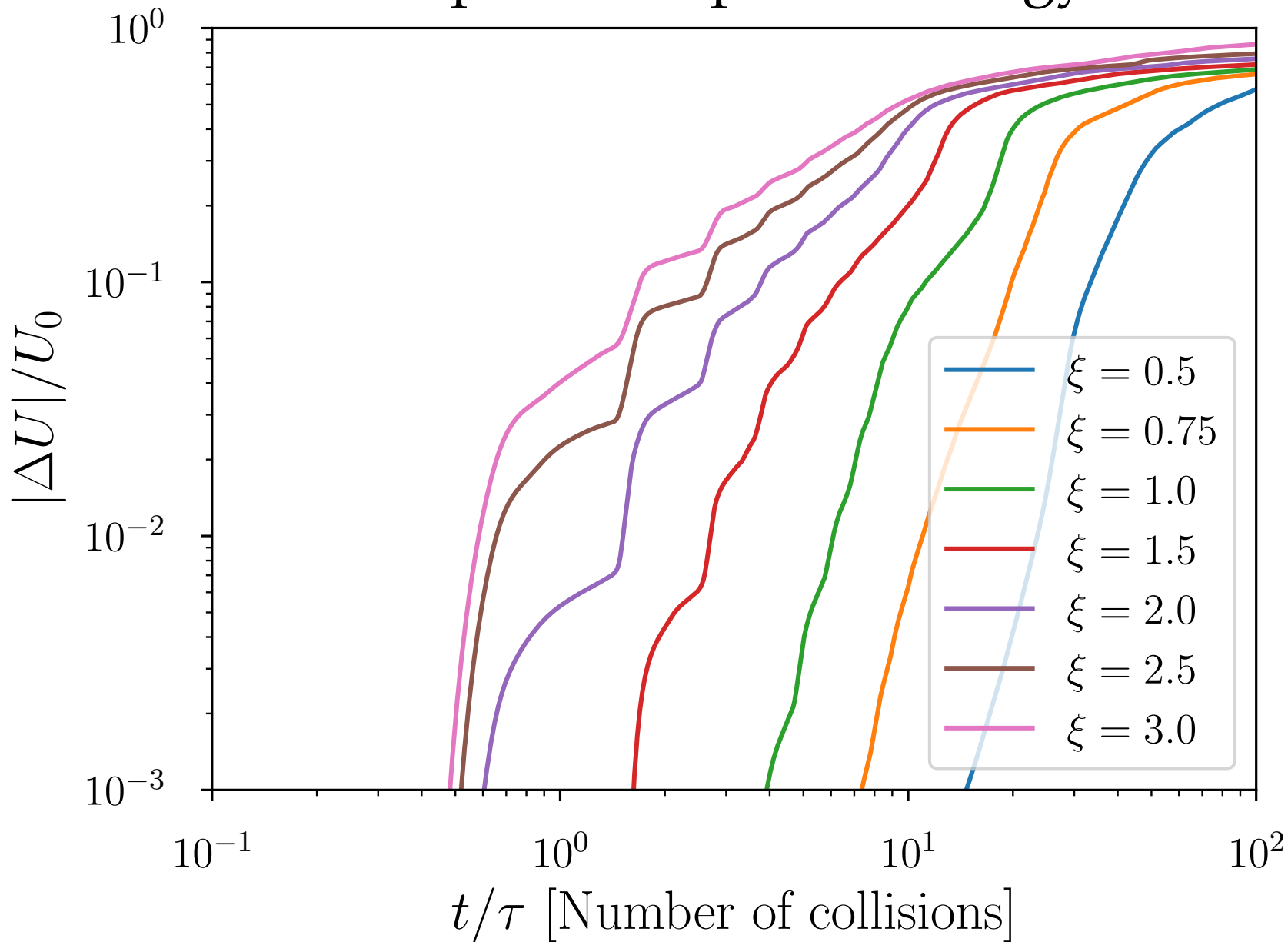
# Development of turbulence cascade



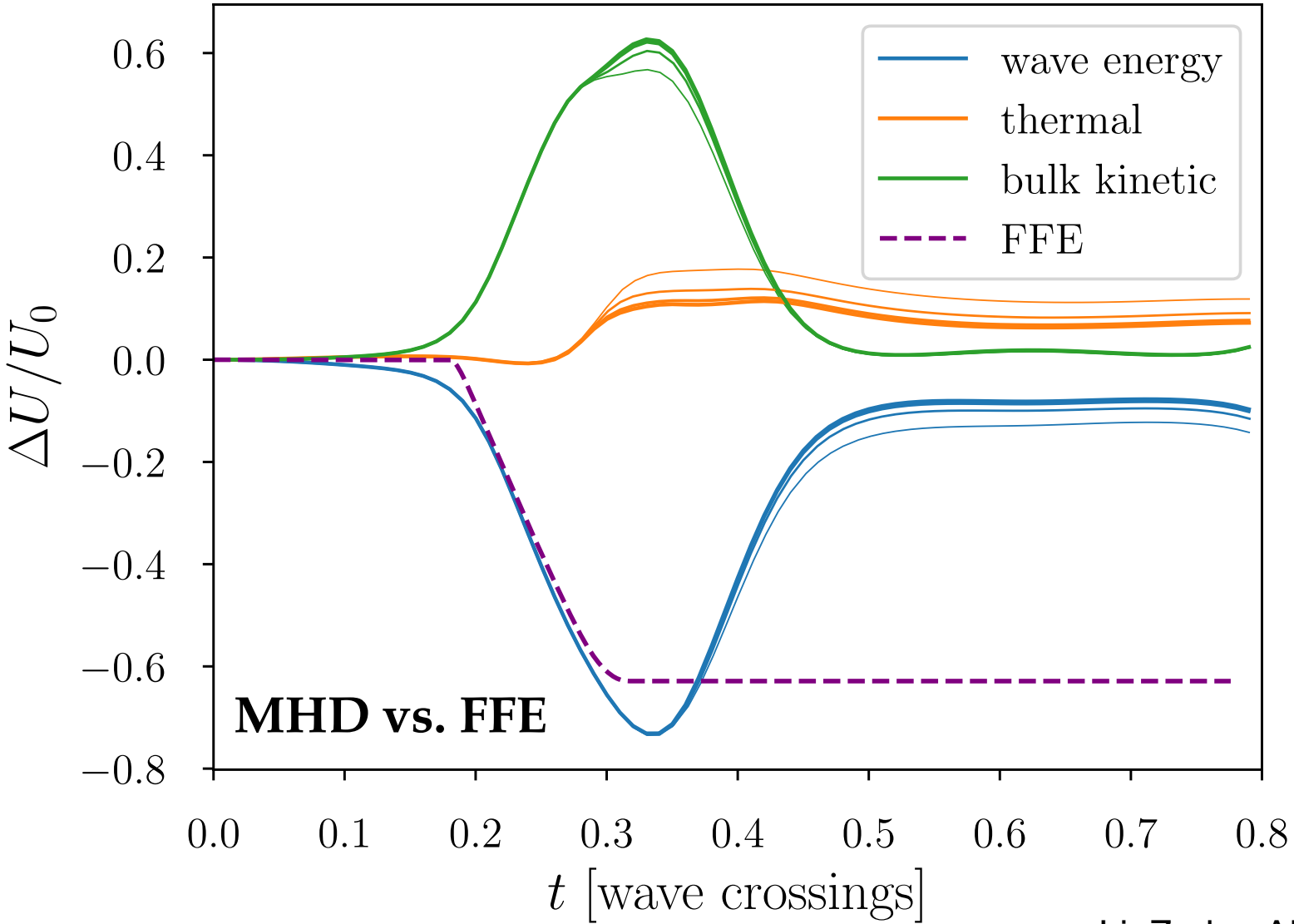
the cascade is dominated by waves with  $k_{\perp} \gg k_{\parallel}$   
( wavevectors almost perpendicular to  $\mathbf{B}_0$  )



# dissipation of packet energy



# Remark 1. Wave polarization and limitations of FFE



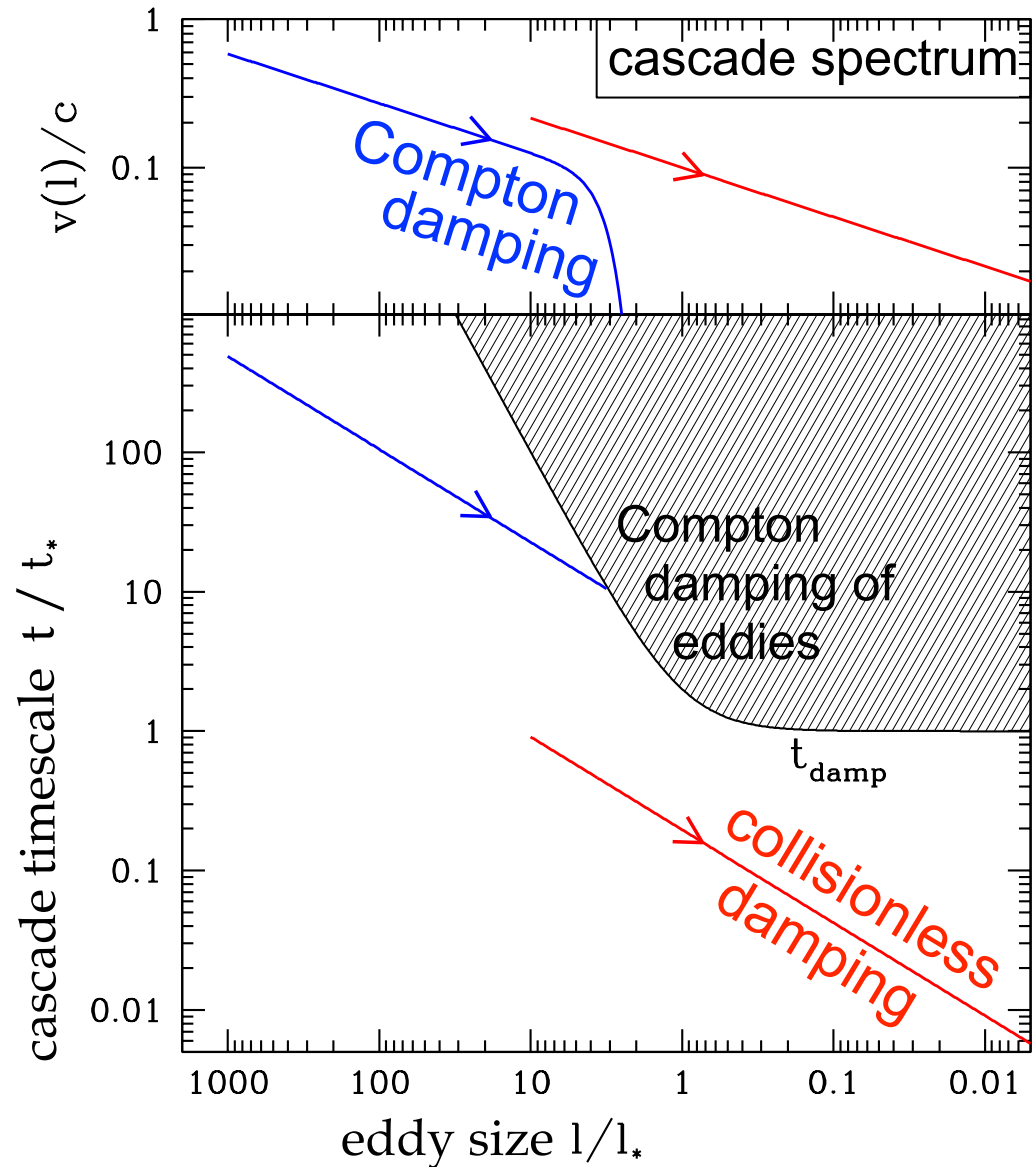
# Remark 2. Cascade is dissipated by Compton drag

$$\tau > \left( \frac{v_0}{c} \frac{1+w}{w} \right)^3$$

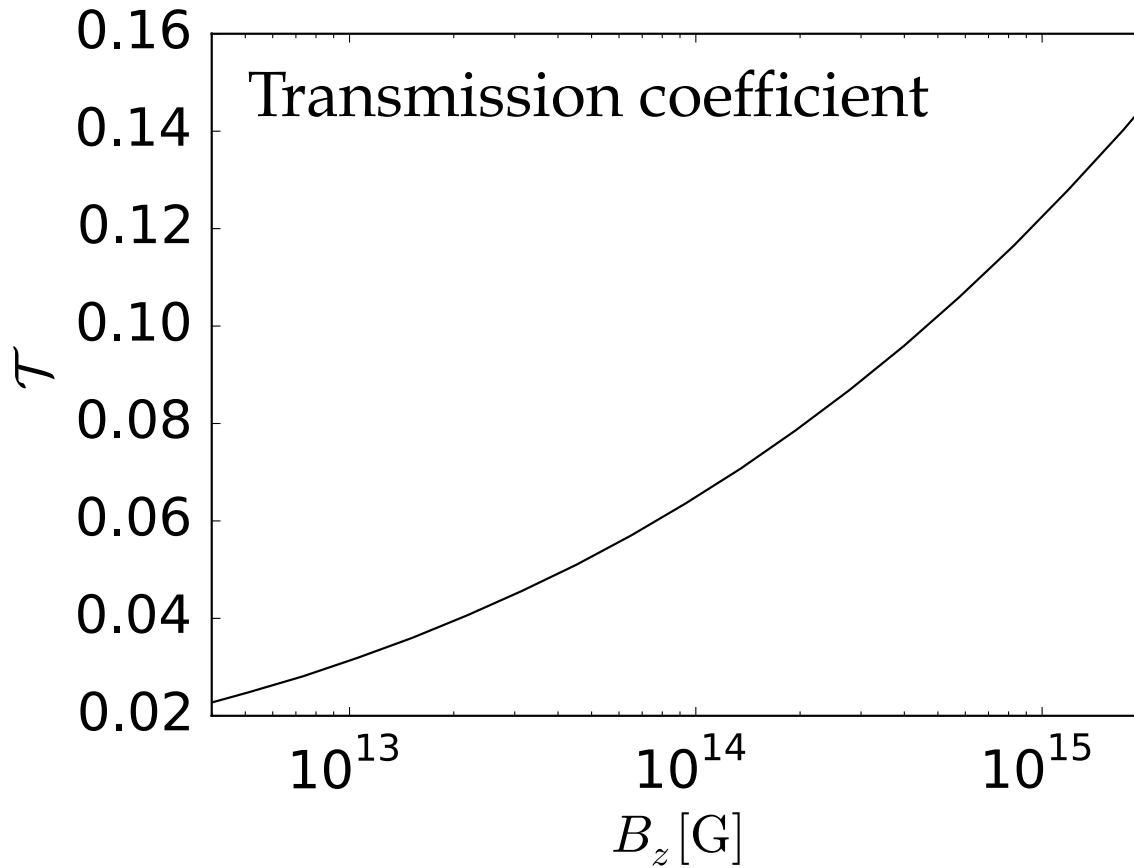
Compton damping

$$\left[ w \equiv \frac{U + P}{\rho c^2} \right]$$

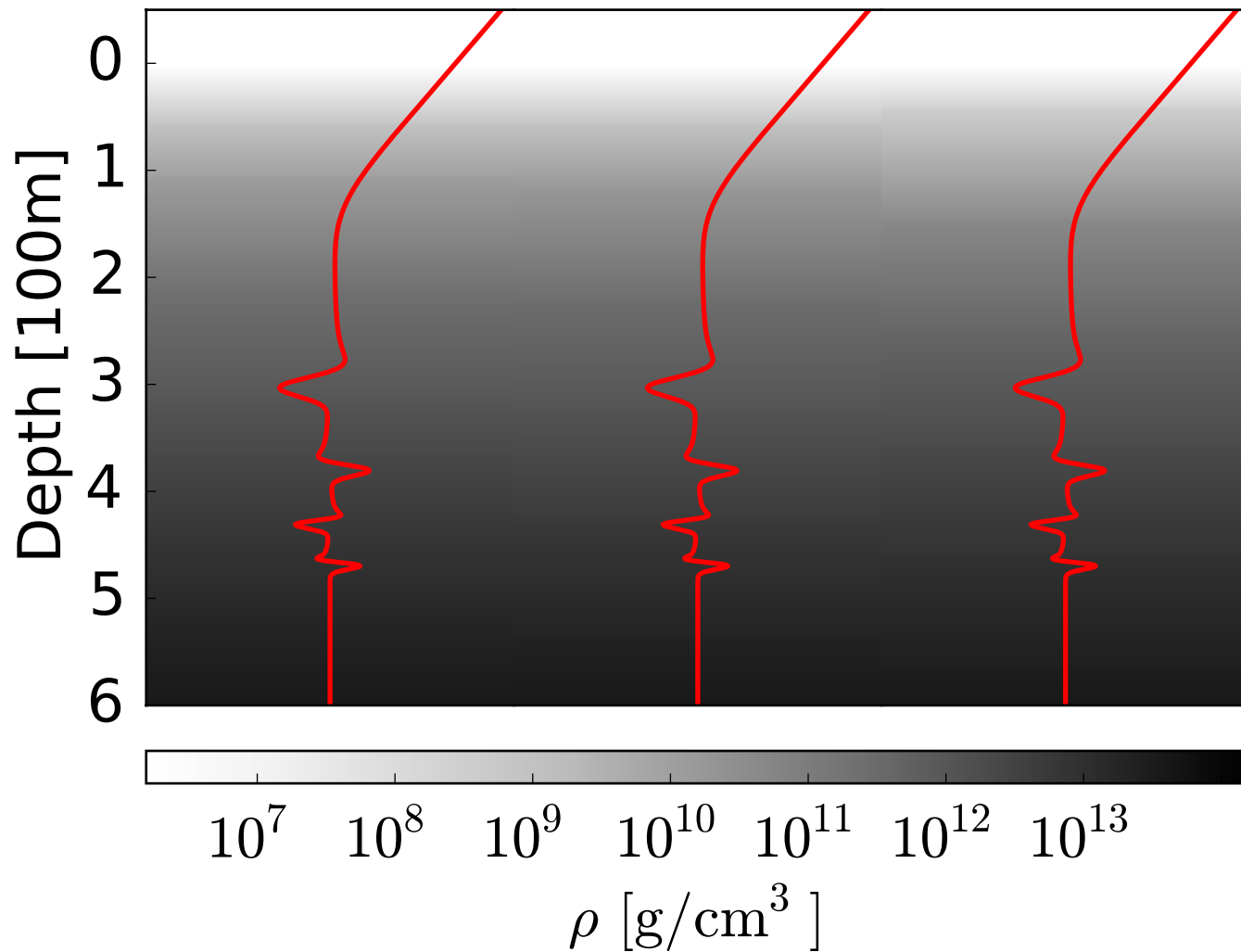
relativistic enthalpy



# Remark 3. Draining turbulence into the star



$$v^2(z) = \frac{B_z^2/4\pi + \mu(z)}{B_z^2/4\pi c^2 + \rho(z)}$$



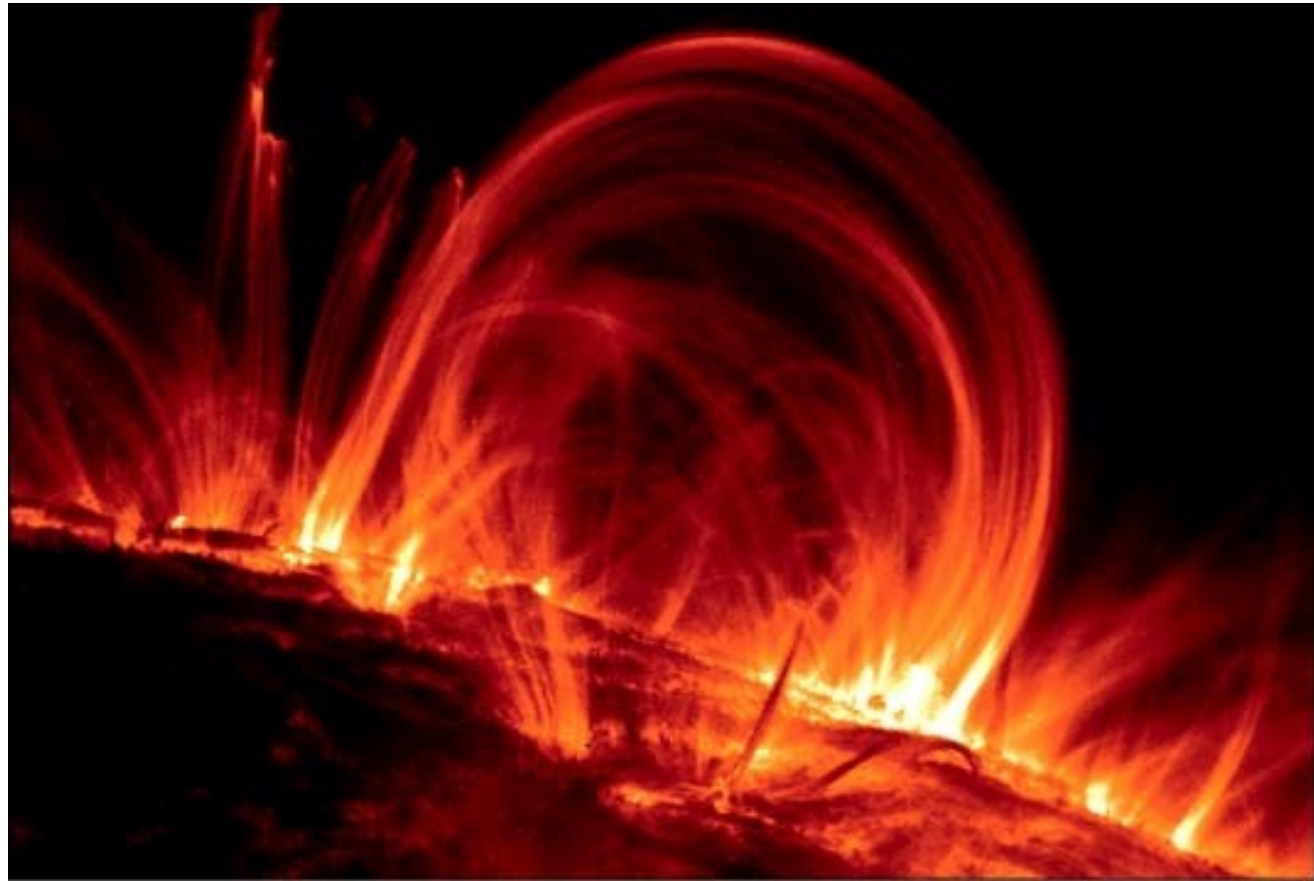
Li & AB 2015

Plastic damping of waves melts the crust; afterglow  $\sim$  month

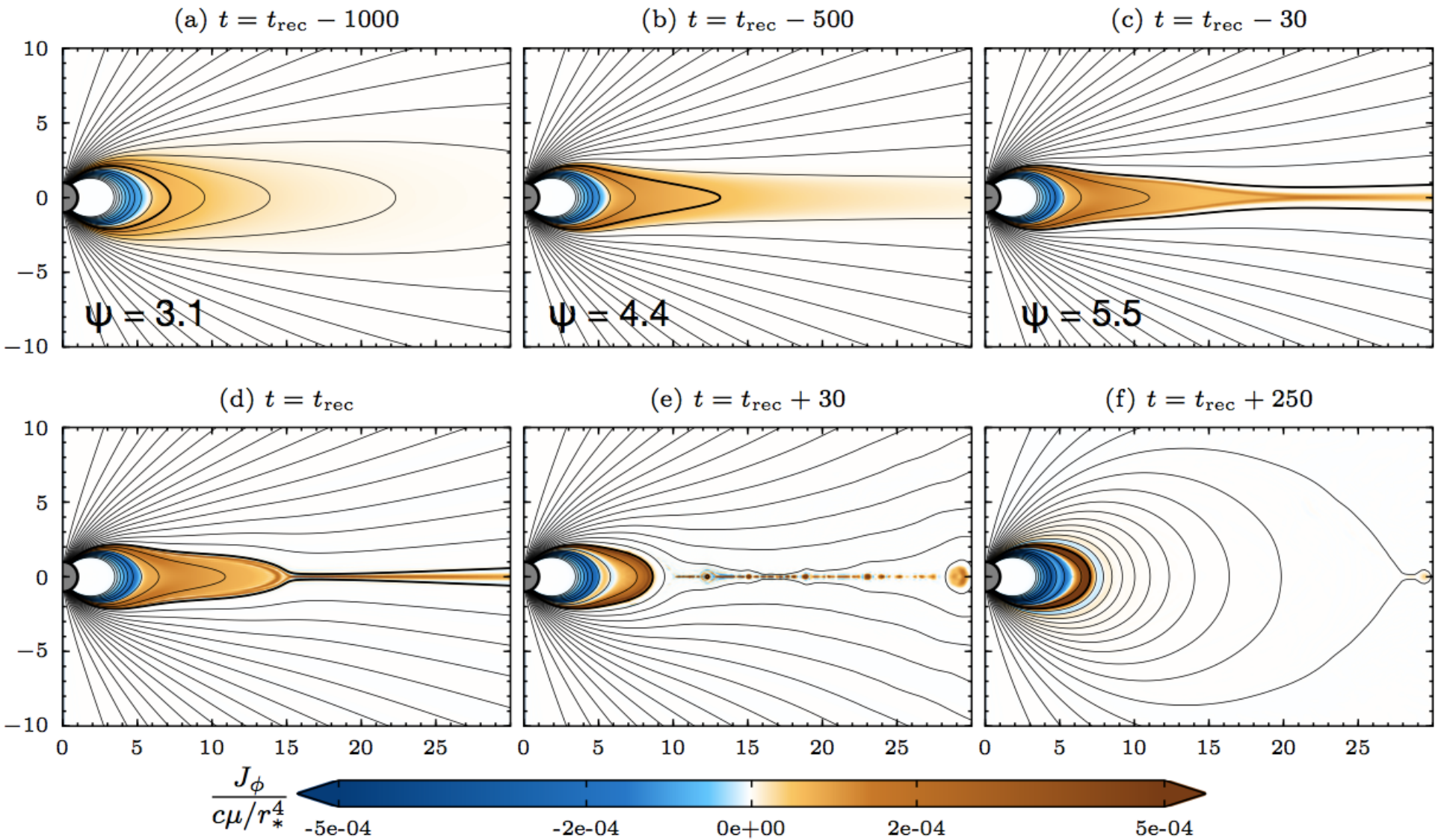
# Heating source for giant flares: II. Magnetic reconnection

Overtwisted  
magnetic loops  
(cf. solar flares)

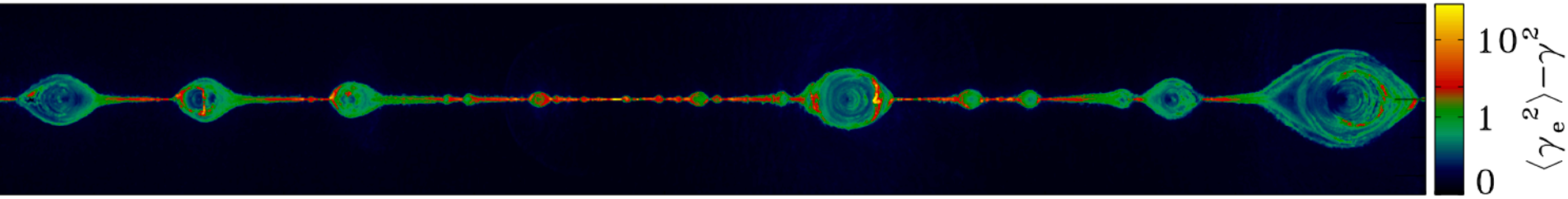
Thompson & Duncan 1996  
Lyutikov 2006  
Parfrey, AB, Hui 2013  
Carrasco et al. 2019



# Loss of magnetic equilibrium and reconnection

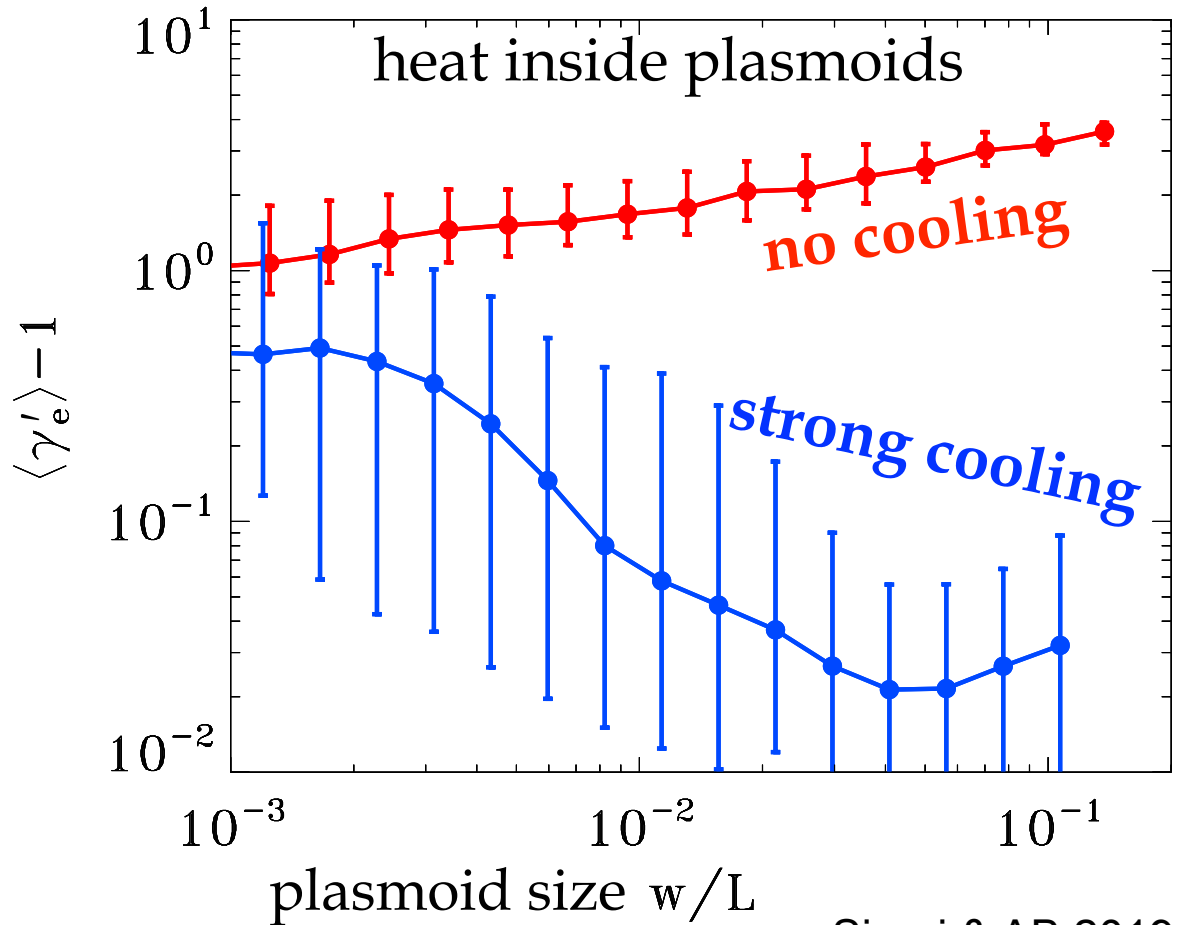


# Radiative reconnection



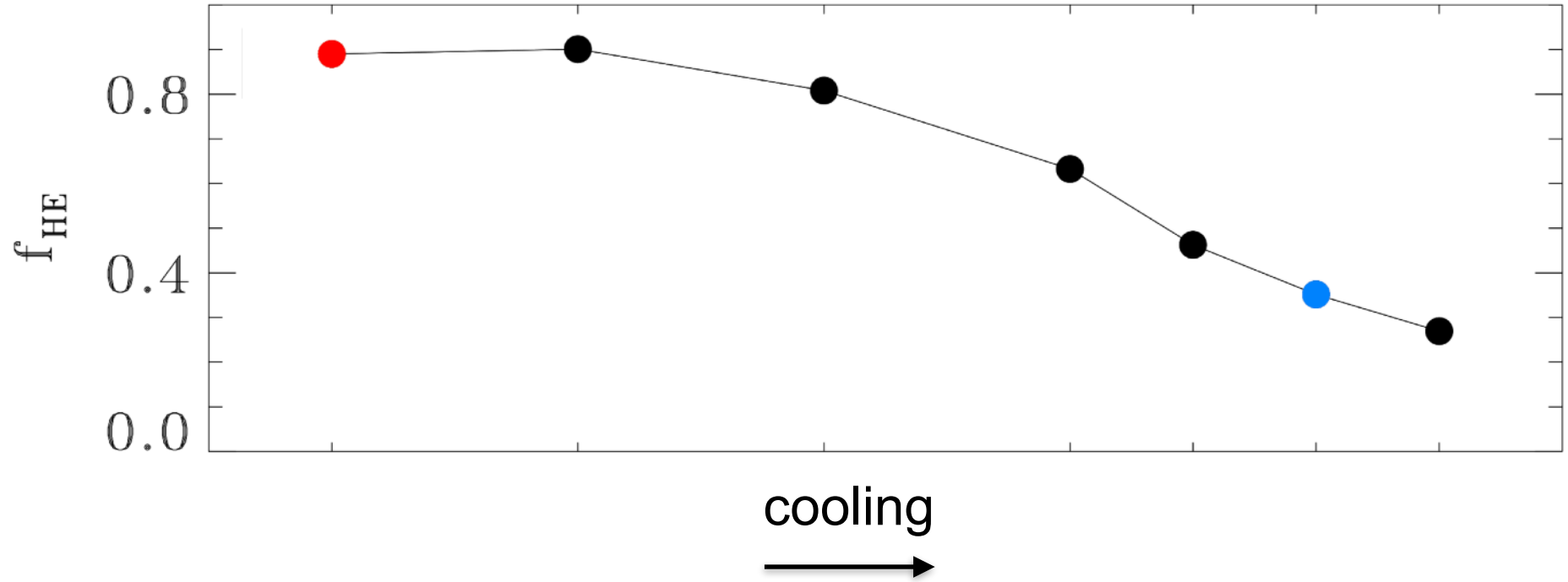
$$t_{\text{cool}} \ll L/c$$

AB 2017  
Werner et al. 2019  
Sironi & AB 2019



Sironi & AB 2019





**Magnetars:**  
“ultra-radiative”  
reconnection

# **Blast waves from magnetar flares and cosmological fast radio bursts (FRBs)**

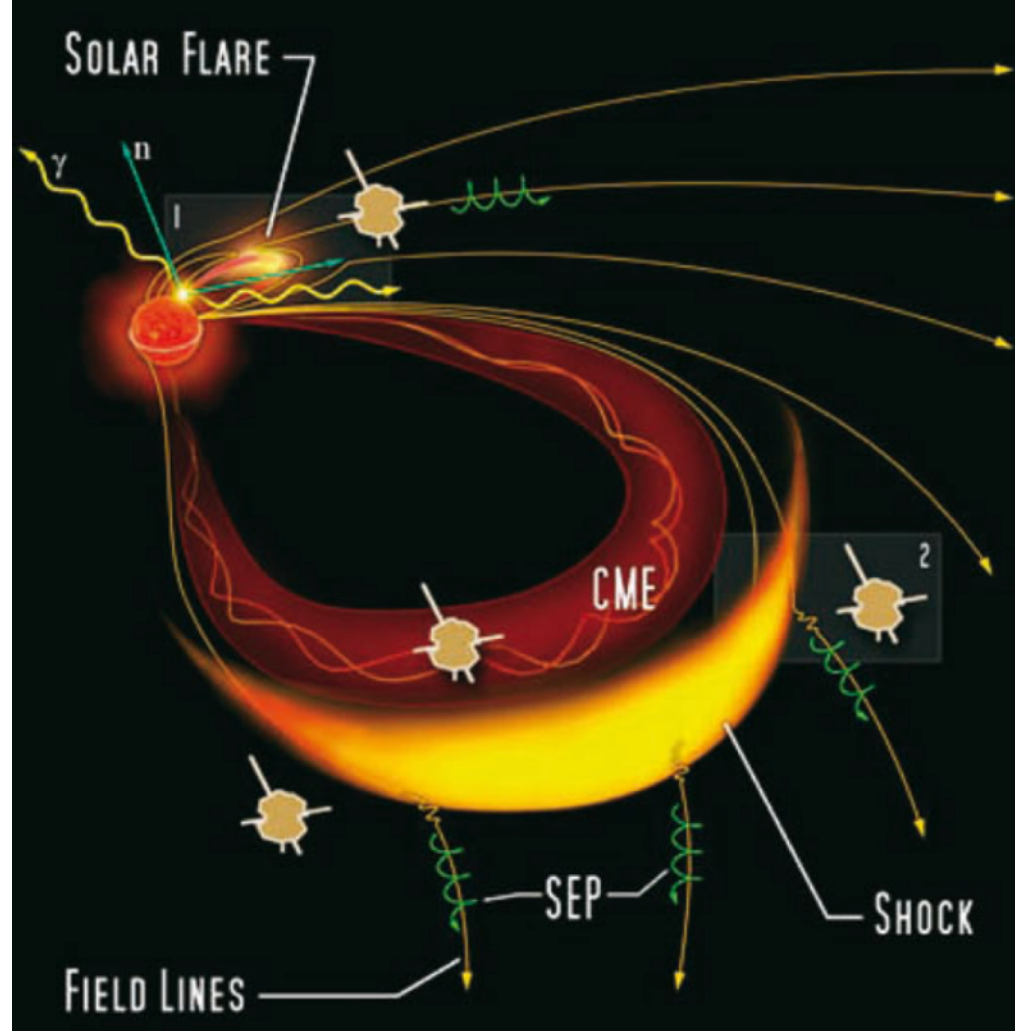
reconnection ejects ultra-relativistic plasmoids  
into the magnetar wind  
(cf. solar flares)

# Blast waves from magnetar flares:

**blast Lorentz factor**

$$\frac{\Gamma}{\Gamma_w} \approx 100 \left( \frac{L_{f,47}}{L_{w,39}} \right)^{1/4}$$

**pre-shock wind Lorentz factor**



cf. solar CME

# Pre-flare wind from the rotating twisted magnetosphere

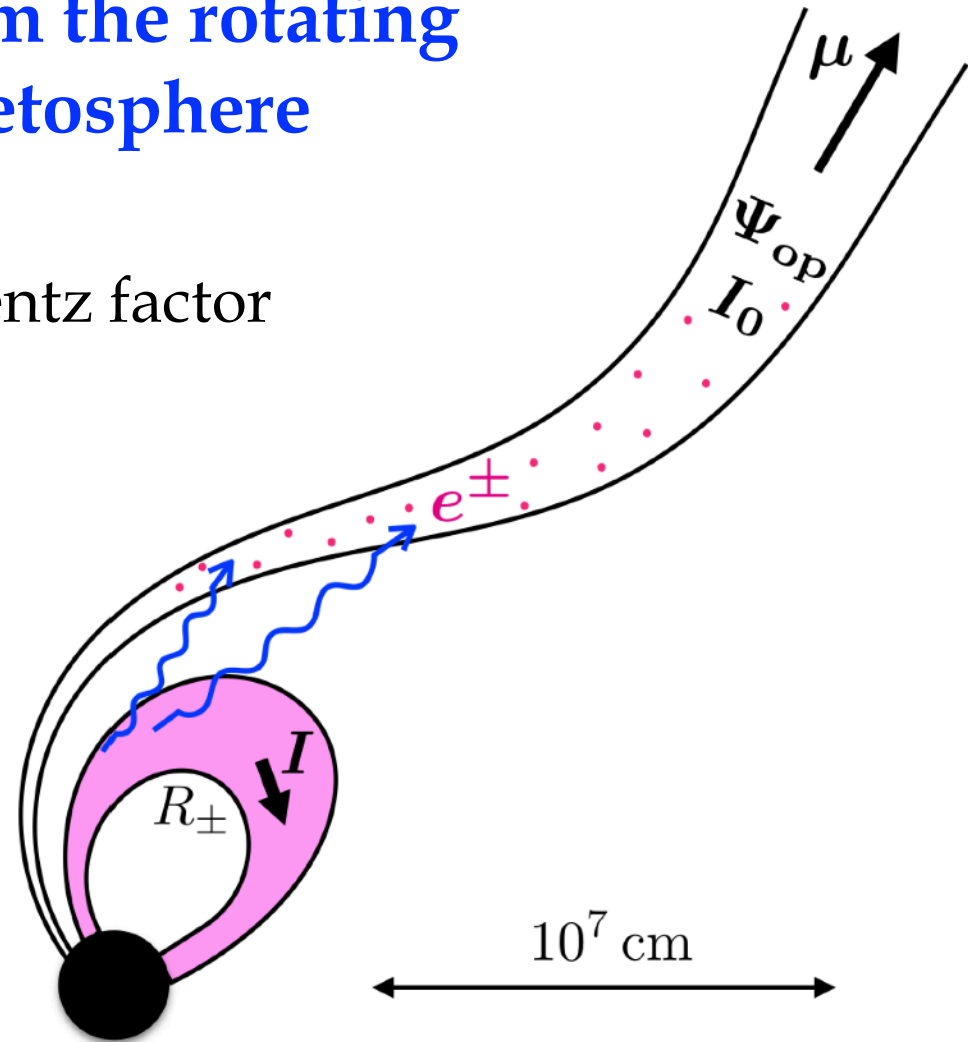
e<sup>±</sup>- loading determines Lorentz factor of the wind:  $\Gamma_w \approx 3\eta^{1/3}$

$$\eta = \frac{L_w}{\dot{N}m_e c^2} \sim 10^2 - 10^4$$

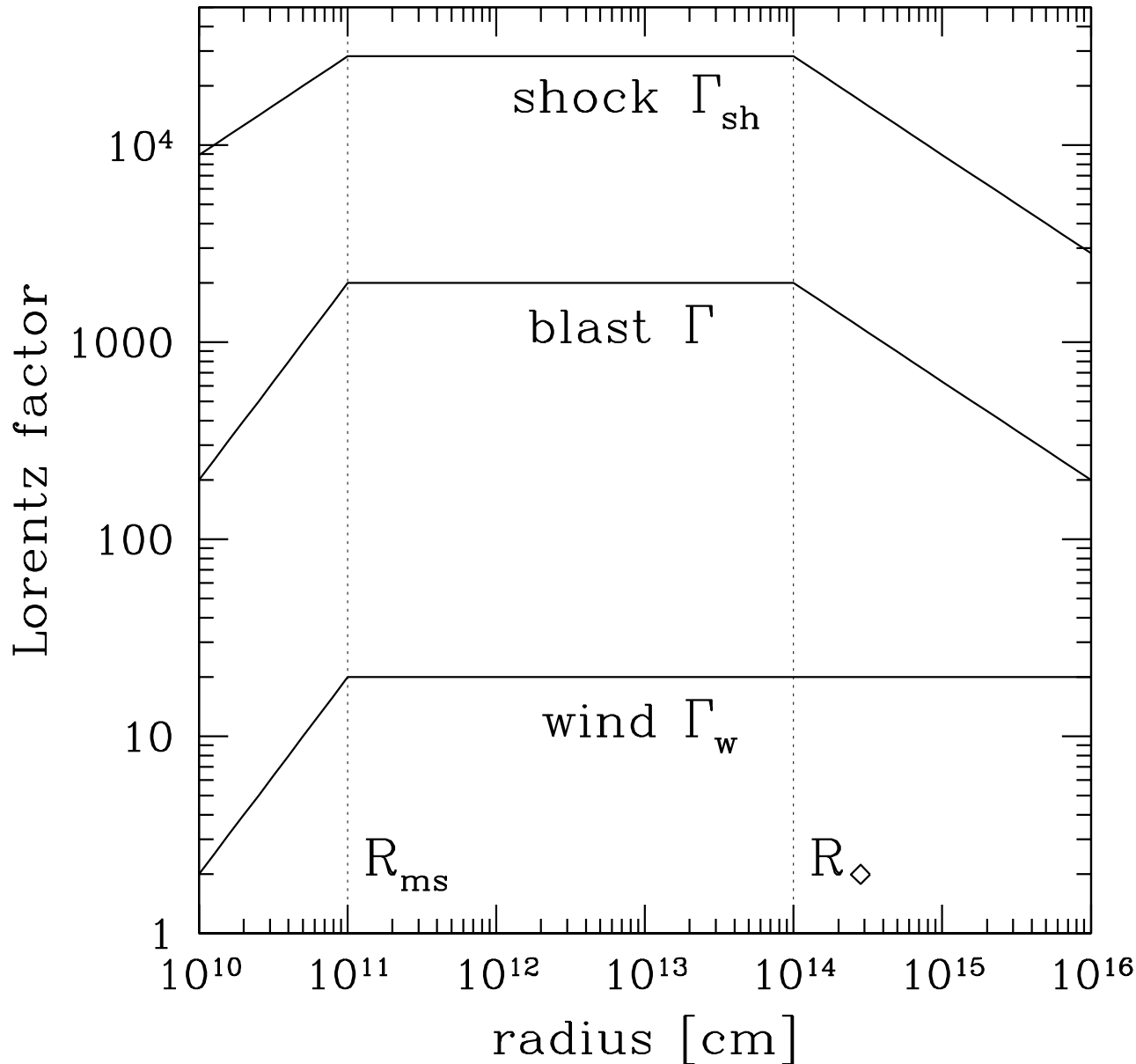
(wind energy per unit rest mass)

$$L_w \sim 10^{38} - 10^{39} \text{ erg/s}$$

$$\sigma_w \Gamma_w = \eta$$



# Flare ejecta drives a blast wave in the wind

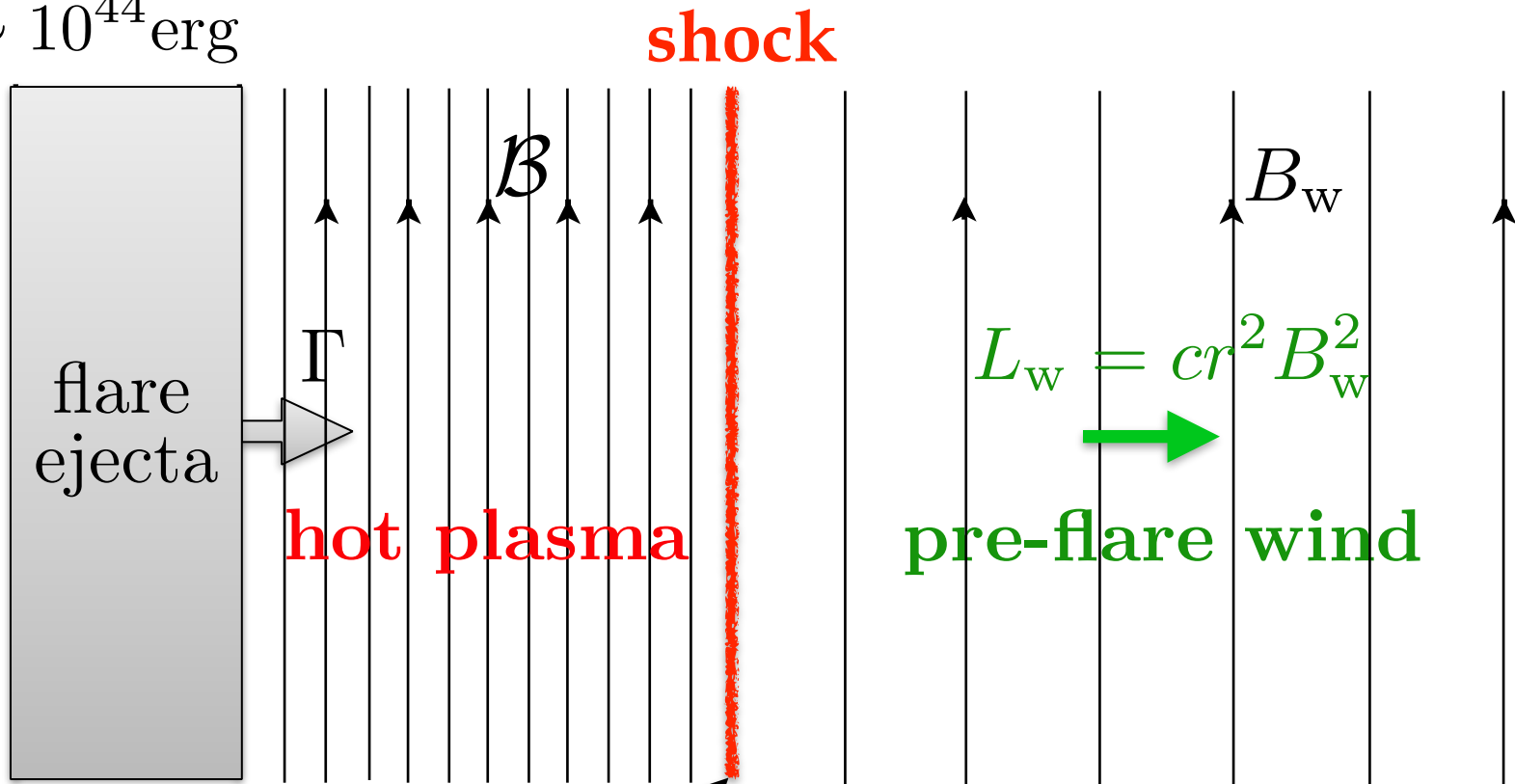


Ejecta deposits all its energy into the blast wave in the wind at

$$R_{\diamond} \sim 10^{14} \text{ cm}$$

Then the blast wave begins to decelerate

$$\mathcal{E} \sim 10^{44} \text{ erg}$$



**maser:**

PIC simulations:

Larmor rotation

$$r_L = \frac{\Gamma_{\text{rel}} m_e c^2}{e\beta}$$

Gallant et al. 1992

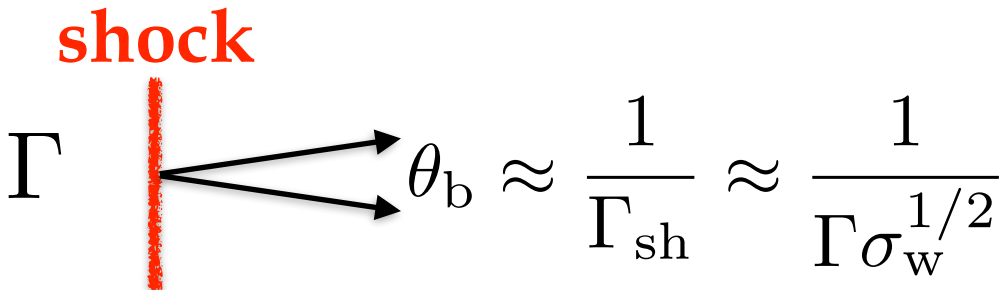
Sironi & Spitkovsky 2009

Iwamoto et al. 2017, 2018

Plotnikov & Sironi 2019

## => Radio burst:

- Frequency:  $\nu_{\text{obs}} \sim \frac{3\Gamma c}{2\pi r_L} \approx \frac{2 \text{ GHz}}{r_{14}} (L_{f,47} L_{w,39})^{1/4}$  drifts high to low

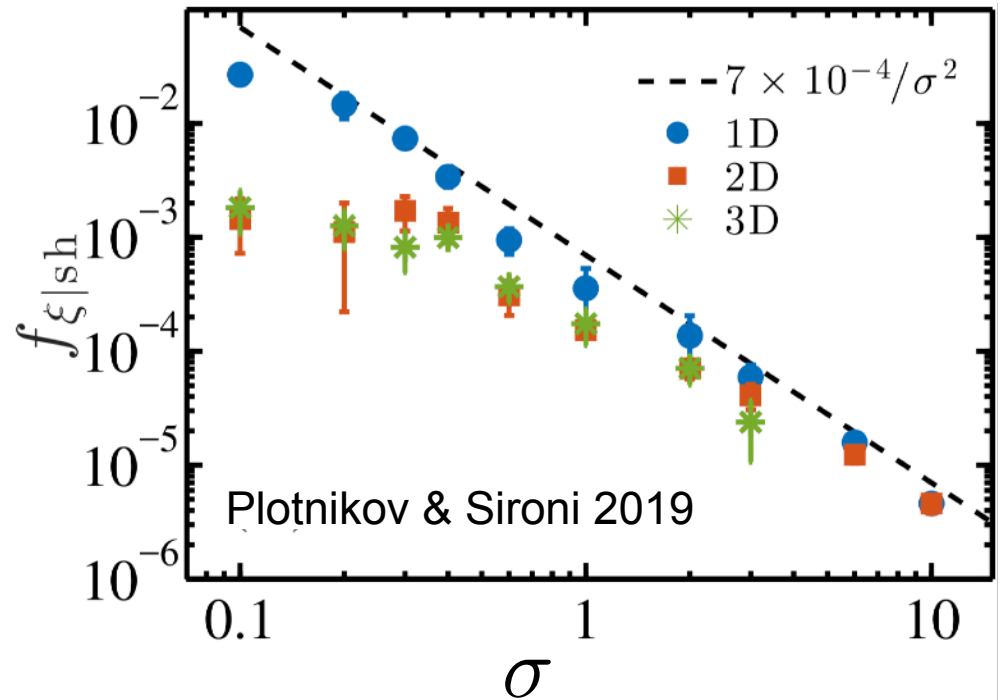
- Beaming:  $\Gamma$    $\theta_b \approx \frac{1}{\Gamma_{\text{sh}}} \approx \frac{1}{\Gamma \sigma_w^{1/2}}$

Doppler factor  $D = \Gamma(1 + \beta \cos \theta) \approx 2\Gamma$

- Duration:  $t_{\text{obs}} \sim \frac{r}{\Gamma_{\text{sh}}^2 c} \lesssim 1 \text{ ms}$

- Radiative efficiency:

$$\frac{\mathcal{E}_{\text{FRB}}}{\mathcal{E}_{\text{bw}}} \sim \frac{10^{-3}}{\sigma_w^2} \sim 10^{-5}$$



- Polarization: linear, position angle is set by the magnetar rotation axis ( $\Rightarrow$  constant)
- Bursts repeat. Energy budget over magnetar lifetime  
 $\sim 10^{-5} \mathcal{E}_B \sim 10^{44}$  erg



Wave strength parameter  $a = \frac{eE}{m_e c \omega} \approx 2 \frac{L_{42}^{1/2}}{r_{13}}$

for the shock maser  $a \sim \frac{1}{\sigma_w} \left( \frac{\mathcal{E}_{44}}{L_{w,39}} \right)^{1/4}$

## Does the radio burst escape?

Thomson optical depth of the wind:  $\tau_T \sim 10^{-9} r_{13}^{-1}$

Induced Compton scattering of the beam:  $\frac{\Delta\nu}{\nu} = -\alpha \mathcal{N} \frac{2h\tilde{\nu}}{m_e c^2} (1 - \cos \tilde{\theta}_b) \tilde{\Omega}_b \tau_T$

$\frac{|\Delta\nu|}{\nu} < 1$  requires  $L_w / \dot{N} m_e c^2 > 50$

# Observed cosmological Fast Radio Bursts (FRBs)

- Rate in the universe  $\sim$  rate of magnetar giant flares
- Some repeat frequently (in particular FRB 121102)
- Durations  $\sim 1$  ms
- Frequency drifts (high to low)
- Energies  $1e38$ - $1e41$  erg
- Linear polarization ( $\sim 100\%$  in FRB 121102, constant angle)
- Persistent radio nebula around FRB 121102 is consistent with magnetar ejecta

# Summary

- Huge sudden dissipation events around magnetars are caused by strong deformations of their magnetospheres
- Dissipation occurs through magnetic reconnection and turbulence cascade. They create bulk motions, which convert to radiation through Compton drag
- Ultra-relativistic ejecta from giant flares drive blast waves into the magnetar wind
- The blast waves emit radio bursts by the shock maser mechanism. The bursts are consistent with observed FRBs

